



Multi-Dimensional Modelling
of
Stellar Interiors

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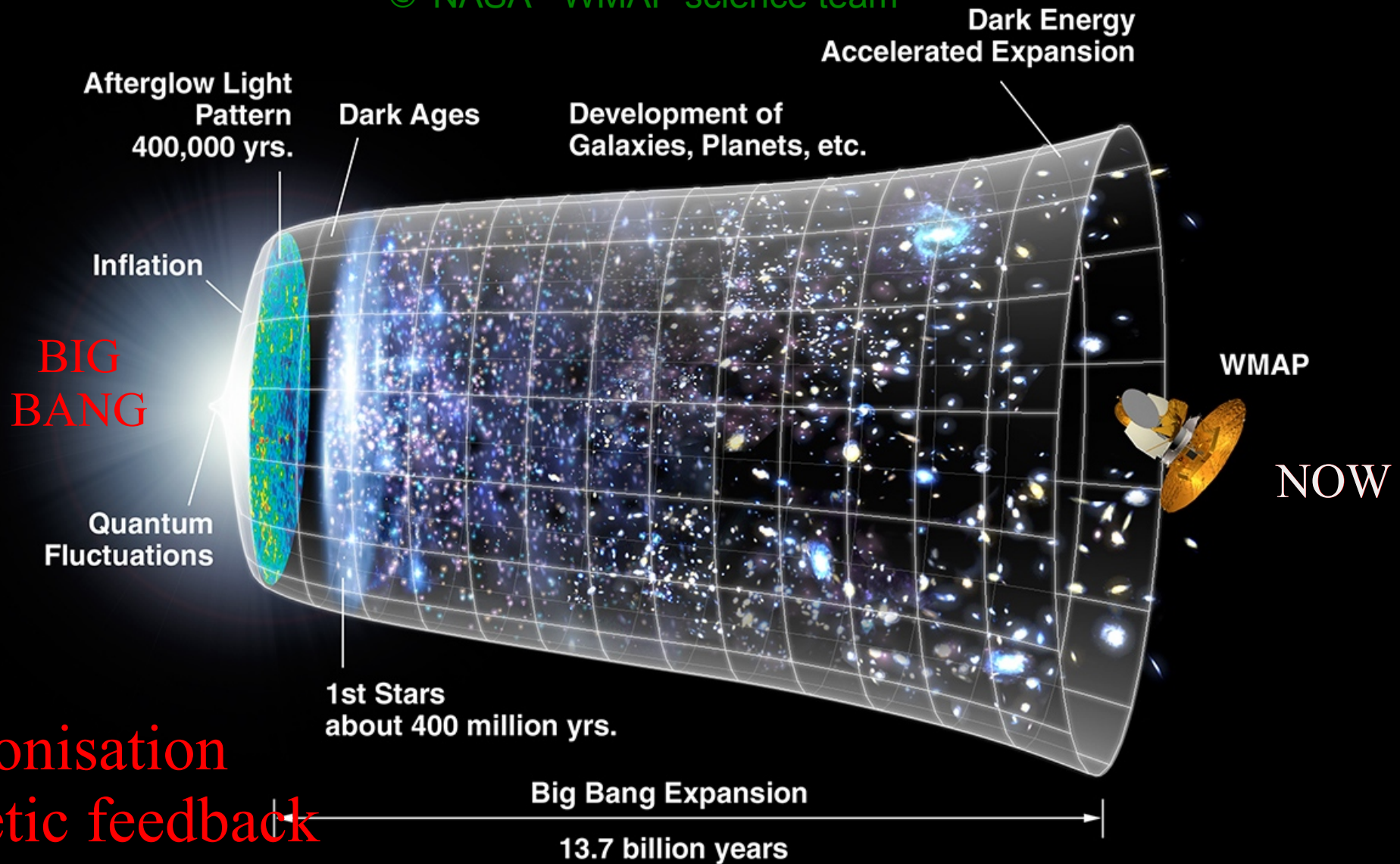


Plan

- Importance, successes & challenges
 - Multi-D models of rotation
 - Multi-D models of convection
 - Conclusions & Outlook
- 

