

LHC-Friendly Avenues for Freeze-in Dark Matter

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Based on 1811.05478, 1906.xxxxx

(w. Patrick Tunney & Bryan Zaldivar)

RISE Multi-Higgs 2019, Helsinki
28/05/19



Outline

- **Freeze-In Production of Dark Matter** (*in a nutshell*)
- **LHC-friendly Freeze-In DM Models**
(A Model of Dark Matter Freeze-In via the Higgs Field)
- **Long-Lived Particle searches: LHC and MATHUSLA**

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→ Freeze-In Production of Dark Matter (*in a nutshell*)

→ LHC-friendly Freeze-In DM Models

(A Model of Dark Matter Freeze-In via the **Higgs Field**)



→ Long-Lived Particle searches: LHC and MATHUSLA

[source: PhD Comics]

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- Long-Lived Particle searches: **LHC** and **MATISLA**



Dark Matter Freeze-In in a Nutshell

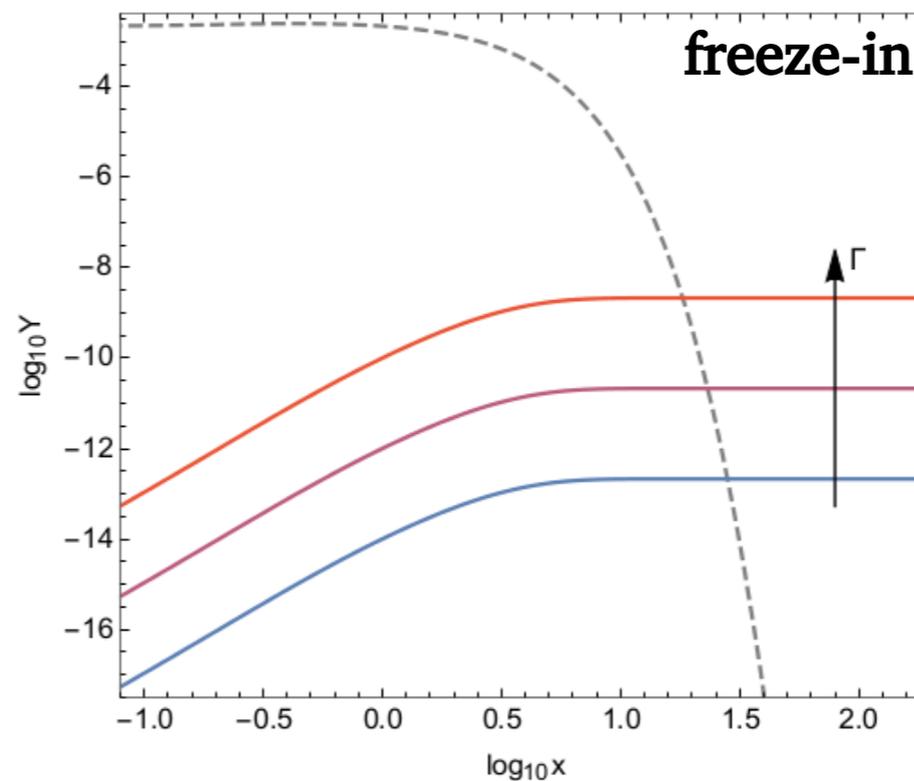
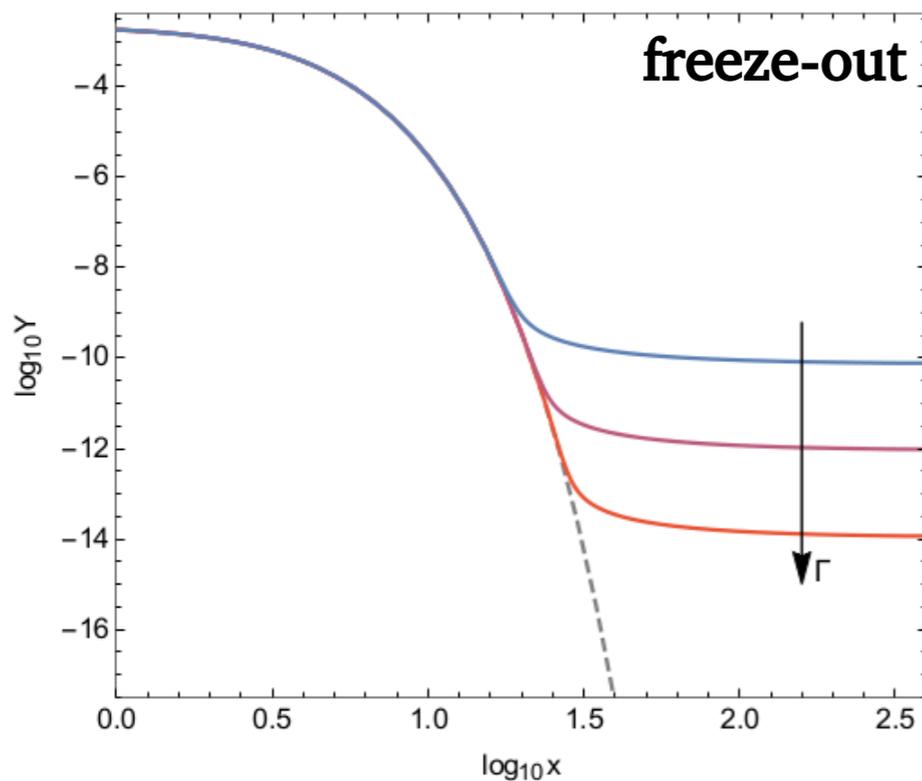
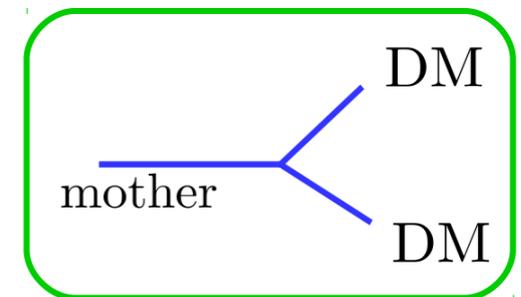
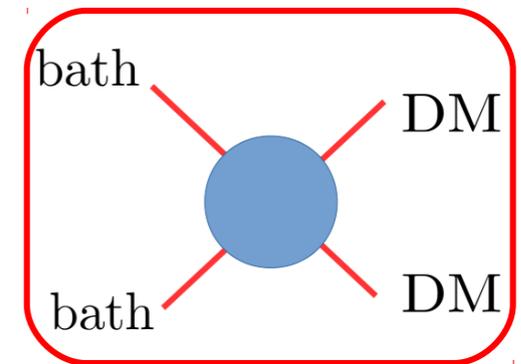
Hall et al, 0911.1120

$$\frac{dn_{\text{DM}}}{dt} + 3Hn_{\text{DM}} = n_{\text{bath}}\Gamma_{\text{prod}} - n_{\text{DM}}\Gamma_{\text{ann}}$$

negligible


interaction rates


- Initial negligible DM abundance (+ DM never thermalizes)
- Production from **scatterings** or **decays**



$$Y = n_{\text{DM}}/s$$

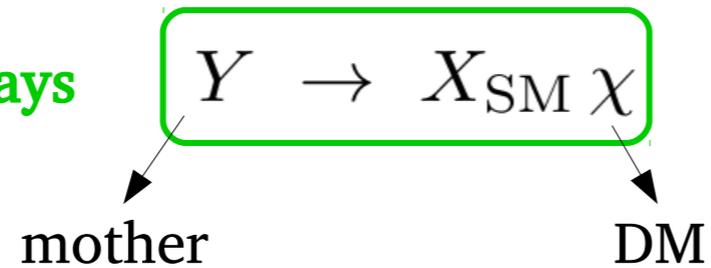
$$x = m/T$$

Bernal et al, 1706.07442

Dark Matter Freeze-In in a Nutshell

Hall et al, 0911.1120

→ Freeze-in from **decays**

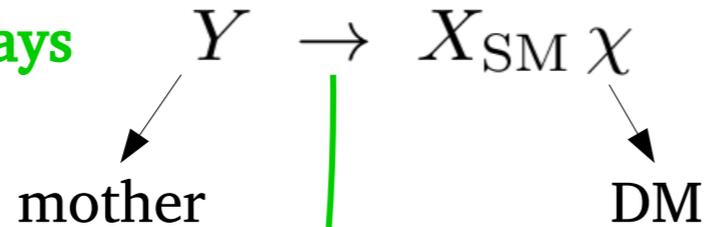


$$\frac{dn_\chi}{dt} + 3H n_\chi = \Gamma_Y n_Y^{\text{eq}} \frac{K_1(m_Y/T)}{K_2(m_Y/T)} \longrightarrow Y_\chi \approx \frac{135 g_Y \xi M_{\text{Pl}} \Gamma_Y}{8\pi^3 m_Y^2} \frac{1}{1.66 g_*(m_Y/3) \sqrt{g_*(m_Y/3)}}$$

Dark Matter Freeze-In in a Nutshell

Hall et al, 0911.1120

→ Freeze-in from **decays**



$$\frac{dn_\chi}{dt} + 3H n_\chi = \Gamma_Y n_Y^{\text{eq}} \frac{K_1(m_Y/T)}{K_2(m_Y/T)} \longrightarrow Y_\chi \approx \frac{135 g_Y \xi M_{\text{Pl}} \Gamma_Y}{8\pi^3 m_Y^2} \frac{1}{1.66 g_*(m_Y/3) \sqrt{g_*(m_Y/3)}}$$

