



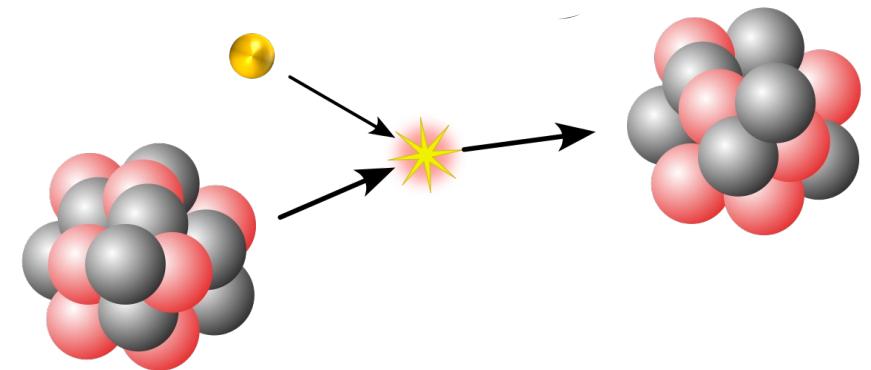
Two-body currents in magnetic dipole moments based on nuclear DFT

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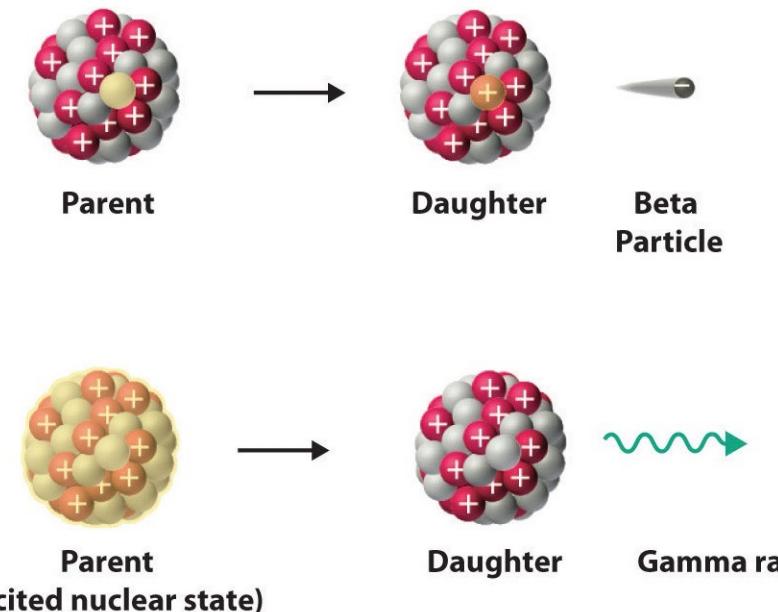
Outline

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 - Transition and moment
 - Deviation between theory and experiment
 - Higher-order corrections
 - Nuclear Density Functional Theory (DFT)
- **2BC in nuclear DFT**
- **Preliminary results**
- **Summary & outlook**

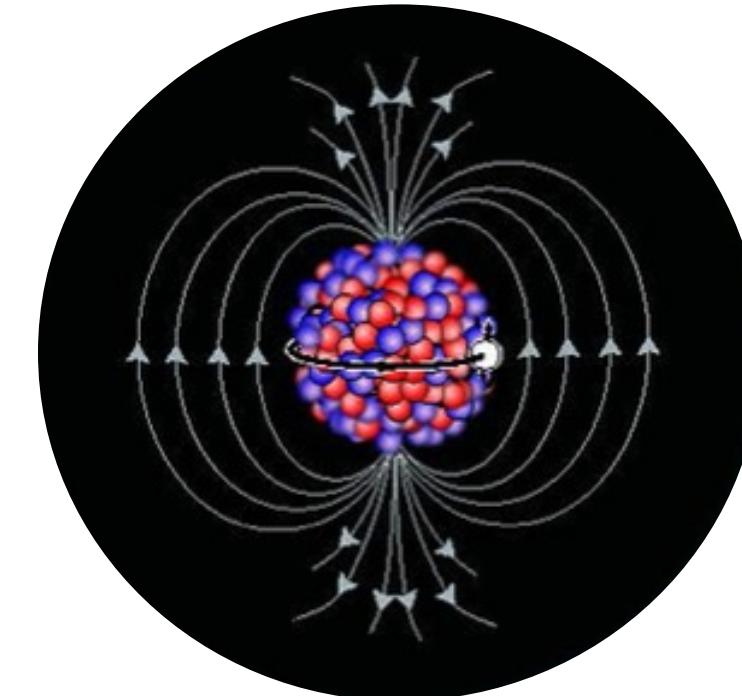


Introduction – Transition and moment

- The transitions and moments are sensitive to nuclear forces and wave functions, which are powerful tools to explore the microscopic structure of nucleus, nuclear forces and even physics beyond the standard model.
- Higher requirements for the precision of transitions and moments are raised both experimentally and theoretically.