Stars' Role in Universe

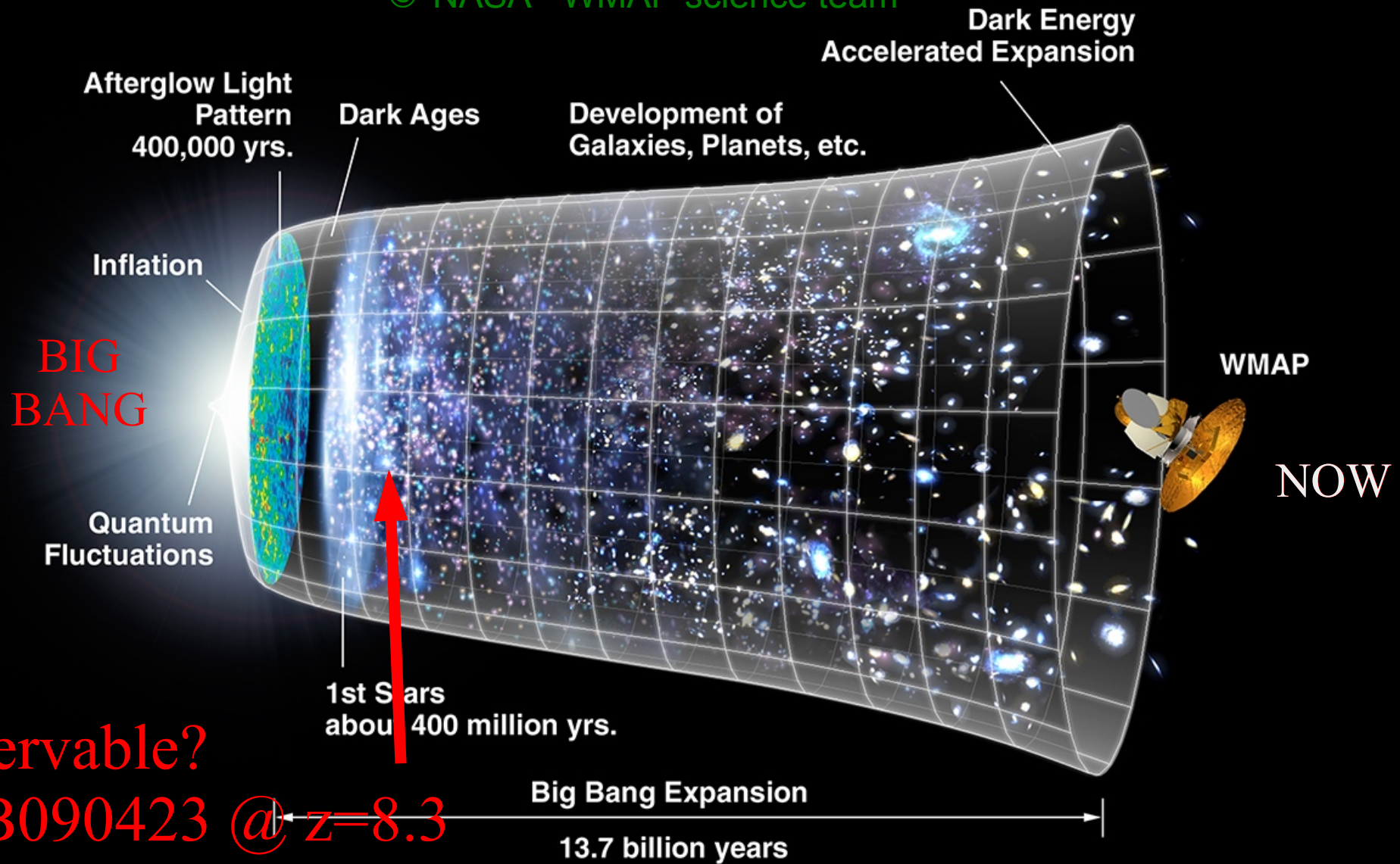
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- Re-ionisation
- Kinetic feedback
- Chemical feedback observed in EMP stars

First Stellar Generations: Importance

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- Observable?