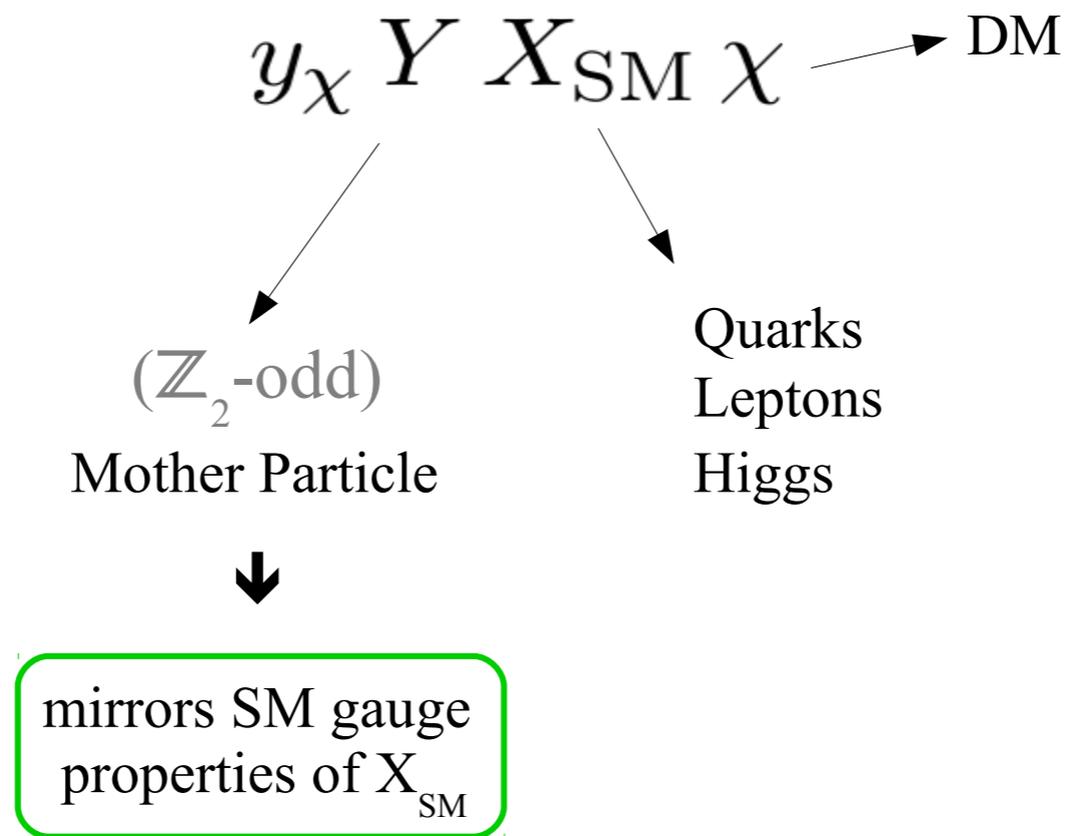
DM Relic Density: $\lambda \sim (10^{-12} - 10^{-13}) \sqrt{\frac{m_Y}{m_\chi}}$

Long-lived particle signatures (e.g. LHC) from Y decays!

$$c\tau_Y \sim 3.6 m \left(\frac{0.12}{\Omega_\chi h^2} \right) \left(\frac{m_\chi}{100 \text{ KeV}} \right) \left(\frac{300 \text{ GeV}}{m_Y} \right)^2$$

Simple Freeze-In DM Models

Belanger et al, 1811.05478



Freeze-In DM from the Higgs

\mathbb{Z}_2 -odd sector

SU(2) Singlet SU(2) Doublet

“Higgs Portal”

$$\mathcal{L} = \mathcal{L}_{\text{SM}} + i \bar{\chi} \gamma^\mu \partial_\mu \chi + i \bar{\psi} \gamma^\mu D_\mu \psi - m_s \bar{\chi} \chi - m_D \bar{\psi} \psi - \boxed{y_\chi \bar{\psi} H \chi + h.c.}$$

→ Simple DM version of complete setups (e.g. Higgsino-Axino, Higgsino-Singlino)

Co et al, 1506.07532

See also: [Calibbi et al, 1505.03867](#) (freeze-out)
[Calibbi et al, 1805.04423](#) (freeze-in)

Freeze-In DM from the Higgs

\mathbb{Z}_2 -odd sector

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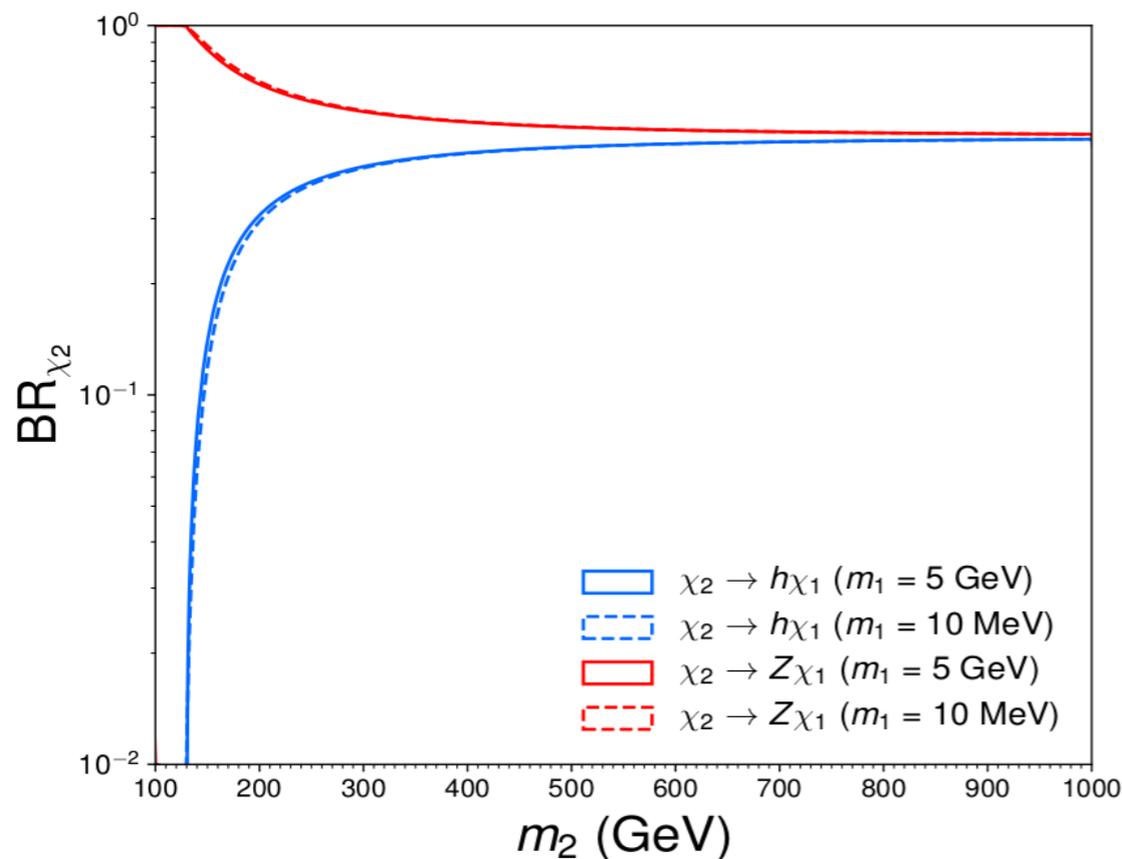
$$\mathcal{L} = \mathcal{L}_{\text{SM}} + i \bar{\chi} \gamma^\mu \partial_\mu \chi + i \bar{\psi} \gamma^\mu D_\mu \psi - m_s \bar{\chi} \chi - m_D \bar{\psi} \psi - \boxed{y_\chi \bar{\psi} H \chi + h.c.}$$

➔ Simple DM version of complete setups (e.g. Higgsino-Axino, Higgsino-Singlino)

Co et al, 1506.07532

➔ Tiny singlet-doublet mixing (DM is singlet-like)

$$\sin \theta \simeq \frac{y_\chi v}{\sqrt{2}(m_2 - m_1)}$$



Decays of χ_2

$$\boxed{\begin{aligned} \chi_2 &\rightarrow h\chi_1 \\ \chi_2 &\rightarrow Z\chi_1 \end{aligned}}$$

Freeze-In DM from the Higgs

\mathbb{Z}_2 -odd sector

SU(2) Singlet SU(2) Doublet “Higgs Portal”

$$\mathcal{L} = \mathcal{L}_{\text{SM}} + i \bar{\chi} \gamma^\mu \partial_\mu \chi + i \bar{\psi} \gamma^\mu D_\mu \psi - m_s \bar{\chi} \chi - m_D \bar{\psi} \psi - y_\chi \bar{\psi} H \chi + h.c.$$

→ Simple DM version of complete setups (e.g. Higgsino-Axino, Higgsino-Singlino)

Co et al, 1506.07532

→ Also note the **charged component** of SU(2) doublet ψ^\pm

Radiative mass splitting

$$\delta m = m_{\psi^\pm} - m_2 \in [260, 340] \text{ MeV}$$

Decays of ψ^\pm

$\psi^\pm \rightarrow \pi^\pm \chi_2$

 $\psi^\pm \rightarrow \ell^\pm \nu \chi_2$
 $\psi^\pm \rightarrow W^\pm \chi_1$

↖ Dominant decay

↘ Relevant for DM production

Freeze-In DM from the Higgs

\mathbb{Z}_2 -odd sector

SU(2) Singlet SU(2) Doublet “Higgs Portal”

$$\mathcal{L} = \mathcal{L}_{\text{SM}} + i \bar{\chi} \gamma^\mu \partial_\mu \chi + i \bar{\psi} \gamma^\mu D_\mu \psi - m_s \bar{\chi} \chi - m_D \bar{\psi} \psi - y_\chi \bar{\psi} H \chi + h.c.$$

→ Simple DM version of complete setups (e.g. Higgsino-Axino, Higgsino-Singlino)

Co et al, 1506.07532

→ Also note the charged component of SU(2) doublet ψ^\pm

Freeze-in Production of DM

(assume standard thermal history)

$$\Omega_1 h^2 \simeq \frac{m_1}{\rho_c/s_0} \frac{135 M_{\text{Pl}} \Gamma_{\text{FI}}}{2 \times 1.66 \pi^3 m_2^2 [g_*(m_2/3)]^{3/2}} \longrightarrow \Gamma_{\text{FI}} = \Gamma(\chi_2 \rightarrow h\chi_1) + \Gamma(\chi_2 \rightarrow Z\chi_1) + \Gamma(\psi^\pm \rightarrow W^\pm\chi_1)$$

