



Indications for a Higgs Bosons below 125 GeV

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- Motivation
- Examples for additional Higgs Bosons
- A Higgs Boson at 96 GeV?!
- Conclusions

1. Motivation: Two Facts:

1: We have a discovery!

2: The SM cannot be the ultimate theory!

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Q': Which model?

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Q: Does the BSM physics have any (relevant) impact on the Higgs?

Q': Which model?

A1: check **changed properties**

A2: check for **additional Higgs bosons**

A2': check for additional Higgs bosons **above** and **below 125 GeV**

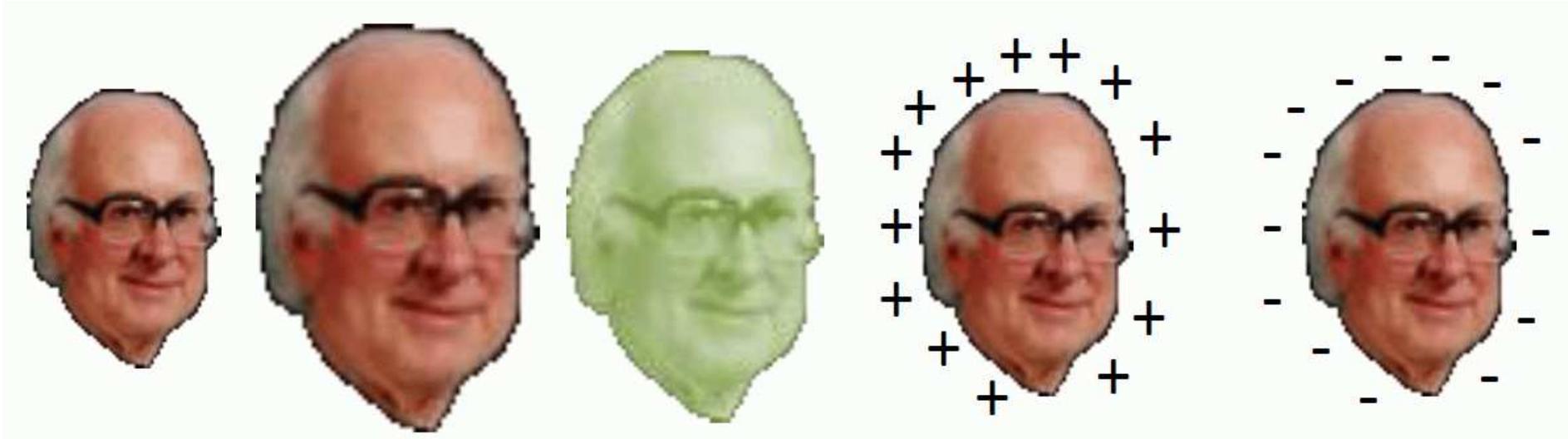
Models with extended Higgs sectors:

1. SM with additional Higgs singlet
 2. Two Higgs Doublet Model (THDM): type I, II, III, IV
 3. N2HDM: 2HDM with one extra singlet: type I, II, III, IV
 4. Minimal Supersymmetric Standard Model (MSSM)
 5. MSSM with one extra singlet (NMSSM)
 6. MSSM with more extra singlets (e.g. $\mu\nu$ SSM)
 7. SM/MSSM with Higgs triplets
 8. . . .
- ⇒ BSM models without extended Higgs sectors still have changed Higgs properties (quantum corrections!)
- ⇒ SM + vector-like fermions, Higgs portal, Higgs-radion mixing, . . .

Models with extended Higgs sectors:

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 4. Minimal Supersymmetric Standard Model (MSSM) ⇐ focus
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- ⇒ BSM models without extended Higgs sectors still have changed Higgs properties (quantum corrections!)
- ⇒ SM + vector-like fermions, Higgs portal, Higgs-radion mixing, ...

2. Examples for additional Higgs Bosons



Search for the MSSM Higgs bosons:

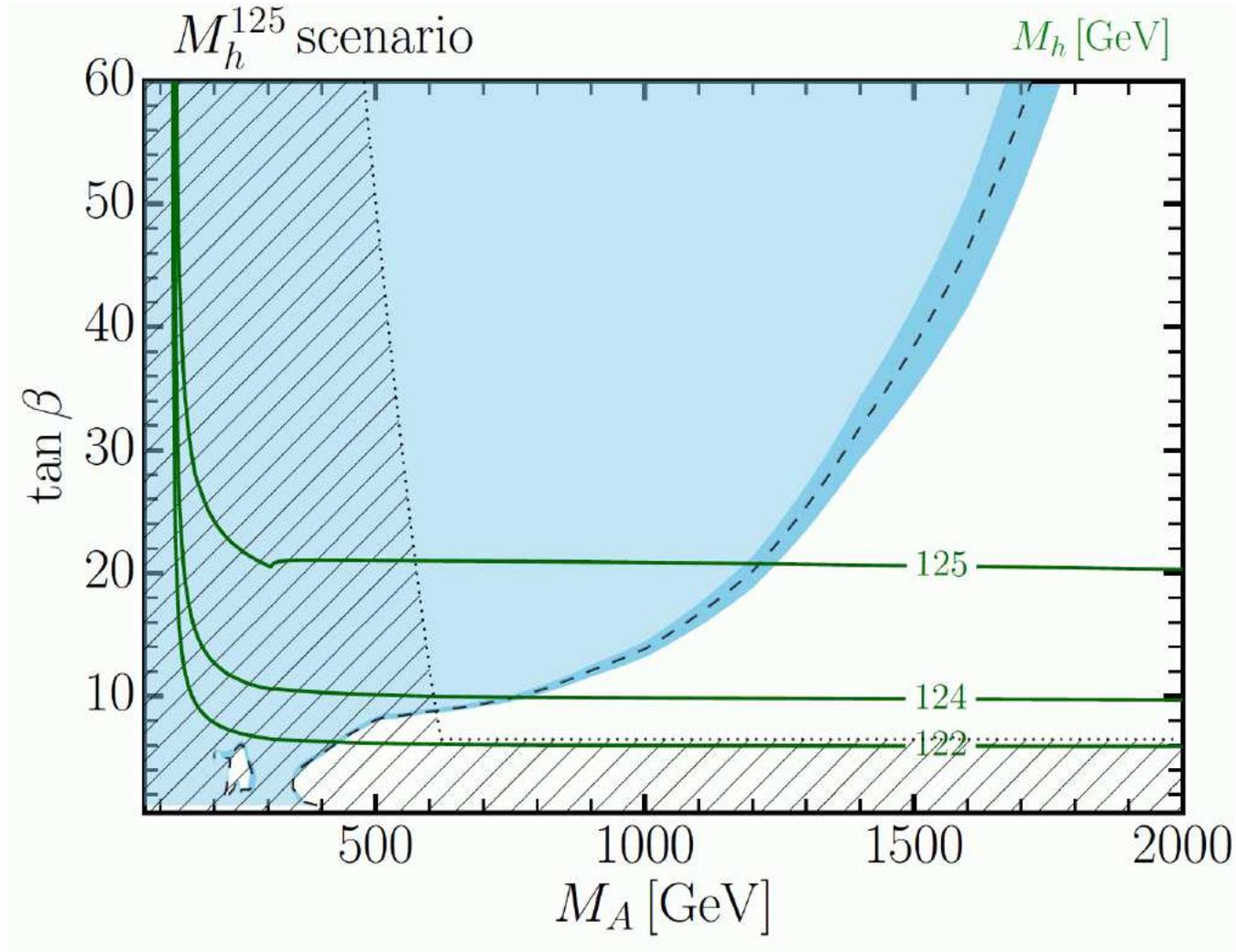
Smart choice of MSSM parameters?

→ investigate benchmark scenarios:

→ Vary only M_A and $\tan \beta$
→ Keep all other SUSY parameters fixed

[E. Bagnaschi, H. Bahl, E. Fuchs, T. Hahn, S.H., S. Liebler, S. Patel, P. Slavich, T. Stefaniak, C. Wagner, G. Weiglein '18]

1. M_h^{125} scenario: 2HDM-like model
2. $M_h^{125}(\tilde{\tau})$ scenario: light staus: $h \rightarrow \gamma\gamma$, $H/A \rightarrow \tilde{\tau}\tilde{\tau}$
3. $M_h^{125}(\tilde{\chi})$ scenario: light EW-inos: $H/A \rightarrow \tilde{\chi}_i^0 \tilde{\chi}_j^0, \tilde{\chi}_k^\pm \tilde{\chi}_l^\mp$
4. M_h^{125} (alignment) scenario: h SM-like for very low M_A
5. M_H^{125} scenario: $M_H \sim 125$ GeV, all Higgses light
6. $M_{h_1}^{125}$ (CPV) scenario: complex phases, h_2 - h_3 interference



$$M_{\tilde{Q}_3} = M_{\tilde{U}_3} = M_{\tilde{D}_3} = 1.5 \text{ TeV}$$

$$M_{\tilde{L}_3} = M_{\tilde{E}_3} = 2 \text{ TeV}$$

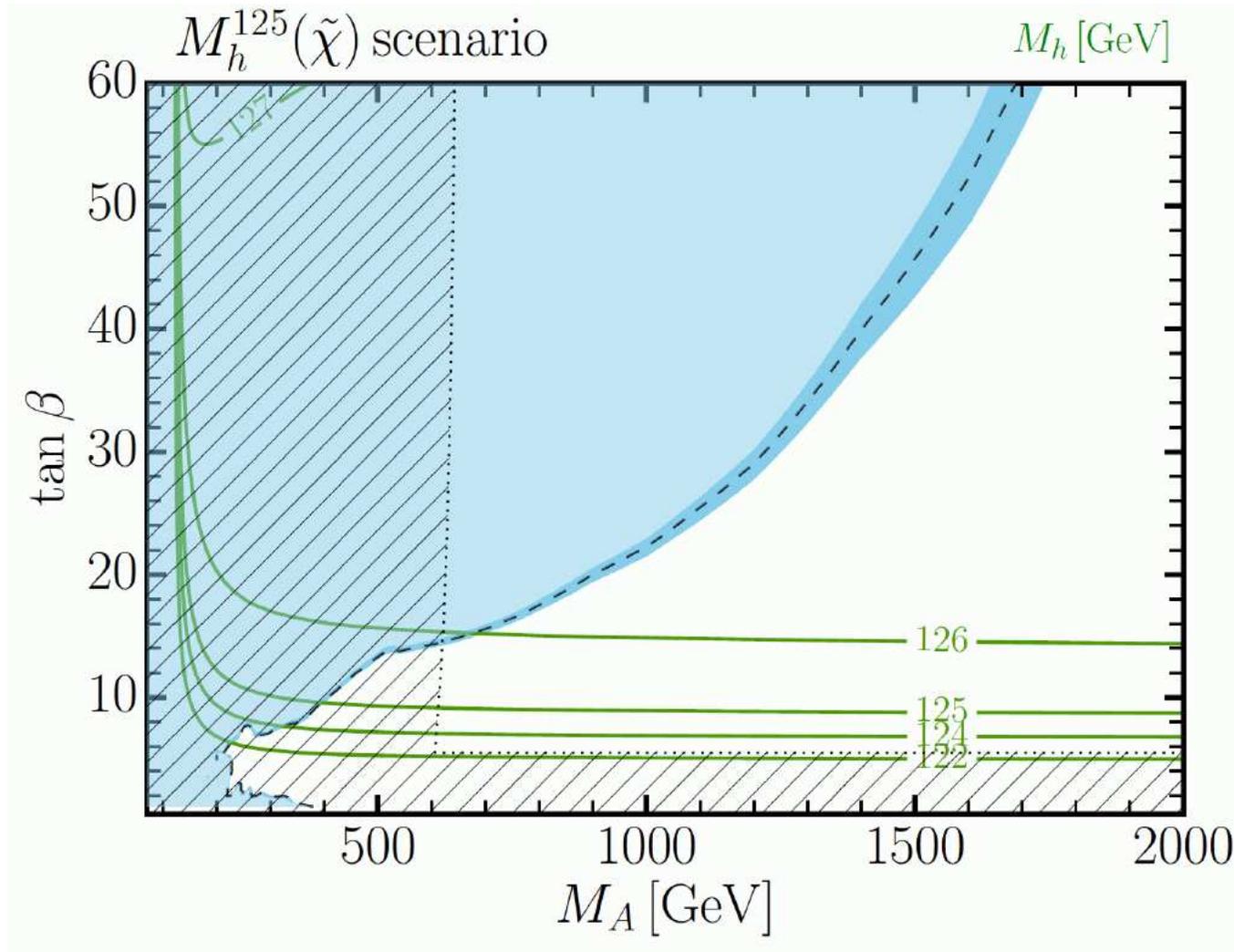
$$\mu = 1 \text{ TeV}, M_1 = 1 \text{ TeV}$$

$$M_2 = 1 \text{ TeV}, M_3 = 2.5 \text{ TeV}$$

$$X_t = 2.8 \text{ TeV}$$

$$A_t = A_b = A_\tau$$

⇒ new vanilla benchmark model



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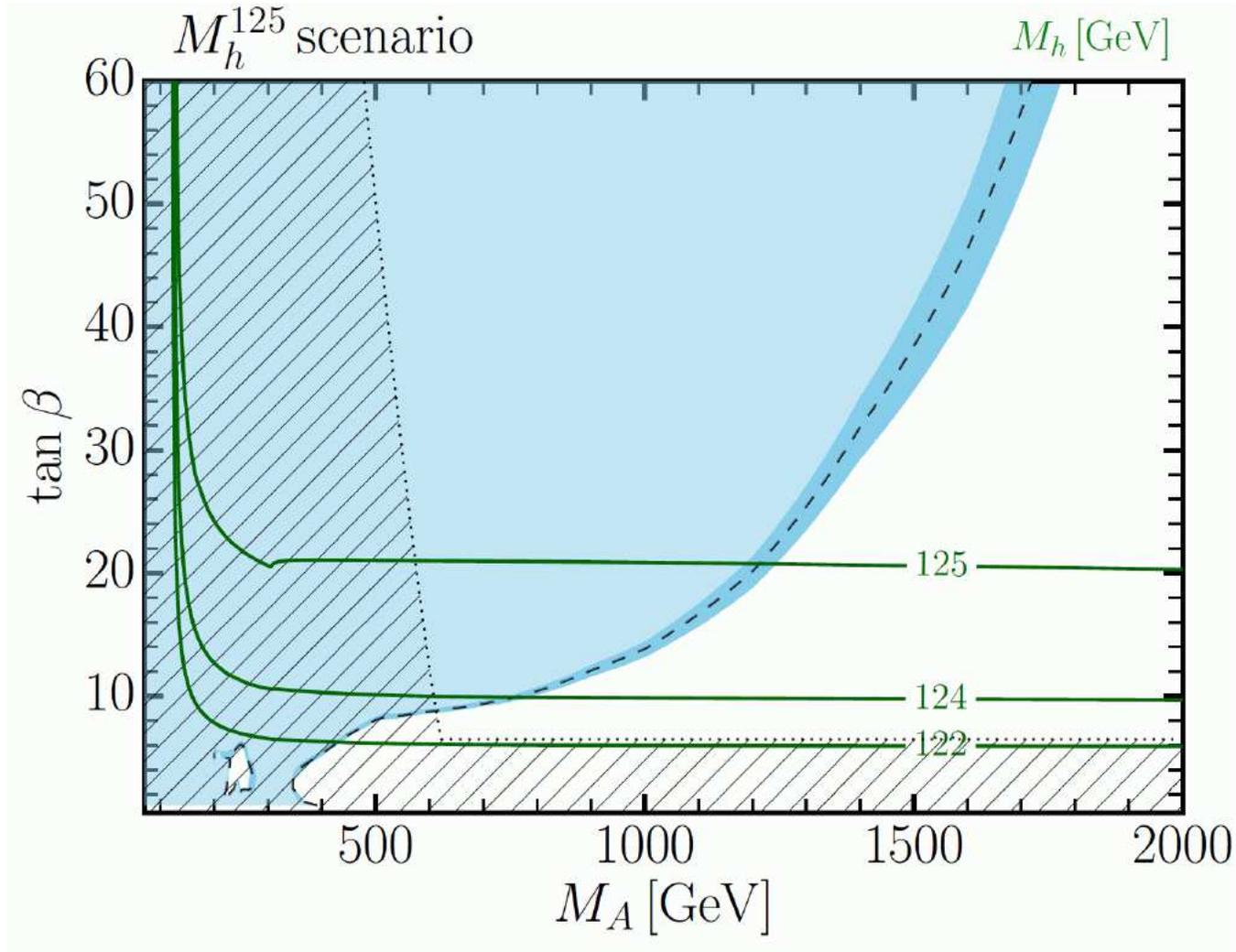
$$\mu = 180 \text{ GeV}, M_1 = 160 \text{ GeV}$$

