Diffusion models for complex Langevin dynamics

Gert Aarts

with Lingxiao Wang, Kai Zhou and Diaa Habibi

JHEP 05 (2024) 060 [2309.17082 [hep-lat]]

NeurIPS workshop 2023 "ML and the Physical Sciences" 2311.03578 [hep-lat]



NeurIPS workshop 2024 "ML and the Physical Sciences" 2410.21212 [hep-lat]

Lattice 2024 2412.01919 [hep-lat]

SIGN 2025, Bern, January 2025

Machine learning is the new playground

many concepts in ML are familiar to theoretical and computational physicists

- o neural networks, say, are systems with many fluctuating degrees of freedom
- training or learning is a minimisation process, typically achieved with stochastic gradient descent (SGD)
- ML parameters are usually contained in matrices

keywords: stochastic dynamics, random matrix theory, non-equilibrium evolution, thermalisation, ...



picture of a playground

Example: Dyson Brownian motion and SGD

• weight matrices are updated using stochastic gradient descent

• stochastic matrix dynamics: random matrix theory

• Coulomb gas and eigenvalue repulsion: implications for training accuracy

• dependence on learning rate (step size) over batch size (size of fluctuations)

GA, Biagio Lucini and Chanju Park, PRE 111 (2025) 1, 015303 [2407.16427 [cond-mat.dis-nn]]

+ Ouraman Hajizadeh, NeurIPS 2024 workshop "ML and the Physical Sciences" [2411.13512 [cond-mat.dis-nn]]

+ Matteo Favoni, Lattice 2024 [2412.20496 [hep-lat]]

This talk: GenAl using Diffusion Models



Diffusion models

stochastic dynamics to generate images (configurations)

start with data set of images

make the images more blurred by applying noise (forward process)



Prior and target distributions

• in pictures: p_0 is target (non-trivial), p_T is the prior (easy)



Outline

- o some comments on diffusion models and stochastic quantisation
- application in lattice scalar field theory in two dimensions
- correlations: higher *n*-point functions and cumulants
- application to sign and complex action problem: complex Langevin dynamics

summary and outlook

- images/configurations are generated during backward process
- stochastic process with time-dependent drift and noise strength

$$\frac{\partial \phi(x,\tau)}{\partial \tau} = g^2(\tau) \nabla_{\phi} \log P(\phi;\tau) + g(\tau) \eta(x,\tau)$$

• write
$$P(\phi; \tau) = rac{e^{-S(\phi, \tau)}}{Z}$$
 such that $abla_\phi \log P(\phi, \tau) = -
abla_\phi S(\phi, \tau)$

$$\circ$$
 then $rac{\partial \phi(x, au)}{\partial au} = -g^2(au)
abla_\phi S(\phi, au) + g(au) \eta(x, au)$

$$\circ$$
 then $rac{\partial \phi(x, au)}{\partial au} = -g^2(au)
abla_\phi S(\phi, au) + g(au) \eta(x, au)$

- very familiar to (lattice) field theorists
- stochastic quantisation (Parisi & Wu 1980)
- path integral quantisation via a stochastic process in fictitious time

$$\frac{\partial \phi(x,\tau)}{\partial \tau} = -\nabla_{\phi} S(\phi) + \sqrt{2} \eta(x,\tau)$$

 \circ stationary solution of associated Fokker-Planck equation $P(\phi) \sim e^{-S(\phi)}$

$$\frac{\partial \phi(x,\tau)}{\partial \tau} = g^2(\tau) \nabla_{\phi} \log P(\phi;\tau) + g(\tau) \eta(x,\tau)$$

$$\frac{\partial \phi(x,\tau)}{\partial \tau} = -\nabla_{\phi} S(\phi) + \sqrt{2} \eta(x,\tau)$$

similarities and differences:

- SQ: fixed drift, determined from known action
 constant noise variance (but can be generalised using kernels)
 thermalisation followed by long-term evolution in equilibrium
- ✓ DM: drift and noise variance time-dependent, learn from data evolution between $0 \le \tau \le T = 1$, many short runs

o diffusion models as an alternative approach to stochastic quantisation



Score matching: learn the drift for backward process

o one degree of freedom, variance-expanding scheme: $\dot{x}(t) = g(t)\eta(t)$ $\eta \sim \mathcal{N}(0,1)$