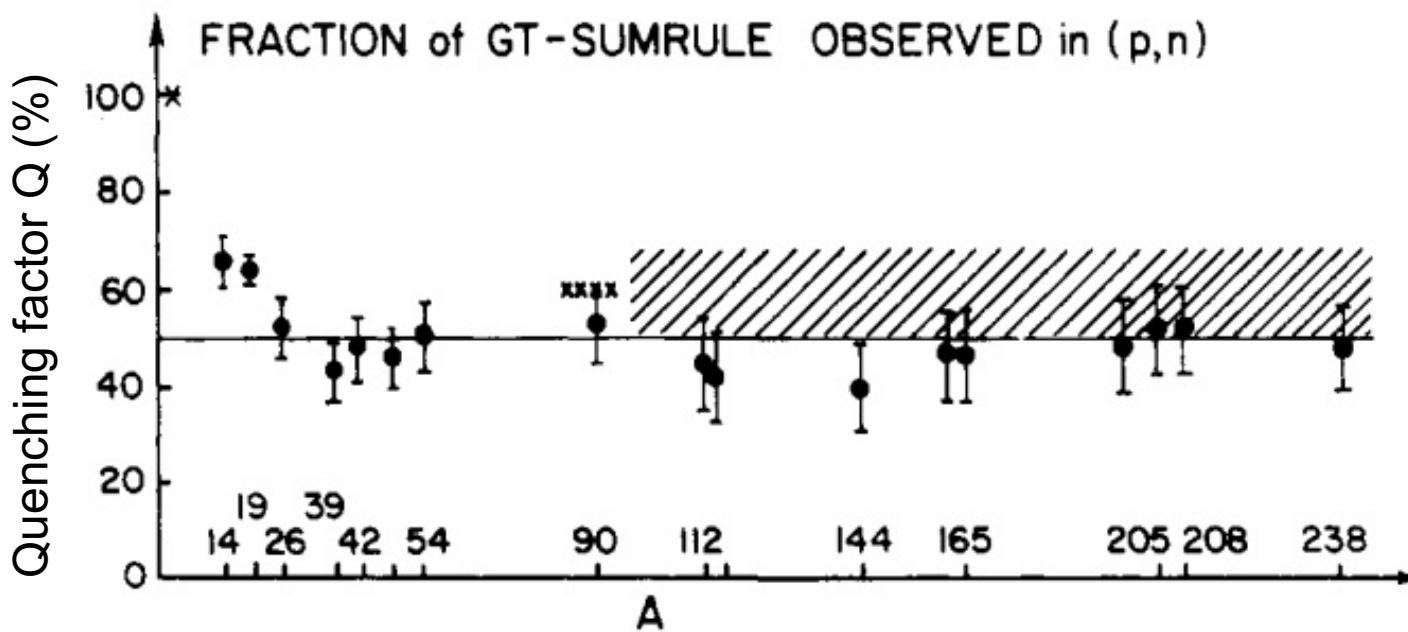
transitions



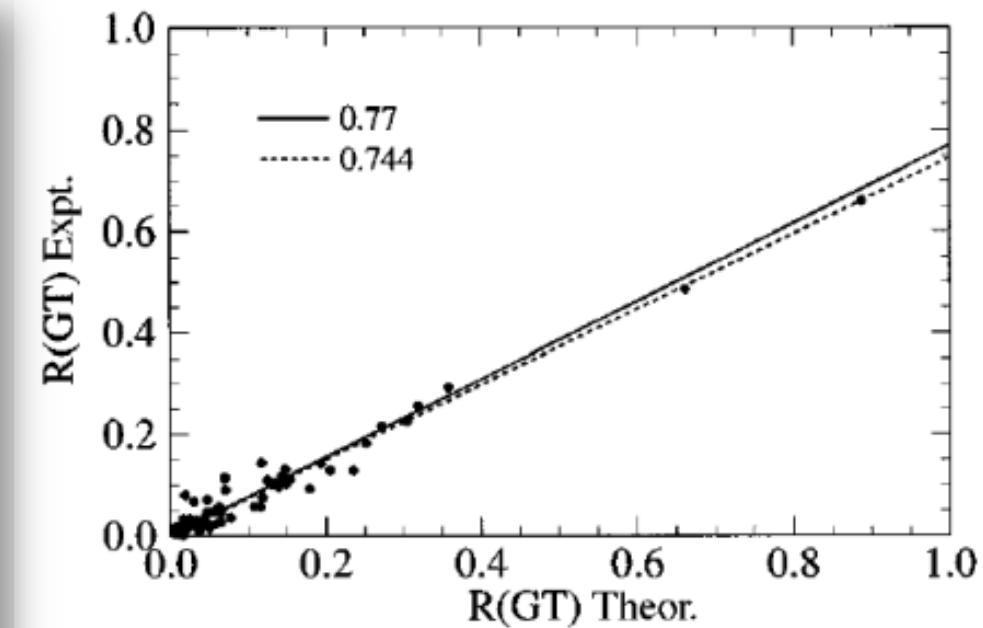
magnetic dipole moment

Introduction – Deviation between theory and experiment

