

Implications of LFU anomalies for high-energy physics

Gino Isidori

[*University of Zürich*]

- ▶ Introduction [*LFU and accidental symmetries*]
- ▶ A closer look to B-physics data
- ▶ EFT and simplified models
- ▶ Speculations on UV completions
- ▶ Conclusions



**University of
Zurich** ^{UZH}



European Research Council
Established by the European Commission

► Introduction

What is **L**epton **F**lavor **U**niversality and why we care about it?

LFU is an accidental approximate symmetry of the SM Lagrangian.
(*exact SM global symmetry in the limit where we neglect lepton Yukawa's*):

$$\mathcal{L}_{\text{SM}} = \mathcal{L}_{\text{gauge}} + \mathcal{L}_{\text{Higgs}} \Big|_{y_l = 0} \longrightarrow 3 \times 3 \text{ (unitary) transformations of lepton fields in flavor space}$$

LFU is badly broken in the Yukawa sector: $y_e \sim 3 \times 10^{-6}$, $y_\mu \sim 3 \times 10^{-4}$, $y_\tau \sim 10^{-2}$

But all the lepton Yukawa couplings are small compared to SM gauge couplings, giving rise to the (*approximate*) universality of decay amplitudes which differ only by the different lepton species involved

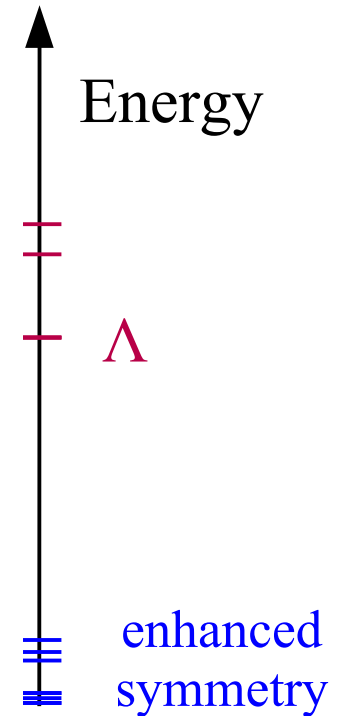
► Introduction

$$\mathcal{L}_{\text{SM-EFT}} = \mathcal{L}_{\text{gauge}} + \mathcal{L}_{\text{Higgs}} + \underbrace{\sum_{d,i} \frac{c_i^{[d]}}{\Lambda^{d-4}} \mathcal{O}_i^{d \geq 5}}_{\text{Violations of accidental symmetries}}$$

Accidental symmetries are excellent probes of high-energy dynamics,

as we know from well-known examples in the past:

Eg: *Low-energy theory:* QED + QCD
Accidental symm.: Flavor [U(1)_{n_f}]
Violated by: Weak interactions → G_F ~ (250 GeV)⁻²



► Introduction

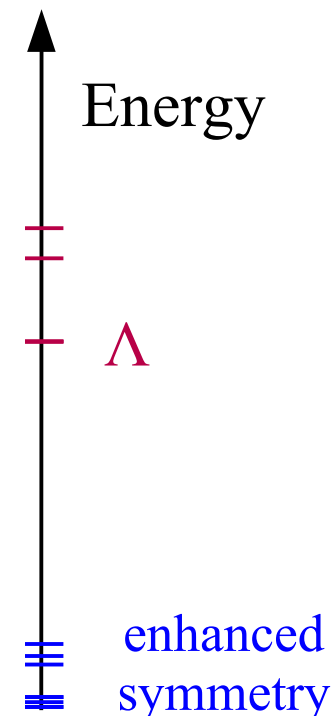
$$\mathcal{L}_{\text{SM-EFT}} = \mathcal{L}_{\text{gauge}} + \mathcal{L}_{\text{Higgs}} + \underbrace{\sum_{d,i} \frac{c_i^{[d]}}{\Lambda^{d-4}} \mathcal{O}_i^{d \geq 5}}_{\text{Violations of accidental symmetries}}$$

Accidental symmetries are excellent probes of high-energy dynamics,

as we know from well-known examples in the past:

Eg: *Low-energy theory:* QED + QCD
Accidental symm.: Flavor [U(1)_{n_f}]
Violated by: Weak interactions → G_F ~ (250 GeV)⁻²

Eg: *Low-energy theory:* SM, 2 generations
Accidental symm.: CP
Violated by: “Super-weak” interactions



(single evidence of the violation in K- \bar{K} mixing for long time)

► Introduction

$$\mathcal{L}_{\text{SM-EFT}} = \mathcal{L}_{\text{gauge}} + \mathcal{L}_{\text{Higgs}} + \underbrace{\sum_{d,i} \frac{c_i^{[d]}}{\Lambda^{d-4}} \mathcal{O}_i^{d \geq 5}}_{\text{Violations of accidental symmetries}}$$

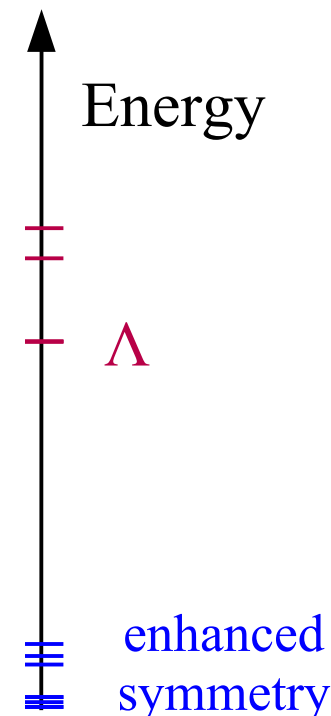
Accidental symmetries are excellent probes of high-energy dynamics,

as we know from well-known examples in the past:

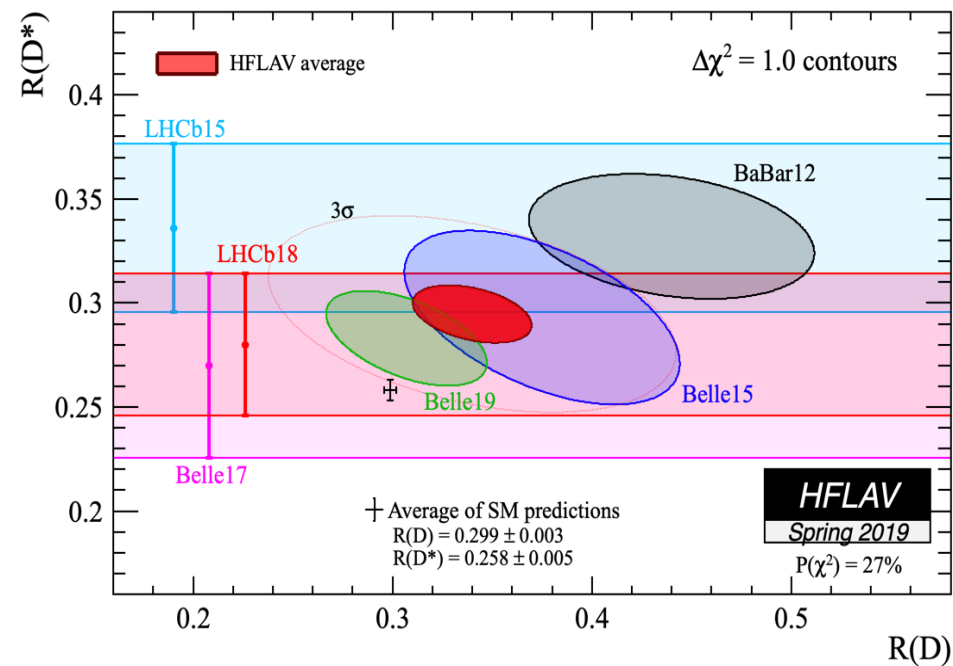
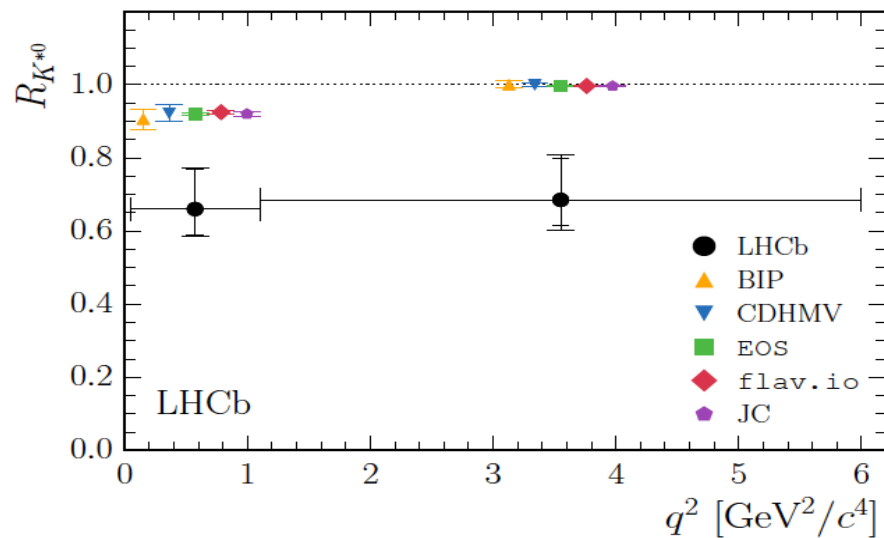
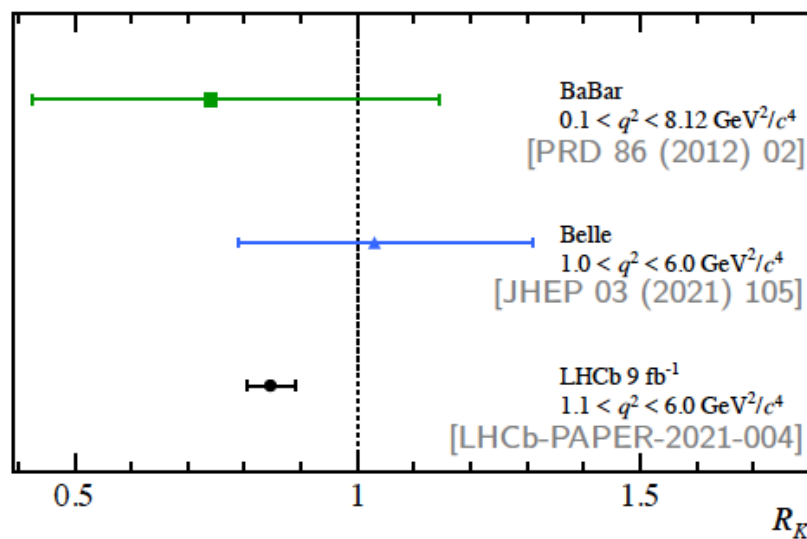
Eg: *Low-energy theory:* QED + QCD
Accidental symm.: Flavor [U(1)_{n_f}]
Violated by: Weak interactions → G_F ~ (250 GeV)⁻²

Eg: *Low-energy theory:* SM, 2 generations
Accidental symm.: CP
Violated by: “Super-weak” interactions → $\frac{(G_F m_t V_{ts} V_{td})^2}{4\pi^2} \sim (10^4 \text{ TeV})^{-2}$

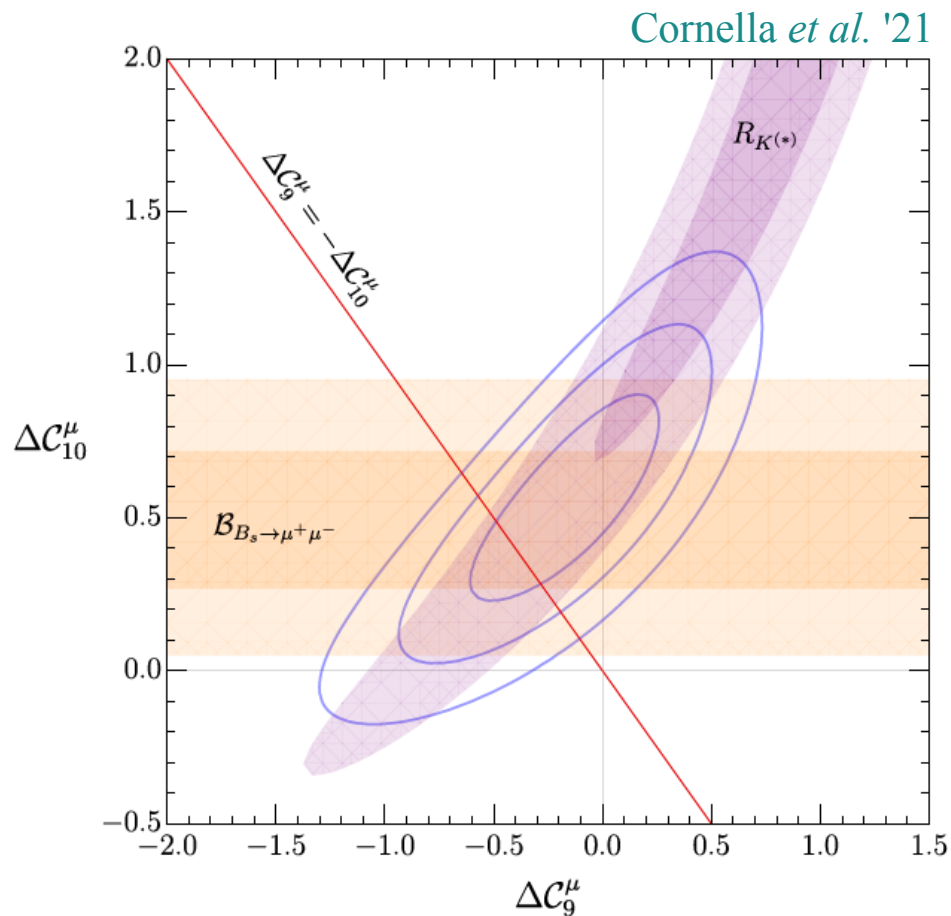
(single evidence of the violation in K- \bar{K} mixing for long time)



A closer look to B-physics data



► A closer look to B-physics data



• $b \rightarrow s l^+ l^-$ (neutral currents)

FCNC operators:

$$\mathcal{O}_{10}^{\ell} = (\bar{s}_L \gamma_{\mu} b_L)(\bar{\ell} \gamma^{\mu} \gamma_5 \ell)$$

$$\mathcal{O}_9^{\ell} = (\bar{s}_L \gamma_{\mu} b_L)(\bar{\ell} \gamma^{\mu} \ell)$$

$$\Delta C_i^{\mu} = C_i^{\mu} - C_i^e$$

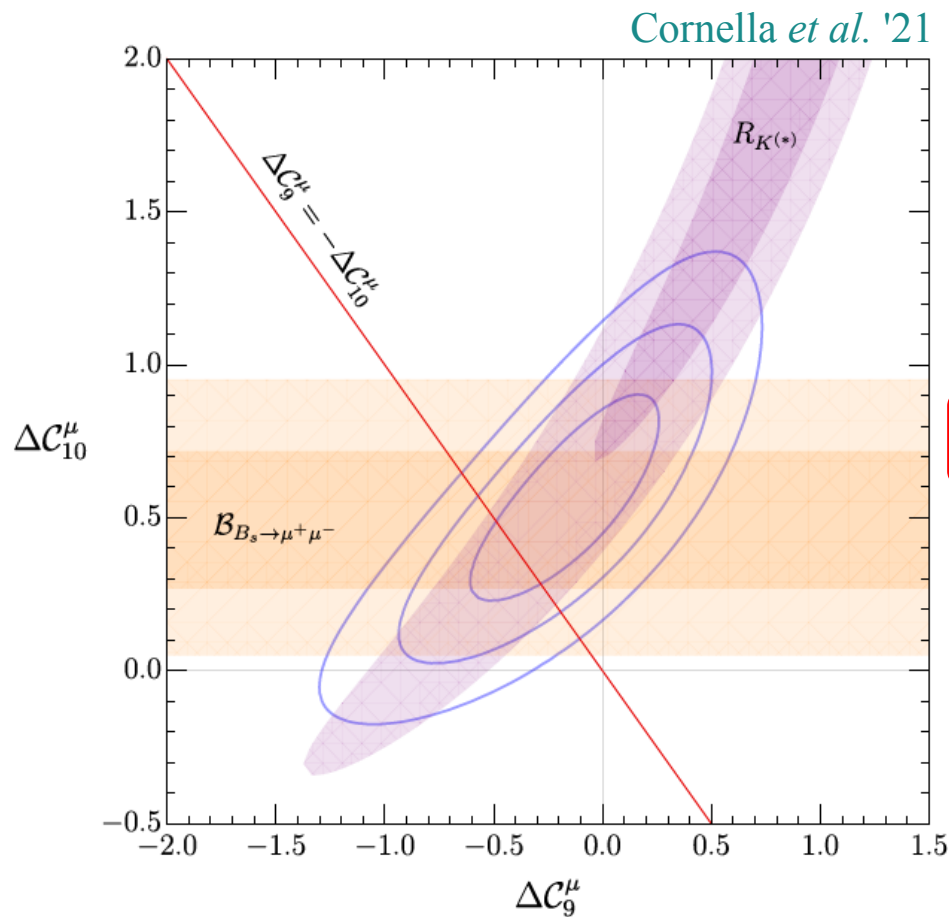
Conservative fit using “clean observables” only” [= R_K , R_{K^*} , $B_s \rightarrow \mu\mu$]

4.6 σ

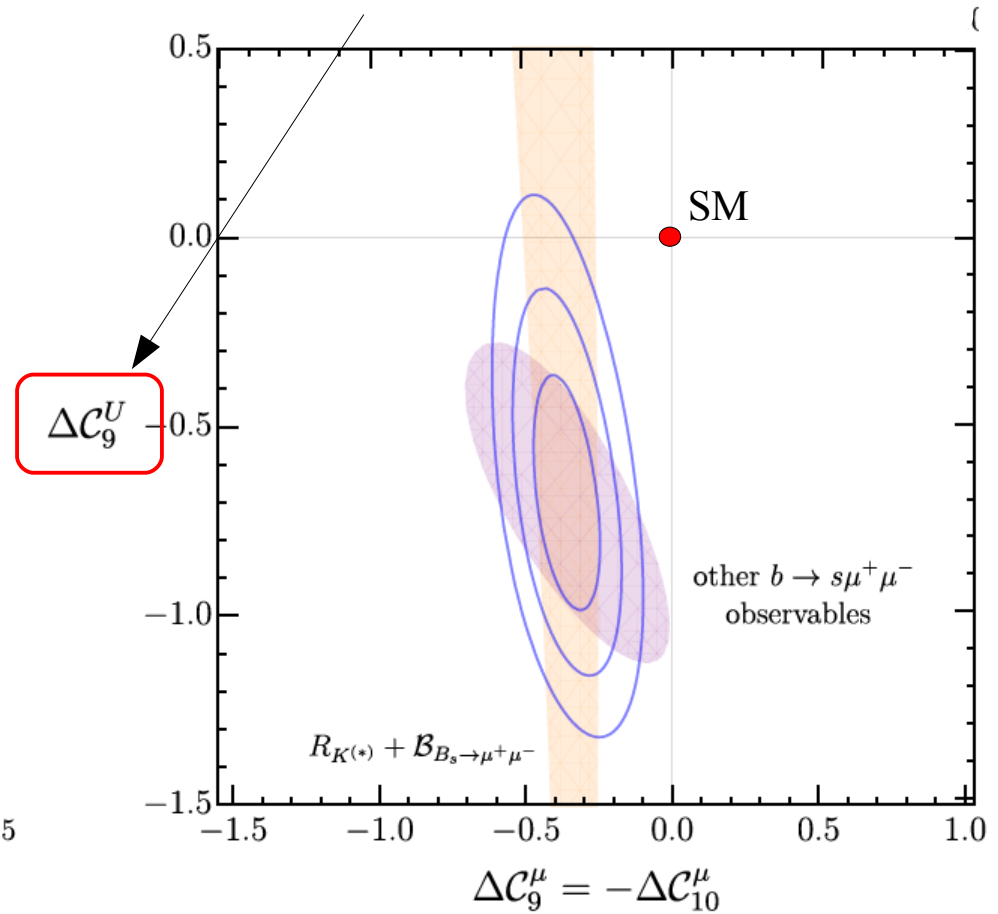
significance of NP hypothesis

$$\Delta C_9^{\mu} = -\Delta C_{10}^{\mu} \text{ vs. SM}$$

► A closer look to B-physics data



Lepton-universal shift to C_9
(sensitive to charm re-scattering)



