

# Surprises and challenges in exclusive $J/\psi$ photoproduction on nuclei



**Vadim Guzey**

University of Jyväskylä & Helsinki Institute of Physics,  
University of Helsinki, Finland



K.J. Eskola, C.A. Flett, V. Guzey, T. Löytäinen, H. Paukkunen, PRC 106 (2022) 035202  
and arXiv:2210.16048 [hep-ph]

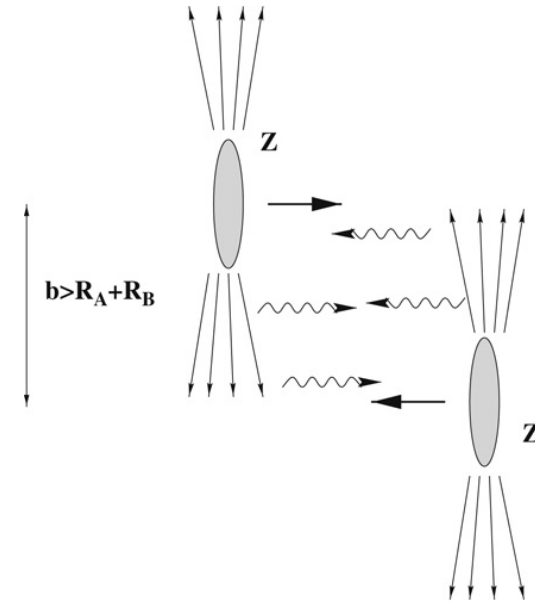
## Outline:

- Exclusive  $J/\psi$  photoproduction on nuclei in ultraperipheral collisions (UPCs) at the LHC: constraints on small- $x$  nuclear gluon density
- Exclusive  $J/\psi$  photoproduction in NLO pQCD: strong scale dependence, uncertainty due to nPDFs, and quark dominance
- Theoretical challenges: small- $x$  resummation, non-relativistic corrections to charmonium wave function

**Particle Physics Day 2022, Helsinki, Nov 24, 2022**

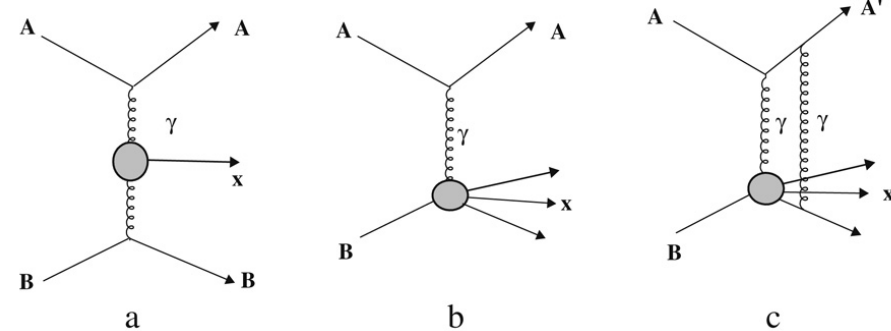
# Ultrapерipheral collisions at the LHC

- **Ultrapерipheral collisions (UPCs)**: ions interact at large impact parameters  $b \gg R_A + R_B \rightarrow$  strong interactions suppressed  $\rightarrow$  interaction via quasi-real photons in Weizsäcker-Williams equivalent photon approximation, Budnev, Ginzburg, Meledin, Serbo, Phys. Rept. 15 (1975) 181



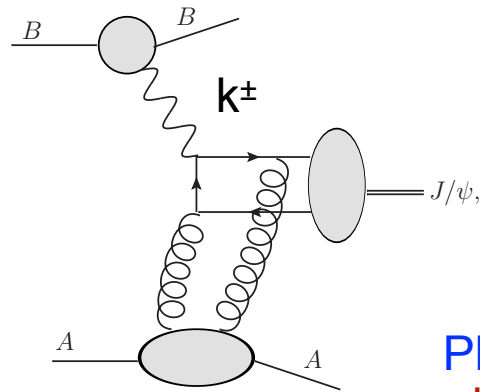
- UPCs@LHC allow one to study  $\gamma\gamma$ ,  $\gamma p$  and  $\gamma A$  interactions at unprecedentedly high energies (energy frontier) reaching:  $W_{\gamma p} = 5 \text{ TeV}$ ,  $W_{\gamma A} = 700 \text{ GeV}/A$ ,  $W_{\gamma\gamma} = 4.2 \text{ TeV}$

- UPCs can be used to study open questions of **proton and nucleus structure in QCD** and search for new physics  $\rightarrow$  e.g., **new info on quark and gluon distributions in nuclei at small x.**



# Exclusive $J/\psi$ photoproduction in UPCs

- Cross section of exclusive, coherent  $J/\psi$  photoproduction in Pb-Pb UPCs  $\rightarrow$  two terms corresponding to high photon mom.  $k^+$  (low  $x$ ) and low  $k^-$  (high  $x$ )



$$\frac{d\sigma^{AB \rightarrow AJ/\psi B}}{dy} = \left[ k \frac{dN_{\gamma/B}}{dk} \sigma^{\gamma A \rightarrow J/\psi A} \right]_{k=k^+} + \left[ k \frac{dN_{\gamma/A}}{dk} \sigma^{\gamma B \rightarrow J/\psi B} \right]_{k=k^-}$$

Photon flux from QED:  
 - high intensity  $\sim Z^2$   
 - high photon energy  $\sim \gamma L$

