

Complex potential at $T > 0$ from fine lattices

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+ HotQCD COLLABORATION

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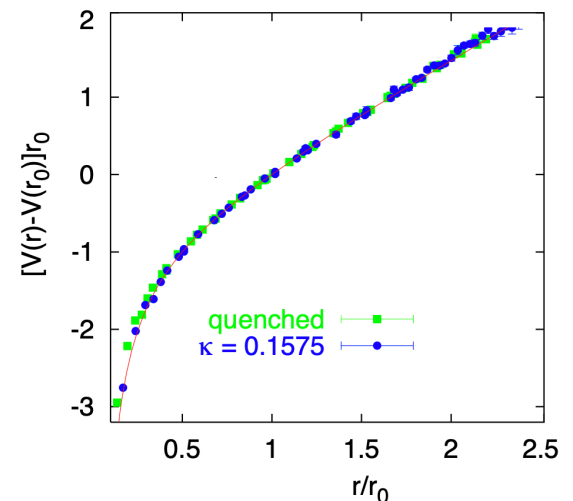
Introduction

Bound states of static quark and anti-quark pair: Probe for existence of Quark Gluon plasma in Heavy ion collisions (Alexander Rothkopf, Heavy quarkonium in extreme conditions, Physics Reports, Volume 858, 2020,).

- Time evolution in Real-Time suffers from sign problem. (QFT suffers from sign problem; see Daniel's talk tomorrow)
- If separation of scales is present: Use EFTs (NRQCD and pNRQCD): describe physics in form of potential?

At $T=0$ Schrodinger like potential picture has been observed (G. Bali Phys.Rept. 343 (2001) 1-136).

$$i \partial_t W_{\square}(t, r) = \Phi(t, r) W_{\square}(t, r)$$
$$V(r) = \lim_{t \rightarrow \infty} \Phi(t, r)$$



Introduction

- It remains to be seen if the potential picture even holds for $T>0$, if it does how is it modified? ; Affirmative indications in Quenched systems (A.Rothkopf and Y Burnier, *Phys.Rev.D* 95 (2017) 5, 054511).
- Spectral function is a link between real and imaginary time: (A Rothkopf ,T Hatsuda, S Sasaki *Phys.Rev.Lett.* 108 (2012) 162001).

$$W_{\square}(r, t) = \int d\omega e^{-i\omega t} \rho_{\square}(r, \omega) \quad \longleftrightarrow \quad W_{\square}(r, \tau) = \int d\omega e^{-\omega\tau} \rho_{\square}(r, \omega)$$

- Computation of Spectral Function : Ill-posed Inverse problem
- Potential linked to the dominant peak position and width of Spectral Function (Yannis Burnier and Alexander Rothkopf 1208.1899).
- In HTL regime there exists a complex potential with screened real part (M.Laine et. al *JHEP* 03 (2007), 054).

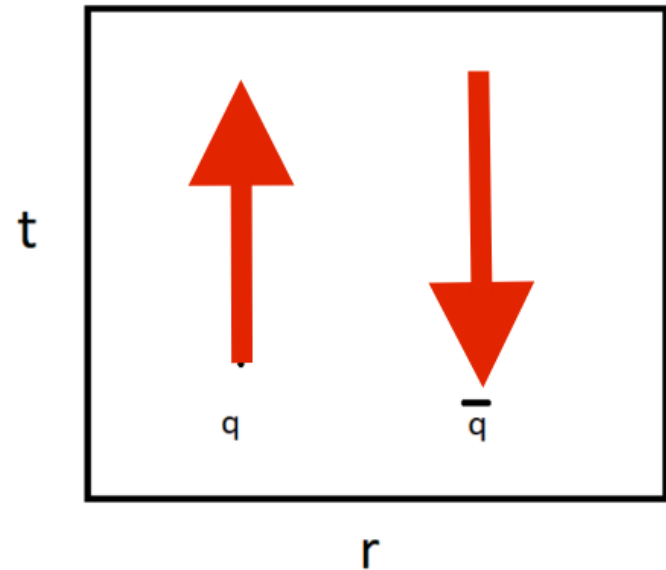
Lattice setup

- (2+1)-flavour QCD - configurations generated by HotQCD and TUMQCD collaborations.

- Using highly improved staggered quark (HISQ) action.