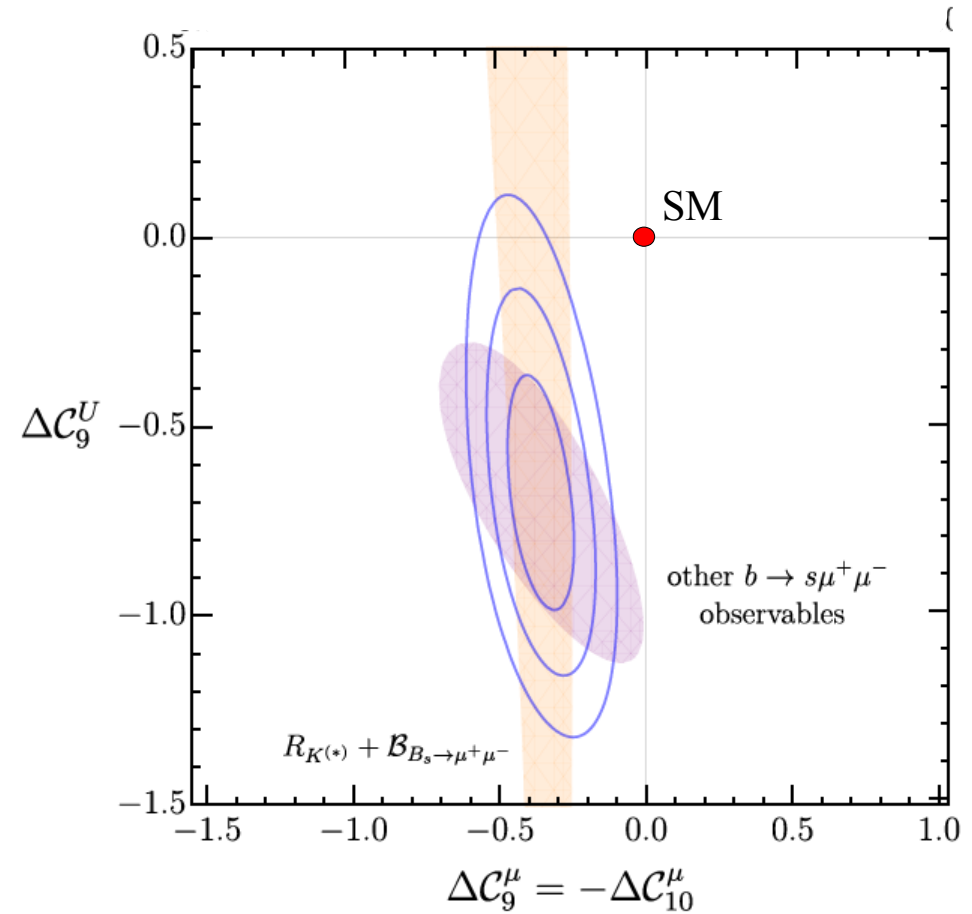
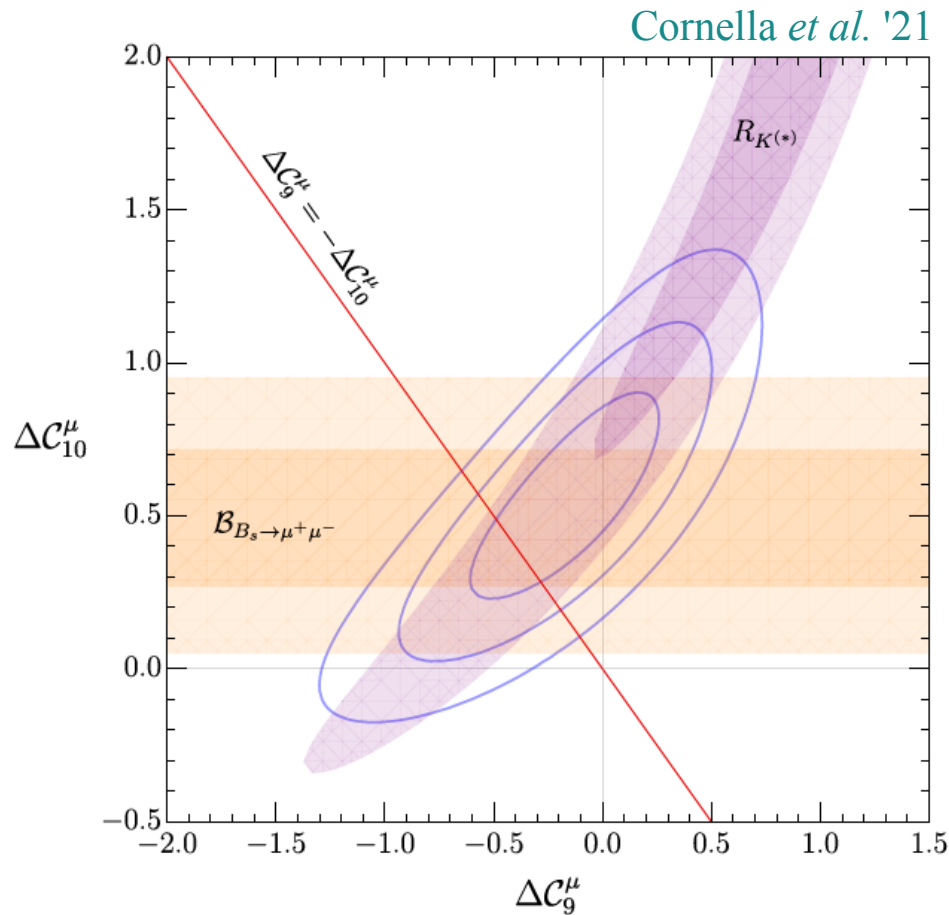
Conservative fit using “clean obs.”
only [$\Delta C_i^\mu = C_i^\mu - C_i^e$]:

4.6 σ significance of NP hypothesis
 $\Delta C_9^\mu = -\Delta C_{10}^\mu$ vs. SM

>> 5 σ with current best estimate of charm contributions

Alguero et al. '19
Ciuchini et al. '20
Li-Sheng Geng et al. '21
Altmanshofer & Stangl '21

► A closer look to B-physics data



Conservative fit using “clean obs.”
only [$\Delta C_i^\mu = C_i^\mu - C_i^e$]:

4.6 σ significance of NP hypothesis
 $\Delta C_9^\mu = -\Delta C_{10}^\mu$ vs. SM

>> 5 σ with current best estimate of charm contributions

3.9 σ global significance of NP
(most conservative estimate)

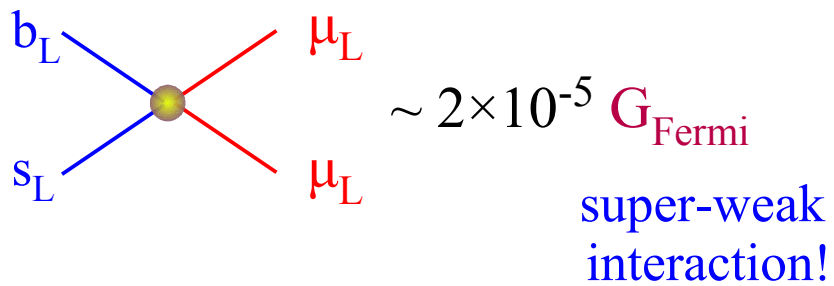
Alguero et al. '19
Ciuchini et al. '20
Li-Sheng Geng et al. '21
Altmanshofer & Stangl '21

Lancierini, GI,
Owen, Serra, '21

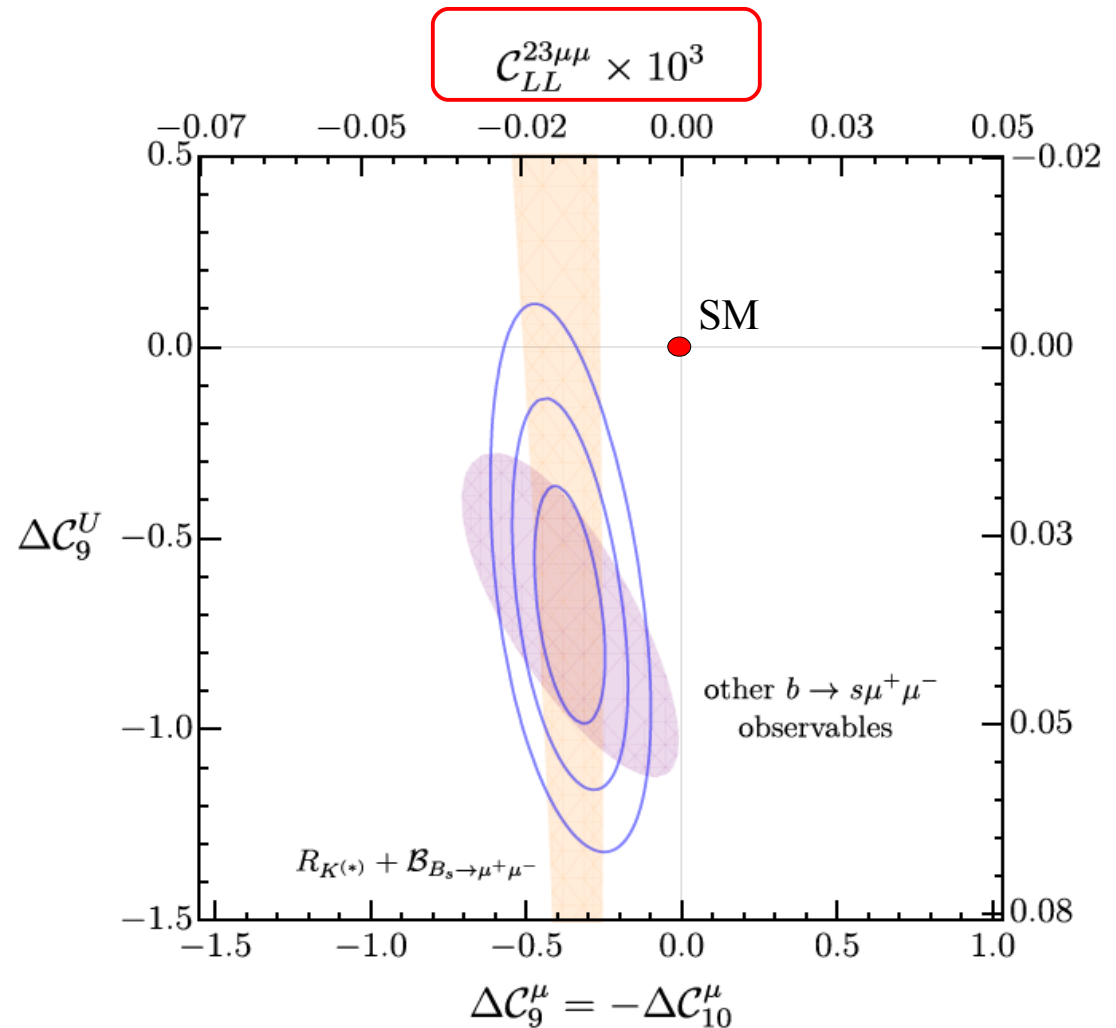
► A closer look to B-physics data

Data point to (short-distance) NP effects in operators of the type

$$\mathcal{O}_{LL}^{ij\alpha\beta} = (\bar{q}_L^i \gamma_\mu \ell_L^\alpha) (\bar{\ell}_L^\beta \gamma_\mu q_L^j)$$



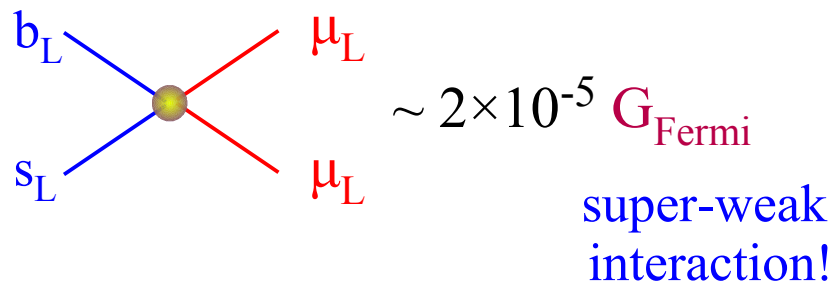
$$C_{LL}^{23\mu\mu} \rightarrow \Delta C_9^\mu = -\Delta C_{10}^\mu$$



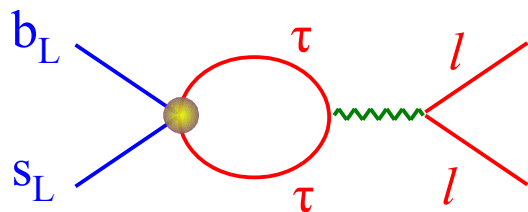
► A closer look to B-physics data

Data point to (short-distance) NP effects in operators of the type

$$\mathcal{O}_{LL}^{ij\alpha\beta} = (\bar{q}_L^i \gamma_\mu \ell_L^\alpha) (\bar{\ell}_L^\beta \gamma_\mu q_L^j)$$



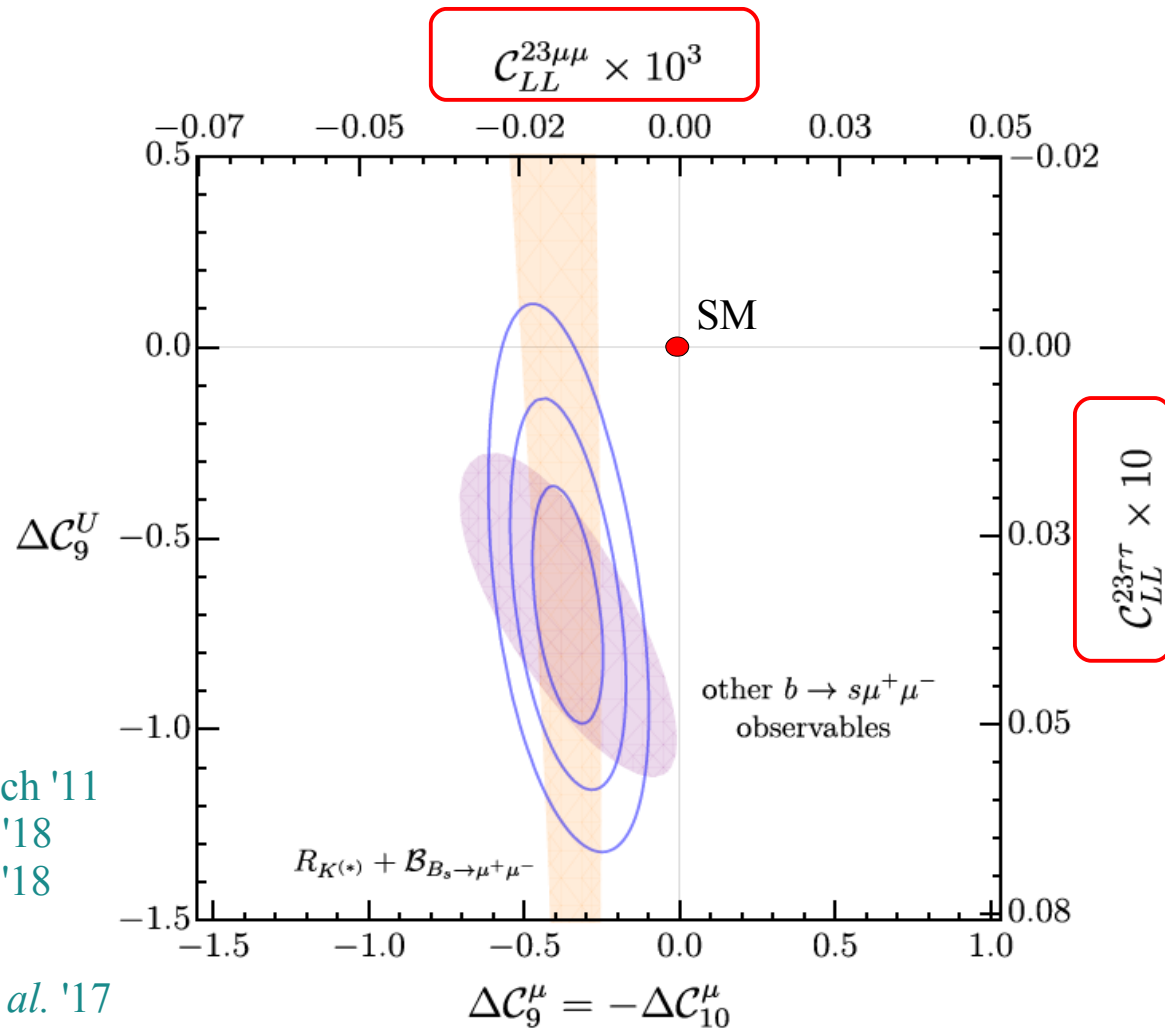
$$C_{LL}^{23\mu\mu} \rightarrow \Delta C_9^\mu = -\Delta C_{10}^\mu$$