Photoproduction cross section

$$k^\pm = \frac{M_{J/\psi}}{2} e^{\pm y}$$

Photon momentum from  $J/\psi$  rapidity  $y$

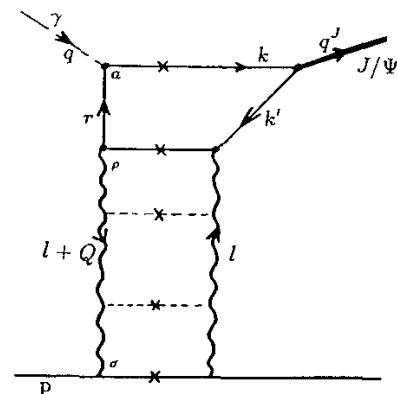
- In leading logarithmic approximation (LLA) of pQCD, Ryskin, Z. Phys. C57 (1993) 89; Frankfurt, Koepf, Strikman, PRD 57 (1998) 512; Frankfurt, McDermott, Strikman, JHEP 03 (2001) 045

$$\frac{d\sigma^{\gamma p \rightarrow J/\psi p}(t=0)}{dt} = \frac{12\pi^3 \Gamma_V M_V^3}{\alpha_{\text{e.m.}} (4m_c^2)^4} [\alpha_s(Q_{\text{eff}}^2) x g(x, Q_{\text{eff}}^2)]^2 C(Q^2=0)$$

$\Gamma_V$  is  $J/\psi$  leptonic decay width

gluon density at  $x=(M_{J/\psi})^2/W^2$  and  $Q_{\text{eff}}^2=2.5-3 \text{ GeV}^2$

depends on details of charmonium distribution amplitude



# Constraints on small-x gluon density in nuclei

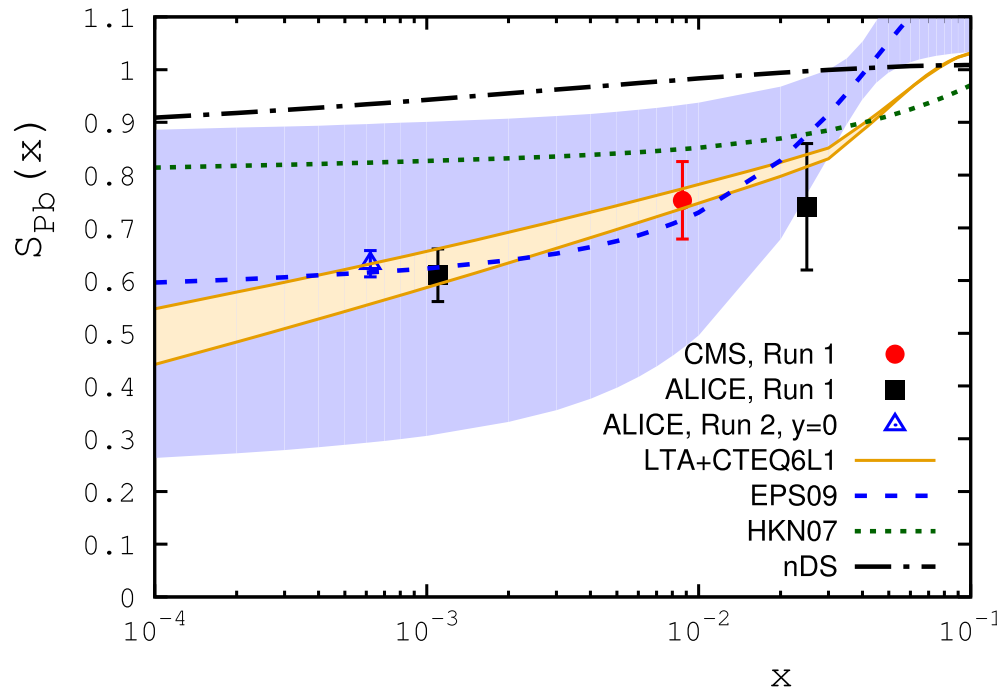
- Ratio of nucleus and proton cross sections  $\rightarrow$  nuclear suppression factor  $S$

$$S(W_{\gamma p}) = \left[ \frac{\sigma_{\gamma Pb \rightarrow J/\psi Pb}}{\sigma_{\gamma Pb \rightarrow J/\psi Pb}^{\text{IA}}} \right]^{1/2} = \kappa_{A/N} \frac{G_A(x, \mu^2)}{AG_N(x, \mu^2)} = \kappa_{A/N} R_g$$

Model-independently using data on Pb-Pb UPCs at the LHC, [Abelev et al. \[ALICE\], PLB718 \(2013\) 1273](#); [Abbas et al. \[ALICE\], EPJ C 73 \(2013\) 2617](#); [\[CMS\] PLB 772 \(2017\) 489](#); [Acharya et al \[ALICE\], arXiv:2101:04577 \[nucl-ex\]](#)

From global QCD fits of nPDFs or leading twist nuclear shadowing model

[Guzey, Kryshen, Strikman, Zhavoronkov, PLB 726 \(2013\) 290](#)  
[Guzey, Zhavoronkov, JHEP 1310 \(2013\) 207](#)



LTA: [Frankfurt, Guzey, Strikman, Phys. Rept. 512 \(2012\) 255](#)

EPS09: [Eskola, Paukkunen, Salgado, JHEP 0904 \(2009\) 065](#)

HKN07: [Hirai, Kumano, Nagai, PRC 76 \(2007\) 065207](#)

nDS: [de Florian, Sassot, PRD 69 \(2004\) 074028](#)

- Good agreement with ALICE data at  $y=0$  (2.76 and 5.02 TeV)  $\rightarrow$  **direct evidence of large gluon shadowing**,  $R_g(x=6 \times 10^{-4} - 0.001) \approx 0.6$   $\rightarrow$  nicely agrees with LTA model and EPS09, EPPS16 nuclear parton distribution functions (nPDFs).