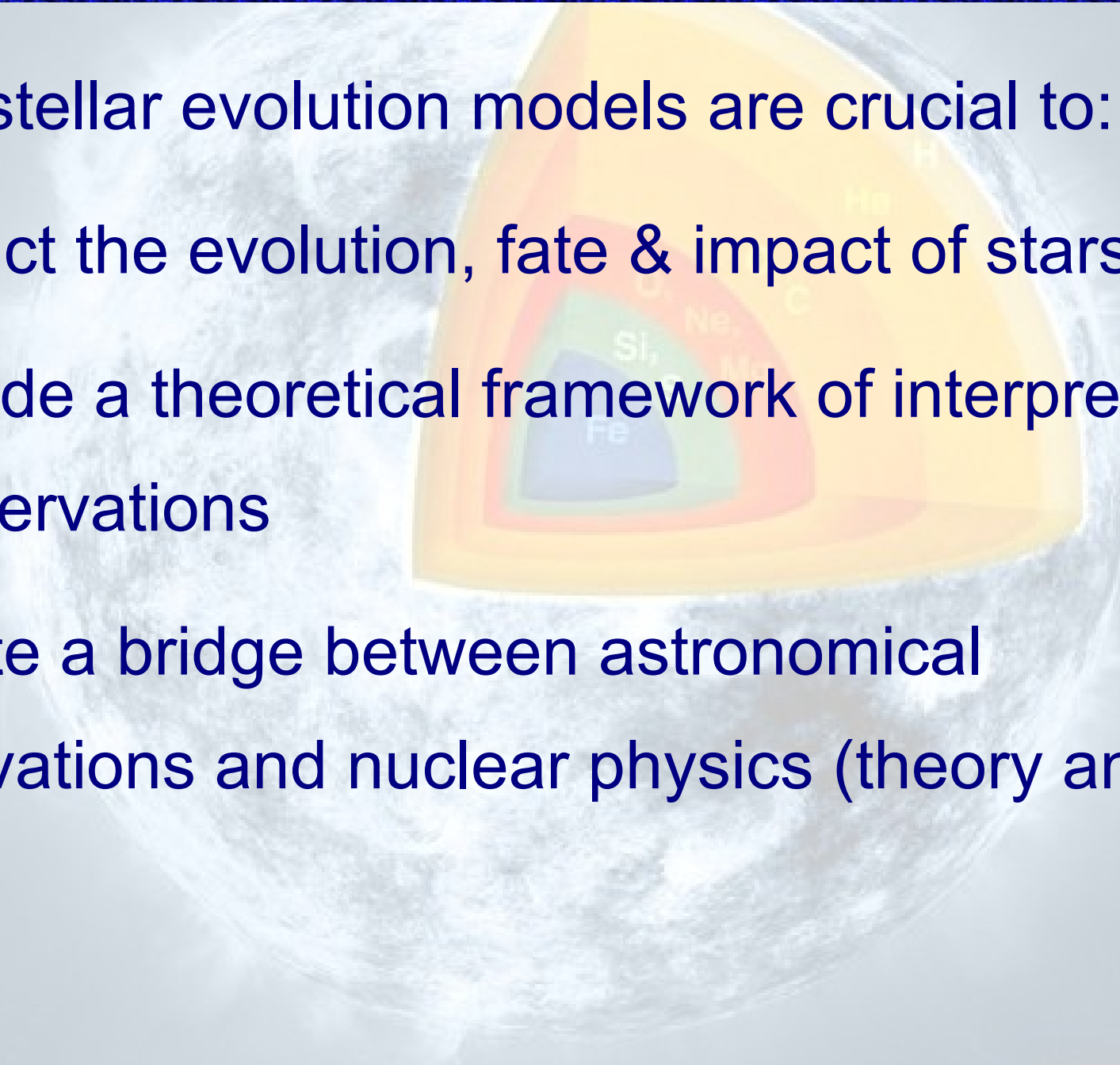
- GRB090423 @ $z=8.3$

Universe age ~ 600 Myr (Tanvir et al 09, Nature: arXiv:0906.1577)

Stellar Evolution Models

Thus stellar evolution models are crucial to:

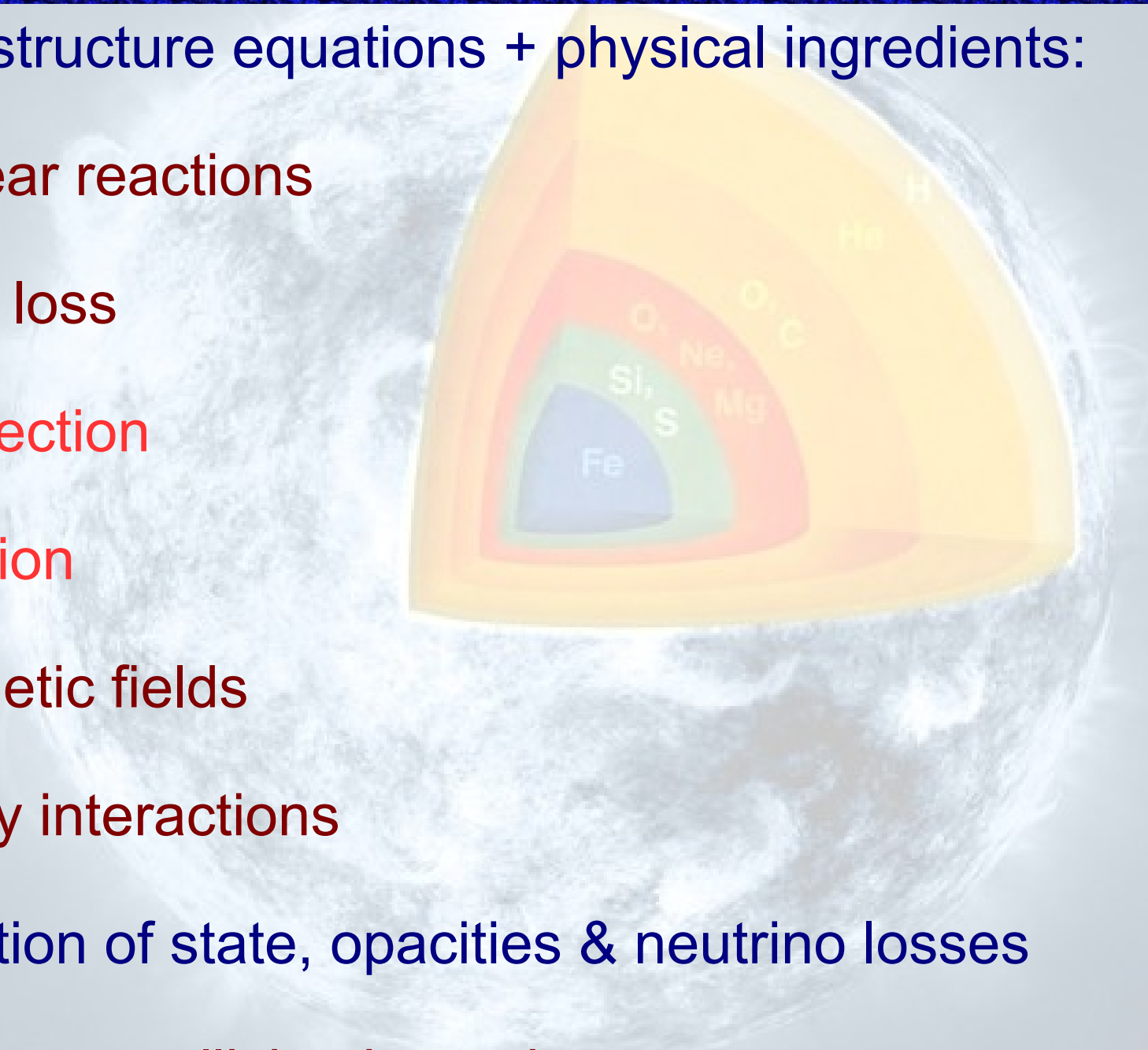
- predict the evolution, fate & impact of stars
- provide a theoretical framework of interpretations of observations
- create a bridge between astronomical observations and nuclear physics (theory and exp.)