$$M_2 = 180 \text{ GeV}, M_3 = 2.5 \text{ TeV}$$

$$X_t = 2.5 \text{ TeV}$$

$$A_t = A_b = A_\tau$$

⇒ strongly reduced heavy Higgs coverage



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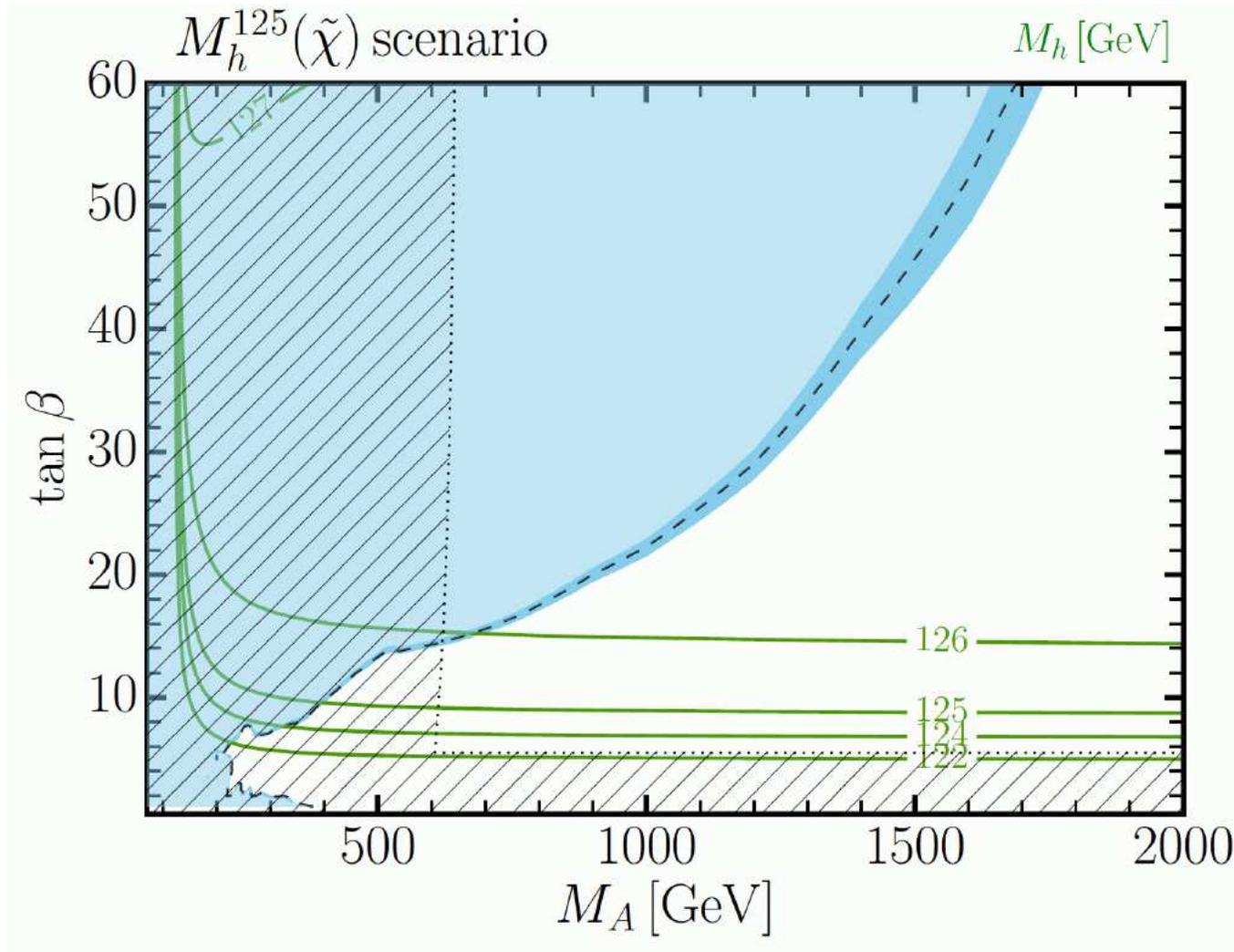
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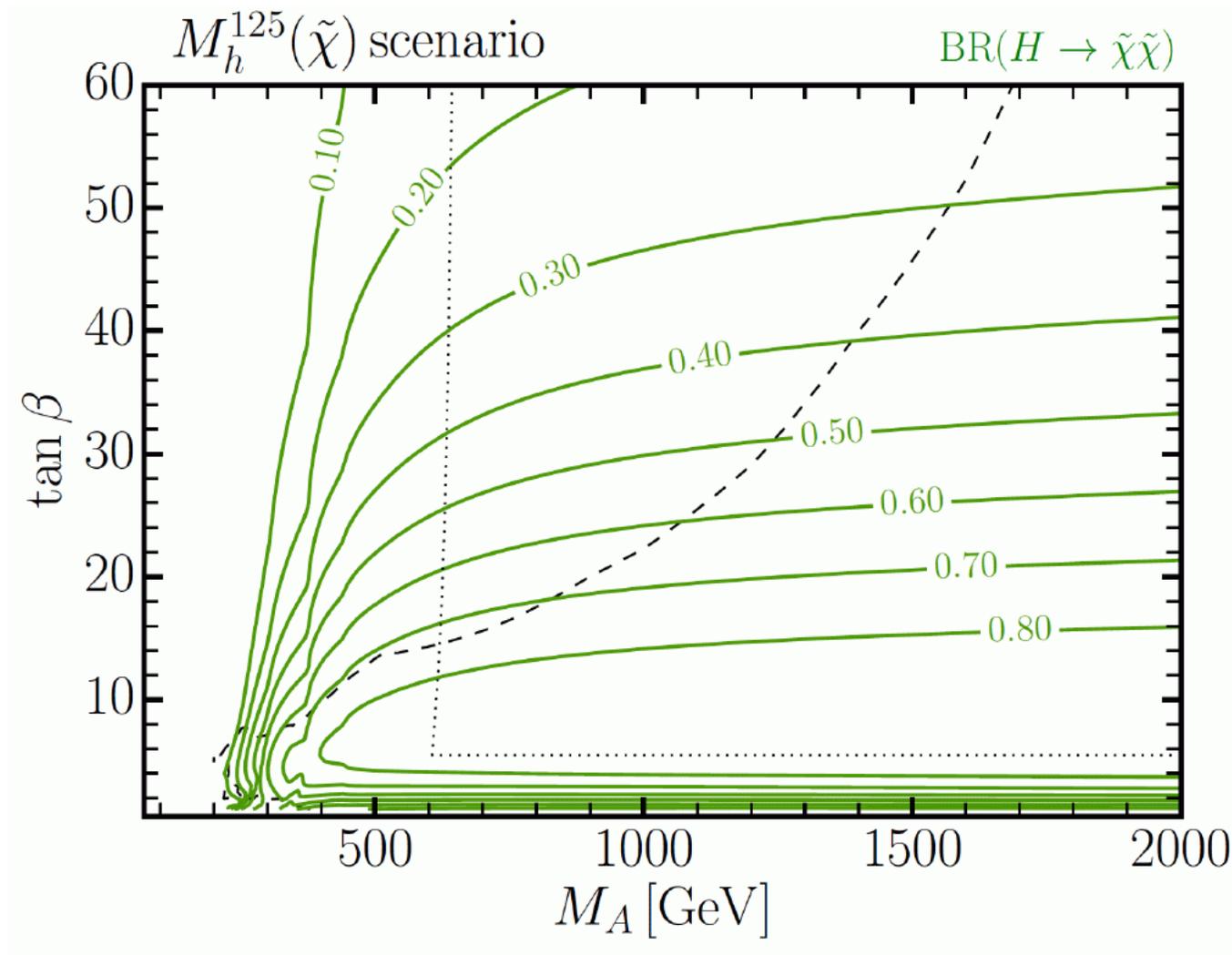
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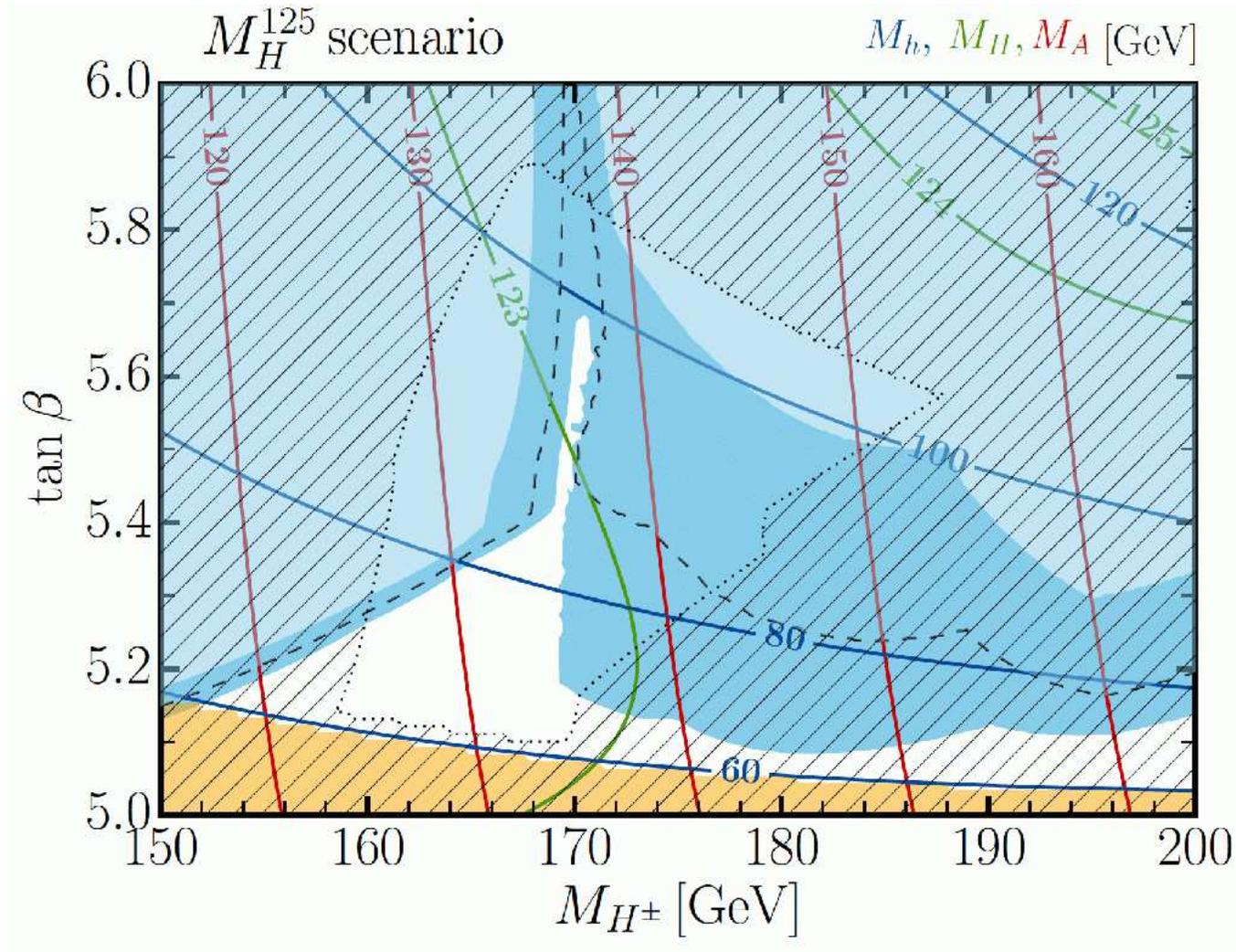
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 $\mu = 180 \text{ GeV}, M_1 = 160 \text{ GeV}$
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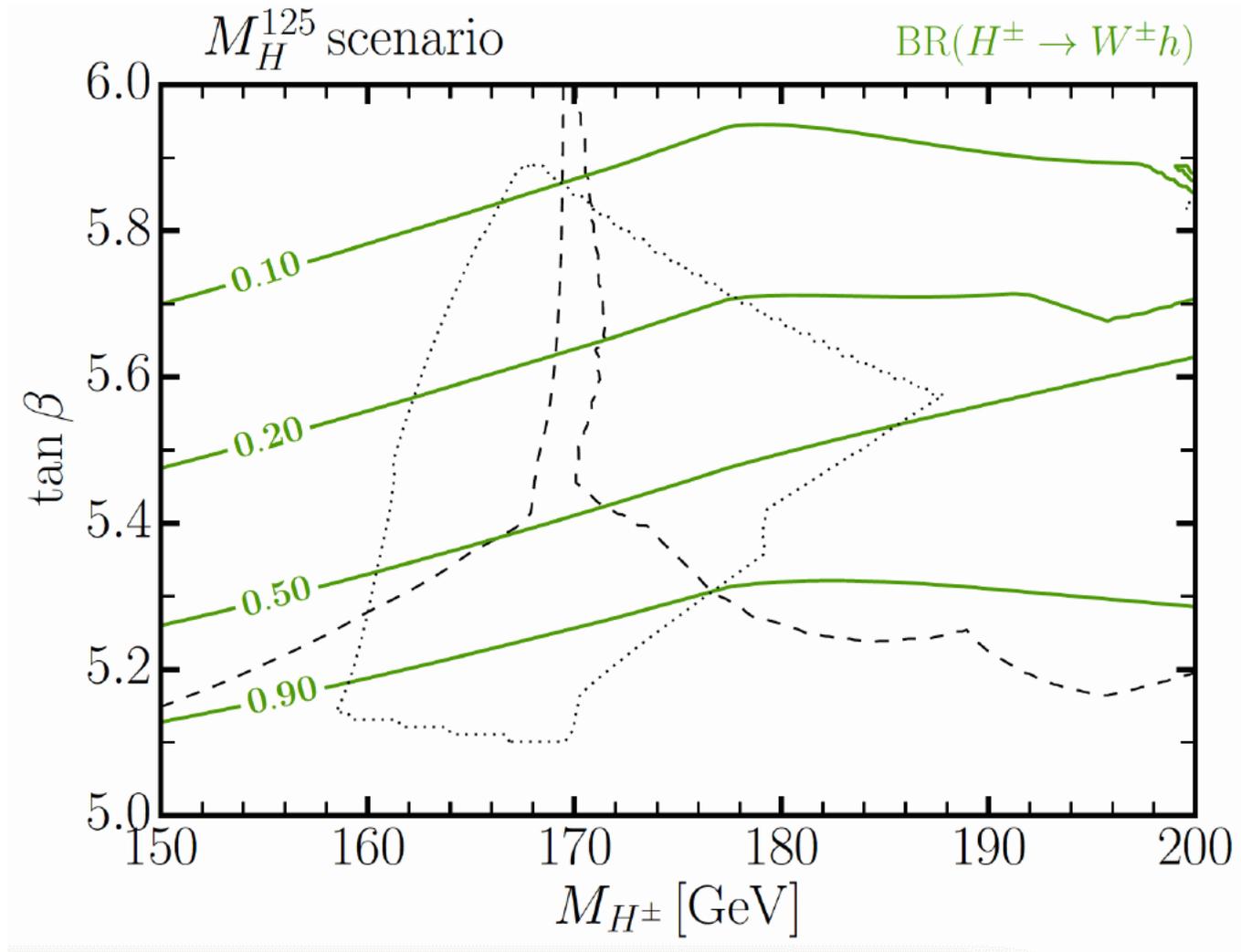
⇒ Huge BR of heavy Higgses to EW-inos



$$\begin{aligned}
 M_{\tilde{Q}_3} &= M_{\tilde{U}_3} = 750 \text{ GeV} \\
 &\quad - 2(M_{H^\pm} - 150 \text{ GeV}) \\
 M_{\tilde{L}_3} &= M_{\tilde{E}_3} = M_{\tilde{D}_3} = 2 \text{ TeV} \\
 \mu &= [5.8 \text{ TeV} \\
 &\quad + 20(M_{H^\pm} - 150 \text{ GeV})] \times \\
 &\quad M_{\tilde{Q}_3} / 750 \text{ GeV} \\
 M_1 &= M_{\tilde{Q}_3} - 75 \text{ GeV} \\
 M_2 &= 1 \text{ TeV}, M_3 = 2.5 \text{ TeV} \\
 A_t &= A_b = A_\tau = 0.65 M_{\tilde{Q}_3}
 \end{aligned}$$

⇒ exotic solution still viable!

⇒ scenario with a Higgs below 125 GeV!



$$\begin{aligned}
 M_{\tilde{Q}_3} &= M_{\tilde{U}_3} = 750 \text{ GeV} \\
 &\quad - 2(M_{H^\pm} - 150 \text{ GeV}) \\
 M_{\tilde{L}_3} &= M_{\tilde{E}_3} = M_{\tilde{D}_3} = 2 \text{ TeV} \\
 \mu &= [5.8 \text{ TeV} \\
 &\quad + 20(M_{H^\pm} - 150 \text{ GeV})] \times \\
 &\quad M_{\tilde{Q}_3} / 750 \text{ GeV} \\
 M_1 &= M_{\tilde{Q}_3} - 75 \text{ GeV} \\
 M_2 &= 1 \text{ TeV}, \quad M_3 = 2.5 \text{ TeV} \\
 A_t &= A_b = A_\tau = 0.65 M_{\tilde{Q}_3}
 \end{aligned}$$

\Rightarrow large $BR(H^\pm \rightarrow W^\pm h)$

3. A Higgs Boson at 96 GeV?!

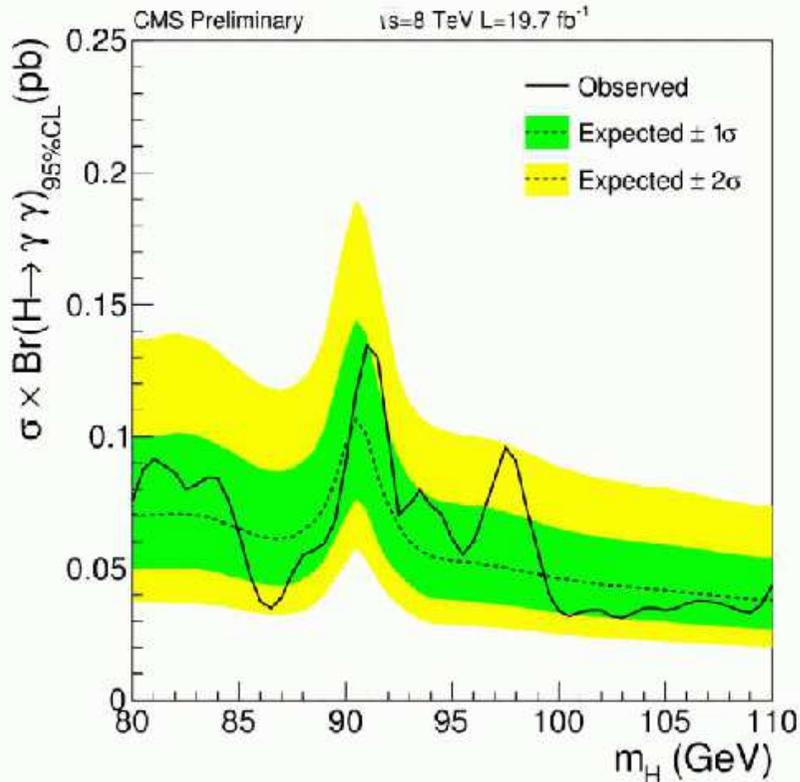
- What was seen in Run I?
- What was seen in Run II?
- What was seen at LEP?
- Should we get excited?
- Which model fits?
- Next project?!

$h \rightarrow \gamma\gamma$ (65-110 GeV) Run 1

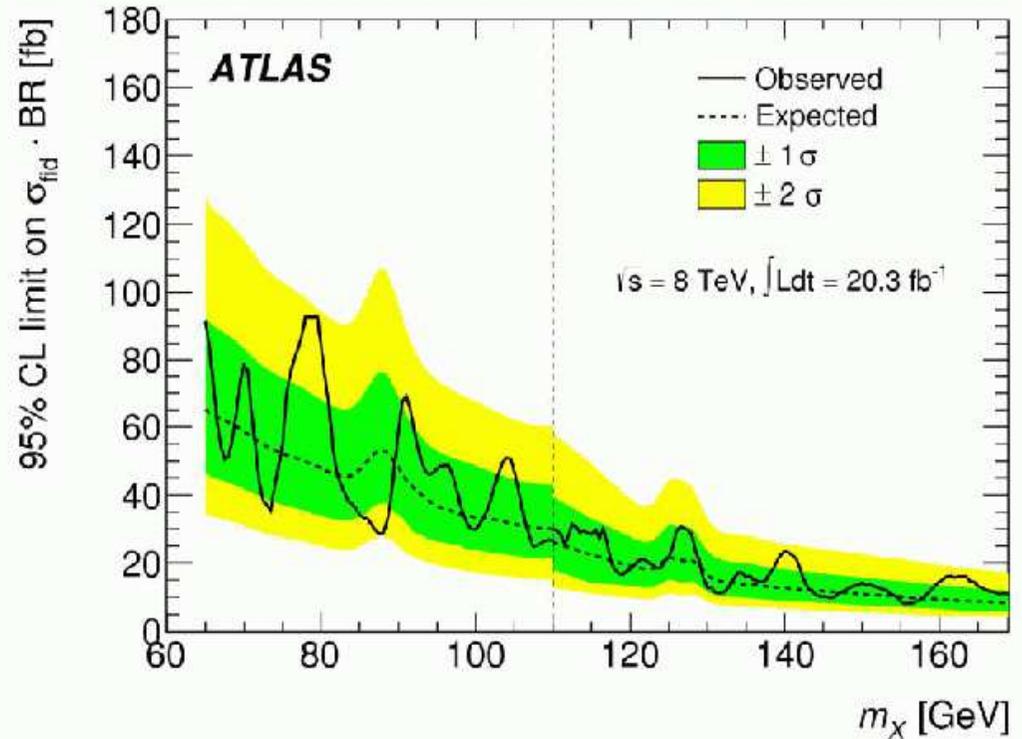


CMS PAS HIG-14-037

PRL 113 171801 (2014)



• $\sim 2\sigma$ excursion @ ~ 97.5 GeV



• $\sim 2\sigma$ excursion @ ~ 80 GeV

S. Gascon-Shotkin HDays17, Santander, ES Sept. 22 2017

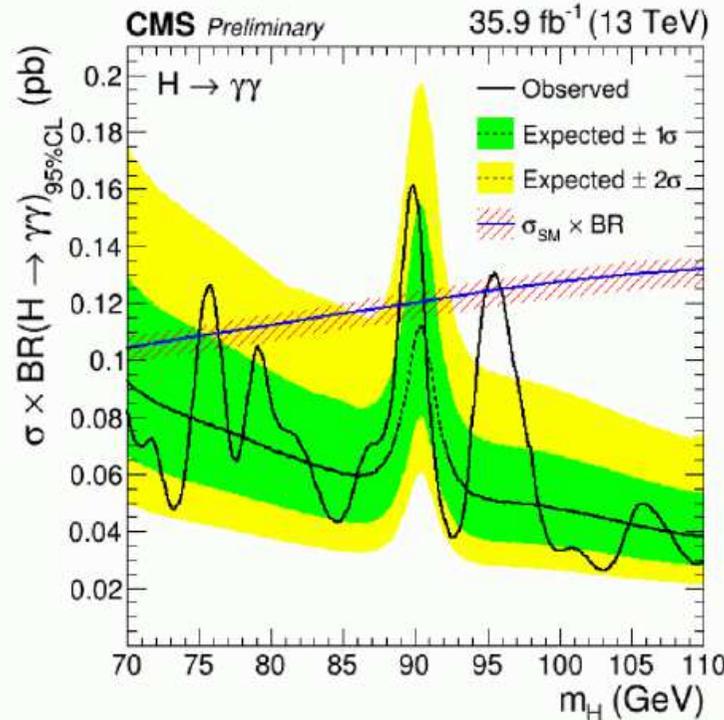
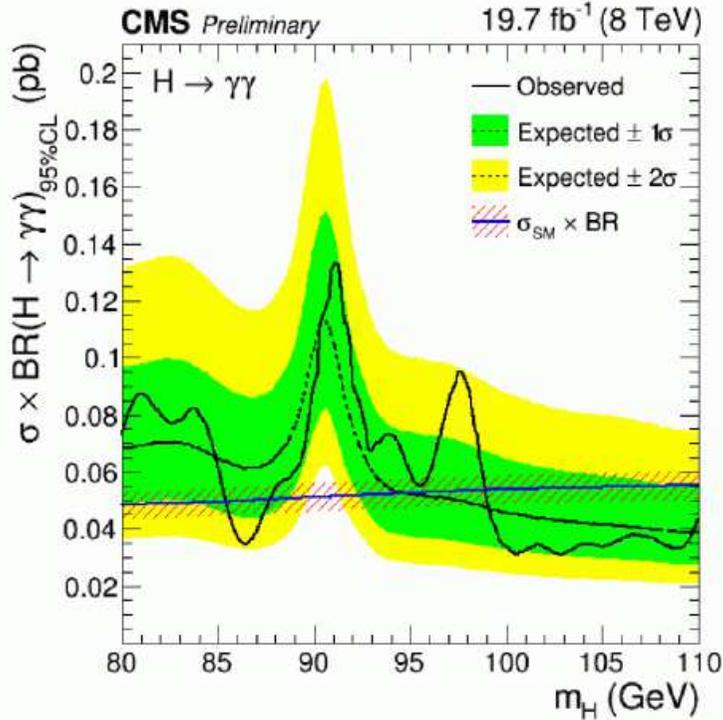
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$h \rightarrow \gamma\gamma$ (70-110 GeV) Runs 1+2



CMS PAS HIG-17-013



8 TeV:
 minimum(maximum)
 limit on $\sigma \times \text{Br}$:
 31(133) fb at
 $m=102.8(91.1)\text{GeV}$

13 TeV:
 minimum(maximum)
 limit on $\sigma \times \text{Br}$:
 26(161) fb at
 $m=103.0(89.9)\text{GeV}$

- 8 TeV limits on $\sigma \times \text{Br}$ redone with 0.1 GeV step. Production processes assumed in SM proportions. No significant excess with respect to expected limits observed.

