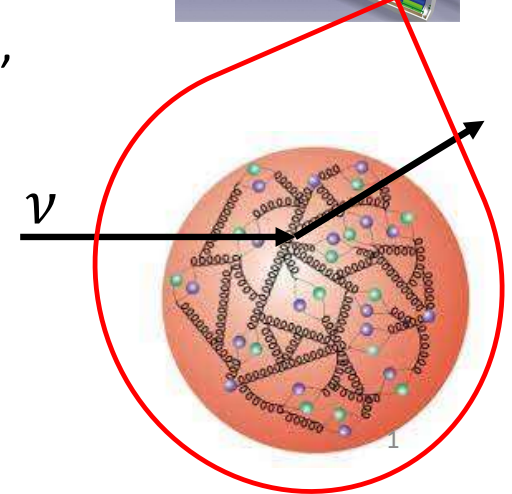
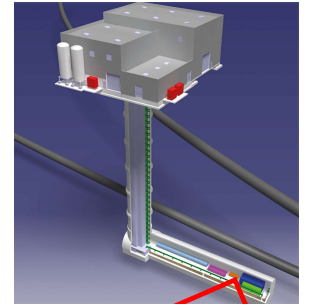


# The LHC as a Neutrino-Ion Collider

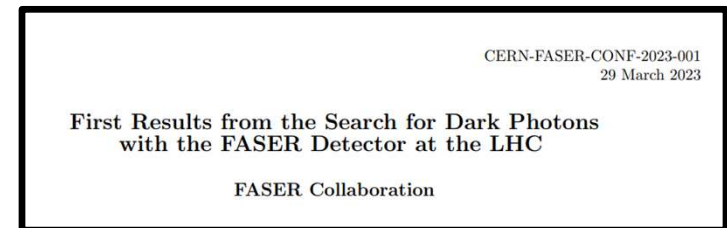
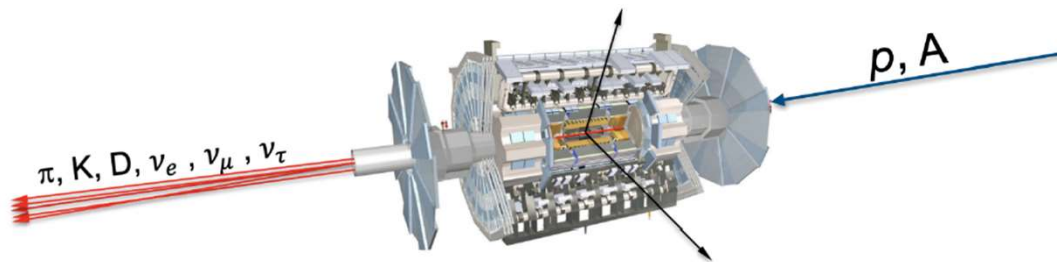
Max Fieg w/

Juan M. Cruz-Martinez , Tommaso Giani , Peter Krack , Toni Makela ,  
Tanjona Rabemananjara , and Juan Rojo

FLArE Forward Physics Working Group 11/7



# Forward Physics



Forward Physics at the LHC has become a very active topic

Many models predict an intense flux of BSM particles in the forward direction

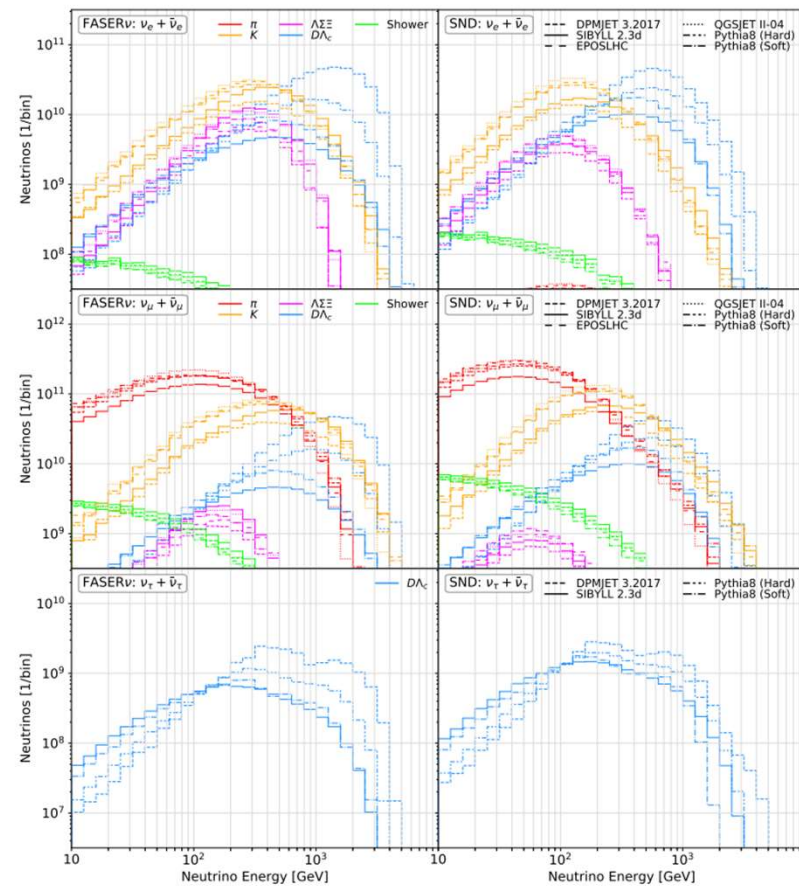
SM predicts intense neutrino beam of all 3 flavors

# Forward Physics – Neutrinos

Neutrinos dominantly produced from hadron decays

Highest energy neutrinos made in a lab setting

In Run 3,  $> 10^{11}$  neutrinos passing through FASERv, and  $O(1000)$  neutrino interactions



# Forward Neutrinos

- Once a background,
- Now a signal, LHC neutrinos have been discovered at FASER and SND
- Can neutrinos be used as a tool?

A. Neutrino-Induced Backgrounds

“...a very forward detector might be able to detect a sizable number of neutrino events...”



First Direct Observation of Collider Neutrinos with FASER at the LHC  
FASER Collaboration



# Forward Physics

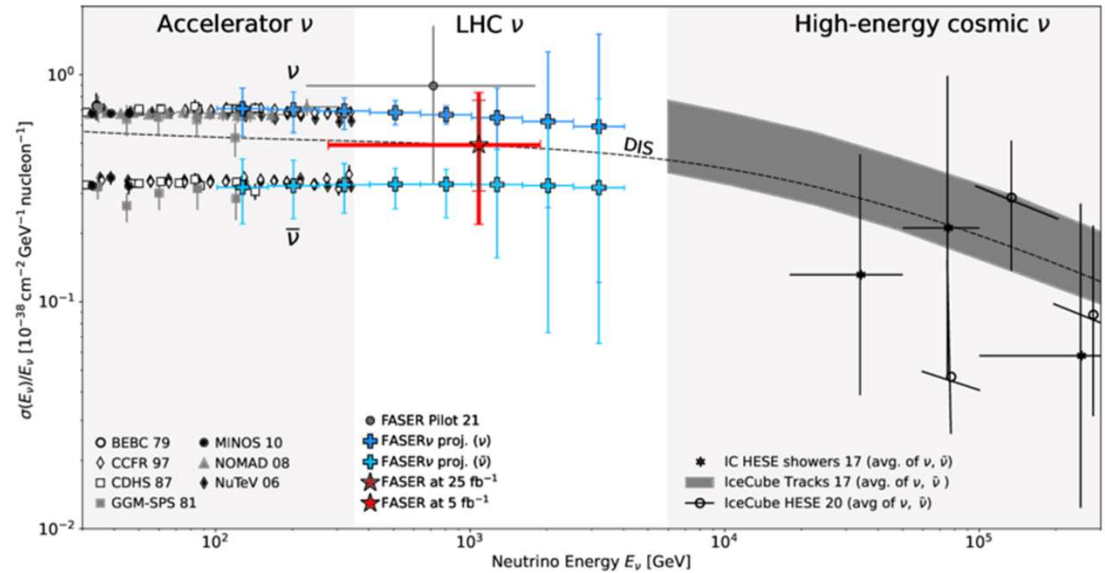
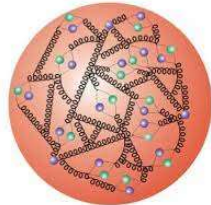
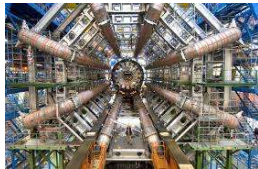
## Neutrino-nucleon cross section

One of the primary targets is the neutrino cross section at uncharted neutrino energies

What more can we do with more statistics?