1D Stellar Evolution Models

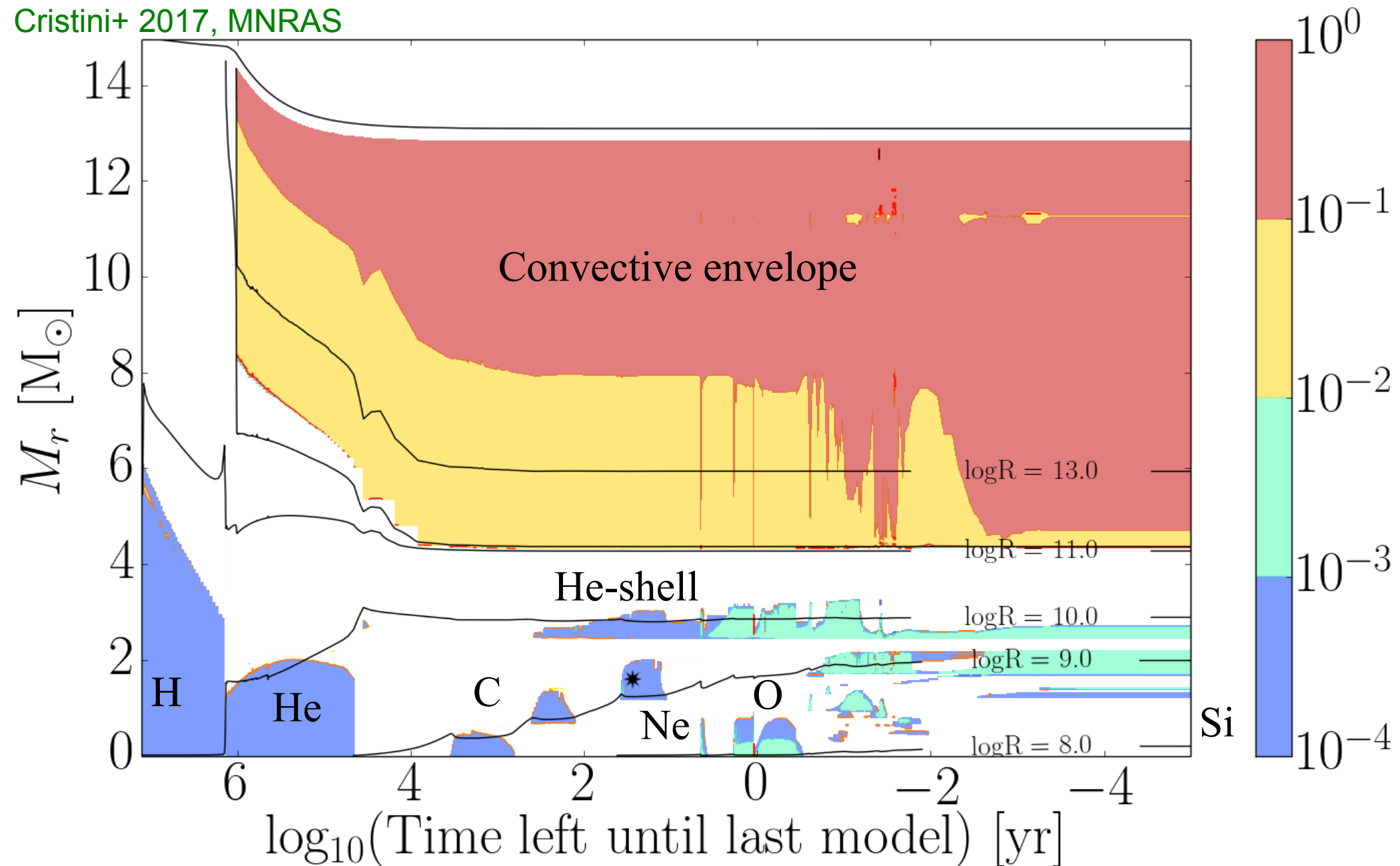
Stellar structure equations + physical ingredients:

- Nuclear reactions
- Mass loss
- Convection
- Rotation
- Magnetic fields
- Binary interactions
- Equation of state, opacities & neutrino losses
including metallicity dependence