Long-standing problem: Quenching of Gamow-Teller strength in nuclei



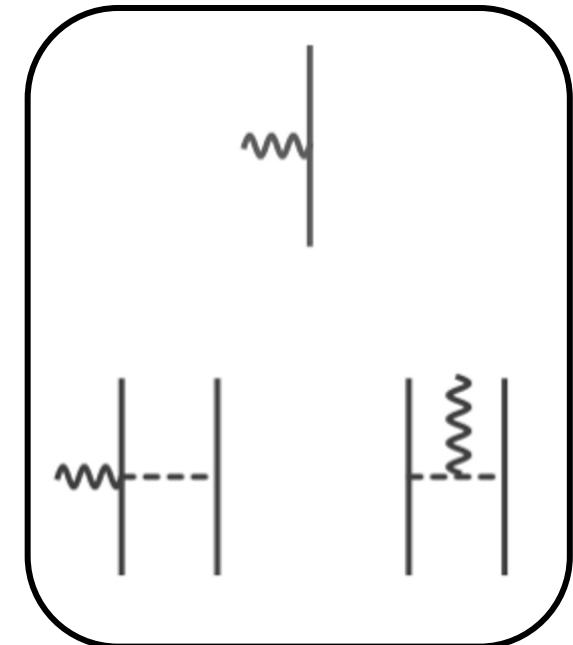
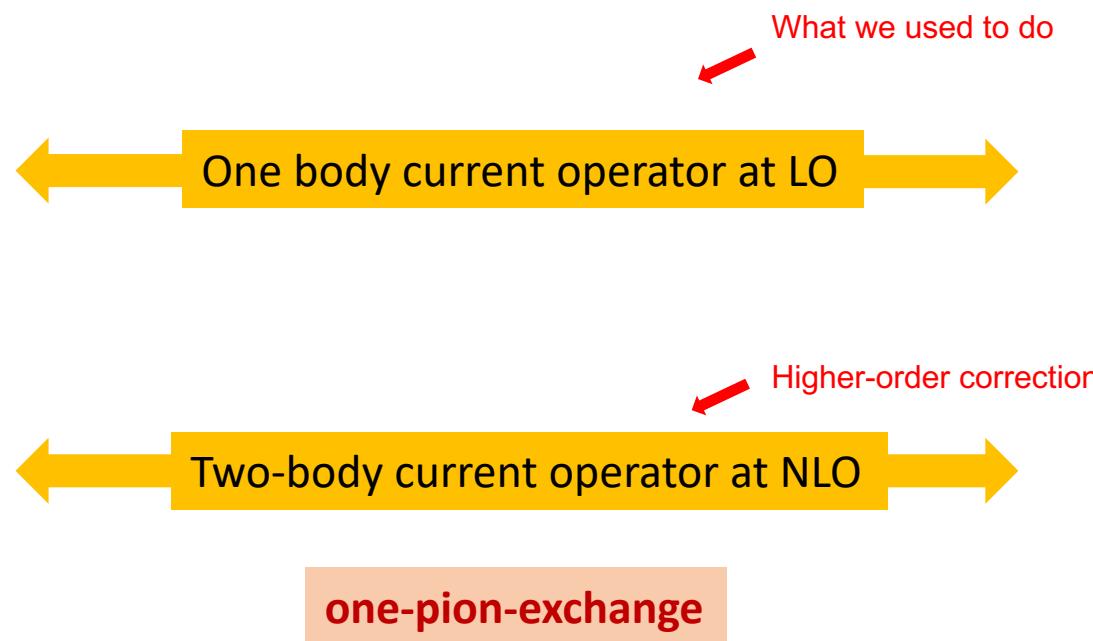
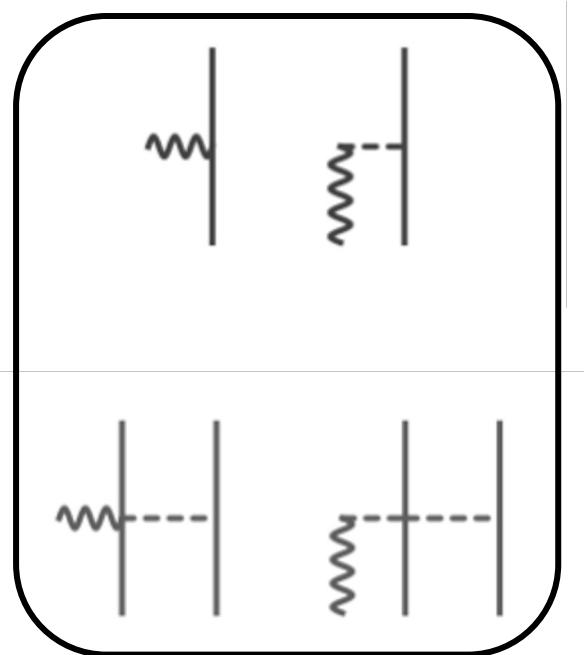
C. Gaarde, Phys. Scr. T5, 55 (1983)



