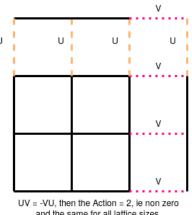
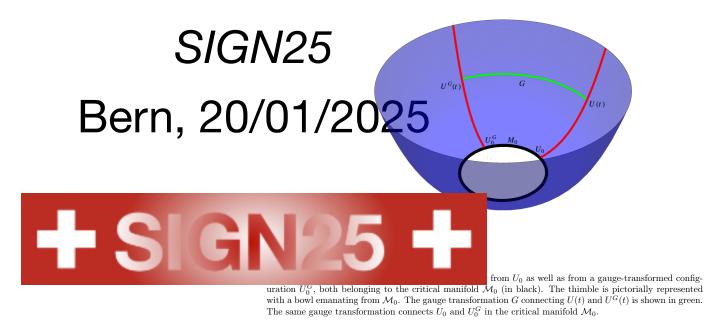
Lattice QCD sign problem as an inverse proble

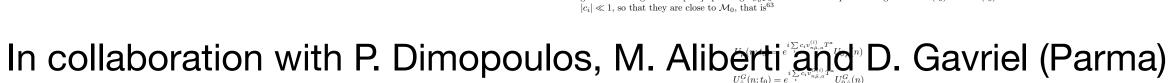


Francesco Di Renzo (University of Parma and INFN)

Figure 3: Untwisted Action again



the n_G Takagi vectors of $H(S;U_0)$ with zero Takagi value⁶² and $N_{U_0}^+\mathcal{M}_0$ spanned by the n_+ Takagi vectors of $H(S;U_0)$ with positive Takagi value. The number of such vectors is $n_+ = n - n_G$, with $n = Vd(N^2 - 1)$ the total number of degrees of freedom and $n_G = V(N^2 - 1)$ the number of gauge degrees of freedom, which means that $n_+ = V(d-1)(N^2 - 1)$. We can easily compute the Takagi vectors $\{v^{G(i)}\}$ spanning $T_{U_0^G}\mathcal{J}_0$ given the Takagi vectors $\{v^{(i)}\}$ spanning $T_{U_0^G}\mathcal{J}_0$. Consider a couple of configurations $U(t_0)$ and $U^G(t_0)$ with $|c_i| \ll 1$, so that they are close to \mathcal{M}_0 , that is⁶³



 $G(n) = e^{i g_{n,a} T^a}$

The previous considerations lead to setting $U^G_{\hat{\mu}}(n;t_0)=G(n)U_{\hat{\mu}}(n;t_0)G^{\dagger}(n+\hat{\mu})$, which imply

⁶²Directions tangent to \mathcal{M}_0 at U_0 represent infinitesimal gauge transformations around U_0 .

⁶³We generically take $|c_i| \ll 1$ in order not to leave $T_U \mathcal{J}_0$ while leaving the critical point U. This condition is automatically ensured for directions corresponding to $\lambda_i > 0$: for these directions $c_i = n_i e^{\lambda_i t_0}$ with $t_0 \to -\infty$, so that we can safely take $n_i = \mathcal{O}(1)$. For directions corresponding to $\lambda_i = 0$, however, the coefficients c_i have to be taken small explicitly.



Bielefeld Parma Collaboration (... K. Zambello ...)

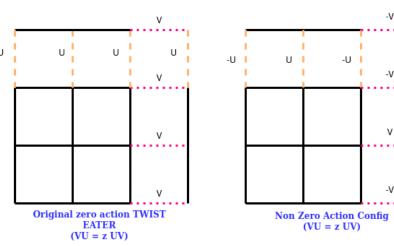
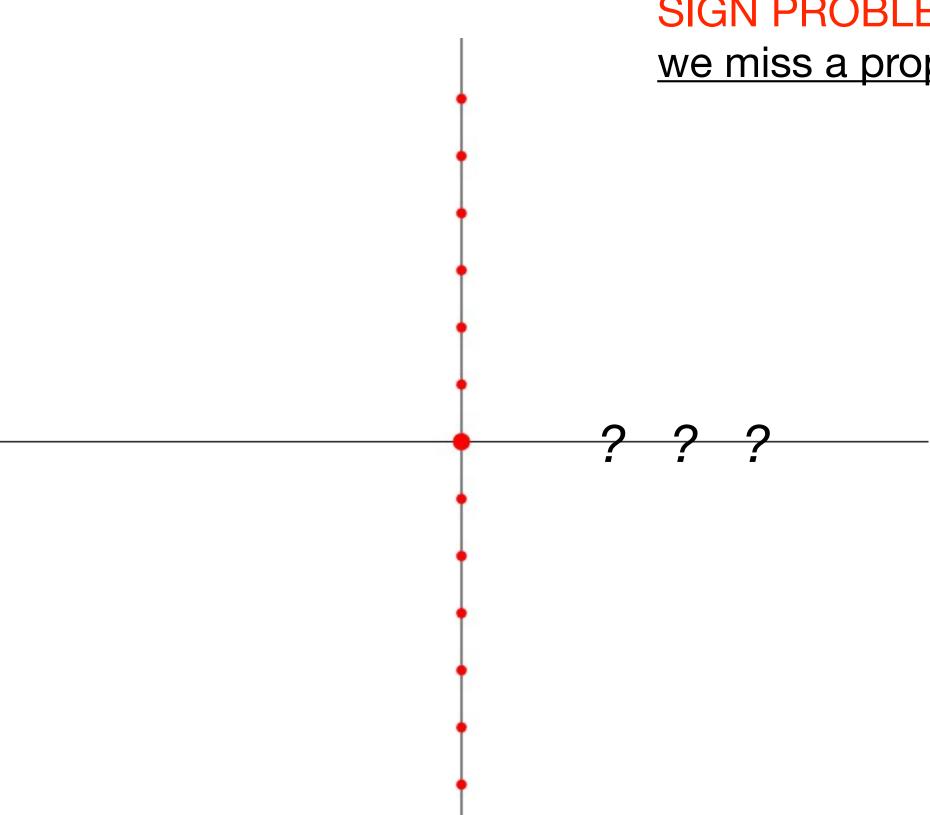


Figure 4: Twisted Action

11



Complex chemical potential plane

SIGN PROBLEM for finite density Lattice QCD:

we miss a properly defined (positive) measure in the path integral! ... no MC simulation

(... but everything is fine on the imaginary axis)



we miss a properly defined (positive) measure in the path integral! ... no MC simulation

(... but everything is fine on the imaginary axis)



Simulating QCD at finite density

Philippe de Forcrand*

Institute for Theoretical Physics, ETH Zürich, CH-8093 Zürich, Switzerland and

CERN, Physics Department, TH Unit, CH-1211 Geneva 23, Switzerland E-mail: forcrand@phys.ethz.ch

2. Sign problem

The sign problem is a necessary evil, unavoidable as soon as one integrates out the fermion fields and expresses the partition function in terms of the gauge fields. Analytic integration over each fermion species gives a factor $\det(D + m + \mu \gamma_0)$, where D is the massless Dirac operator and the last term appears when the chemical potential μ is non-zero. Now, D satisfies γ_5 -hermiticity: $\gamma_5 D \gamma_5 = D^{\dagger}$, so that

$$\gamma_5(\not D + m + \mu \gamma_0)\gamma_5 = \not D^{\dagger} + m - \mu \gamma_0 = (\not D + m - \mu^* \gamma_0)^{\dagger}$$
 (2.1)

Taking the determinant on both sides gives $\det(D + m + \mu \gamma_0) = \det^*(D + m - \mu^* \gamma_0)$, which constrains the determinant to be real only if μ is zero or pure imaginary.



we miss a properly defined (positive) measure in the path integral! ... no MC simulation

(... but everything is fine on the imaginary axis)

$$\gamma_5(D + m + \mu \gamma_0)\gamma_5 = D^{\dagger} + m - \mu \gamma_0 = (D + m - \mu^* \gamma_0)^{\dagger}$$

$$\det(D \!\!\!/ + m + \mu \gamma_0) = \det^*(D \!\!\!/ + m - \mu^* \gamma_0)$$

Fermionic determinant is REAL only for ZERO or IMAGINARY values of the chemical potential!



we miss a properly defined (positive) measure in the path integral! ... no MC simulation

(... but everything is fine on the imaginary axis)

Mainly two working solutions:

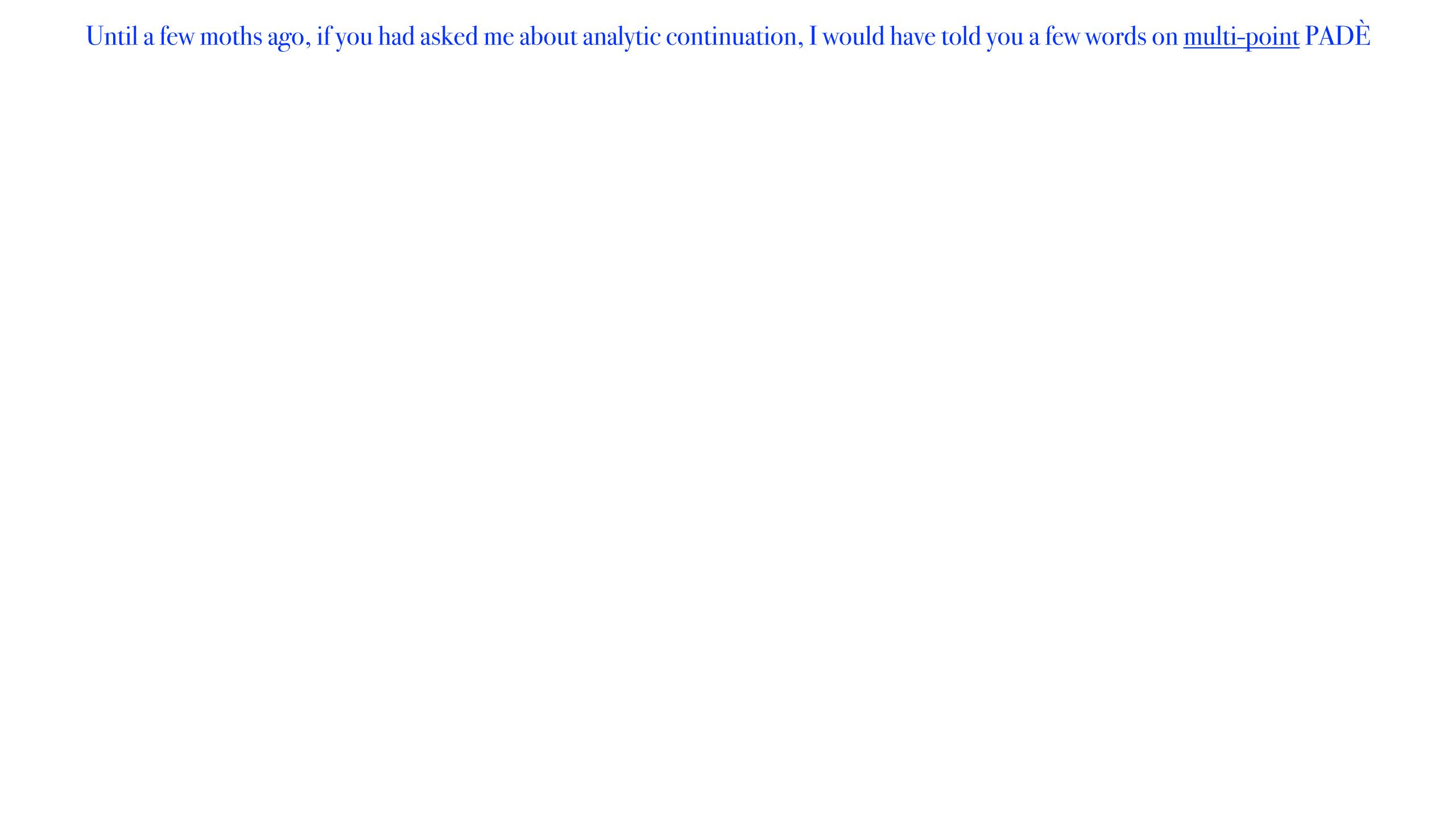
- Compute Taylor expansions at $\mu_B = 0$
- Compute on the imaginary axis $\mu_B = i\mu_I$

The two solutions are obviously related ... and both imply (strictly speaking) an ANALYTIC CONTINUATION

There are tensions in between differente results for Taylor coefficients in the literature...

Agenda

- An invitation (1. sign problem...)
- An invitation (2. analytic continuation from multi-point Padé)
- The sign problem as an inverse problem ...
- ... and something more on other inverse problems ...
- WORK IN PROGRESS!



Suppose you know the values of a function (and of its derivatives) at a number of points

...,
$$f(z_k), f'(z_k), \ldots, f^{(s-1)}(z_k), \ldots \qquad k = 1 \ldots N$$

If you want to approximate the function with a rational function

$$R_n^m(z) = \frac{P_m(z)}{\tilde{Q}_n(z)} = \frac{P_m(z)}{1 + Q_n(z)} = \frac{\sum_{i=0}^m a_i z^i}{1 + \sum_{i=1}^n b_i z^i}$$

the obvious requirement is that

$$R_n^{m(j)}(z_k) = f^{(j)}(z_k)$$
 $k = 1...N, j = 0...s - 1$

Suppose you know the values of a function (and of its derivatives) at a number of points

...,
$$f(z_k), f'(z_k), \ldots, f^{(s-1)}(z_k), \ldots$$
 $k = 1 \ldots N$

If you want to approximate the function with a rational function

$$R_n^m(z) = \frac{P_m(z)}{\tilde{Q}_n(z)} = \frac{P_m(z)}{1 + Q_n(z)} = \frac{\sum_{i=0}^m a_i z^i}{1 + \sum_{i=1}^n b_i z^i}$$

the obvious requirement is that

$$R_n^{m(j)}(z_k) = f^{(j)}(z_k)$$
 $k = 1...N, j = 0...s - 1$

This is the starting point for a *multi-point Padè approximation*: solve the linear system

• •

$$P_m(z_k) - f(z_k)Q_n(z_k) = f(z_k)$$

$$P'_m(z_k) - f'(z_k)Q_n(z_k) - f(z_k)Q'_n(z_k) = f'(z_k)$$

• •

from which we want to get the unknown

$$\{a_i \mid i = 0 \dots m\} \quad \{b_i \mid j = 1 \dots n\} \quad n+m+1=Ns$$

So, we want to approximate a function with a rational function

$$R_n^m(z) = \frac{P_m(z)}{\tilde{Q}_n(z)} = \frac{P_m(z)}{1 + Q_n(z)} = \frac{\sum_{i=0}^m a_i z^i}{1 + \sum_{j=1}^n b_j z^j}$$

Any useful ...?

Yes! LATTICE QCD at IMAGINARY values of the baryonic chemical potential

PHYSICAL REVIEW D 105, 034513 (2022)

Contribution to understanding the phase structure of strong interaction matter: Lee-Yang edge singularities from lattice QCD

P. Dimopoulos, ¹ L. Dini, ² F. Di Renzo, ¹ J. Goswami, ² G. Nicotra, ² C. Schmidt, ² S. Singh, ^{1,*} K. Zambello, ¹ and F. Ziesch, ²

... where we computed and "multi-point Padè approximated"

$$\chi_n^B(T, V, \mu_B) = \left(\frac{\partial}{\partial \hat{\mu}_B}\right)^n \frac{\ln Z(T, V, \mu_l, \mu_s)}{VT^3}$$
$$= \left(\frac{1}{3} \frac{\partial}{\partial \hat{\mu}_l} + \frac{1}{3} \frac{\partial}{\partial \hat{\mu}_s}\right)^n \frac{\ln Z(T, V, \mu_l, \mu_s)}{VT^3}$$

So, we want to approximate a function with a rational function

$$R_n^m(z) = \frac{P_m(z)}{\tilde{Q}_n(z)} = \frac{P_m(z)}{1 + Q_n(z)} = \frac{\sum_{i=0}^m a_i z^i}{1 + \sum_{j=1}^n b_j z^j}$$

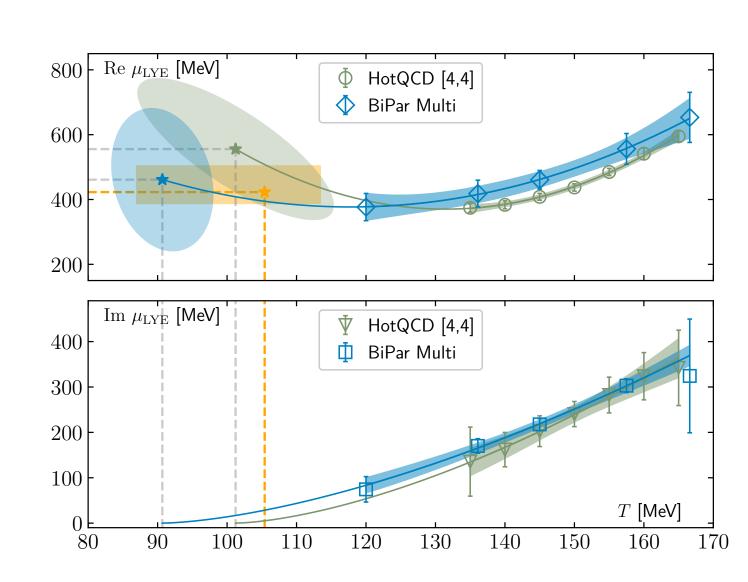


FIG. 4. Scaling fits for the LYE singularities related to the CEP. Green data come from a [4,4] Padé from Ref. [7]. Blue data come from the multi-point Padé. *Top*: Scaling of the real part. *Bottom*: Scaling of the imaginary part. The ellipses shown in the top panel represent the 68% confidence region deduced from the covariance matrix of the fit. The orange box indicates the AIC weighted estimate (6).

Any useful ...?