$$C_{LL}^{23\tau\tau} \rightarrow \Delta C_9^{\text{Univ}}$$

Bobeth & Haisch '11
Crivellin *et al.* '18
Alguero *et al.* '18

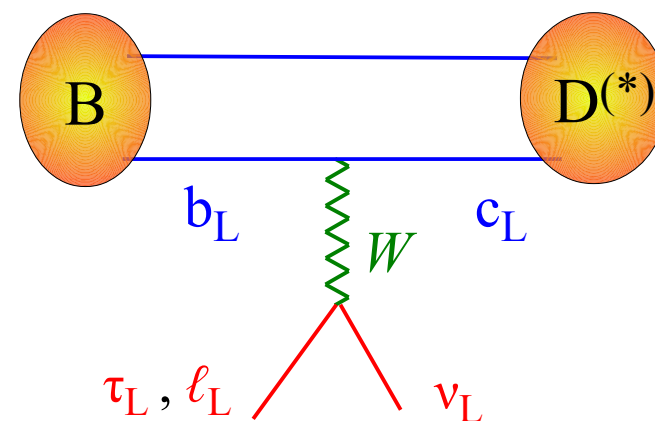
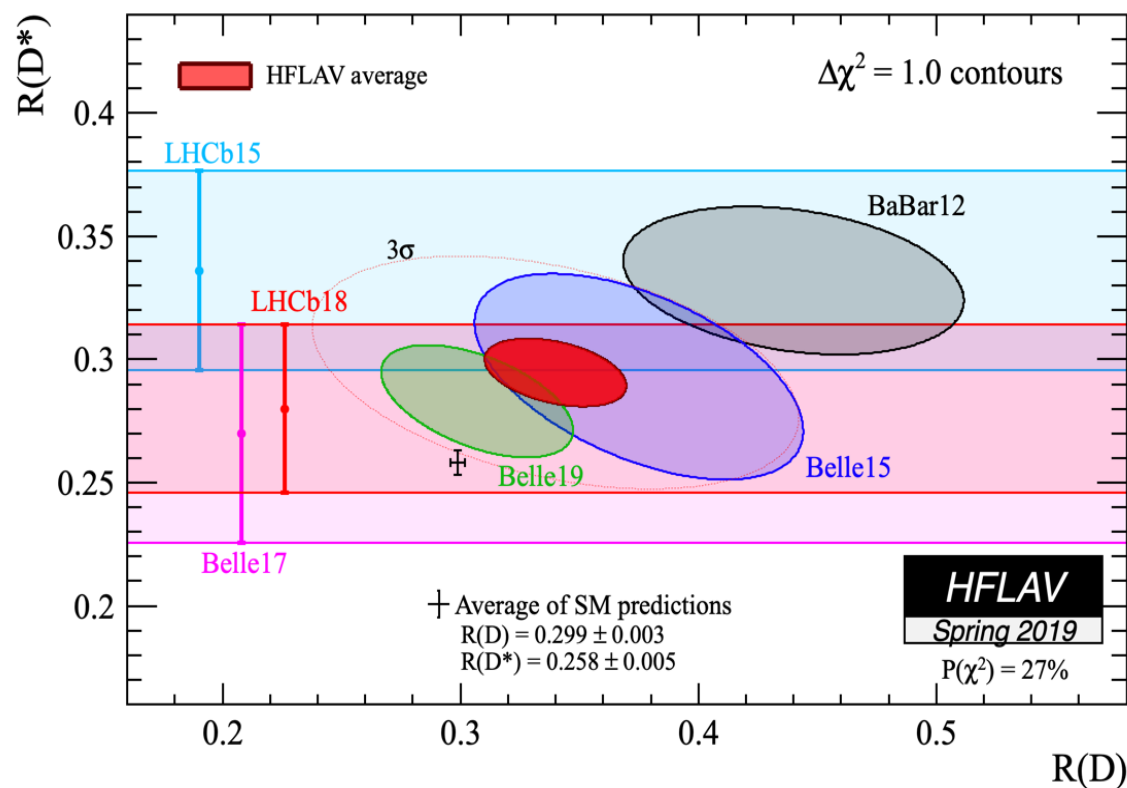
Link to CC anomaly Greljo *et al.* '17



► A closer look to B-physics data

- $b \rightarrow c \ell \bar{\nu}$ (charged currents): τ vs. light leptons (μ, e)

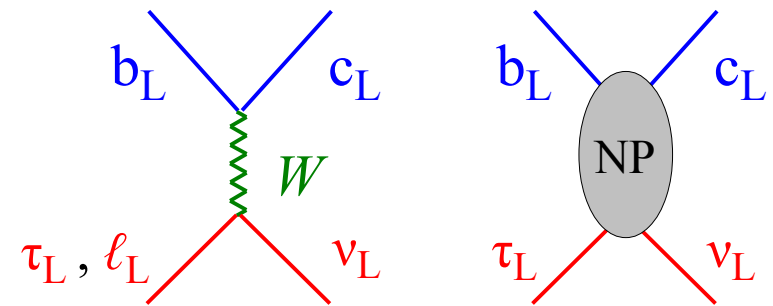
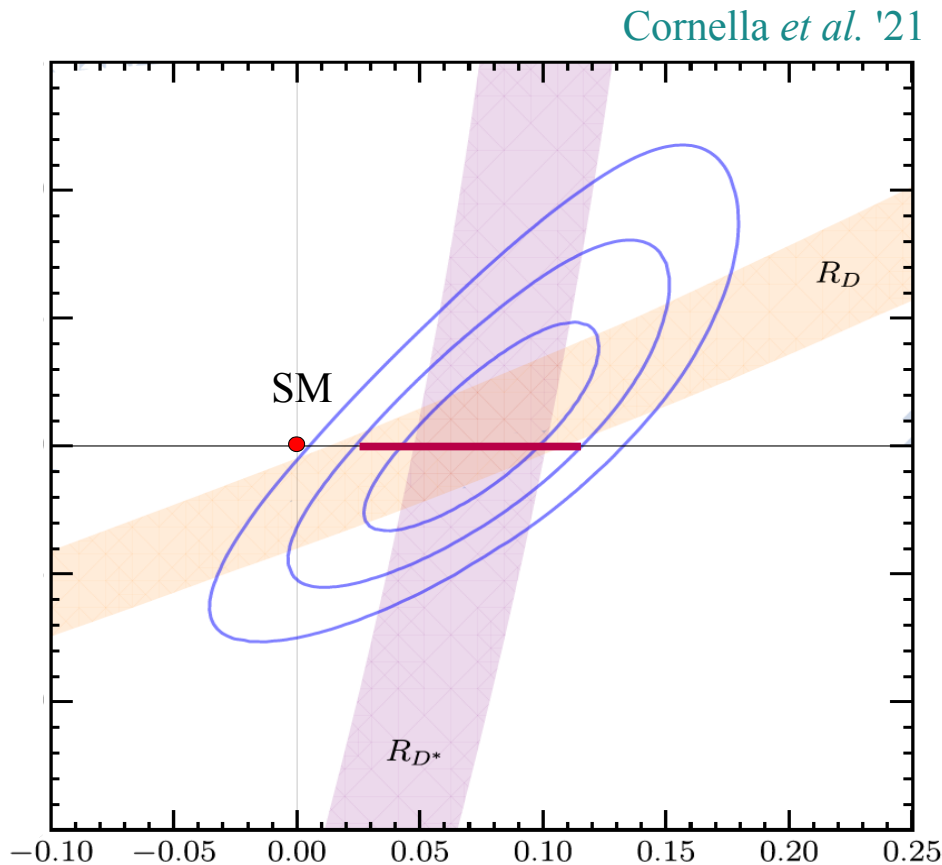
$$R(X) = \frac{\Gamma(B \rightarrow X \tau \bar{\nu})}{\Gamma(B \rightarrow X \ell \bar{\nu})} \quad X = D \text{ or } D^*$$



- Consistent results by three different expts. $\sim 3.1\sigma$ excess over SM (D and D^* combined)
- SM predictions quite “clean”: hadronic uncertainties cancel (to large extent) in the ratios

► A closer look to B-physics data

- $b \rightarrow c l \nu$ (charged currents): τ vs. light leptons (μ, e)



Data consistent with a universal enhancement (10-20%) of τ modes

C_{LL}^c

$$\frac{V_{cb} \mathcal{C}_{LL}^{33\tau\tau} + V_{cs} \mathcal{C}_{LL}^{23\tau\tau}}{V_{cb}}$$

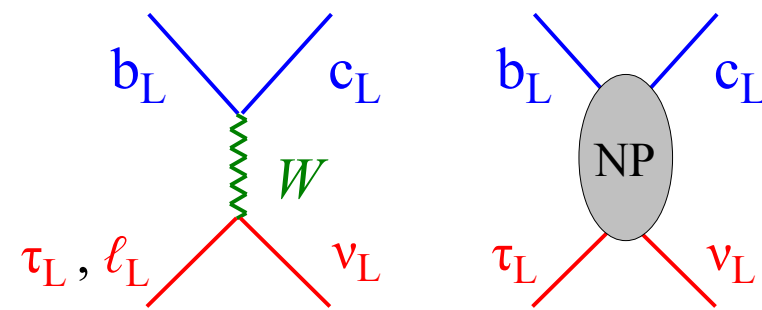
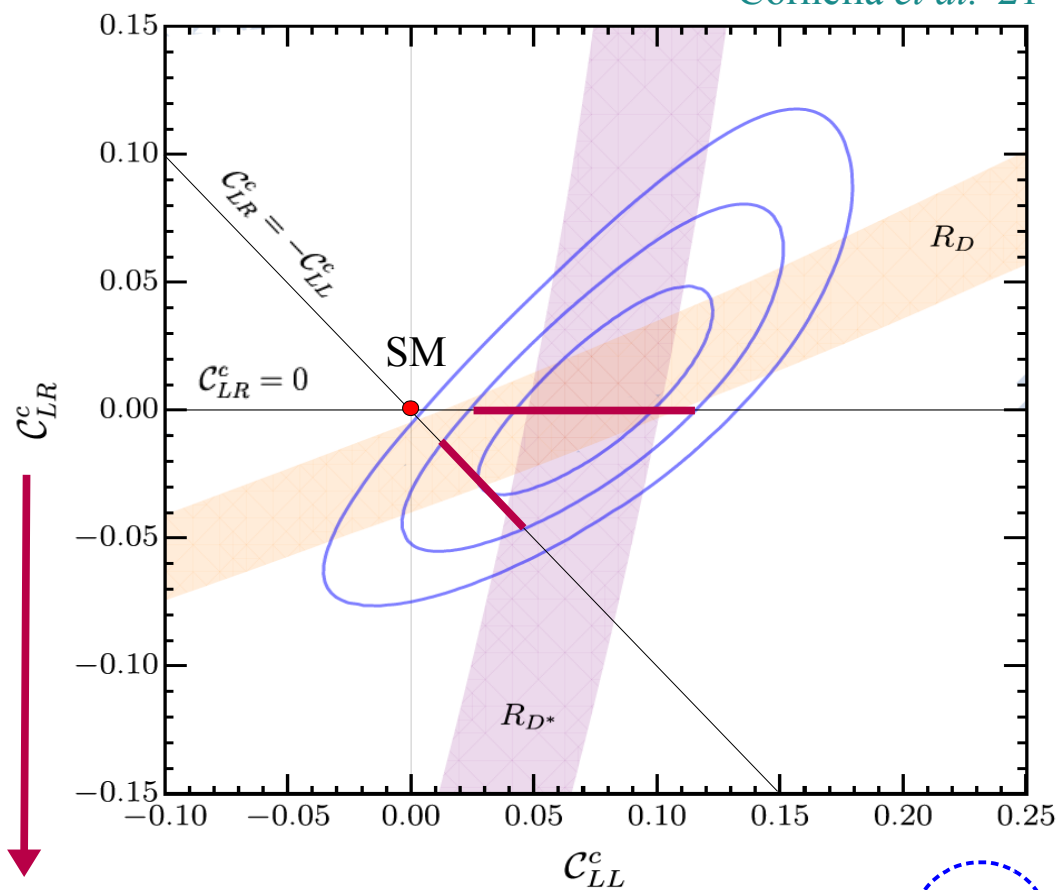
Same operator contributing to $b \rightarrow s ll$

all 3rd gen. (contribute via CKM rotation)

► A closer look to B-physics data

• $b \rightarrow c \ell \nu$ (charged currents): τ vs. light leptons (μ, e)

Cornella et al. '21



Data consistent with a universal enhancement (10-20%) of τ modes
 But other options (*RH currents*) possible

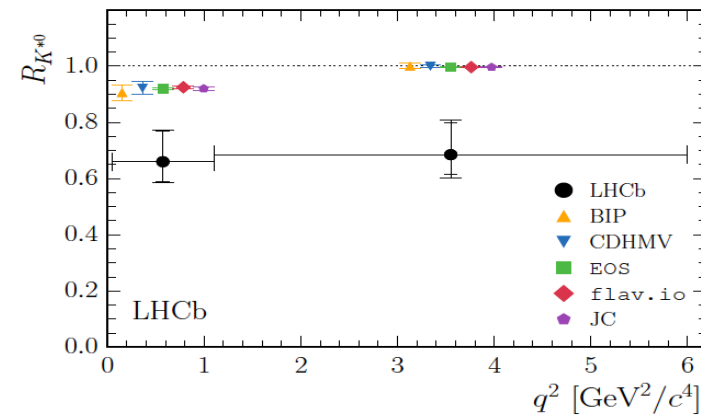
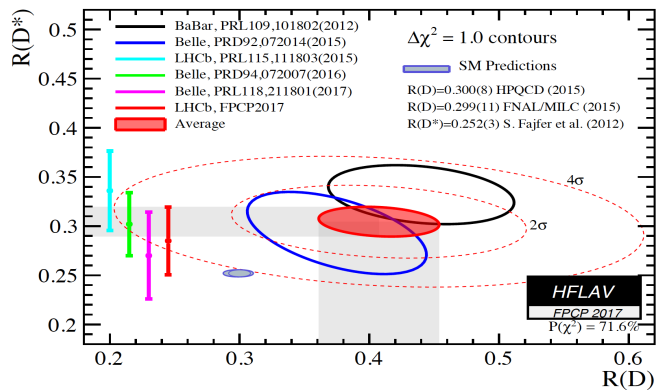
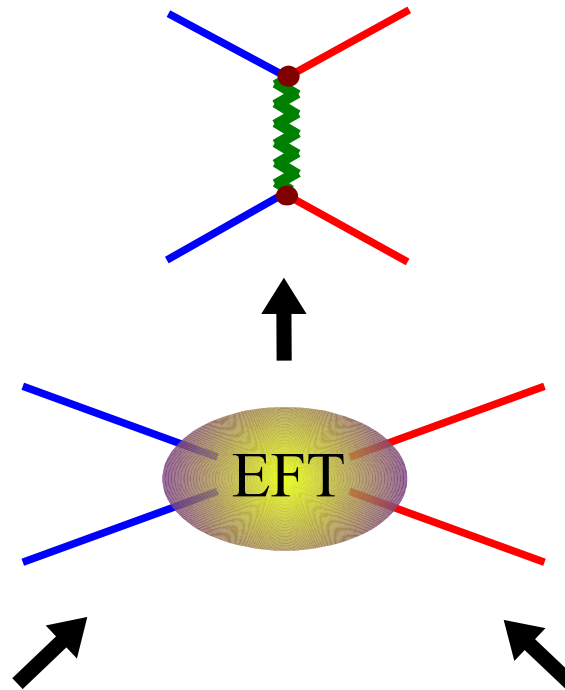
$b_R \rightarrow c_L \tau_R \nu_L$

$$\frac{V_{cb} \mathcal{C}_{LL}^{33\tau\tau} + V_{cs} \mathcal{C}_{LL}^{23\tau\tau}}{V_{cb}}$$

Same operator contributing to $b \rightarrow s \ell \ell$

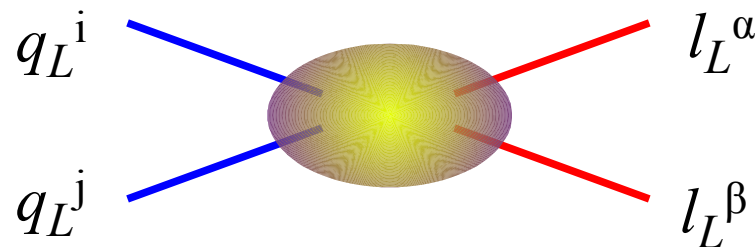
all 3rd gen. (contribute via CKM rotation)

From EFT to simplified models



► EFT considerations

- Anomalies are seen only in semi-leptonic (**quark**×**lepton**) operators
- We definitely need non-vanishing **left-handed** current-current operators although other contributions are also possible



Bhattacharya *et al.* '14
 Alonso, Grinstein, Camalich '15
 Greljo, GI, Marzocca '15
 (+many others...)