- time-dependent distribution $P(x,t) = P_t(x)$ describes forward and backward process
- o so-called score $\nabla \log P_t(x)$ is not known, needs to be "learnt" during forward process



Score matching: learn the drift for backward process

o one degree of freedom, variance-expanding scheme: $\dot{x}(t) = g(t)\eta(t)$ $\eta \sim \mathcal{N}(0,1)$

- time-dependent distribution $P(x,t) = P_t(x)$ describes forward and backward process
- o so-called score $\nabla \log P_t(x)$ is not known, needs to be "learnt"

$$\circ \quad \text{loss function } \mathcal{L}(\theta) = \frac{1}{2} \int_0^T dt \, \mathbb{E}_{P_t(x)} \left[\sigma^2(t) \, \|s_\theta(x,t) - \nabla \log P_t(x)\|^2 \right] \qquad \qquad \sigma^2(t) = \int_0^t ds \, g^2(s) \, d$$

 $\circ s_{\theta}(x,t)$ approximates score, vector field learnt by some neural network

o introduce conditional distribution $P_t(x) = \int dx_0 P_t(x|x_0) P_0(x_0)$ initial data $P_0(x_0)$

$$P_t(x) = \int dx_0 \, P_t(x|x_0) P_0(x_0)$$

Score matching: learn the drift

$$\circ$$
 loss function $\mathcal{L}(\theta) = \frac{1}{2} \int_0^T dt \mathbb{E}_{P_t(x)} \left[\sigma^2(t) \| s_\theta(x,t) - \nabla \log P_t(x) \|^2 \right]$

o diffusion process $\dot{x}(t) = g(t)\eta(t)$ easily solved $x(t) = x_0 + \sigma(t)\eta(t)$ $\sigma^2(t) = \int_0^t ds \, g^2(s)$

$$\circ$$
 conditional distribution $P_t(x|x_0) = \mathcal{N}(x;x_0,\sigma^2(t)) = rac{1}{\sqrt{2\pi\sigma^2(t)}}e^{-(x-x_0)^2/(2\sigma^2(t))}$

 \circ and hence $\nabla \log P_t(x_t|x_0) = -(x_t - x_0)/\sigma^2(t)$

$$o \text{ loss function } \mathcal{L}(\theta) = \frac{1}{2} \int_0^T dt \, \mathbb{E}_{P_t(x_t)} \left[\left\| \sigma(t) s_{\theta}(x_t, t) + \frac{x_t - x_0}{\sigma(t)} \right\|^2 \right]$$
$$= \frac{1}{2} \int_0^T dt \, \mathbb{E}_{P_t(x_t)} \left[\left\| \sigma(t) s_{\theta}(x_t, t) + \eta(t) \right\|^2 \right]$$

tractable, computable

Diffusion model for 2d ϕ^4 scalar theory

- \circ 32² lattice, choice of action parameters in symmetric and broken phase
- training data set generated using Hybrid Monte Carlo (HMC)
- variance expanding DM trained using
 U-Net architecture

generating configurations:

- o broken phase
- "denoising" (backward process)
- large-scale clusters emerge, as expected

L Wang, GA, K Zhou, JHEP 05 (2024) 060 [2309.17082 [hep-lat]]

 $\tau = 1$

15

 $\tau = 0.25$ $\tau = 0.5$ $\tau = 0.75$

 $\tau = 0$

Diffusion models

ok, so it seems to work: many questions

- correlations: how are they destroyed and rebuilt?
- usually attention is on two-point function or variance
- higher *n*-point functions contain interactions in field theory
- essential for applications in field theory, correlations = interactions
- focus on moments and cumulants

discuss forward and backward process in more detail



GA, D Habibi, L Wang, K Zhou [2410.21212 [hep-lat]]

Diffusion models in more detail

$$\circ$$
 forward process $\dot{x}(t) = K(x(t),t) + g(t)\eta(t)$ $0 \le t \le T$

noise profile $g(t)=\sigma^{t/T}$

backward process

$$x'(\tau) = -K(x(\tau), T - \tau) + g^2(T - \tau)\partial_x \log P(x, T - \tau) + g(T - \tau)\eta(\tau)$$

score
$$\tau = T - t$$

two main schemes

- \circ variance-expanding (VE): no drift K(x,t) = 0
- variance-preserving (VP) or denoising diffusion probabilistic models (DDPMs):

linear drift $K(x(t),t) = -rac{1}{2}k(t)x(t)$

assume first moment vanishes $x_0 o x_0 - \mathbb{E}_{P_0}[x_0]$

Solve forward process

o forward process $\dot{x}(t) = K(x(t),t) + g(t)\eta(t)$ $K(x(t),t) = -\frac{1}{2}k(t)x(t)$

o initial data from target ensemble $x_0 \sim P_0(x_0)$

• solution
$$x(t) = x_0 f(t, 0) + \int_0^t ds f(t, s) g(s) \eta(s)$$
 $f(t, s) = e^{-\frac{1}{2} \int_s^t ds' \, k(s')}$