↙

$$\rho_c/s_0 = 3.6 \times 10^{-9} \text{ GeV}$$

Constraints from Cosmology

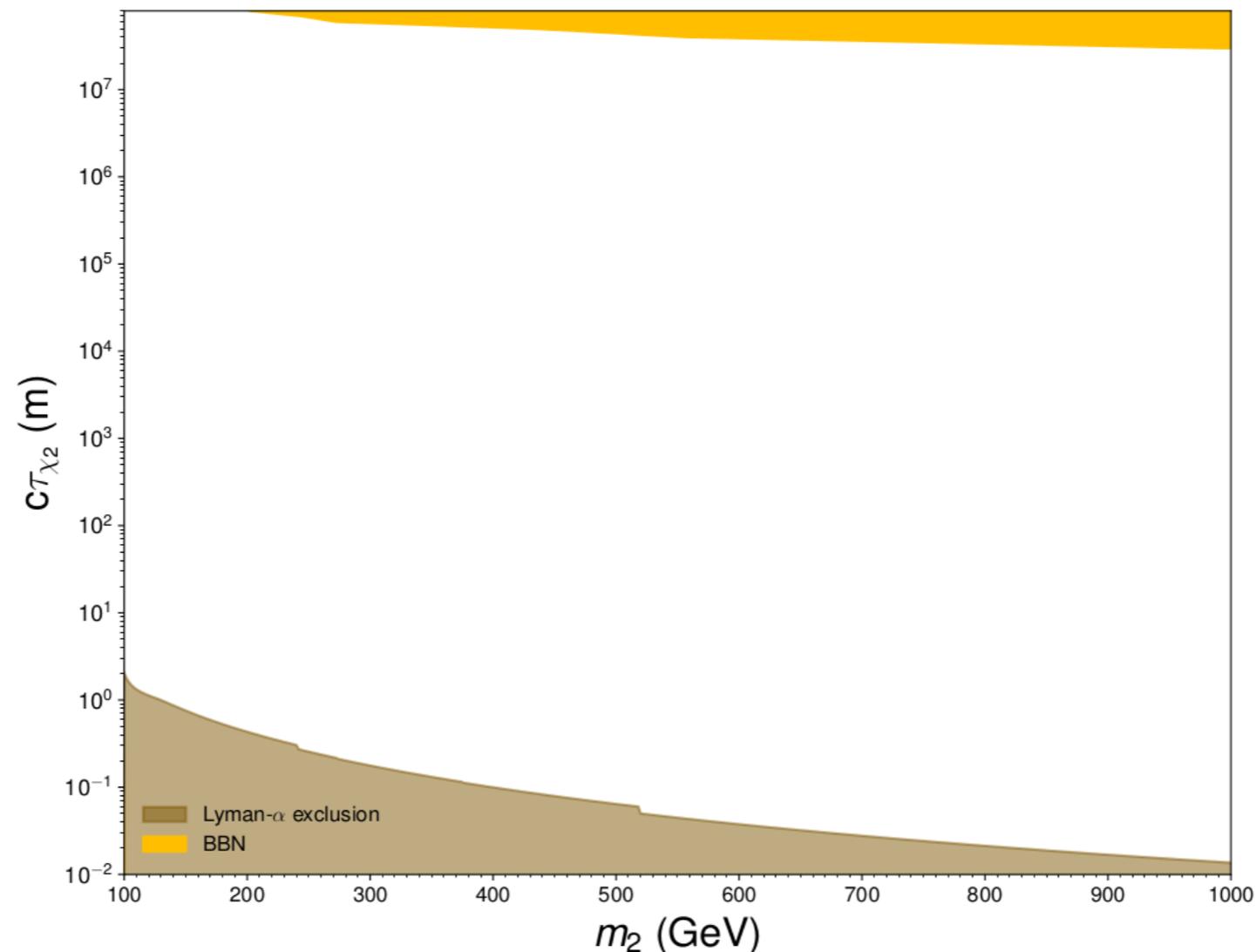
→ Lyman- α forest observations

Constrain washout of small-scale structure by partially relativistic (“*warm*”) DM

→ Big-bang nucleosynthesis

If long-lived, χ_2 visible decay products can affect BBN predictions

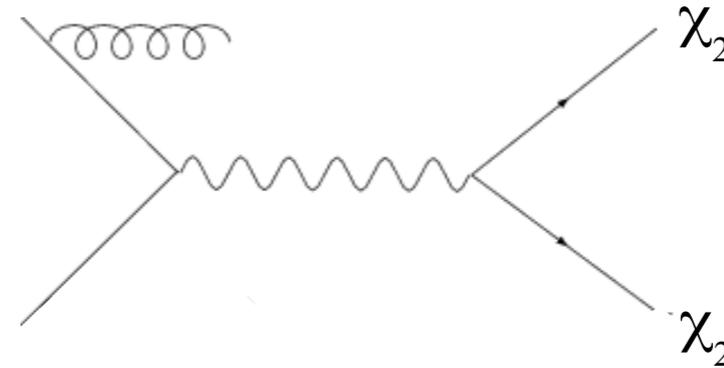
Constrain (hadronic) energy injection from the decay of χ_2



Probing Freeze-in at LHC

→ Mono-X (mono-jet)

Consider χ_2 as effectively stable on LHC



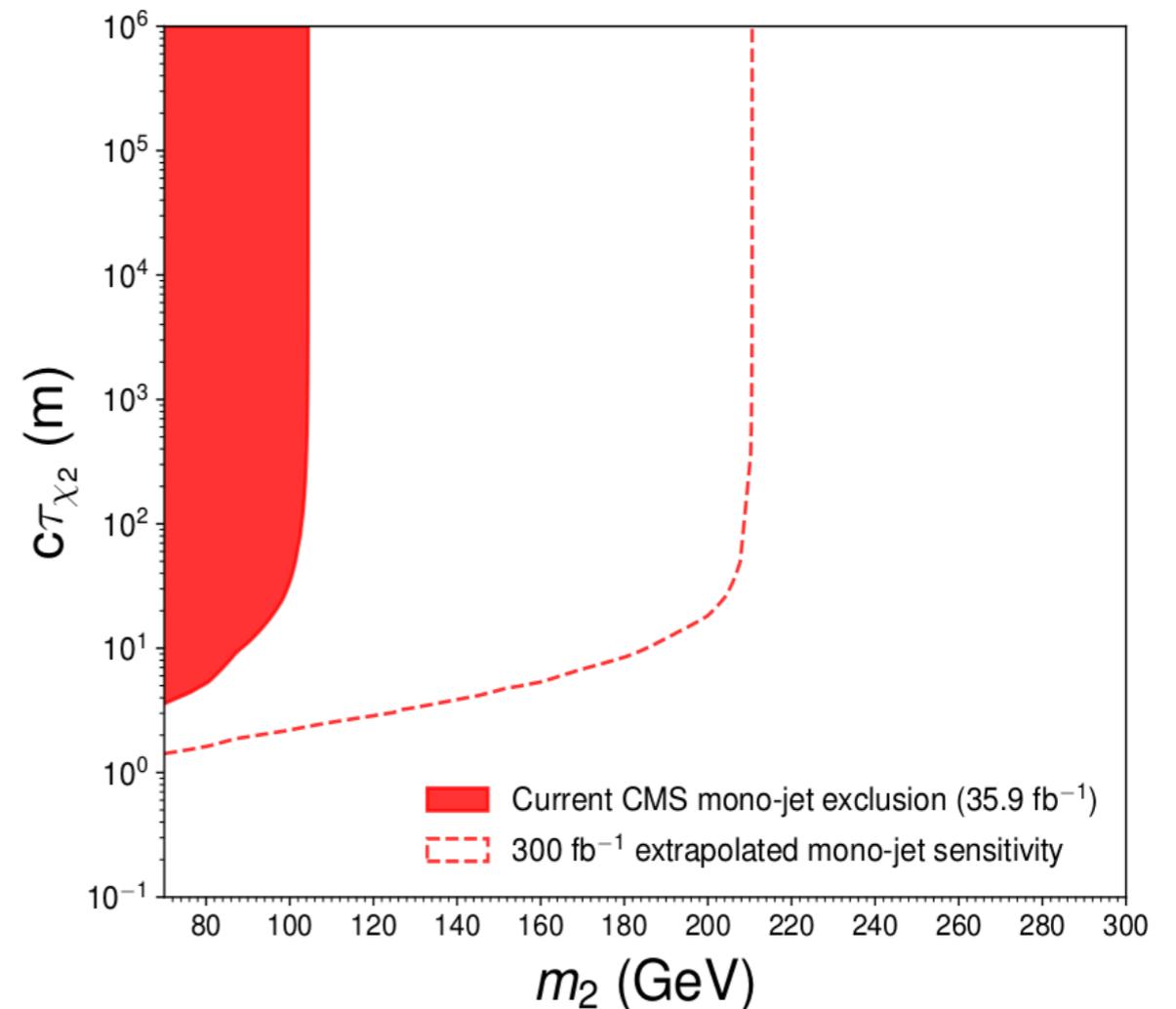
(Re)interpret CMS bounds for DM simplified models with vector mediator with $m = m_Z$

CMS [1712.02345](#) (35.9 fb^{-1})

Search for new physics in final states with an energetic jet or a hadronically decaying W or Z boson and transverse momentum imbalance at $\sqrt{s} = 13 \text{ TeV}$

The CMS Collaboration*

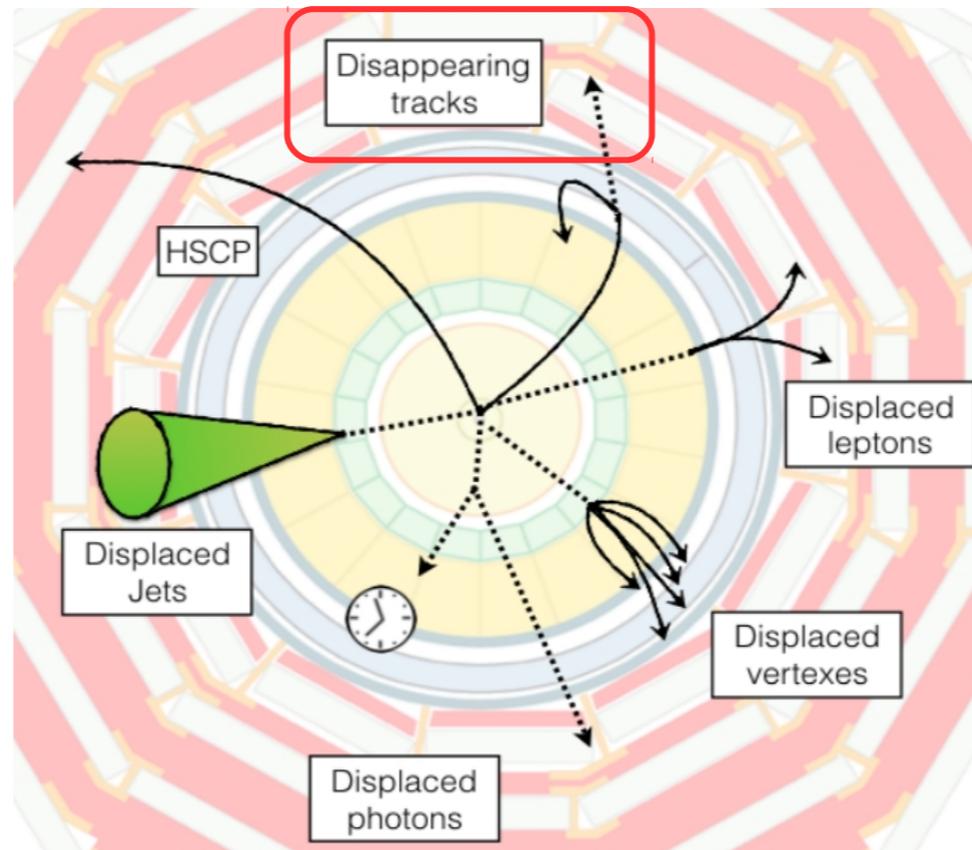
$$\mathcal{L} \supset Z'_\mu (g_{\text{SM}} \bar{q} \gamma^\mu q + g_\chi \bar{\chi} \gamma^\mu \chi)$$



Probing Freeze-in at LHC

→ Disappearing tracks: $\psi^\pm \rightarrow \pi^\pm \chi_2$ $\delta m = m_{\psi^\pm} - m_2 \in [260, 340]$ MeV
few *cm* decay length

ATLAS disappearing track analysis for pure Higgsino production
[ATL-PHYS-PUB-2017-019](#) (36.1 fb⁻¹) $m_{\psi^\pm} > 145$ GeV