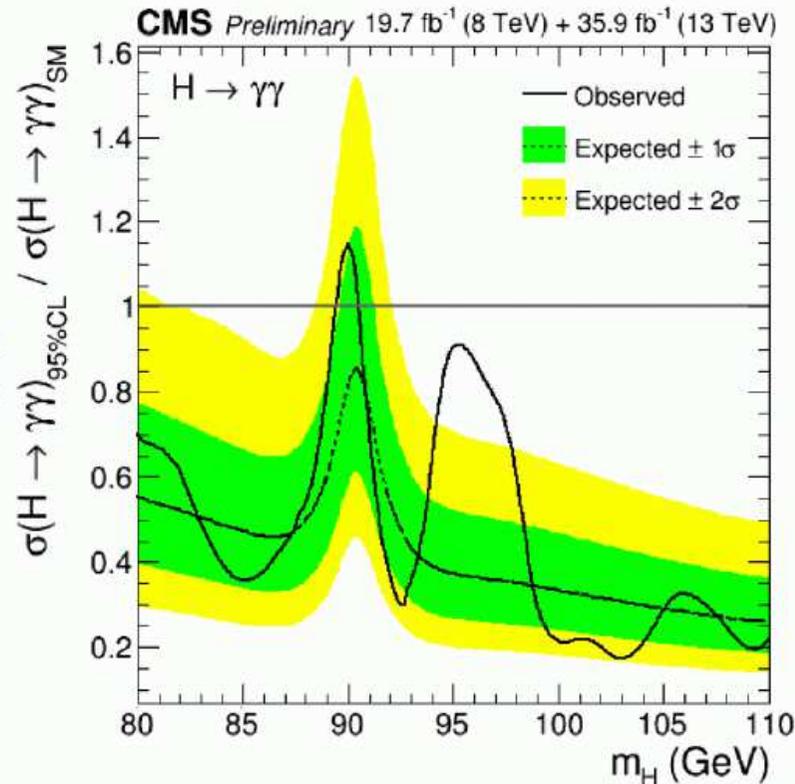


$h \rightarrow \gamma\gamma$ (70-110 GeV) Runs 1+ 2



CMS PAS HIG-17-013

All experimental + theoretical systematic uncertainties assumed uncorrelated except for those on signal acceptance due to scale variations + those on production cross sections (assumed 100% correlated).



8 TeV+13 TeV:
 minimum(maximum) limit
 on $(\sigma \times Br) / (\sigma \times Br)_{SM}$:
 0.17(1.15) at
 $m=103.0(90.0)\text{GeV}$

- Combined 8 TeV+13 TeV $\sigma \times Br$ limit normalized to SM expectation (production processes assumed in SM proportions). No significant excess with respect to expected limits observed.

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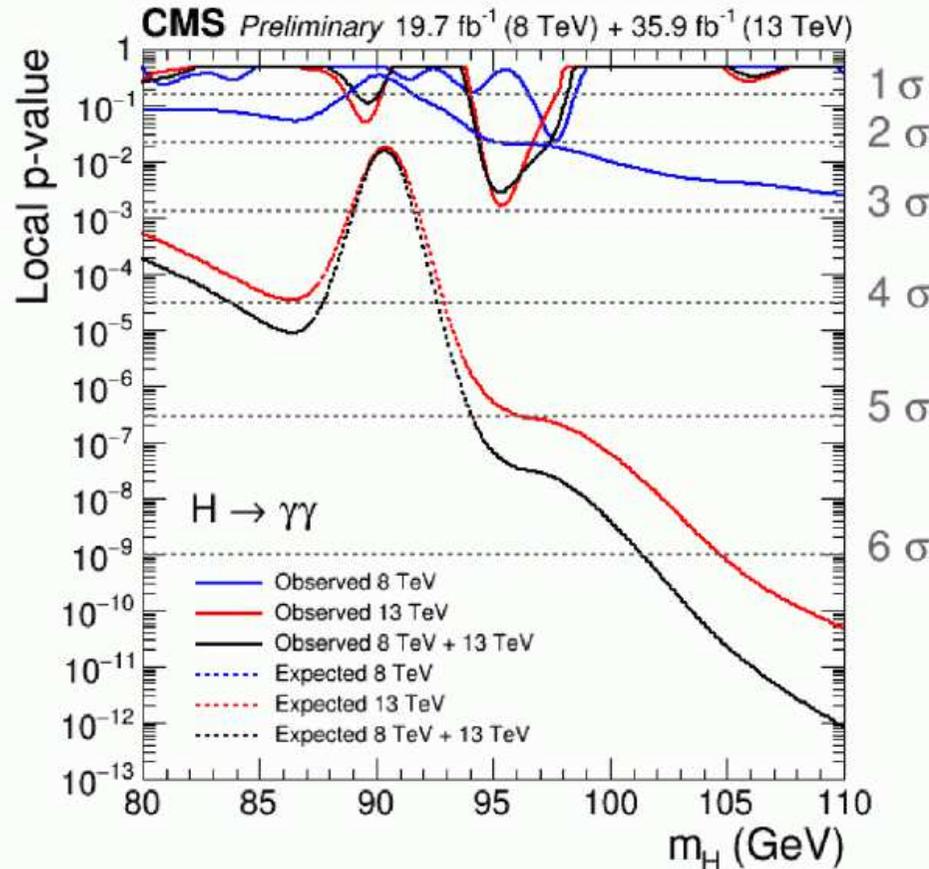
S. Gascon-Shotkin HDays17, Santander, ES Sept. 22 2017



$h \rightarrow \gamma\gamma$ (70-110 GeV) Runs 1+ 2



CMS PAS HIG-17-013



8 TeV: Excess with $\sim 2.0 \sigma$ local significance at $m=97.6$ GeV

13 TeV: Excess with $\sim 2.9 \sigma$ local (1.47σ global) significance at $m=95.3$ GeV

8TeV+13 TeV: Excess with $\sim 2.8 \sigma$ local (1.3σ global) significance at $m=95.3$ GeV

More data are required to ascertain the origin of this excess

- Expected and observed local p-values for 8 TeV, 13 TeV and their combination

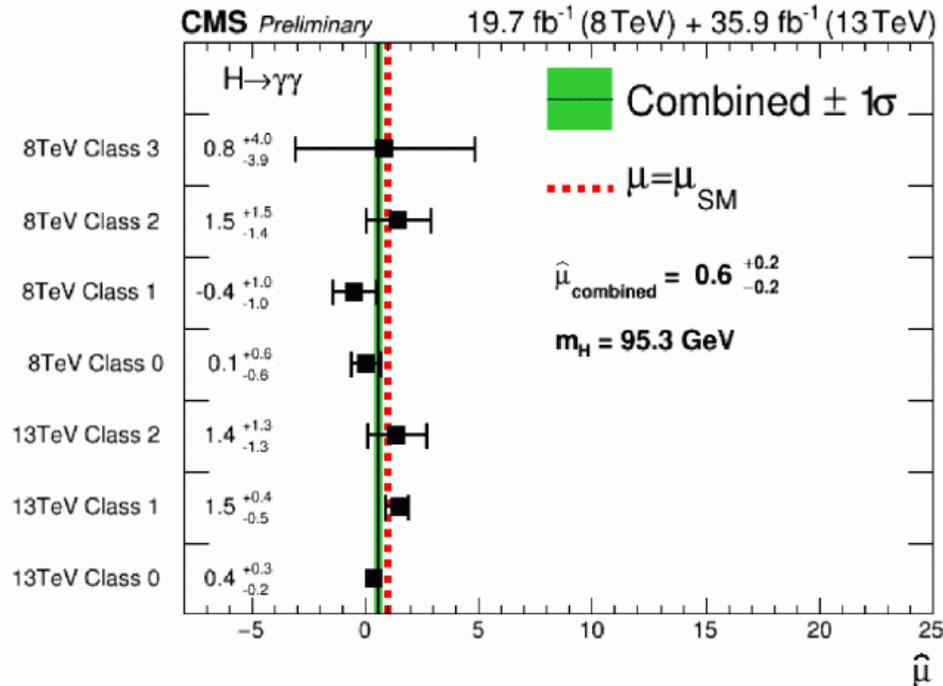
S. Gascon-Shotkin HDays17, Santander, ES Sept. 22 2017



$h \rightarrow \gamma\gamma$ (70-110 GeV) **Runs 1+2**



CMS PAS HIG-17-013



Excess here mostly driven by class 1 (&2) at 13 TeV

χ^2 probability for the seven individual values to be compatible with a single signal hypothesis: 41%

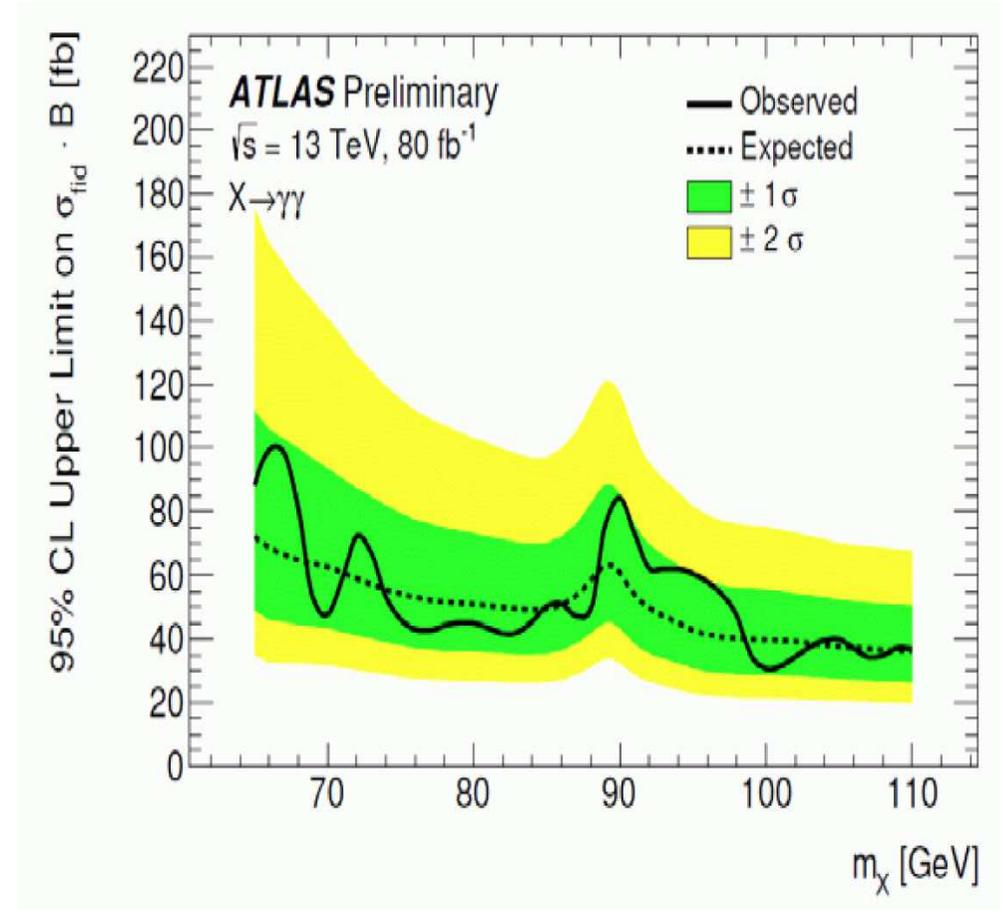
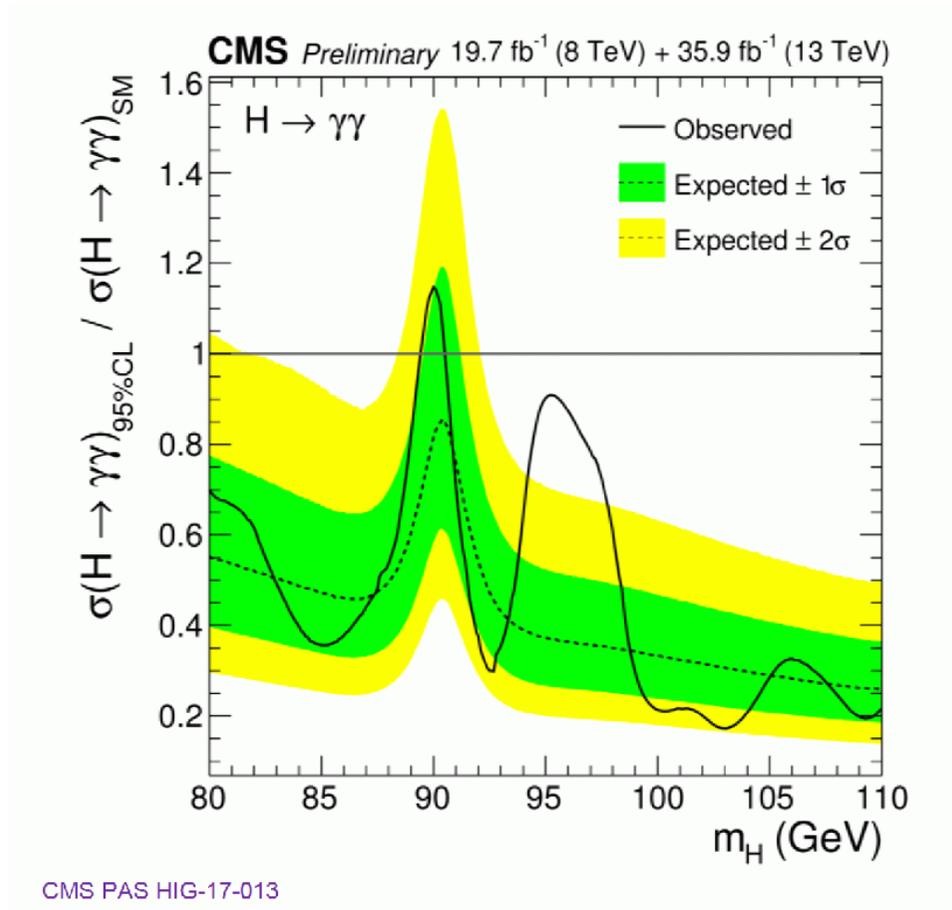
- ‘Signal’ strengths for the 7 event classes and overall, in the 8 TeV+13TeV combination, fixing $m_H=95.3 \text{ GeV}$
- More data are required to ascertain the origin of this excess

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$$\mu_{\text{CMS}}(96 \text{ GeV}) = [\sigma(pp \rightarrow h_1) \times \text{BR}(h_1 \rightarrow \gamma\gamma)]_{\text{exp/SM}} = 0.6 \pm 0.2$$

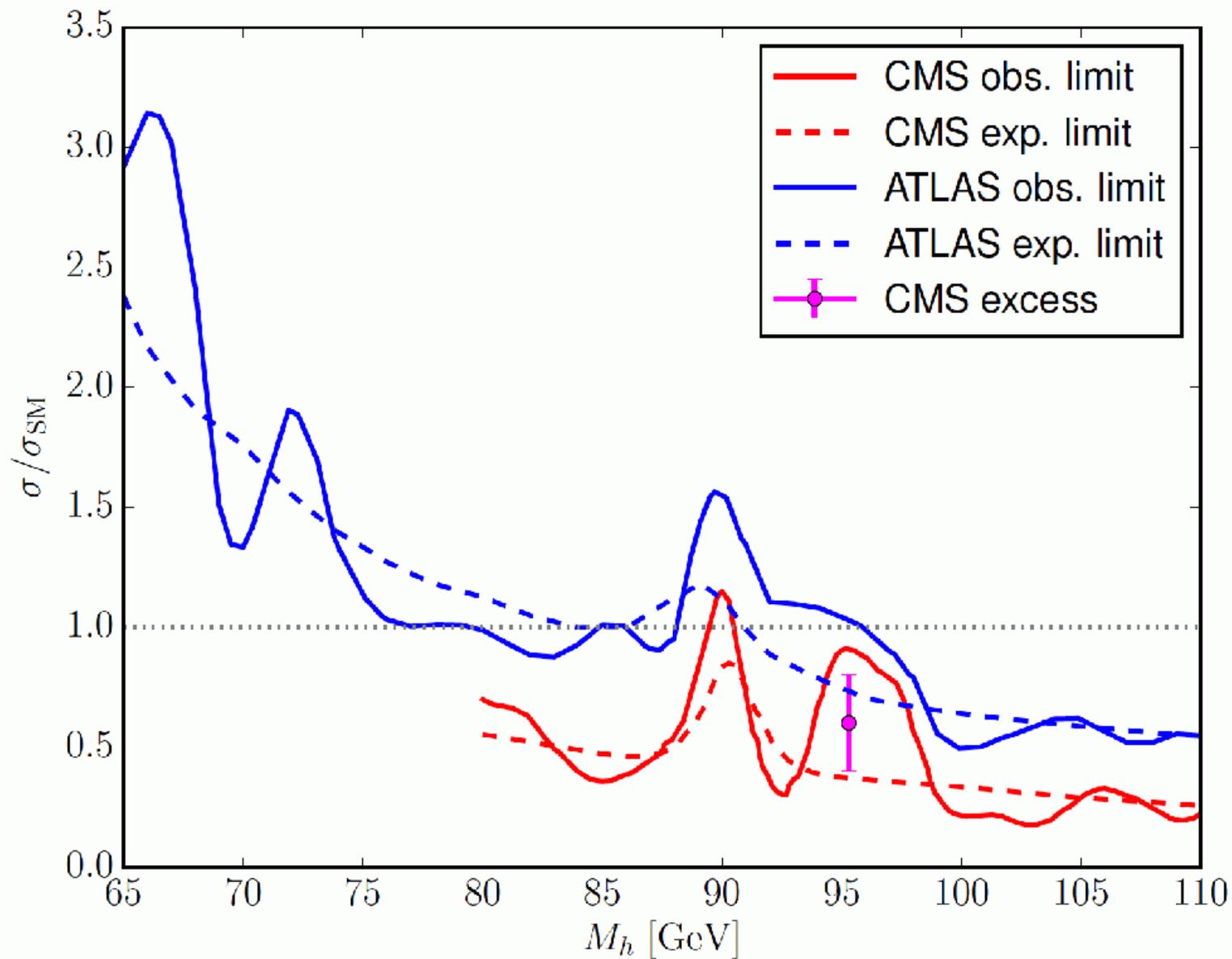
What about ATLAS?