• second moment/cumulant/variance $\kappa_2(t) = \mu_2(t) = \mu_2(0) f^2(t,0) + \Xi(t)$

$$\Xi(t) = \int_0^t ds \int_0^t ds' f(t,s) f(t,s') g(s) g(s') \mathbb{E}_{\eta}[\eta(s)\eta(s')] = \int_0^t ds f^2(t,s) g^2(s)$$

 $f(t,s) = e^{-\frac{1}{2}\int_s^t ds' \, k(s')}$

Higher-order moments and cumulants

 \circ moments $\mu_n(t) = \mathbb{E}[x^n(t)]$ and cumulants $\kappa_n(t)$: straightforward algebra

 $\kappa_3(t) = \mu_3(t) = \kappa_3(0)f^3(t,0)$

$$\mu_4(t) = \mu_4(0)f^4(t,0) + 6\mu_2(0)f^2(t,0)\Xi(t) + 3\Xi^2(t)$$

 $\kappa_4(t) = \mu_4(t) - 3\mu_2^2(t) = \left[\mu_4(0) - 3\mu_2^2(0)\right] f^4(t,0) = \kappa_4(0)f^4(t,0)$

$$\kappa_5(t) = \left[\mu_5(0) - 10\mu_3(0)\mu_2(0)\right] f^5(t,0) = \kappa_5(0) f^5(t,0)$$

variance-expanding scheme: no drift

f(t,0) = 1

higher cumulants conserved!

$\Xi(t)=\int_0^T ds\, f^2(t,s)g^2(s)$

Proof to all orders

• generating functionals: average over both noise and target distributions

$$\begin{array}{ll} \text{moments} & Z[J] = \mathbb{E}[e^{J(t)x(t)}] & \text{cumulants} & W[J] = \log Z[J] \\ \\ \text{o noise average} & Z_{\eta}[J] = \mathbb{E}_{\eta}[e^{J(t)x(t)}] = \frac{\int D\eta \, e^{-\frac{1}{2}\int_{0}^{t} ds \, \eta^{2}(s) + J(t) \left[x_{0}f(t,0) + \int_{0}^{t} ds \, f(t,s)g(s)\eta(s)\right]}}{\int D\eta \, e^{-\frac{1}{2}\int_{0}^{t} ds \, \eta^{2}(s)} } \\ \\ \text{o full average} & Z[J] = \mathbb{E}[e^{J(t)x(t)}] = e^{\frac{1}{2}J^{2}(t)\Xi(t)} \int dx_{0} \, P_{0}(x_{0})e^{J(t)x_{0}f(t,0)} \end{array}$$

• cumulant generator
$$W[J] = \log Z[J] = \frac{1}{2}J^2(t)\Xi(t) + \log \int dx_0 P_0(x_0)e^{J(t)x_0f(t,0)}$$

$f(t,s) = e^{-\frac{1}{2}\int_{s}^{t} ds' k(s')}$ E(t) = $\int_{0}^{T} ds f^{2}(t,s)g^{2}(s)$

• cumulant generator
$$W[J] = \log Z[J] = \frac{1}{2}J^2(t)\Xi(t) + \log \int dx_0 P_0(x_0) e^{J(t)x_0 f(t,0)}$$

• 2nd cumulant
$$\kappa_2(t) = \frac{d^2 W[J]}{dJ(t)^2}\Big|_{J=0} = \Xi(t) + \mathbb{E}_{P_0}[x_0^2]f^2(t,0)$$

• higher-order
cumulants
$$\kappa_{n>2}(t) = \frac{d^n W[J]}{dJ(t)^n}\Big|_{J=0} = \frac{d^n}{dJ(t)^n} \log \mathbb{E}_{P_0}[e^{J(t)x_0f(t,0)}]\Big|_{J=0} = \kappa_n(0)f^n(t,0)$$