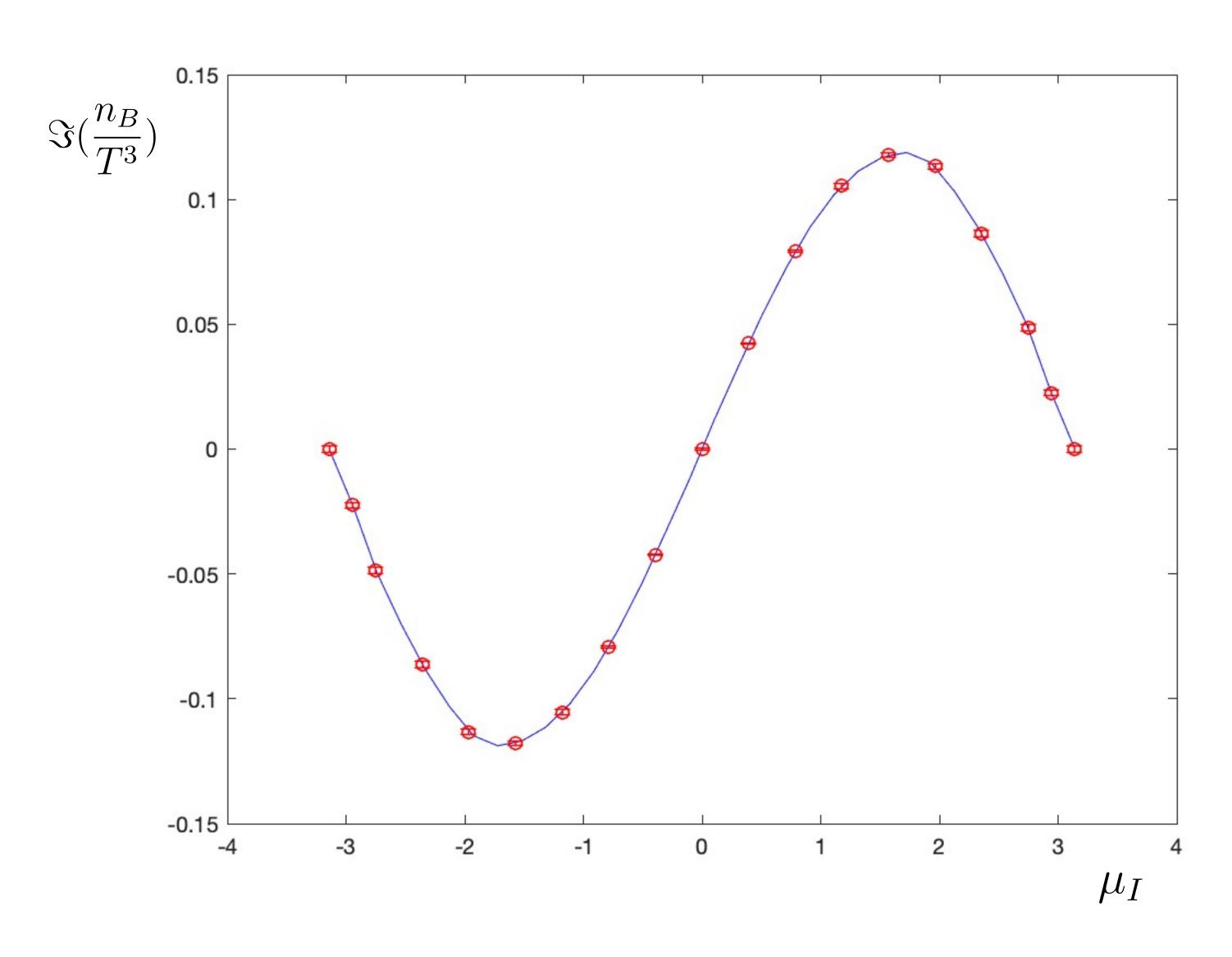
Yes! LATTICE QCD at IMAGINARY values of the baryonic chemical potential

... a natural analytic continuation to real chemical potential!

... and not only that: singularities!

In the following we will play with a few Bielefeld Parma Collaboration data 2+1 HISQ at physical quarks mass, at fixed cutoff ($N_{\tau}=6$)

$$T = 157.5 \ (\sim 155) \ \mathrm{MeV}$$

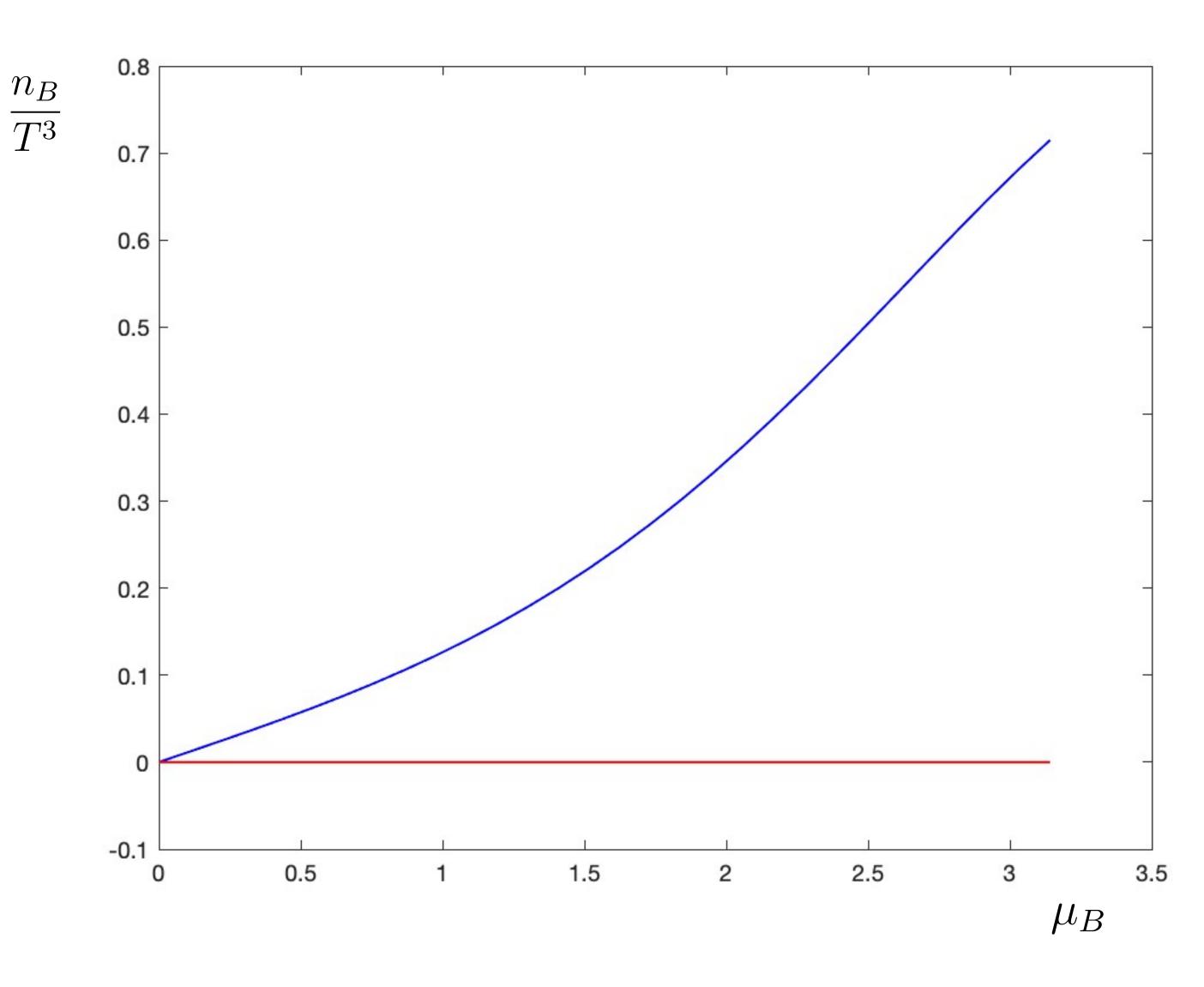


... which is pretty simple (we will be concerned with the number density):

you take your rational function, which describes very well data at IMAGINARY VALUES of μ_B

CAVEAT: errors on data points are there ... no error shown on the interpolating function (as for now ...)

$$T = 157.5 \ (\sim 155) \ \mathrm{MeV}$$



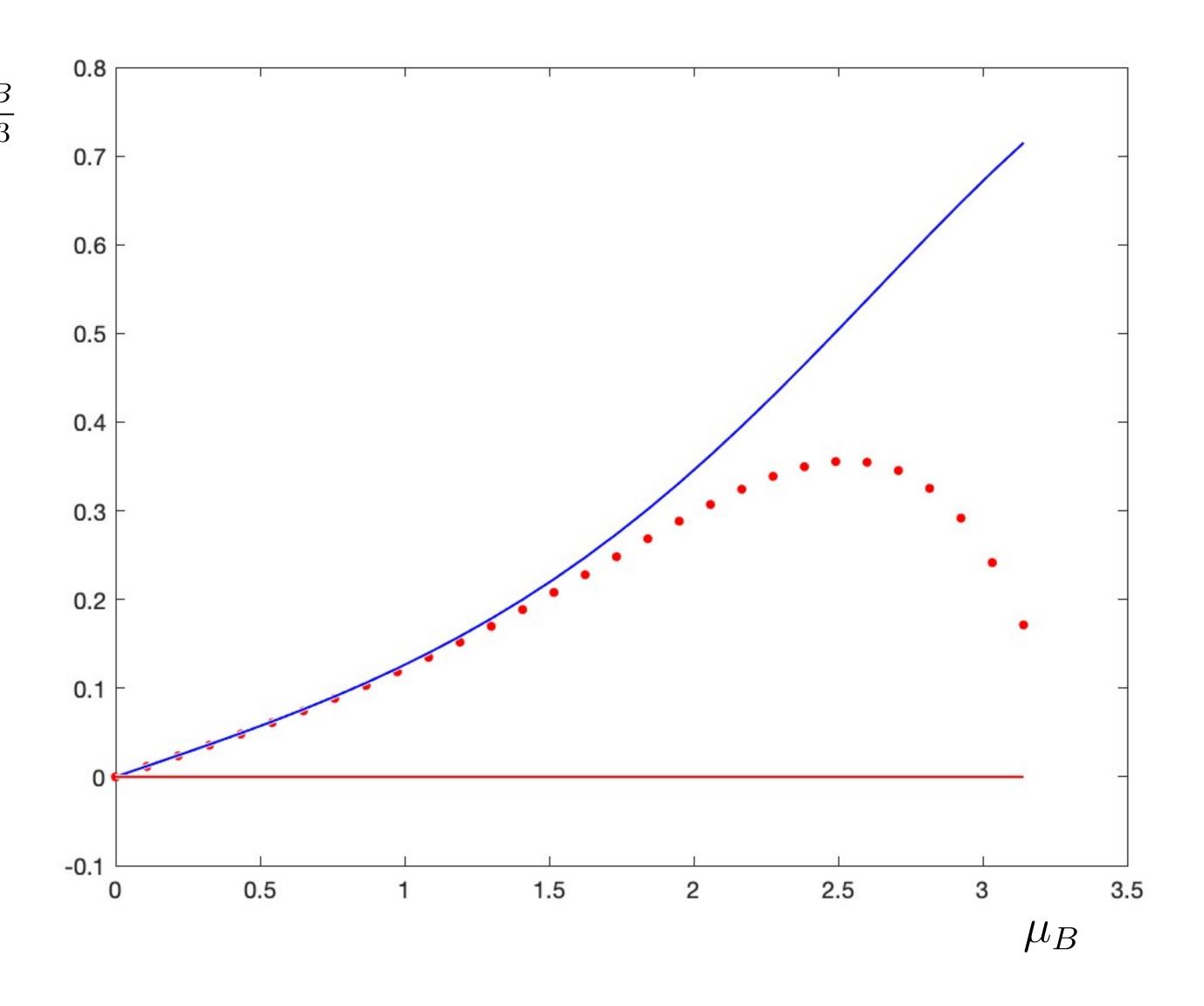
... which is pretty simple (we will be concerned with the number density):

you take your rational function, which describes very well data at IMAGINARY VALUES of μ_B

... and you simply compute it for REAL VALUES of μ_B

CAVEAT: no error shown as for now ... here we are concerned with *trends*... FIXED CUTOFF!

$$T = 157.5 \ (\sim 155) \ \mathrm{MeV}$$



... which is pretty simple (we will be concerned with the number density):

you take your rational function, which describes very well data at IMAGINARY VALUES of μ_B

... and you simply compute it for REAL VALUES of μ_B

You can compare the result with HotQCD results

PHYSICAL REVIEW D **105**, 074511 (2022)

Taylor expansions and Padé approximants for cumulants of conserved charge fluctuations at nonvanishing chemical potentials

D. Bollweg, ¹ J. Goswami, ² O. Kaczmarek, ² F. Karsch, ² Swagato Mukherjee, ³ P. Petreczky, ³ C. Schmidt, ² and P. Scior, ³ (HotQCD Collaboration)

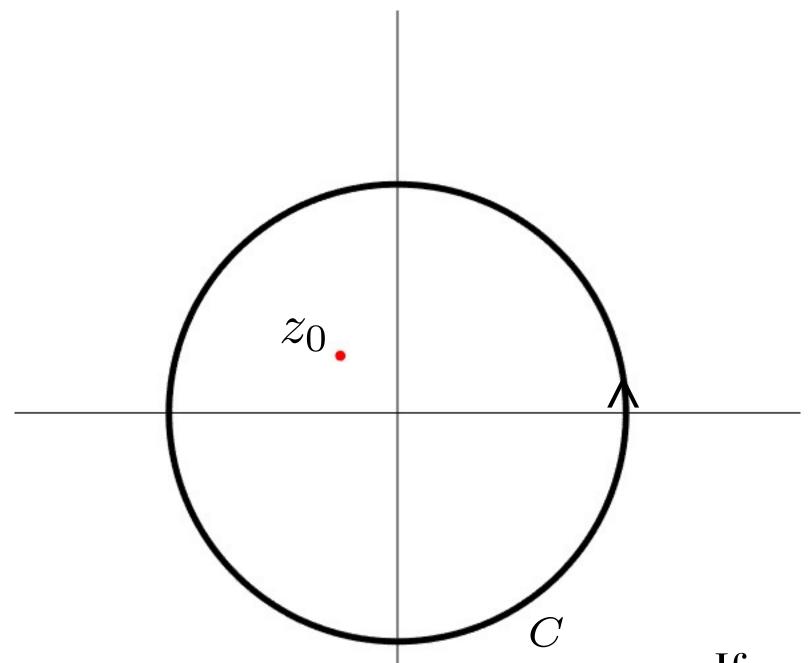
CAVEAT: no error shown as for now ... here we are concerned with *trends*... FIXED CUTOFF!

Finite density QCD as an inverse problem

Finite density QCD as an inverse problem

... aka How to trade a difficult problem for another (even more?) difficult one ...

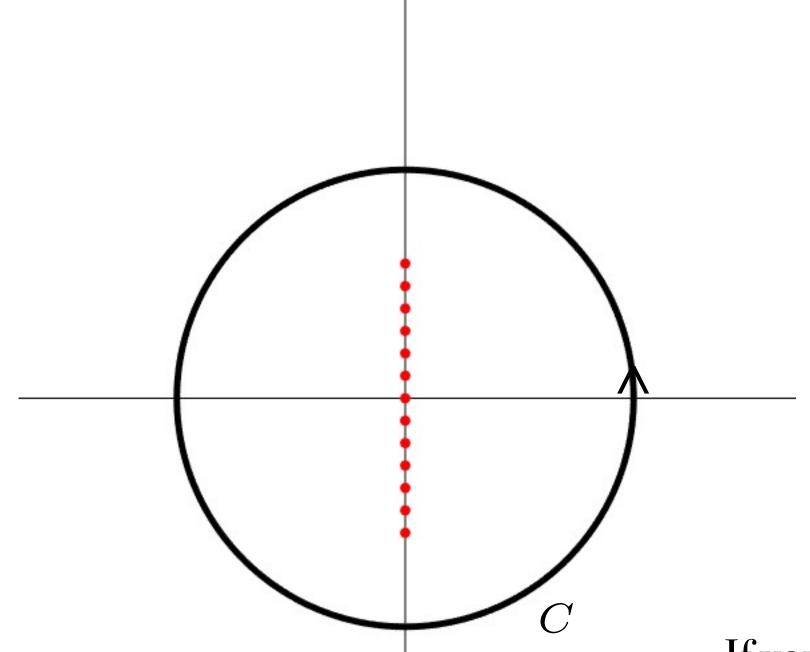
One simple way of thinking of it is that you can perfectly know such functions from an apparently limited amount of information.



CAUCHY FORMULA

$$f(z_0) = \frac{1}{2\pi i} \oint_C \frac{f(z)}{z - z_0} dz$$

If you know the function on the contour, you can compute it at any point inside... sounds good!