$N_\sigma^3 \times N_\tau$ lattices. $N_\tau = 10, 12, 16$ and $N_\sigma/N_\tau = 4$

- Calculate Wilson Line correlator in Coulomb Gauge.
- Fix box approach; temp range - 140MeV to 2GeV.
- Pion mass 160MeV, Kaon mass physical.

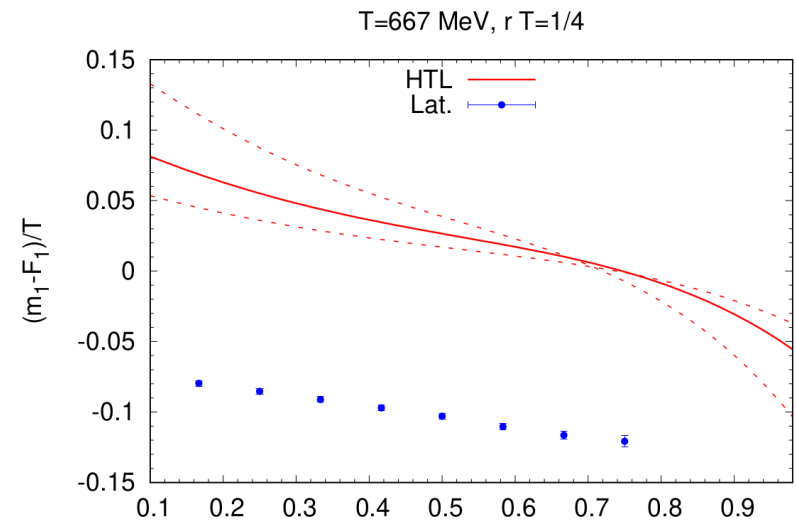
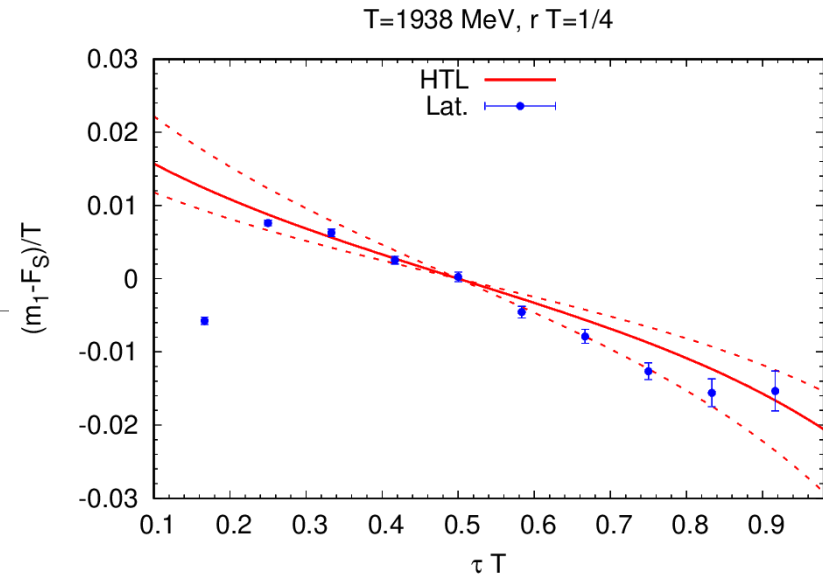


HTL comparison

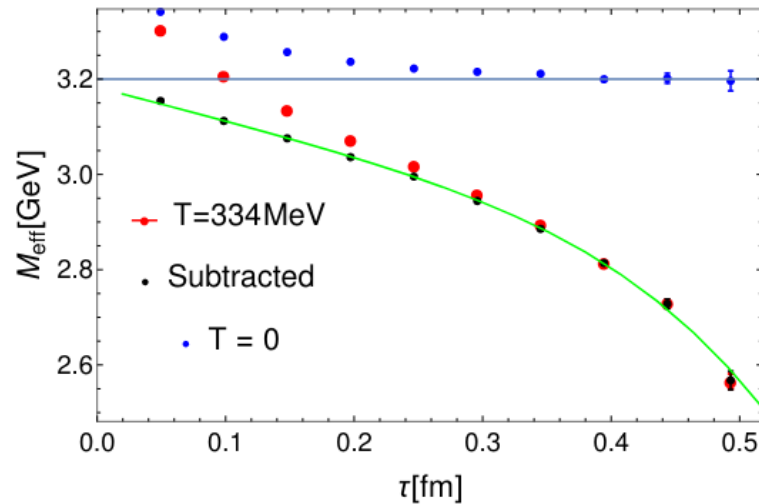
Define effective mass of the correlation function:

$$m_{eff}(r, \tau, T) = -\partial_{\tau} \ln W(r, \tau, T)$$

- Subtracting UV part using $T=0$ correlator results in linear behavior at small tau.
- Plot shows effective mass subtracted from Free Energies.
- HTL does not quantitatively fit the data except for some specific T and separation distance.