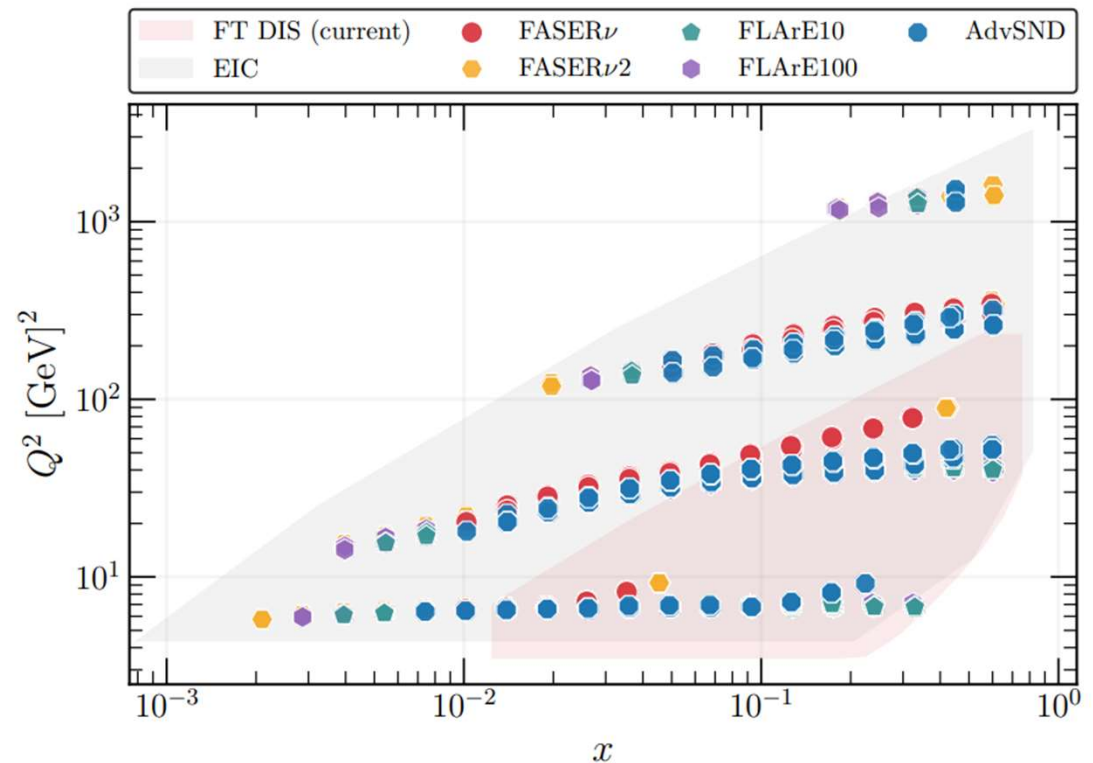


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# Forward Physics at the LHC : A Neutrino Ion collider

- TeV neutrino energy offers new measurements in small Bjorken- $x$  and high- $Q^2$  momentum transfer regime
- Complementary coverage as the proposed Electron-Ion Collider
- After boosting, the forward neutrino program at the LHC can be viewed as a neutrino-ion collider



# Outline

1. DIS pseudodata generation
2. Forward Physics Experiments
3. (n)PDF fitting results
4. Phenomenology implications

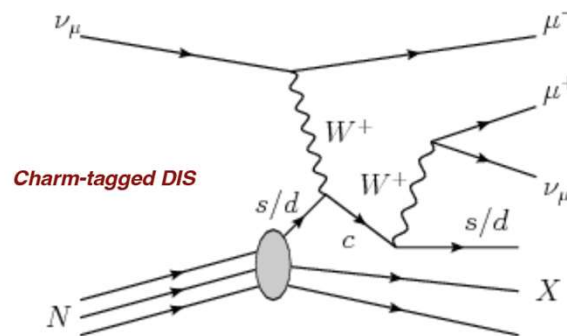
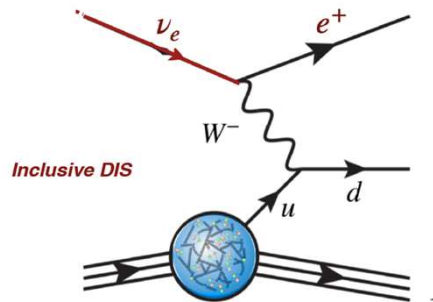
# Deep-inelastic Scattering

Neutrino scattering off a nucleon can be described via

- Bjorken- $x$  momentum fraction,  $Q^2$  momentum transfer,  $E_\nu$

Two processes of interest

- Inclusive DIS
- Charm-tagged DIS





# Pseudodata Generation

- In bins of  $(x, Q^2, E_\nu)$  can write the binned event rate as

$$N_{\text{ev}}^{(i)} = n_T L_T \int_{Q_{\text{min}}^{2(i)}}^{Q_{\text{max}}^{2(i)}} \int_{x_{\text{min}}^{(i)}}^{x_{\text{max}}^{(i)}} \int_{E_{\text{min}}^{(i)}}^{E_{\text{max}}^{(i)}} \frac{dN_\nu(E_\nu)}{dE_\nu} \left( \frac{d^2\sigma(x, Q^2, E_\nu)}{dx dQ^2} \right) \mathcal{A}(x, Q^2, E_\nu) dQ^2 dx dE_\nu$$

Incoming  $\nu$  flux

$\nu$ -nucleon cross-section

Experimental acceptance

# Pseudodata Generation (charm)

- In bins of  $(x, Q^2, E_\nu)$  can write the binned event rate as

$$N_{\text{ev,c}}^{(i)} = n_T L_T \int_{Q_{\text{min}}^{2(i)}}^{Q_{\text{max}}^{2(i)}} \int_{x_{\text{min}}^{(i)}}^{x_{\text{max}}^{(i)}} \int_{E_{\text{min}}^{(i)}}^{E_{\text{max}}^{(i)}} \frac{dN_\nu(E_\nu)}{dE_\nu} \left( \frac{d^2\sigma^{\nu N \rightarrow \ell+c+X}(x, Q^2, E_\nu)}{dx dQ^2} \right) \mathcal{A}(x, Q^2, E_\nu) dQ^2 dx dE_\nu$$

Incoming  $\nu$  flux

$\nu$ -nucleon cross-section

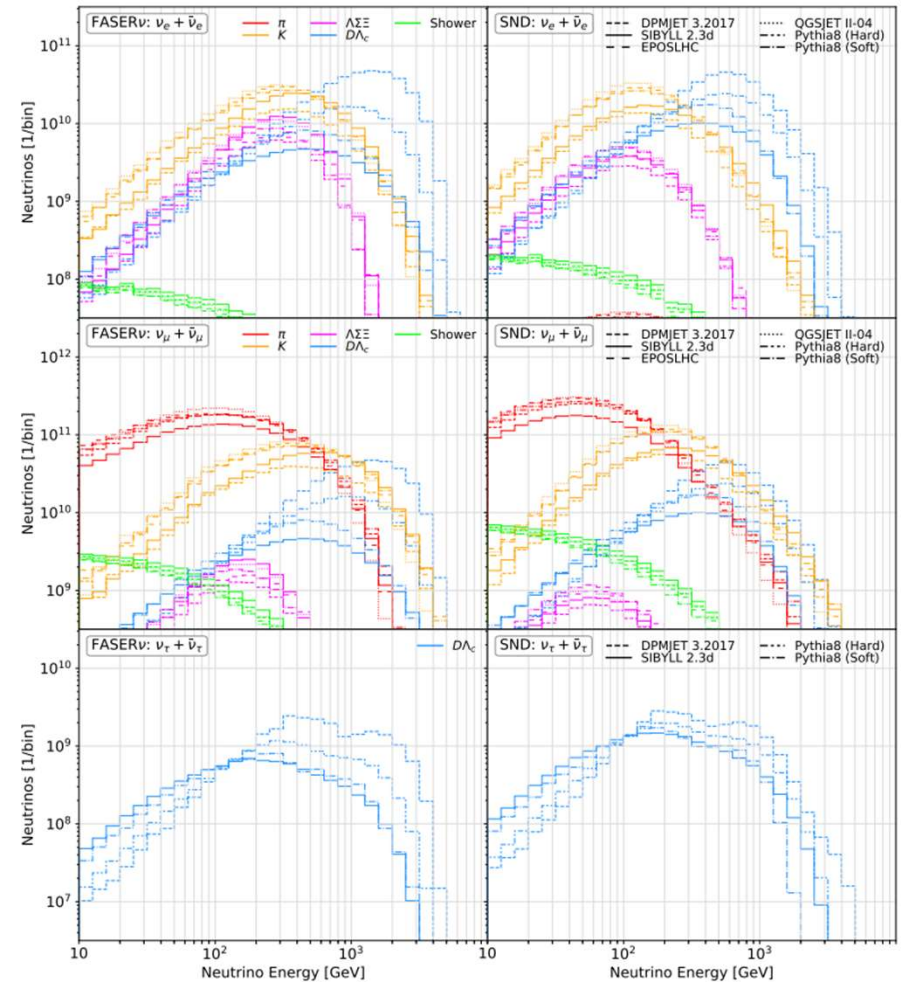
Experimental acceptance

# Neutrino Flux

- Well known that there are flux uncertainties in forward neutrino distribution
  - $\nu_\mu$  is smallest at  $\sim 10 - 20\%$
- Measurements constrain  $\frac{dN}{dE_\nu} \times \sigma$
- Future measurements and modeling will improve on this leading to better predictions of the flux

2309.10417

2309.12793



# DIS structure functions

- Double differential cross-section for (anti-)neutrino free-nucleon CC scattering can be expressed with structure functions,  $F_i^{\nu A}$

$$\frac{d^2\sigma^{\nu A}(x, Q^2, y)}{dxdy} = \frac{G_F^2 s / 4\pi}{(1 + Q^2/m_W^2)^2} [Y_+ F_2^{\nu A}(x, Q^2) - y^2 F_L^{\nu A}(x, Q^2) + Y_- x F_3^{\nu A}(x, Q^2)]$$

$$\frac{d^2\sigma^{\bar{\nu} A}(x, Q^2, y)}{dxdy} = \frac{G_F^2 s / 4\pi}{(1 + Q^2/m_W^2)^2} [Y_+ F_2^{\bar{\nu} A}(x, Q^2) - y^2 F_L^{\bar{\nu} A}(x, Q^2) - Y_- x F_3^{\bar{\nu} A}(x, Q^2)]$$

- Structure functions contain PDF information
- Different quark flavor dependencies between NC and CC

## DIS structure functions – NC vs CC

- For illustration, with  $n_f = 4$  and diagonal CKM, structure functions for CC neutrino NC lepton scattering can be expressed as (with  $Q^2 \ll M_Z^2$ )

- CC: 
$$F_2^{\nu p}(x, Q^2) = 2x (f_{\bar{u}} + f_d + f_s + f_{\bar{c}})(x, Q^2)$$

$$f_{q^\pm} = f_q \pm f_{\bar{q}}$$

- NC: 
$$F_2^{\ell p}(x, Q^2) = x \left( \frac{4}{9} [f_{u^+} + f_{c^+}] + \frac{1}{9} [f_{d^+} + f_{s^+}] \right) (x, Q^2)$$