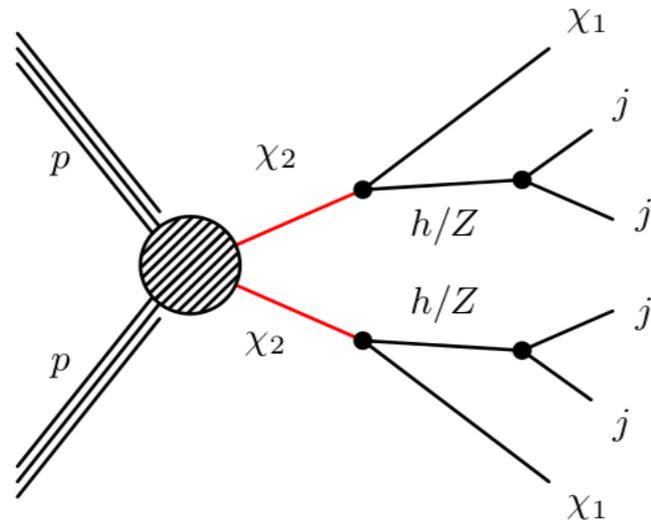


[Courtesy of Livia Soffi, CMS Collaboration]

Probing Freeze-in at LHC

→ Displaced vertices (jets) + MET

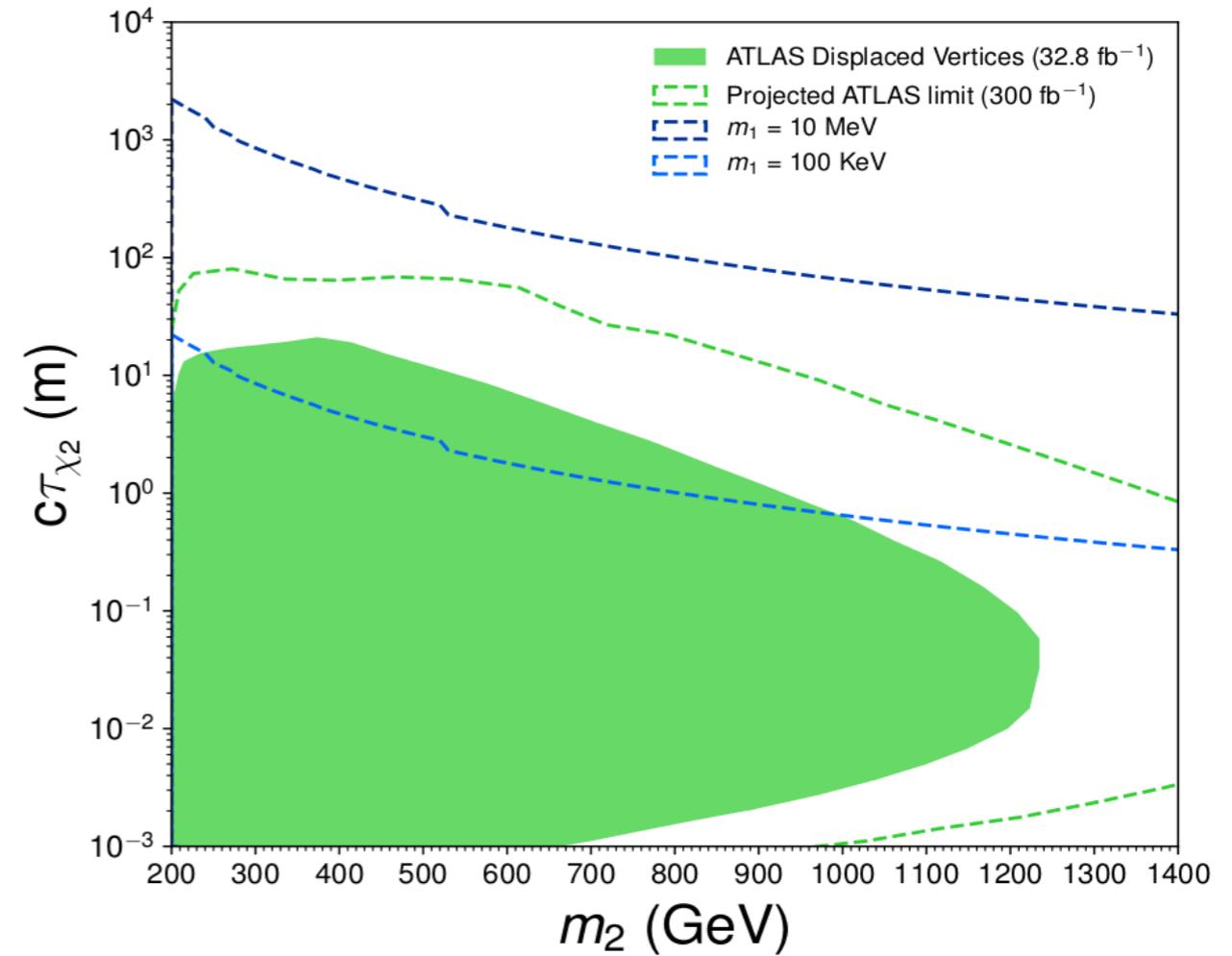
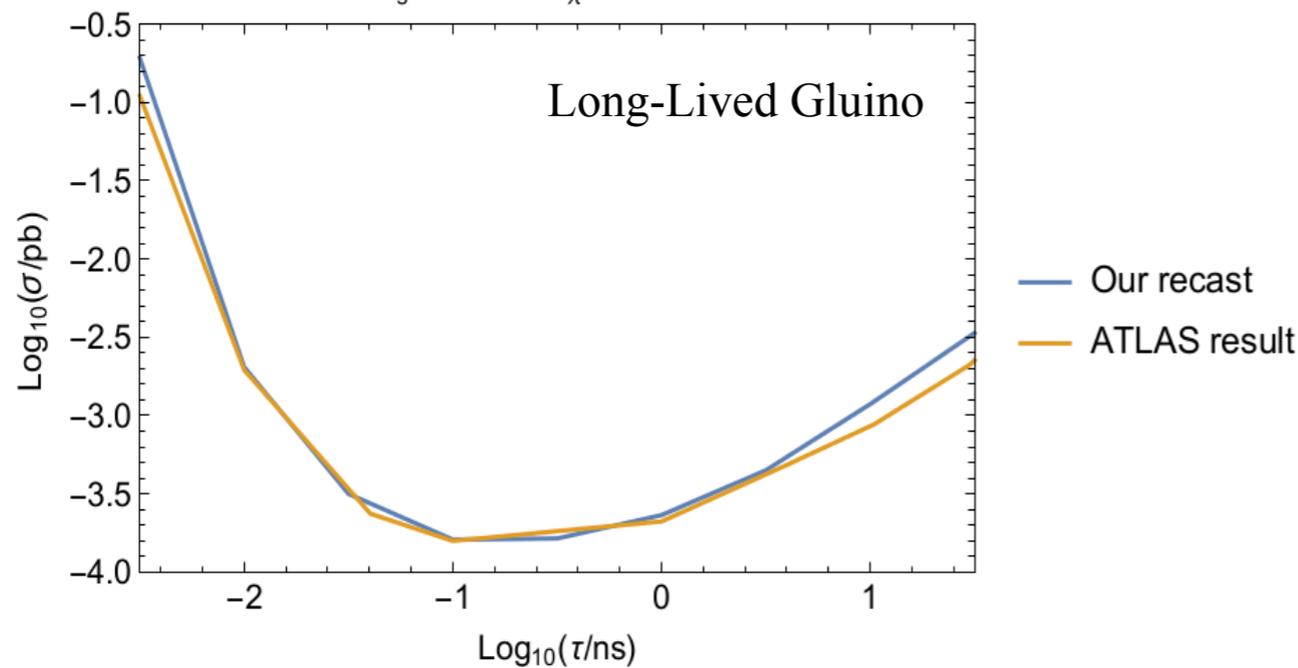
ATLAS 1710.04901 (32.8 fb⁻¹)



Calibbi et al, 1805.04423

JMN, Tunney, Zaldivar, 1906.xxxxx

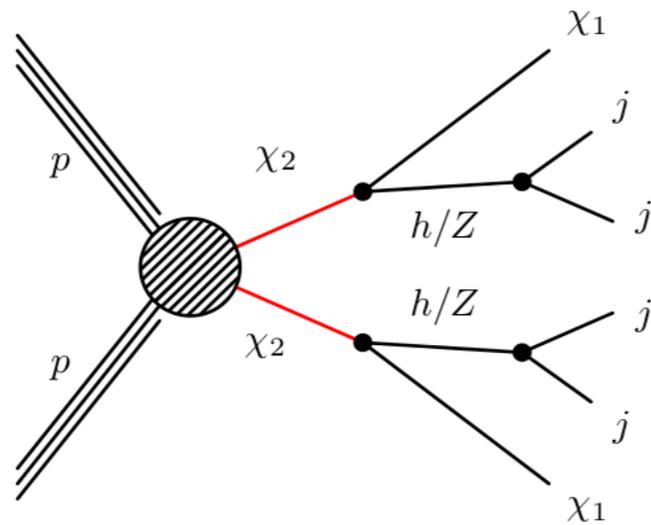
$m_{\tilde{g}} = 2 \text{ TeV}, m_{\tilde{\chi}} = 100 \text{ GeV}$