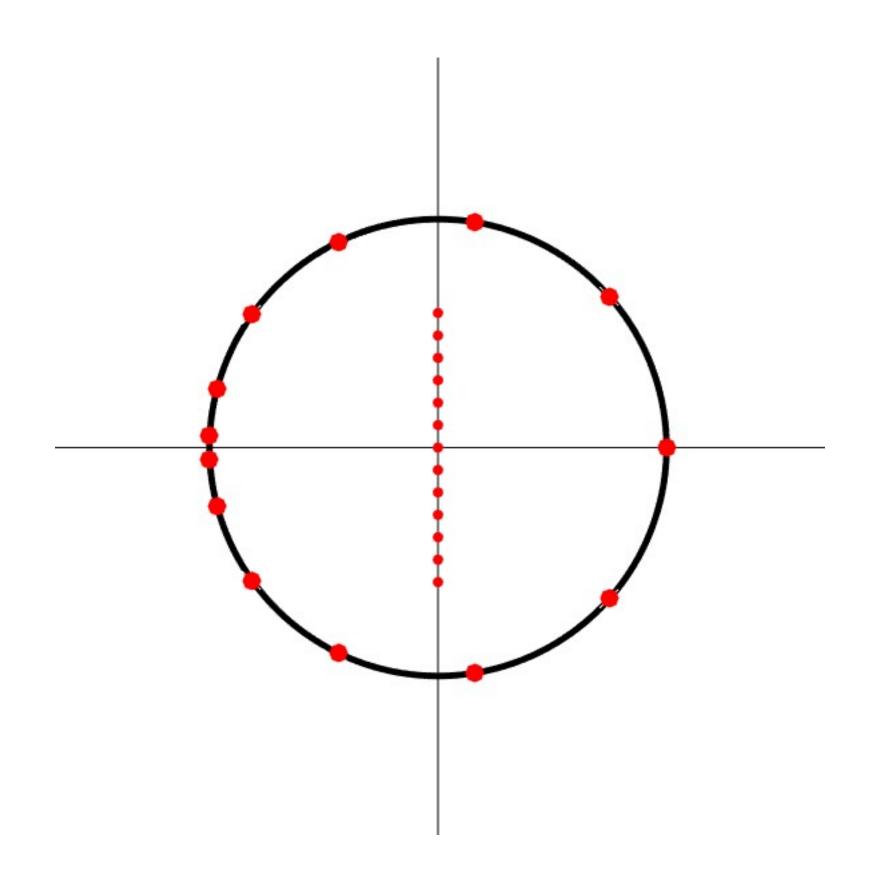


CAUCHY FORMULA

$$f(z_0) = \frac{1}{2\pi i} \oint_C \frac{f(z)}{z - z_0} dz$$

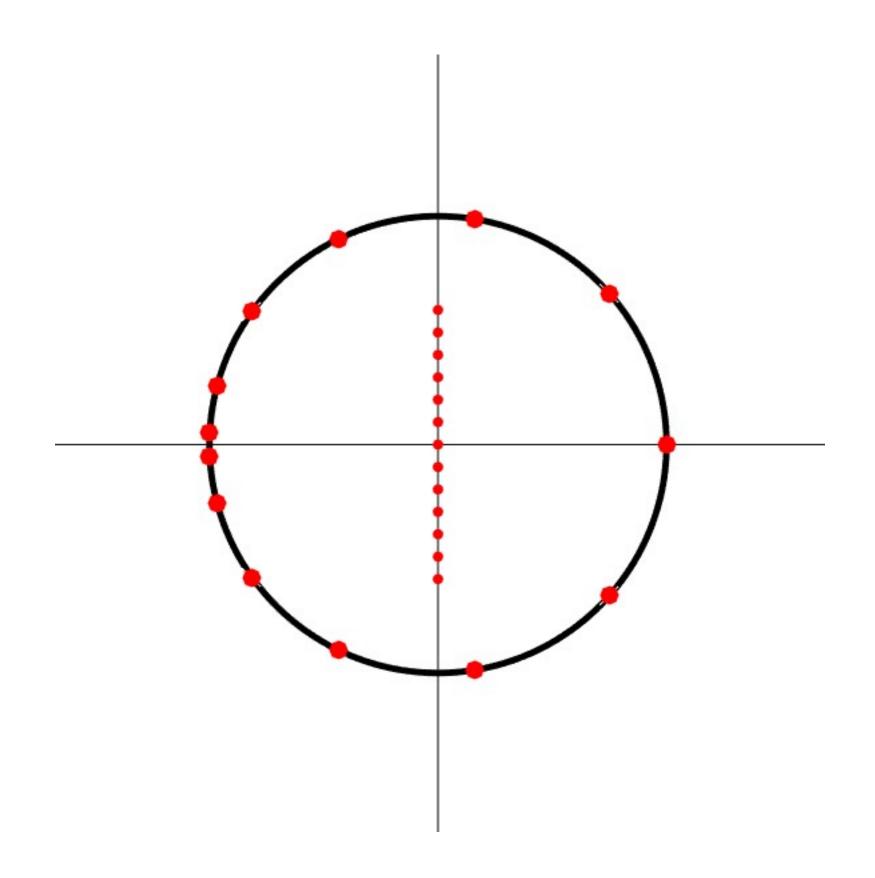
If you know the function on the contour, you can compute it at any point inside... sounds good!

... at any point, including the (only) ones we can compute (on the imaginary axis) in our case...



$$f(z_0) = \frac{1}{2\pi} \int_0^{2\pi} \frac{f(Re^{i\theta}) Re^{i\theta}}{Re^{i\theta} - z_0} d\theta \simeq \frac{1}{2\pi} \sum_{k=1}^n w_k \frac{f(Re^{i\theta_k}) Re^{i\theta_k}}{Re^{i\theta_k} - z_0}$$

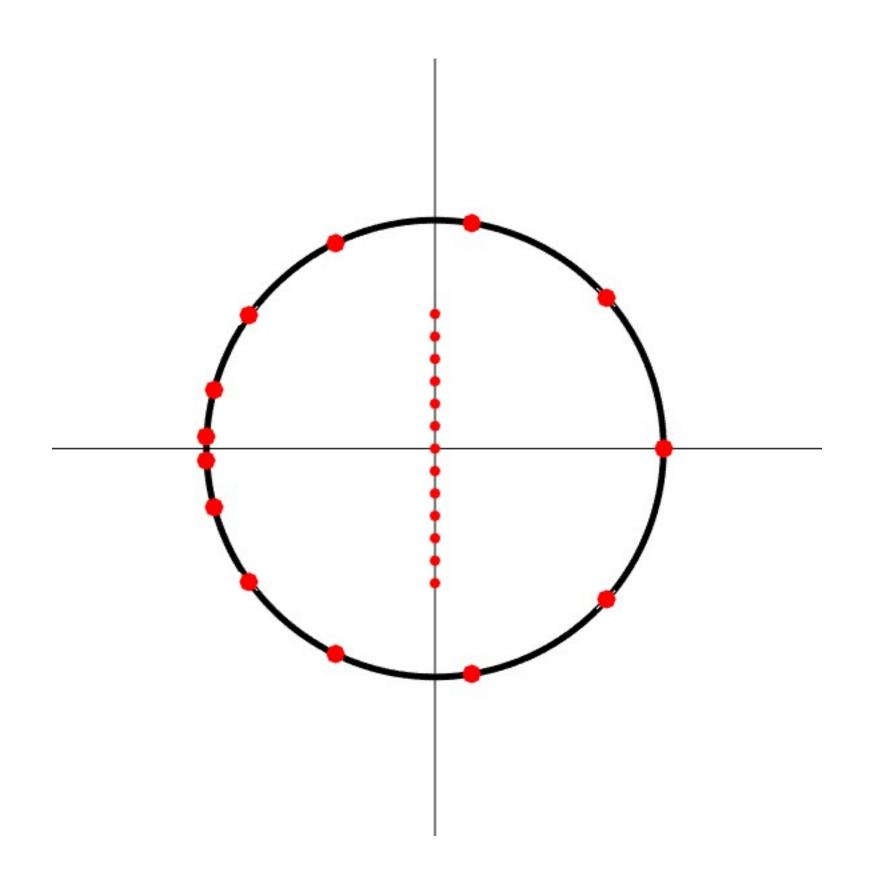
$$y_i = \frac{1}{2\pi} \sum_{k=1}^{n} w_k \frac{R e^{i\theta_k}}{R e^{i\theta_k} - z_i} \hat{f}_k, \ i = 1, 2, \dots, n$$



$$f(z_0) = \frac{1}{2\pi} \int_0^{2\pi} \frac{f(Re^{i\theta}) Re^{i\theta}}{Re^{i\theta} - z_0} d\theta \simeq \frac{1}{2\pi} \sum_{k=1}^n w_k \frac{f(Re^{i\theta_k}) Re^{i\theta_k}}{Re^{i\theta_k} - z_0}$$

$$y_i = \frac{1}{2\pi} \sum_{k=1}^{n} w_k \frac{R e^{i\theta_k}}{R e^{i\theta_k} - z_i} \hat{f}_k, \ i = 1, 2, \dots, n$$

... and then you are ready for your (BRAVE) INVERSE PROBLEM!



$$A_{ik} = \frac{1}{2\pi} w_k \frac{R e^{i\theta_k}}{R e^{i\theta_k} - z_i}$$

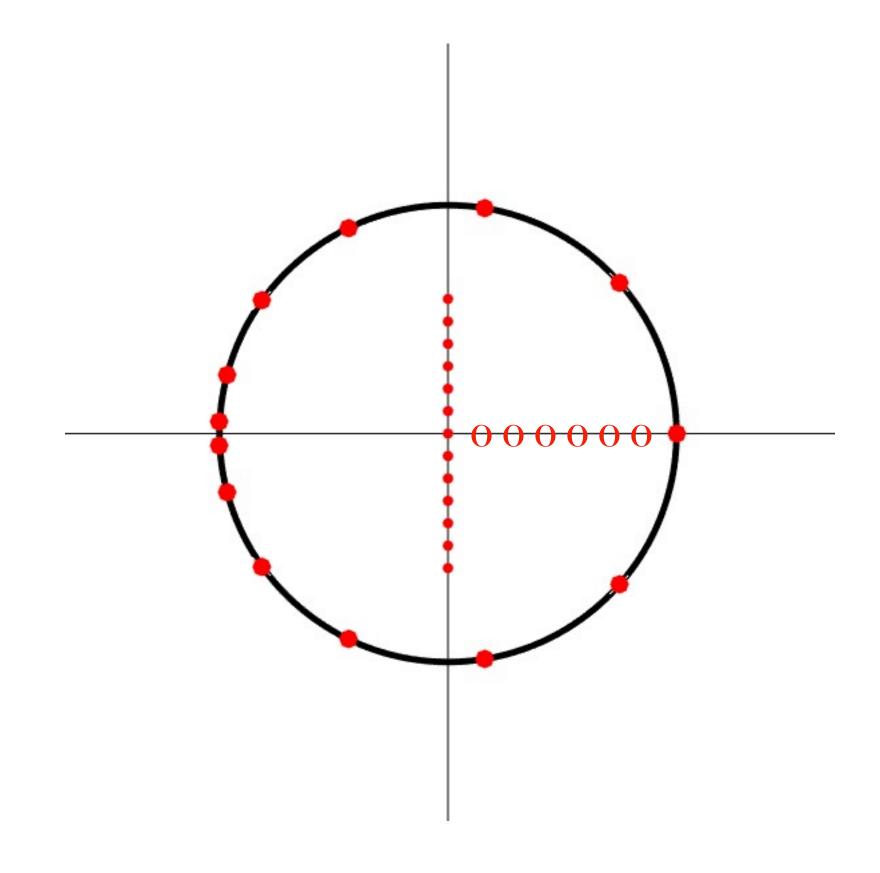
$$f(z_0) = \frac{1}{2\pi} \int_0^{2\pi} \frac{f(Re^{i\theta}) Re^{i\theta}}{Re^{i\theta} - z_0} d\theta \simeq \frac{1}{2\pi} \sum_{k=1}^n w_k \frac{f(Re^{i\theta_k}) Re^{i\theta_k}}{Re^{i\theta_k} - z_0}$$

$$y_i = \frac{1}{2\pi} \sum_{k=1}^{n} w_k \frac{R e^{i\theta_k}}{R e^{i\theta_k} - z_i} \hat{f}_k, \ i = 1, 2, \dots, n$$

... and then you are ready for your (BRAVE) INVERSE PROBLEM!

$$A\mathbf{x} = \mathbf{b}$$

SOLVE for the \hat{f}_k !



$$f(z_0) = \frac{1}{2\pi} \int_0^{2\pi} \frac{f(Re^{i\theta}) Re^{i\theta}}{Re^{i\theta} - z_0} d\theta \simeq \frac{1}{2\pi} \sum_{k=1}^n w_k \frac{f(Re^{i\theta_k}) Re^{i\theta_k}}{Re^{i\theta_k} - z_0}$$

$$y_i = \frac{1}{2\pi} \sum_{k=1}^{n} w_k \frac{R e^{i\theta_k}}{R e^{i\theta_k} - z_i} \hat{f}_k, \ i = 1, 2, \dots, n$$

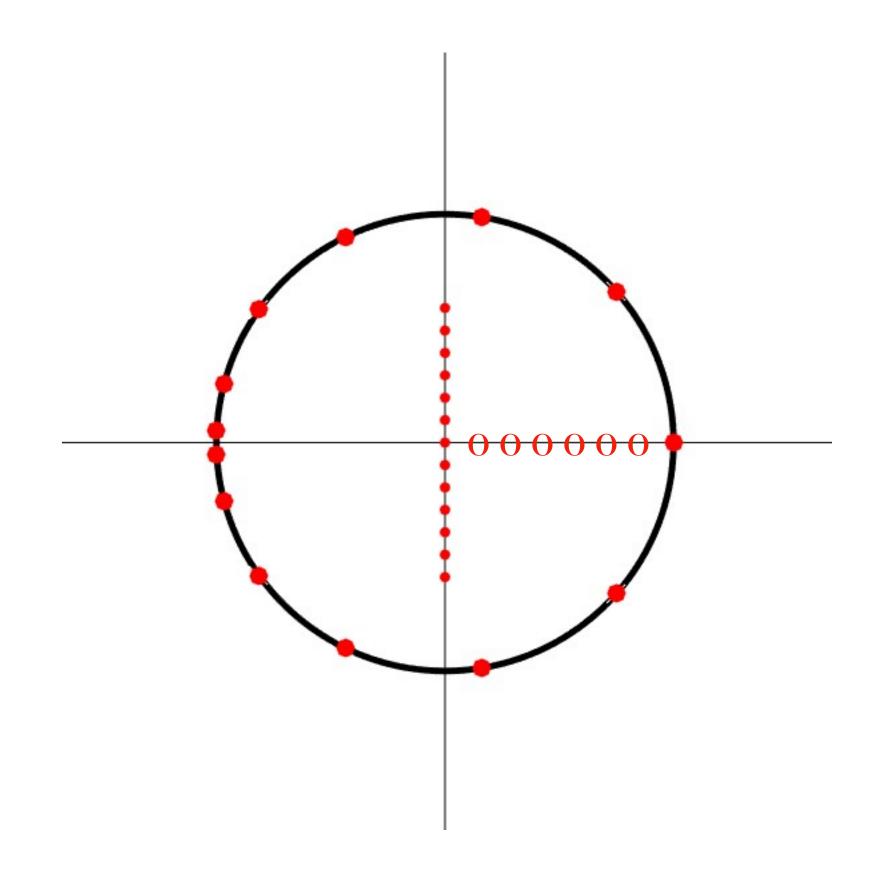
... and when you have solved your (BRAVE) INVERSE PROBLEM(!) ...

$$A_{ik} = \frac{1}{2\pi} w_k \frac{R e^{i\theta_k}}{R e^{i\theta_k} - z_i}$$

$$A\mathbf{x} = \mathbf{b}$$

SOLVE for the \hat{f}_k !

Once you have the values on the contour, you can compute on the REAL AXIS! (ooooo)



$$A_{ik} = \frac{1}{2\pi} w_k \frac{R e^{i\theta_k}}{R e^{i\theta_k} - z_i}$$

With your favourite QUADRATURE method ... you can go numeric!

De facto, you would like to think of Legendre quadrature

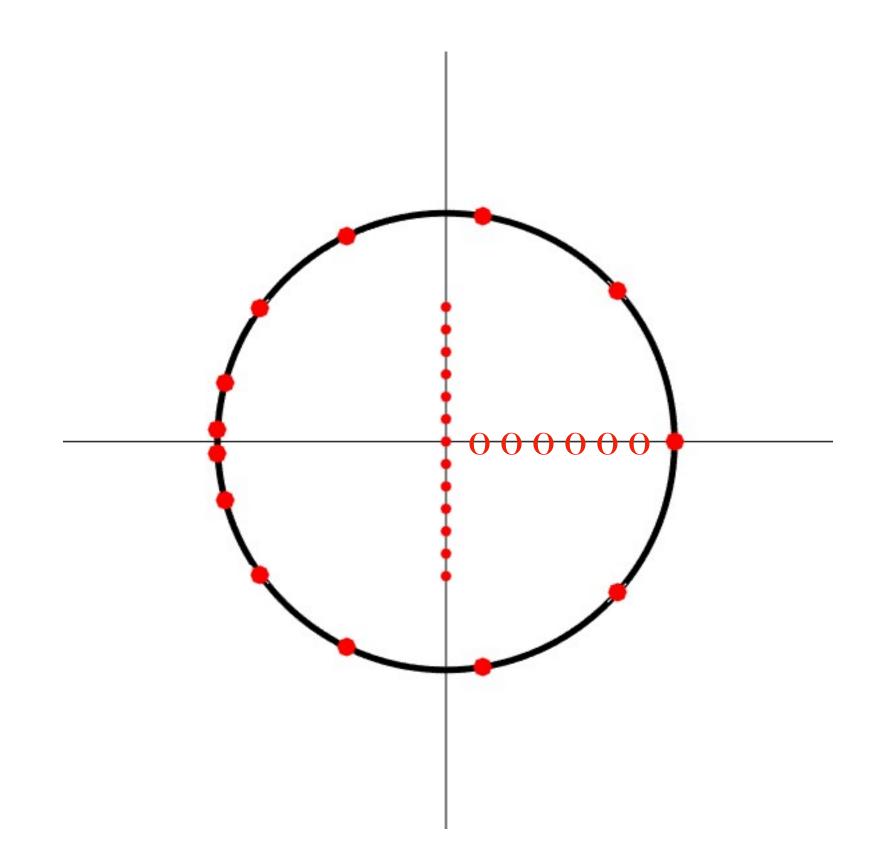
$$f(z_0) = \frac{1}{2\pi} \int_0^{2\pi} \frac{f(Re^{i\theta}) Re^{i\theta}}{Re^{i\theta} - z_0} d\theta \simeq \frac{1}{2\pi} \sum_{k=1}^n w_k \frac{f(Re^{i\theta_k}) Re^{i\theta_k}}{Re^{i\theta_k} - z_0}$$

$$y_i = \frac{1}{2\pi} \sum_{k=1}^{n} w_k \frac{R e^{i\theta_k}}{R e^{i\theta_k} - z_i} \hat{f}_k, \ i = 1, 2, \dots, n$$

... a BRAVE INVERSE PROBLEM!...

$$A\mathbf{x} = \mathbf{b}$$

SOLVE for the \hat{f}_k !



$$f(z_0) = \frac{1}{2\pi} \int_0^{2\pi} \frac{f(Re^{i\theta}) Re^{i\theta}}{Re^{i\theta} - z_0} d\theta \simeq \frac{1}{2\pi} \sum_{k=1}^n w_k \frac{f(Re^{i\theta_k}) Re^{i\theta_k}}{Re^{i\theta_k} - z_0}$$

$$y_i = \frac{1}{2\pi} \sum_{k=1}^{n} w_k \frac{R e^{i\theta_k}}{R e^{i\theta_k} - z_i} \hat{f}_k, \ i = 1, 2, \dots, n$$

$$A\mathbf{x} = \mathbf{b}$$

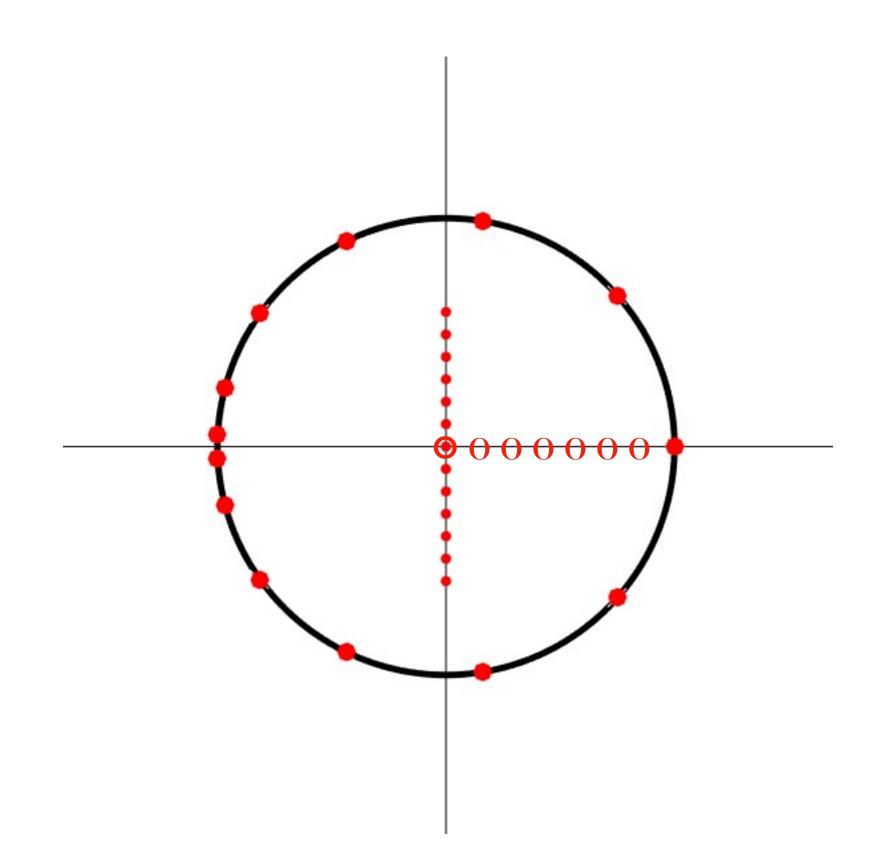
SOLVE for the \hat{f}_k !