→ Best sensitivity on quark flavor separation is to use CC from FPF and NC from EIC

# FPF Experiments

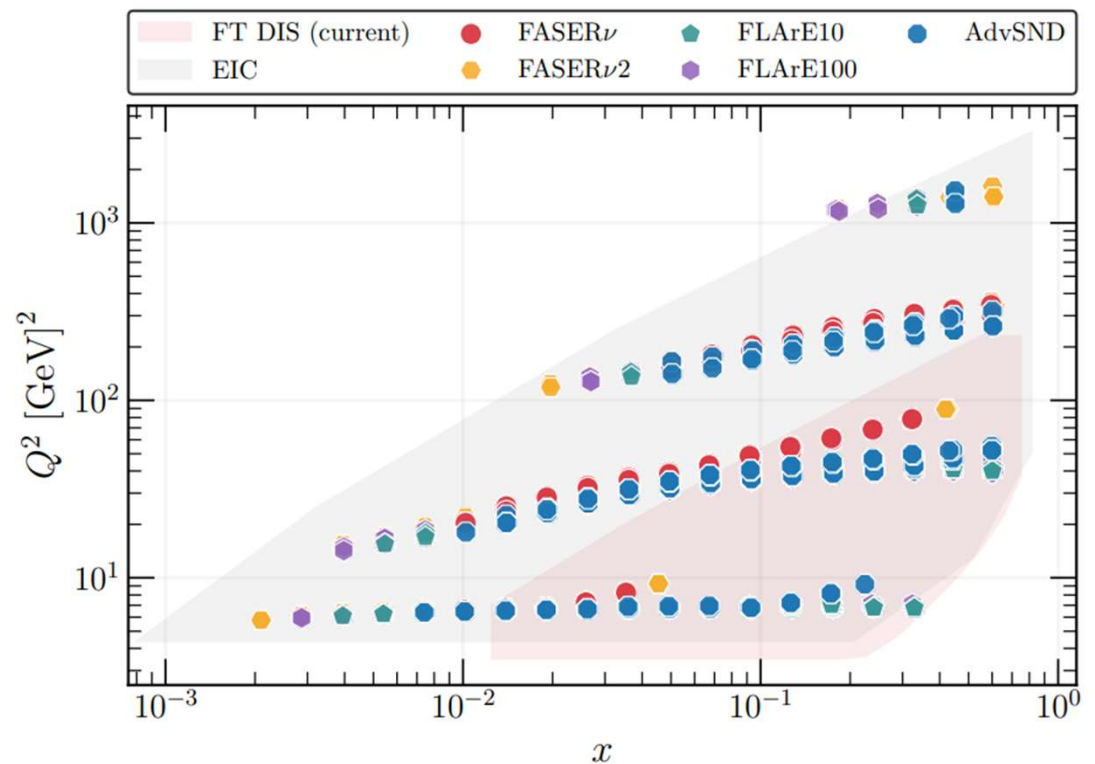
We consider experiments for Run 3 and HL-LHC

- Run 3 : FASERv and SND
  - HL-LHC / FPF : FASERv2, AdvSND, FLArE-10, FLArE-100
- 
- We assess the uncertainties and acceptances related with reconstructing  $(x, Q^2, E_\nu)$  or equivalently  $(E_\ell, E_h, \theta_\ell)$

$$\begin{aligned} E_\nu &= E_h + E_\ell, \\ Q^2 &= 4(E_h + E_\ell)E_\ell \sin^2(\theta_\ell/2) \\ x &= \frac{4(E_h + E_\ell)E_\ell \sin^2(\theta_\ell/2)}{2m_N E_h} \end{aligned}$$

# FPF Experiments - kinematic coverage

- Increase reach in  $x$  and  $Q^2$  by almost an order of magnitude in each direction
- Comparable reach of the highest energy proposed at EIC
- Small- $x$  region relevant for heavy boson production in central region



# FPF Experiments – systematics

- We study DIS event rate for electron and muon neutrinos using some benchmark systematics
  - Cut on  $E_h, E_\ell, \theta_\ell$
- We also study the impact of charge ID and charm ID
- For FLArE-10(0), we measure muon charge and energy using FASER magnet
  - Imposes angular cut

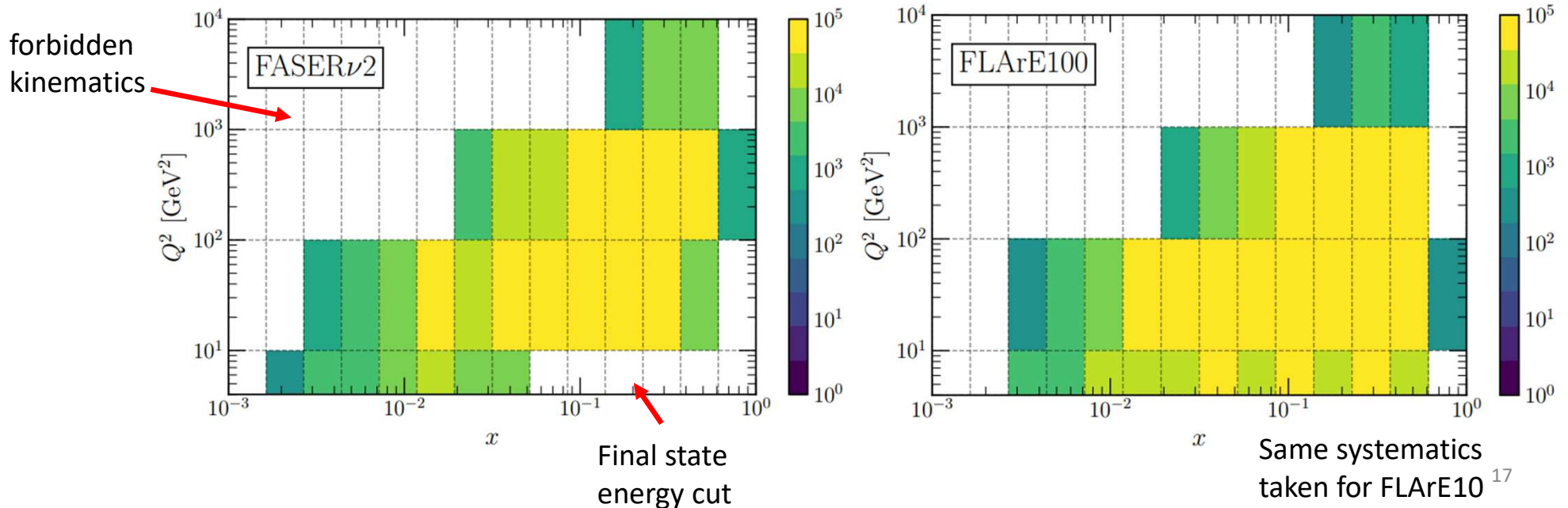
Detector	Rapidity	Target	Charge ID	Acceptance	Performance
FASER $\nu$	$\eta_\nu \geq 8.5$	Tungsten (1.1 tonnes)	muons	$E_\ell, E_h \gtrsim 100$ GeV $\tan \theta_\ell \lesssim 0.025$ (charge ID) reco $E_h$ & charm ID	$\delta E_\ell \sim 30\%$ $\delta \theta_\ell \sim 0.06$ mrad $\delta E_h \sim 30\%$
SND@LHC7	$2 \leq \eta_\nu \leq 8.4$	Tungsten (0.83 tonnes)	n/a	$E_\ell, E_h \gtrsim 20$ GeV $\theta_\mu \lesssim 0.15, \theta_e \lesssim 0.5$	n/a
FASER $\nu 2$	$\eta_\nu \geq 8.5$	Tungsten (20 tonnes)	muons	$E_\ell, E_h \gtrsim 100$ GeV $\tan \theta_\ell \lesssim 0.05$ (charge ID) reco $E_h$ & charm ID	$\delta E_\ell \sim 30\%$ $\delta \theta_\ell \sim 0.06$ mrad $\delta E_h \sim 30\%$
AdvSND-fair2	$2 \leq \eta_\nu \leq 8.4$	Tungsten (5 tonnes)	muons	$E_\ell, E_h \gtrsim 20$ GeV $\theta_\mu \lesssim 0.15, \theta_e \lesssim 0.5$ reco $E_h$	n/a
FLArE (*)	$\eta_\nu \geq 7.5$	LAr (10, 100 tonnes)	muons	$E_\ell, E_h \gtrsim 2$ GeV, $E_e \lesssim 2$ TeV $\theta_\mu \lesssim 0.025, \theta_e \lesssim 0.5$ reco $E_h$	$\delta E_e \sim 5\%, \delta E_\mu \sim 30\%$ $\delta \theta_\ell \sim 15$ mrad $\delta E_h \sim 30\%$

$$N_{\text{ev}}^{(i)} = n_T L_T \int_{Q_{\text{min}}^{2(i)}}^{Q_{\text{max}}^{2(i)}} \int_{x_{\text{min}}^{(i)}}^{x_{\text{max}}^{(i)}} \int_{E_{\text{min}}^{(i)}}^{E_{\text{max}}^{(i)}} \frac{dN_\nu(E_\nu)}{dE_\nu} \left( \frac{d^2\sigma(x, Q^2, E_\nu)}{dx dQ^2} \right) \mathcal{A}(x, Q^2, E_\nu) dQ^2 dx dE_\nu$$



# FPF experiments – $\nu_\mu$ event distribution at FASER $\nu$ 2 and FLArE\*

- O(1M) total event rate  $\rightarrow$   $\sim$ 500k after acceptance cuts
- Reach  $x \sim 10^{-3}$
- Lower energy thresholds at FLArE