# Exclusive $J/\psi$ photoproduction in NLO pQCD

- Collinear factorization for hard exclusive processes, [Collins, Frankfurt, Strikman, PRD 56 \(1997\) 2982](#):  $\gamma A \rightarrow J/\psi A$  amplitude in terms of generalized parton distribution functions (GPDs), [Ji, PRD 55 \(1997\) 7114](#); [Radyushkin PRD 56 \(1997\) 5524](#); [Diehl, Phys. Rept. 388 \(2003\) 41](#)

- To next-to-leading order (NLO) of perturbative QCD, [Ivanov, Schafer, Szymanowski, Krasnikov, EPJ C 34 \(2004\) 297, 75 \(2015\) 75 \(Erratum\)](#); [Jones, Martin, Ryskin, Teubner, J. Phys. G: Nucl. Part. Phys. 43 \(2016\) 035002](#)

$$\mathcal{M}^{\gamma A \rightarrow J/\psi A} \propto \sqrt{\langle O_1 \rangle_{J/\psi}} \int_{-1}^1 dx \left[ T_g(x, \xi) F_A^g(x, \xi, t, \mu_F) + T_q(x, \xi) F_A^q(x, \xi, t, \mu_F) \right]$$

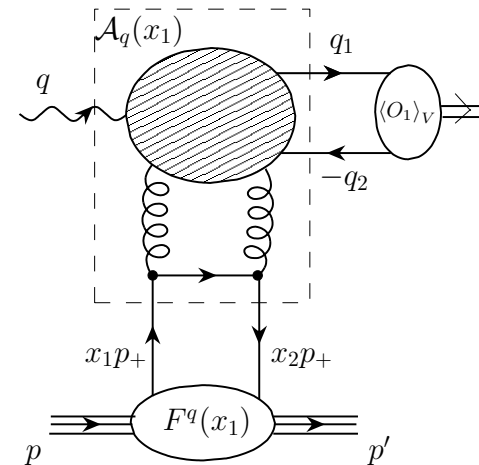
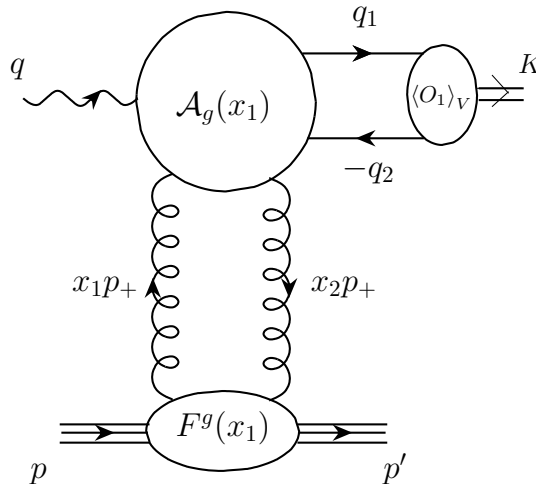
NRQCD matrix element from  $J/\psi$  leptonic decay

pQCD coefficient function

Gluon GPD

Quark contribution

- To leading order (LO), only gluons; both quarks and gluons at NLO.



# Exclusive $J/\psi$ photoproduction in NLO pQCD (2)

- In the limit of **high  $W$**  corresponding to **small  $\xi=(1/2)(M_{J/\psi})^2/W^2 \ll 1$**

$$\begin{aligned} \mathcal{M}^{\gamma A \rightarrow J/\psi A} \propto & i\sqrt{\langle O_1 \rangle_{J/\psi}} \left[ F_A^g(\xi, \xi, t, \mu_F) + \frac{\alpha_s N_c}{\pi} \ln \left( \frac{m_c^2}{\mu_F^2} \right) \int_{\xi}^1 \frac{dx}{x} F^g(x, \xi, t) \right. \\ & \left. + \frac{\alpha_s C_F}{\pi} \ln \left( \frac{m_c^2}{\mu_F^2} \right) \int_{\xi}^1 dx (F^{q,S}(x, \xi, t) - F^{q,S}(-x, \xi, t)) \right] \end{aligned}$$

→ helps to qualitatively understand the features of our numerical calculations.

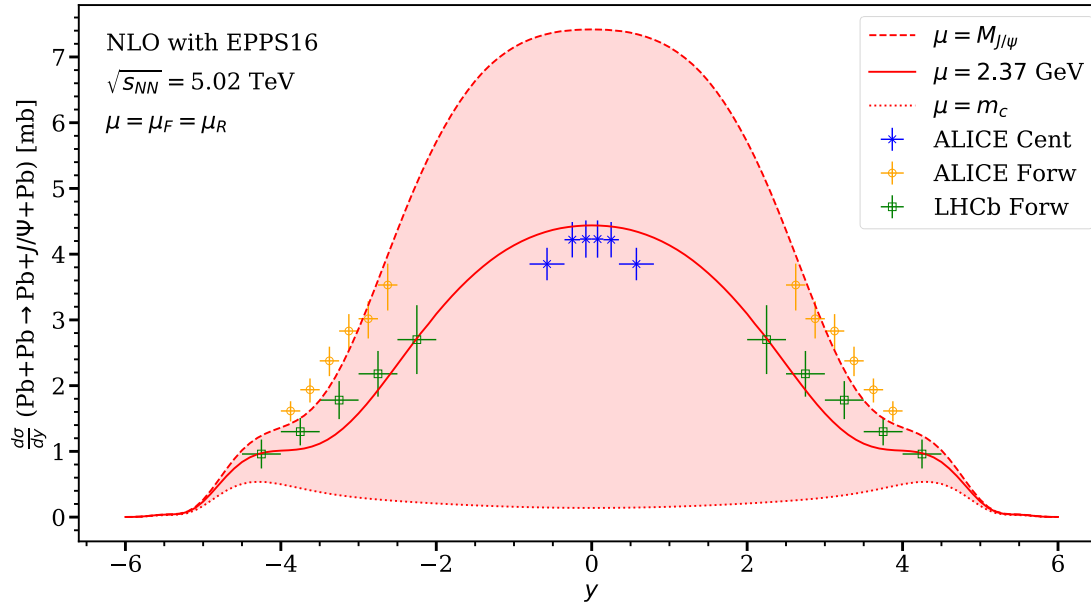
- GPDs are hybrid distributions interpolating between **usual PDFs** and **form factors** → depend on momentum fractions  **$x$**  and  **$\xi$**  and momentum transfer  **$t$** .
- Connection between GPDs is necessarily model-dependent. In our analysis, we neglect dependence of GPDs on  **$\xi$**  and used the **forward model**, Freund, McDermott, Strikman, PRD 67 (2003) 036001. For gluons (quarks are similar):

$$F_A^g(x, \xi, t, \mu_F) = x g_A(x, \mu_F) F_A(t)$$

Nuclear PDFs: EPPS16, nCTEQ15,  
nNNPDF2.0 + update with EPPS21,  
nCTEQ15WZSIH, nNNPDF3.0