Note: ATLAS gives fiducial cross section! Conversion factor: 1/0.45

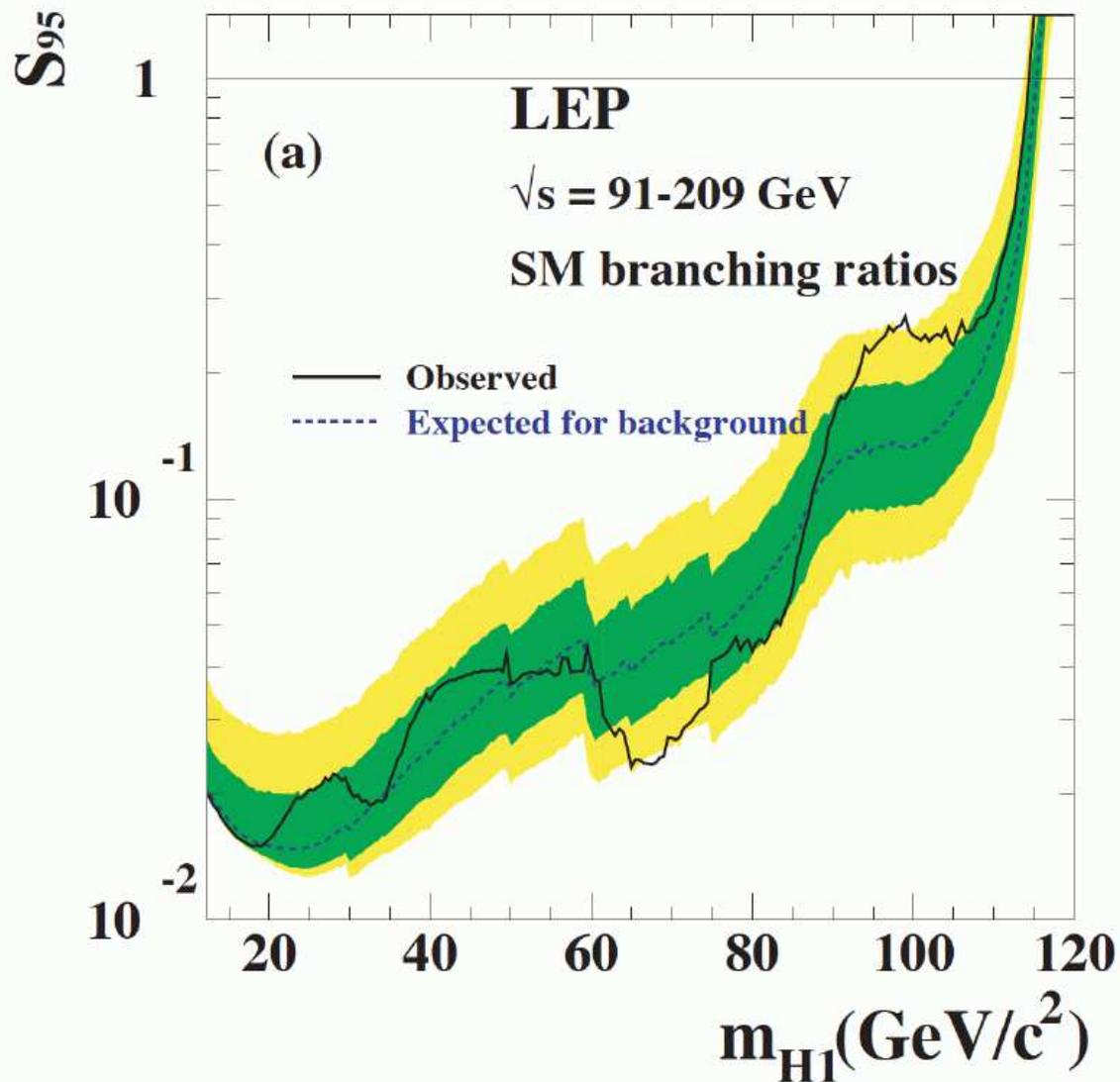
⇒ ATLAS exclusion limit even weaker than CMS!

Q: why does ATLAS has same sensitivity with twice amount of data?



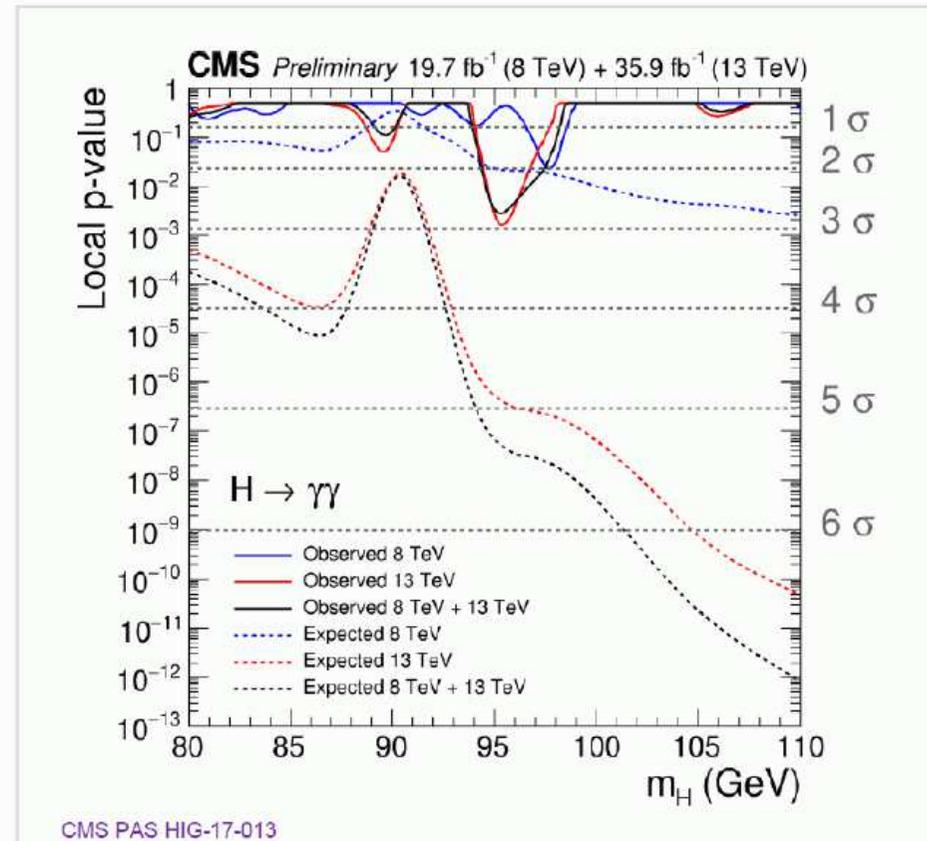
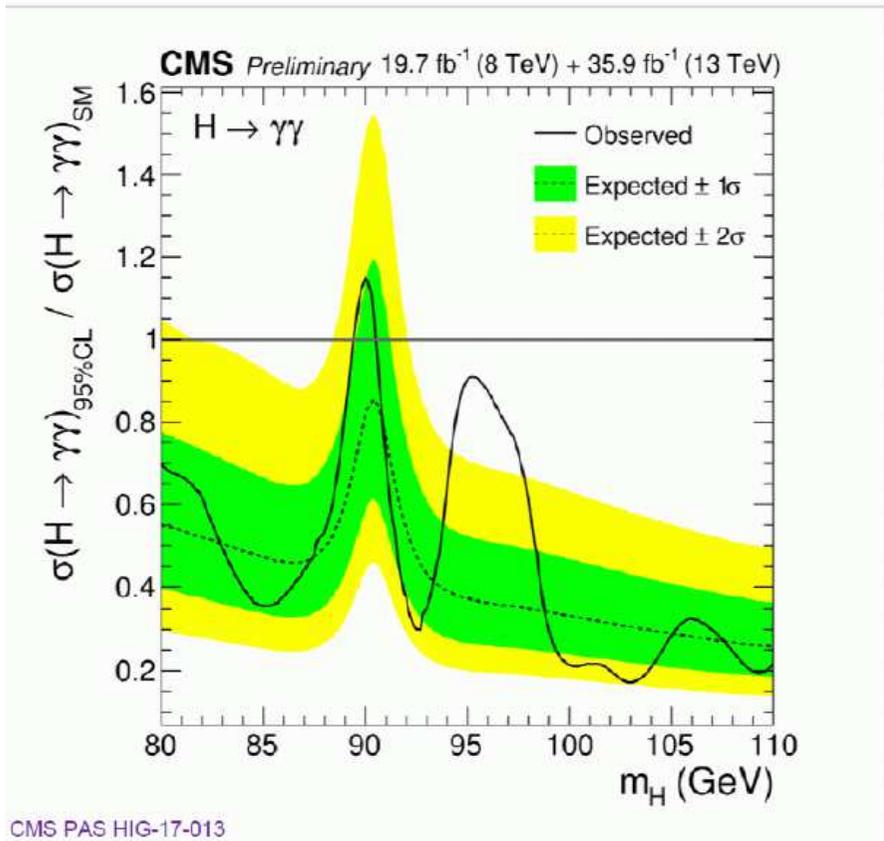
⇒ everything well compatible with the excess!

What was seen at LEP?



$$\mu_{\text{LEP}}(98 \text{ GeV}) = \left[\sigma(e^+e^- \rightarrow Zh_1) \times \text{BR}(h_1 \rightarrow b\bar{b}) \right]_{\text{exp/SM}} = 0.117 \pm 0.057$$

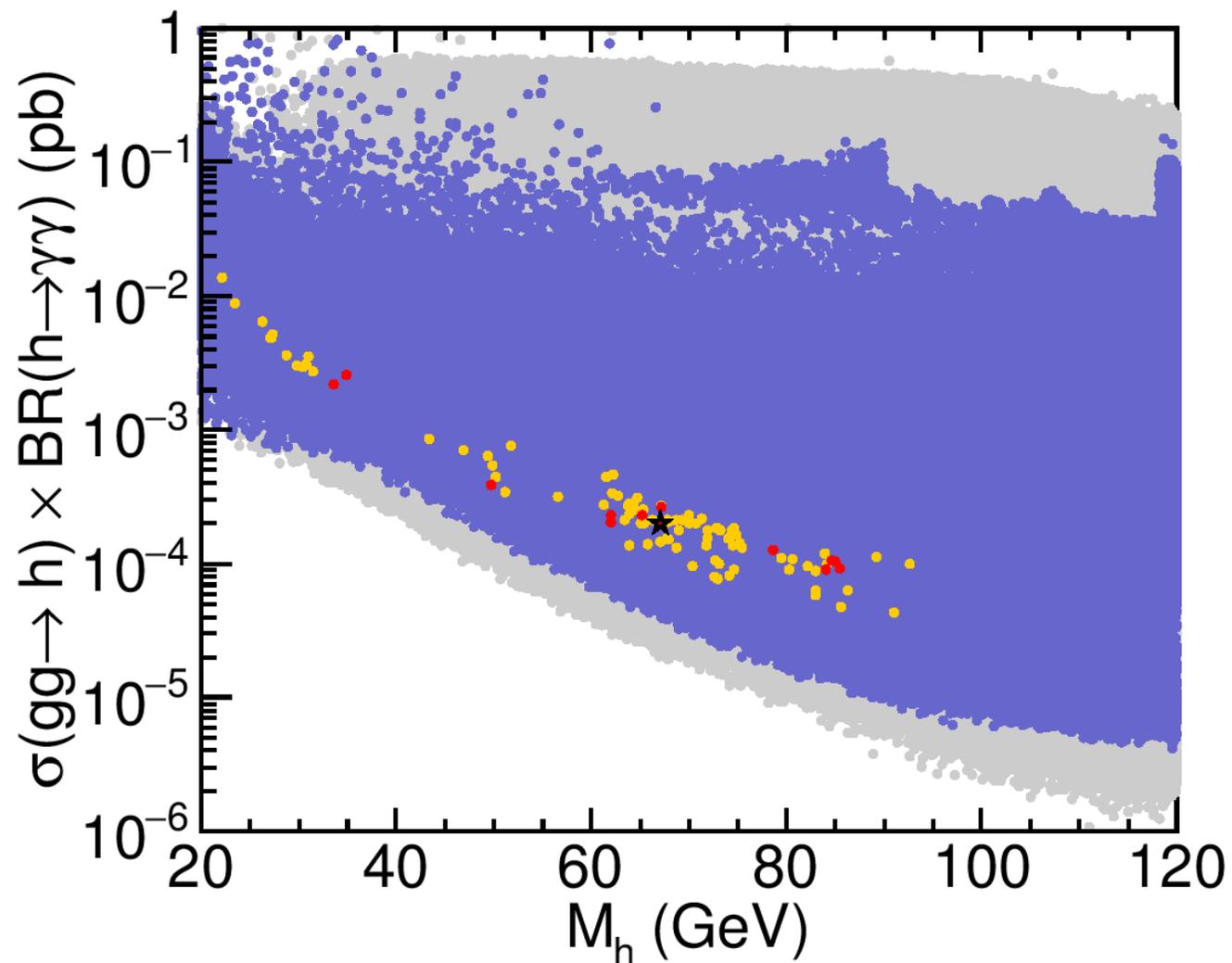
- **Combined 8 TeV + 13 TeV** $\sigma \times \text{BR}$ limit normalized to SM expectation:
 - Production processes assumed in SM proportions
 - **No significant excess** with respect to background expectations
- Expected and observed local p-values for **8 TeV**, **13 TeV** and their **combination**



Q: When do you dare to something “significant” ?

What about the MSSM?

[P. Bechtle, H. Haber, S.H., O. Stål, T. Stefaniak, G. Weiglein, L. Zeune '16]



⇒ too small rates!

⇒ problem: 2HDM structure too “rigid”

More general Ansatz: **N2HDM**

[*T. Biekötter, M. Chakraborti, S.H. '19*]

Fields:

$$\Phi_1 = \begin{pmatrix} \phi_1^+ \\ \frac{1}{\sqrt{2}}(v_1 + \rho_1 + i\eta_1) \end{pmatrix}, \quad \Phi_2 = \begin{pmatrix} \phi_2^+ \\ \frac{1}{\sqrt{2}}(v_2 + \rho_2 + i\eta_2) \end{pmatrix}, \quad \Phi_S = v_S + \rho_S$$

Potential:

$$\begin{aligned} V = & m_{11}^2 |\Phi_1|^2 + m_{22}^2 |\Phi_2|^2 - m_{12}^2 (\Phi_1^\dagger \Phi_2 + h.c.) + \frac{\lambda_1}{2} (\Phi_1^\dagger \Phi_1)^2 + \frac{\lambda_2}{2} (\Phi_2^\dagger \Phi_2)^2 \\ & + \lambda_3 (\Phi_1^\dagger \Phi_1) (\Phi_2^\dagger \Phi_2) + \lambda_4 (\Phi_1^\dagger \Phi_2) (\Phi_2^\dagger \Phi_1) + \frac{\lambda_5}{2} [(\Phi_1^\dagger \Phi_2)^2 + h.c.] \\ & + \frac{1}{2} m_S^2 \Phi_S^2 + \frac{\lambda_6}{8} \Phi_S^4 + \frac{\lambda_7}{2} (\Phi_1^\dagger \Phi_1) \Phi_S^2 + \frac{\lambda_8}{2} (\Phi_2^\dagger \Phi_2) \Phi_S^2 \end{aligned}$$

Z_2 symmetry: $\Phi_1 \rightarrow \Phi_1$, $\Phi_2 \rightarrow -\Phi_2$, $\Phi_S \rightarrow \Phi_S$

Physical states: h_1, h_2, h_3 (CP -even), A (CP -odd), H^\pm (charged)

Extension of the Z_2 symmetry to fermions determines four types:

	u -type	d -type	leptons
type I	Φ_2	Φ_2	Φ_2
type II	Φ_2	Φ_1	Φ_1
type III (lepton-specific)	Φ_2	Φ_2	Φ_1
type IV (flipped)	Φ_2	Φ_1	Φ_2

\Rightarrow exactly as in 2HDM

Three neutral \mathcal{CP} -even Higgses:

$$\begin{pmatrix} h_1 \\ h_2 \\ h_3 \end{pmatrix} = R \begin{pmatrix} \rho_1 \\ \rho_2 \\ \rho_S \end{pmatrix}, \quad R = \begin{pmatrix} c_{\alpha_1} c_{\alpha_2} & s_{\alpha_1} c_{\alpha_2} & s_{\alpha_2} \\ -(c_{\alpha_1} s_{\alpha_2} s_{\alpha_3} + s_{\alpha_1} c_{\alpha_3}) & c_{\alpha_1} c_{\alpha_3} - s_{\alpha_1} s_{\alpha_2} s_{\alpha_3} & c_{\alpha_2} s_{\alpha_3} \\ -c_{\alpha_1} s_{\alpha_2} c_{\alpha_3} + s_{\alpha_1} s_{\alpha_3} & -(c_{\alpha_1} s_{\alpha_3} + s_{\alpha_1} s_{\alpha_2} c_{\alpha_3}) & c_{\alpha_2} c_{\alpha_3} \end{pmatrix}$$

Coupling to massive gauge bosons: (identical for all four types)

$$c_{h_i VV} = c_\beta R_{i1} + s_\beta R_{i2}$$

$$h_1 \quad c_{\alpha_2} c_{\beta - \alpha_1}$$

$$h_2 \quad -c_{\beta - \alpha_1} s_{\alpha_2} s_{\alpha_3} + c_{\alpha_3} s_{\beta - \alpha_1}$$

$$h_3 \quad -c_{\alpha_3} c_{\beta - \alpha_1} s_{\alpha_2} - s_{\alpha_3} s_{\beta - \alpha_1}$$

Coupling to fermions: (same pattern as in 2HDM)

	u -type ($c_{h_i tt}$)	d -type ($c_{h_i bb}$)	leptons ($c_{h_i \tau\tau}$)
type I	$\frac{R_{i2}}{s_\beta}$	$\frac{R_{i2}}{s_\beta}$	$\frac{R_{i2}}{s_\beta}$
type II	$\frac{R_{i2}}{s_\beta}$	$\frac{R_{i1}}{c_\beta}$	$\frac{R_{i1}}{c_\beta}$
type III (lepton-specific)	$\frac{R_{i2}}{s_\beta}$	$\frac{R_{i2}}{s_\beta}$	$\frac{R_{i1}}{c_\beta}$
type IV (flipped)	$\frac{R_{i2}}{s_\beta}$	$\frac{R_{i1}}{c_\beta}$	$\frac{R_{i2}}{s_\beta}$

“Physical” input parameters:

$$\alpha_{1,2,3}, \quad \tan \beta, \quad v, \quad v_S, \quad m_{h_{1,2,3}}, \quad m_A, \quad M_{H^\pm}, \quad m_{12}^2$$

Needed to fit the two excesses: $m_{h_1} \sim 96 \text{ GeV}$, $m_{h_2} \sim 125 \text{ GeV}$

- $c_{h_1 VV}^2$ strongly reduced for μ_{LEP}
- $c_{h_1 bb}$ reduced to enhance $\text{BR}(h_1 \rightarrow \gamma\gamma)$
- $c_{h_1 tt}$ not reduced for μ_{CMS}
- $c_{h_1 \tau\tau}$ possibly reduced to enhance $\text{BR}(h_1 \rightarrow \gamma\gamma)$

	Decrease $c_{h_1 b\bar{b}}$	No decrease $c_{h_1 t\bar{t}}$	No enhancement $c_{h_1 \tau\bar{\tau}}$
type I	$(\frac{R_{12}}{s_\beta}) :-)$	$(\frac{R_{12}}{s_\beta}) :-)$	$(\frac{R_{12}}{s_\beta}) :-)$
type II	$(\frac{R_{11}}{c_\beta}) :-)$	$(\frac{R_{12}}{s_\beta}) :-)$	$(\frac{R_{11}}{c_\beta}) :-)$
type III	$(\frac{R_{12}}{s_\beta}) :-)$	$(\frac{R_{12}}{s_\beta}) :-)$	$(\frac{R_{11}}{c_\beta}) :-)$
type IV	$(\frac{R_{11}}{c_\beta}) :-)$	$(\frac{R_{12}}{s_\beta}) :-)$	$(\frac{R_{12}}{s_\beta}) :-)$

Type II and IV: $c_{h_1 bb}$ and $c_{h_1 tt}$ independent

Type II bonus: $c_{h_1 \tau\tau}$ can be suppressed (together with $c_{h_1 bb}$)

\Rightarrow only type II and IV can fit CMS and LEP excesses

⇒ Parameter scan ⇒ ScannerS

Constraints:

- Tree-level perturbativity ⇒ ScannerS
- Minimum of potential is global minimum ⇒ ScannerS
- Higgs searches at LEP, Tevatron, LHC ⇒ HiggsBounds (N2HDECAY)
- SM-like Higgs properties ⇒ HiggsSignals (N2HDECAY, SusHi)
 $\chi_{\text{red}}^2 := \chi^2/n_{\text{obs}}$
- Flavor physics (mainly $\text{BR}(B_s \rightarrow X_s \gamma)$, ΔM_{B_s}) ⇒ SuperIso bounds
- Electroweak precision data (T and S) ⇒ ScannerS