Spectral Function Model Fits



- Lattice data sensitive only to peak position (Ω) and effective width (Γ).

- Parametrize Correlator as: $C_{sub}(\tau, T) \sim \exp(-\Omega\tau + \frac{1}{2}\Gamma^2\tau^2 + O(\tau^3))$

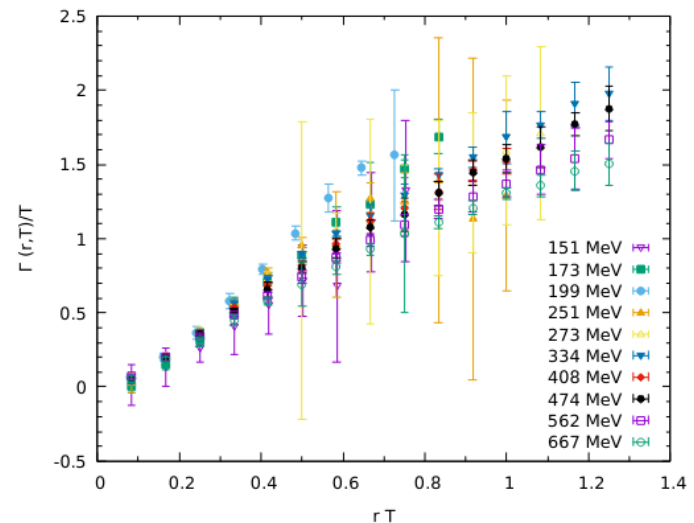
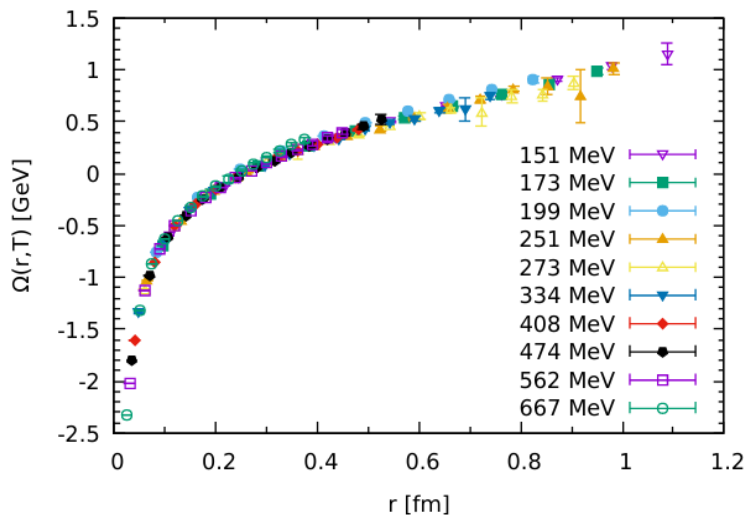
$$\rho_r(\omega, T) = A(T) \exp\left(-\frac{[\omega - \Omega(T)]^2}{2\Gamma^2(T)}\right) + A^{\text{cut}}(T) \delta(\omega - \omega^{\text{cut}}(T))$$