Can we hope it could work? ... are we afraid it should not?

... and it could(?)/should(?) not for a combination of

- (a) bad condition number of the linear system
- (b) the quadrature formula being NOT exact

Much care is needed ... and so we will perform some tests...



$$f(z_0) = \frac{1}{2\pi} \int_0^{2\pi} \frac{f(Re^{i\theta}) Re^{i\theta}}{Re^{i\theta} - z_0} d\theta \simeq \frac{1}{2\pi} \sum_{k=1}^n w_k \frac{f(Re^{i\theta_k}) Re^{i\theta_k}}{Re^{i\theta_k} - z_0}$$

$$y_i = \frac{1}{2\pi} \sum_{k=1}^{n} w_k \frac{R e^{i\theta_k}}{R e^{i\theta_k} - z_i} \hat{f}_k, \ i = 1, 2, \dots, n$$

$$A\mathbf{x} = \mathbf{b}$$

SOLVE for the \hat{f}_k !

Not the end of the story ... Obviously, it is a good (CHEAP) idea to also remember Cauchy formula for derivatives

$$f^{(n)}(z_0) = \frac{n!}{2\pi i} \oint_C \frac{f(z)}{(z-z_0)^{n+1}} dz = \frac{n!}{2\pi i} \int_0^{2\pi} \frac{f(R\exp(i\theta)) R\exp(i\theta)}{(R\exp i\theta - z_0)^{n+1}} d\theta$$

LOOK! ... we tackled a funny inverse problem, a non-standard one ... other are very much investigated!

LOOK! ... we tackled a funny inverse problem, a non-standard one ... other are very much investigated!

SOMETHING ELSE you can do with the inverse problem machinery: we can play the same game for inverse Laplace transform ...

You know this ...
$$f(s) = \int_0^\infty e^{-ts} \, F(t) \, dt$$
 You want this ...

LOOK! ... we tackled a funny inverse problem, a non-standard one ... other are very much investigated!

SOMETHING ELSE you can do with the inverse problem machinery: we can play the same game for inverse Laplace transform ...

You know this ...

$$f(s) = \int_0^\infty e^{-ts} F(t) dt$$
You want this ...

We slightly rephrase the problem ...

$$f(s) = \int_0^\infty e^{-t} e^{-t(s-1)} F(t) dt$$

... and we can play the same game ...

$$f(s) = \int_0^\infty e^{-t} e^{-t(s-1)} F(t) dt \sim \sum_j w_j e^{-t_j(s-1)} F(t_j)$$

This time, Laguerre quadratures ...

Can we hope it could work? ... are we afraid it should not?

... A brave attempt, but with a rational ...

... and it could(?)/should(?) not for a combination of

- (a) bad condition number of the linear system
- (b) the quadrature formula being NOT exact

... A brave attempt, but with a rational ...

Can we hope it could work? ... are we afraid it should not?

... and it could(?)/should(?) not for a combination of

- (a) bad condition number of the linear system
- (b) the quadrature formula being NOT exact

You know this ...

know this ...
$$f(s_j) = \int_0^\infty e^{-ts} F(t) dt$$
 You want this ...

HARD ... from a limited amount of information you want to get an infinite one...

Can we hope it could work? ... are we afraid it should not?

... and it could(?)/should(?) not for a combination of

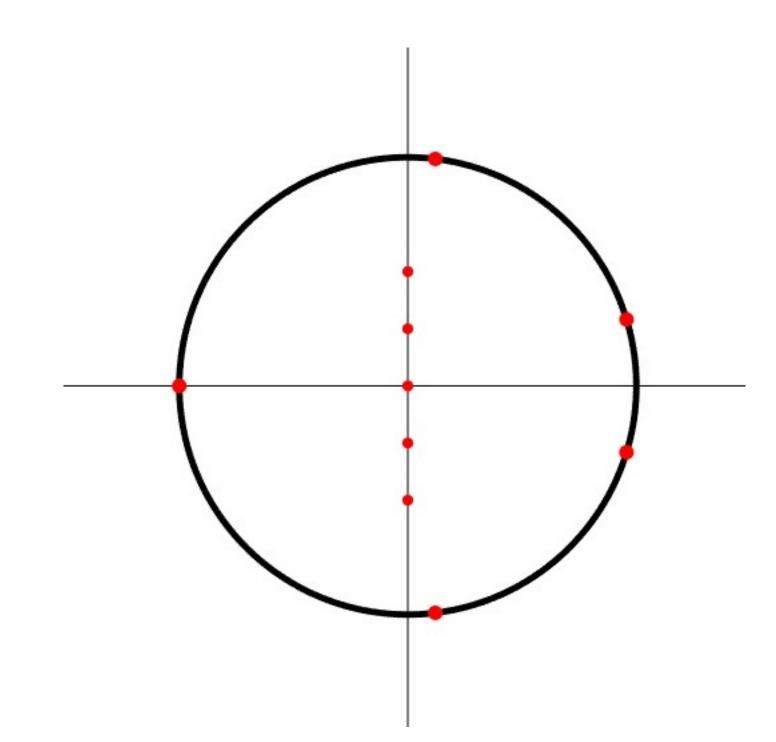
- (a) bad condition number of the linear system
- (b) the quadrature formula being NOT exact

You know this ...

know this ...
$$f(s_j) = \int_0^\infty e^{-ts} F(t) dt$$
 You want this ...

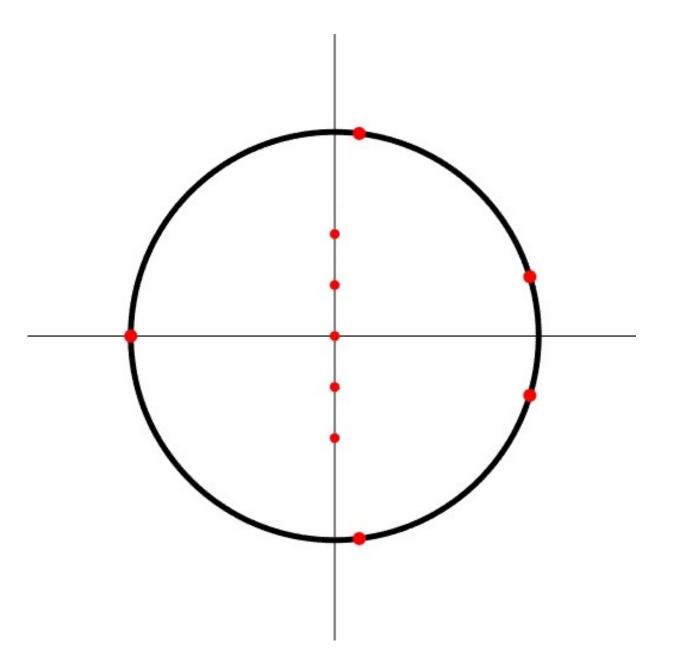
HARD ... from a limited amount of information you want to get an infinite one...

Actually, both for (anti)Laplace and for the Cauchy formula, we provide n values and try to get n as well ... e.g.



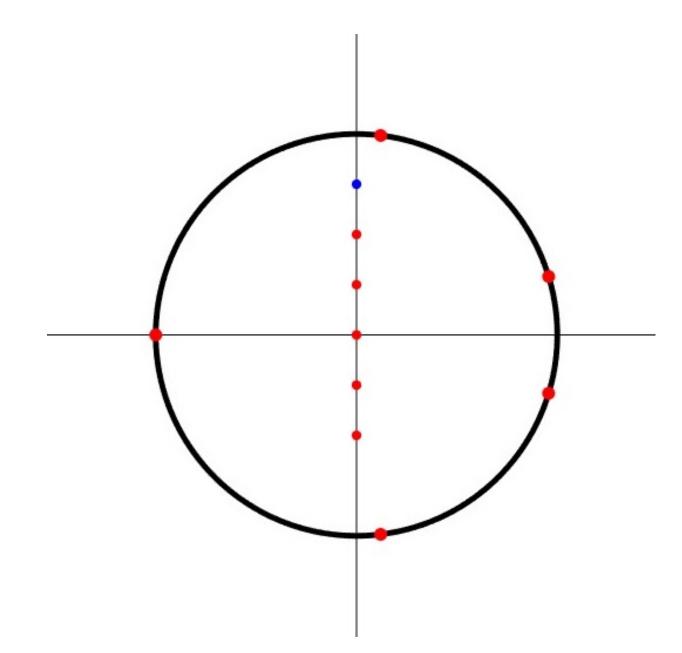
Actually, both for (anti)Laplace and for the Cauchy formula, we provide n values and try to get n as well ... e.g.

... A brave attempt, but with a rational ...

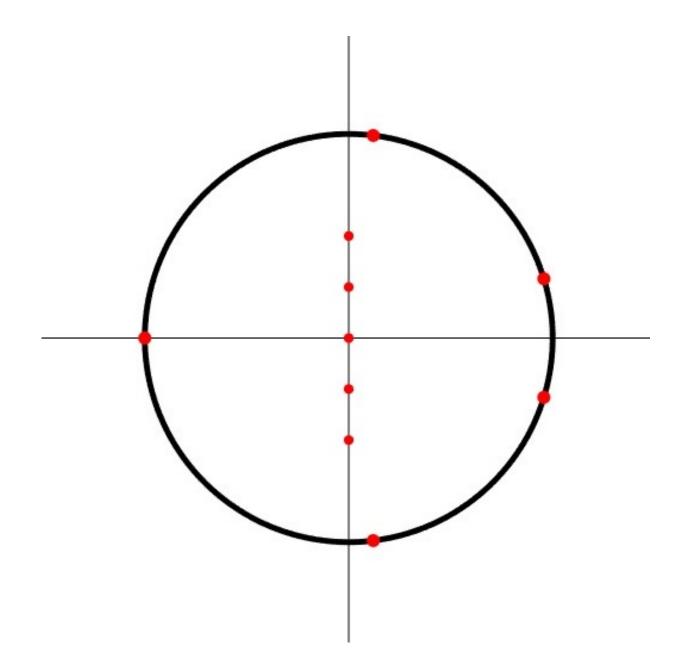


Actually, both for (anti)Laplace and for the Cauchy formula, we provide n values and try to get n as well ... e.g.

... but then ... simply add one (single) point ...



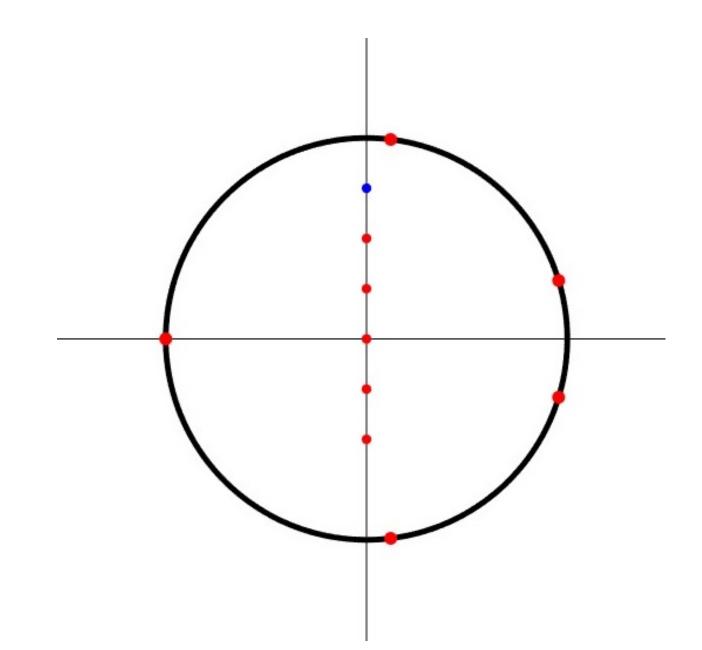
... A brave attempt, but with a rational ...

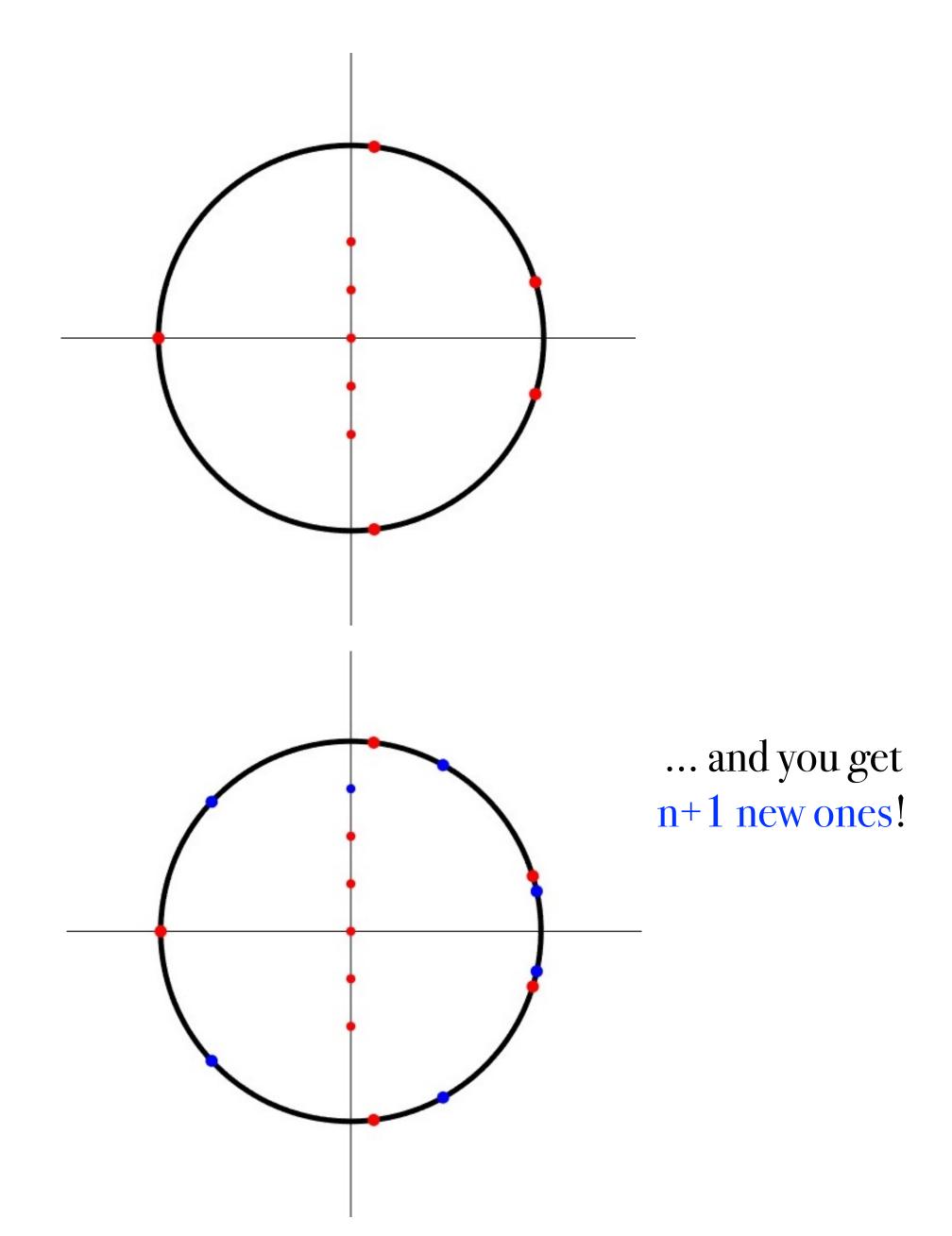


... A brave attempt, but with a rational ...

Actually, both for (anti)Laplace and for the Cauchy formula, we provide n values and try to get n as well ... e.g.