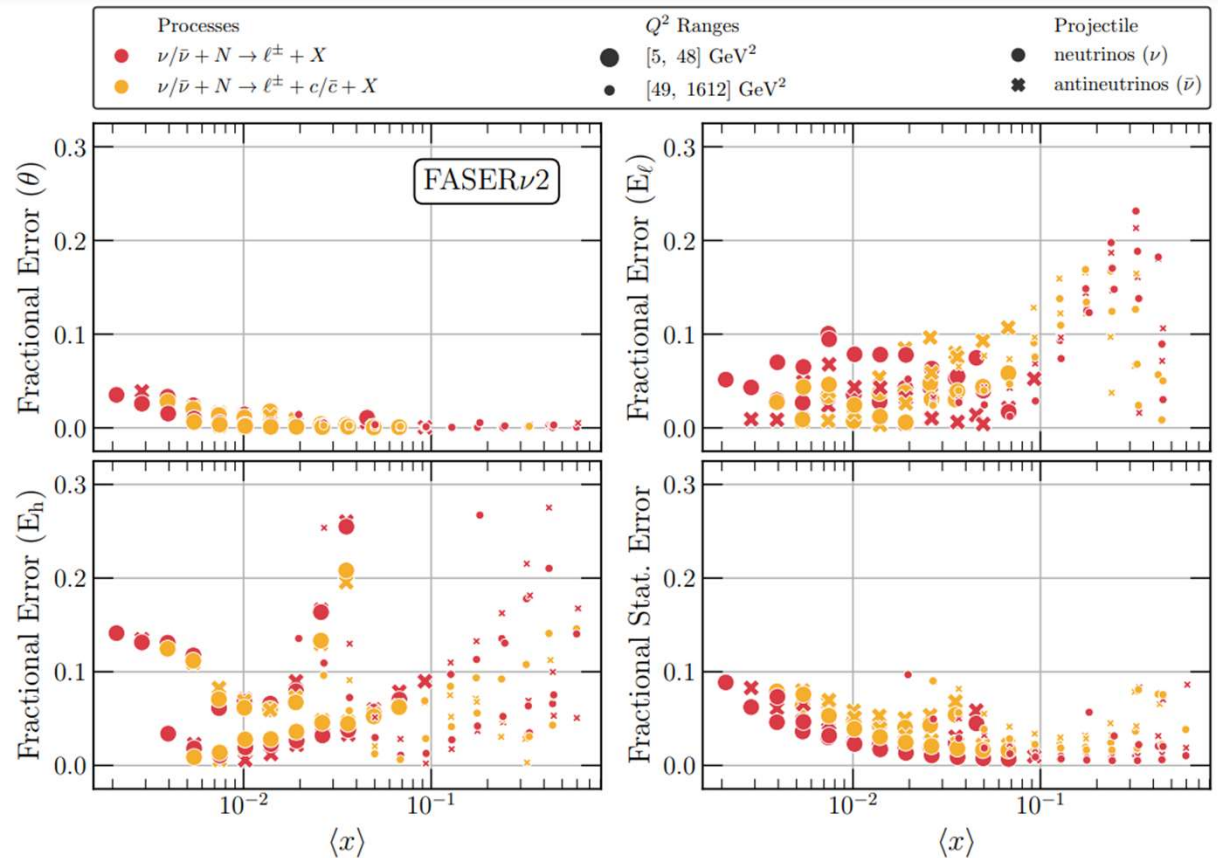
# FPF experiments – inclusive (charm) event rates

Detector	$N_{\nu_e}$	$N_{\bar{\nu}_e}$	$N_{\nu_e} + N_{\bar{\nu}_e}$	$N_{\nu_\mu}$	$N_{\bar{\nu}_\mu}$	$N_{\nu_\mu} + N_{\bar{\nu}_\mu}$
FASER $\nu$	400 (62)	210 (38)	610 (100)	1.3k (200)	500 (90)	1.8k (290)
SND@LHC	180 (22)	76 (11)	260 (32)	510 (59)	190 (25)	700 (83)
FASER $\nu$ 2	116k (17k)	56k (9.9k)	170k (27k)	380k (53k)	133k (23k)	510k (76k)
AdvSND-far	12k (1.5k)	5.5k (0.82k)	18k (2.3k)	40k (4.8k)	16k (2.2k)	56k (7k)
FLArE10	44k (5.5k)	20k (3.0k)	64k (8.5k)	76k (10k)	38k (5.0k)	110k (15k)
FLArE100	290k (35k)	130k (19k)	420k (54k)	440k (60k)	232k (30k)	670k (90k)

- $\nu_\mu$  dominates the event rate over  $\nu_e$  by a factor of 2-3
- HL/FPF brings  $\sim 250x$  more events beyond Run 3 counterpart
- Charm production is  $\sim 15\%$  of event rate  $\rightarrow$  strange PDF

# Systematics uncertainties at FASERv2

- Uncertainties of  $\sim 10 - 30\%$
- Dominated by energy reconstruction
- Statistical errors  $\lesssim 10\%$



# PDF fitting strategy

- PDF4LHC21 - proton
  - *xFitter* with Hessian profiling with prior proton PDF sets
- NNPDF4.0 - proton
  - Direct inclusion into global PDF fit
  - Cross-check for robustness and stability with PDF4LHC21
- EPPS21 – Tungsten Nucleus
  - Nucleus correction



We take two cases

1. Statistical uncertainty only
  - $\delta N = \sqrt{N}$
2. With systematic uncertainties added in quadrature

\*See paper for PDF fitting details

# PDF fit summary

We make lots of comparisons in the paper

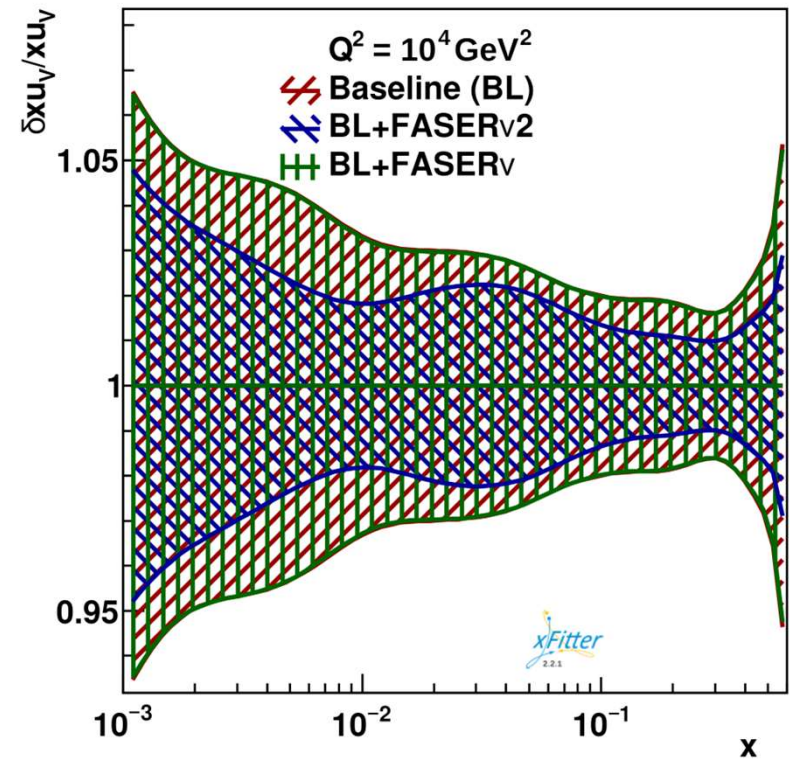
- FASERv(2)
  - Statistics vs systematics
  - Charm ID vs no charm ID
  - Charge ID vs no charge ID
- Experimental comparison
  - FASERv2 vs AdvSND, FLArE10
  - Total FPF data
- NNPDF
  - Consistent and robust results
- EPPS21
  - Tungsten Nucleus
  - qualitatively similar to proton PDF improvement



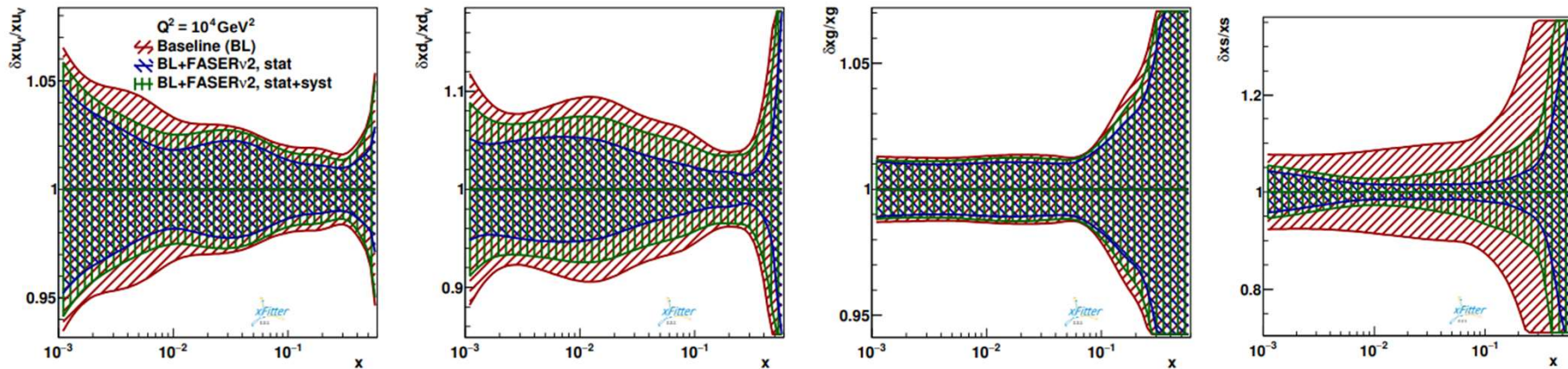
I will highlight a few

# PDF fits – FASER $\nu$ vs FASER $\nu$ 2

- Run 3 statistics too small to be sensitive to PDF fit...
- Not a surprise, let's look at FPF