- Large coupling [*competing with SM tree-level*] in **bc** → $l_3 \nu_3$ [$\mathbf{R}_D, \mathbf{R}_{D^*}$]
- Small coupling [*competing with SM loop-level*] in **bs** → $l_2 l_2$ [$\mathbf{R}_K, \mathbf{R}_{K^*}, \dots$]

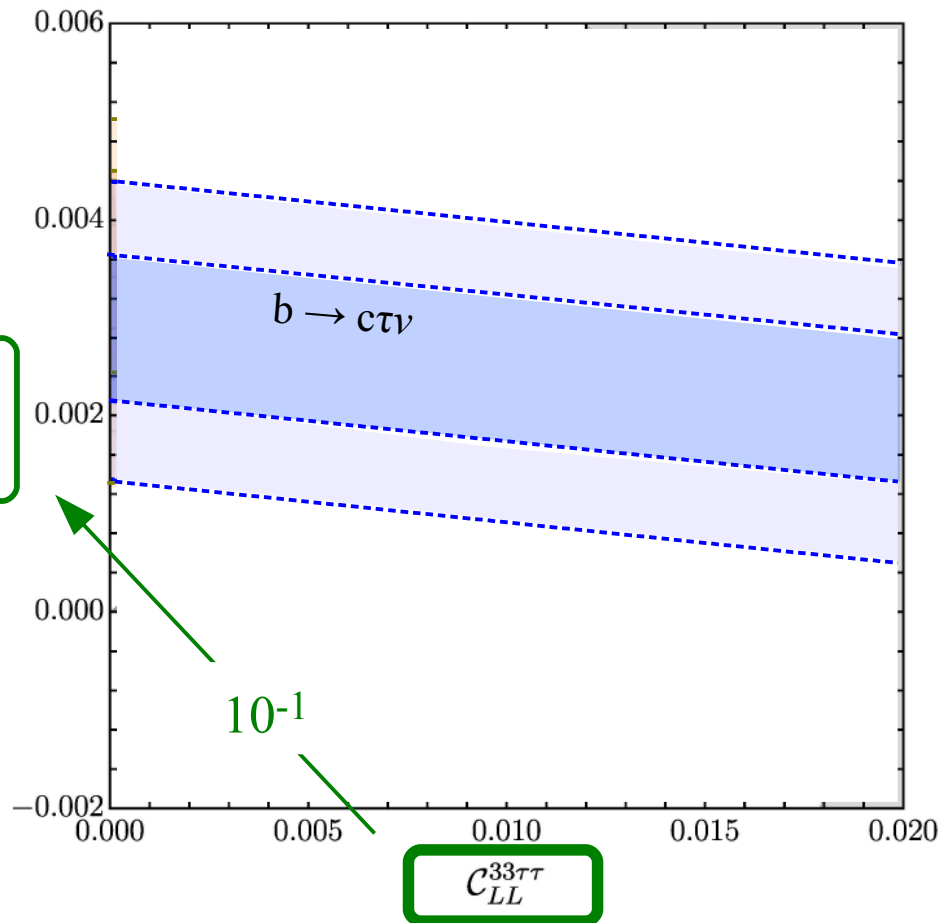


$$C_{ij\alpha\beta} = \begin{array}{l} \text{large for} \\ 3^{\text{rd}} \text{ generation} \\ \text{fields} \end{array} + \begin{array}{l} \text{small terms} \\ \text{for } 2^{\text{nd}} \text{ (& } 1^{\text{st}}) \\ \text{generations} \end{array}$$



*Link to pattern
of the Yukawa
couplings !*

► EFT considerations



charged-currents only:

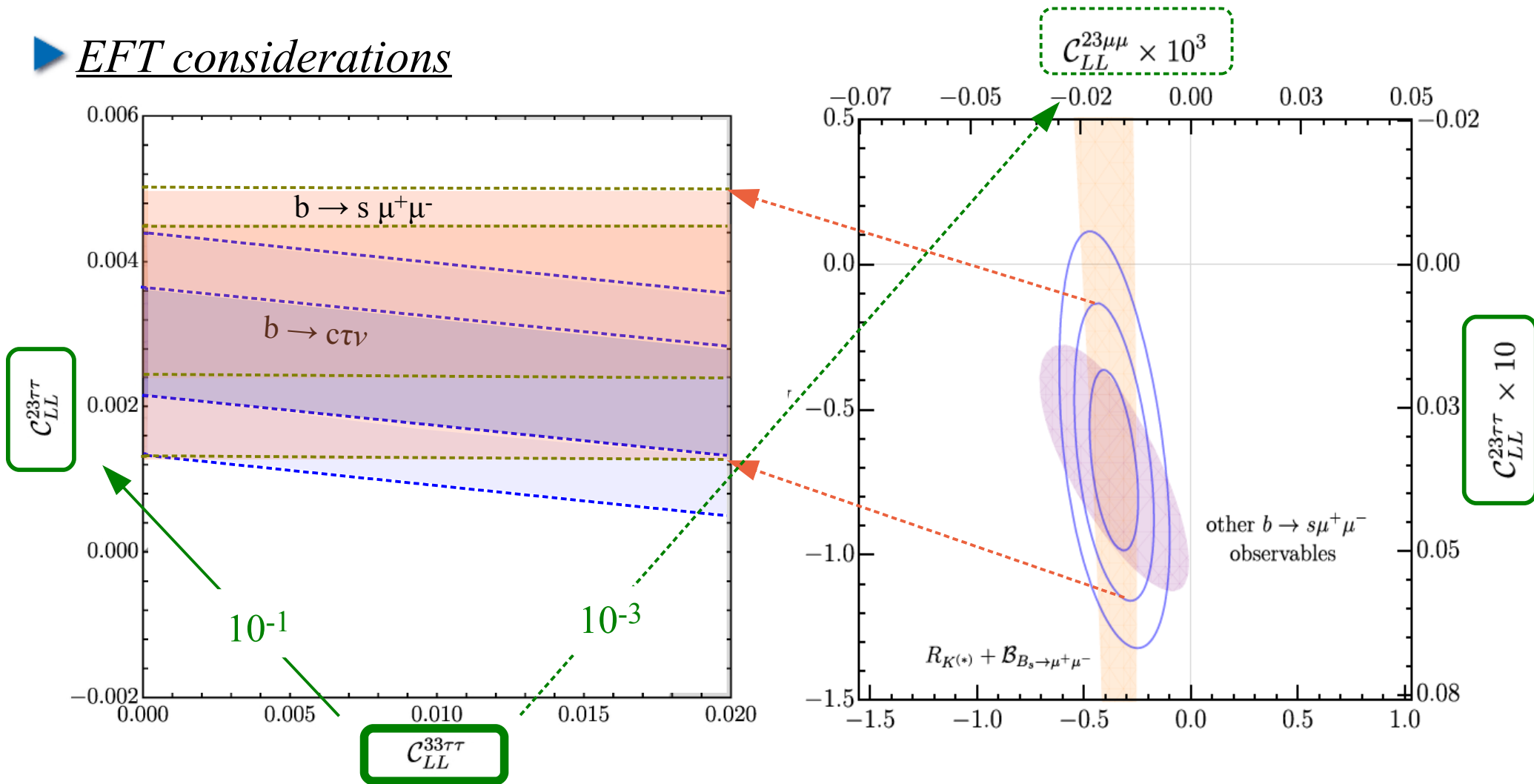
$$\frac{V_{cb} C_{LL}^{33\tau\tau} + V_{cs} C_{LL}^{23\tau\tau}}{V_{cb}}$$

Pattern emerging from data:

$$\mathcal{O}_{LL}^{ij\alpha\beta} = (\bar{q}_L^i \gamma_\mu \ell_L^\alpha) (\bar{\ell}_L^\beta \gamma_\mu q_L^j)$$

- ✓ $O(10^{-1})$ suppress. for each 2nd gen. q_L or l_L [recall $|V_{ts}| \sim 0.4 \times 10^{-1}$]

► EFT considerations



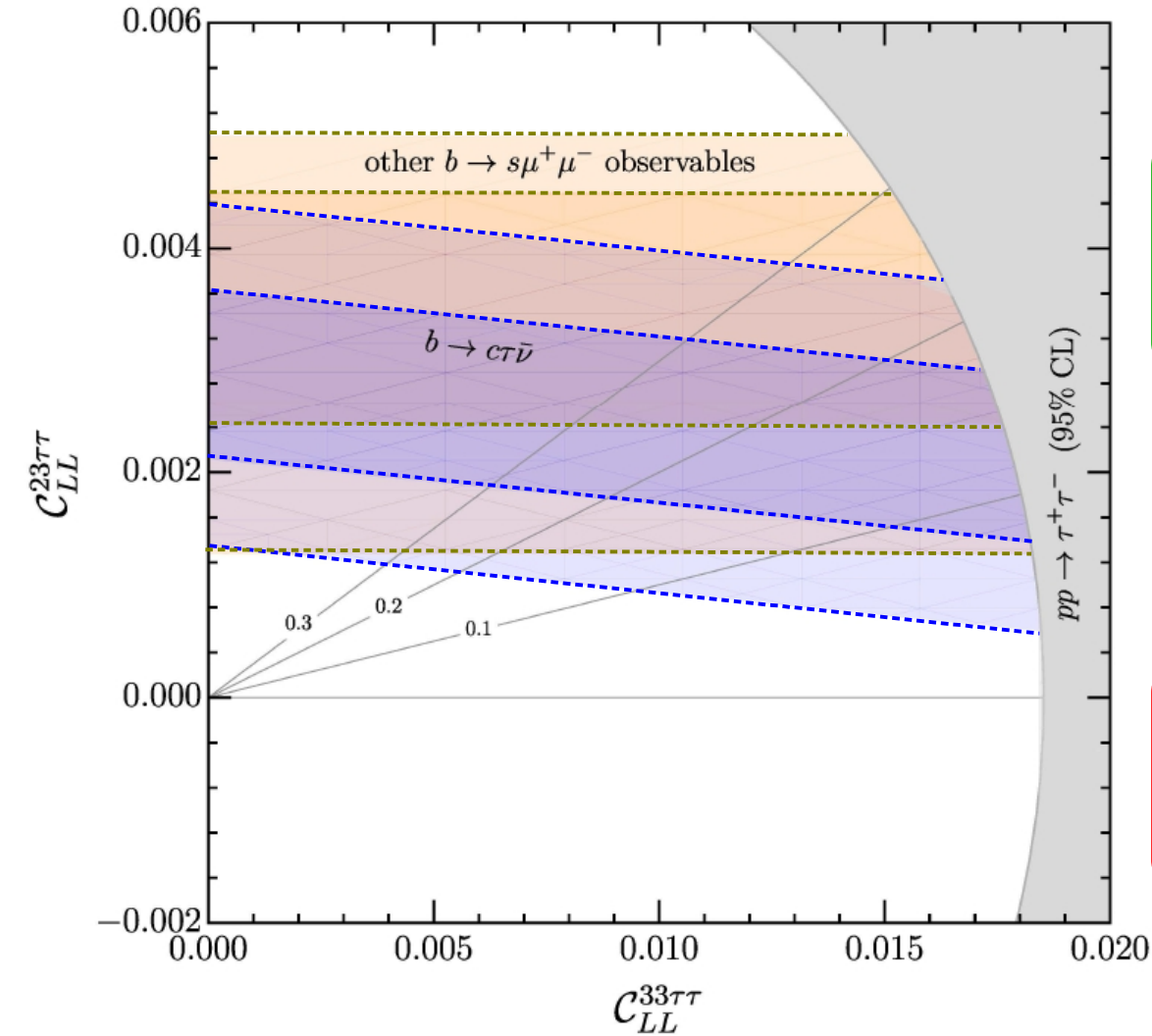
Pattern emerging from data:

$$\mathcal{O}_{LL}^{ij\alpha\beta} = (\bar{q}_L^i \gamma_\mu \ell_L^\alpha) (\bar{\ell}_L^\beta \gamma_\mu q_L^j)$$

- ✓ $O(10^{-1})$ suppress. for each 2nd gen. q_L or l_L [recall $|V_{ts}| \sim 0.4 \times 10^{-1}$]
- ✓ Nice consistency among the 2 sets of anomalies

► EFT considerations

$$\mathcal{O}_{LL}^{ij\alpha\beta} = (\bar{q}_L^i \gamma_\mu \ell_L^\alpha) (\bar{\ell}_L^\beta \gamma_\mu q_L^j) = \frac{1}{2} \left[\mathcal{O}_{\ell q}^{(1)} + \mathcal{O}_{\ell q}^{(3)} \right]^{ij\alpha\beta}$$



Pattern emerging from data:

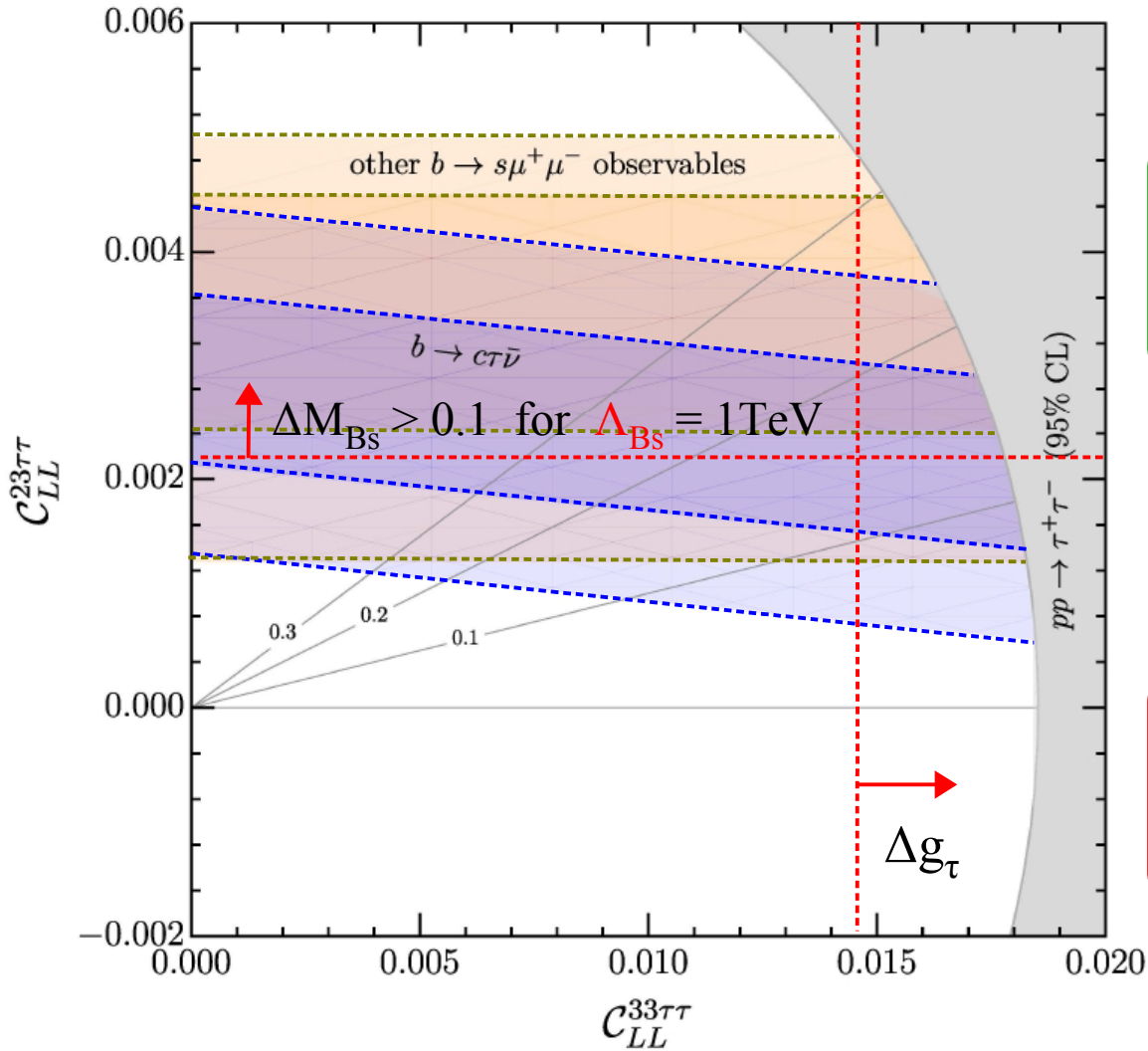
- ✓ $O(10^{-1})$ for each 2nd gen. q_L or l_L
- ✓ Nice consistency among the two sets of anomalies

What we do not see (*seem to call for an additional \sim loop suppression*):

- ✗ Four-quarks ($\Delta F=2$)
- ✗ Four-leptons ($\tau \rightarrow \mu\nu\nu$)
- ✗ Semi-leptonic $O^{(1-3)}$ ($b \rightarrow s\nu\nu$)

► EFT considerations

$$\mathcal{O}_{LL}^{ij\alpha\beta} = (\bar{q}_L^i \gamma_\mu \ell_L^\alpha) (\bar{\ell}_L^\beta \gamma_\mu q_L^j) = \frac{1}{2} \left[\mathcal{O}_{\ell q}^{(1)} + \mathcal{O}_{\ell q}^{(3)} \right]^{ij\alpha\beta}$$

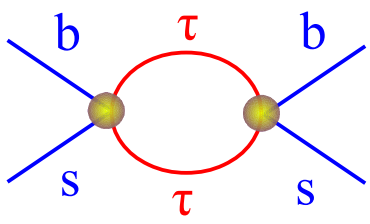


Pattern emerging from data:

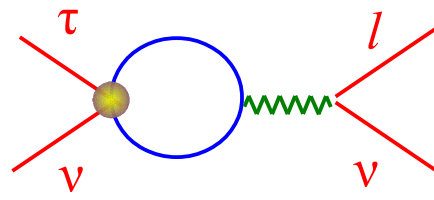
- ✓ $O(10^{-1})$ for each 2nd gen. q_L or l_L
- ✓ Nice consistency among the two sets of anomalies

What we do not see (*seem to call for an additional ~ loop suppression*):

- ✗ Four-quarks ($\Delta F=2$)
- ✗ Four-leptons ($\tau \rightarrow \mu \nu \nu$)
- ✗ Semi-leptonic $O^{(1-3)}$ ($b \rightarrow s \nu \nu$)



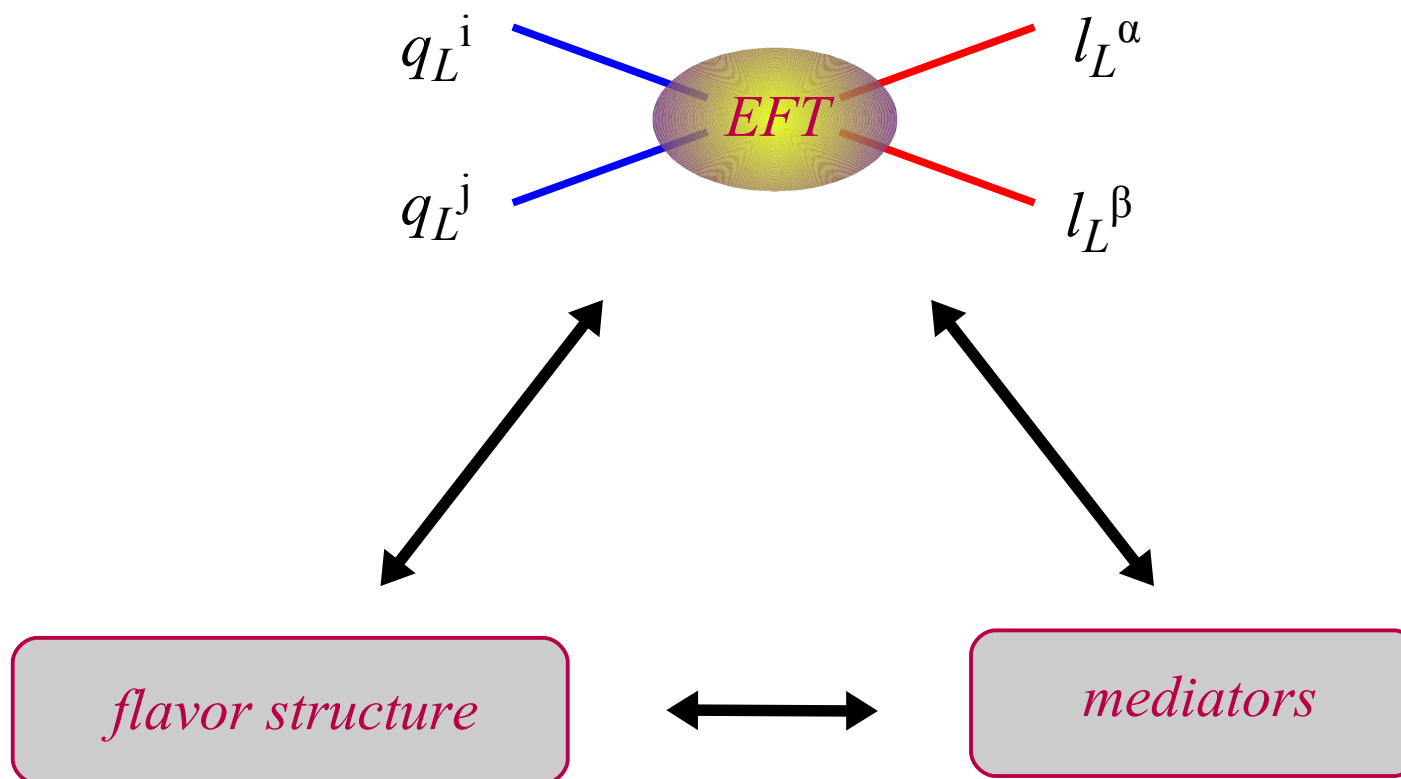
$$\Delta M_{B_s} \sim (C^{23\tau\tau})^2 \Lambda_{B_s}^2$$