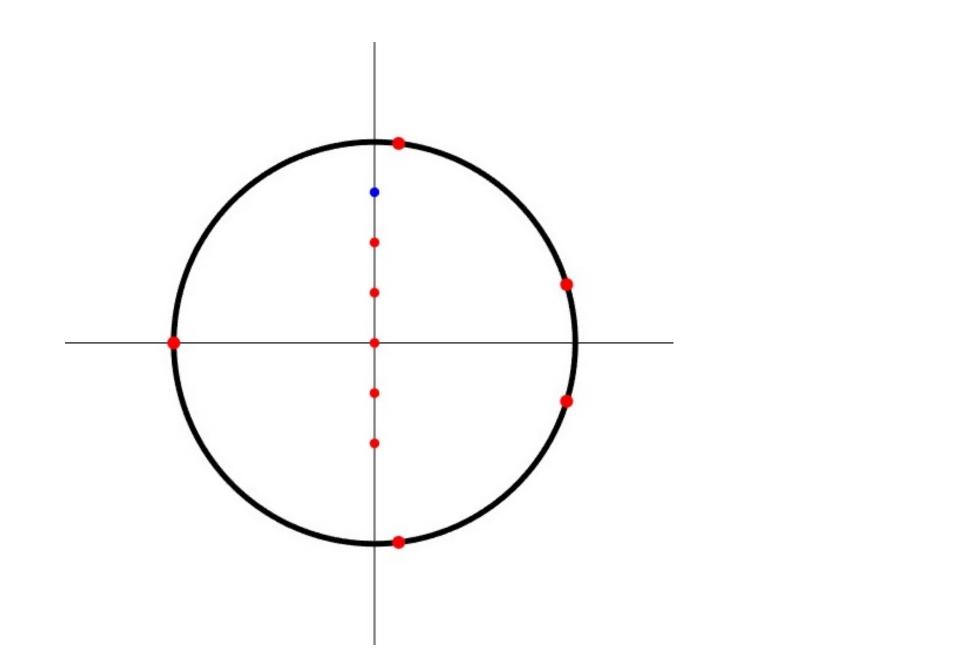
... but then ... simply add one (single) point ...

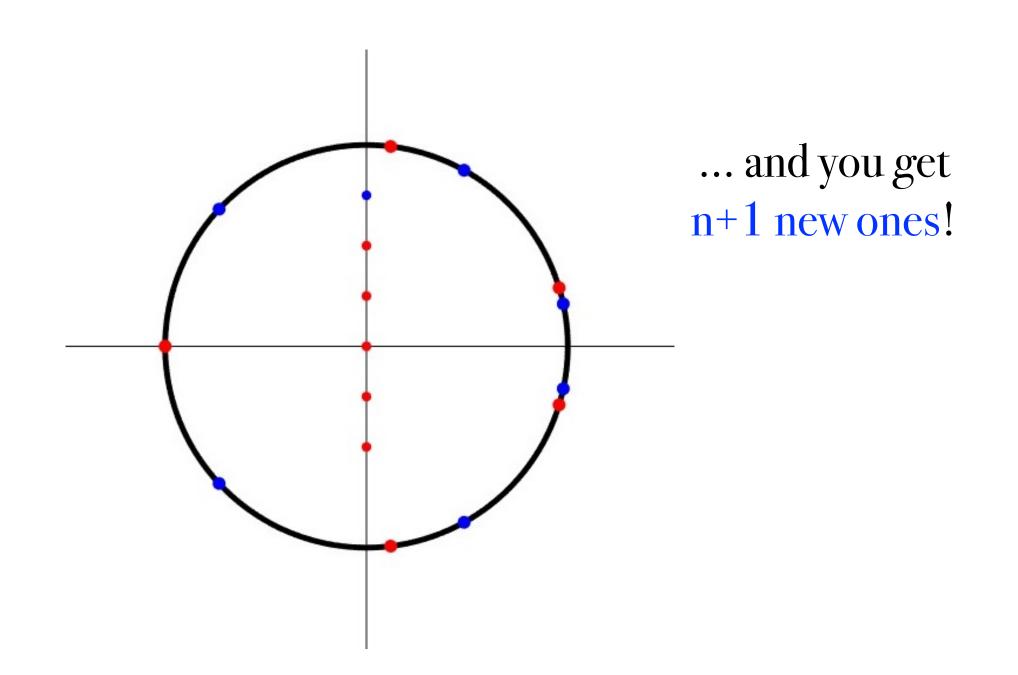




... A brave attempt, but with a rational ...

... simply add one (single) point ...





... extra tricks we can play ... (Laplace) reparametrization!

$$f(s_j) = \int_0^\infty e^{-ts} F(t) dt = t_0 \int_0^\infty e^{-tt_0 s} F(tt_0) dt = t_0 \int_0^\infty e^{-t} e^{-t(t_0 s - 1)} F(tt_0) dt$$

... we can get quite a number of values! ... actually all those for which the machinery works ...

... (Laplace) reparametrization!

... In the end ... DOES IT WORK?!

$$f(s_j) = \int_0^\infty e^{-ts} F(t) dt = t_0 \int_0^\infty e^{-tt_0 s} F(tt_0) dt = t_0 \int_0^\infty e^{-t} e^{-t(t_0 s - 1)} F(tt_0) dt$$

... we can get quite a number of values! ... actually all those for which the machinery works ...

Not only we can possibly get quite a number of values ... we get a CONSISTENCY ARGUMENT!

$$f(s_j) \sim \sum_k e^{-t_k(t_0 s - 1)} F(t_k t_0) dt$$

... (Laplace) reparametrization!

... In the end ... DOES IT WORK?!

$$f(s_j) = \int_0^\infty e^{-ts} F(t) dt = t_0 \int_0^\infty e^{-tt_0 s} F(tt_0) dt = t_0 \int_0^\infty e^{-t} e^{-t(t_0 s - 1)} F(tt_0) dt$$

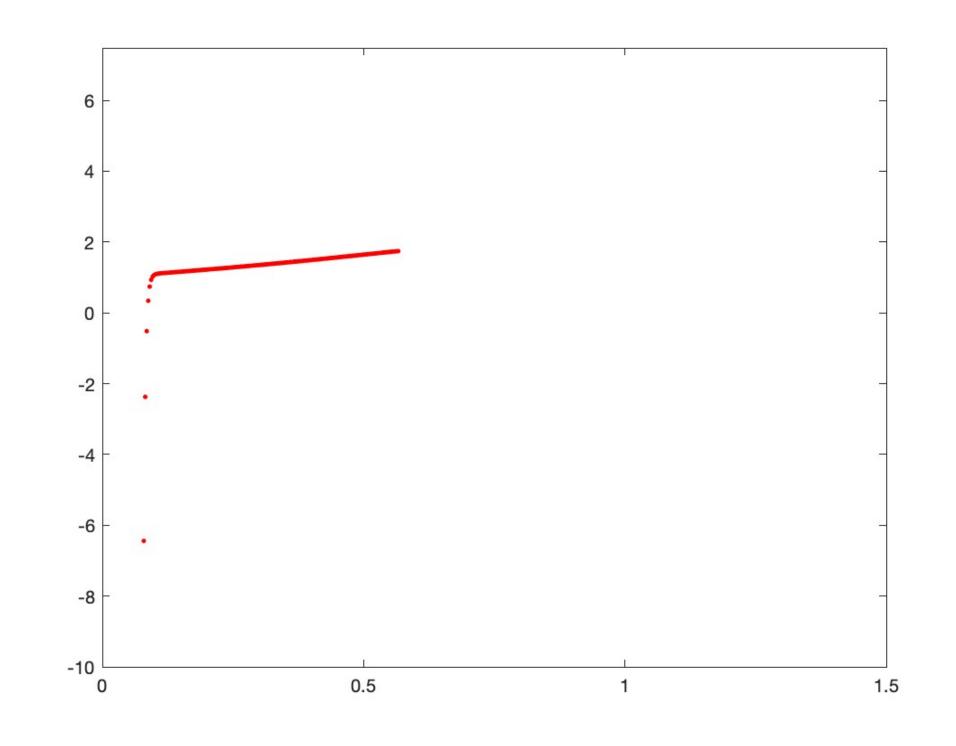
... we can get quite a number of values! ... actually all those for which the machinery works ...

Not only we can possibly get quite a number of values ... we get a CONSISTENCY ARGUMENT!