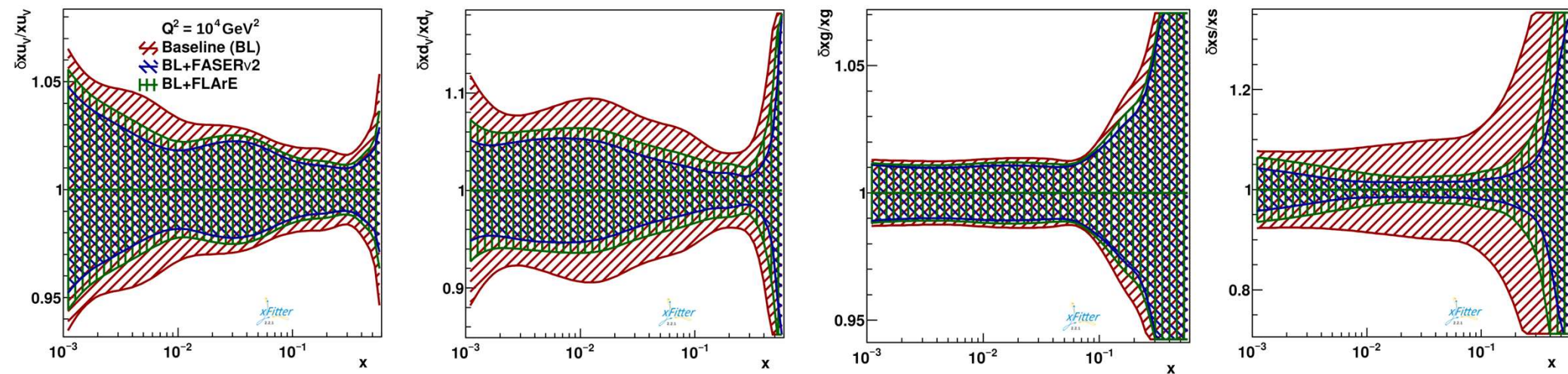
# PDF fits – FASERv2 stat vs systematics



- PDF4LHC21 includes existing neutrino DIS measurements
  - FPF still manages to improve!
- Gluon PDF unaffected  $\rightarrow$  expected for a neutrino scattering experiment
- Huge improvement in strange quark
  - Consequence of charm tagging!



# PDF fits – FASERv2 vs FLArE-10 (stat only)

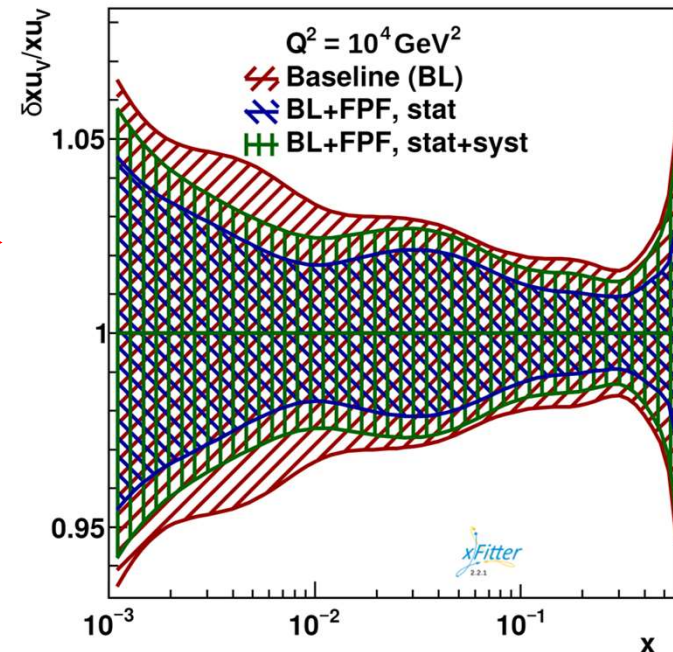
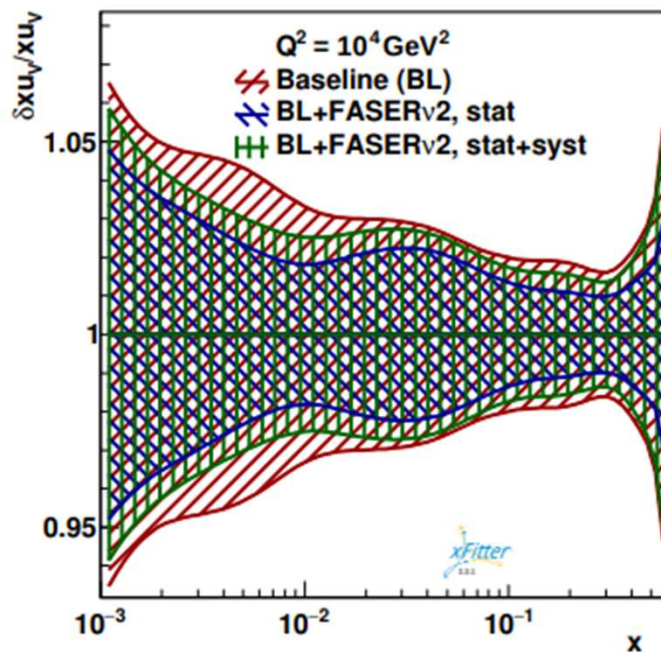


- At  $Q^2 = 10^4 \text{ GeV}^2$ , have somewhat comparable sensitivities
- FLArE would have sensitivity at small recoils

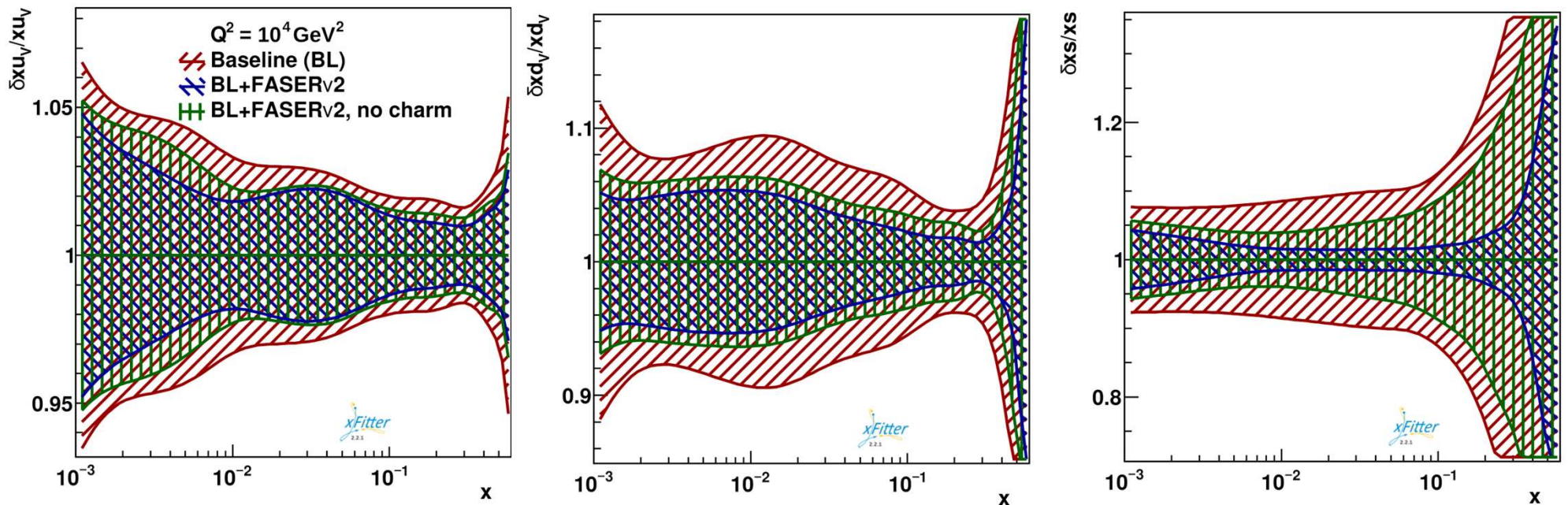


# PDF fits – total FPF

- Results marginally improved with inclusion of additional data  
→ Experiment with largest statistics dominates

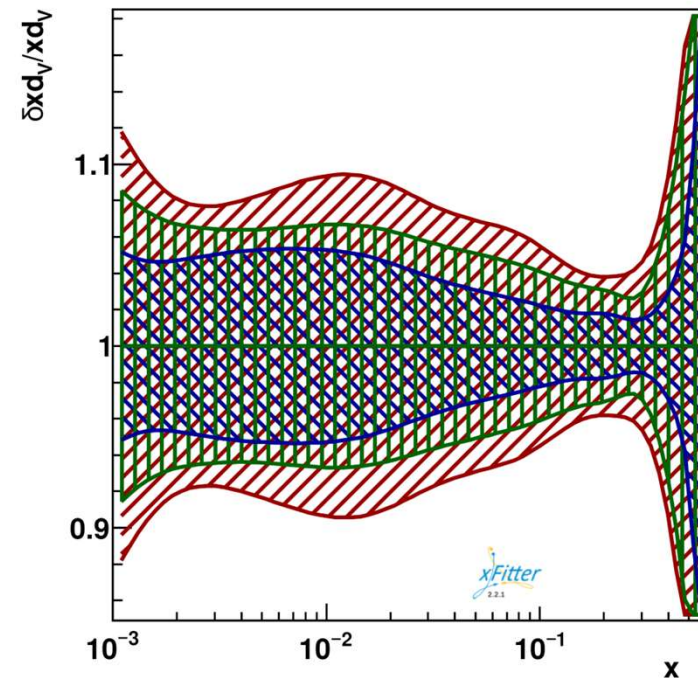
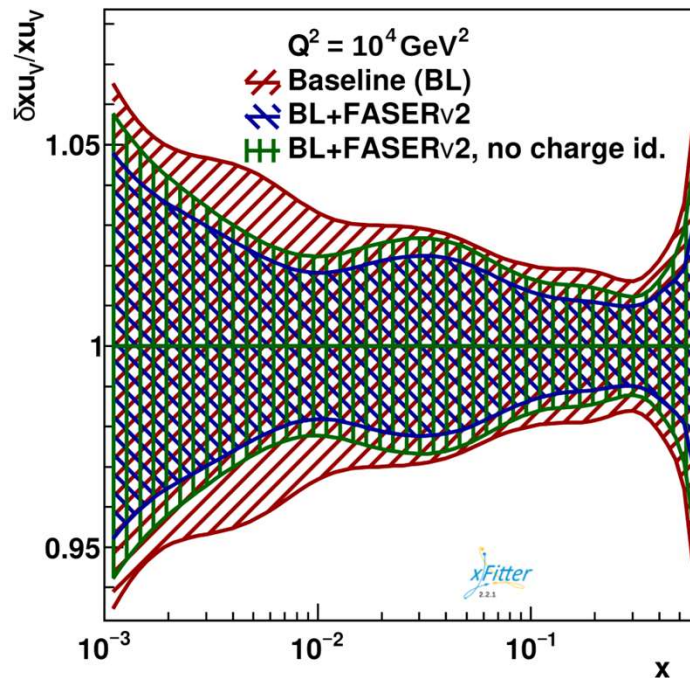