$$\Delta g_\tau \sim (C^{33\tau\tau}) \log(\Lambda/m_t)$$

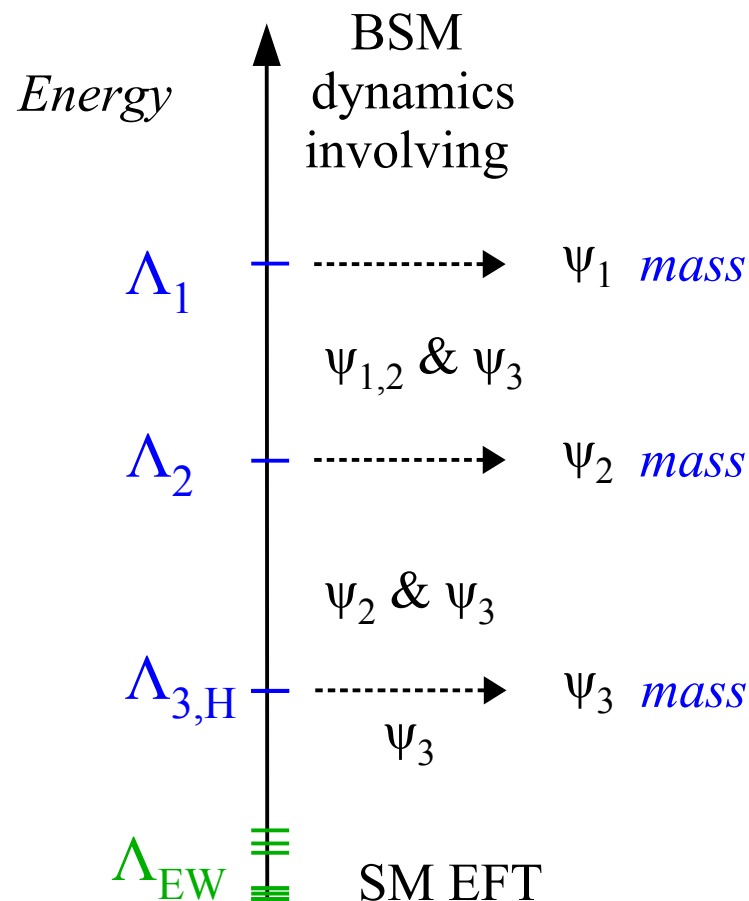
► From EFT to simplified models

To move from the EFT toward more complete/ambitious models, we need to address two general aspects: the *flavor structure* of the underlying theory, and the nature of the possible *mediators*



► From EFT to simplified models [the flavor structure]

Multi-scale picture @ origin of flavor:



Barbieri '21
 Allwicher, GI, Thomsen '20
 ⋮
 Bordone *et al.* '17
 Panico & Pomarol '16
 ⋮
 Dvali & Shifman '00

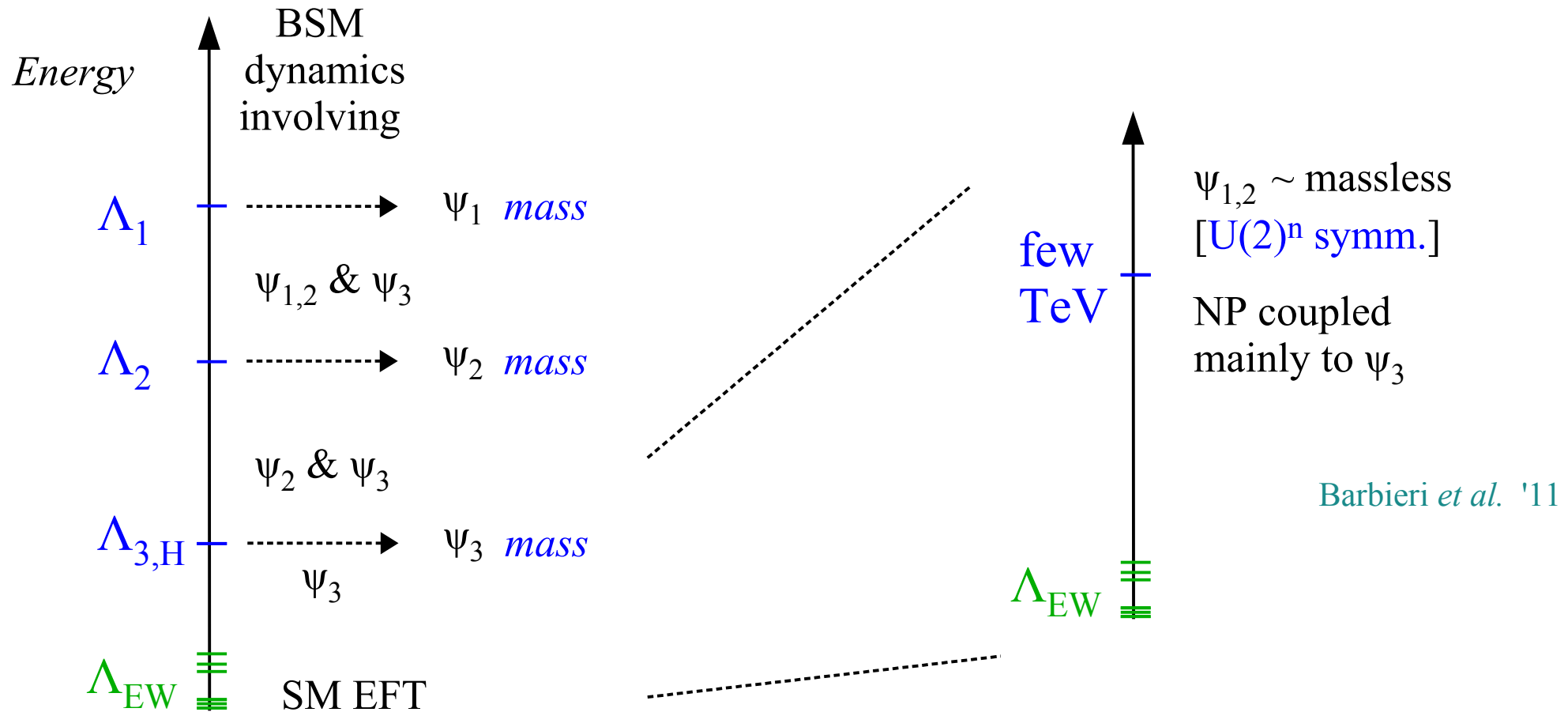
Main idea:

- Flavor **non-universal interactions** already at the **TeV scale**:
- **1st & 2nd gen.** have small masses because they are coupled to **NP at heavier scales**

~~3 gen. = "identical copies"
 up to high energies~~

► From EFT to simplified models [the flavor structure]

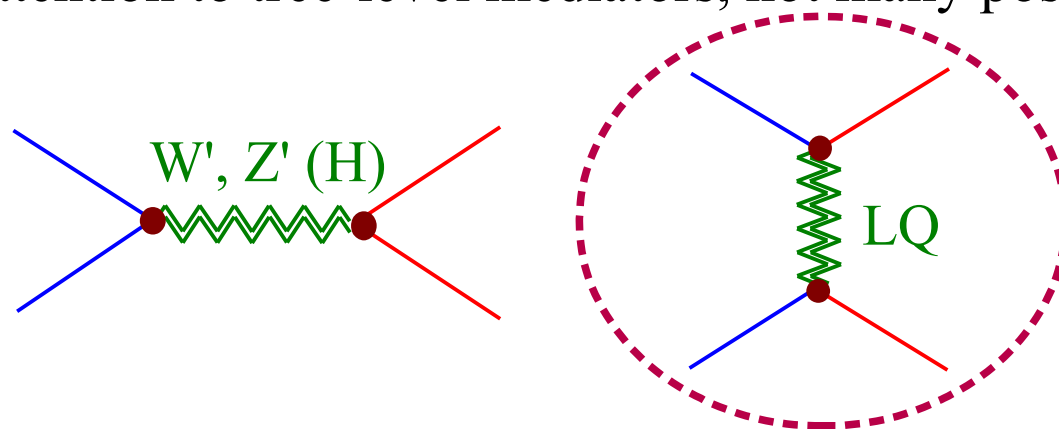
Multi-scale picture @ origin of flavor:



$$\mathcal{L}_{\text{SM-EFT}} = \mathcal{L}_{\text{gauge}} + \mathcal{L}_{\text{Higgs}} + \underbrace{\mathcal{L}_Y + \sum_i \frac{1}{\Lambda_i^{d-4}} \mathbf{O}_i^{d \geq 5}}_{\text{Non-trivial UV imprints}}$$

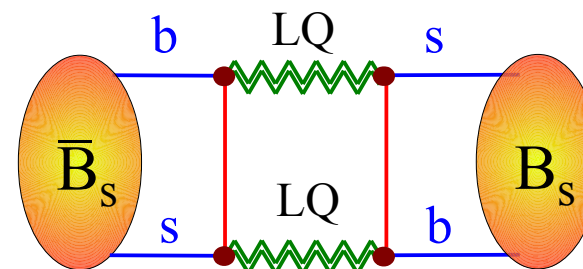
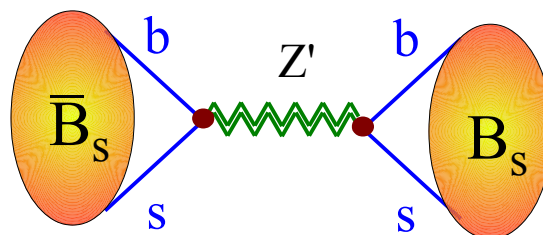
► From EFT to simplified models [the possible mediators]

Which mediators can generate the effective operators required for by the EFT fit?
If we restrict the attention to tree-level mediators, not many possibilities...



LQ (both scalar and vectors) have two general strong advantages with respect to the other mediators:

I. $\Delta F=2$ &
 $\tau \rightarrow l\nu\nu$



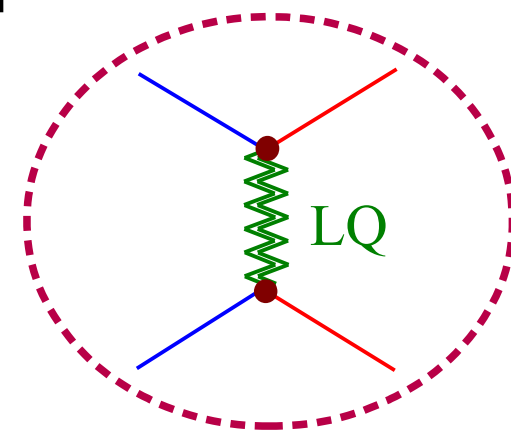
II. Direct searches:

3rd gen. LQ are also in better shape as far as direct searches are concerned (*contrary to Z'...*).

► From EFT to simplified models [the possible mediators]

Leptoquarks suffered of an (*undeserved*) “bad reputation” for two main reasons:

- Could mediate proton decay → **not a general feature of the LQ: it depends on the model...!**
[*e.g. not the case in the Pati-Salam model*]
- Severe bounds from processes involving μ & e (such as $K_L \rightarrow \mu e$)
→ **avoided with non-trivial flavor structure** [*e.g. non-univ. interactions*]



On the other hand, they are a “natural” feature in many SM extensions
→ “Renaissance” of LQ models (*to explain the anomalies, but not only...*):

- | | | |
|--|--|---|
| <ul style="list-style-type: none"> • Scalar LQ as PNG
Gripaios, '10
Gripaios, Nardecchia, Renner, '14
Marzocca '18 | <ul style="list-style-type: none"> • Scalar LQ from GUTs & \mathcal{R} SUSY
Hiller & Schmaltz, '14; Becirevic <i>et al.</i> '16,
Fajfer <i>et al.</i> '15-'17; Dorsner <i>et al.</i> '17;
Crivellin <i>et al.</i> '17; Altmannshofer <i>et al.</i> '17
Trifinopoulos '18, Becirevic <i>et al.</i> '18 + ... | <ul style="list-style-type: none"> • Vector LQ in GUT gauge models

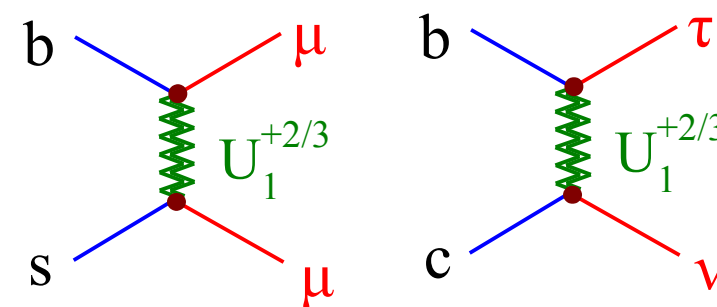
Assad <i>et al.</i> '17
Di Luzio <i>et al.</i> '17
Bordone <i>et al.</i> '17
Heeck & Teresi '18
+ ... |
| <ul style="list-style-type: none"> • Vector LQ as techni-fermion resonances
Barbieri <i>et al.</i> '15; Buttazzo <i>et al.</i> '16,
Barbieri, Murphy, Senia, '17 + ... | <ul style="list-style-type: none"> • LQ as Kaluza-Klein excit.
Megias, Quiros, Salas '17
Megias, Panico, Pujolas, Quiros '17
Blanke, Crivellin, '18 + ... | |

► From EFT to simplified models [the possible mediators]

Which LQ explains which anomaly?

	Model	$R_{K(*)}$	$R_{D(*)}$	$R_{K(*)}$ & $R_{D(*)}$
Scalars	$S_1 = (\mathbf{3}, \mathbf{1})_{-1/3}$	✗	✓	✗
	$R_2 = (\mathbf{3}, \mathbf{2})_{7/6}$	✗	✓	✗
	$\tilde{R}_2 = (\mathbf{3}, \mathbf{2})_{1/6}$	✗	✗	✗
Vector	$S_3 = (\mathbf{3}, \mathbf{3})_{-1/3}$	✓	✗	✗
	$U_1 = (\mathbf{3}, \mathbf{1})_{2/3}$	✓	✓	✓
	$U_3 = (\mathbf{3}, \mathbf{3})_{2/3}$	✓	✗	✗

Angelescu, Becirevic, DAF, Sumensari [1808.08179]



LQ of the Pati-Salam gauge group:

$SU(4) \times SU(2)_L \times SU(2)_R$

Barbieri, GI, Patteri, Senia '15

- mediator: U_1
- flavor structure: $U(2)^n$

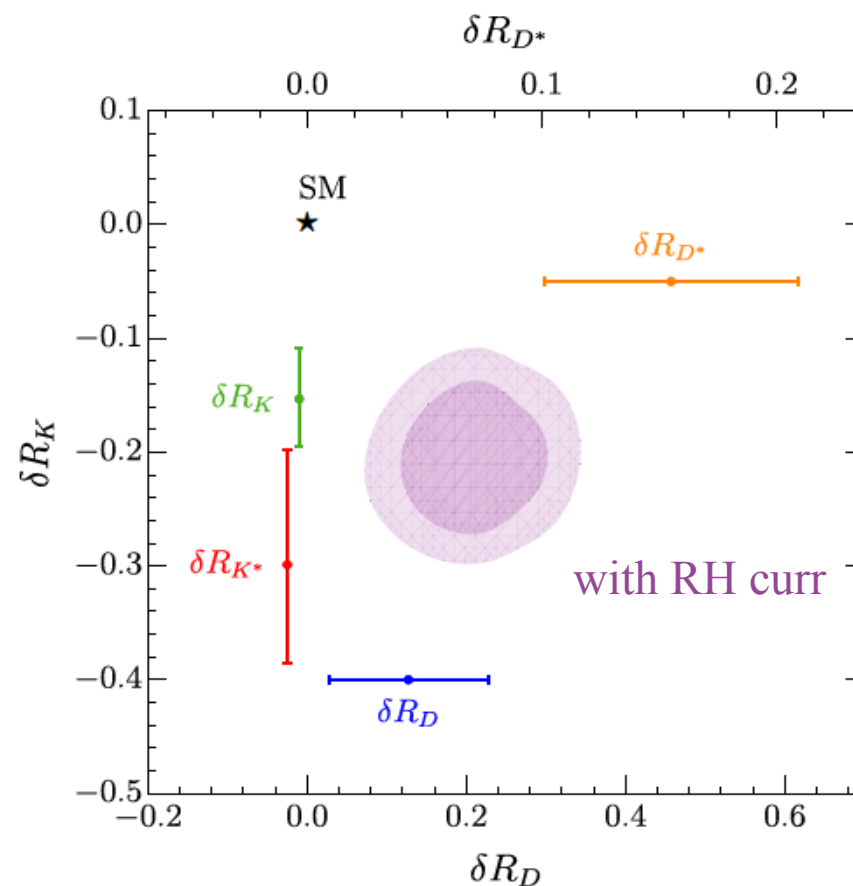
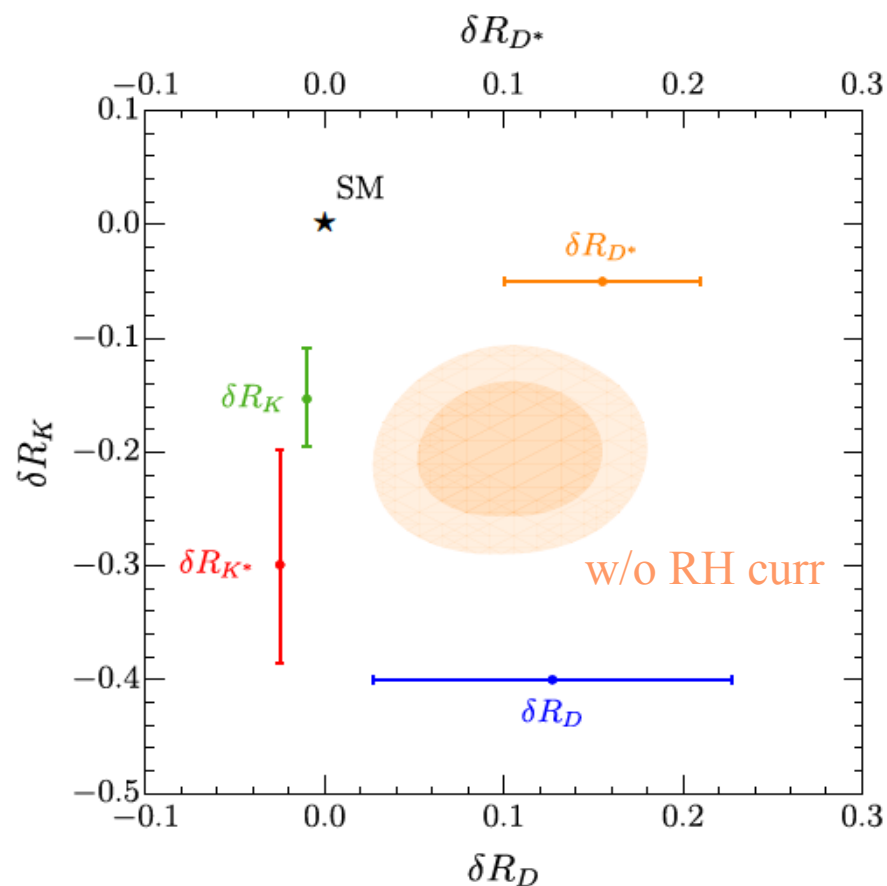
(approx. flavor symmetry resulting from the multi-scale picture)