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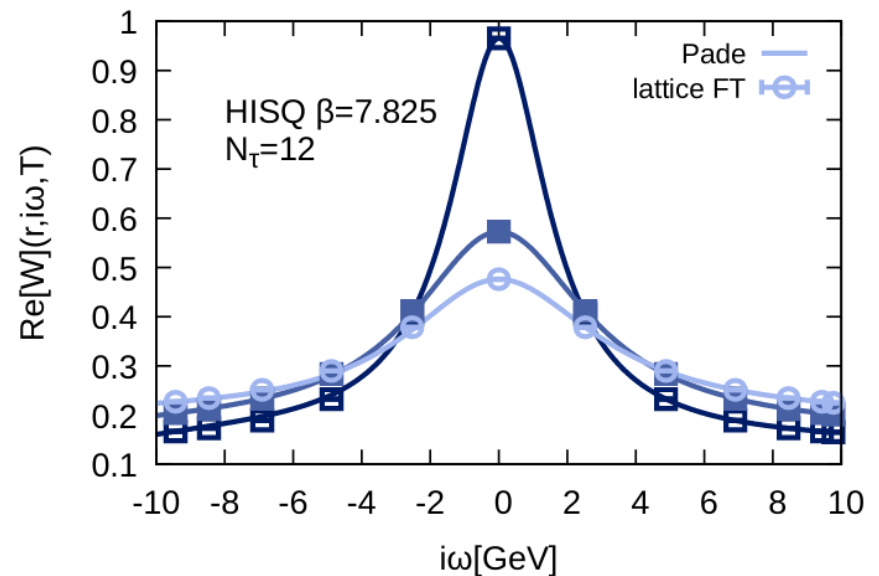
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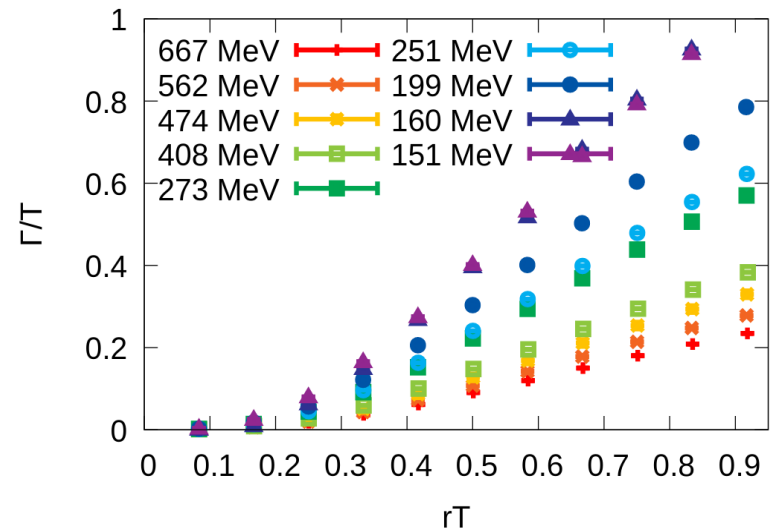
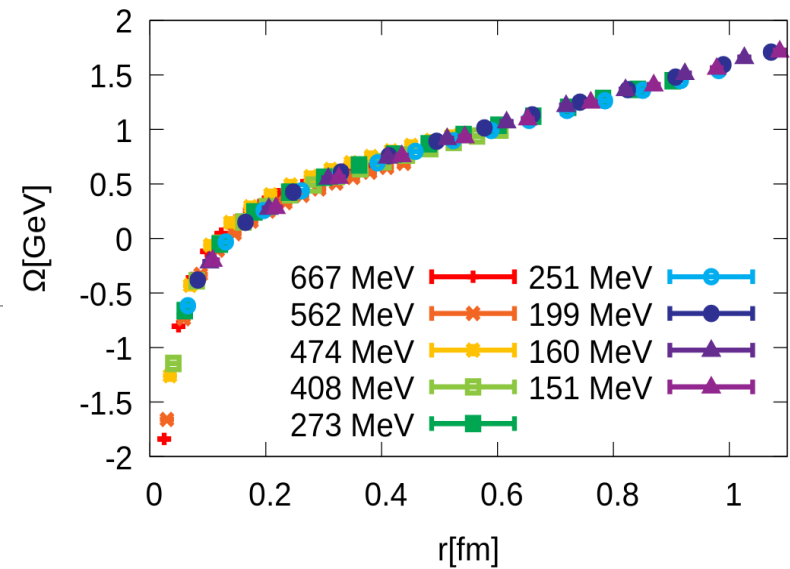
Pade' Interpolation

- Transform the Euclidean correlator into Matsubara frequency space.
- Implement Pade approximation in the form of continued fraction according to Schlessinger prescription (L. Schlessinger, Phys. Rev. 167, 1411 (1968)).
- This is interpolation of data and not fitting. Does not require minimization.
- Obtain pole structure from rational function: Directly related to the peak position (Ω) and width (Γ).