# PDF fits –charm ID



- Small improvement in up quark PDF
- Massive improvement in strange PDF!

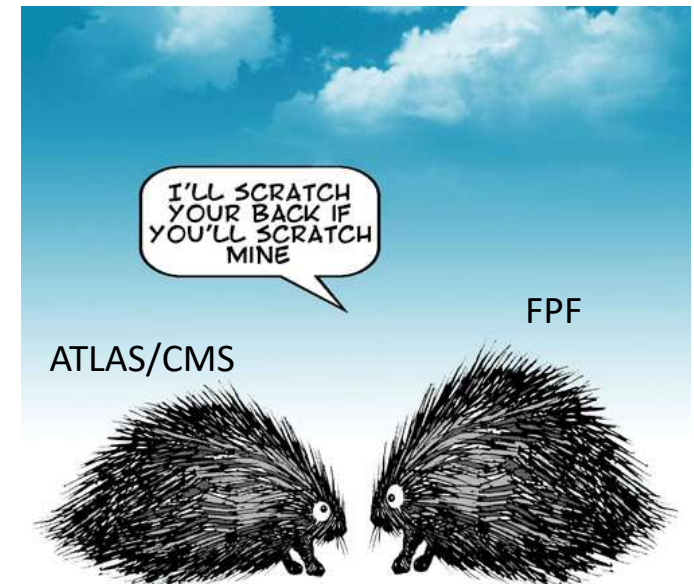
# PDF fits –charge ID



- Small improvement up, down
- Negligible improvement in remaining  $\rightarrow$  charge ID not important

# Phenomenology Implications

- What can we do with improved PDF fits from FPF data?
  - With collected data, have improved understanding of proton quark content
- Look at quark-initiated processes of heavy bosons at LHC!
  - Excellent complementarity

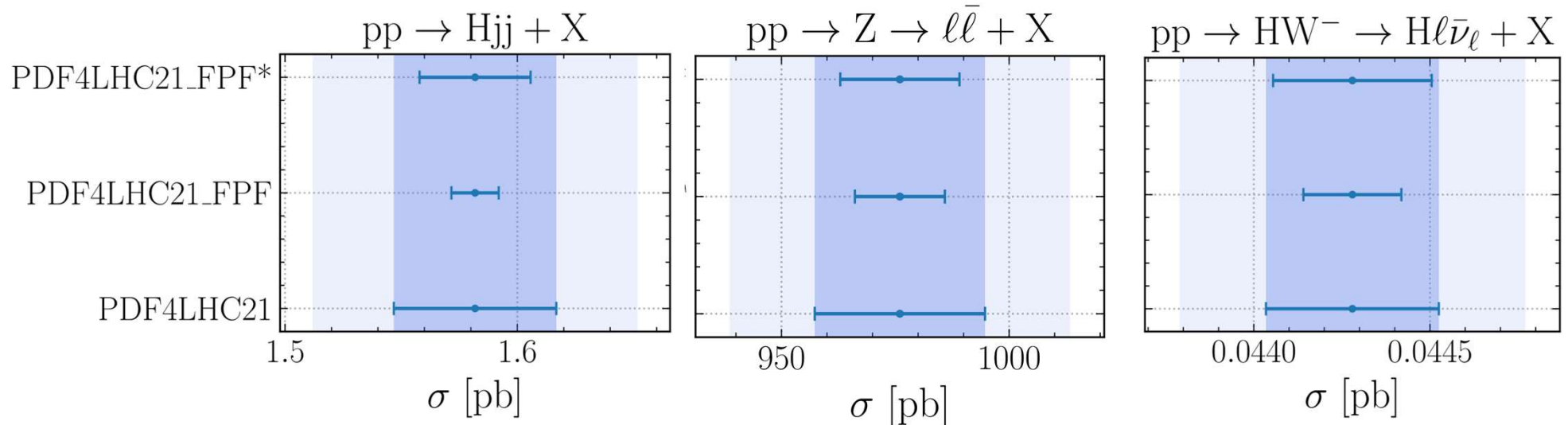


LYNCH



# Phenomenology Implications

- Baseline vs systematics (FPF\*) vs stat only (FPF)
- Forward measurements improve central predictions!
  - Including process relevant for  $m_W$ , and  $\sin^2\theta_W$  measurements



# Conclusions

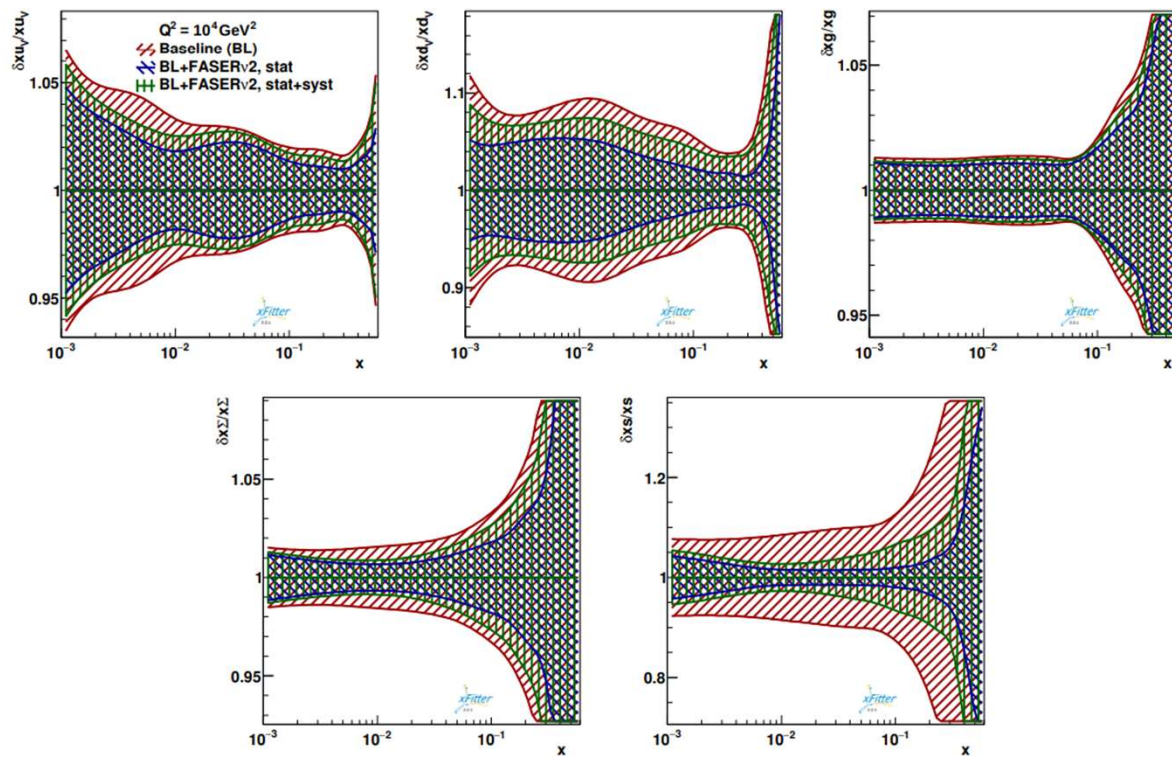
- We calculate neutrino scattering rates at the FPF, with detection systematics folded in
  - Neutrinos are a target and a tool for FPF
- We explore the impact that DIS measurements at Run 3 and the FPF can improve PDF fits
  - Despite wealth of existing data, FPF still manages to improve PDF
  - Greatest gains in strange content due charm tagging
  - Charge ID has small improvement
- Fits from FPF can improve predictions at ATLAS
- Future work includes using gluon PDF at small-x to constrain production

Thank you!