$$f(s_j) \sim \sum_k e^{-t_k(t_0 s - 1)} F(t_k t_0) dt$$

$$f(s) = \frac{1}{s-1} \longrightarrow F(t) = e^t$$

$$k=2$$

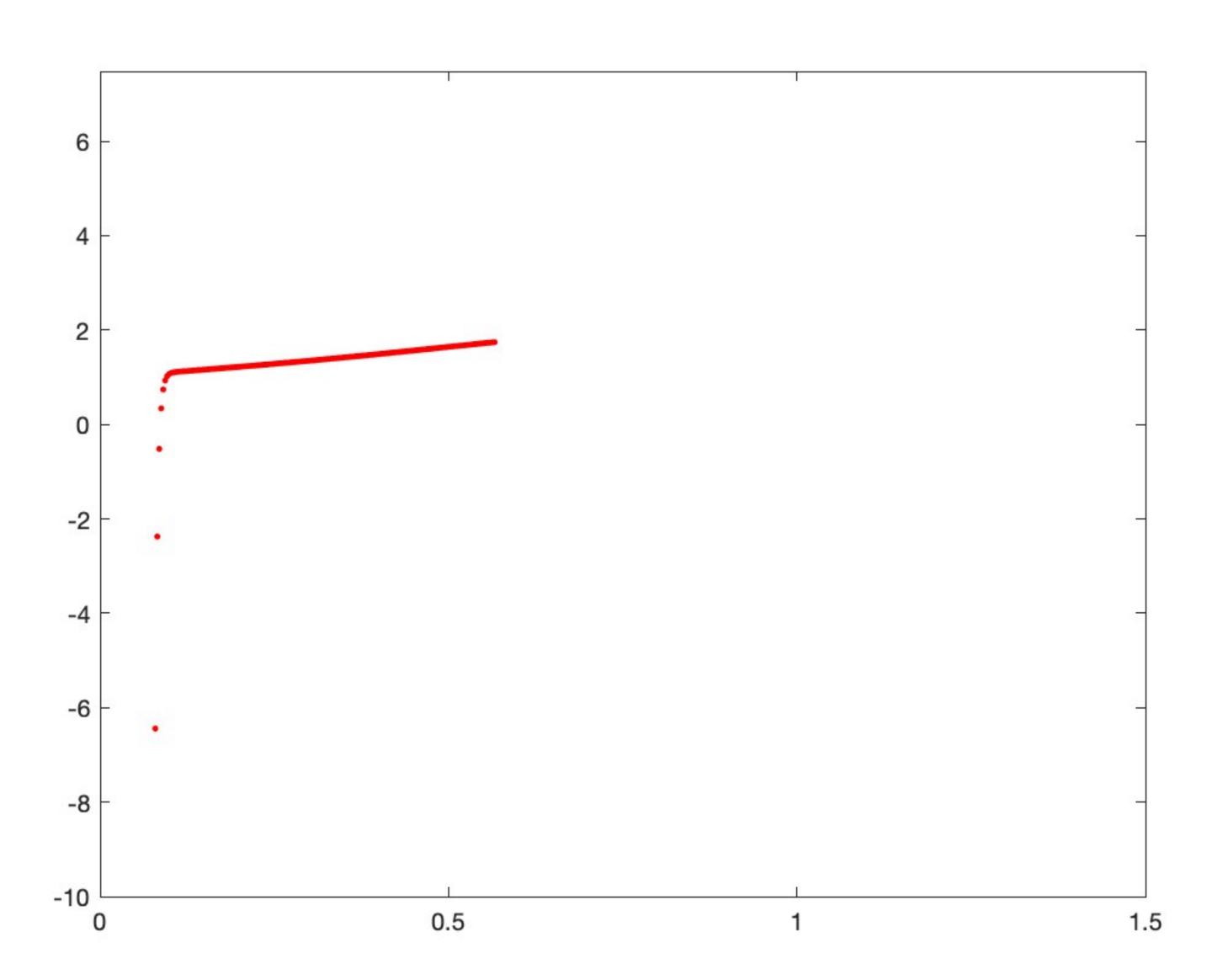


Some CRAZY points
... and ...
a number falling on
a SMOOTH line

$$f(s_j) \sim \sum_k e^{-t_k(t_0 s - 1)} F(t_k t_0) dt$$

$$f(s) = \frac{1}{s - 1} \quad \to \quad F(t) = e^t$$

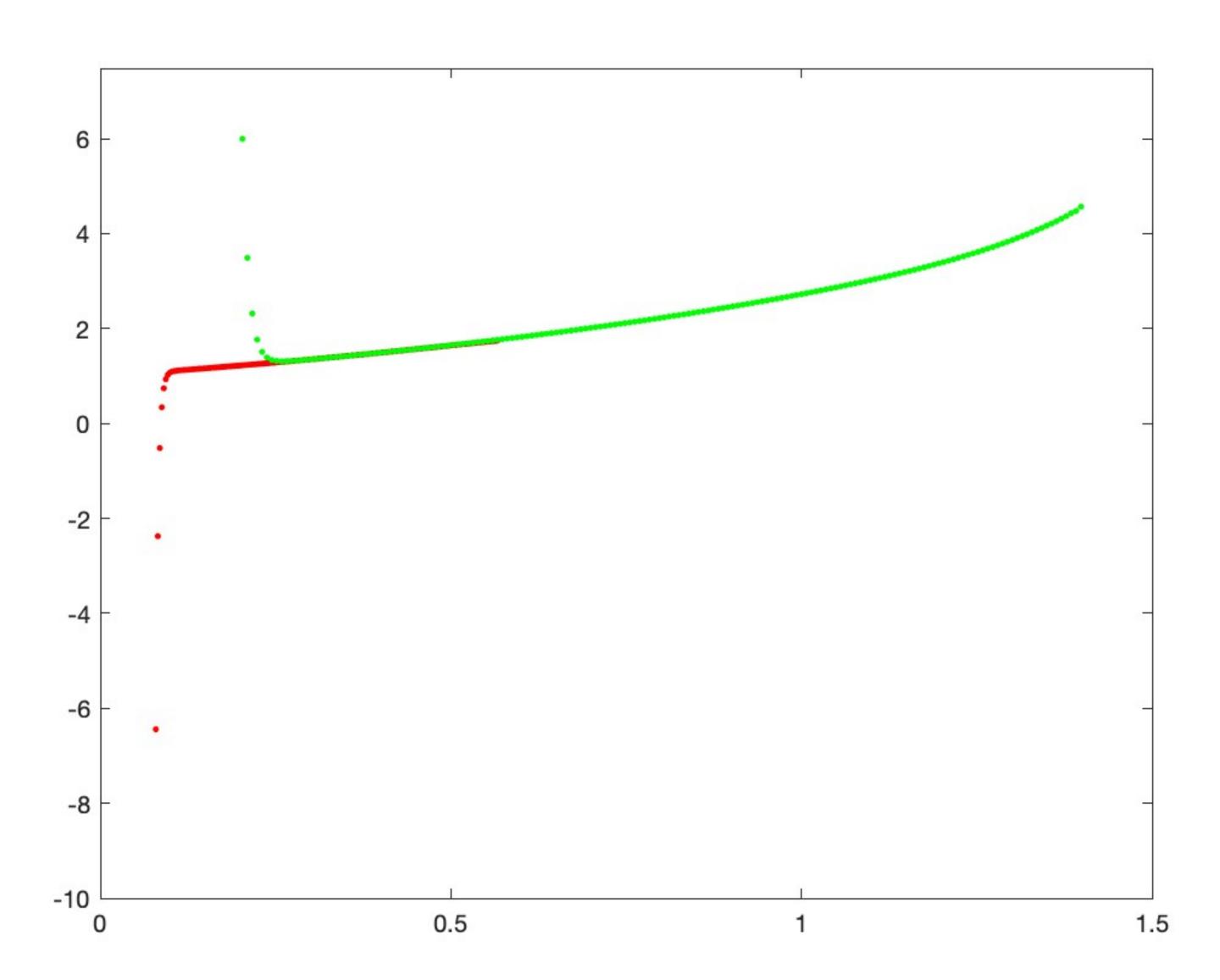
$$k = 2$$



$$f(s_j) \sim \sum_k e^{-t_k(t_0 s - 1)} F(t_k t_0) dt$$

$$f(s) = \frac{1}{s - 1} \longrightarrow F(t) = e^t$$

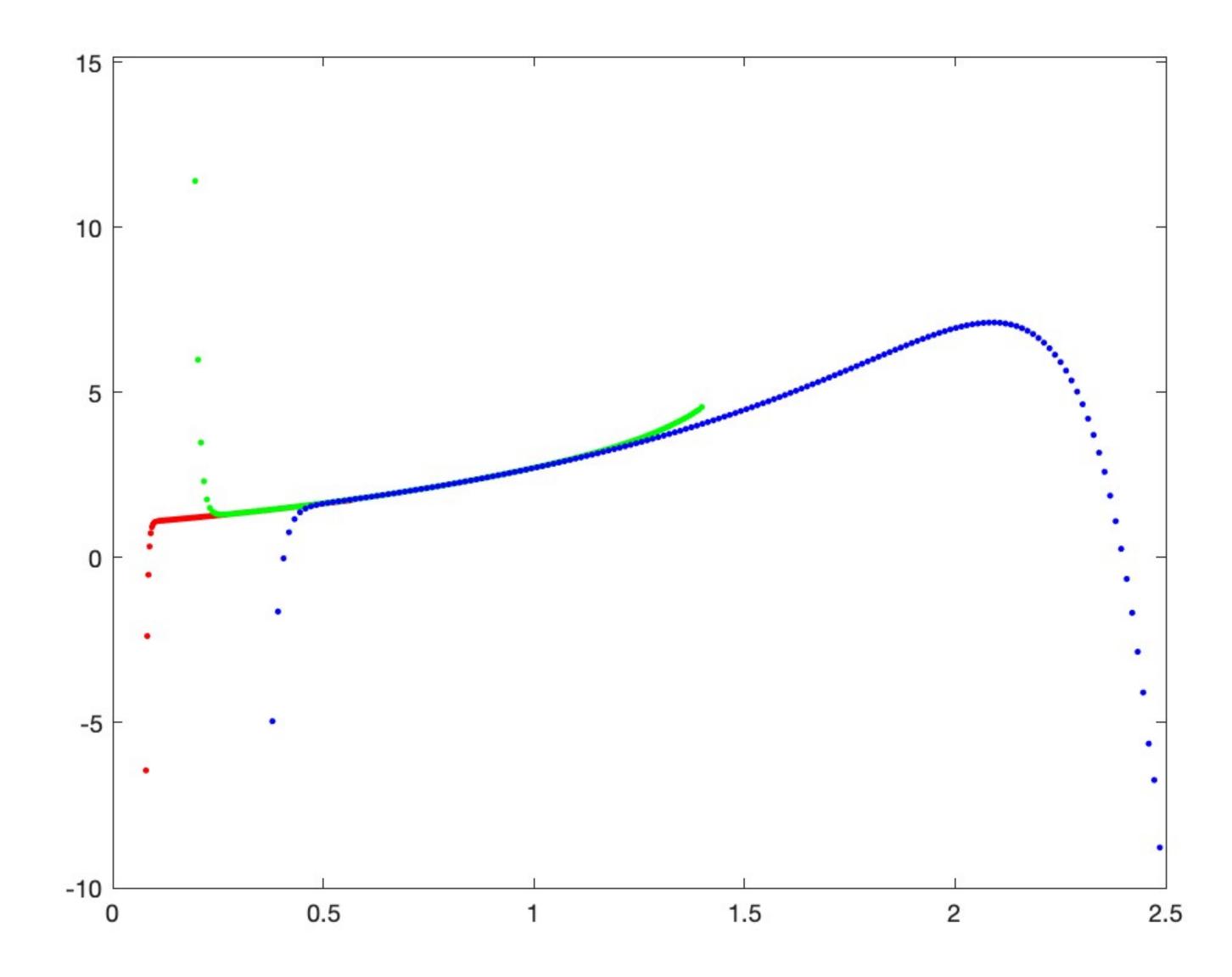
$$k = 2, 3$$



$$f(s_j) \sim \sum_k e^{-t_k(t_0 s - 1)} F(t_k t_0) dt$$

$$f(s) = \frac{1}{s - 1} \rightarrow F(t) = e^t$$

$$k = 2, 3, 4$$

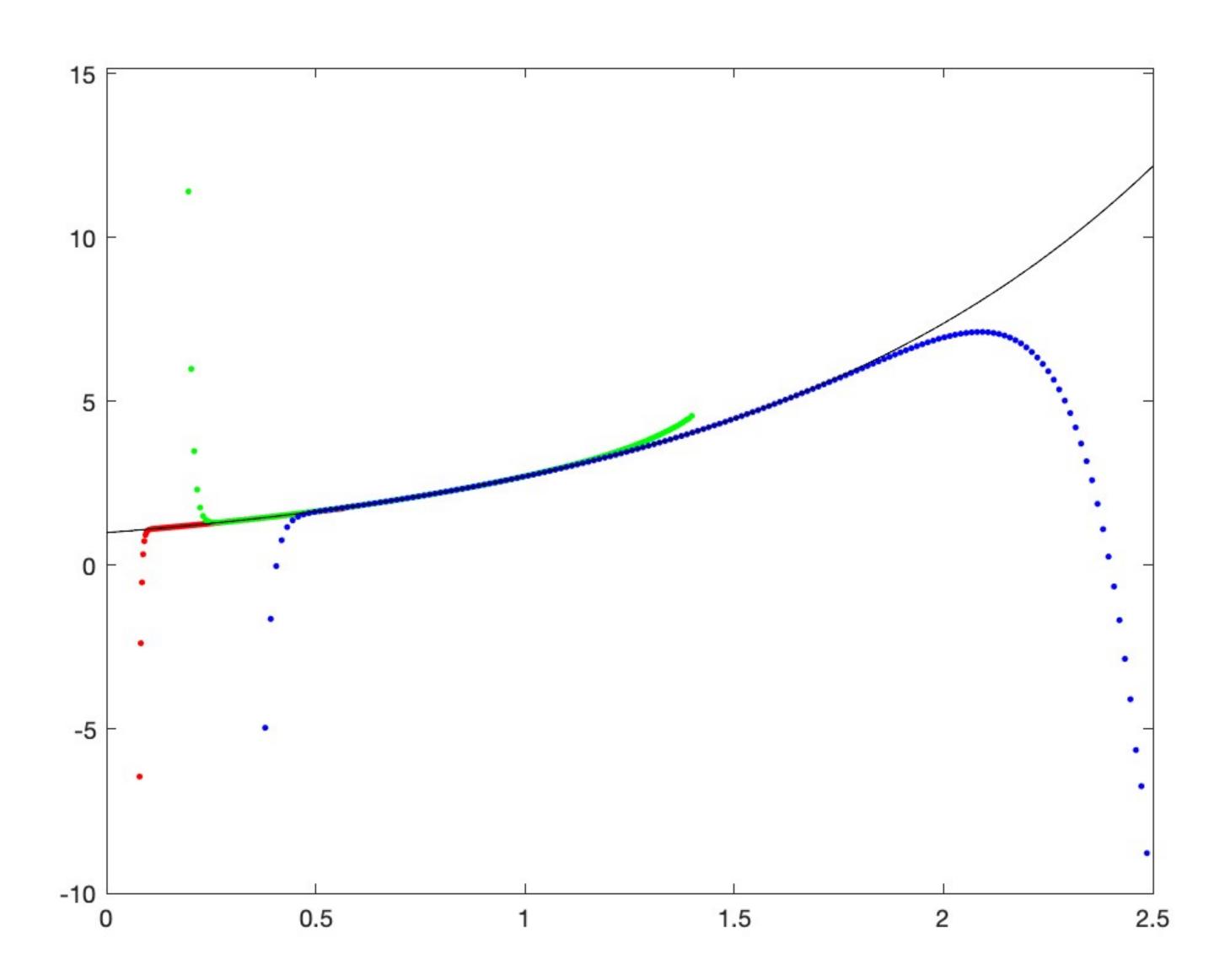


$$f(s_j) \sim \sum_k e^{-t_k(t_0 s - 1)} F(t_k t_0) dt$$

$$f(s) = \frac{1}{s - 1} \quad \rightarrow \quad F(t) = e^t$$

$$k = 2, 3, 4$$

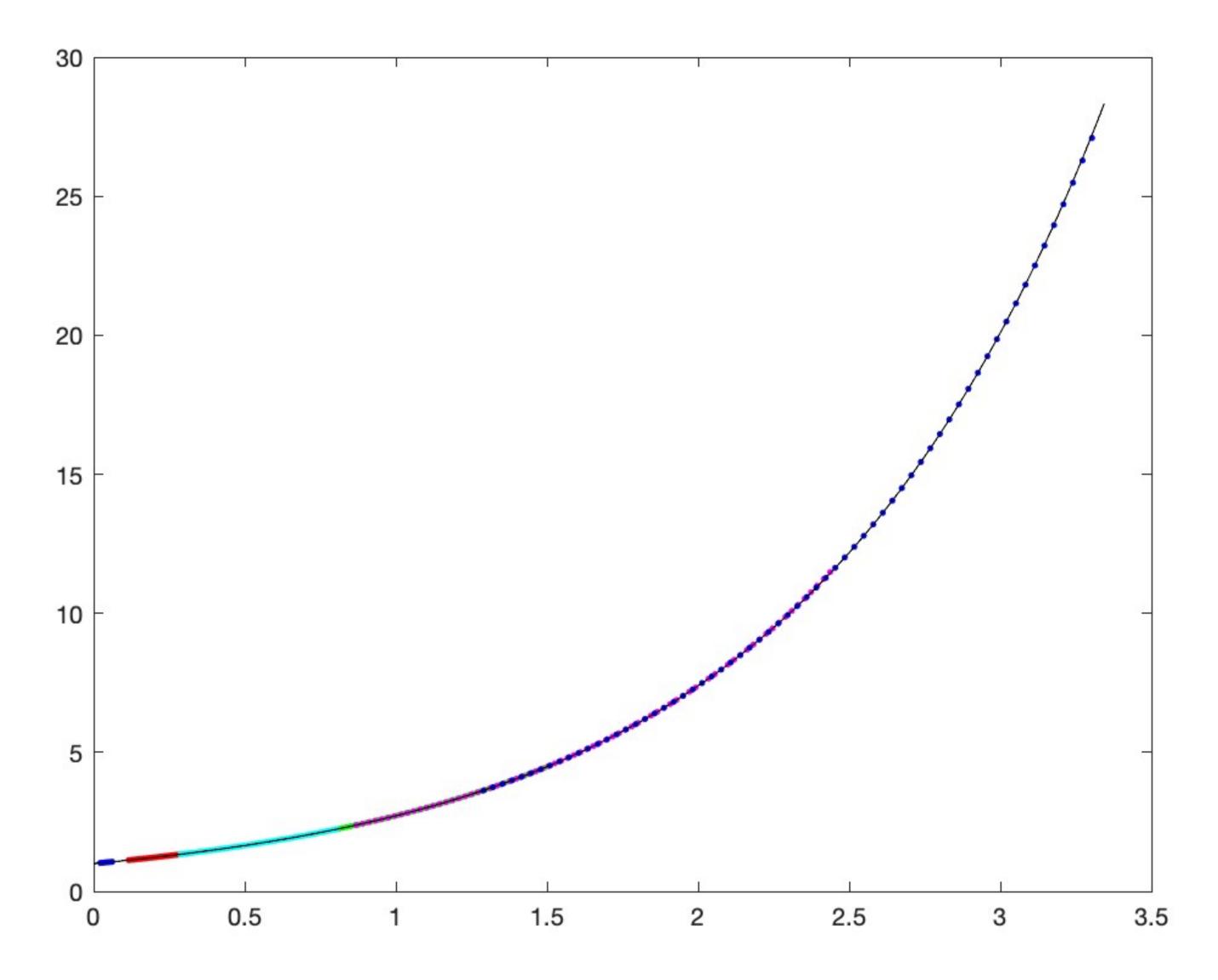
The black line is the **EXPONENTIAL** ...!



... In the end ... DOES IT WORK?!

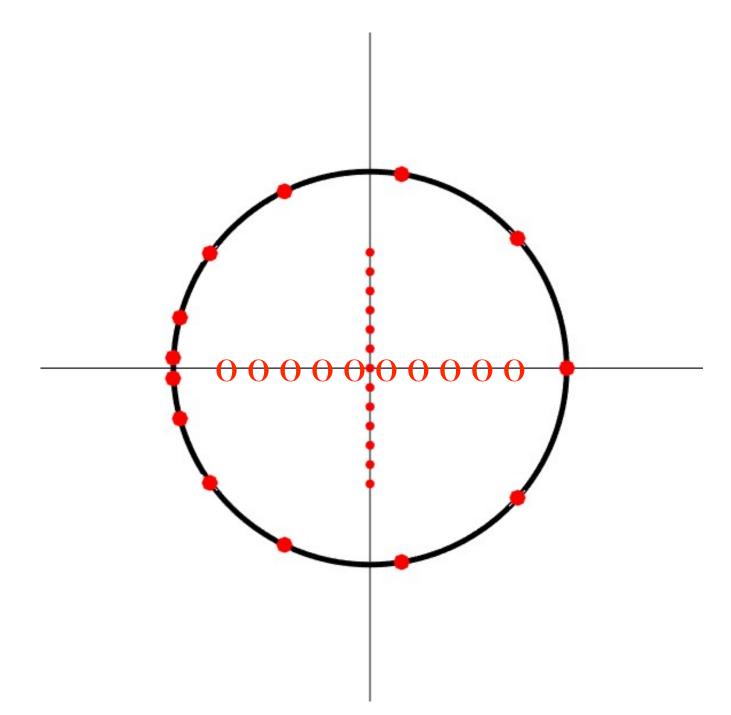
... It is fair to say... IT WORKS!

... provided we look for <u>smooth overlaps</u> ...! (... of course, provided they show up ...)



... but what about the CAUCHY FORMULA?!

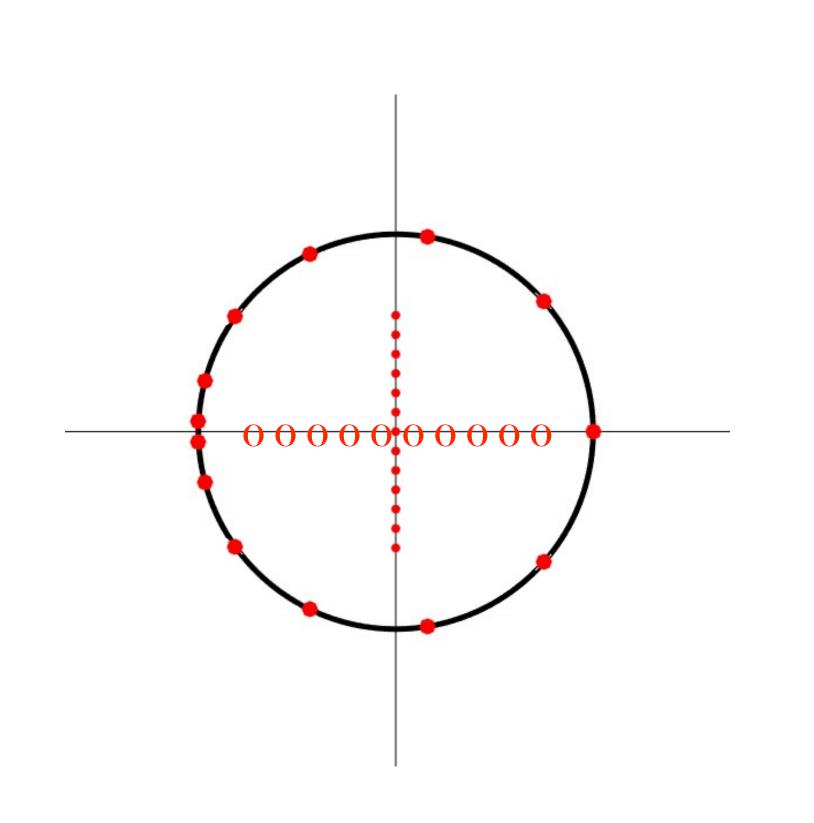
This is the cartoon to remember ...

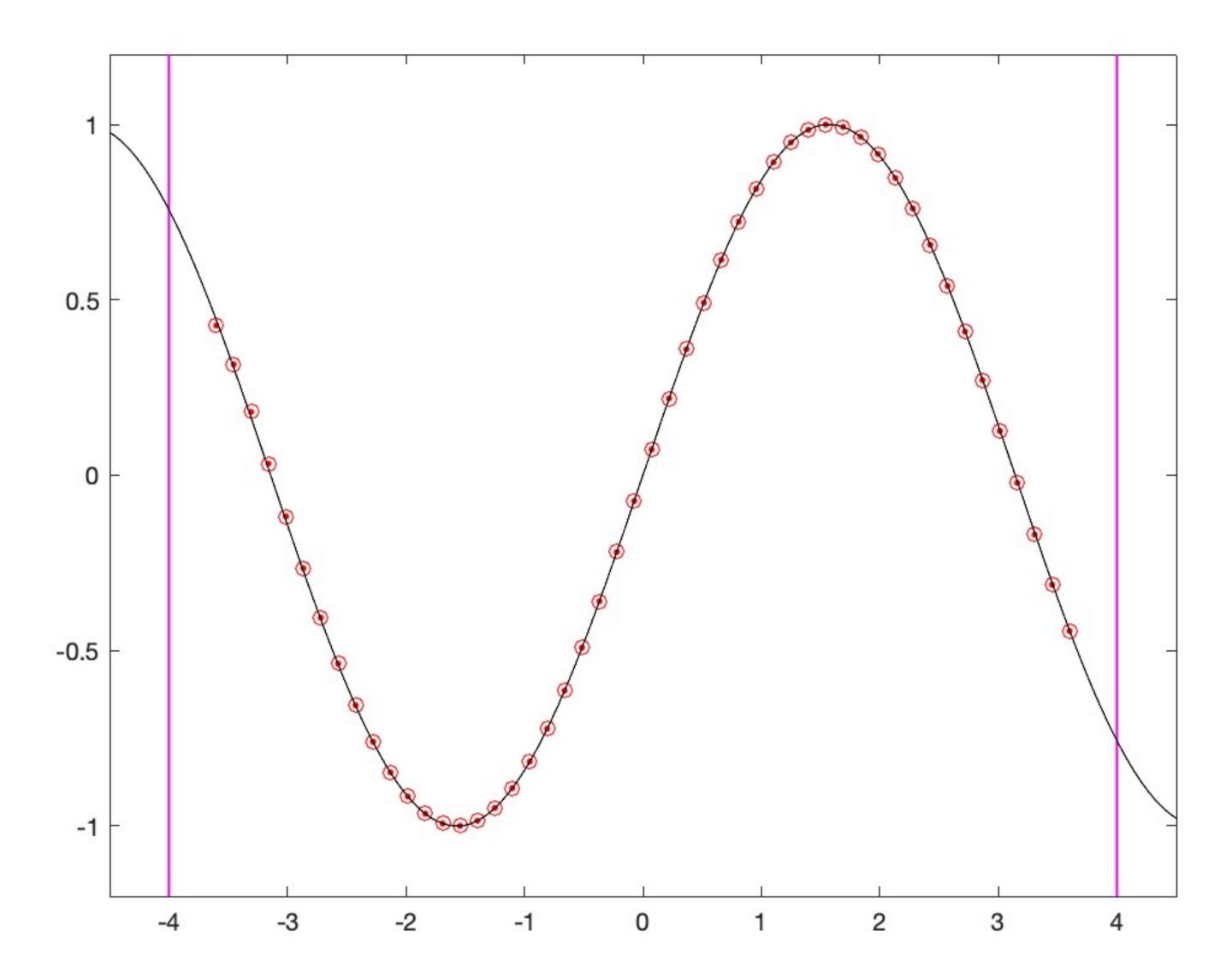


... but what about the CAUCHY FORMULA?!

... but what about the CAUCHY FORMULA?!

Here we go!... computing the sin function on the real axis knowing values on the imaginary axis!

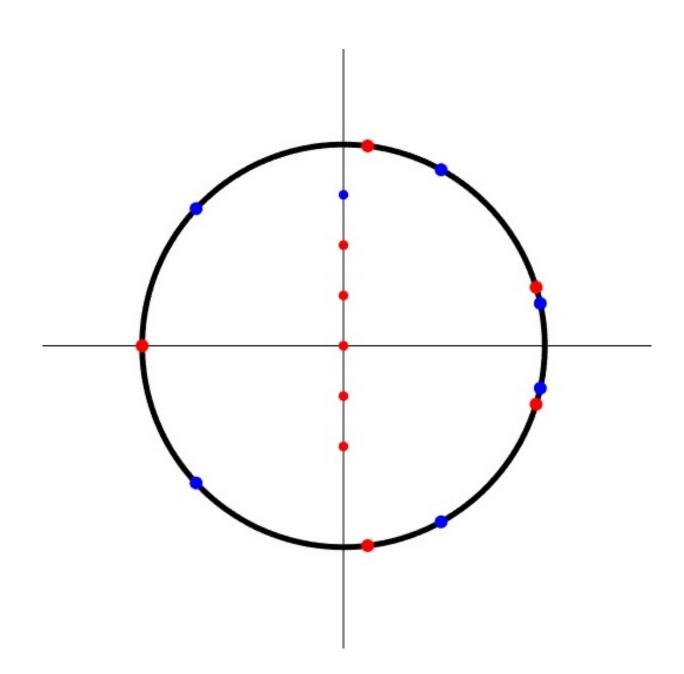




... but what about the CAUCHY FORMULA?! Can we trust this (apparently) good results?