Fitting the excesses:

$$\mu_{\text{LEP}} = 0.117 \pm 0.057, \quad \mu_{\text{CMS}} = 0.6 \pm 0.2$$

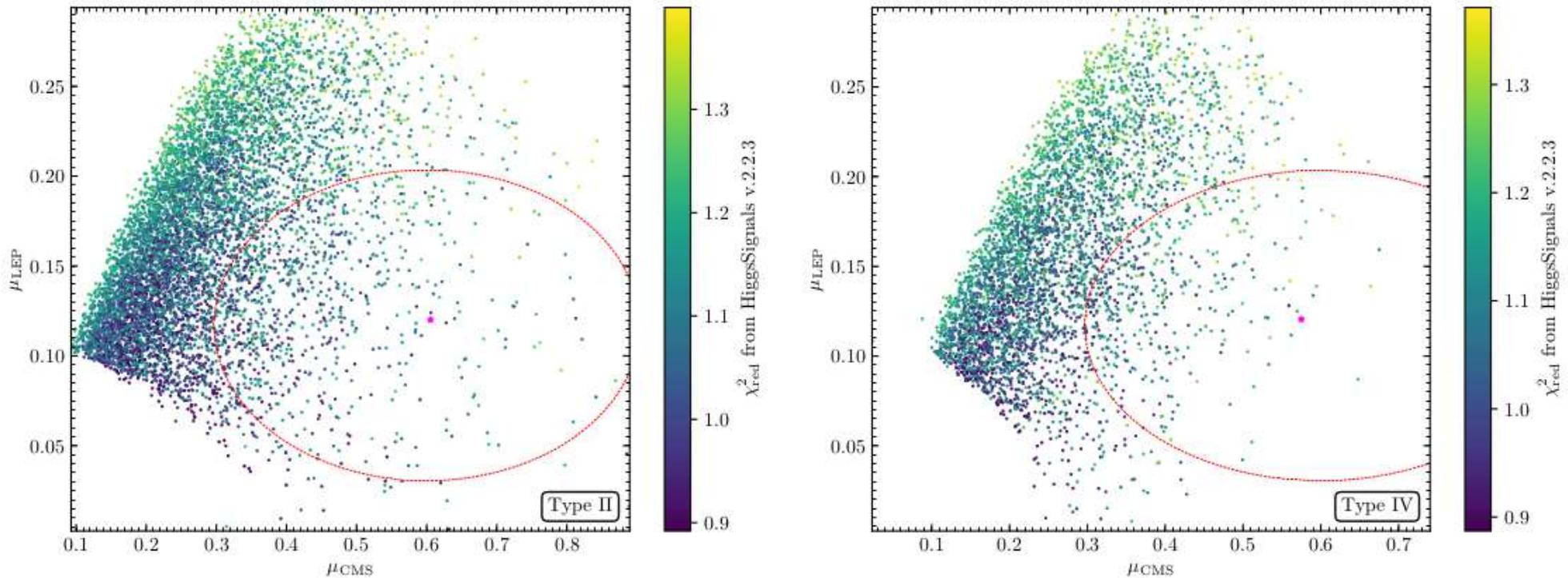
$$\begin{aligned} \mu_{\text{LEP}} &= \frac{\sigma_{\text{N2HDM}}(e^+e^- \rightarrow Zh_1)}{\sigma_{\text{SM}}(e^+e^- \rightarrow ZH)} \cdot \frac{\text{BR}_{\text{N2HDM}}(h_1 \rightarrow b\bar{b})}{\text{BR}_{\text{SM}}(H \rightarrow b\bar{b})} \\ &= |c_{h_1VV}|^2 \frac{\text{BR}_{\text{N2HDM}}(h_1 \rightarrow b\bar{b})}{\text{BR}_{\text{SM}}(H \rightarrow b\bar{b})} \end{aligned}$$

$$\begin{aligned} \mu_{\text{CMS}} &= \frac{\sigma_{\text{N2HDM}}(gg \rightarrow h_1)}{\sigma_{\text{SM}}(gg \rightarrow H)} \cdot \frac{\text{BR}_{\text{N2HDM}}(h_1 \rightarrow \gamma\gamma)}{\text{BR}_{\text{SM}}(H \rightarrow \gamma\gamma)} \\ &= |c_{h_1tt}|^2 \frac{\text{BR}_{\text{N2HDM}}(h_1 \rightarrow \gamma\gamma)}{\text{BR}_{\text{SM}}(H \rightarrow \gamma\gamma)} \end{aligned}$$

$$\chi_{\text{CMS-LEP}}^2 = \frac{(\mu_{\text{LEP}} - 0.117)^2}{(0.057)^2} + \frac{(\mu_{\text{CMS}} - 0.6)^2}{(0.2)^2}$$

⇒ “best-fit point”

Fitting the excesses: [T. Biekötter, M. Chakraborti, S.H. '19]



⇒ excesses well fitted, with good χ_{red}^2

⇒ preferred M_{H^\pm} : 650 GeV – 950 GeV (lower limit: flavor constr.)

⇒ preferred $\tan \beta$: 0.8 – 3.8

Best-fit point in type II:

m_{h_1}	m_{h_2}	m_{h_3}	m_A	M_{H^\pm}	
96.5263	125.09	535.86	712.578	737.829	
$\tan \beta$	α_1	α_2	α_3	m_{12}^2	v_S
1.26287	1.26878	-1.08484	-1.24108	80644.3	272.72
$\text{BR}_{h_1}^{bb}$	$\text{BR}_{h_1}^{gg}$	$\text{BR}_{h_1}^{\tau\tau}$	$\text{BR}_{h_1}^{\gamma\gamma}$	$\text{BR}_{h_1}^{WW}$	$\text{BR}_{h_1}^{ZZ}$
0.5048	0.2682	$5.09 \cdot 10^{-2}$	$2.582 \cdot 10^{-3}$	$1.37 \cdot 10^{-2}$	$1.753 \cdot 10^{-3}$
$\text{BR}_{h_2}^{bb}$	$\text{BR}_{h_2}^{gg}$	$\text{BR}_{h_2}^{\tau\tau}$	$\text{BR}_{h_2}^{\gamma\gamma}$	$\text{BR}_{h_2}^{WW}$	$\text{BR}_{h_2}^{ZZ}$
0.5916	0.0771	$6.36 \cdot 10^{-2}$	$2.153 \cdot 10^{-3}$	0.2087	$2.610 \cdot 10^{-3}$

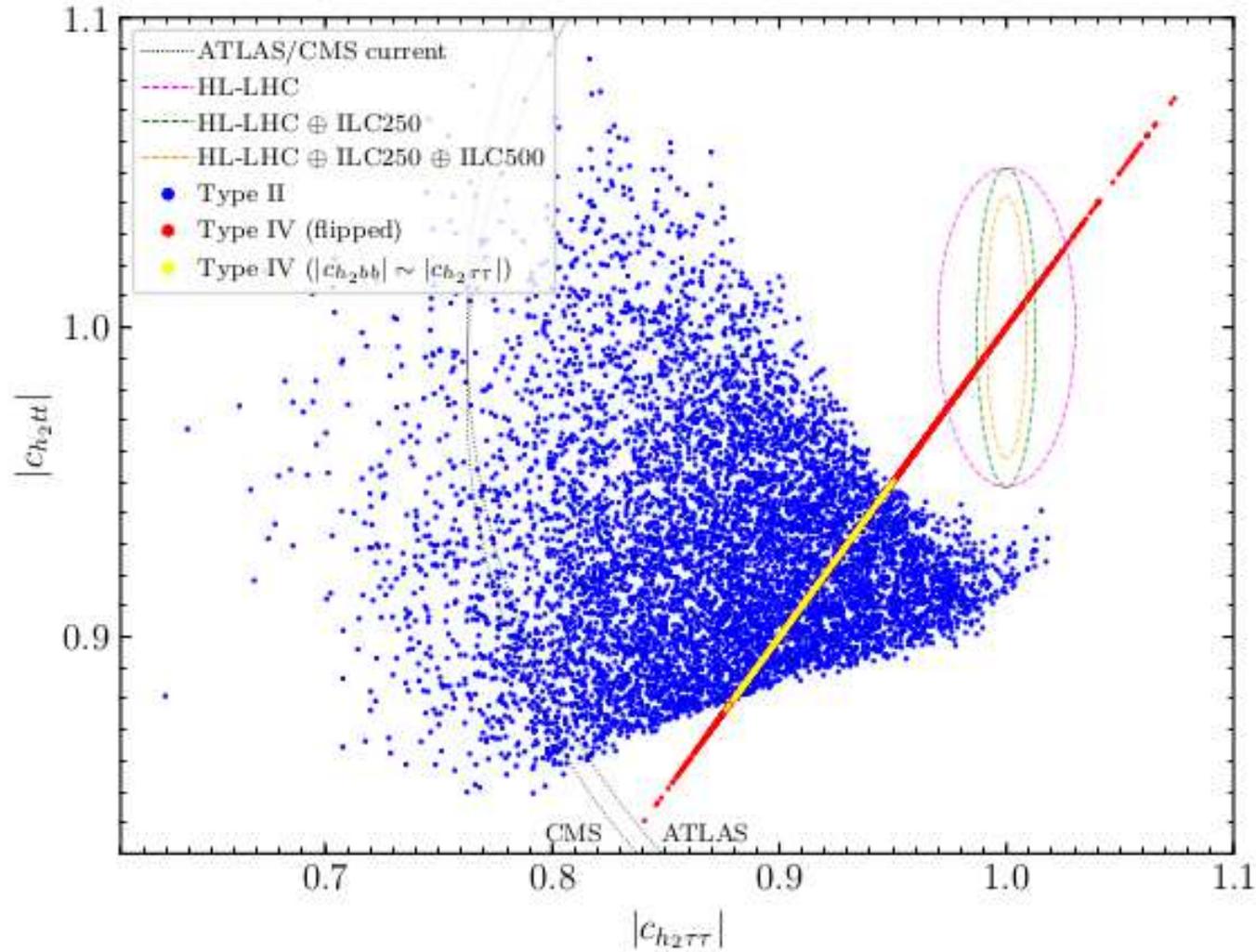
⇒ surprisingly large $\text{BR}_{h_1}^{\gamma\gamma}$

Best-fit point in type IV:

m_{h_1}	m_{h_2}	m_{h_3}	m_A	M_{H^\pm}	
97.8128	125.09	485.998	651.502	651.26	
$\tan \beta$	α_1	α_2	α_3	m_{12}^2	v_S
1.3147	1.27039	-1.02829	-1.32496	41034.1	647.886
$\text{BR}_{h_1}^{bb}$	$\text{BR}_{h_1}^{gg}$	$\text{BR}_{h_1}^{\tau\tau}$	$\text{BR}_{h_1}^{\gamma\gamma}$	$\text{BR}_{h_1}^{WW}$	$\text{BR}_{h_1}^{ZZ}$
0.4074	0.20714	0.248324	$2.139 \cdot 10^{-3}$	$1.347 \cdot 10^{-2}$	$1.579 \cdot 10^{-3}$
$\text{BR}_{h_2}^{bb}$	$\text{BR}_{h_2}^{gg}$	$\text{BR}_{h_2}^{\tau\tau}$	$\text{BR}_{h_2}^{\gamma\gamma}$	$\text{BR}_{h_2}^{WW}$	$\text{BR}_{h_2}^{ZZ}$
0.5363	0.09388	$7.58 \cdot 10^{-2}$	$2.247 \cdot 10^{-3}$	0.2267	$2.836 \cdot 10^{-2}$

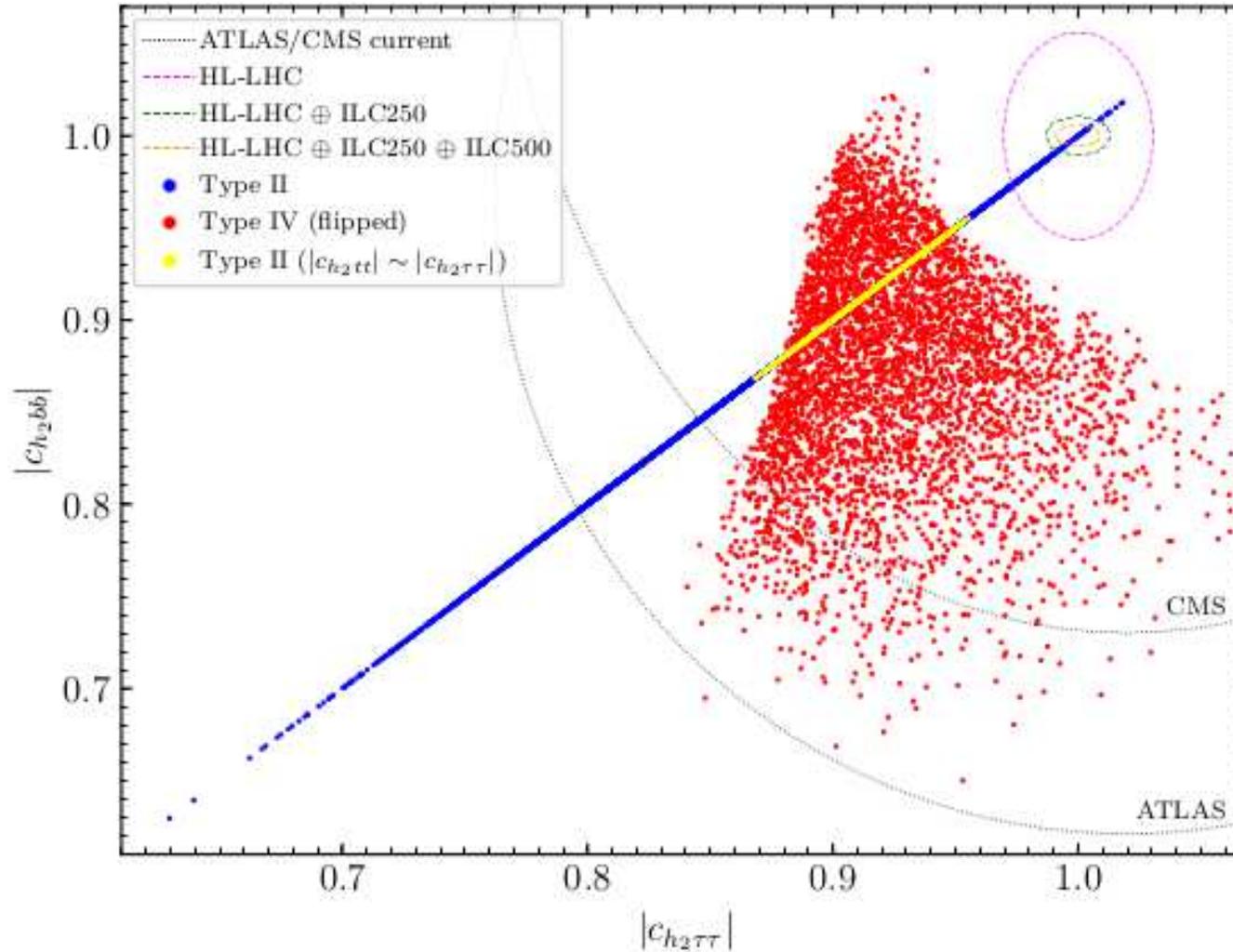
⇒ substantially larger $\text{BR}_{h_1}^{\tau\tau}$ than in type II

Next project? \Rightarrow ILC Higgs coupling measurements



\Rightarrow type II shows deviation from SM

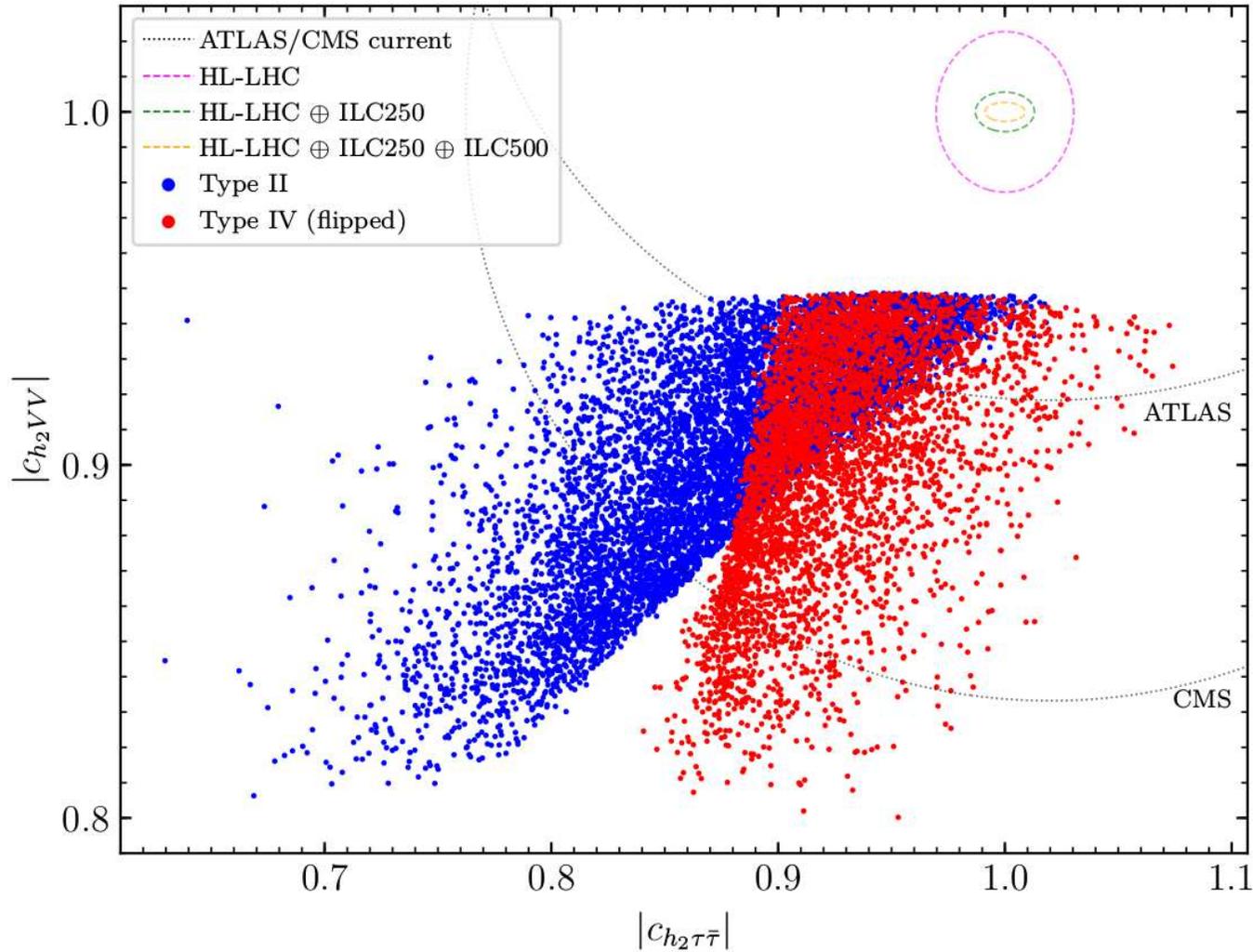
Next project? \Rightarrow ILC Higgs coupling measurements



\Rightarrow type IV shows deviations from SM

\Rightarrow N2HDM can always be distinguished from SM!

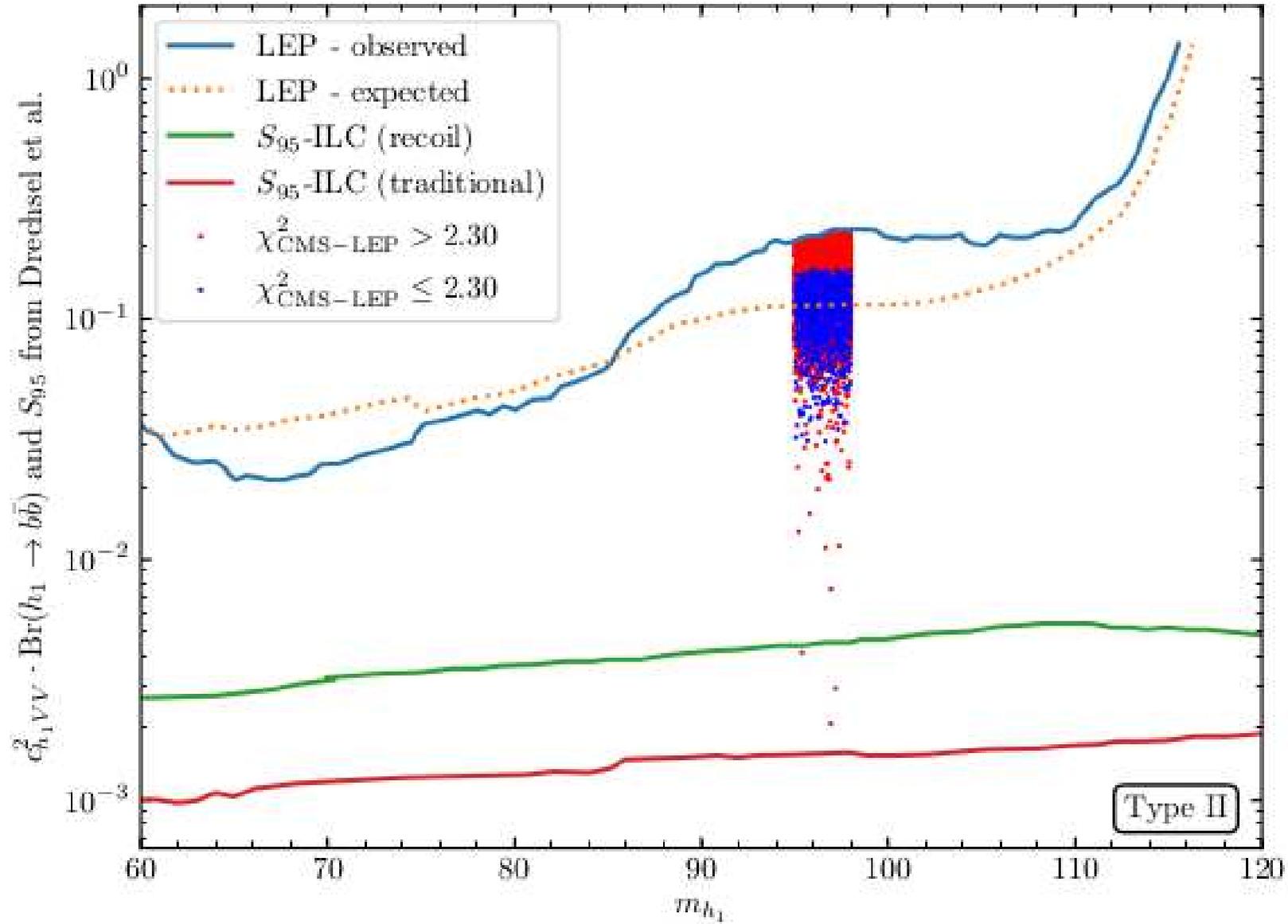
Next project? \Rightarrow ILC Higgs coupling measurements



\Rightarrow type II and IV show strong deviations from SM

\Rightarrow N2HDM can always be distinguished from SM!