G. Martinez-Pinedo et al, Phys. Rev. C 53, R2602 (1996)

Introduction – Higher-order corrections

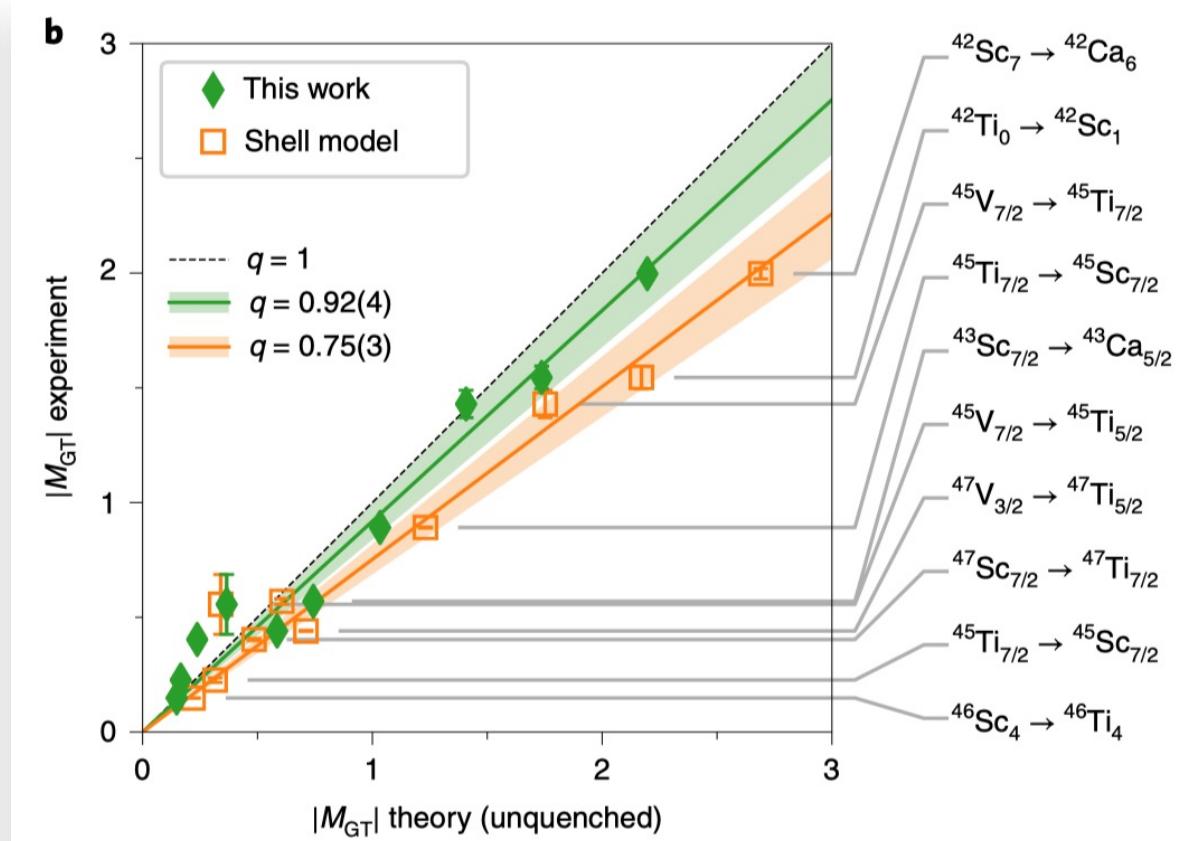
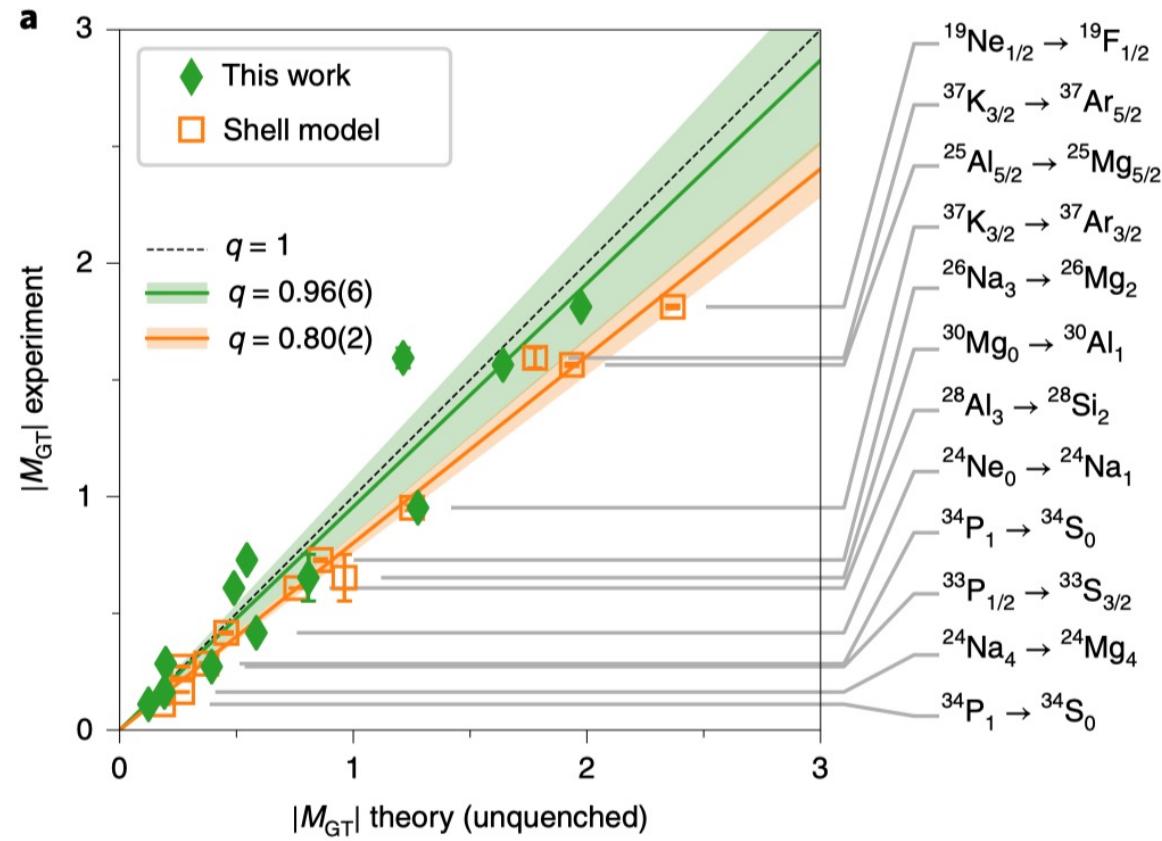
E.g. transitions/moment operators