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Spectral Function Extraction using Bayesian Method

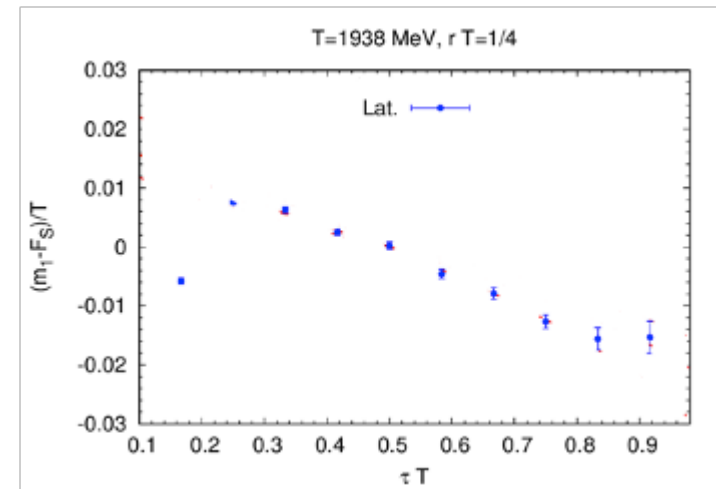
$$P[\rho|D, I] \propto P[D|\rho, I]P[\rho|I] = \exp[-L + \alpha S_{BR}]$$

L is the usual quadratic distance used in chi-square fitting.

The prior probability $P(\rho|I) = \exp(\alpha S_{BR})$ acts as a regulator

$$S_{BR} = \int d\omega \left(1 - \frac{\rho(\omega)}{m(\omega)} + \log \left[\frac{\rho(\omega)}{m(\omega)} \right] \right).$$

- Look for the most probably spectrum by locating the extremum of the posterior.
- Effective masses at high T show non-monotonicity at small tau; non-positive spectral function.--- cannot use Bayesian Methods.



HTL inspired fits

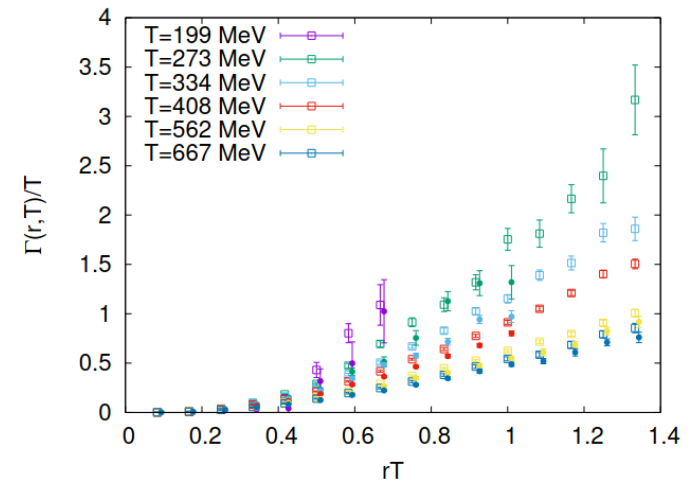
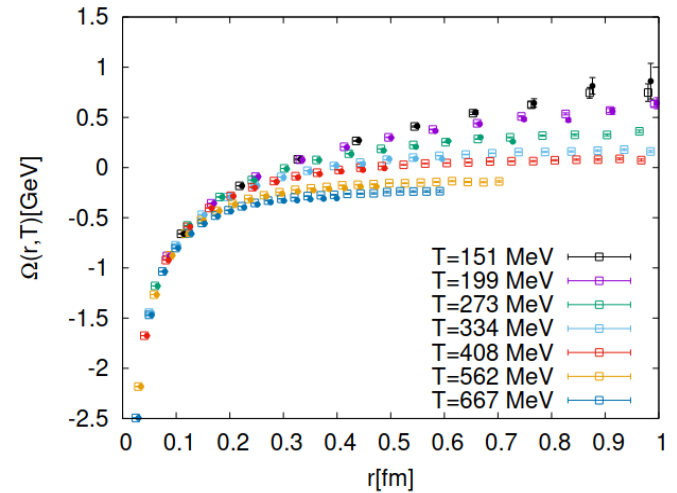
Peak position (Ω) and width (Γ) interpreted as the real and imaginary part of thermal static energy E_s (D. Bala and S. Datta, Phys. Rev. D 101, 034507(2020)).