► From EFT to simplified models [the possible mediators]

Considering the U_1 only

$$\mathcal{L} \supset \frac{g_U}{\sqrt{2}} U_1^\mu \left[\beta_{i\alpha}^L (\bar{q}_L^i \gamma_\mu \ell_L^\alpha) - \beta_{i\alpha}^R (\bar{d}_R^i \gamma_\mu e_R^\alpha) \right] + \text{h.c.}$$

and fitting all low-energy data leads to an excellent description of present data:

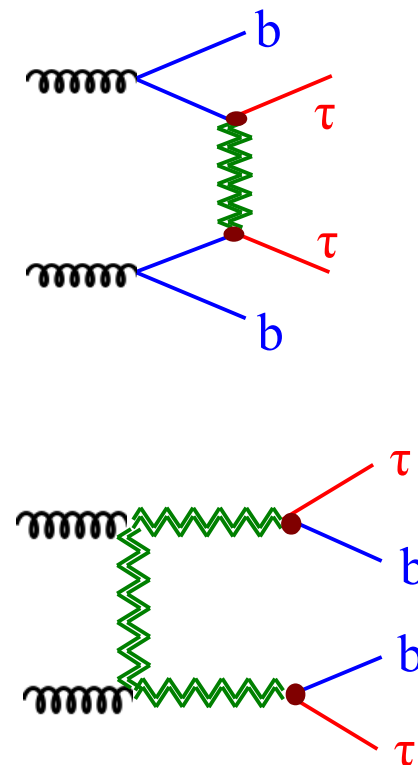
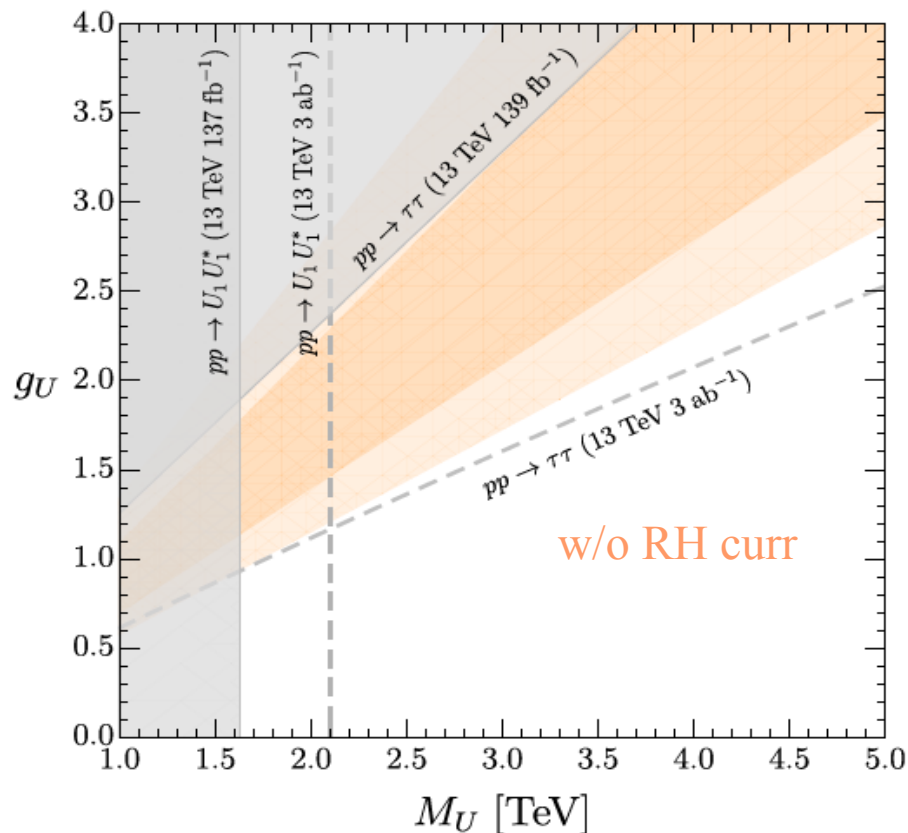


► From EFT to simplified models [the possible mediators]

Considering the U_1 only

$$\mathcal{L} \supset \frac{g_U}{\sqrt{2}} U_1^\mu \left[\beta_{i\alpha}^L (\bar{q}_L^i \gamma_\mu \ell_L^\alpha) - \beta_{i\alpha}^R (\bar{d}_R^i \gamma_\mu e_R^\alpha) \right] + \text{h.c.}$$

and fitting all low-energy data leads to an excellent description of present data which is fully consistent with high-pT searches [*within the reach of HL-LHC*]:



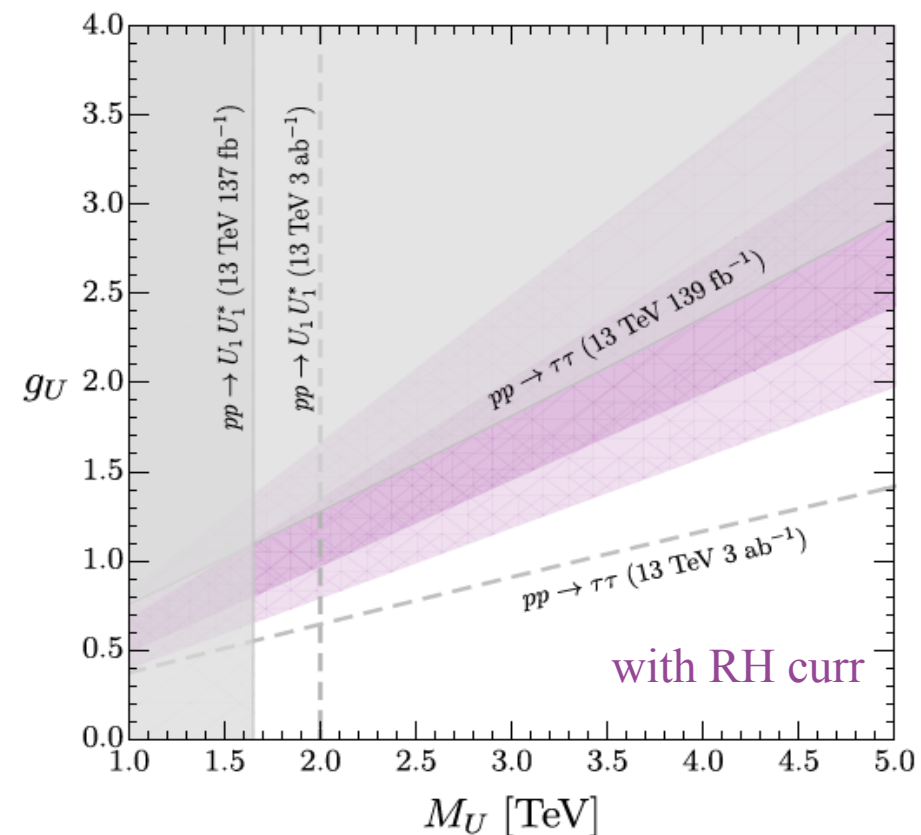
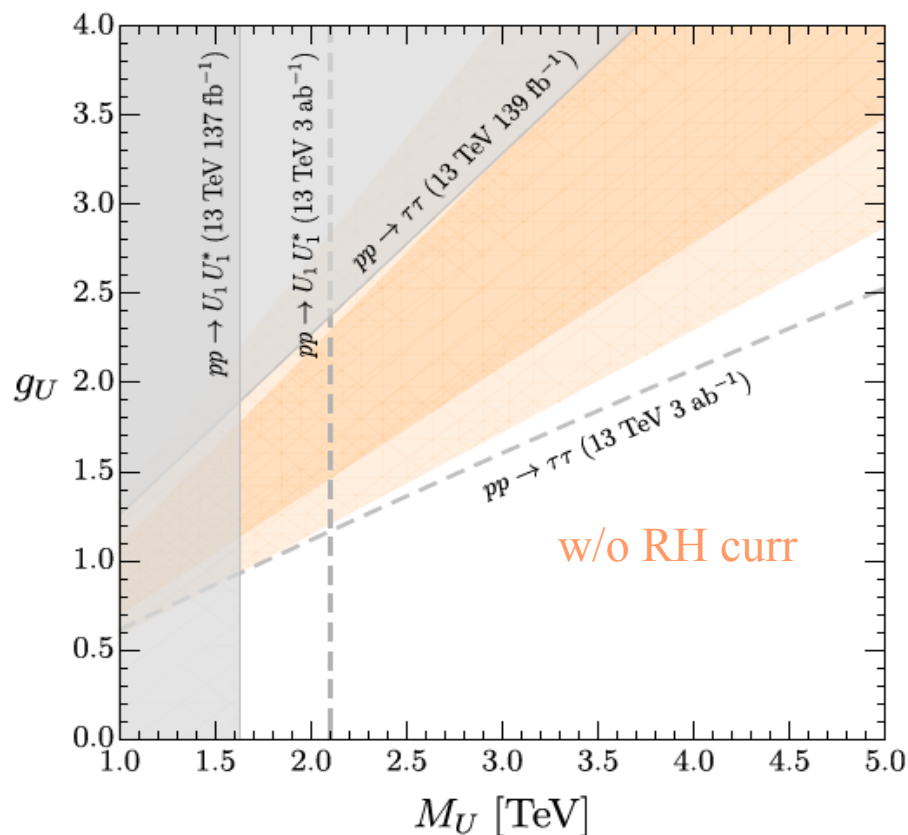
► From EFT to simplified models [the possible mediators]

Considering the U_1 only

$$\mathcal{L} \supset \frac{g_U}{\sqrt{2}} U_1^\mu \left[\beta_{i\alpha}^L (\bar{q}_{L\mu}^i \gamma_\mu \ell_L^\alpha) - \beta_{i\alpha}^R (\bar{d}_{R\mu}^i \gamma_\mu e_R^\alpha) \right] + \text{h.c.}$$

and fitting all low-energy data leads to an excellent description of present data which is fully consistent with high-pT searches [*within the reach of HL-LHC*]:

Cornella, Fuentes-Martin, Faroughi, GI, Neubert, '21



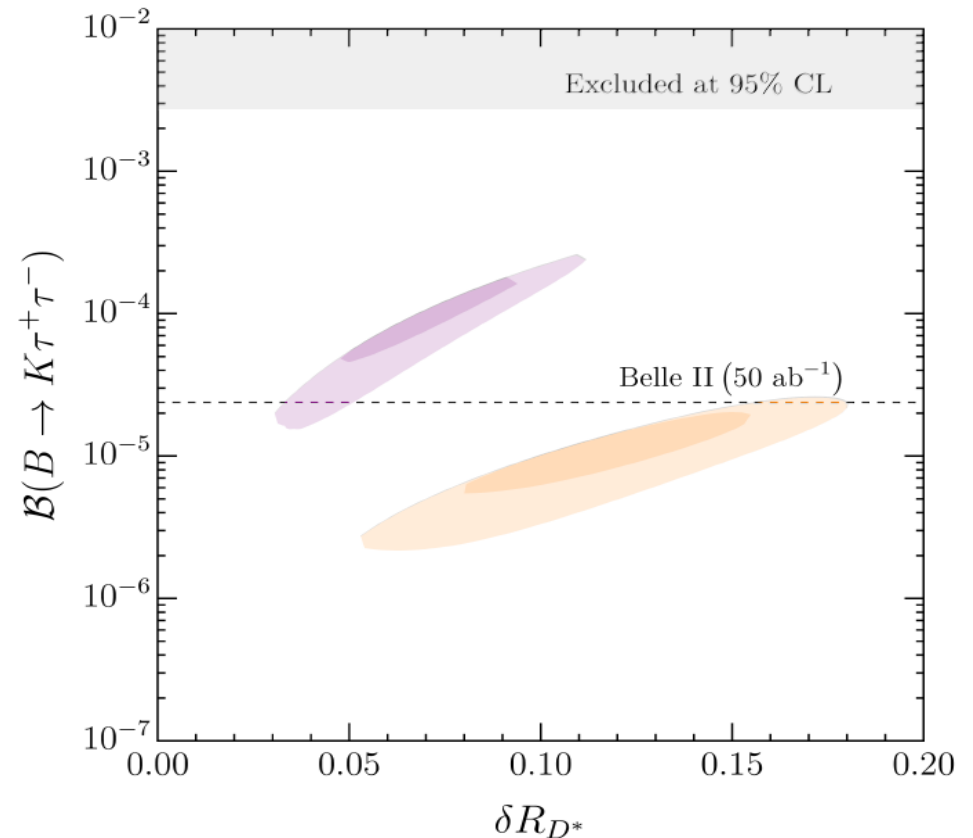
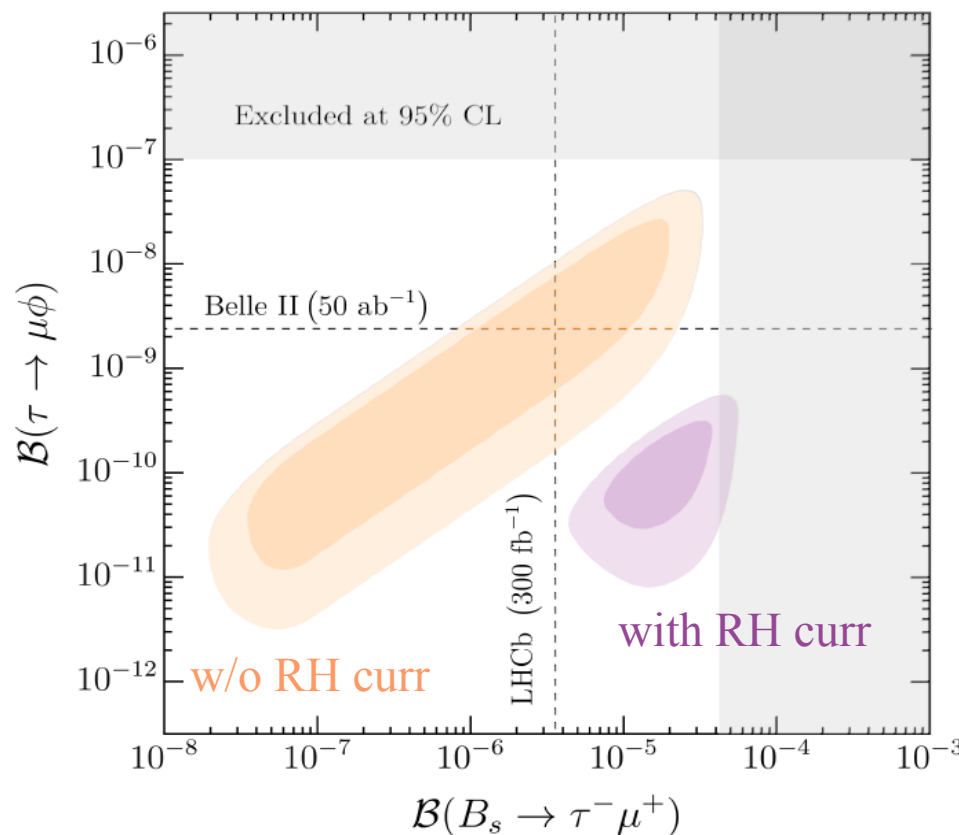
► From EFT to simplified models [the possible mediators]

Considering the U_1 only

$$\mathcal{L} \supset \frac{g_U}{\sqrt{2}} U_1^\mu \left[\beta_{i\alpha}^L (\bar{q}_{L\mu}^i \gamma_\mu \ell_L^\alpha) - \beta_{i\alpha}^R (\bar{d}_{R\mu}^i \gamma_\mu e_R^\alpha) \right] + \text{h.c.}$$

and fitting all low-energy data leads to an excellent description of present data which is fully consistent with high-pT searches, and has interesting implications for future low-energy searches:

Cornella, Fuentes-Martin, Faroughi, GI, Neubert, '21



Speculations on UV completions



► Speculations on UV completions

First observation: the Pati & Salam group, proposed in the 70's to unify quarks & leptons predicts the only massive LQ that is a good mediator for both anomalies:

Pati-Salam group: $SU(4) \times SU(2)_L \times SU(2)_R$

Fermions in SU(4):

$$\begin{bmatrix} Q_L^\alpha \\ Q_L^\beta \\ Q_L^\gamma \\ L_L \end{bmatrix} \quad \begin{bmatrix} Q_R^\alpha \\ Q_R^\beta \\ Q_R^\gamma \\ L_R \end{bmatrix}$$

Main Pati-Salam idea:
Lepton number as “the 4th color”

The massive LQ [U_1] arise from the breaking $SU(4) \rightarrow SU(3)_C \times U(1)_{B-L}$

$$SU(4) \sim \left[\begin{array}{c|c} SU(3)_C & 0 \\ \hline 0 & 0 \end{array} \right] \quad \left[\begin{array}{c|c} 0 & LQ \\ \hline LQ & \end{array} \right] \quad \left[\begin{array}{c|c} \frac{1}{3} & 0 \\ \hline 0 & -1 \end{array} \right]$$

► Speculations on UV completions

First observation: the Pati & Salam group, proposed in the 70's to unify quarks & leptons predicts the only massive LQ that is a good mediator for both anomalies:

Pati-Salam group: $SU(4) \times SU(2)_L \times SU(2)_R$

Fermions in $SU(4)$:

$$\begin{bmatrix} Q_L^\alpha \\ Q_L^\beta \\ Q_L^\gamma \\ L_L \end{bmatrix} \quad \begin{bmatrix} Q_R^\alpha \\ Q_R^\beta \\ Q_R^\gamma \\ L_R \end{bmatrix}$$

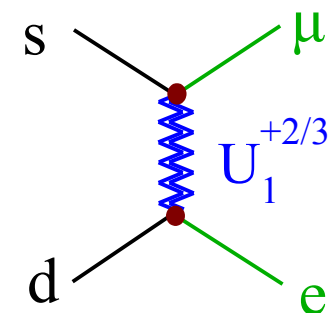
Main Pati-Salam idea:
Lepton number as “the 4th color”