Probing Freeze-in at LHC

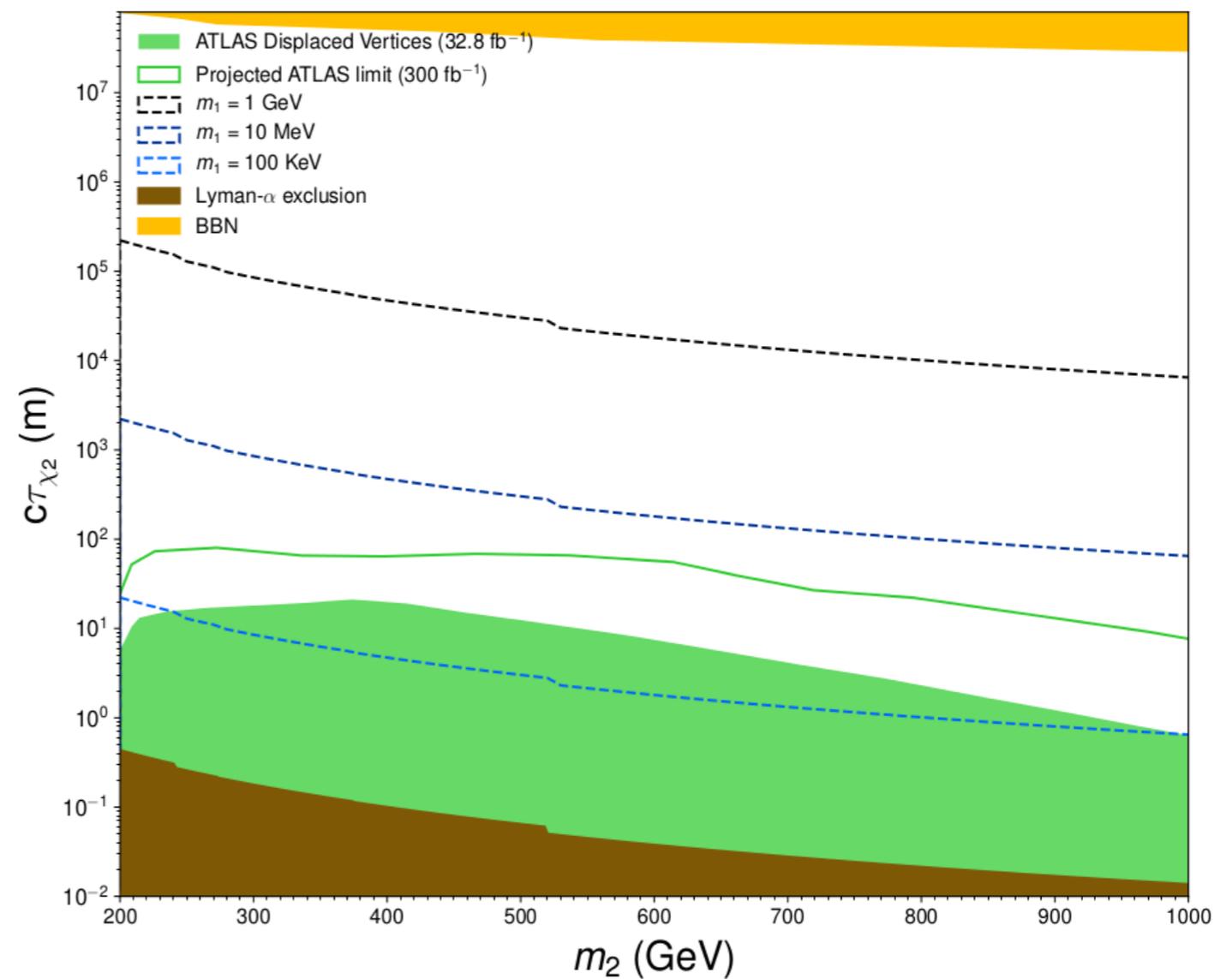
→ Displaced vertices (jets) + MET

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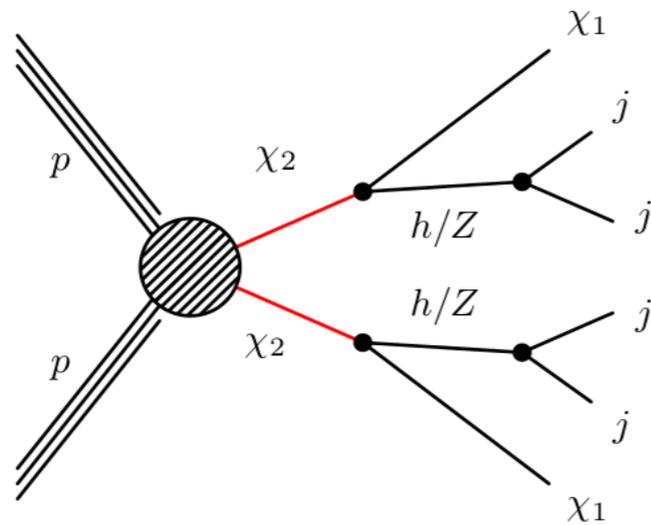
JMN, Tunney, Zaldivar, 1906.xxxxx



Probing Freeze-in at LHC

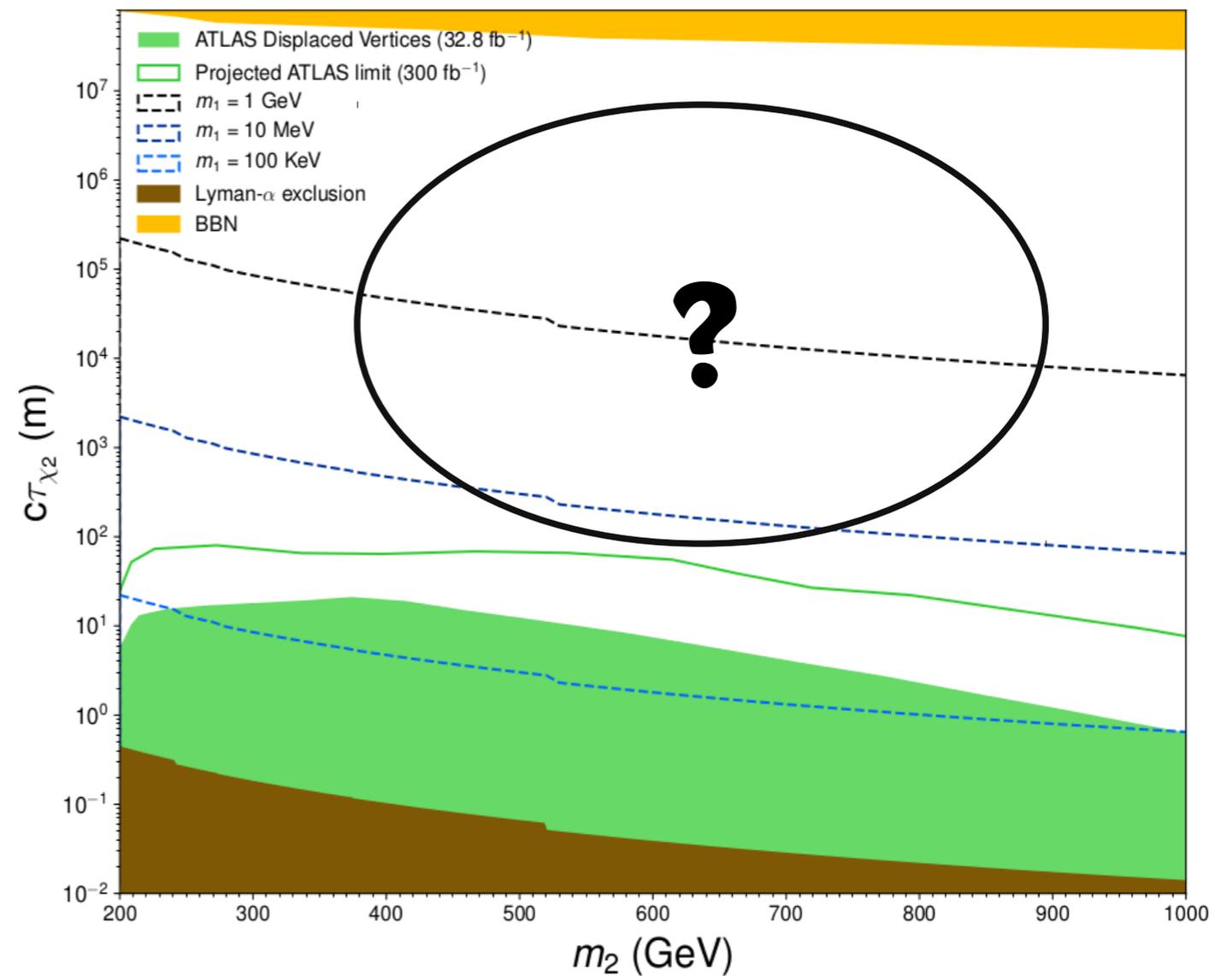
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Calibbi et al, 1805.04423

JMN, Tunney, Zaldivar, 1906.xxxxx

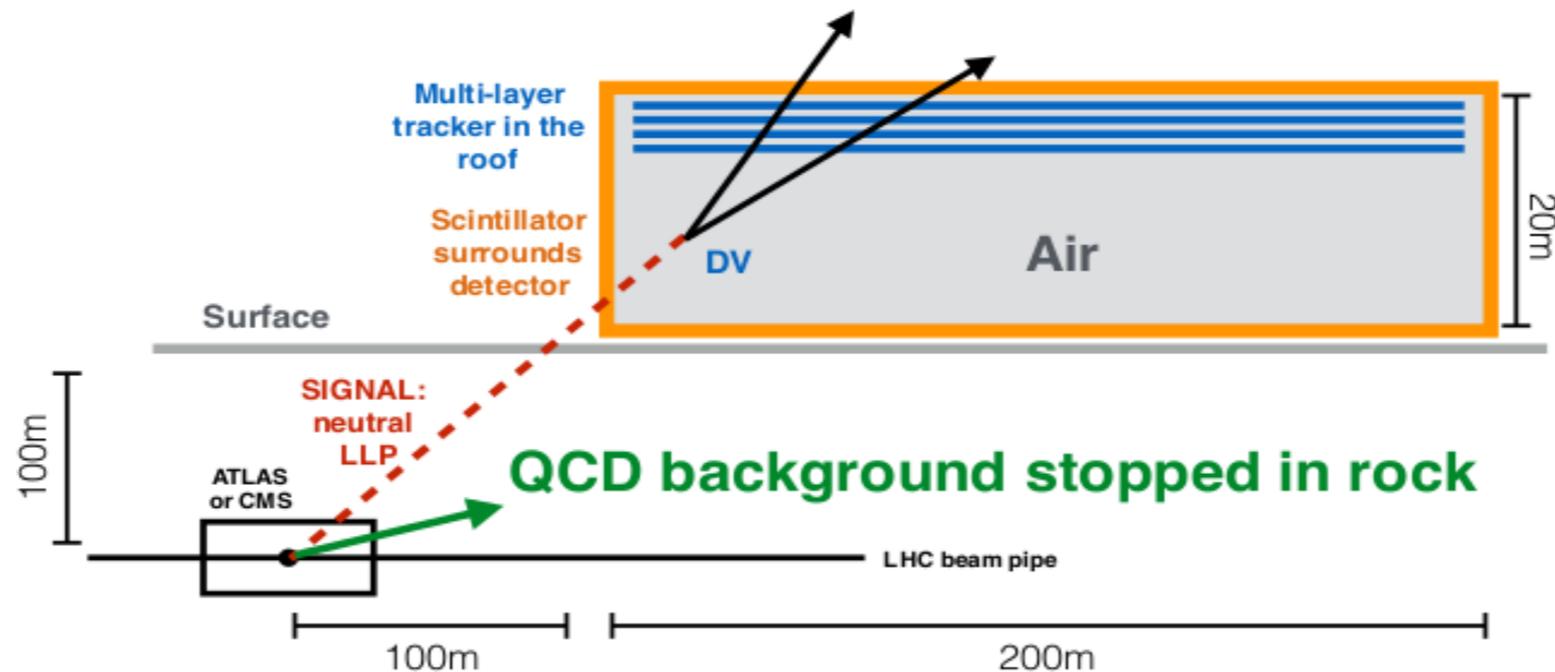


Probing Freeze-in at MATHUSLA

Courtesy of
David Curtin

MATHUSLA

An external LLP detector for the HL- or HE-LHC



Chou, DC, Lubatti
1606.06298

Reliance on well-understood technology (RPC, plastic scintillators) means this could be implemented in time for the HL-LHC. **But design not set in stone, will explore other options!**