Next project? \Rightarrow ILC production of the light scalar



\Rightarrow new state easily in the reach of the ILC

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⇒ type II fits best, type II is needed for SUSY

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- $\mu\nu$ SSM
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Q: Can the models fit the excesses **despite** the additional SUSY constraints on the Higgs sector **???**

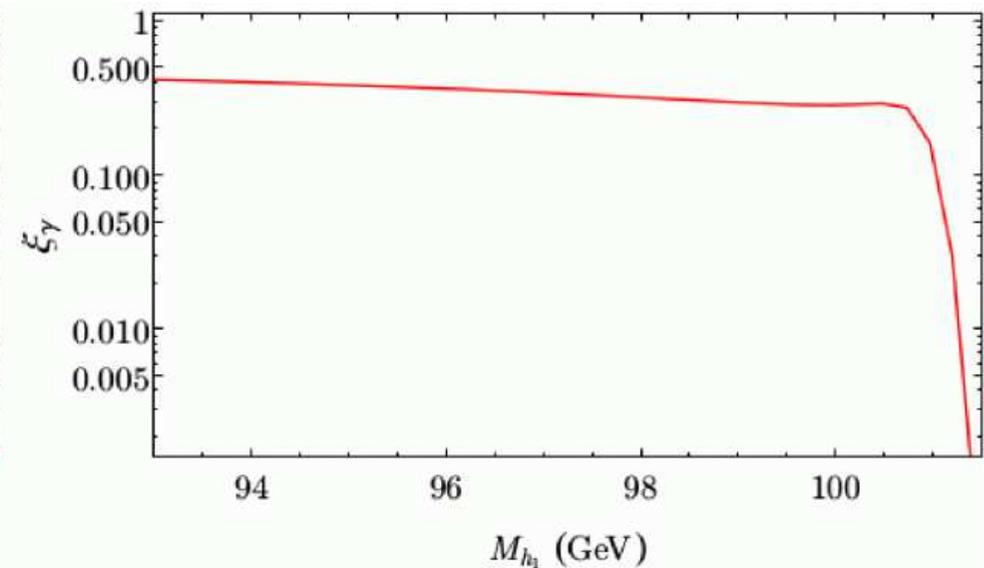
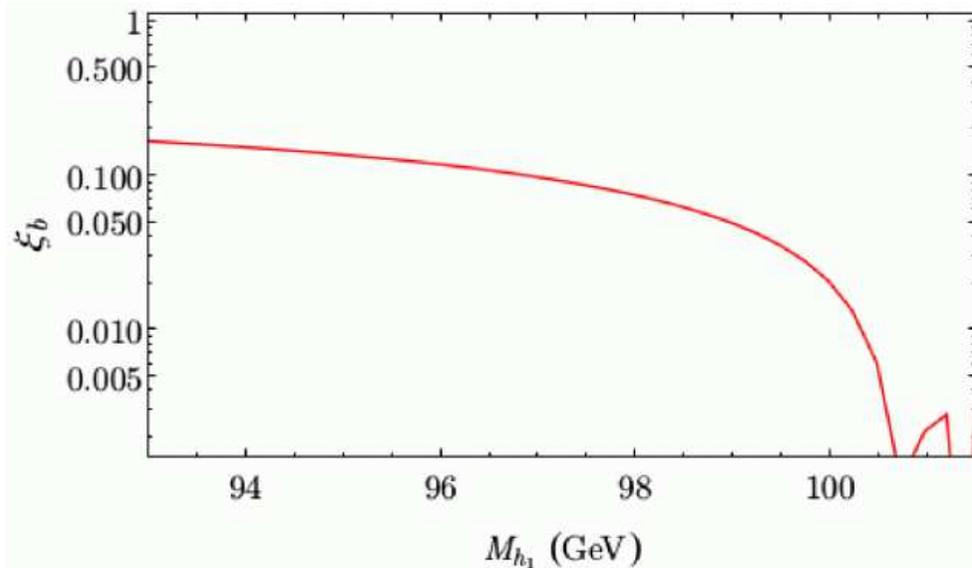
What about the NMSSM?

[F. Domingo, S.H., S. Passehr, G. Weiglein '18]

Parameters:

$\lambda = 0.6$, $\kappa = 0.035$, $\tan\beta = 2$, $\mu_{\text{eff}} = (397 + 15x)$ GeV, $M_{H^\pm} = 1$ TeV,
 $A_\kappa = -325$ GeV, $M_{\text{SUSY}} = 1$ TeV, $A_t = A_b = 0$

$$\xi_b \equiv \frac{\Gamma[h_1 \rightarrow ZZ] \cdot \text{BR}[h_1 \rightarrow b\bar{b}]}{\Gamma[H_{\text{SM}}(M_{h_1}) \rightarrow ZZ] \cdot \text{BR}[H_{\text{SM}}(M_{h_1}) \rightarrow b\bar{b}]} \sim \frac{\sigma[e^+e^- \rightarrow Z(h_1 \rightarrow b\bar{b})]}{\sigma[e^+e^- \rightarrow Z(H_{\text{SM}}(M_{h_1}) \rightarrow b\bar{b})]}$$
$$\xi_\gamma \equiv \frac{\Gamma[h_1 \rightarrow gg] \cdot \text{BR}[h_1 \rightarrow \gamma\gamma]}{\Gamma[H_{\text{SM}}(M_{h_1}) \rightarrow gg] \cdot \text{BR}[H_{\text{SM}}(M_{h_1}) \rightarrow \gamma\gamma]} \sim \frac{\sigma[gg \rightarrow h_1 \rightarrow \gamma\gamma]}{\sigma[gg \rightarrow H_{\text{SM}}(M_{h_1}) \rightarrow \gamma\gamma]}.$$



⇒ both excesses can be fitted simultaneously!

What about the $\mu\nu$ SSM?

$\mu\nu$ SSM: [*D. Lopez-Fogliani, C. Muñoz '06*]

$\mu\nu$ SSM: NMSSM + well motivated RPV (in simple terms)
 \Rightarrow EW scale seesaw to reproduce the neutrino data

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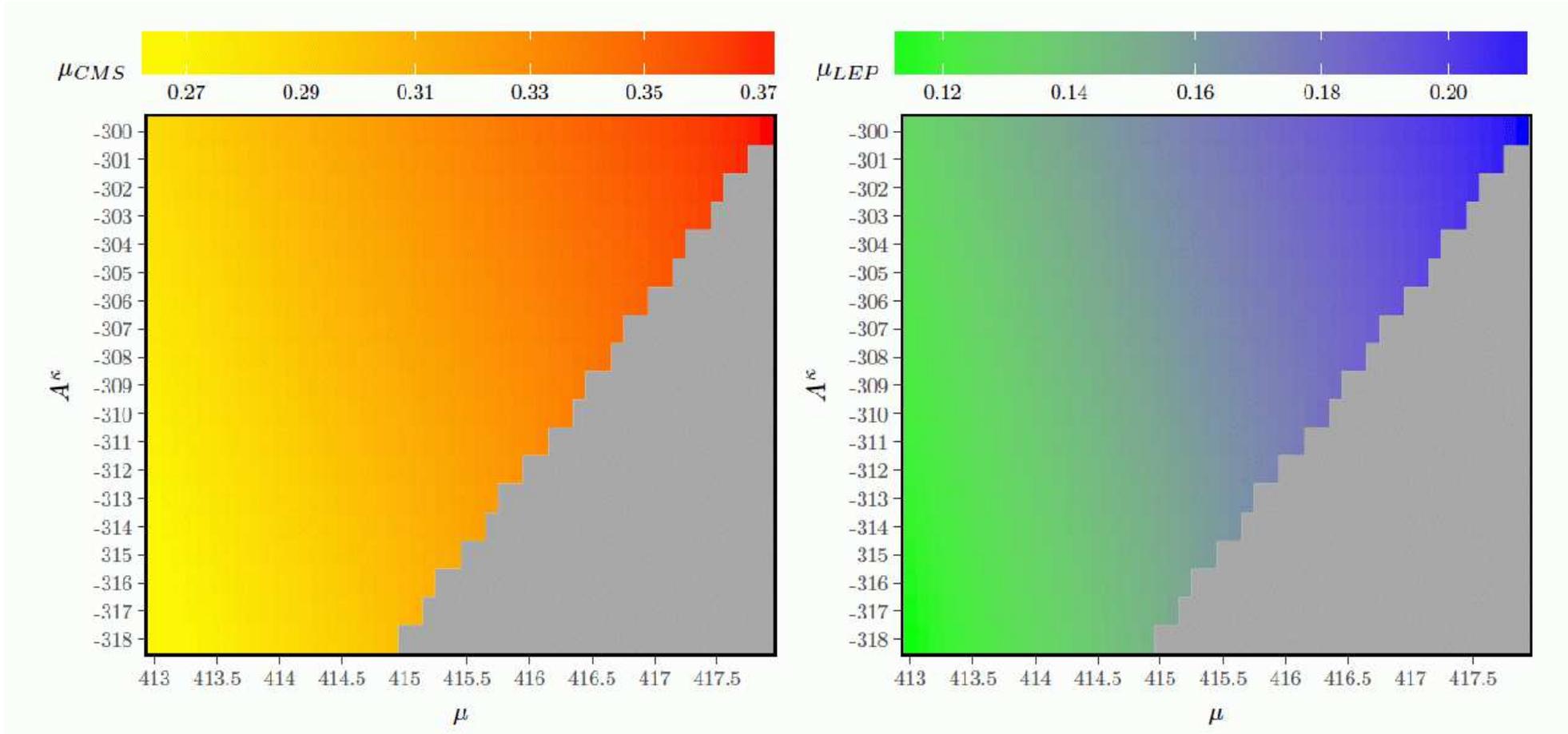
Can the $\mu\nu$ SSM explain the two excesses?

[T. Biekötter, S.H., C. Muñoz '17]

v_{iL}	Y_i^ν	A_i^ν	$\tan\beta$	μ	λ	A^λ	κ	A^κ	M_1
$\sqrt{2} \cdot 10^{-5}$	10^{-7}	-1000	2	[413; 418]	0.6	956.035	0.035	[-300; -318]	100
M_2	M_3	$m_{\tilde{Q}_{iL}}^2$	$m_{\tilde{u}_{iR}}^2$	$m_{\tilde{d}_{iR}}^2$	A_1^u	$A_{2,3}^{u,d}$	$(m_e^2)_{ii}$	A_{33}^e	$A_{11,22}^e$
200	1500	800^2	800^2	800^2	0	0	800^2	0	0

Can the $\mu\nu$ SSM explain the two excesses?

[*T. Biekötter, S.H., C. Muñoz '17*]



⇒ YES, WE CAN! :-)
(at the 1 – 1.5 σ level)

4. Conclusinos

- The Higgs boson discovered at the LHC **cannot be the SM Higgs!**
- New MSSM Higgs benchmark proposal:
 - M_h^{125} scenario: 2HDM-like model
 - $M_h^{125}(\tilde{\chi})$ scenario: light EW-inos: $H/A \rightarrow \tilde{\chi}_i^0 \tilde{\chi}_j^0, \tilde{\chi}_k^\pm \tilde{\chi}_l^\mp$
 - M_H^{125} scenario: $M_H \sim 125$ GeV, all Higgses light
- A light Higgs at 96 GeV?

new CMS/ATLAS result: $pp \rightarrow \phi_{96} \rightarrow \gamma\gamma, \mu_{\text{CMS}} = 0.6 \pm 0.2$

old LEP result: $e^+e^- \rightarrow Z \phi_{96} \rightarrow Z b\bar{b}, \mu_{\text{LEP}} = 0.117 \pm 0.057$

 - **MSSM** cannot explain the excesses
 - **N2HDM** easily fits the excesses
 - \Rightarrow **type II** favored (as predicted by SUSY)
 - **NMSSM** can explain CMS(/ATLAS) and LEP excesses
 - **$\mu\nu$ SSM** can explain CMS(/ATLAS) and LEP excesses
 - \Rightarrow **perfect physics case for the ILC:** 96 GeV direct \oplus 125 GeV coupl.

Katharsis of Ultimate Theory Standards

10th meeting: 08.-10. April 2019 (Dresden Univ.)

Precise Calculation of

(N)

Higgs Boson masses

Local organizer: D. Stoeckinger

Organized by:
M. Carena, H. Haber
R. Harlander, S. Heinemeyer
W. Hollik, P. Slavich, G. Weiglein

⇒ next meeting: 11/2019 in Munich

Workshop announcement:



Opportunities at Future High Energy Colliders

Workshop dates: June 11 - July 05 2019 (IFT, Madrid, Spain)

The workshop will bring together key theorists and experimentalists to address these questions, aiming at a more coherent, global view of the opportunities and rationale for the next generation of high energy colliders.

Program of the workshop:

- first week: dark matter and implications from cosmology
- second week: origin of lepton and quark flavour structure; fundamental symmetry tests
- third week: electroweak symmetry breaking; naturalness
- final week: discussion of complementary of the different collider opportunities as pertains to the physics themes.

Higgs Days at Santander 2019

Theory meets Experiment

16.-20. September



Contact: Sven.Heinemeyer@cern.ch
Local: Alicia.Calderon@cern.ch
Gervasio.Gomez@cern.ch
<http://hdays.csic.es>

Further Questions?



Obtaining a light Higgs with SM-like couplings

[J. Gunion, H. Haber, hep-ph/0207010]

→ \mathcal{CP} conserving 2HDM in the Higgs basis ($\langle H_1 \rangle = v/\sqrt{2}$, $\langle H_2 \rangle = 0$)

$$\mathcal{V} = \dots + \frac{1}{2}Z_1(H_1^\dagger H_1)^2 + \dots + \left[\frac{1}{2}Z_5(H_1^\dagger H_2)^2 + Z_6(H_1^\dagger H_1)(H_1^\dagger H_2) + \text{h.c.} \right] + \dots$$

⇒ \mathcal{CP} -even mass matrix:

$$\mathcal{M}^2 = \begin{pmatrix} Z_1 v^2 & Z_6 v^2 \\ Z_6 v^2 & M_A^2 + Z_5 v^2 \end{pmatrix}$$

with mixing angle $\cos(\beta - \alpha) \equiv c_{\beta-\alpha}$

Decoupling limit: $M_A^2 \gg Z_i v^2$
⇒ $m_h^2 \sim Z_1 v^2$, $|c_{\beta-\alpha}| \ll 1$, h is SM-like

Alignment limit: $Z_6 = 0$ and $Z_1 < Z_5 + M_A^2/v^2$
⇒ h is identical to the SM Higgs, $c_{\beta-\alpha} = 0$
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⇒ H is identical to the SM Higgs, $c_{\beta-\alpha} = 1$

Alignment limit: see e.g.

[M. Carena, I. Low, N. Shah, C. Wagner '13][M. Carena, H. Haber, I. Low, N. Shah, C. Wagner '14]

In the **MSSM** $Z_6 = 0$ can be obtained through an “accidental” cancellation between tree-level and loop contribution, roughly at:

$$\tan \beta \sim \left[M_h^2 + M_Z^2 + \frac{3m_t^2 \mu^2}{4\pi^2 v^2 M_S^2} \left(\frac{A_t^2}{2M_S^2} - 1 \right) \right] / \left[\frac{3m_t^2}{4\pi^2 v^2} \frac{\mu A_t}{M_S^2} \left(\frac{A_t^2}{6M_S^2} - 1 \right) \right]$$

Compare: $m_h^{\text{mod+}}$ and m_h^{alt} :

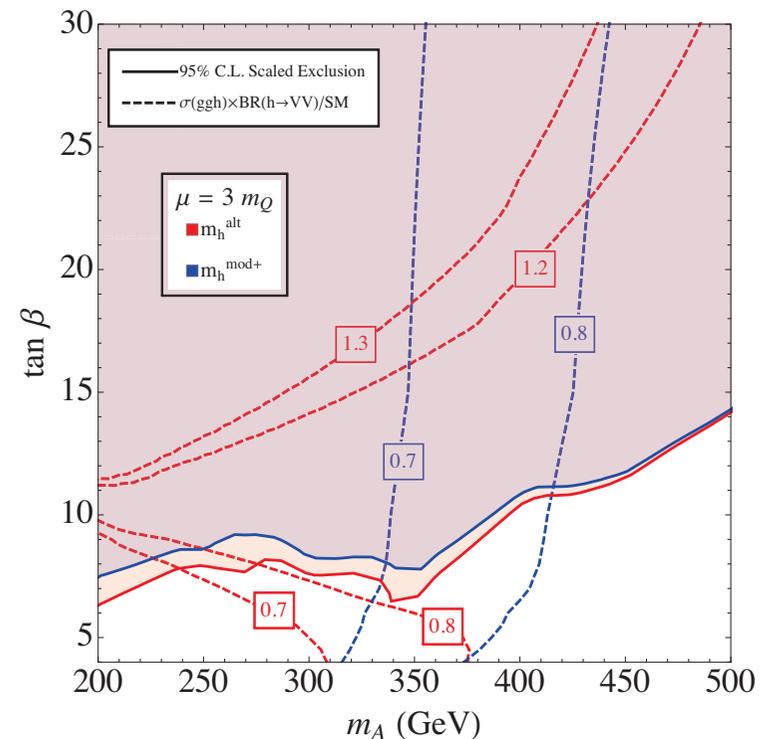
$$A_t/M_S = 2.45, \quad A_t = A_f,$$

$$M_S = m_{\tilde{f}} \geq 1 \text{ TeV}, \quad m_{\tilde{g}} = 1.5 \text{ TeV},$$

$$M_2 = 2 M_1 = 200 \text{ GeV}, \quad \mu \text{ adjustable}$$

(low M_A and $\tan \beta$: tune $M_S \geq 1 \text{ TeV}$ to obtain $M_h \geq 122 \text{ GeV}$)

\Rightarrow SM-like Higgs for all M_A



Alignment limit: see e.g.

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m_h^{alt} : **HiggsSignals** [P. Bechtle et al. '15]

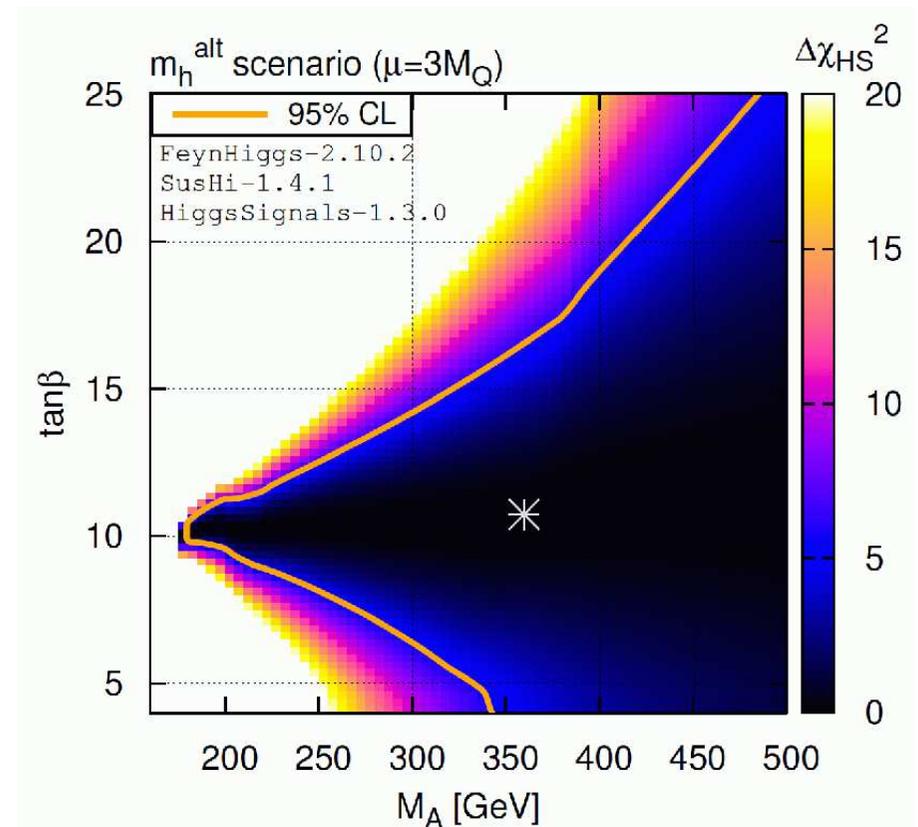
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Interesting case: light singlet

Singlet does not couple to SM particles!