Toy model: two-peak distribution



o test the predictions in simple zero-dimensional model

• sum of two Gaussians
$$P_0(x) = \frac{1}{2} \left[\mathcal{N}(x; \mu_0, \sigma_0^2) + \mathcal{N}(x; -\mu_0, \sigma_0^2) \right]$$

- o exactly solvable, all even cumulants non-zero, time-dependent score is known analytically
- o quickly show higher-order cumulants, see paper for details

4th, 6th, 8th cumulant with drift (DDPM)



 $\kappa_{n>2}(t) = \kappa_n(0) f^n(t,0)$

 $f(t,0) \rightarrow 0$

$\kappa_{n>2}(t) = \kappa_n(0)$

4th, 6th, 8th cumulant without drift



forward

backward

Higher-order cumulants

• with drift (DDPM): cumulants go to zero, distribution becomes normal

- without drift (variance-expanding): higher-order cumulants are conserved,
 up to numerical cancellations, required between moments which increase in time
- o initial conditions for backward process taken from normal distribution
- score has higher-order cumulants encoded: cumulants are reconstructed



Comparison between schemes

	κ_2	κ_4	κ_6	κ_8
Exact	1.0625	-2	16	-272
Data	1.0624(5)	-2.000(2)	16.00(2)	-272.0(6)
Variance expanding	1.0692(6)	-2.001(2)	16.03(3)	-272.7(6)
Variance preserving (DDPM)	1.0609(5)	-1.976(2)	15.72(2)	-265.6(6)

expectation values at the end of the backward process

✓ variance-expanding scheme slightly outperforms variance-preserving scheme

Two-dimensional scalar fields

extension to scalar fields trivial: each lattice point is treated separately

$$\circ$$
 forward $\partial_t \phi(x,t) = K[\phi(x,t),t] + g(t)\eta(x,t)$

o backward
$$\partial_{\tau}\phi(x,\tau) = -K[\phi(x,\tau),T-\tau] + g^2(T-\tau)\nabla_{\phi}\log P(\phi,T-\tau) + g(T-\tau)\eta(x,\tau)$$

• two-point function $G(x,y;t) \equiv \mathbb{E}[\phi(x,t)\phi(y,t)] = \mathbb{E}_{P_0}[\phi_0(x)\phi_0(y)]f^2(t,0) + \Xi(t)\delta(x-y)$

 \circ moments $\mu_n(x,t) = \mathbb{E}[\phi^n(x,t)]$ independent of x

$\Xi(t) = \int_0^T ds \, f^2(t,s) g^2(s)$

full path integral

with sources

Generating functionals

moment generating

$$Z[J] = \mathbb{E}[e^{J(x,t)\phi(x,t)}] = e^{\frac{1}{2}J^2(x,t)\Xi(t)} \int D\phi_0 P_0[\phi_0] e^{J(x,t)\phi_0(x)f(t,0)}$$

variance preserving

o cumulant generating

$$W[J] = \log Z[J] = \frac{1}{2}J^2(x,t)\Xi(t) + \log \int D\phi_0 P_0[\phi_0] e^{J(x,t)\phi_0(x)f(t,0)}$$

variance

 $f(t,0) \rightarrow 0$

expanding

f(t,0)) = 1

• higher-order cumulants

$$\kappa_{n>2}(t) = \frac{\delta^n W[J]}{\delta J(x,t)^n} \Big|_{J=0} = \frac{\delta^n}{\delta J(x,t)^n} \log \mathbb{E}_{P_0}[e^{J(x,t)\phi_0(x)f(t,0)}] \Big|_{J=0}$$

2nd, 4th, 6th cumulant without drift





forward

backward

Comparison: trained diffusion model

	κ_2	κ_4	κ_6	κ_8
HMC (normalised)	0.39597(4)	-0.29453(6)	0.90108(28)	-5.8689(25)
Diffusion model	0.39598(4)	-0.29454(7)	0.90113(32)	-5.8694(28)

 ϕ^4 theory: $32^2, \kappa = 0.4, \lambda = 0.022, 10^5$ configurations

expectation values at the end of the backward process

excellent agreement



Best Paper Awards 🏅

Sponsored by Apple. Awardees get an iPhone 16 Pro.