Unofficial cost estimates of current design:
~ 50 million USD

Probing Freeze-in at MATHUSLA

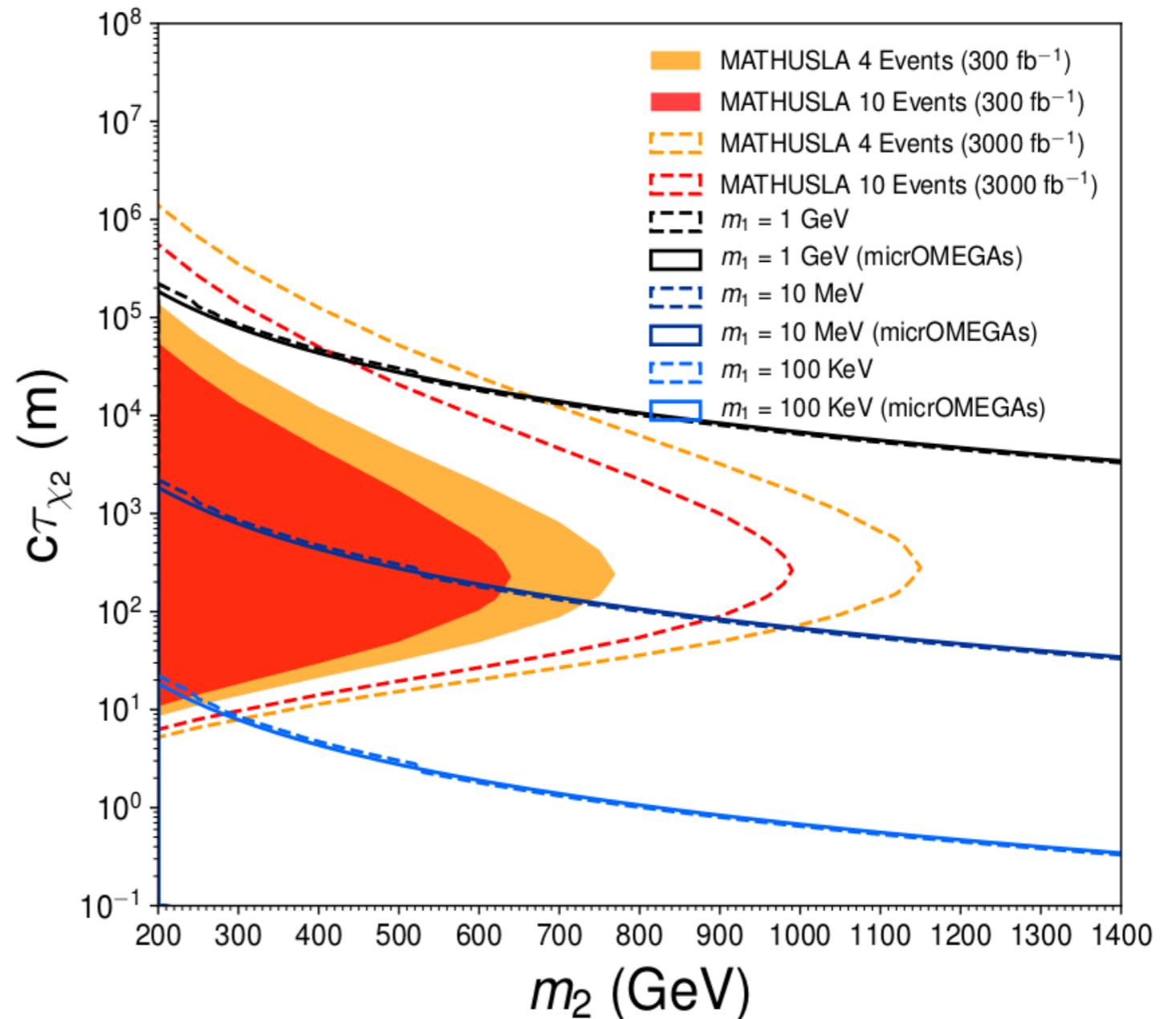
$$c\tau_2 \simeq 3500 \text{ m} \left(\frac{m_1}{100 \text{ MeV}} \right) \left(\frac{500 \text{ GeV}}{m_2} \right)^2$$

MATHUSLA solid angle
from interaction point

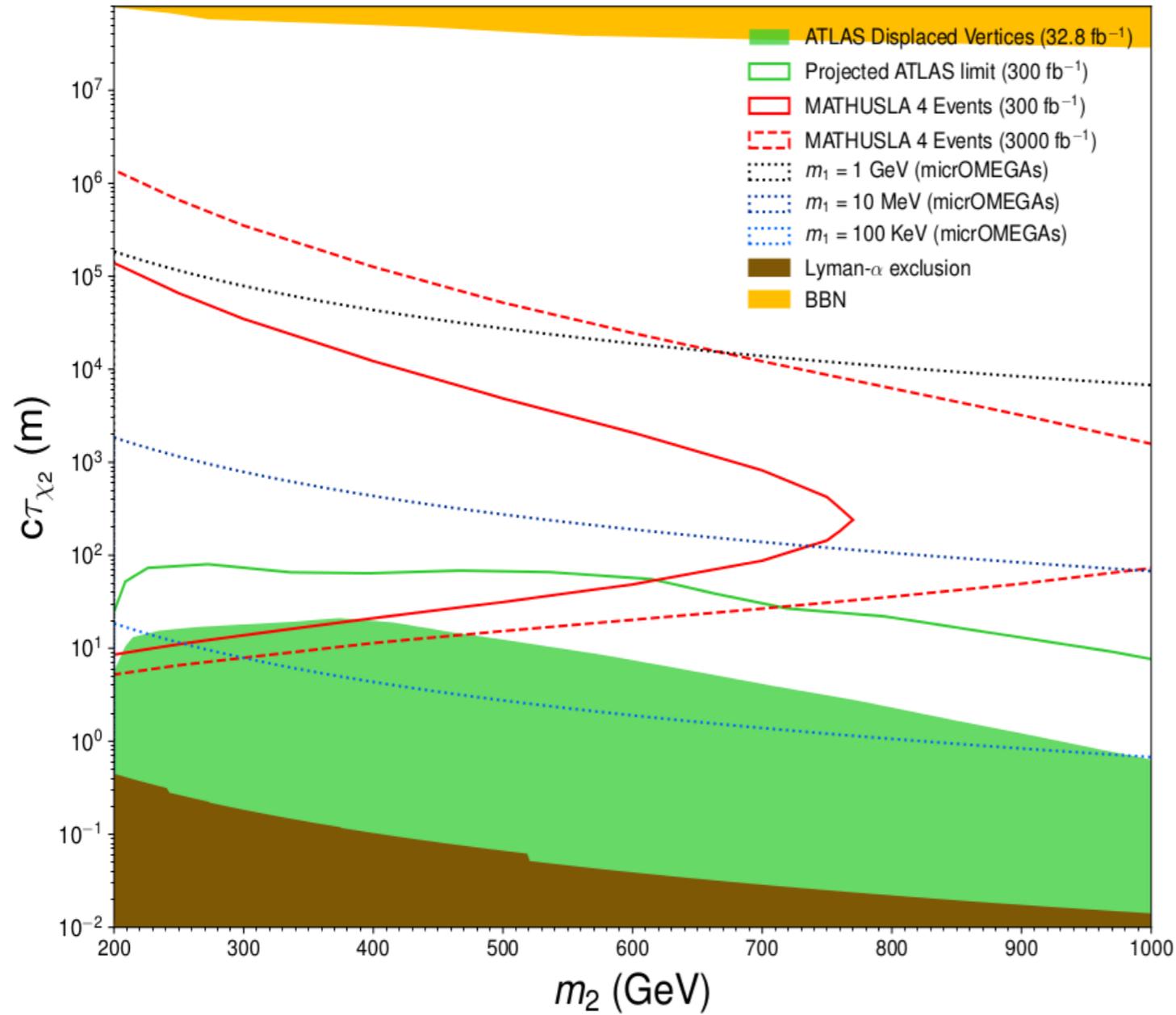
$$P_{\text{decay}}^{\text{MATH}} = \epsilon_{\text{geometric}} \times P_{\text{decay}}(\beta c\tau_2, L_a, L_b)$$

$$P_{\text{decay}}(\beta c\tau_2, L_a, L_b) = e^{-\frac{L_a}{\beta c\tau_2}} - e^{-\frac{L_b}{\beta c\tau_2}}$$

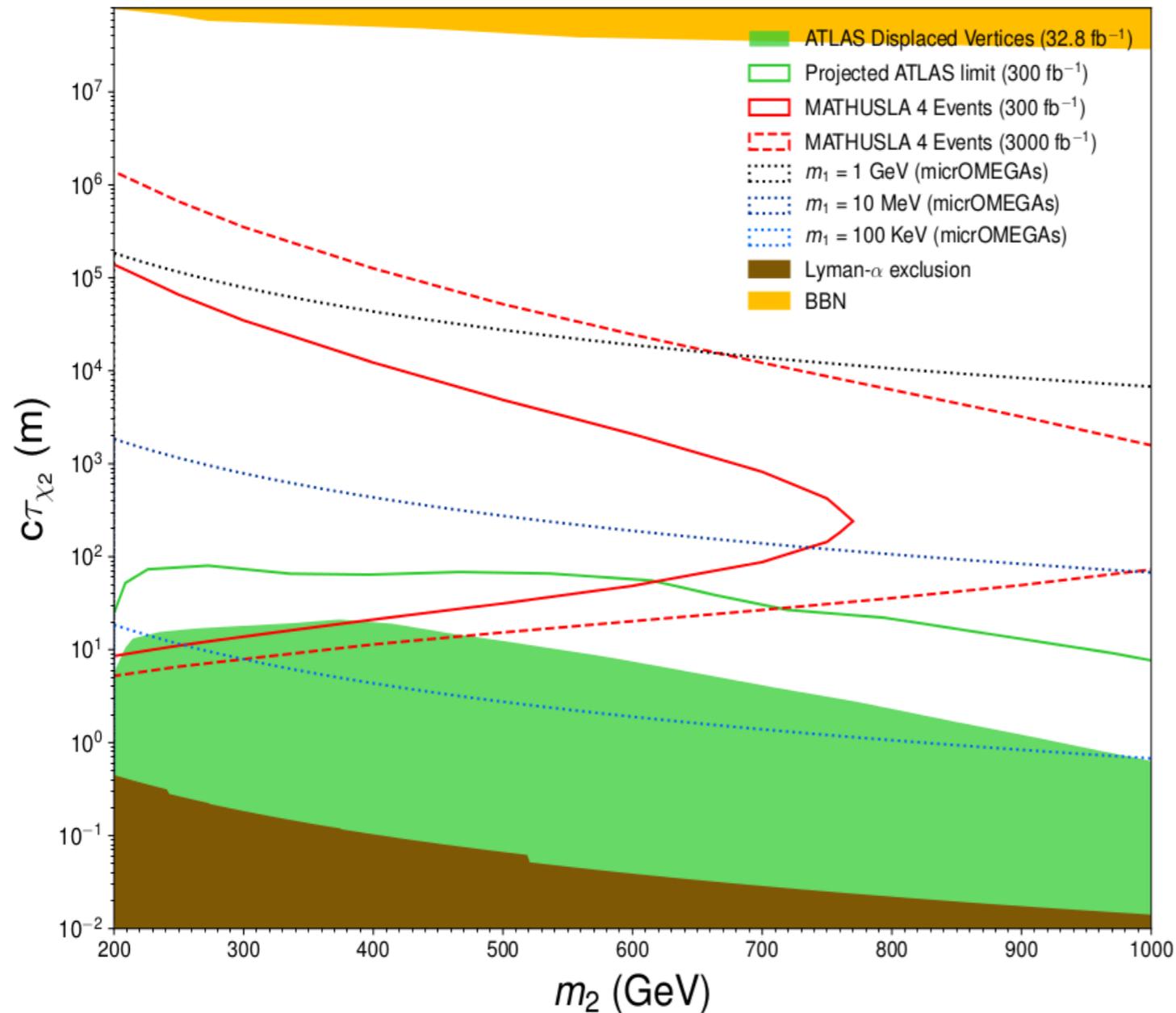
$$|\vec{p}_{\chi_2}|/m_2$$



Putting it all together & wrapping up...