Structure Evolution of Massive Stars

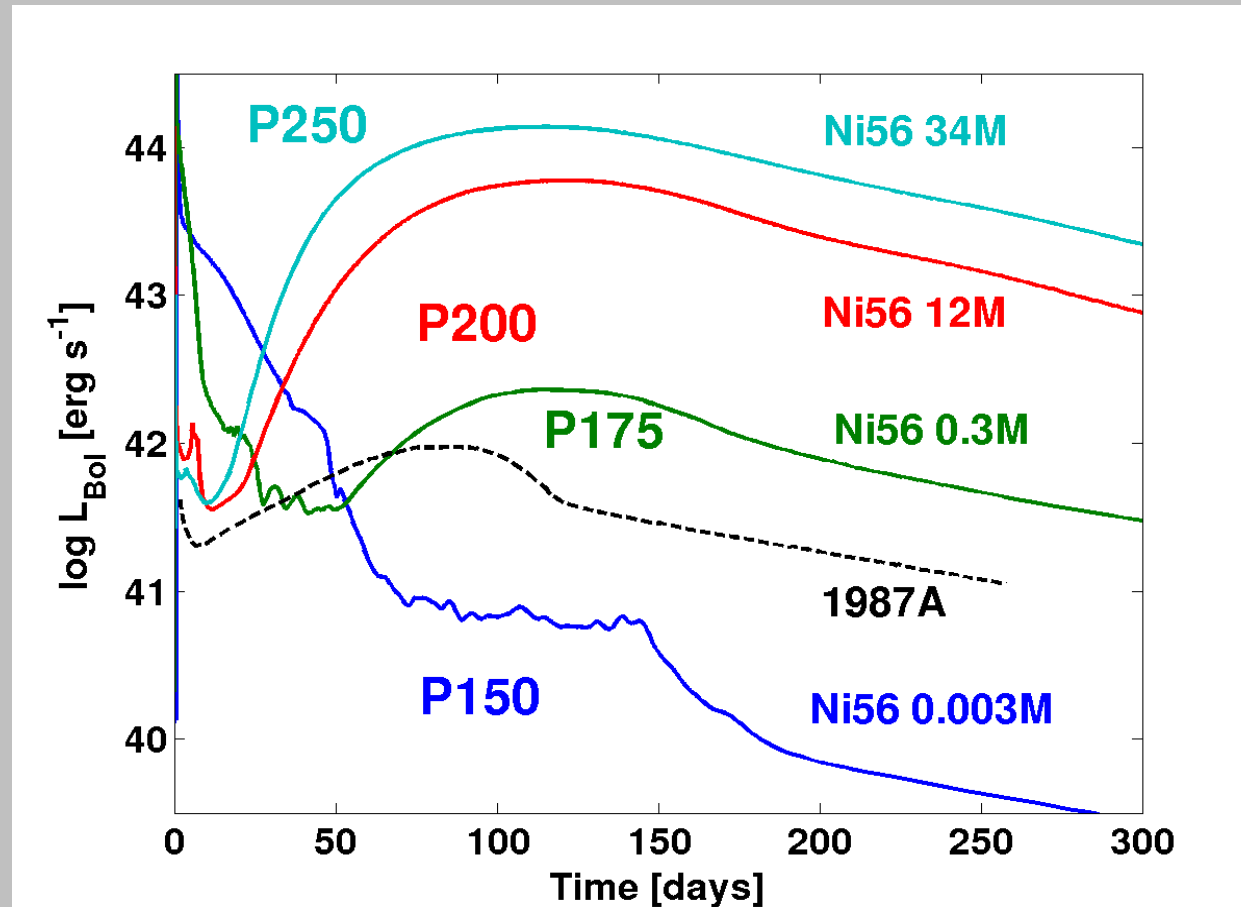
Cristini+ 2017, MNRAS



Convection takes place during most burning stages

PISNe Models: from ZAMS to LCs (at $Z=0.001$)

(Kozyreva+RH+ 2017MNRAS.464.2854K, Gilmer+RH+ 2017ApJ.846.100, ArXiv170607454G)