BEST 'AI FOR PHYSICS' PAPER AWARD 🐰

Robust Emulator for Compressible Navier-Stokes using Equivariant Geometric Convolutions Wilson G. Gregory, David W Hogg, Kaze W. K. Wong, Soledad Villar [paper]

BEST 'PHYSICS FOR AI' PAPER AWARD 🕇

Higher-order cumulants in diffusion models

Gert Aarts, Diaa Eddin Habibi, Lingxiao Wang, Kai Zhou

[paper] [poster]

Sign problem, complex Langevin dynamics and diffusion models

Diaa Habibi, GA, Lingxiao Wang, Kai Zhou

Lattice 2024 2412.01919 [hep-lat] and in preparation

Stochastic quantisation: complex actions

- stochastic quantisation not limited to real-valued distributions/actions
- extend Langevin process to complex manifold: complex Langevin dynamics (Parisi 1981)

$$z \sim \rho(z) \in \mathbb{C} \quad \Rightarrow \quad x, y \sim P(x, y) \in \mathbb{R}$$

- o convergence not guaranteed, no general solution of Fokker-Planck equation
- o a posteriori justification (GA, Seiler, Stamatescu 2009, Nagata, Nishimura, Shimasaki 2016)
- many talks at this meeting

(Complex) Langevin dynamics

o observables
$$\langle O(x) \rangle = \int dx \, \rho(x) O(x) \qquad \rho(x) = \frac{1}{Z} \exp[-S(x)] \qquad Z = \int dx \, \rho(x)$$

- Langevin equation and drift $\dot{x}(t) = K[x(t)] + \eta(t)$ $K(x) = \frac{d}{dx} \log \rho(x) = -\frac{dS(x)}{dx}$
- Fokker-Planck equation (FPE) $\partial_t \rho(x;t) = \partial_x \left[\partial_x K(x)\right] \rho(x;t)$
- what if weight is complex? drift is complex, FPE only formal
- \circ complexify degrees of freedom $z \rightarrow x + iy$

Complex Langevin dynamics

- \circ complexify degrees of freedom $z \rightarrow x + iy$
- Langevin equation and drift $\dot{z}(t) = K[z(t)] + \eta(t),$ $K(z) = \frac{d}{dz} \log \rho(z) = -\frac{dS(z)}{dz}$
- take real and imaginary part

$$N_x - N_y = 1$$

$$\dot{x}(t) = K_x + \eta_x(t), \qquad K_x = \operatorname{Re}\frac{d}{dz}\log\rho(z), \qquad \langle\eta_x(t)\eta_x(t')\rangle = 2N_x\delta(t-t')$$
$$\dot{y}(t) = K_y + \eta_y(t), \qquad K_y = \operatorname{Im}\frac{d}{dz}\log\rho(z), \qquad \langle\eta_y(t)\eta_y(t')\rangle = 2N_y\delta(t-t')$$

• FPE
$$\partial_t P(x, y; t) = [\partial_x (N_x \partial_x - K_x) + \partial_y (N_y \partial_y - K_y)] P(x, y; t)$$
 $P(x, y; t) \ge 0$

o observables
$$\langle O[x(t)+iy(t)] \rangle_\eta = \int dx dy P(x,y;t) O(x+iy)$$

Complex Langevin dynamics

- FPE $\partial_t P(x, y; t) = \left[\partial_x \left(N_x \partial_x K_x\right) + \partial_y \left(N_y \partial_y K_y\right)\right] P(x, y; t)$
- \circ cannot be solved, non-integrable $\partial_x K_y \neq \partial_y K_x$

• formal justification
$$\int dx dy P(x, y) O(x + iy) = \int dx \rho(x) O(x)$$

- \circ relation (cannot be verified in practice) $ho(x) = \int dy \, P(x-iy,y)$
- o instead, a posteriori criteria for correctness

GA, E Seiler, IO Stamatescu, *Phys. Rev. D* **81** (2010) 054508 [0912.3360] GA, F James, E Seiler, IO Stamatescu, *Eur. Phys. J. C* **71** (2011) 1756 [1101.3270]

Complex Langevin distributions

• FPE
$$\partial_t P(x, y; t) = [\partial_x (N_x \partial_x - K_x) + \partial_y (N_y \partial_y - K_y)] P(x, y; t)$$
 real noise:
 $N_x = 1, N_y = 0$

- o want to describe/understand this distribution
 o further sampling
 o criteria for correctness
 o (modify process)
- o use diffusion model, learn from CL generated data
- diffusion model does not care what the origin of the data is
- note: no solution to the sign problem if CL fails