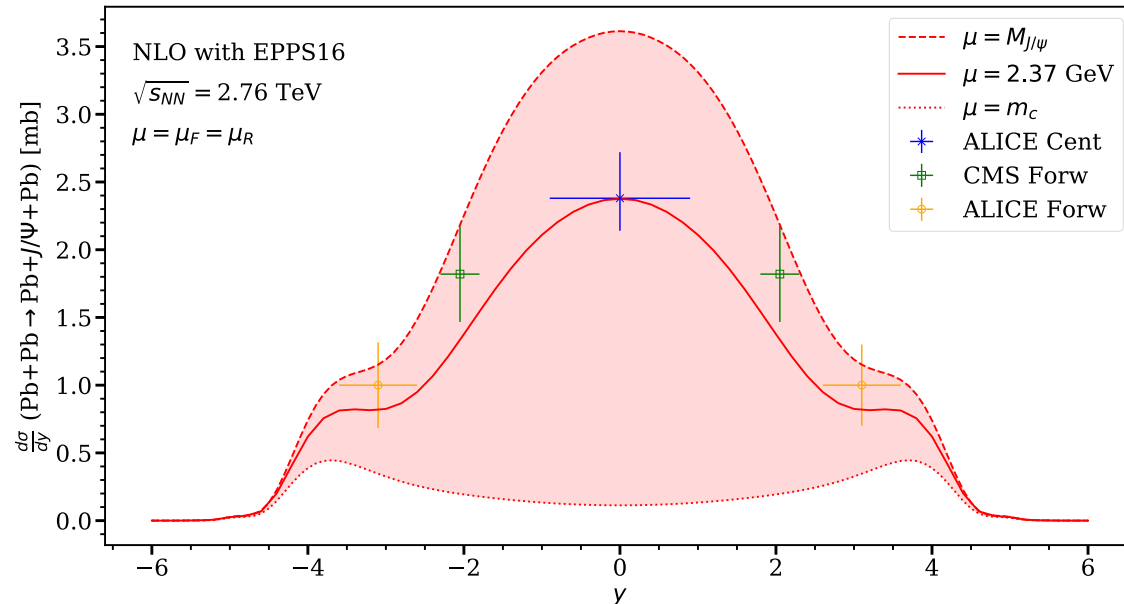
Nucleus form factor  
(Woods-Saxon form)

# Scale dependence and comparison to data on $J/\psi$ photoproduction in Pb-Pb UPCs (Runs 1&2)



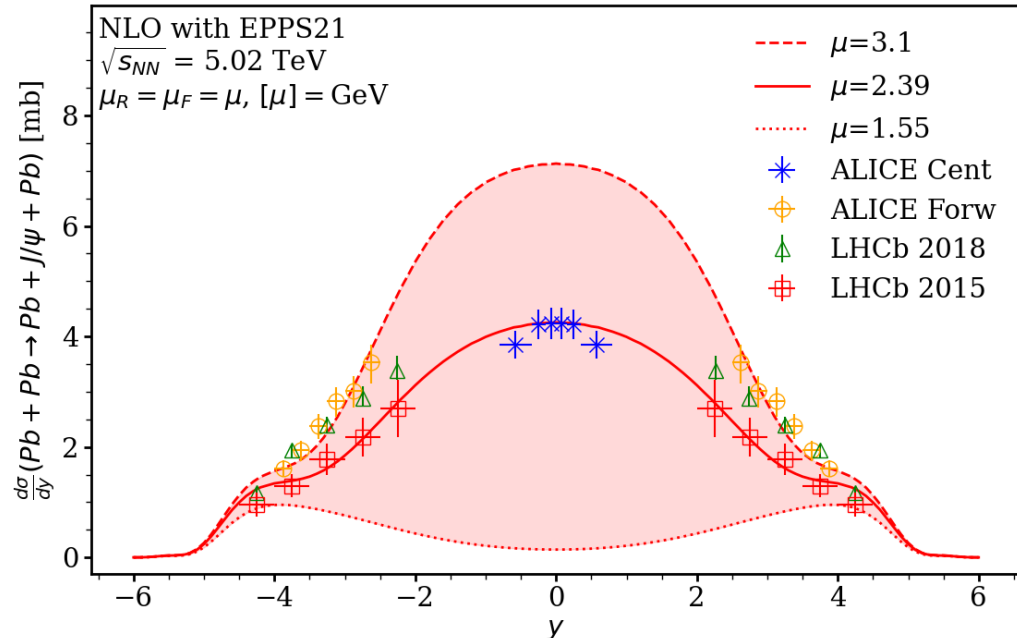
- Scale dependence of our NLO pQCD results for  $m_c \leq \mu_F \leq M_{J/\psi}$  is very strong.

- One can find an “optimal scale” giving simultaneous good description of Run 1&2 UPC data.