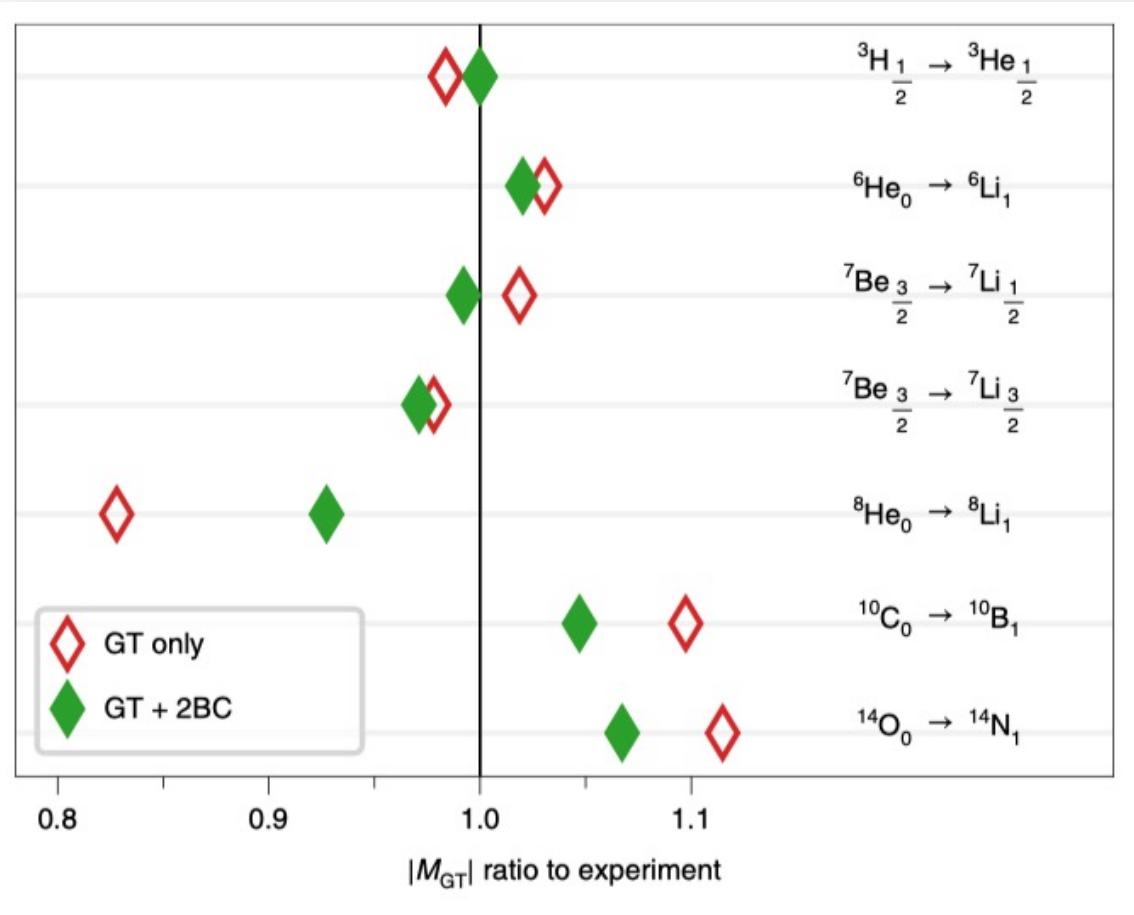
GT transition, ...

EM observables, ...