Gaussian model (solvable)

$$\circ$$
 complex quadratic action $S(x) = rac{1}{2} \sigma_0 x^2$ $\sigma_0 = A + i B$

• CL equations $\dot{x} = K_x + \eta$, $K_x = -Ax + By$, $\dot{y} = K_y$, $K_y = -Ay - Bx$

• here FPE can be solved
$$P(x,y) = N \exp\left[-\alpha x^2 - \beta y^2 - 2\gamma xy\right], \qquad N = \frac{1}{\pi}\sqrt{\alpha\beta - \gamma^2}$$

$$\circ$$
 with coefficients $lpha=A,eta=A(1+2A^2/B^2),\gamma=A^2/B$

$$\rho(x) = \int dy \, P(x-iy,y)$$

○ note: score ≠ CL drift $\partial_x \log P(x, y) = -$

$$(y) = -2\alpha x - 2\gamma y, \qquad \partial_y \log P(x, y) = -2\beta y - 2\gamma x$$

Flow from CL and from score: Gaussian model

A = B = 1



CL dynamics:

$$\partial_x K_y \neq \partial_y K_x$$



 $\partial_x \partial_y \log P(x, y) = \partial_y \partial_x \log P(x, y)$ 39

Trained diffusion model: Gaussian case



Quartic model

$$\circ$$
 simple model with quartic coupling $S=rac{1}{2}\sigma_0x^2+rac{1}{4}\lambda x^4$ $\sigma_0=A+iB$

- o detailed analysis in GA, Giudice, Seiler, Annals Phys. 337 (2013) 238 [1306.3075]
- \circ CL converges, provided $3A^2 B^2 > 0$, dynamics is contained inside a strip, $-y_- < y < y_-$
- o this follows from CL drift

 $y_{-}^{2} = \frac{A}{2\lambda} \left(1 - \sqrt{1 - \frac{B^{2}}{3A^{2}}} \right)$

- FPE can be solved (approximately) using double expansion in Hermite polynomials
- train diffusion model on CL generated data



solution of FPE using double expansion in Hermite polynomials

solution obtained by sampling from trained diffusion model

Trained diffusion model: quartic model



complex Langevin drift

 $A = B = \lambda = 1$

 $y_{-} \approx 0.3029$

Comparison

cumulants in the quartic model

n	2		4		6		8	
	re	—im	re	—im	re	—im	re	—im
Exact	0.428142	0.148010	-0.060347	-0.100083	-0.00934	0.19222	0.41578	-0.5923
CL	0.4277(5)	0.1478(2)	-0.0597(6)	-0.0991(6)	-0.010(1)	0.188(2)	0.406(4)	-0.57(1)
DM	0.4267(6)	0.1459(2)	-0.0582(6)	-0.0981(5)	-0.008(1)	0.188(2)	0.400(5)	-0.58(1)

expectation values at the end of the backward process

note: diffusion model learns from CL data, not the "exact" value

Trained diffusion model: quartic model

very different processes

complex Langevin:

- non-integrable drift
- noise in real direction
- attractor at origin

diffusion model:

- integrable score
- noise in both directions
- saddle at origin



different Fokker-Planck equations

yet same distributions are created for data generation

have obtained access to $\nabla \log P(x, y)$

Summary and outlook

- o diffusion models offer a new approach for ensemble generation to explore in LFT
- learn from data: requires high-quality ensembles
- close relation to stochastic quantisation
- moment- and cumulant-generating functionals

higher *n*-point functions important in LFT applications

- apply to complex actions/complex Langevin: DMs learn elusive real-valued distributions
- in progress: apply to theories with fermions

auxiliary field bosonic models, or DMs learn presence of fermions implicitly

BACKUP SLIDES

f(t,s) = 1

2nd cumulant without drift

 \circ variance-expanding scheme $\kappa_2(t) = \kappa_2(0) + \Xi(t)$

$$\Xi(t) = \int_0^t ds \, g^2(s) \sim \sigma^{2t/T}$$





2nd cumulant with drift (DDPM)

$$f(t,s) = e^{-\frac{1}{2}u(t) + \frac{1}{2}u(s)}$$
$$u(t) = \int_0^t ds \, g^2(s)$$

• variance-preserving scheme $\kappa_2(t) = \mu^2(t) + \sigma^2(t) = \left(\mu_0^2 + \sigma_0^2 - 1\right) f^2(t,0) + 1$



analytic = analytic score