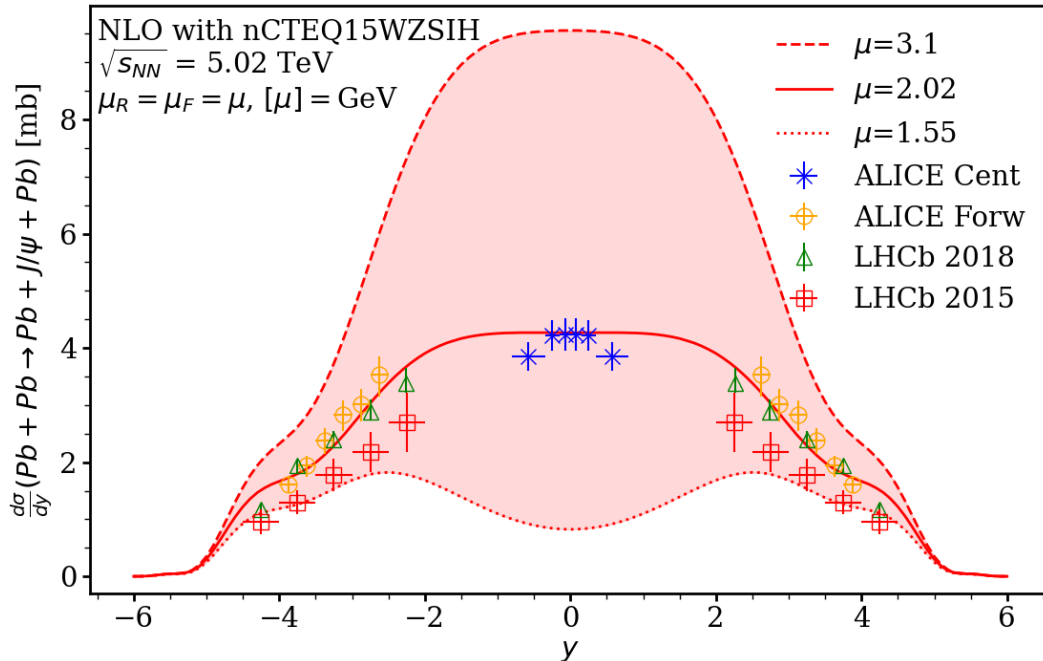
- With this choice of scale, the  $\gamma+p \rightarrow J/\psi+p$  proton data is somewhat overestimated, but within large scale uncertainties.

# Scale dependence and comparison to data on $J/\psi$ photoproduction in Pb-Pb UPCs: update



- Repeated our calculations using state-of-the-art EPPS21, nNNPDF3.0 and nCTEQ15WZSIH nPDFs.

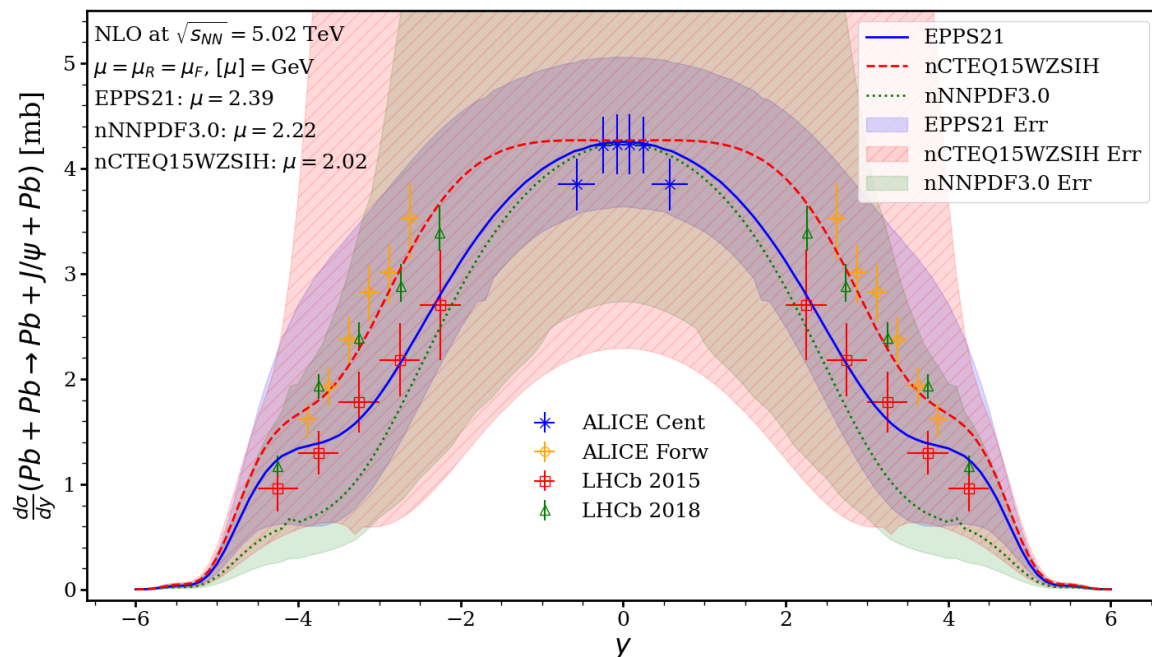
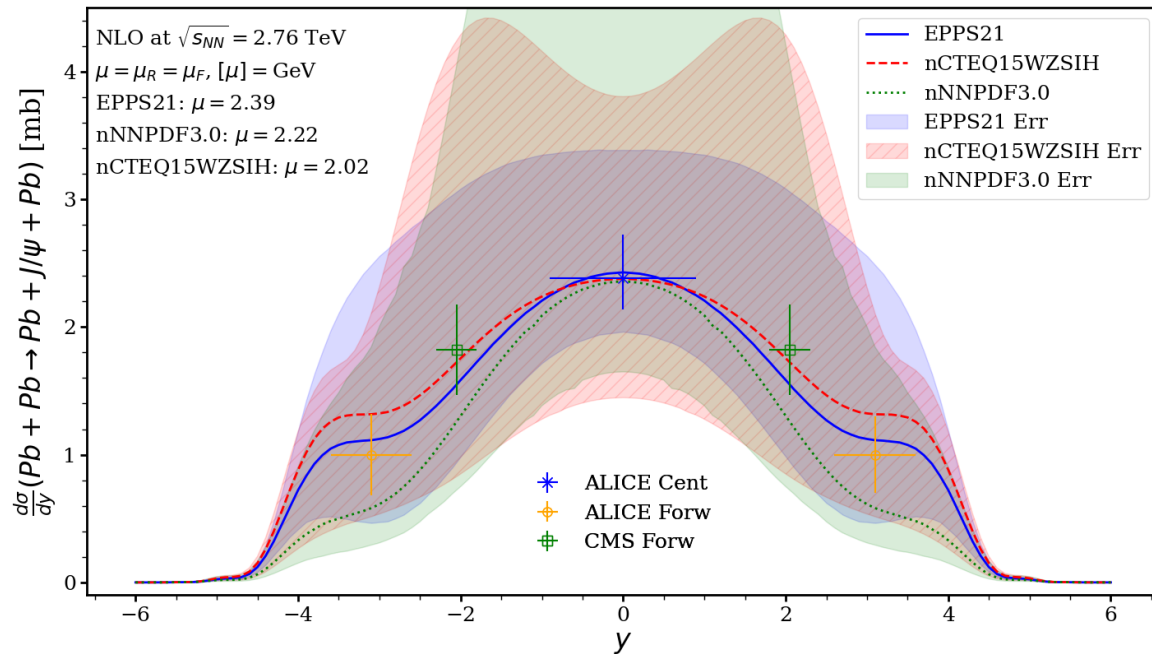
- Note that updated LHCb data have moved up worsening the agreement with **EPPS21**.