$$E_s(r, T) = \lim_{t \rightarrow \infty} i \frac{\partial \log W(r, t, T)}{\partial t} = \Omega(r, T) - i\Gamma(r, T).$$

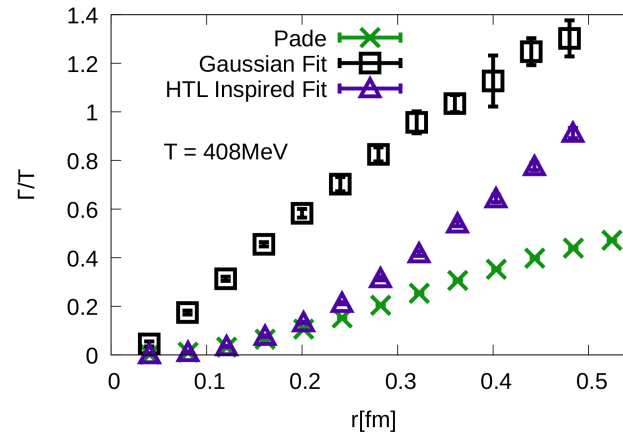
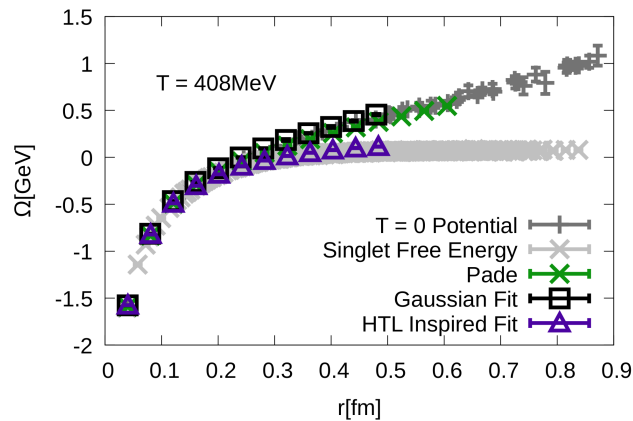
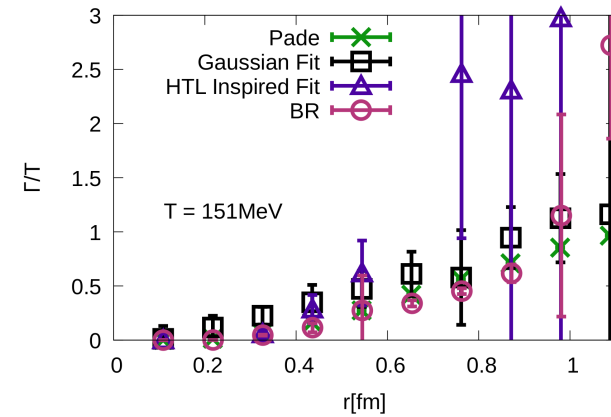
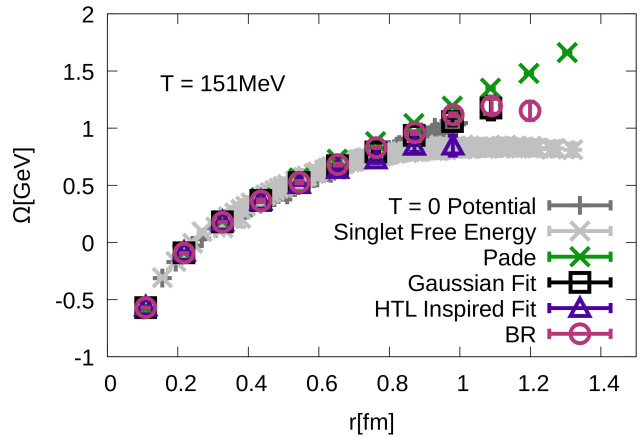
$W(r, t, T)$ is the Fourier transform of the spectral function $\rho_r(r, \omega)$

$$m_{eff}(r, n_\tau = \tau/a)a = \log \left(\frac{W(r, n_\tau, N_\tau)}{W(r, n_\tau + 1, N_\tau)} \right)$$

$$= \Omega(r, T)a - \frac{\Gamma(r, T)aN_\tau}{\pi} \log \left[\frac{\sin(\pi n_\tau / N_\tau)}{\sin(\pi(n_\tau + 1) / N_\tau)} \right]$$