... but what about the CAUCHY FORMULA?! Can we trust this (apparently) good results?

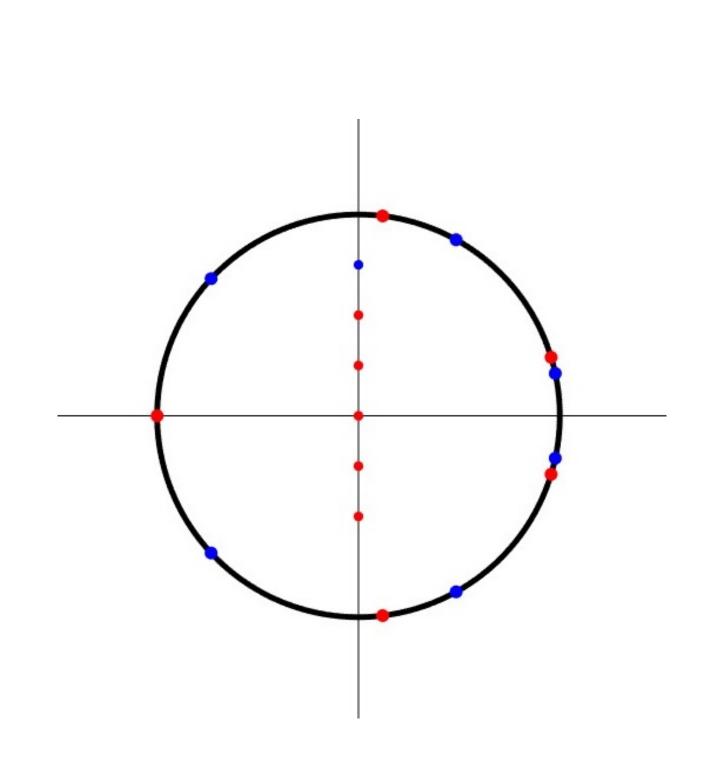
Now this is the cartoon to remember ...

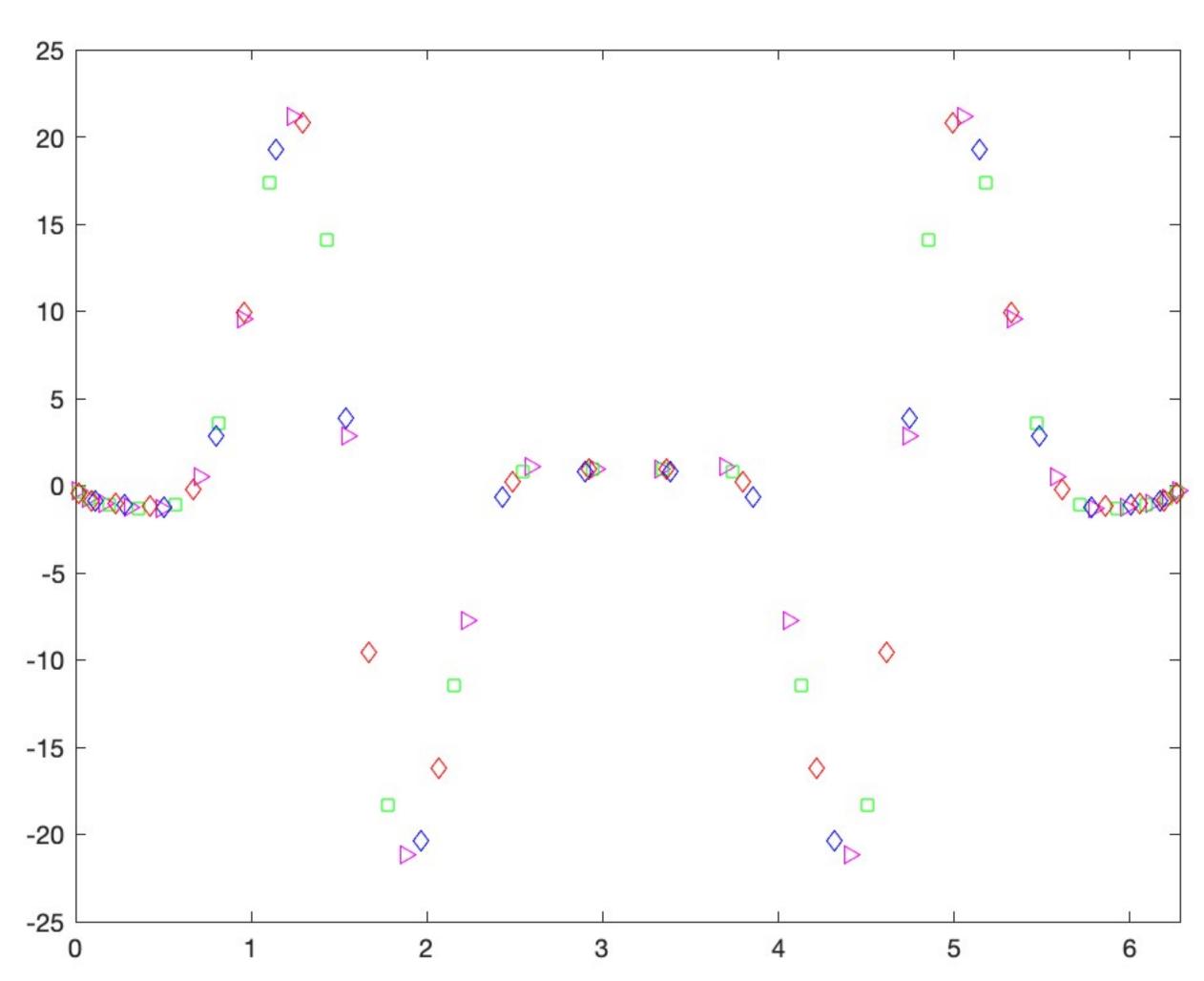


... but what about the CAUCHY FORMULA?! Can we trust this (apparently) good results?

Looking at the solution of the inverse problem, we get a smooth curve showing up!

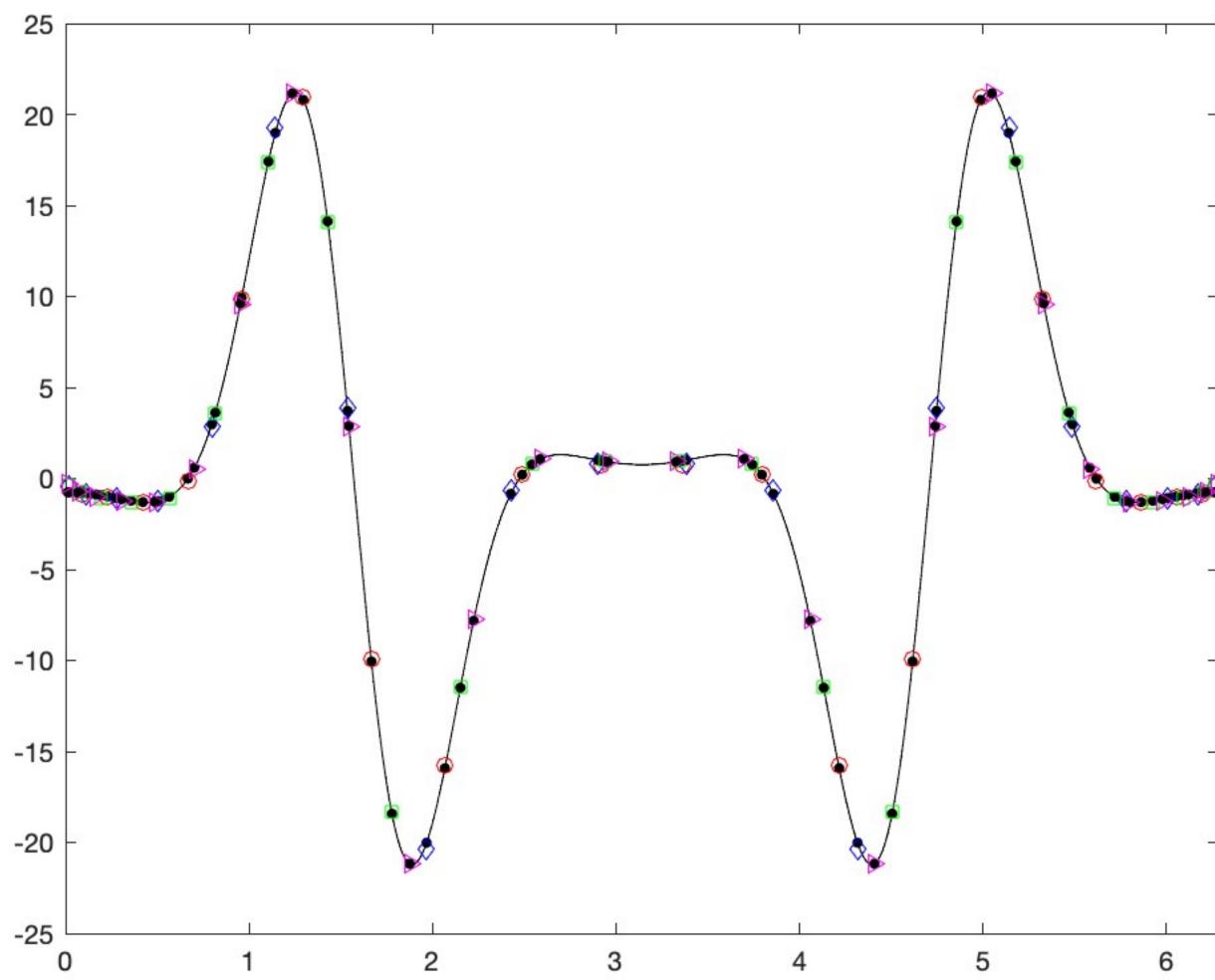
(This is the real part of the sin function at the quadrature points)





(This is the real part of the sin function at the quadrature points)

It is indeed what we should have found ...

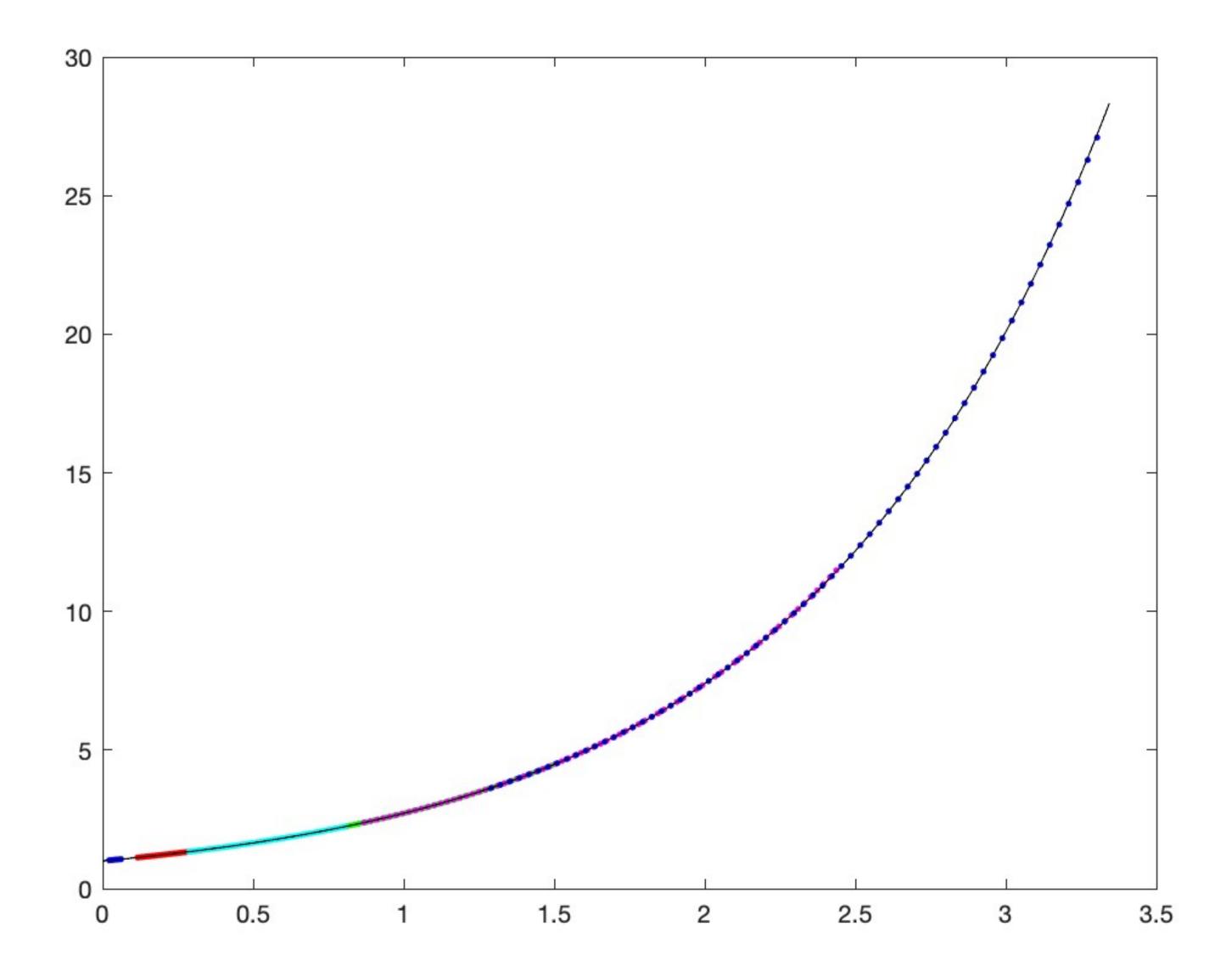


I have not yet told you the entire story ...

I have not yet told you the entire story ...

Do you remember the great success?

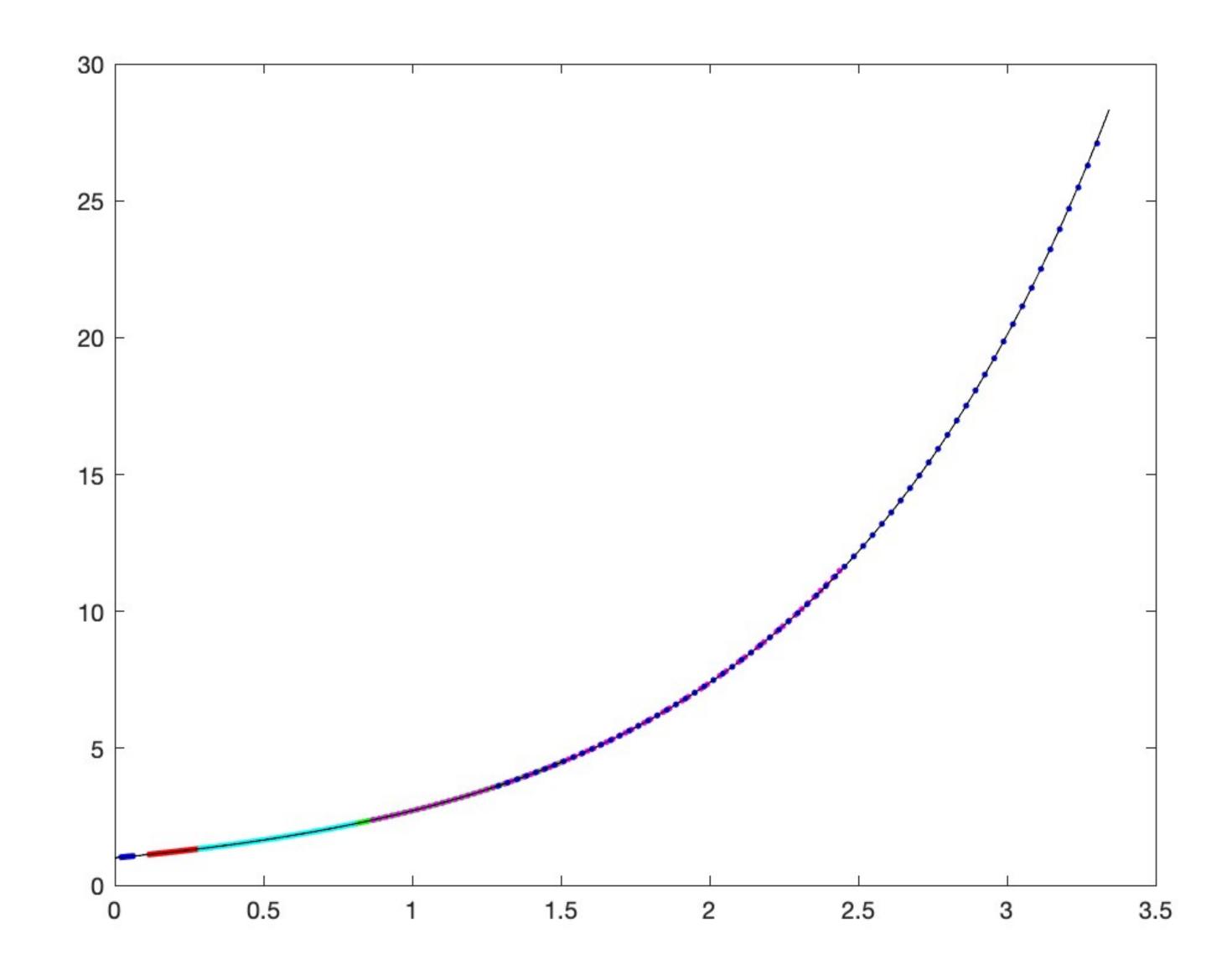
$$f(s_j) \sim \sum_k e^{-t_k(t_0 s - 1)} F(t_k t_0) dt$$
$$f(s) = \frac{1}{s - 1} \longrightarrow F(t) = e^t$$



Do you remember the great success?

$$f(s_j) \sim \sum_k e^{-t_k(t_0 s - 1)} F(t_k t_0) dt$$
$$f(s) = \frac{1}{s - 1} \longrightarrow F(t) = e^t$$