- However, the agreement is restored by using **nCTEQ15WZSIH** nPDFs, which are characterized by large strange quark density  $\rightarrow$  **sensitivity to strange quarks in nuclei?**



# Uncertainties due to nuclear PDFs



- Uncertainties due to nPDFs are quite significant → **opportunity to reduce** them using the data on  $J/\psi$  photoproduction in AA UPCs.

- Compared to our original calculations, abnormally large uncertainty associated with **EPPS16** disappears when using more recent **EPPS21**.

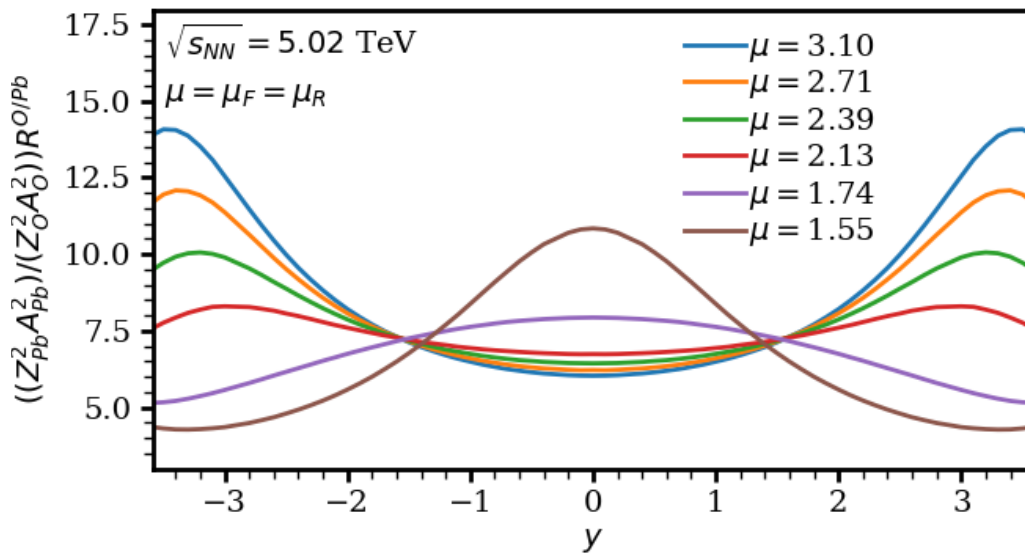
- The **nNNPDF3.0** nPDFs correspond to much less constrained fit → large uncertainties.

# Reduction of uncertainties using O/Pb ratio

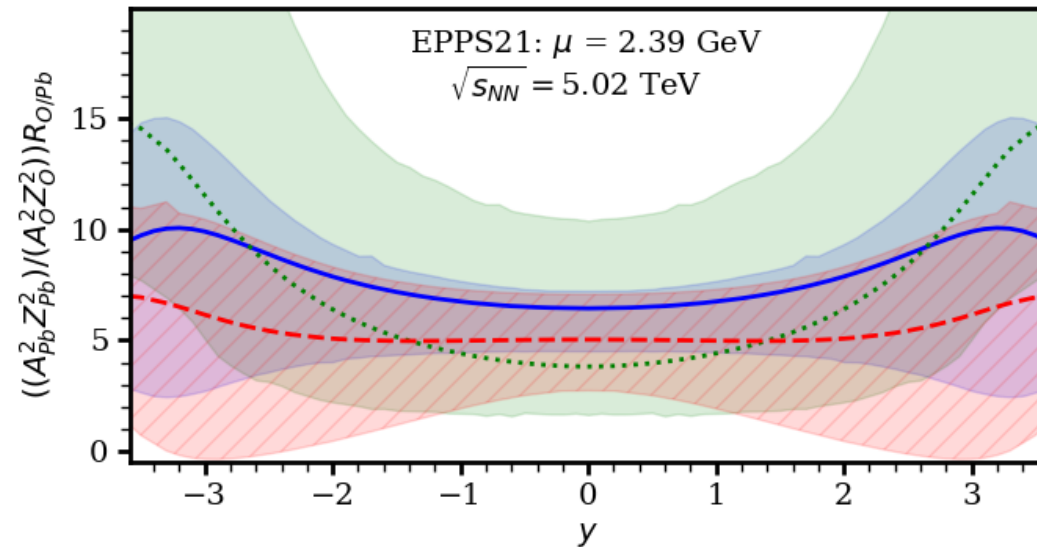
- One can reduce the significant scale  $\mu_F$  and nPDF uncertainties by considering the ratio of oxygen to lead UPC cross sections:

$$R^{O/Pb} = \left( \frac{208Z_{Pb}}{16Z_O} \right)^2 \frac{d\sigma(O + O \rightarrow O + J/\psi + O)/dy}{d\sigma(Pb + Pb \rightarrow Pb + J/\psi + Pb)/dy}$$