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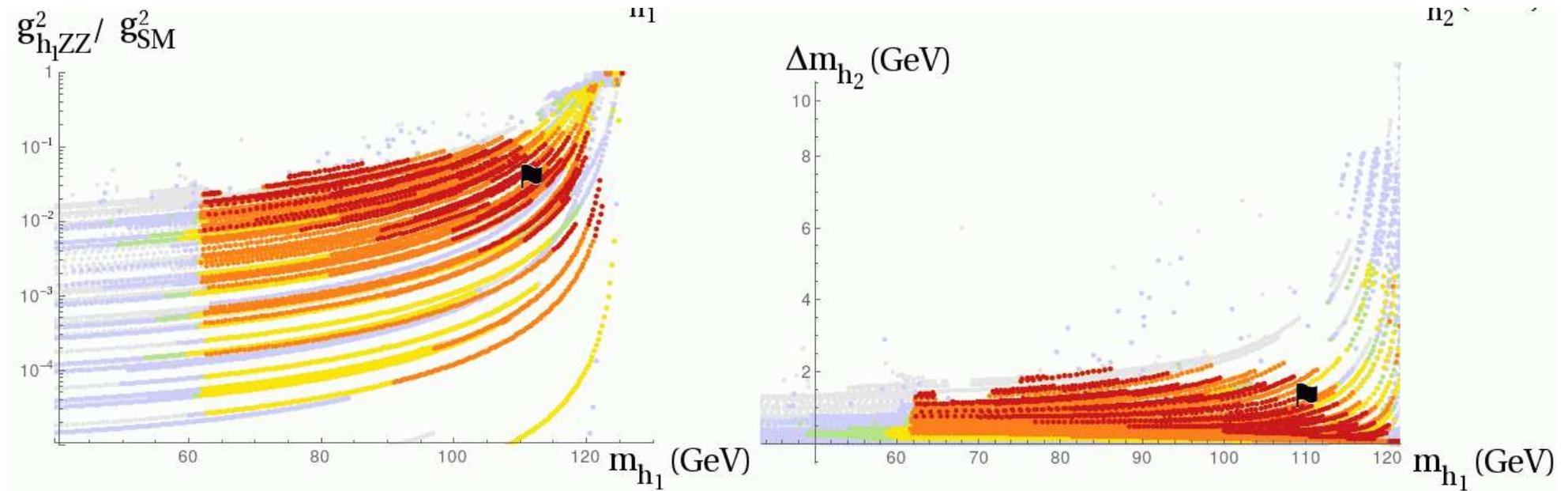
“Easily” possible in the NMSSM:

Light, singlet-like Higgs below 125 GeV

Which collider can find them?

Parameters:

$\tan \beta = 8$, $M_A = 1$ TeV, $A_\kappa = -2 \dots 0$ TeV, $\mu = 120 \dots 2000$ GeV,
 $2M_1 = M_2 = 500$ GeV, $M_3 = 1.5$ TeV, $m_{\tilde{Q}_3} = 1$ TeV, $m_{\tilde{Q}_{1,2}} = 1.5$ TeV,
 $A_t = -2$ TeV, $A_{b,\tau} = -1.5$ TeV



⇒ light Higgs below 125 GeV

⇒ strongly reduced couplings to gauge bosons!

⇒ possibly within ILC reach!

Data to be taken into account:

- Higgs boson mass (LHC) \Rightarrow FeynHiggs

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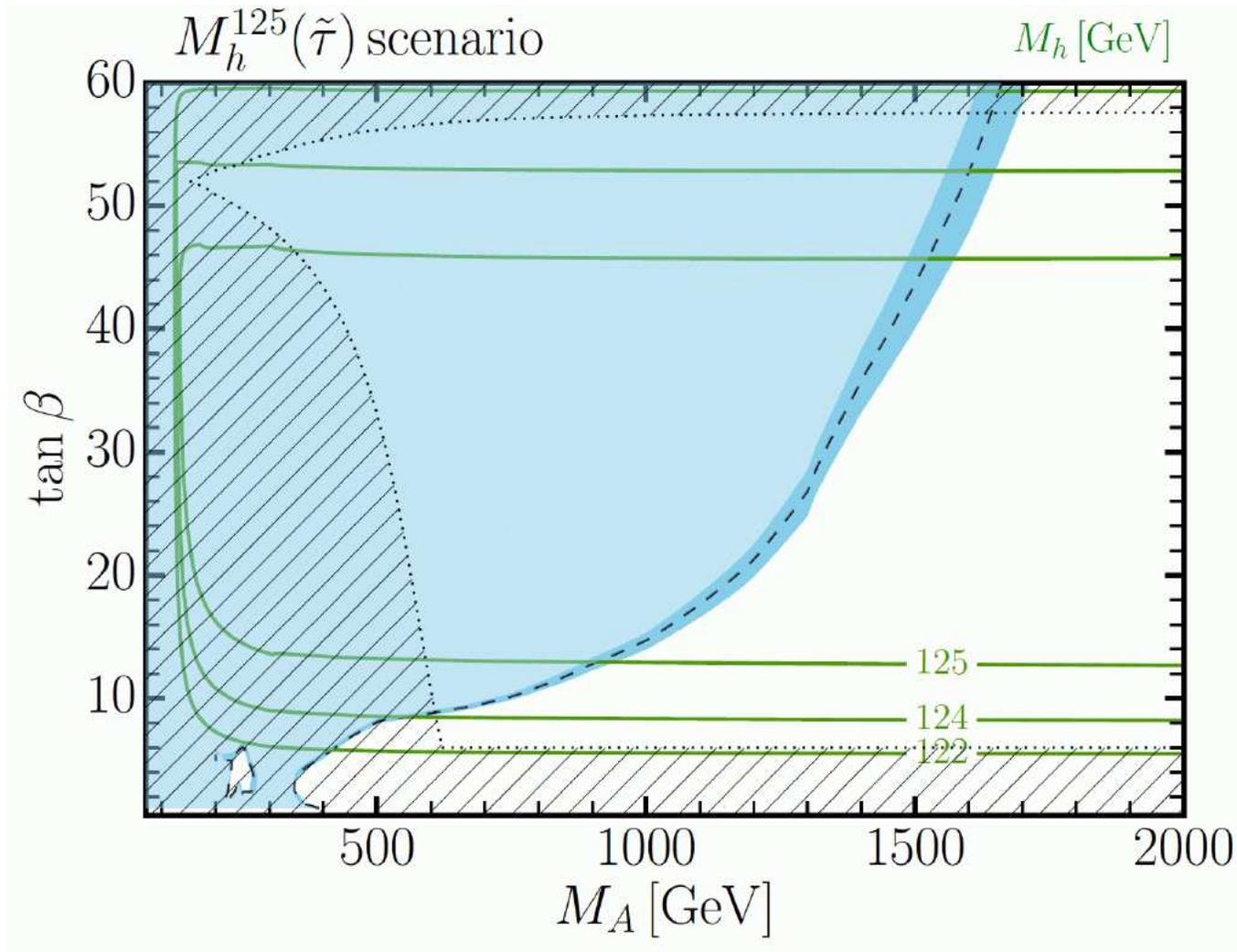
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Data on purpose not taken into account:

- electroweak precision data
- flavor data
- astrophysical data (DM properties)



$$M_{\tilde{Q}_3} = M_{\tilde{U}_3} = M_{\tilde{D}_3} = 1.5 \text{ TeV}$$

$$M_{\tilde{L}_3} = M_{\tilde{E}_3} = 350 \text{ GeV}$$

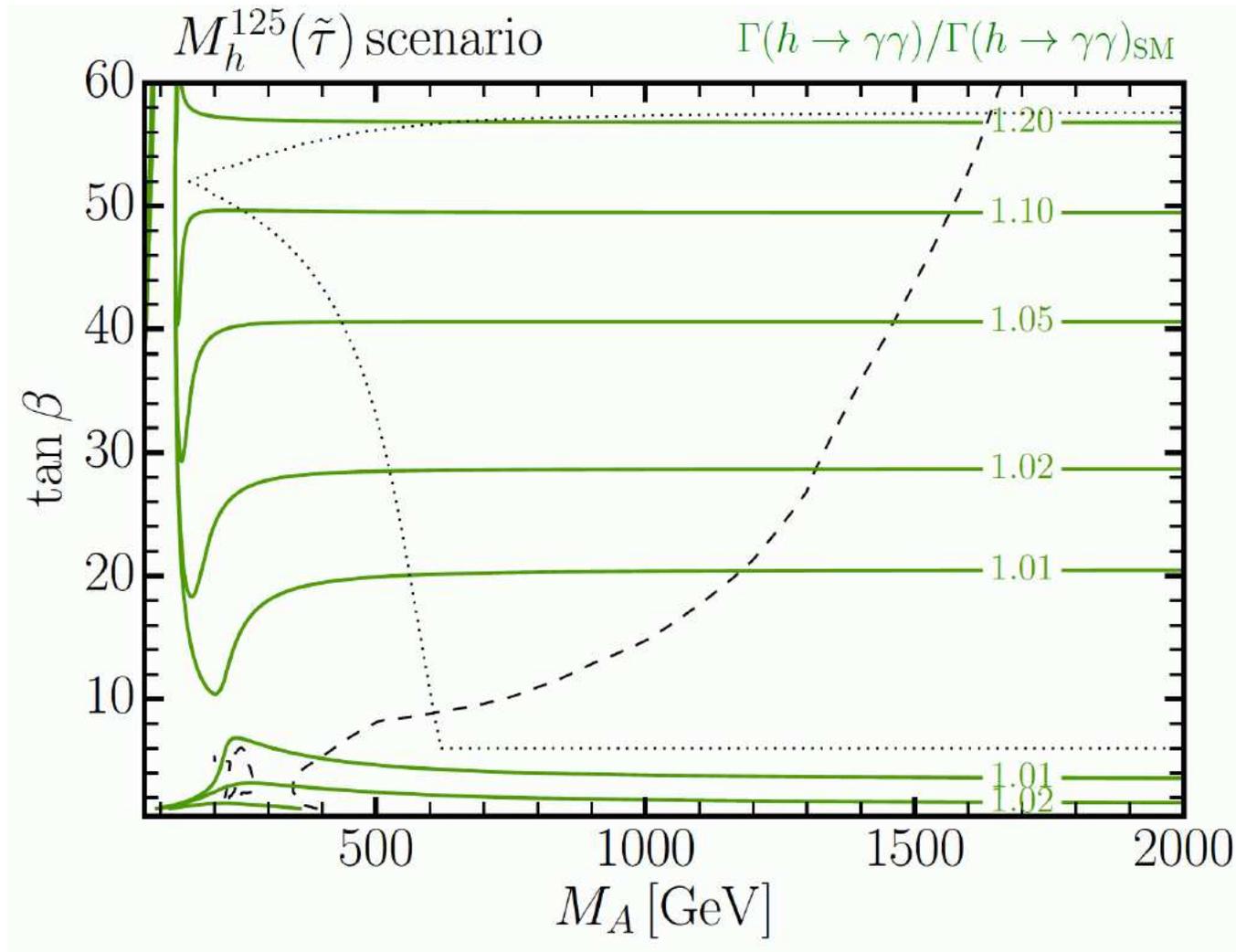
$$\mu = 1 \text{ TeV}, M_1 = 180 \text{ GeV}$$

$$M_2 = 300 \text{ GeV}, M_3 = 2.5 \text{ TeV}$$

$$X_t = 2.8 \text{ TeV}$$

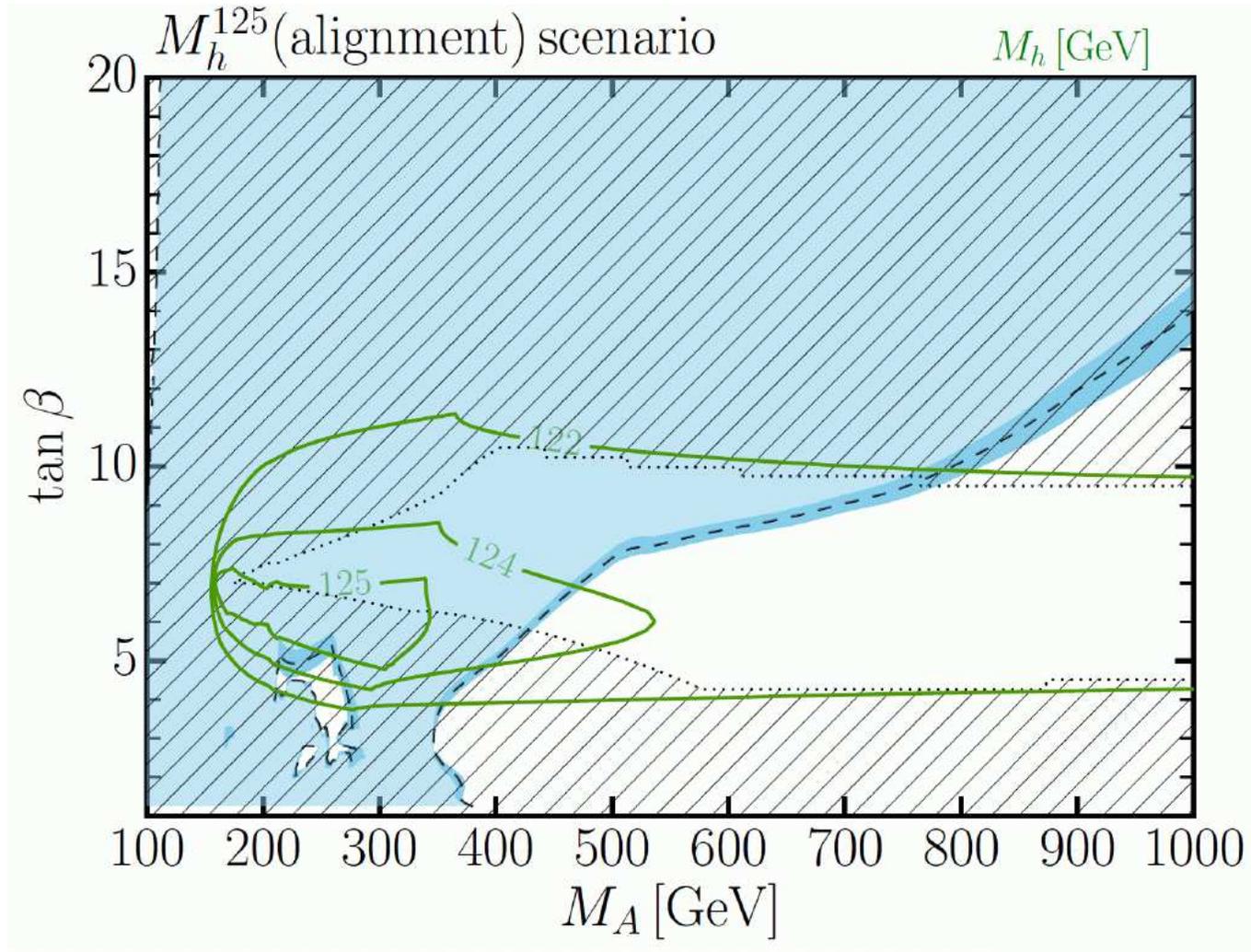
$$A_t = A_b, A_\tau = 800 \text{ GeV}$$

⇒ slightly reduced heavy Higgs coverage



$M_{\tilde{Q}_3} = M_{\tilde{U}_3} = M_{\tilde{D}_3} = 1.5 \text{ TeV}$
 $M_{\tilde{L}_3} = M_{\tilde{E}_3} = 350 \text{ GeV}$
 $\mu = 1 \text{ TeV}, M_1 = 180 \text{ GeV}$
 $M_2 = 300 \text{ GeV}, M_3 = 2.5 \text{ TeV}$
 $X_t = 2.8 \text{ TeV}$
 $A_t = A_b, A_\tau = 800 \text{ GeV}$

\Rightarrow strong impact on $\Gamma(h \rightarrow \gamma\gamma)$



$$M_{\tilde{Q}_3} = M_{\tilde{U}_3} = M_{\tilde{D}_3} = 2.5 \text{ TeV}$$

$$M_{\tilde{L}_3} = M_{\tilde{E}_3} = 2 \text{ TeV}$$

$$\mu = 7.5 \text{ TeV}, M_1 = 500 \text{ GeV}$$

$$M_2 = 1 \text{ TeV}, M_3 = 2.5 \text{ TeV}$$

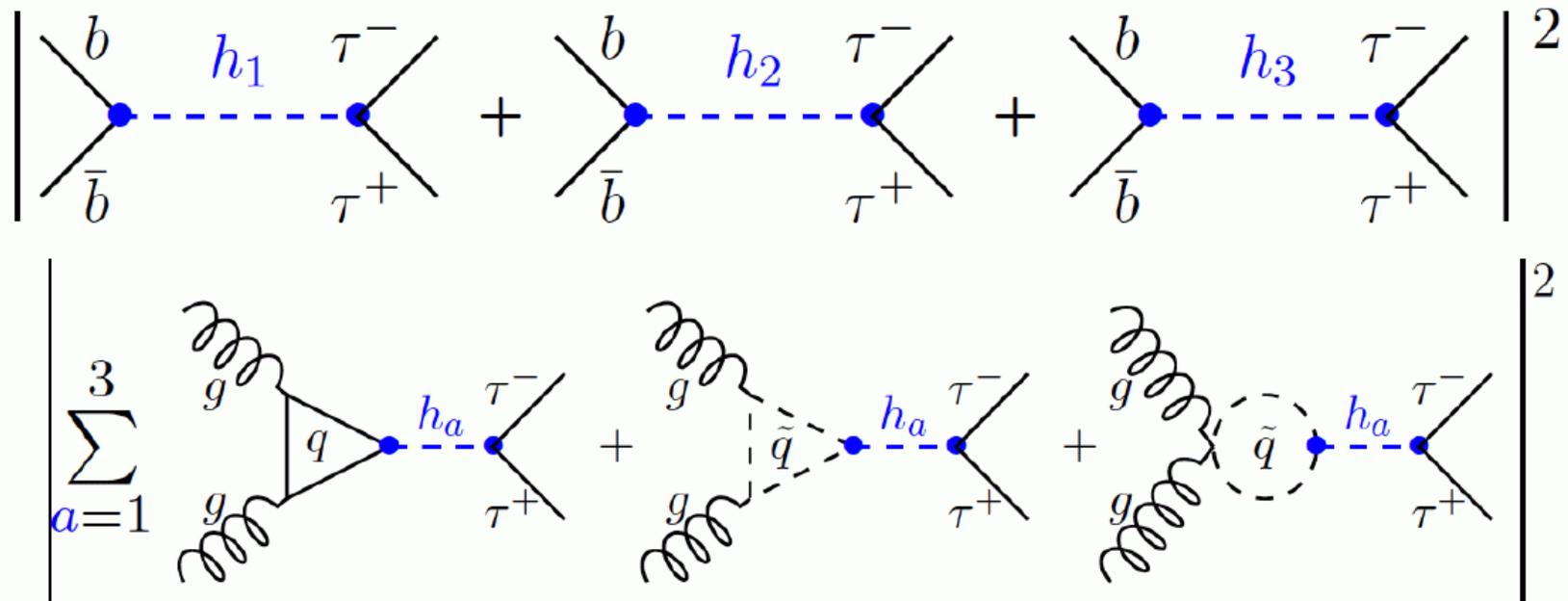
$$A_t = A_b = A_\tau = 6.25 \text{ TeV}$$