Putting it all together & wrapping up...



➔ Freeze-in well-motivated alternative to WIMP

➔ Explore freeze-in DM probes @LHC with simple (minimal) models e.g. **Higgs “portal”**

➔ Possible to probe freeze-in DM with long-lived particle searches!

LHC
MATHUSLA (*advantage of dedicated LLP detector*)

Thank you!



KEEP
CALM
AND
BACKUP
YOUR
WORK

Constraints from Cosmology

→ Lyman- α forest observations

Constrain washout of small-scale structure by partially relativistic (“warm”) DM

$$m_{\text{DM}} \gtrsim 4.65 \text{ keV}$$

Thermal warm DM

Baur et al, 1706.03118

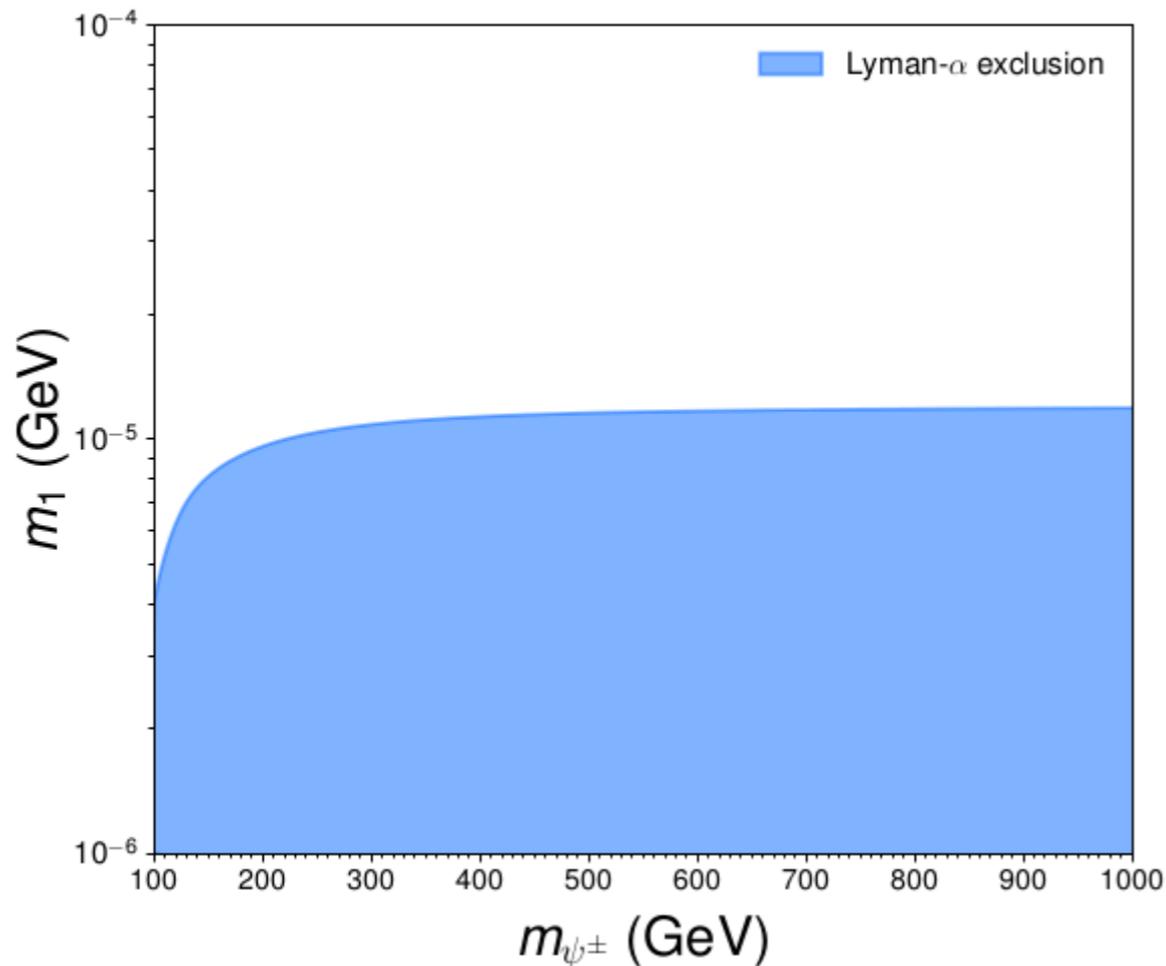
$$m_{\text{DM}} \gtrsim 12 \text{ keV} \left(\frac{\sum_{ij} g_i \Gamma_{ij} \Delta_{ij}^\eta}{\sum_{ij} g_i \Gamma_{ij}} \right)^{1/\eta}$$

Freeze-in DM from 2-body decay

Heeck et al, 1706.09909

$$\Gamma(A_i \rightarrow B_j \chi_1)$$

$$1 - \frac{m_{B_j}^2}{m_{A_i}^2}$$



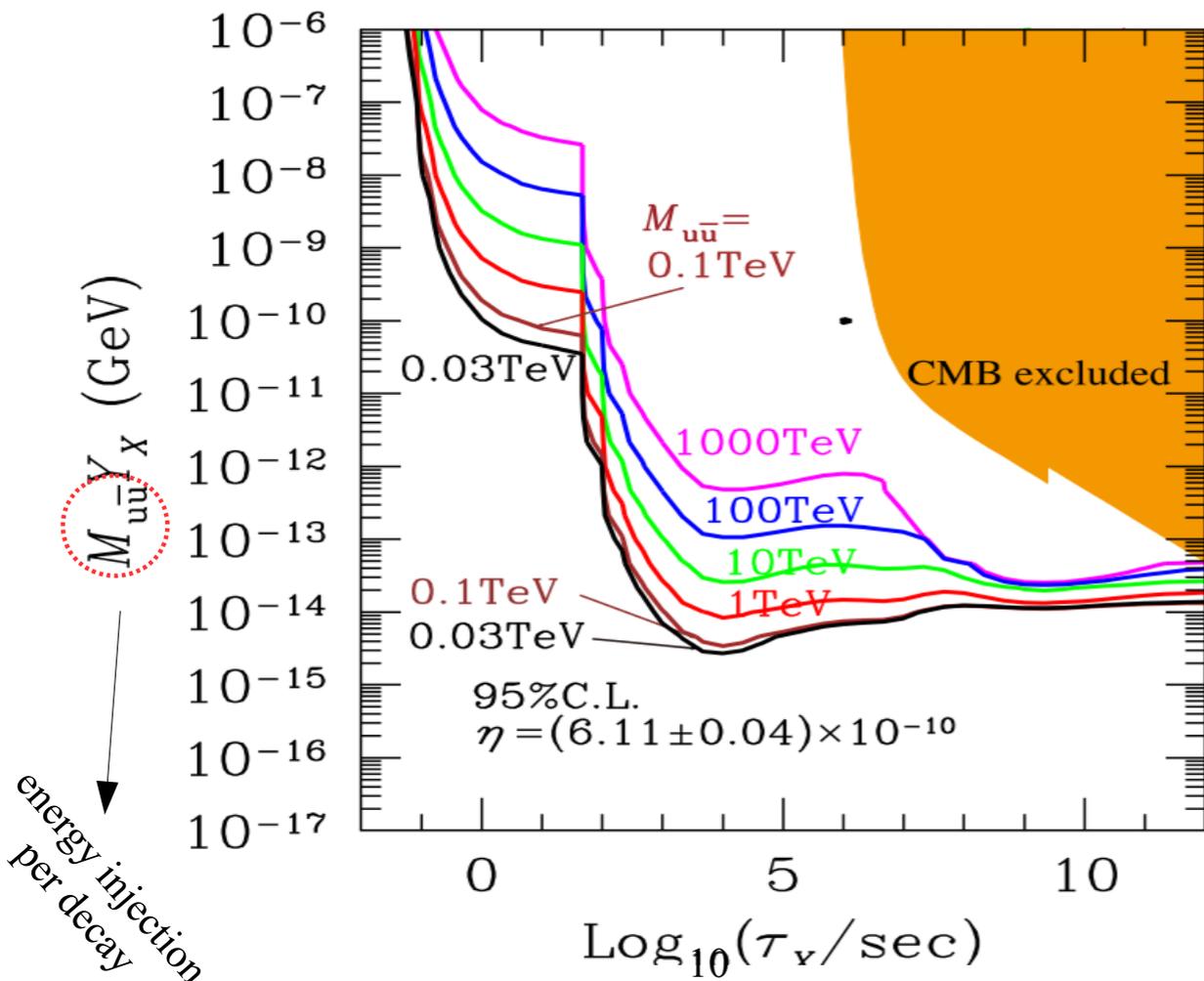
Assumes DM relic abundance is saturated!