Scale uncertainty of R



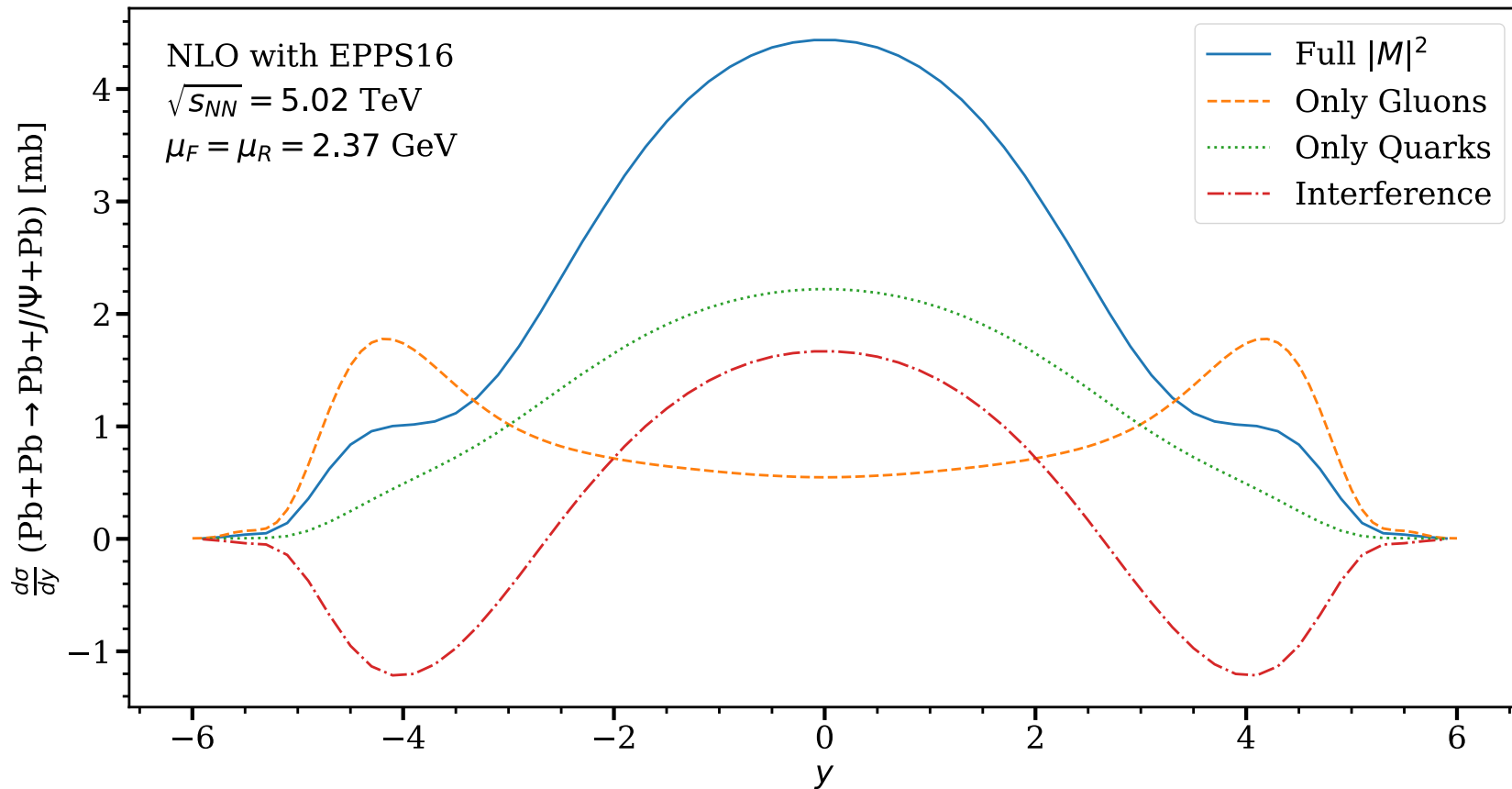
nPDF uncertainty of R



- Hard scattering coefficient functions for O and Pb are the same  $\rightarrow$  the **scale dependence** comes from nPDFs  $\rightarrow$  **reduced by factor of 10** compared to individual UPC cross sections.
- Reduction of **nPDF uncertainties** is even larger due to additional partial cancellation of uncertainties associated with proton PDFs.

# Surprise: dominance of quark contribution

- The most striking result is strong cancellations between LO and NLO gluons → **dominance of quark contribution** at central rapidities.