Comparison of Results



Quenched QCD

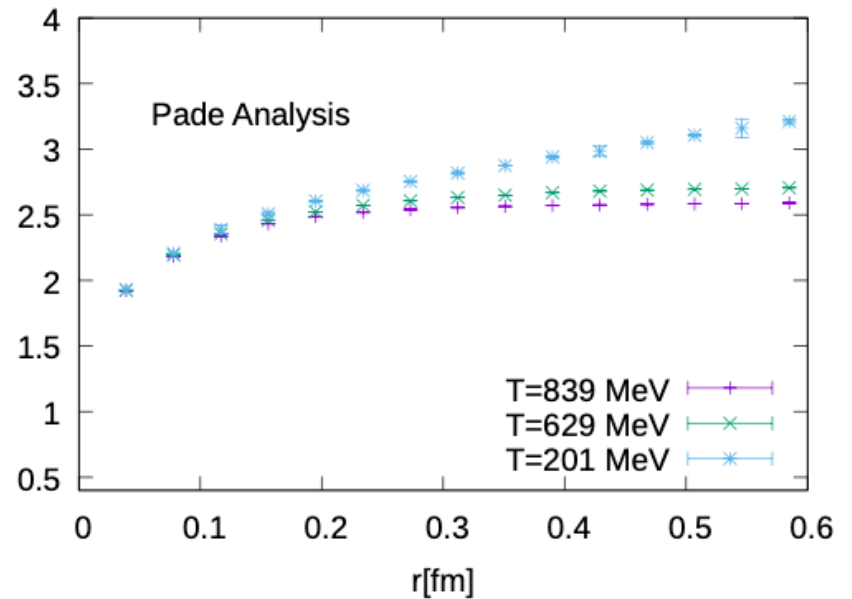
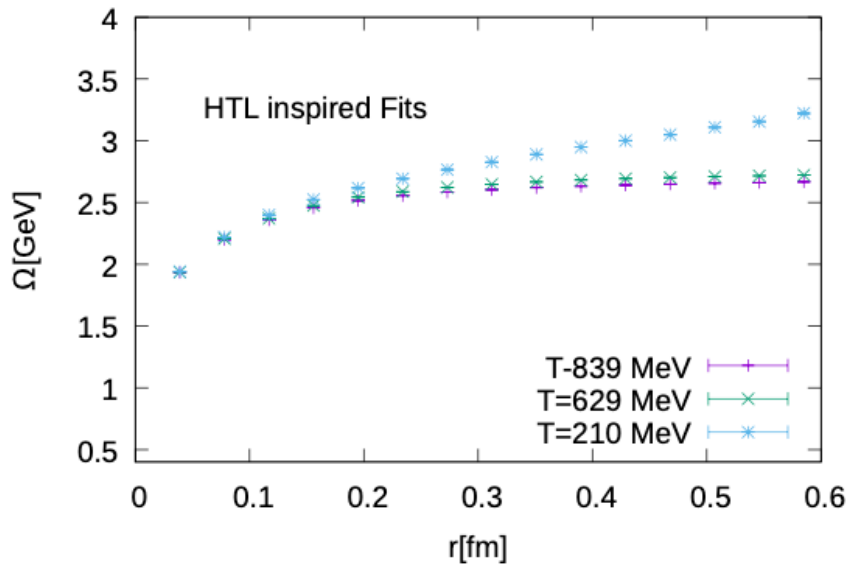
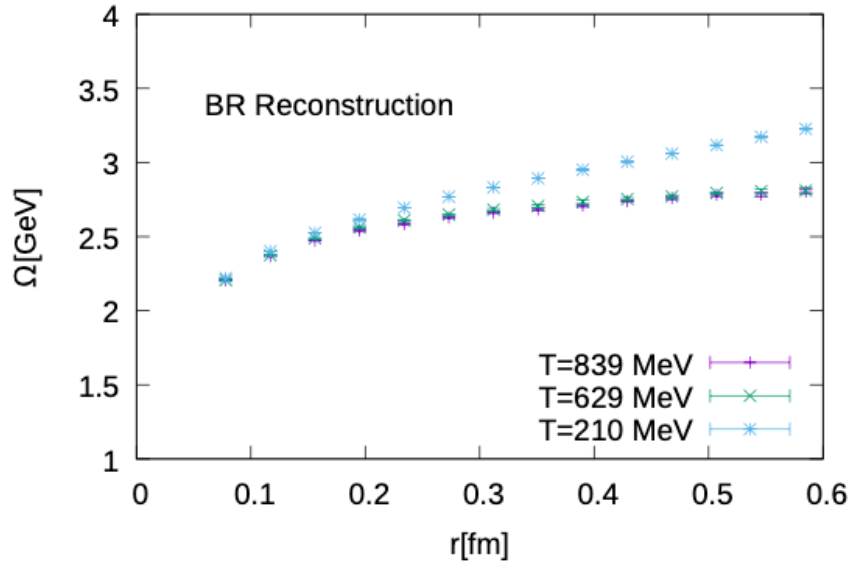
- We see that from the Gaussian fits and Pade' peak position (Ω) for HISQ Lattices is temperature independent;; Quite puzzling.
- Results obtained are very different from previous studies of Quenched Lattices (A Rothkopf, Y. Burnier 1607.04049);; Different methods used.
- Need further investigation:: Check new methods with Bayesian reconstruction.
- Check robustness of methods with new Quenched QCD lattices.