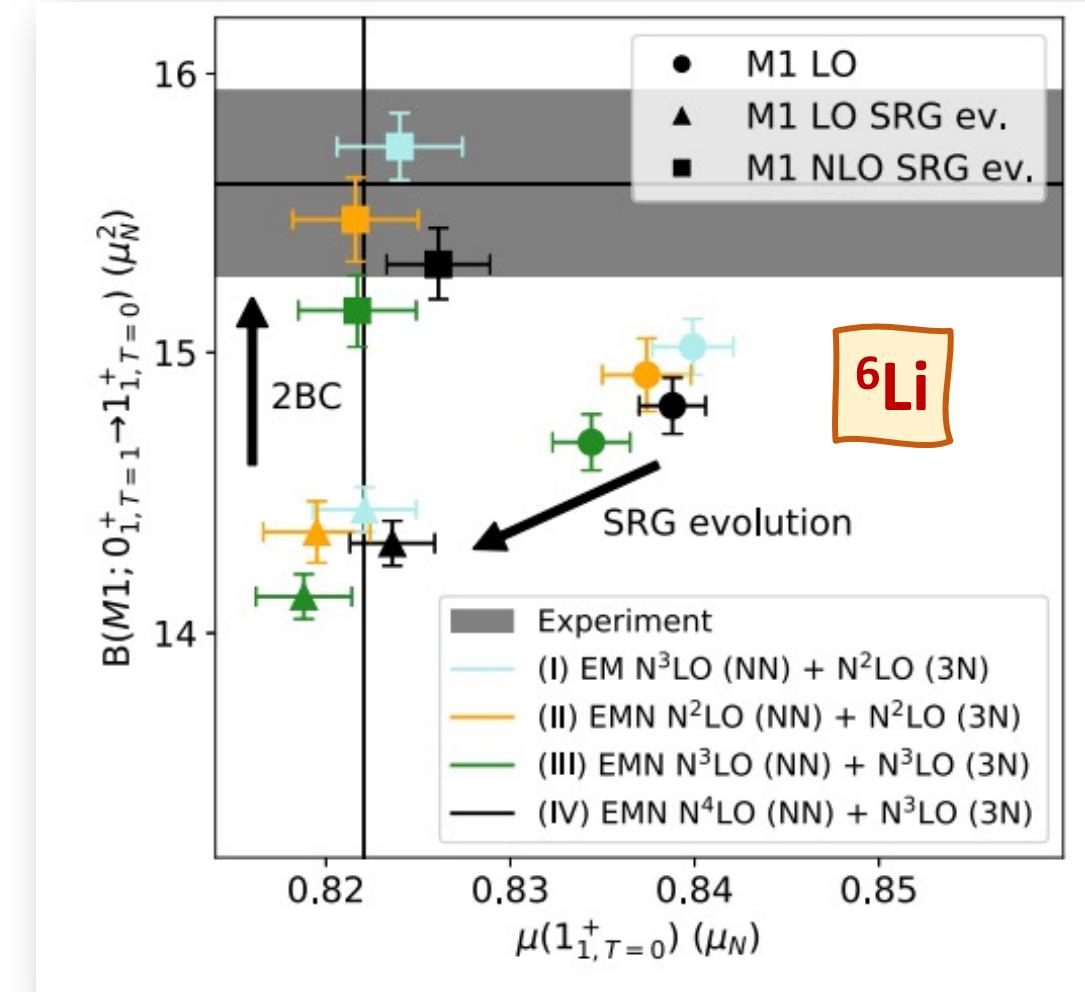
Introduction – Higher-order corrections



Introduction – Higher-order corrections

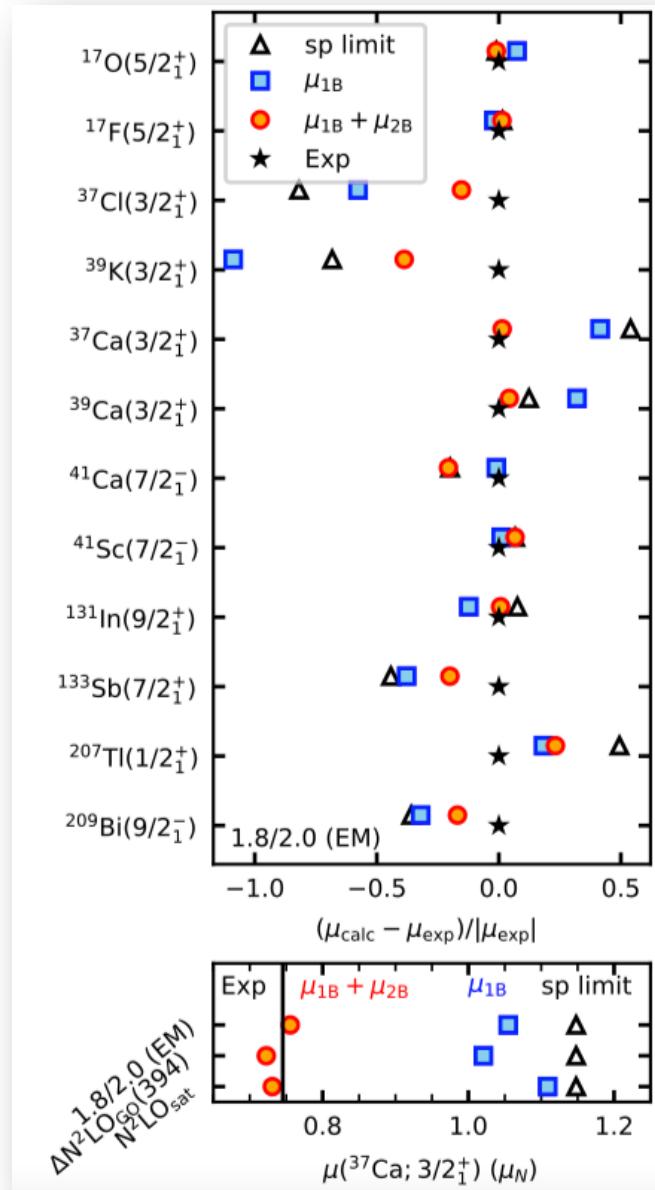


P. Gysbers et al., Nat. Phys. 15, 428 (2019)



U. Friman-Gayer et al., Phys. Rev. Lett. 126, 102501 (2021)

Introduction – Higher-order corrections



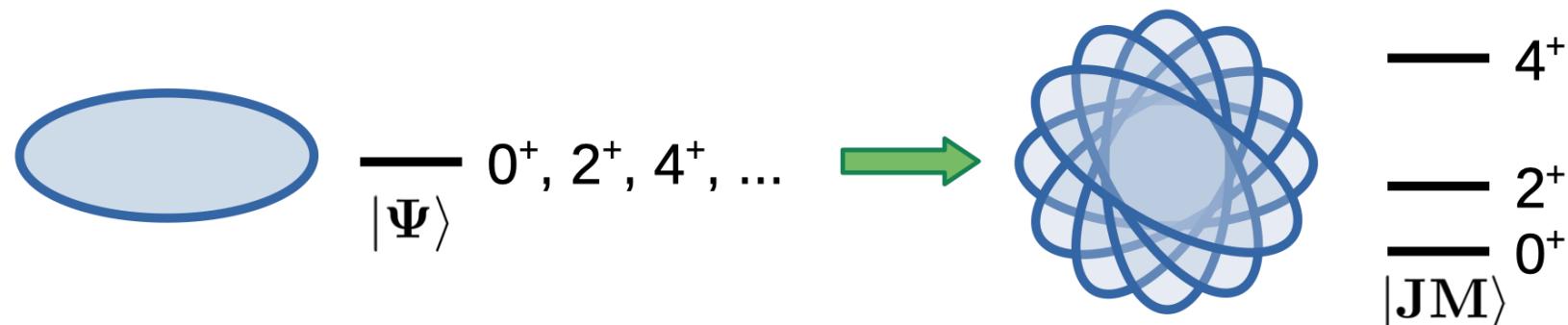
- The above configuration interaction calculations suggest that the NLO corrections could improve the description of transitions and moments.