- At the face value, **this totally changes** the interpretation of data on coherent  $J/\psi$  photoproduction in heavy-ion UPCs as a probe of small-x nuclear gluons → **but this requires overcoming certain theoretical challenges.**

# Exclusive $J/\psi$ photoproduction in NLO pQCD.

## Challenge 1: small-x resummation

- NLO corrections and, hence, the scale dependence are very large  $\rightarrow$  large theoretical uncertainties in phenomenological applications.
- The reason is well understood  $\rightarrow$  large  $\ln(Q^2) \ln(1/\xi)$  terms for  $2\xi \approx (M_{J/\psi})^2/W^2 \ll 1$

$$\mathcal{M}^{\gamma A \rightarrow J/\psi A} \propto i \sqrt{\langle O_1 \rangle_{J/\psi}} \left[ F_A^g(\xi, \xi, t, \mu_F) + \frac{\alpha_s N_c}{\pi} \ln \left( \frac{m_c^2}{\mu_F^2} \right) \int_{\xi}^1 \frac{dx}{x} F^g(x, \xi, t) \right. \\ \left. + \frac{\alpha_s C_F}{\pi} \ln \left( \frac{m_c^2}{\mu_F^2} \right) \int_{\xi}^1 dx (F^{q,S}(x, \xi, t) - F^{q,S}(-x, \xi, t)) \right]$$

- Possible solution: so-called  $Q_0$  subtraction based on subtraction of  $k_T < Q_0 \sim m_c$  contribution from NLO coefficient functions, [Jones, Martin, Ryskin, Teubner, EPJC 76 \(2016\) 11, 633](#); [Flett, Jones, Martin, Ryskin, Teubner, PRD 101 \(2020\) 9, 094011](#)
- Other (related) option to tame small-x behavior: BFKL resummation, [Ivanov, arXiv:0712.31983 \[hep-ph\]](#); [Ivanov, Pire, Szymanowski, Wagner, EPJ Web. Conf. 112 \(2016\) 01020](#)

$$\text{Im} \mathcal{M}^g \sim H^g(\xi, \xi) + \int_{\xi}^1 \frac{dx}{x} H^g(x, \xi) \sum_{n=1}^L C_n(L) \frac{\bar{\alpha}_s^n}{(n-1)!} \log^{n-1} \frac{x}{\xi} \quad \text{where} \quad L = \ln(M_V^2/\mu_F^2)$$

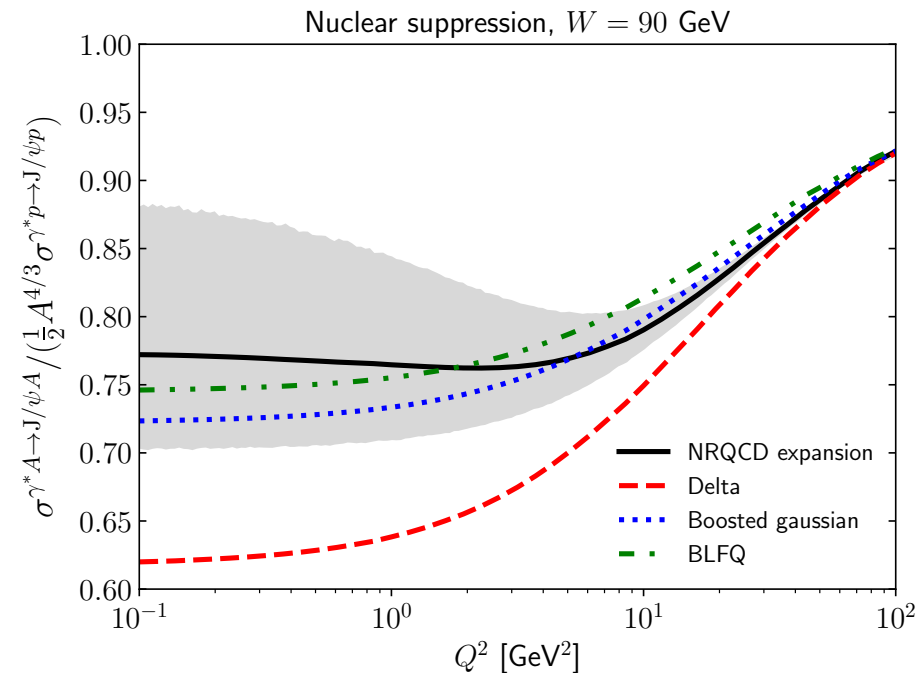
 Need to calculate  $C_n(L)$  and include in our NLO pQCD analysis.

# Exclusive $J/\psi$ photoproduction in NLO pQCD.

## Challenge 2: non-relativistic corrections to charmonium wf

- Our analysis assumes a static (non-relativistic) limit for  $J/\psi$  vertex.

- Recent analyses have shown that relativistic  $v/m_c$  corrections are sizable, [Eskobedo, Lappi, PRD 101 \(2020\) 3, 034030](#); [Lappi, Mantysaari, Penttala, PRD 102 \(2020\) 5, 054020](#)  
→ affect interpretation of nuclear suppression in AA UPCs@LHC.



- There is also a related issue of D-wave (spin rotation) of the charmonium wave function, [Krelina, Nemchik, Pasechnik, EPJ C 80 \(2020\) 2, 92](#)

Need to couple non-relativistic  $v/m_c$  corrections to our NLO pQCD calculations.

# Summary

- First NLO pQCD calculation of exclusive  $J/\psi$  photoproduction in Pb-Pb and O-O UPCs@LHC in the framework of collinear factorization.
- Our analysis confirmed strong scale dependence noticed earlier, quantified uncertainty due to nuclear PDFs, observed the dominance of the quark contribution, and provided simultaneous description of Run 1&2 LHC data.
- From phenomenology point of view, the ultimate goal is to use these data to obtain new information on nuclear PDFs at small  $x$ , e.g., by using the UPC data in global QCD fits.
- In the present form, this is challenging. Possible solutions:
  - ❖ Consider ratio of AA to OO/pp UPC cross sections, where most of complications (scale dependence, uncertainties of nPDFs, details of GPD modeling, relativistic corrections to the charmonium wave function) partially cancel.
  - ❖ Small- $x$  BFKL resummation and  $Q_0$  subtraction to tame the large scale dependence.
  - ❖ Even in the case of the UPC cross section ratios, nonrelativistic corrections to charmonium wf do not cancel exactly and should be taken into account.