Lattice setup

- Using anisotropic Wilson action configurations generated using OpenQCD code from FASTSUM collaboration.
- $N_s^3 \times N_\tau$ lattices. $N_\tau = 24, 32, 48, 96, 192$ and $N_s = 64$.
- Use renormalized anisotropy : $a_t / a_s = 4$ ($\xi=3.5$)
- Calculate Wilson Line correlator in Coulomb Gauge (simulatedQCD code).
- Fix box approach; temp range 105-839 MeV .

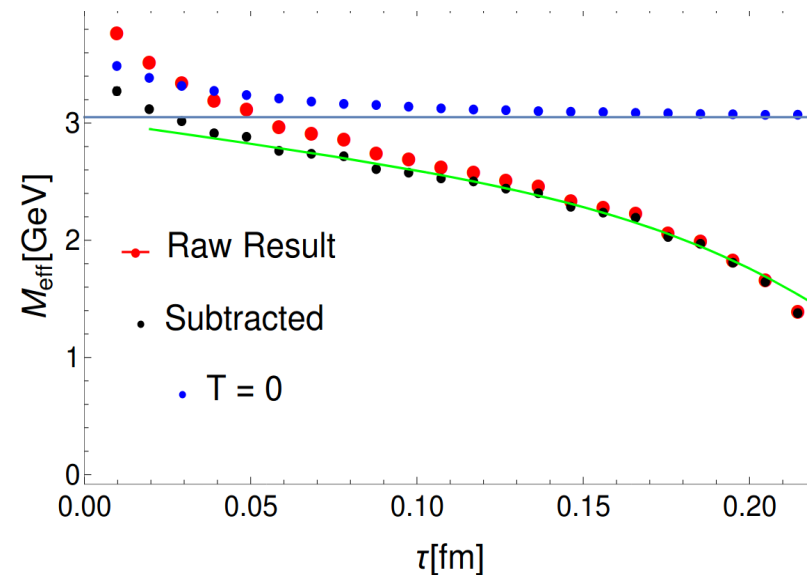
Preliminary Results

$T_c = 270$ MeV



Effective masses and correlator

- M_{eff} at $r/a = 12$ for $N_\tau = 96$ (blue), $N_\tau = 24$ ($T=839\text{MeV}$).
- Continuum extracted from $N_\tau = 96$ and removed from $N_\tau = 24$ (black).
- Lines shows exponential and gaussian plus delta function fits.
- Small τ behaviour shows T dependence.
- Cannot do Gaussian fits.



Summary

- Spectral functions of Wilson line correlators encode the real and imaginary part of the complex potential between static quark-antiquark pairs.
- We show analysis of spectral structure with four different methods.
- We see that from the Gaussian fits and Pade' peak position (Ω) is temperature independent in HISQ lattices.
- Results obtained are very different from previous quenched QCD studies.
- Peak position of new quenched study shows screening above T_c
- Quenched study is still work in progress.
- Need to understand origin of differences between quenched and dynamical QCD.