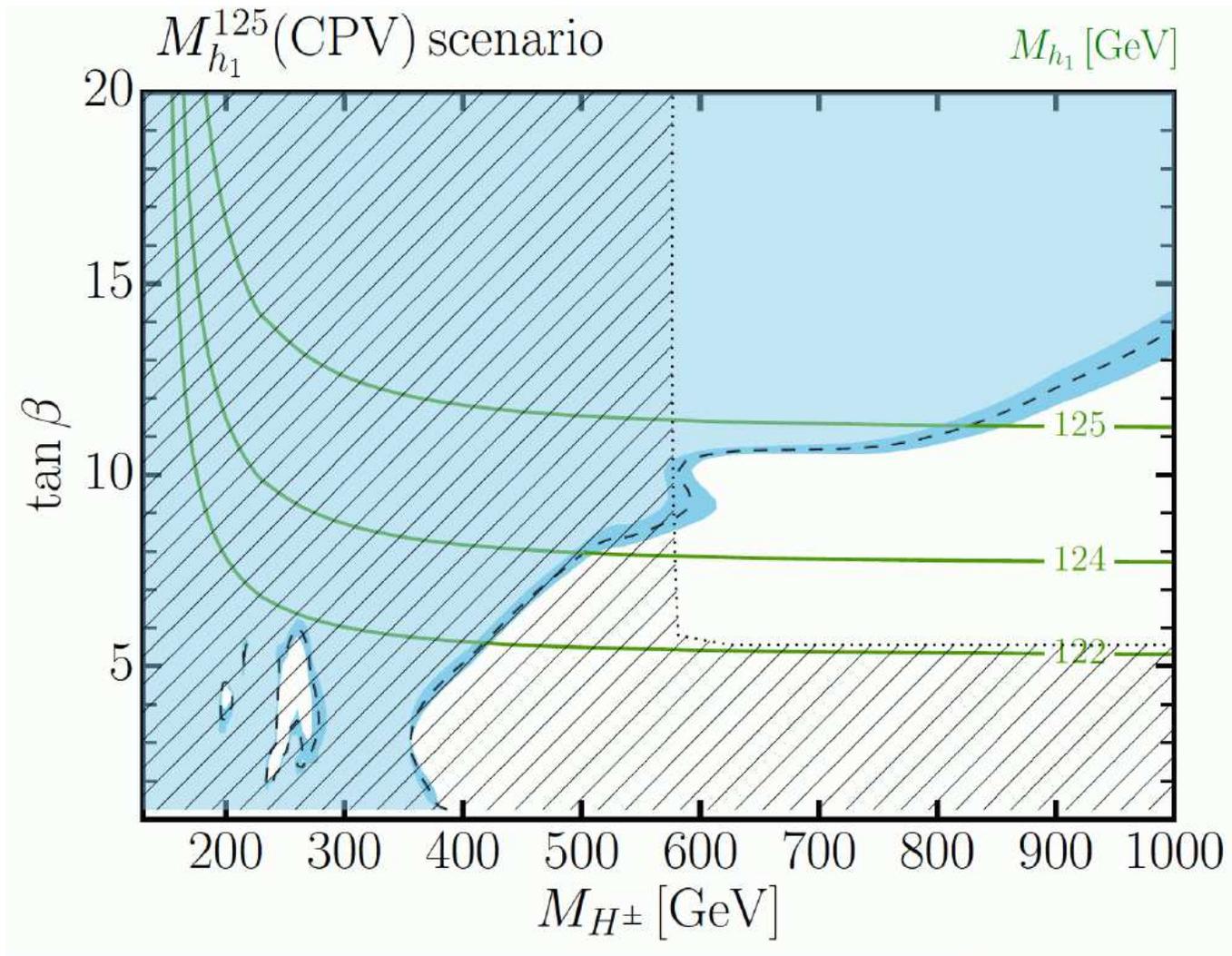
$\Rightarrow h$ SM-like for very low M_A

LHC Higgs searches for complex parameters:

$h_1 \sim H_{125}$, $M_{h_2} \approx M_{h_3}$, CPV: large h_2 - h_3 mixing possible:

Higgs bosons as intermediate states in $\{b\bar{b}, gg\} \rightarrow h_a \rightarrow \tau\tau$





$$M_{\tilde{Q}_3} = M_{\tilde{U}_3} = M_{\tilde{D}_3} = 2 \text{ TeV}$$

$$M_{\tilde{L}_3} = M_{\tilde{E}_3} = 2 \text{ TeV}$$

$$\mu = 1.65 \text{ TeV}, M_1 = 1 \text{ TeV}$$

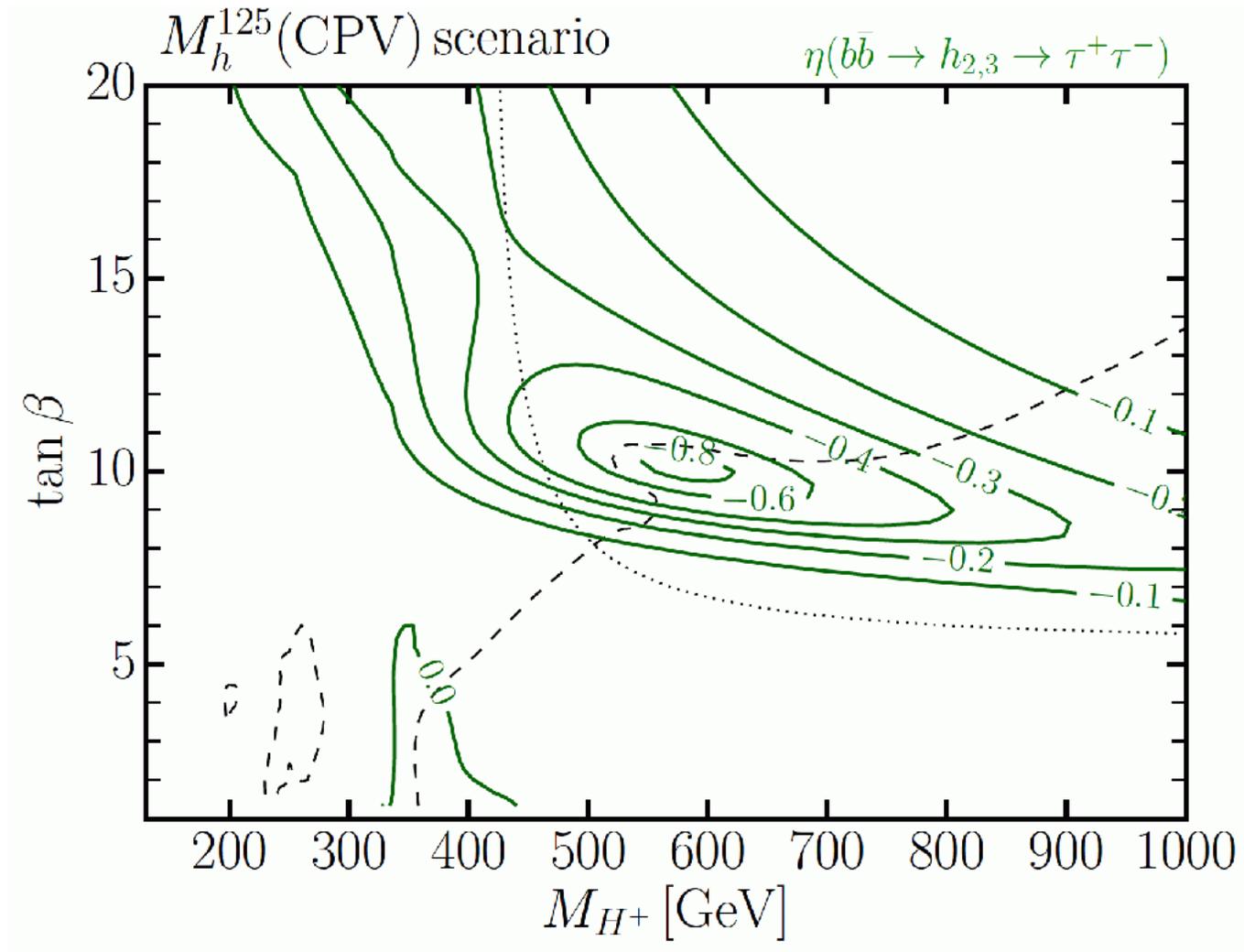
$$M_2 = 1 \text{ TeV}, M_3 = 2.5 \text{ TeV}$$

$$|A_t| = \mu / \tan \beta + 2.8 \text{ TeV}$$

$$\phi_{A_i} = 2/15 \pi$$

$$|A_t| = A_b = A_\tau$$

⇒ reduced coverage due to h_2 - h_3 interference



$$M_{\tilde{Q}_3} = M_{\tilde{U}_3} = M_{\tilde{D}_3} = 2 \text{ TeV}$$

$$M_{\tilde{L}_3} = M_{\tilde{E}_3} = 2 \text{ TeV}$$

$$\mu = 1.65 \text{ TeV}, M_1 = 1 \text{ TeV}$$

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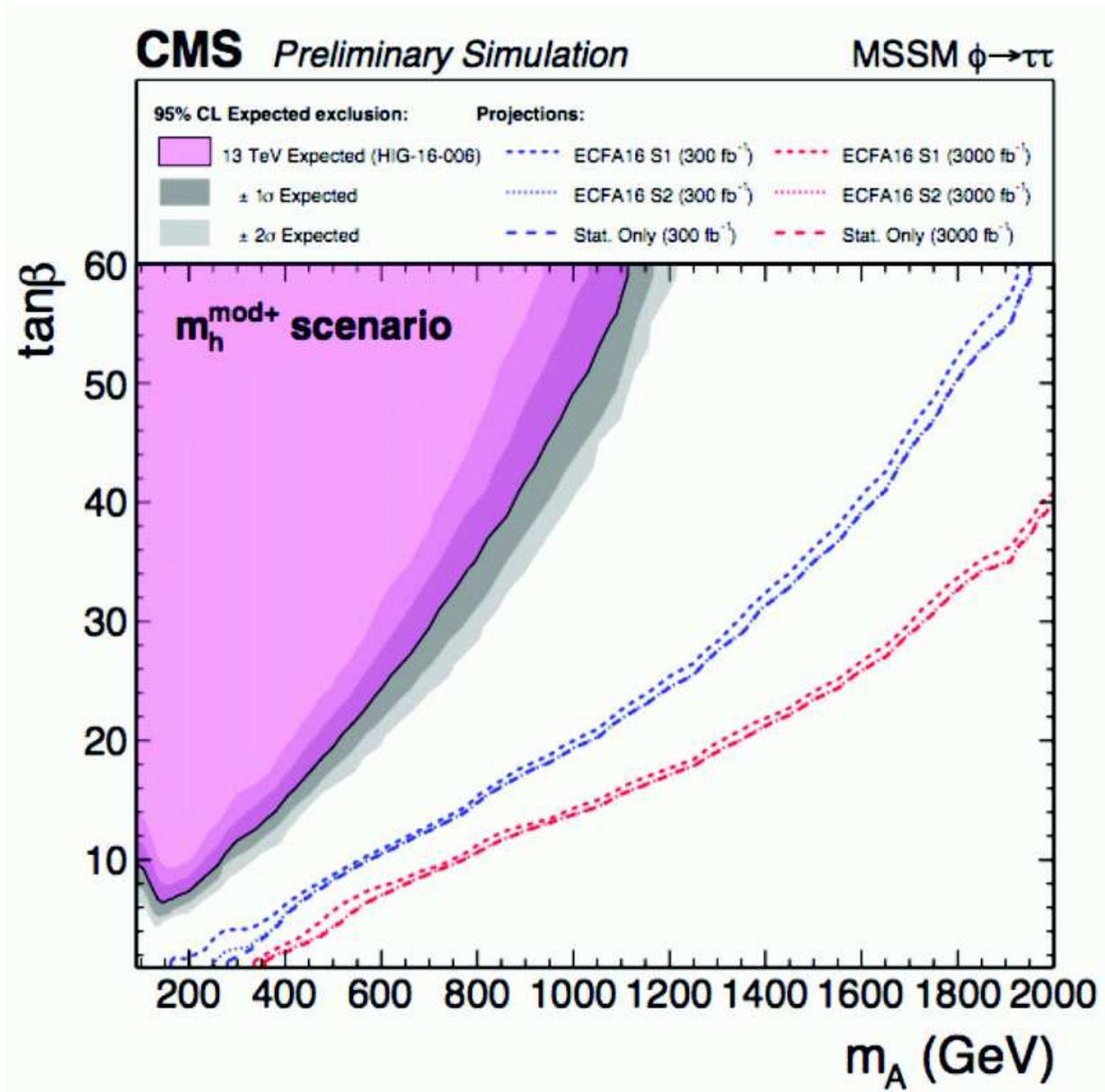
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Future (HL-)LHC projections:

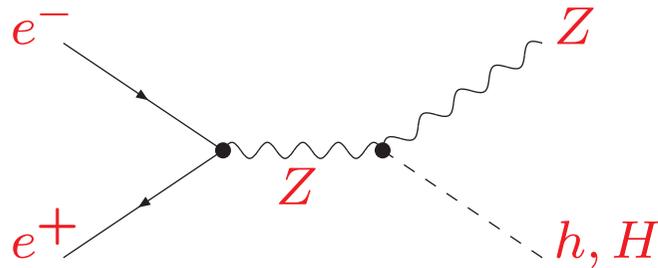


⇒ strong (HL-)LHC limits

Sum rule in the MSSM with h SM-like: $\sin(\beta - \alpha) \approx 1$, $\cos(\beta - \alpha) \approx 0$

Search for neutral SUSY Higgs bosons:

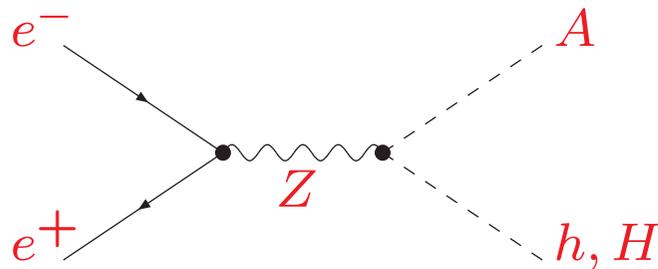
$e^+e^- \rightarrow Zh, ZH$



$$\sigma_{hZ} \approx \sin^2(\beta - \alpha_{\text{eff}}) \sigma_{hZ}^{\text{SM}}$$

$$\sigma_{HZ} \approx \cos^2(\beta - \alpha_{\text{eff}}) \sigma_{hZ}^{\text{SM}}$$

$e^+e^- \rightarrow Ah, AH$

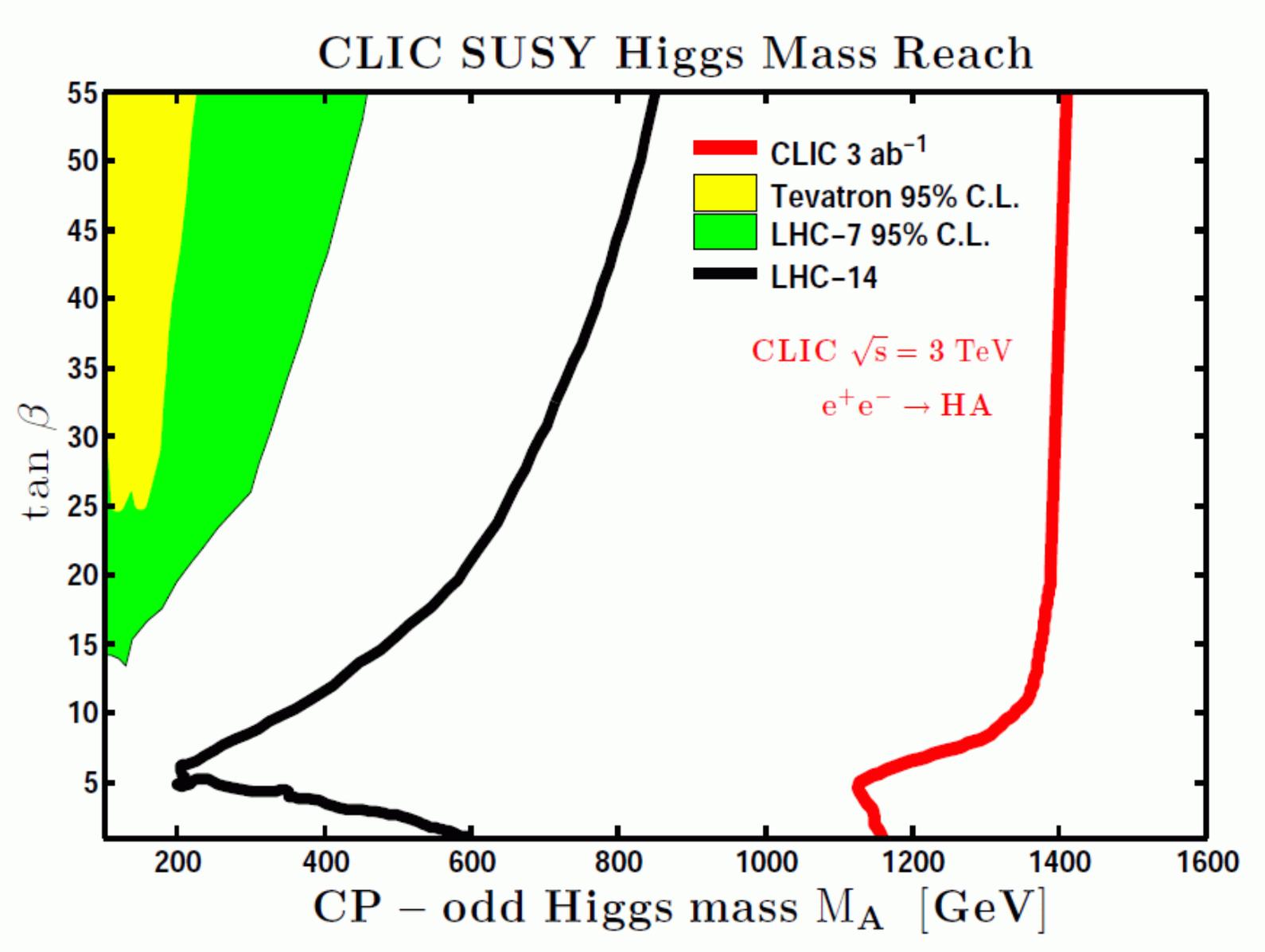


$$\sigma_{hA} \propto \cos^2(\beta - \alpha_{\text{eff}}) \sigma_{hZ}^{\text{SM}}$$

$$\sigma_{HA} \propto \sin^2(\beta - \alpha_{\text{eff}}) \sigma_{hZ}^{\text{SM}}$$

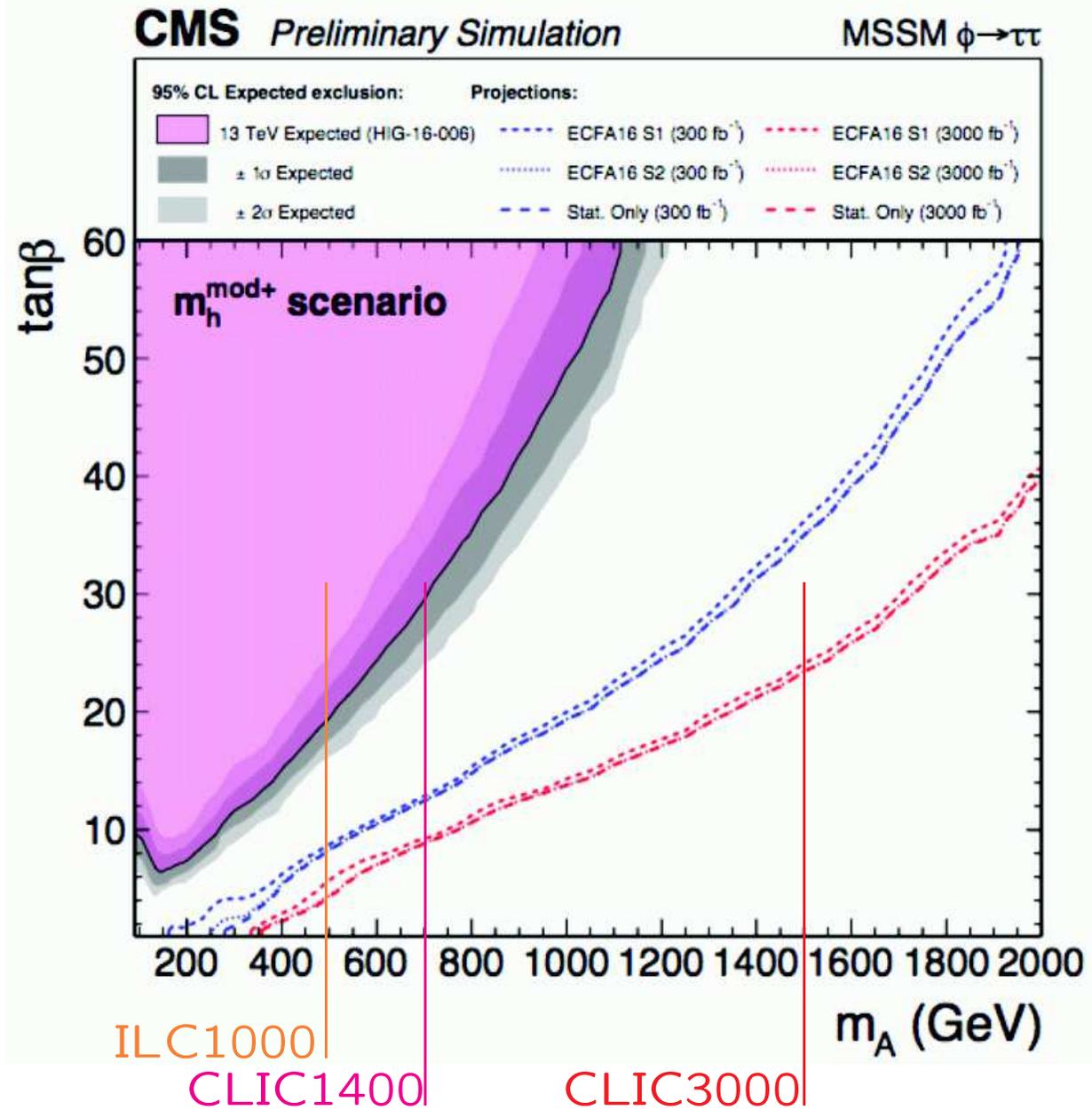
\Rightarrow only pair production of heavy Higgs bosons!

reach: $M_A \lesssim \sqrt{s}/2$



⇒ close to kinematic limit

“Simple” LC reach in the MSSM (neglecting $t\bar{t}$ final states)



⇒ unique opportunities!