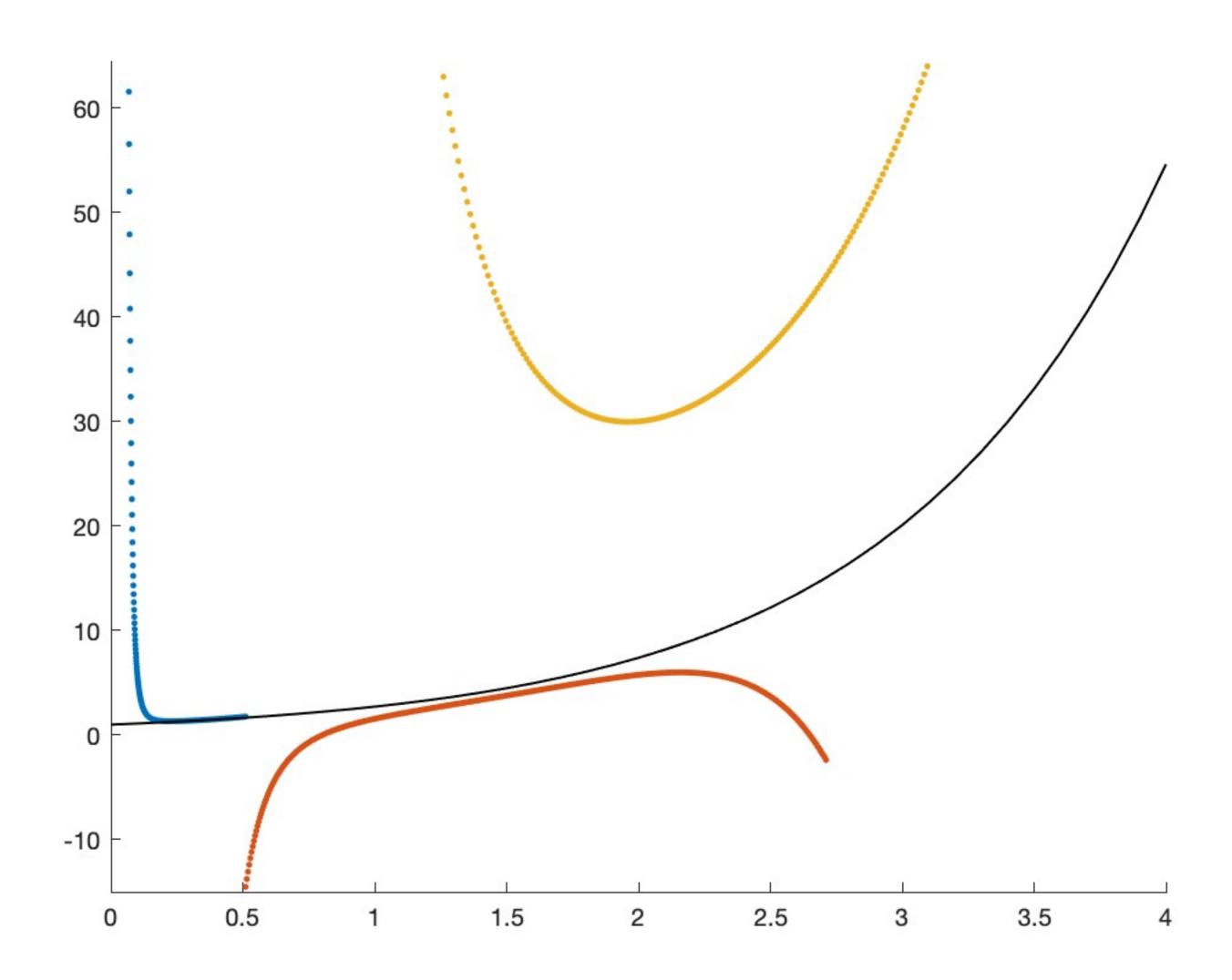
... so, what if ERRORS show up?!?



I have not yet told you the entire story ...

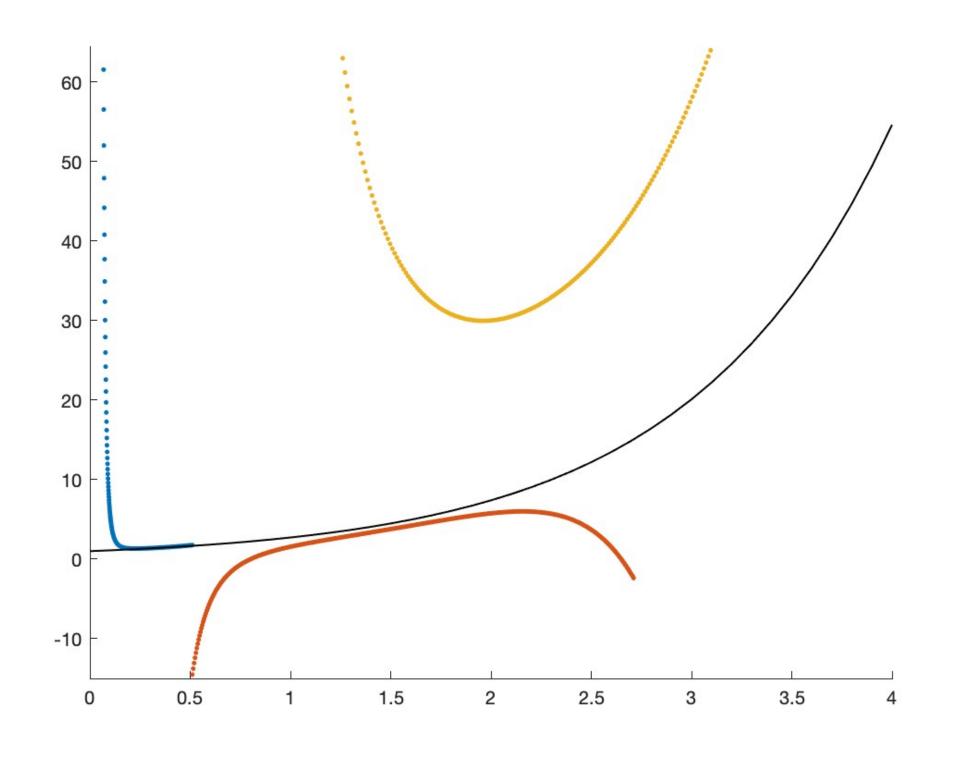
... so, what if ERRORS show up?!?

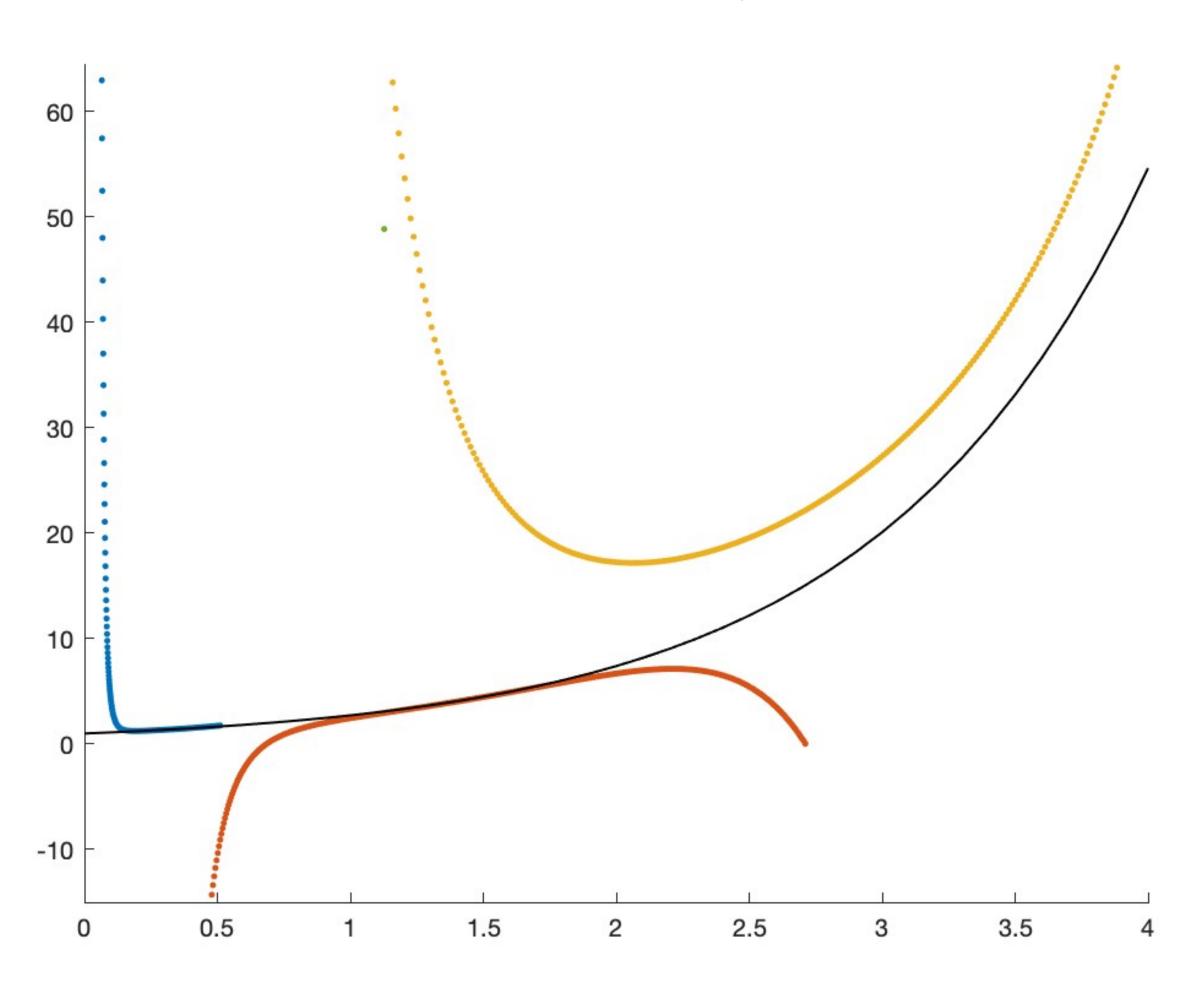
Apparently, you loose everything!



If ERRORS show up, apparently, you loose everything!







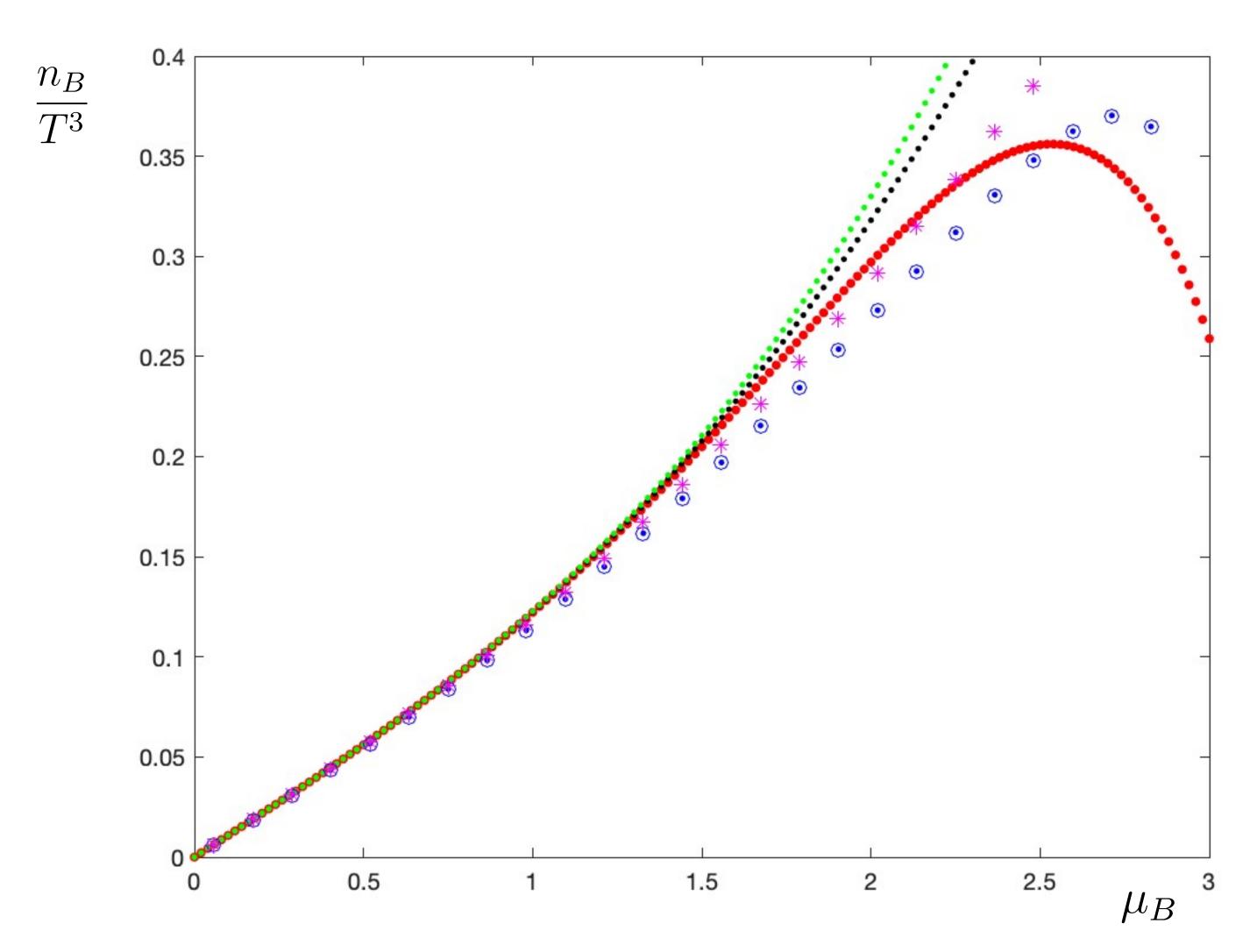
The BIG ISSUE and the (dream?) project: moving points within error-bars, the picture can change quite a lot! This example was found by accident(!), but one can systematically go hunting for a a smooth solution of the inverse problem...

WORK IN PROGRESS



... and still, let's go back to the main point: the lattice QCD sign problem as an inverse problem ...





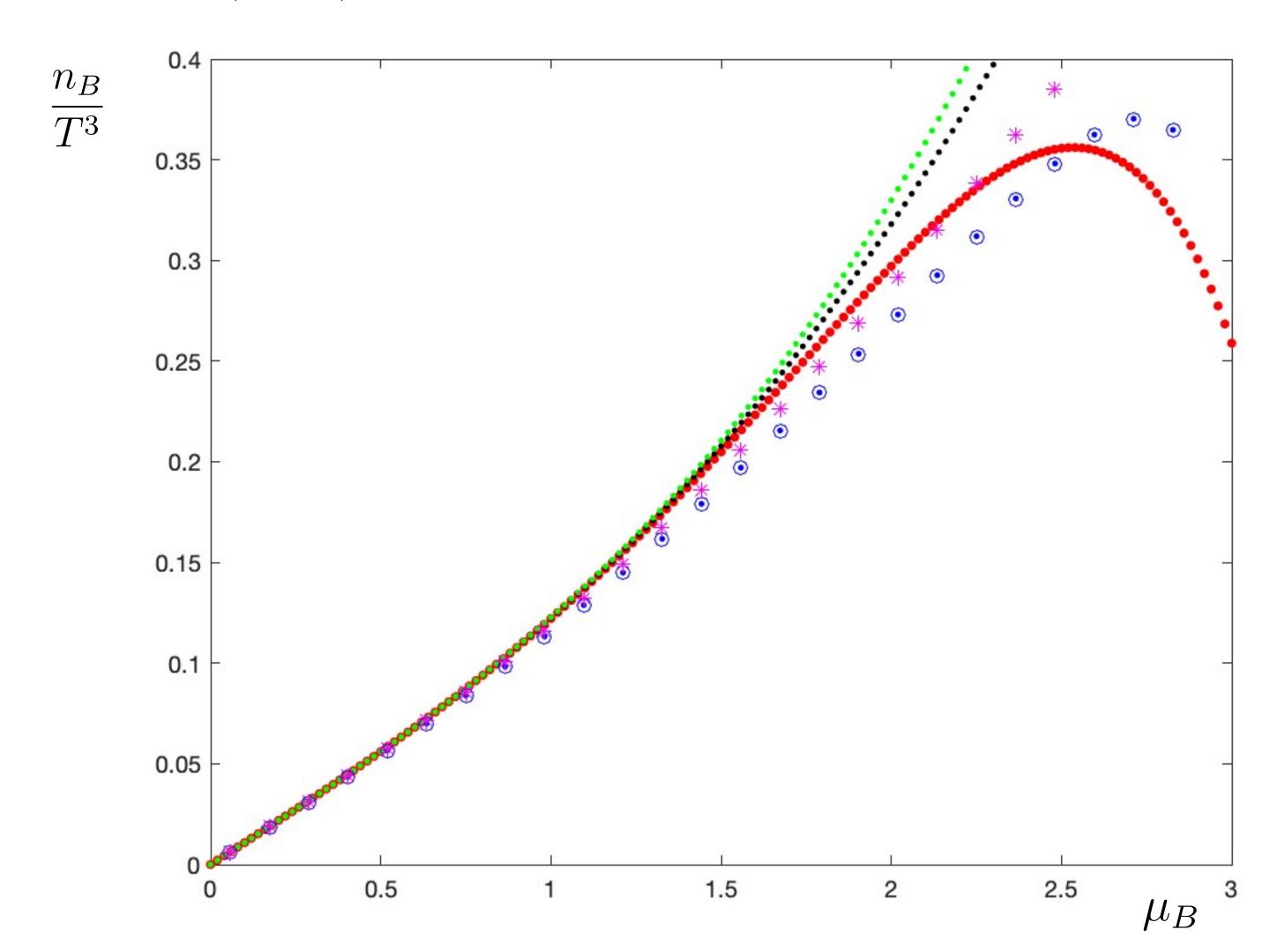
(VERY!) Preliminary results from the Cauchy formula inverse problem:

DOTS are results coming from Taylor coefficients computations (at different orders)

other symbols come
from the Cauchy formula
(still quite sizeable systematic effects,
i.e. sensitivity to the input we provided,
but this is very preliminary

... the dream project still under its way!)

$$T = 157.5 \ (\sim 155) \ \mathrm{MeV}$$



(VERY!) Preliminary results from the Cauchy formula inverse problem:

Symbols come
from the Cauchy formula
(still quite sizeable systematic effects,
i.e. sensitivity to the input we provided,
but this is very preliminary

... the dream project still under its way!)



STAY TUNED!



- Working on (inverse) Cauchy formula for Lattice QCD sign problem
- ...but also on (anti)Laplace transforms (spectral functions and all that...)
- In the end, many applications in (general) inverse problems ...
- ... which also means interpolation and extrapolation problems