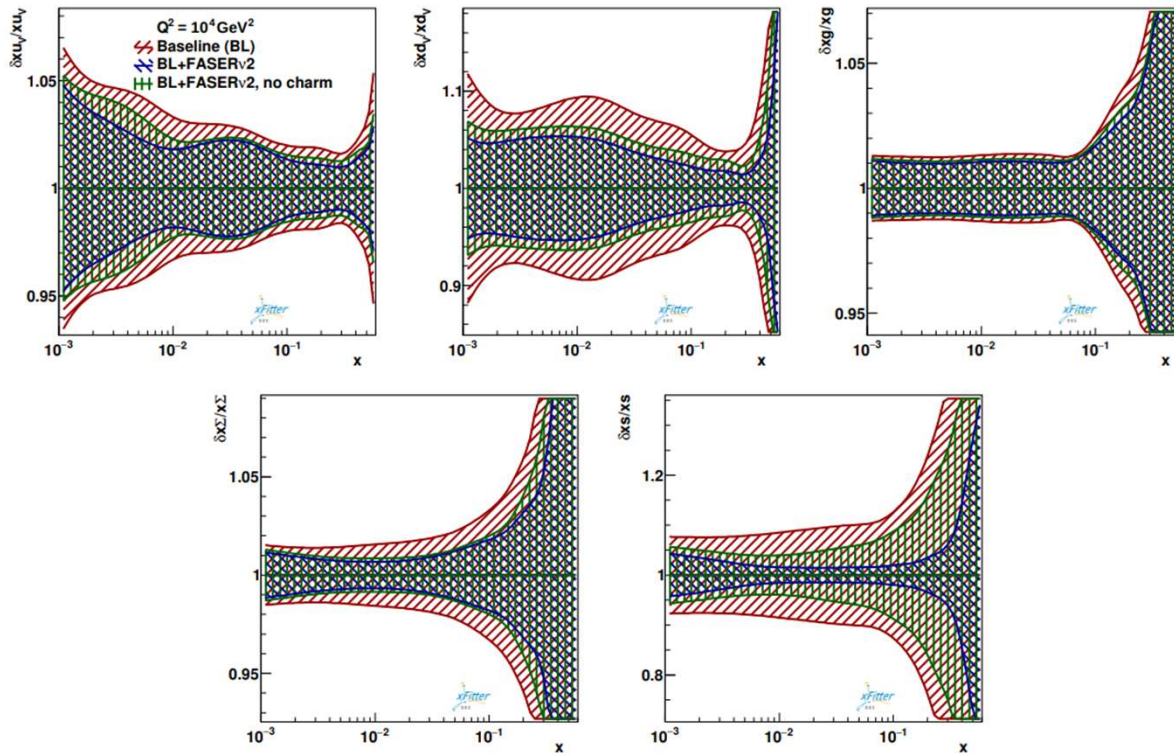
# Backup



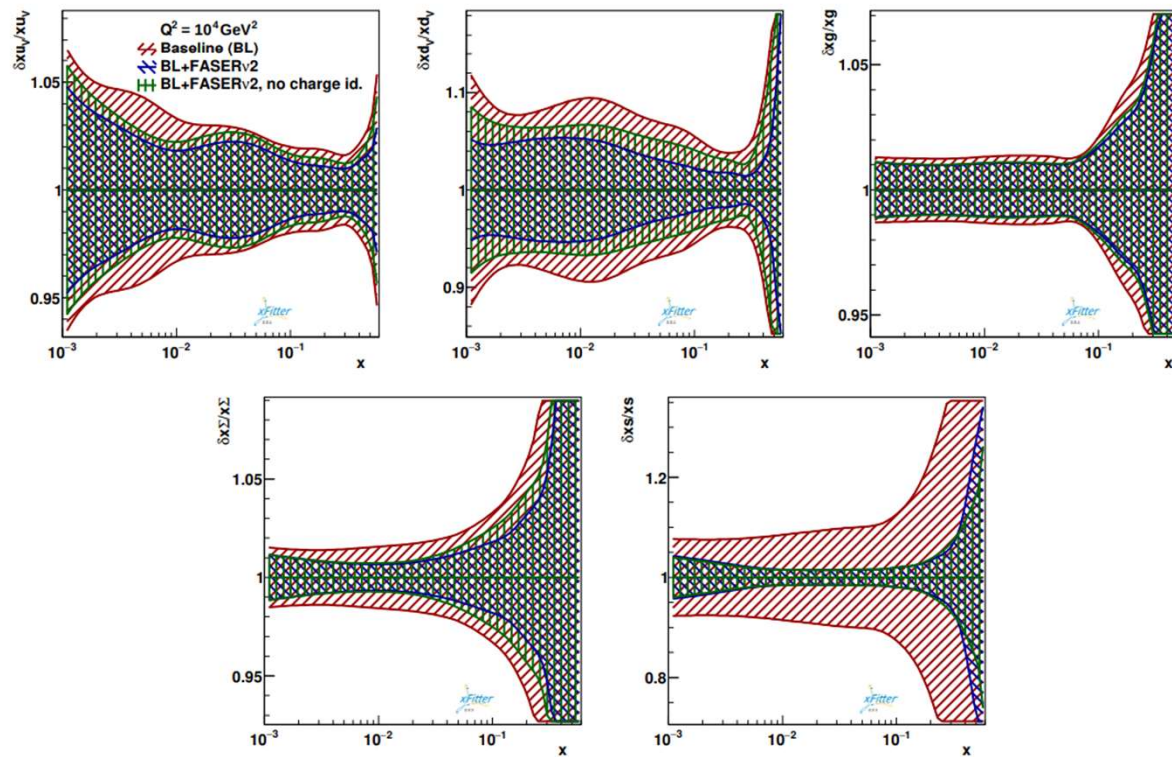
# Faserv2 stat vs sys



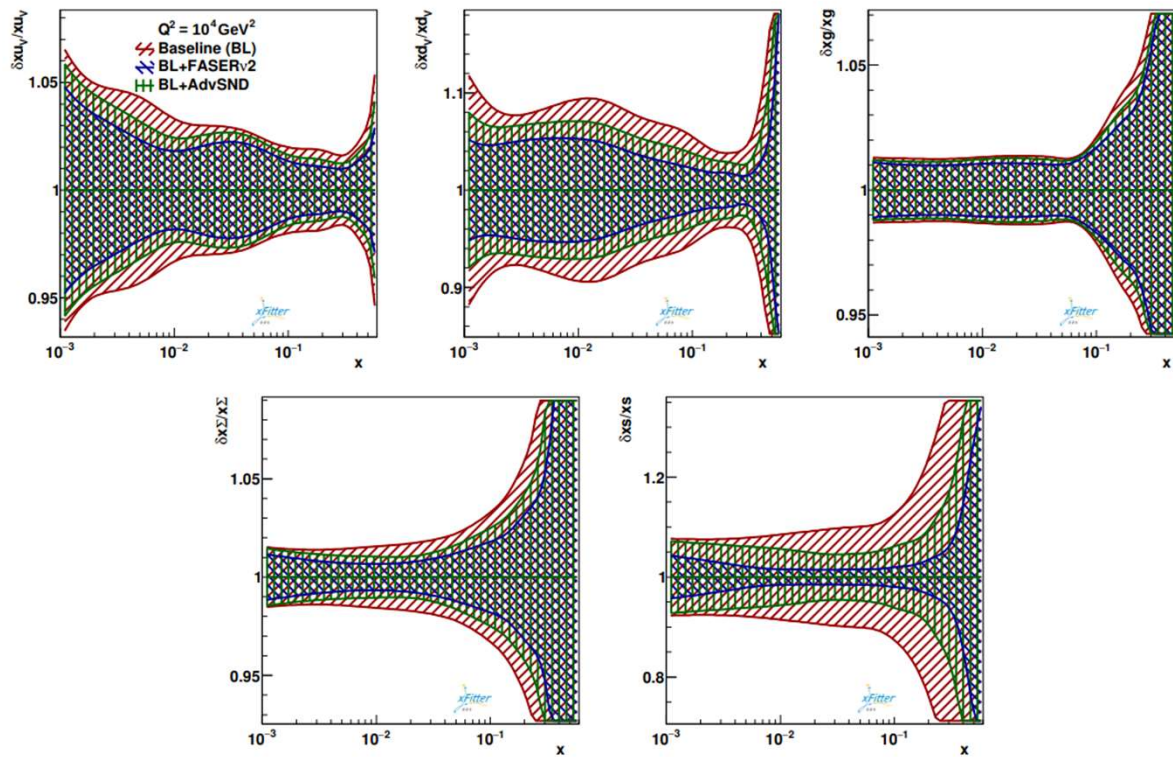
# Faserv2 charm vs no charm



# Faserv2 charge vs no charge

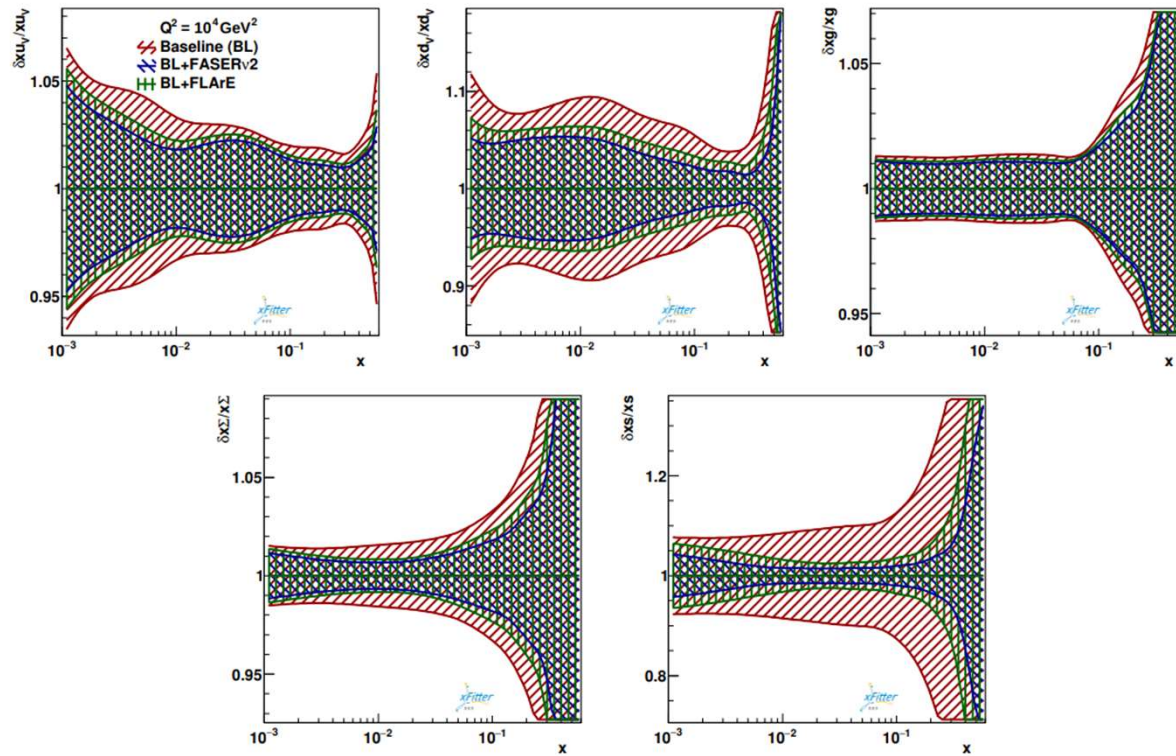