The massive LQ [U_1] arise from the breaking $SU(4) \rightarrow SU(3)_C \times U(1)_{B-L}$

The problem of the “original PS model” are the strong bounds on the LQ couplings to 1st & 2nd generations [e.g. $M > 200 \text{ TeV}$ from $K_L \rightarrow \mu e$]

Attempts to solve this problem simply adding extra fermions or scalars

Calibbi, Crivellin, Li, '17;
Fornal, Gadam, Grinstein, '18
Heeck, Teresi, '18

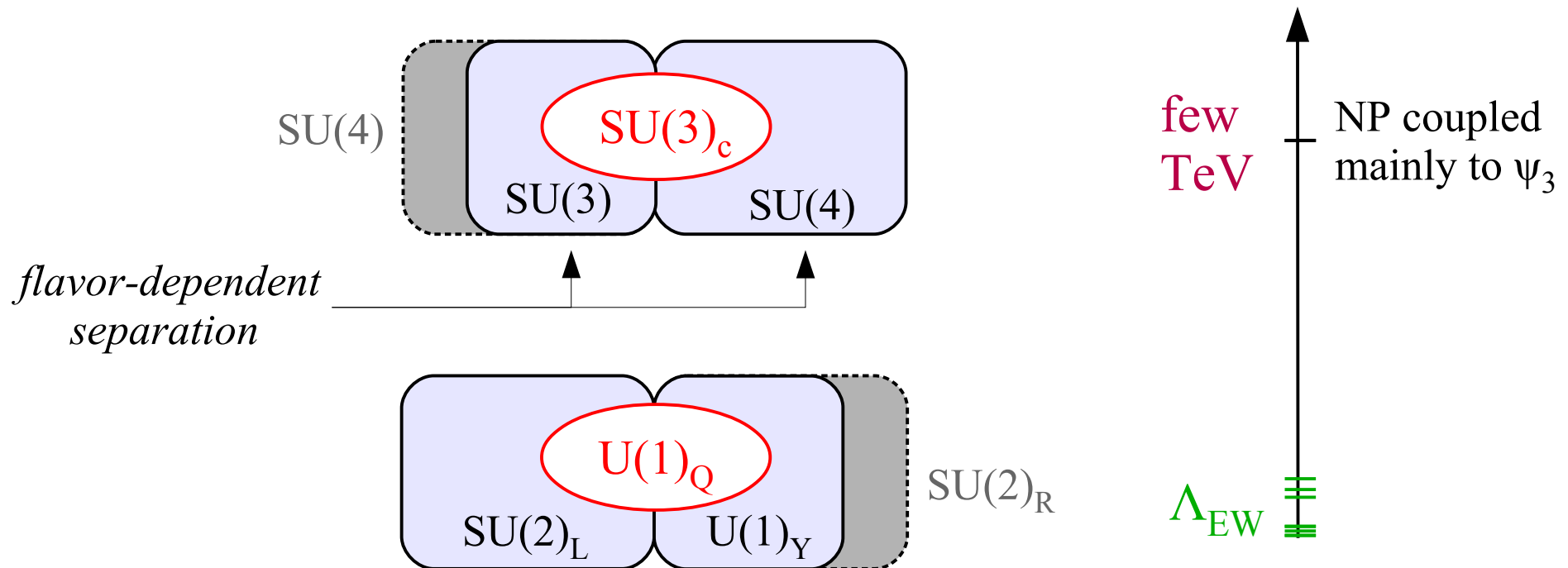


► Speculations on UV completions

Second observation: we can “protect” the light families charging under SU(4) only the 3rd gen. or, more generally, “separating” the universal SU(3) component

PS group: $SU(4) \times SU(2)_L \times SU(2)_R$ • *flavor universality*

4321 models: $SU(4) \times SU(3) \times G_{EW} = \begin{cases} SU(2)_L \times SU(2)_R \\ SU(2)_L \times U(1)_Y \end{cases}$



► Speculations on UV completions

Second observation: we can “protect” the light families charging under SU(4) only the 3rd gen. or, more generally, “separating” the universal SU(3) component

PS group: $SU(4) \times SU(2)_L \times SU(2)_R$ • *flavor universality*

4321 models:

$$SU(4) \times SU(3) \times G_{EW} = \begin{cases} SU(2)_L \times SU(2)_R \\ SU(2)_L \times U(1)_Y \end{cases}$$

• *Non-universality via mixing*

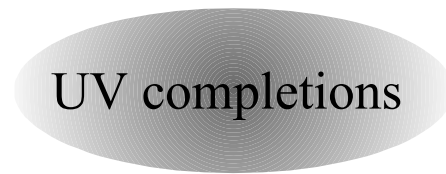
$$SU(4) \times SU(3)$$

$$SU(4)_3 \times SU(3)_{1,2}$$

• *Accidental $U(2)^5$ flavor symm. in the gauge sect.*

$$SU(3) \times G_{EW} \times G_{HC}$$

Barbieri, Tesi '17



$$SU(4)_h \times SU(4)_l \times G_{EW} \times G_{HC}$$

Fuentes-Martin & Stangl '20

$$SU(4) \times SU(3) \times G_{EW}$$

Di Luzio, Greljo, Nardecchia, '17

$$[PS]^3 = [SU(4) \times G_{EW}]^3$$

Bordone et al. '17

$$[PS]_{\text{warped-5d, 3-branes}}$$

Fuentes-Martin et al. '20 + work in prog.

► Speculations on UV completions

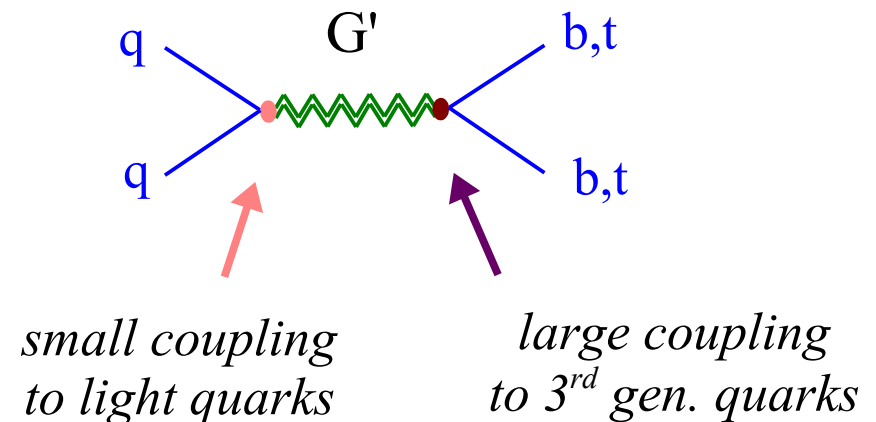
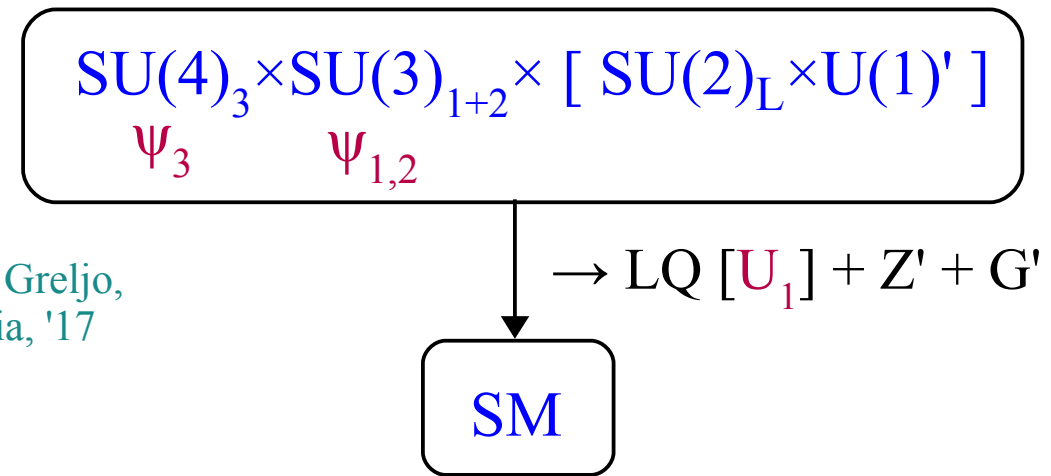
In most *PS-extended models* collider and low-energy pheno are controlled by the effective 4321 gauge group that rules TeV-scale dynamics

Di Luzio, Greljo, Nardecchia, '17

Despite the apparent complexity, the construction is highly constrained

- Positive features the EFT reproduced
 - Calculability of $\Delta F=2$ processes
 - Precise predictions for **high-pT data**
- consistent with present data !

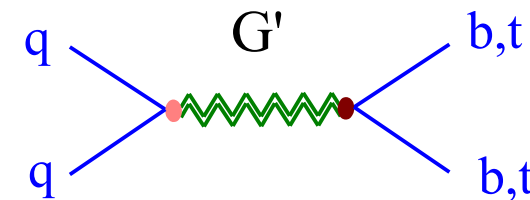
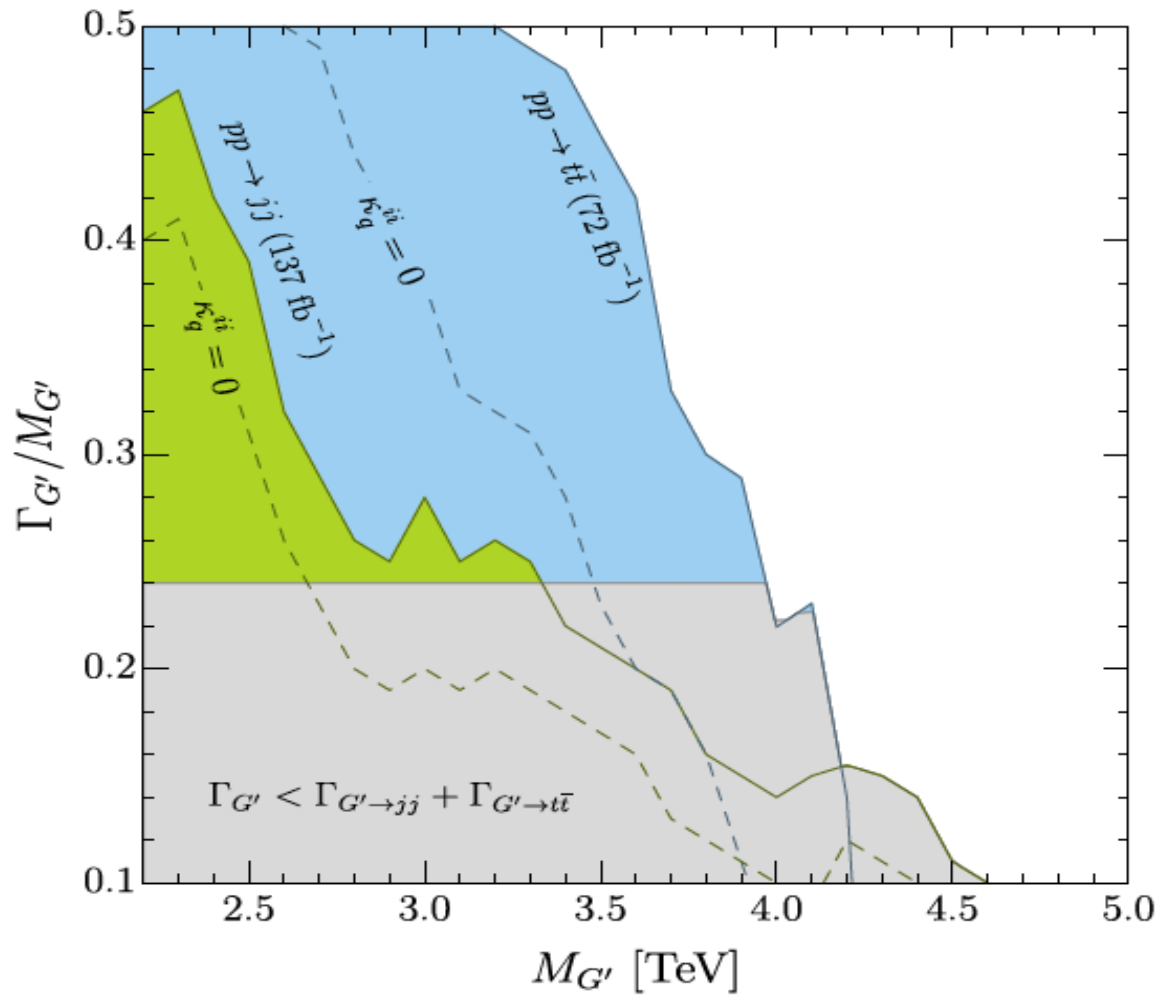
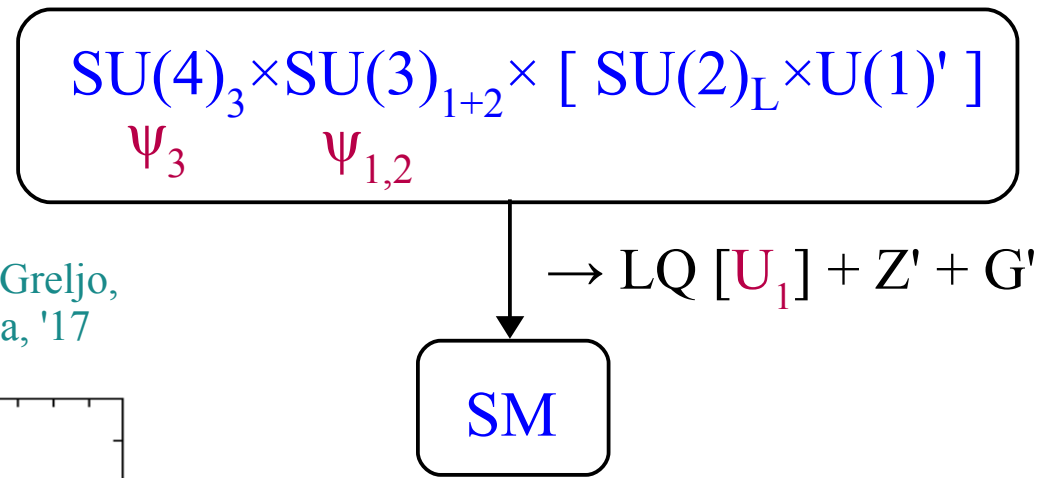
New striking collider signature:
 G' (“coloron” = heavy color octet)



► Speculations on UV completions

In most *PS-extended models* collider and low-energy pheno are controlled by the effective 4321 gauge group that rules TeV-scale dynamics

Di Luzio, Greljo, Nardecchia, '17

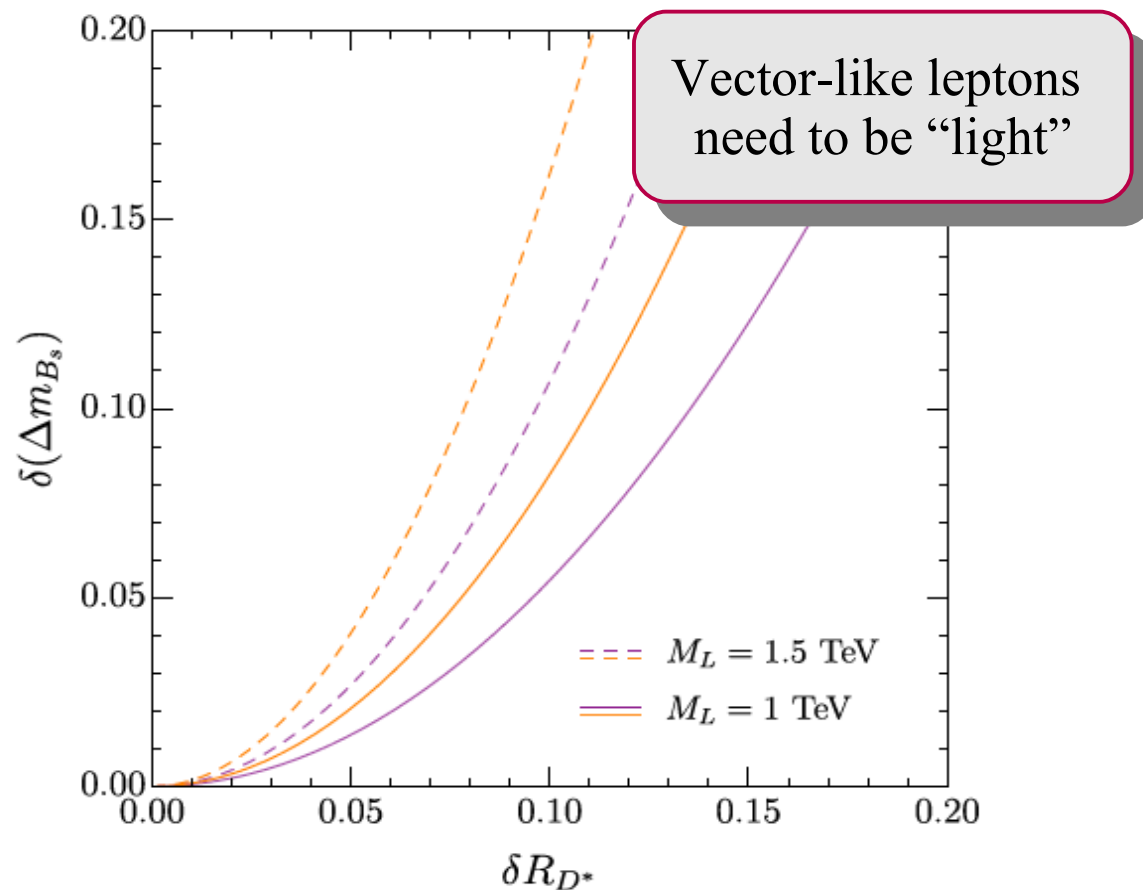
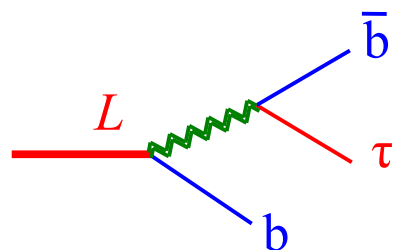
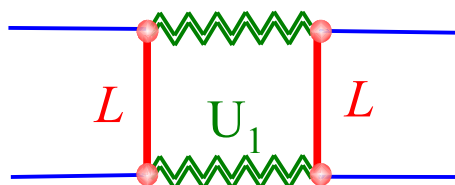


Presently, the strongest constraint on the energy scale of the model [from $pp \rightarrow t\bar{t}$]

► Speculations on UV completions

A second important ingredient of this class of UV models are **vector-like quarks** and **vector-like leptons**, which play a key role in both low- and high-energy observables