Constraints from Cosmology

→ Big-bang nucleosynthesis

If long-lived, χ_2 visible decay products can affect BBN predictions

Constrain (hadronic) energy injection from the decay of χ_2



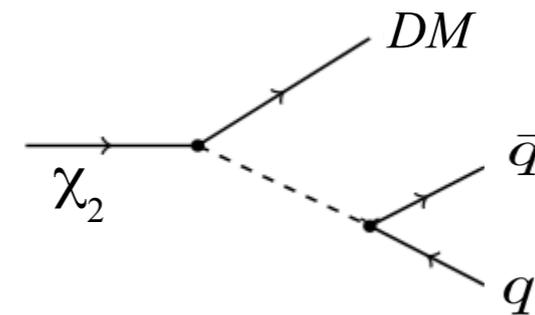
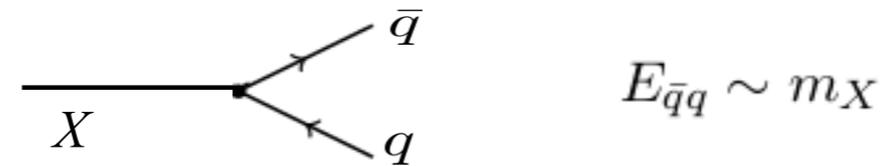
Kawasaki et al, 1709.01211

PHYSICAL REVIEW D **97**, 023502 (2018)

Editors' Suggestion

Revisiting big-bang nucleosynthesis constraints on long-lived decaying particles

Masahiro Kawasaki,^{1,2} Kazunori Kohri,^{3,4,5} Takeo Moroi,^{6,2} and Yoshitaro Takaesu^{6,7}



$$E_{\bar{q}q} \simeq \sqrt{m_Z^2 + \frac{m_2^2}{4} \left[(1 - \xi_z - \xi_1)^2 - 4\xi_z\xi_1 \right]} \quad (\chi_2 \rightarrow Z\chi_1)$$

$$\text{Constrain } \frac{\Omega_2 h^2}{m_2} \times \sum_{a=h,Z} E_{\bar{q}q}^a \times \text{BR}(\chi_2 \rightarrow a\chi_1) \text{BR}(a \rightarrow \text{hadrons})$$

Freeze-In DM Production

→ Production of DM in early Universe (assume standard thermal history)

Freeze-In:
$$\Omega_1 h^2 \simeq \frac{m_1}{\rho_c/s_0} \frac{135 M_{\text{Pl}} \Gamma_{\text{FI}}}{2 \times 1.66 \pi^3 m_2^2 [g_*(m_2/3)]^{3/2}}$$

Super-WIMP:

Feng et al, hep-ph/0302215

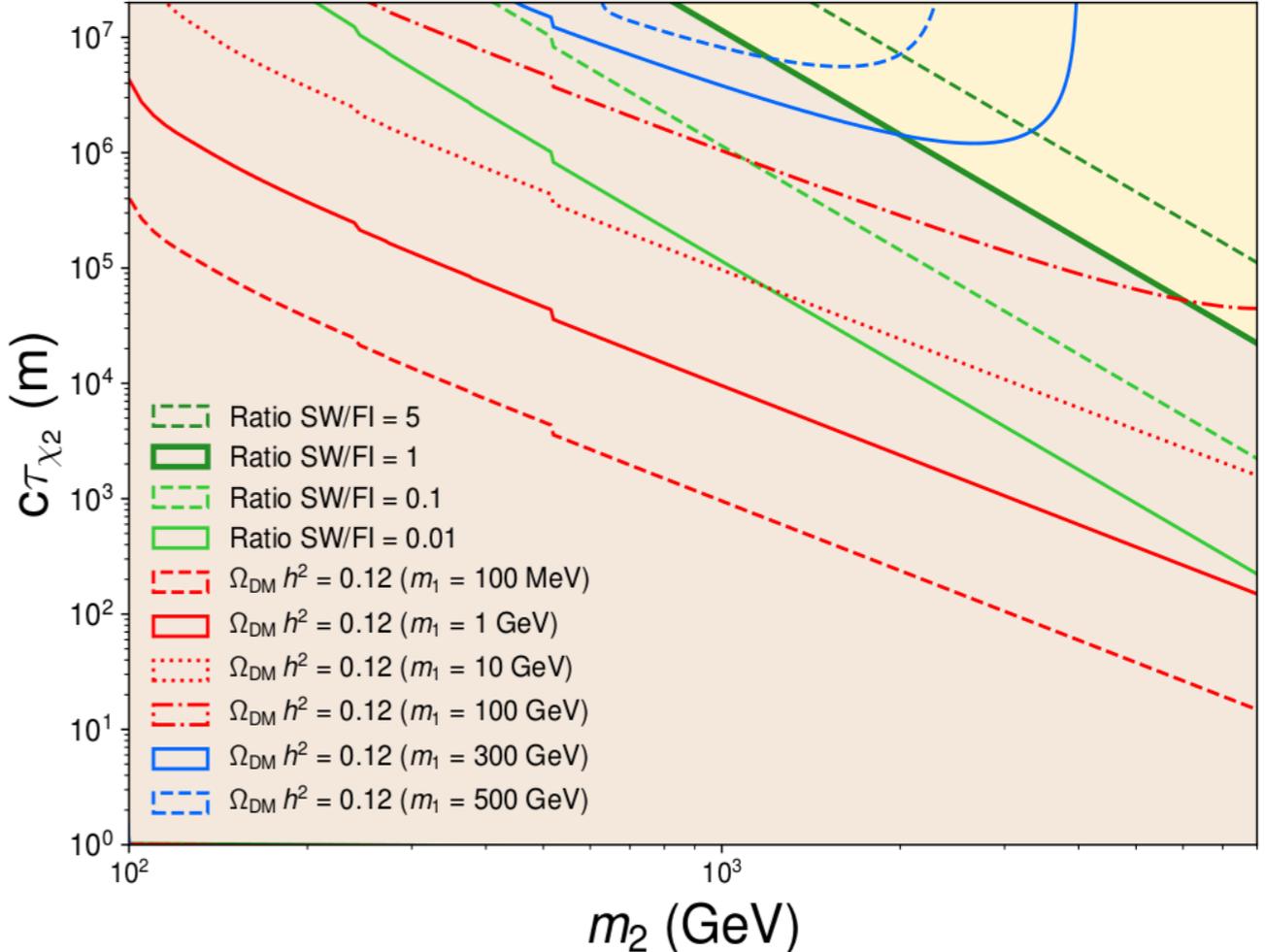
χ_2 freezes-out and then decays

$$\Omega_2 h^2 \simeq 0.1 \left(\frac{m_2}{\text{TeV}} \right)^2$$

$$\Omega_{\text{DM}} h^2 = \Omega_1 h^2 + \Omega_2 h^2 \times \frac{m_1}{m_2}$$

Note however that if Super-WIMP component dominant and mass ratio large, DM would be quite relativistic!

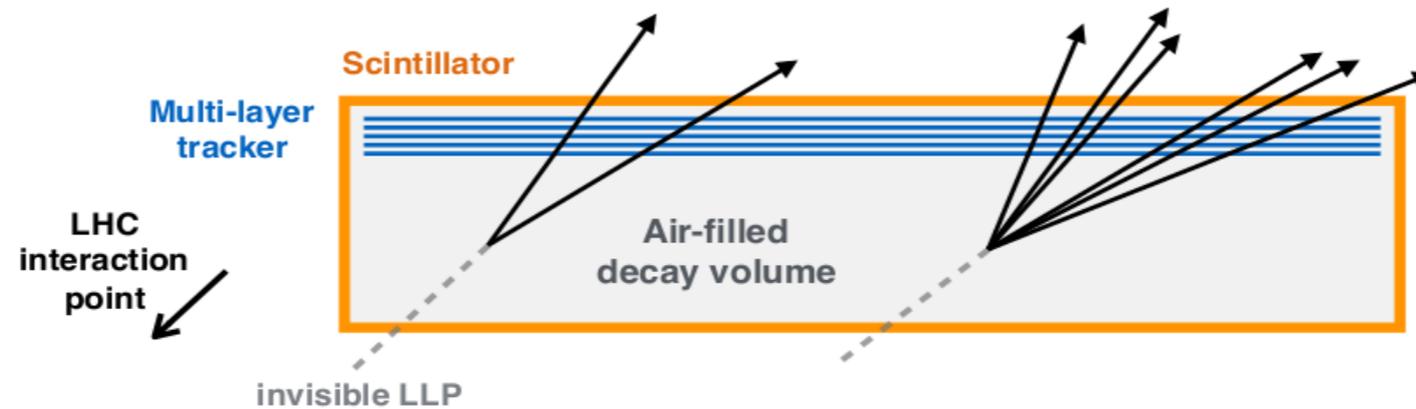
Garny et al, 1809.10135



Signal Reconstruction

Chou, DC, Lubatti
1606.06298

Courtesy of
David Curtin

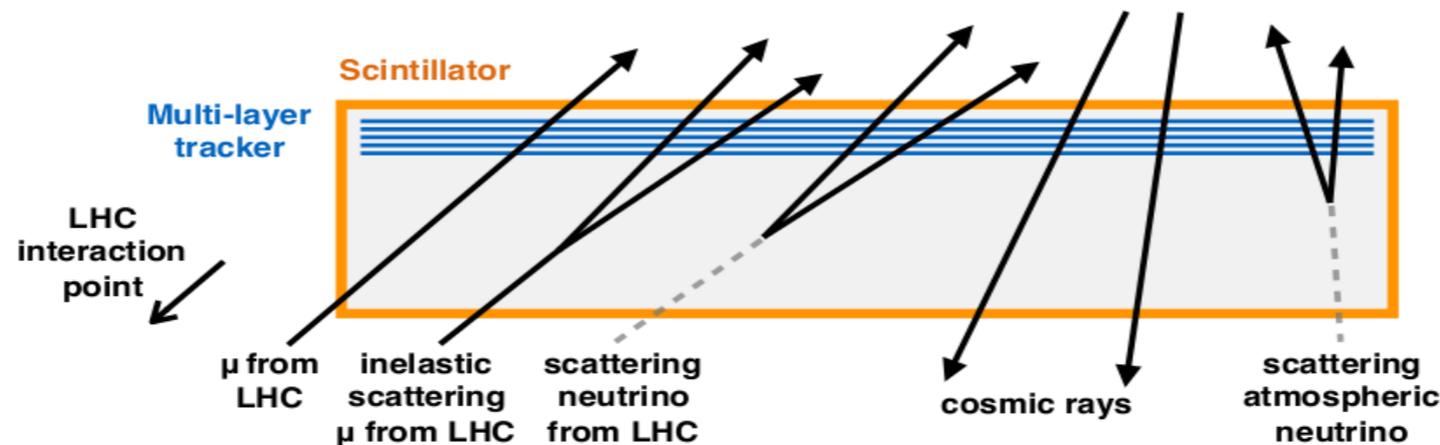


~5% geometric coverage, but much deeper than ATLAS/CMS
→ similar geometric acceptance for LLP decays (but no BG!)

Charged particle tracks are reconstructed with
~cm spatial resolution and ~ns timing resolution.

→ determine charged particle *speed* with ~0.05c precision.

LLP decays are reconstructed as Displaced Vertices (DV)
in both **space and time**, with strict **geometric requirements** and vetoes.



Reject using tight DV signal requirements, geometry & timing.

~Zero background regime can be reached!

*Cosmic backgrounds can be measured and studied during
beam down-time to verify rejection strategies.*