# Faserv2 vs AdvSND

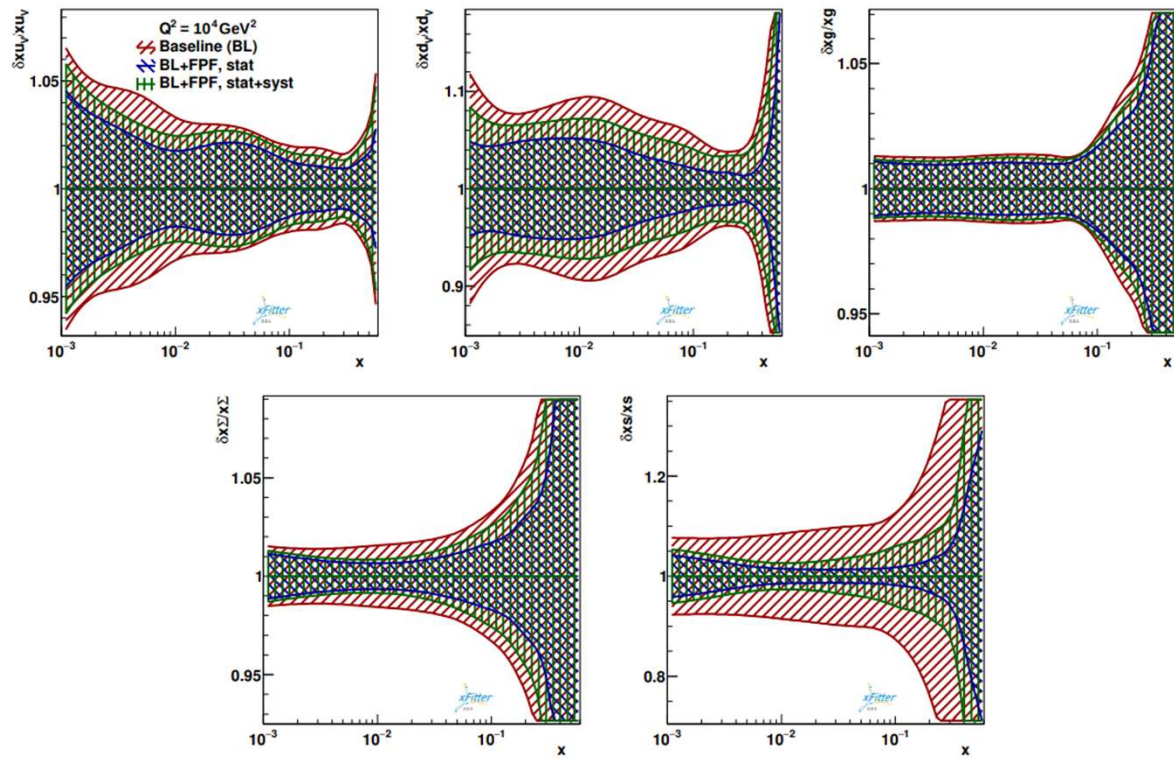




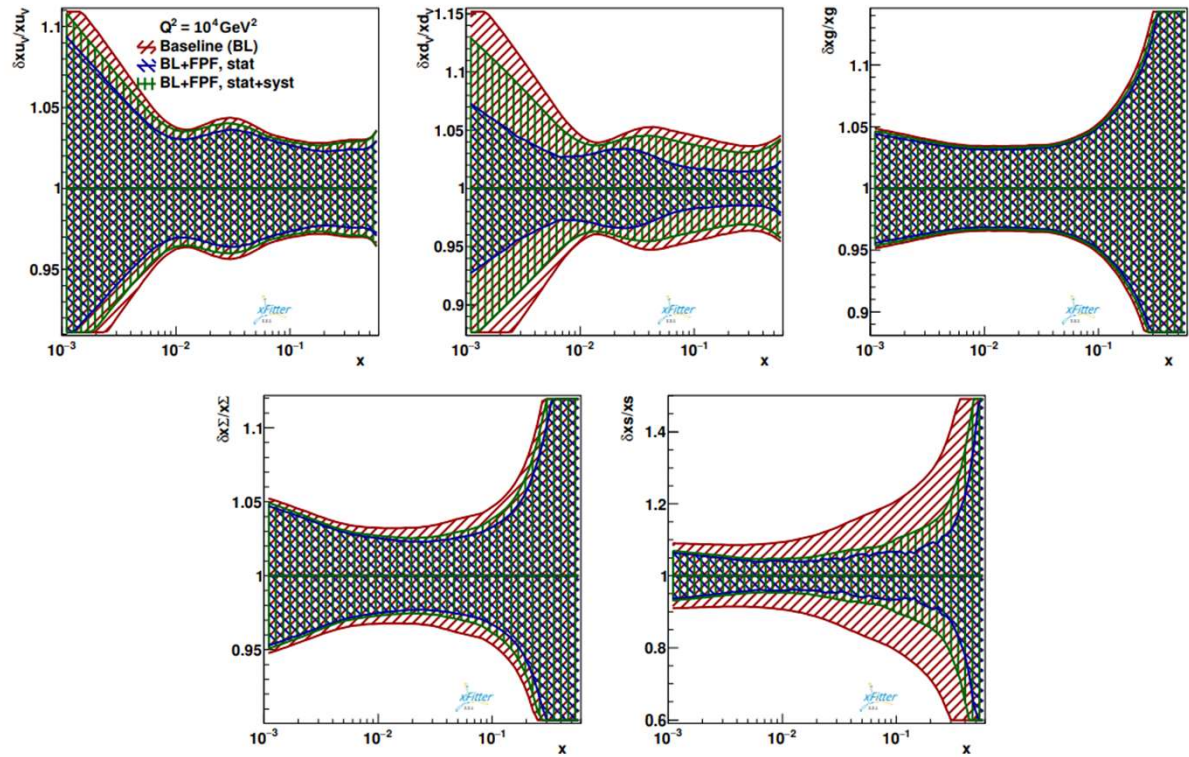
# Faserv2 vs flare



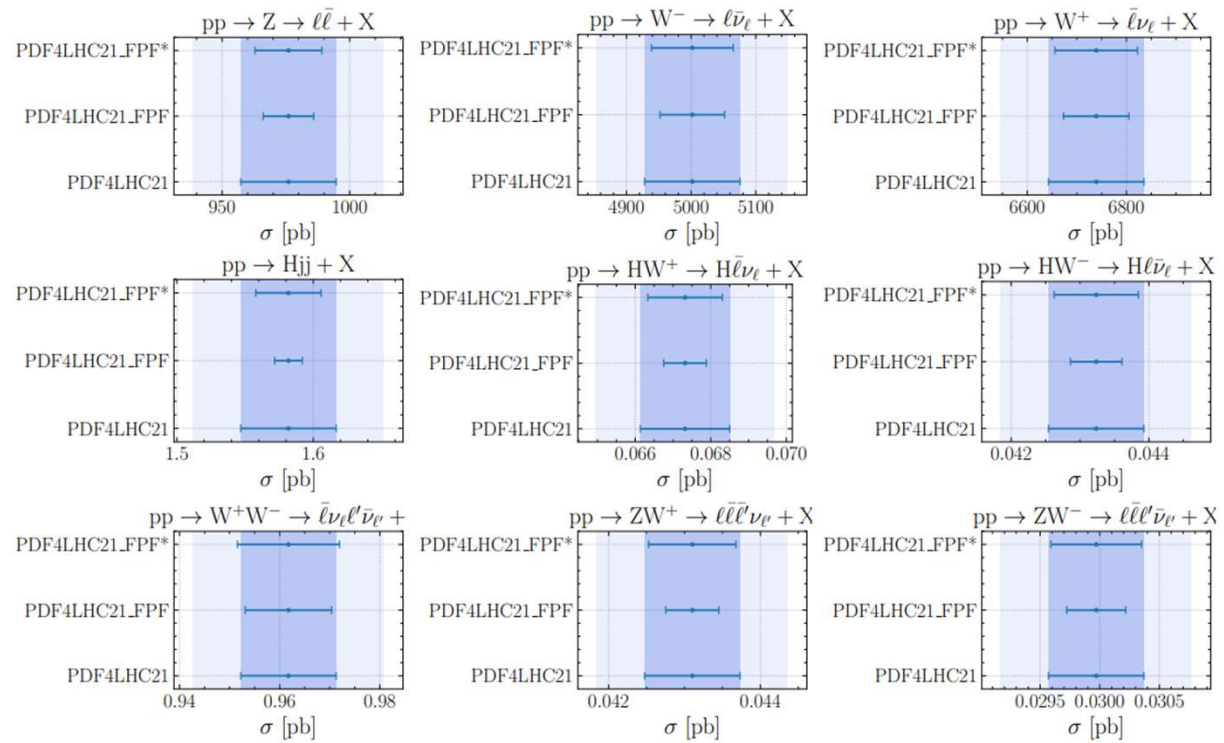
# Fpf total



# Tungsten



# Pheno pdf4lhc21







# PDF fits – FPF