Introduction – Nuclear Density Functional Theory (DFT)

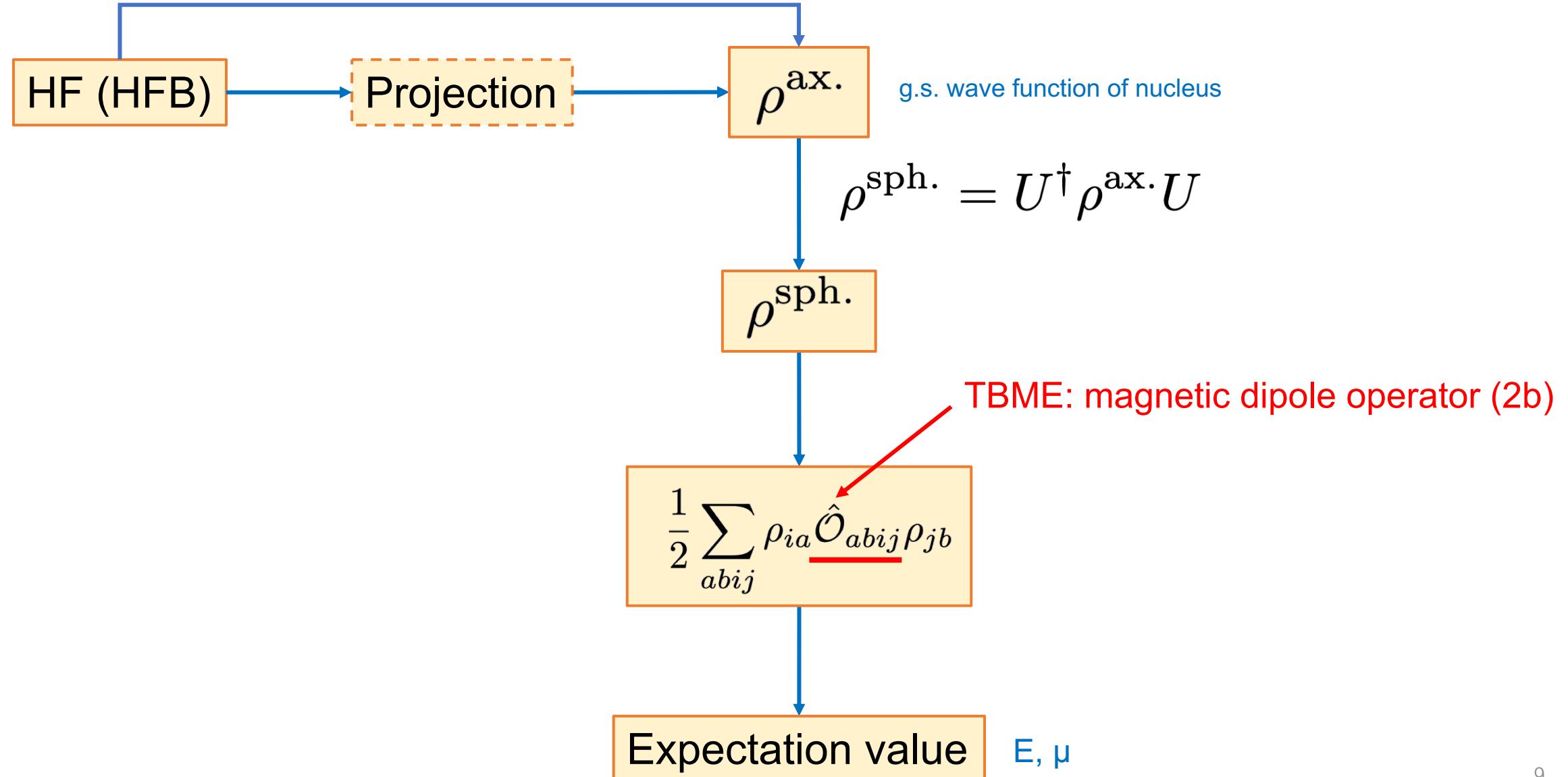
- DFT: Density matrix, variational theorem
- HFB equation (Bogoliubov-de Gennes equations):

$$\begin{pmatrix} h & \Delta \\ -\Delta^* & -h^* \end{pmatrix} \begin{pmatrix} U_\alpha \\ V_\alpha \end{pmatrix} = \begin{pmatrix} U_\alpha \\ V_\alpha \end{pmatrix} E_\alpha$$

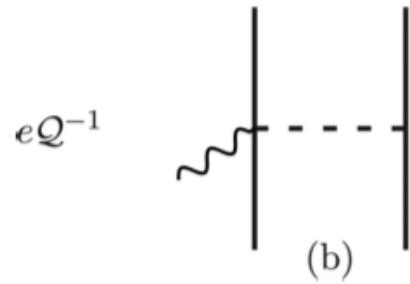
- Collective mode, large deformation, heavy nuclei.
- Angular momentum projection (PAV):