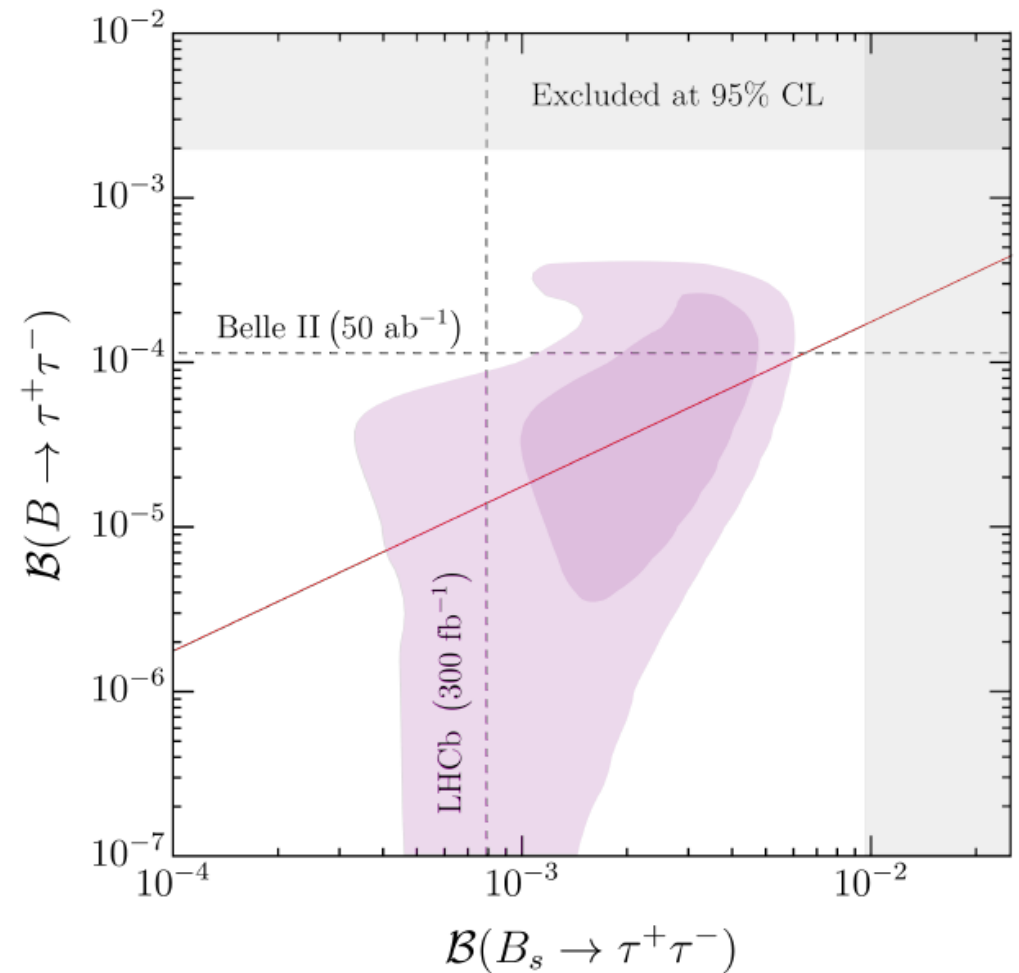
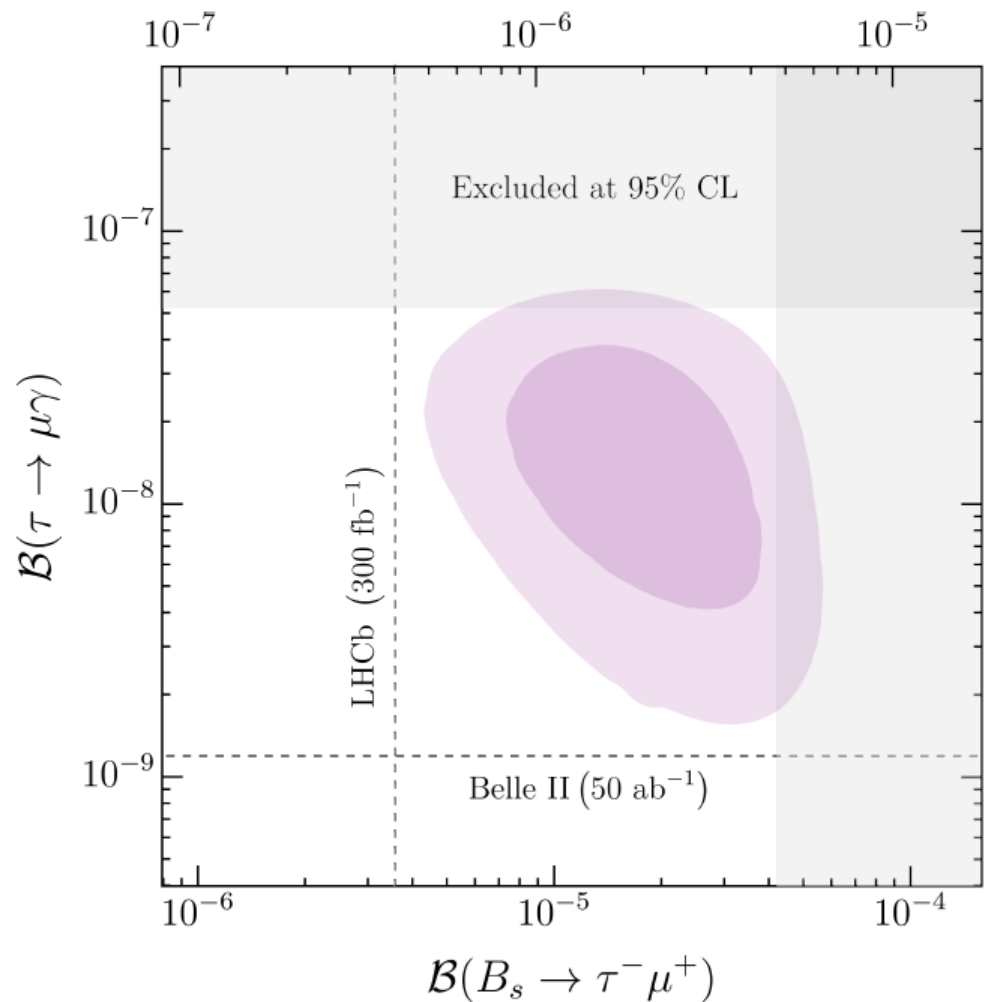
E.g.: B_s mixing [$\Delta F=2$]



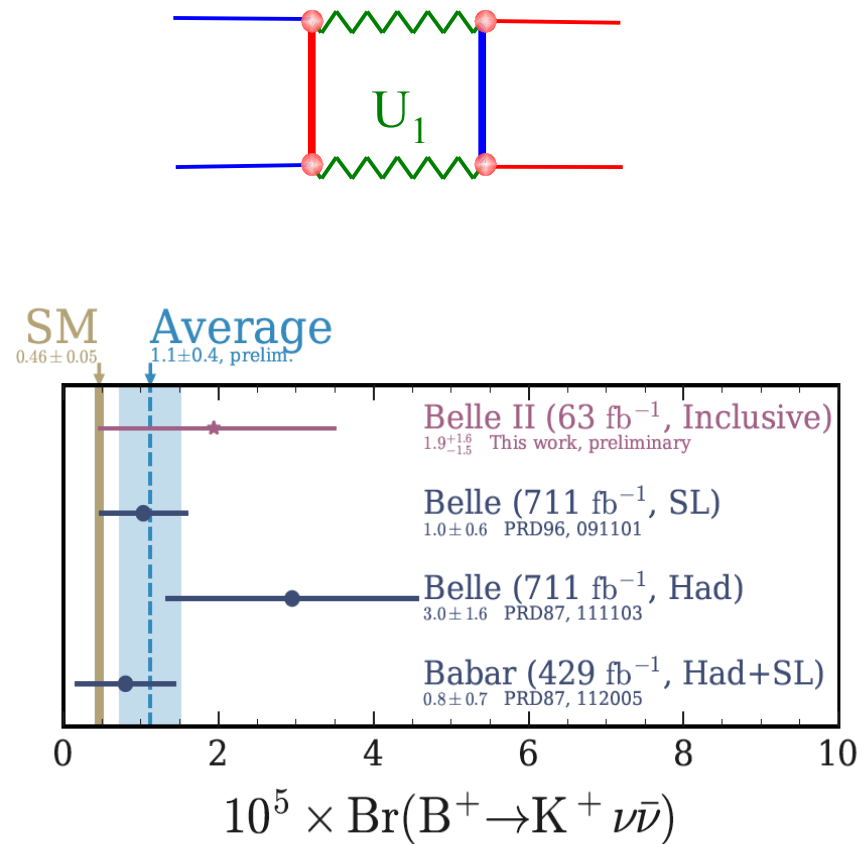
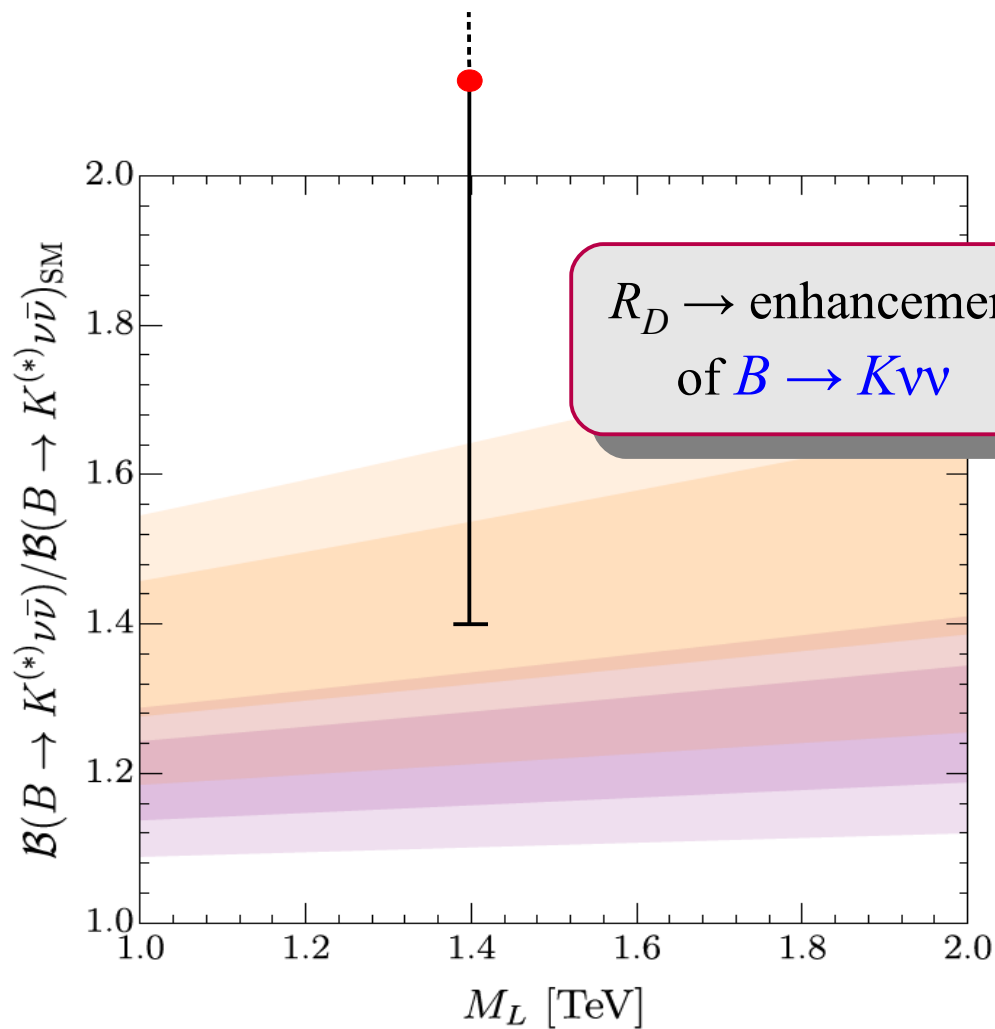
Fuentes-Martin, GI, Konig, Selimovic, '20
Cornella, Fuentes-Martin, Faroughi, GI, Neubert, '21



► Other low-energy observables



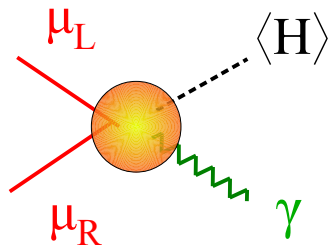
► Other low-energy observables



Fuentes-Martin, GI, Konig, Selimovic, '20
 Cornella, Fuentes-Martin, Faroughi, GI, Neubert, '21

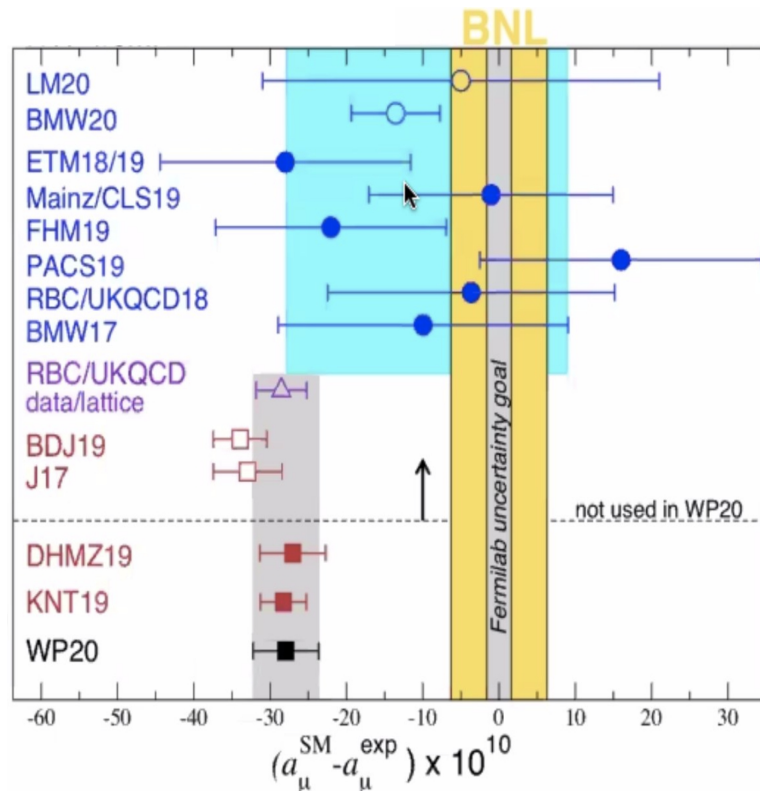
► What about $(g-2)_\mu$?

Not obvious how to reconcile the $(g-2)_\mu$ anomaly with both flavor anomalies and, more generally, with models with a “natural” flavor structure ($\leftrightarrow Y_{SM}$).



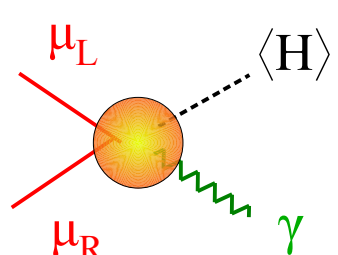
$$\Delta a_\mu = (a_\mu^{\text{exp}} - a_\mu^{\text{SM}}) \approx (a_\mu^{\text{SM}})^{\text{EW}} \rightarrow \Lambda \sim$$

100 GeV (helicity suppression $\sim m_\mu$ as in SM)
 10 TeV (remove helicity suppression. $m_\mu \rightarrow m_t$)



► What about $(g-2)_\mu$?

Not obvious how to reconcile the $(g-2)_\mu$ anomaly with both flavor anomalies and, more generally, with models with a “natural” flavor structure ($\leftrightarrow Y_{SM}$).



A Feynman diagram showing a muon loop. Two red lines represent muons, labeled μ_L and μ_R . A dashed line represents a Higgs boson, labeled $\langle H \rangle$. A green wavy line represents a photon, labeled γ . The loop is shaded orange.

$$\Delta a_\mu = (a_\mu^{\text{exp}} - a_\mu^{\text{SM}}) \approx (a_\mu^{\text{SM}})^{\text{EW}} \rightarrow \Lambda \sim \begin{cases} 100 \text{ GeV} & \text{(helicity suppression } \sim m_\mu \text{ as in SM)} \\ 10 \text{ TeV} & \text{(remove helicity suppression, } m_\mu \rightarrow m_t) \end{cases}$$

Main difficulty: strong flavor alignment needed to avoid bounds from $\mu \rightarrow e \gamma$ (10^{-5} alignment in the $1 \leftrightarrow 2$ sector) & $\tau \rightarrow \mu \gamma$ (10^{-1} alignment)

Example of recent *attempts to combine LFU anomalies & g-2*

→ $a_\mu \oplus R_K$ with special role of muons [$U(1)_{B-3L_\mu} \subset G$]

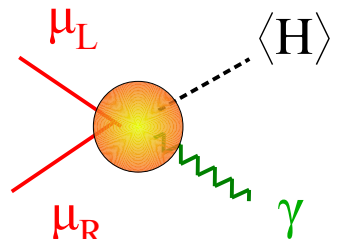
Greljo, Stangl, Thomsen '21

→ $a_\mu \oplus R_K \oplus R_D$ with 2 scalars [$S_1 + \phi^+$] and peculiar flavor struct.

Marzocca, Trifinopoulos '21

► What about $(g-2)_\mu$?

Not obvious how to reconcile the $(g-2)_\mu$ anomaly with both flavor anomalies and, more generally, with models with a “natural” flavor structure ($\leftrightarrow Y_{SM}$).



The diagram shows a muon loop with two external muon lines labeled μ_L and μ_R . A dashed line representing a Higgs boson $\langle H \rangle$ is inserted into the loop. A wavy line representing a photon γ is emitted from the loop.

$$\Delta a_\mu = (a_\mu^{\text{exp}} - a_\mu^{\text{SM}}) \approx (a_\mu^{\text{SM}})^{\text{EW}} \rightarrow \Lambda \sim \begin{cases} 100 \text{ GeV} & \text{(helicity suppression } \sim m_\mu \text{ as in SM)} \\ 10 \text{ TeV} & \text{(remove helicity suppression, } m_\mu \rightarrow m_t) \end{cases}$$

Main difficulty: strong flavor alignment needed to avoid bounds from $\mu \rightarrow e \gamma$ (10^{-5} alignment in the $1 \leftrightarrow 2$ sector) & $\tau \rightarrow \mu \gamma$ (10^{-1} alignment)

Example of recent *attempts to combine LFU anomalies & g-2*

→ $a_\mu \oplus R_K$ with special role of muons [$U(1)_{B-3L_\mu} \subset G$]

Greljo, Stangl, Thomsen '21

→ $a_\mu \oplus R_K \oplus R_D$ with 2 scalars [$S_1 + \phi^+$] and peculiar flavor struct.

Marzocca, Trifinopoulos '21

However... $(g-2)_\mu$ is more “flexible” (no generation change, necessary loop-level) → could come from light NP. *No obvious connection to the flavor anomalies*