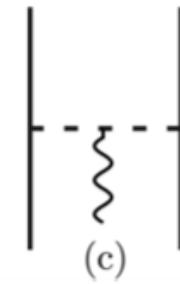
2BC in nuclear DFT



2BC in nuclear DFT – magnetic dipole moment operators



“seagull”



π -in-flight

R. Seutin et al., Phys. Rev. C 108, 054005 (2023)

$$\mathbf{j}_{\text{seagull}}^{\text{NLO}}(\mathbf{q}_1, \mathbf{q}_2) = -ie \frac{g_A^2}{F_\pi^2} G_E^V(Q^2) (\boldsymbol{\tau}_1 \times \boldsymbol{\tau}_2)_z \boldsymbol{\sigma}_1 \frac{\boldsymbol{\sigma}_2 \cdot \mathbf{q}_2}{\omega_{q_2}^2} + 1 \rhd 2,$$

$$\mathbf{j}_{\pi-\text{in-flight}}^{\text{NLO}}(\mathbf{q}_1, \mathbf{q}_2) = ie \frac{g_A^2}{F_\pi^2} G_E^V(Q^2) (\boldsymbol{\tau}_1 \times \boldsymbol{\tau}_2)_z \frac{\mathbf{q}_1 - \mathbf{q}_2}{\omega_{q_1}^2 \omega_{q_2}^2} \boldsymbol{\sigma}_1 \cdot \mathbf{q}_1 \boldsymbol{\sigma}_2 \cdot \mathbf{q}_2,$$

$$\boldsymbol{\mu}_{2b}^{\text{NLO}} = \boldsymbol{\mu}_{2b}^{\text{NLO,int}} + \boldsymbol{\mu}_{2b}^{\text{NLO,Sachs}}.$$

$$\boldsymbol{\mu}_{2b,ij}^{\text{NLO,int}}(\mathbf{r}_{ij}) = -\frac{g_A^2 m_\pi}{32\pi F_\pi^2} (\boldsymbol{\tau}_i \times \boldsymbol{\tau}_j)_z \{f(r_{ij}) [(\boldsymbol{\sigma}_i \times \boldsymbol{\sigma}_j) \cdot \hat{\mathbf{r}}_{ij}] \hat{\mathbf{r}}_{ij} - (\boldsymbol{\sigma}_i \times \boldsymbol{\sigma}_j)\} e^{-m_\pi r_{ij}}$$

Only depends on intrinsic coordinate

$$\boldsymbol{\mu}_{2b,ij}^{\text{NLO,Sachs}}(\mathbf{r}_{ij}, \mathbf{R}_{ij}) = -\frac{1}{2} (\boldsymbol{\tau}_i \times \boldsymbol{\tau}_j)_z V_{1\pi}(\mathbf{r}_{ij}) \mathbf{R}_{ij} \times \mathbf{r}_{ij}$$

Contains c.o.m. dependence

Preliminary results

without AMP, UNEDF1, HF, Nsh=7

Nucleus	State	$\mu_{2b_intr} (\mu_N)$	$\mu_{2b_sach} (\mu_N)$	$\mu_{2b} (\mu_N)$
3He	1/2+	-0.061823	0	-0.061823
5He	3/2-	0.006372	-0.021636	-0.015264
17O	5/2+	0.015166	-0.05781	-0.042644
39K	3/2+	0.025387	0.088248	0.113635
41Ca	7/2+	0.020146	-0.100212	-0.080066
131In	9/2+	-0.030637	0.229622	0.198985
133Sb	7/2+	0.01435	0.159426	0.173776

39K: with AMP, UNEDF1, Nsh=6

$$\mu_{exp} = 0.391 \mu_N$$

$$\mu_{cal} (1b) = 0.125834 \mu_N$$

$$\mu_{cal} (1b+2b) = 0.231247 \mu_N$$

Summary & outlook

- Introduce the implementation of 2BC based on nuclear DFT to explore the contribution of higher-order corrections to magnetic dipole moments.
- 2BC corrections improve the magnetic dipole moment generally.
- Outlook:
 - Improve the computational efficiency and compute in larger shells and heavier nuclei.
 - Transitions: e.g. M1 transitions.

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- TU Darmstadt: T. Miyagi

References

- [1] C. Gaarde, Phys. Scr. T5, 55 (1983).
- [2] A. Arima, Adv. Nucl. Phys. 18, 1 (1987).
- [3] I. Towner, Phys. Rep. 155, 263 (1987).
- [4] G. Martinez-Pinedo et al, Phys. Rev. C 53, R2602 (1996).
- [5] P. Gysbers et al., Nature Phys. 15, 428 (2019).
- [6] E. Epelbaum, arXiv:1908.09349 (2019).
- [7] U. Friman-Gayer et al., Phys. Rev. Lett. 126, 102501 (2021).
- [8] A. Vernon et al., Nature 607, 260 (2022).
- [9] P. Sassarini et al., J. Phys. G Nucl. Part. Phys. 49, 11LT01 (2022).
- [10] J. Bonnard et al., Phys. Lett. B 843, 138014 (2023).
- [11] R. Seutin et al., Phys. Rev. C 108, 054005 (2023).
- [12] T. Miyagi et al., Phys. Rev. Lett. 132, 232503 (2024).