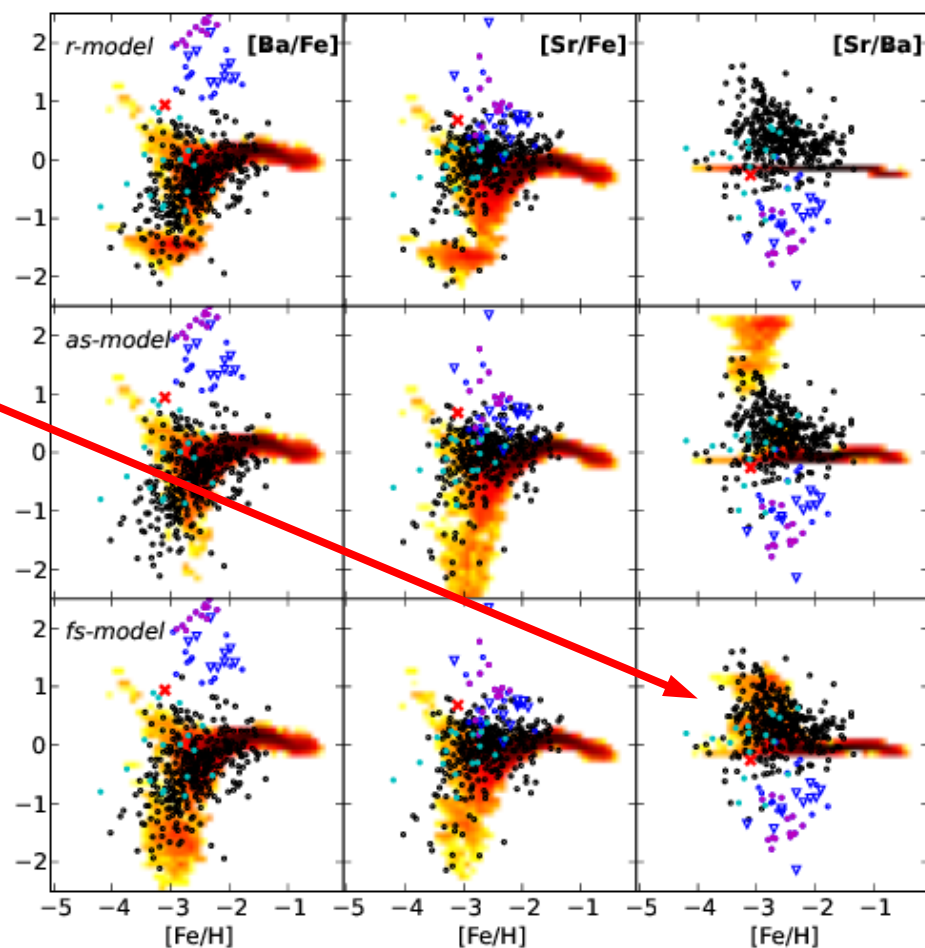


Some (many?) PISNe and most PPISNe are faint!!

Successes: e.g. [Sr/Ba] in EMP stars

Strong variations in $[\text{Sr}/\text{Ba}] > 1$ matches well observed range for EMP stars (black circles)!

(no main s-process included so cannot explain CEMP-s stars in blue)



Cescutti *et al.* (2013),
see also 2015

Model name	panels in Fig. 5	s-process	r-process
r-	Upper	No s-process from massive stars	standard + extended r-process site (8 - 30 M_{\odot})
as-	middle	average rotators ($v_{\text{rot}}/v_{\text{critic}} = 0.4$)	standard r-process site (8 - 10 M_{\odot})
fs-	lower	fast rotators ($v_{\text{rot}}/v_{\text{critic}} = 0.5$) and 1/10 for $^{17}\text{O}(\alpha, \gamma)$ reaction rate	standard r-process site (8 - 10 M_{\odot})

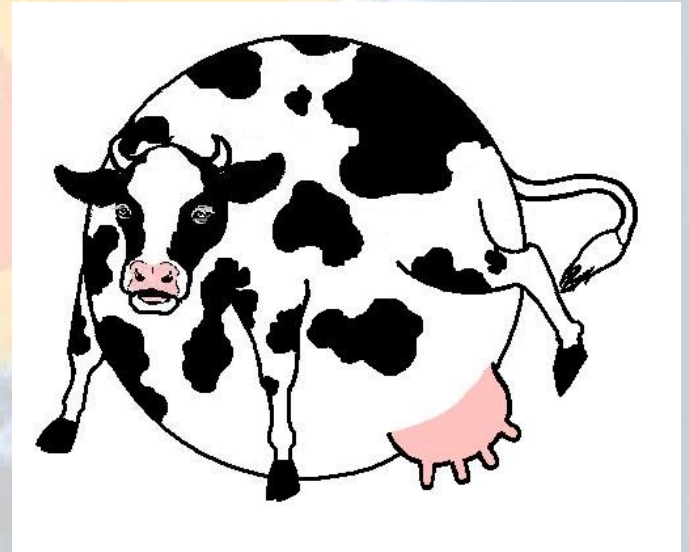
Advantages:

- model entire evolution ($\Delta t \sim 10^3$ yrs)
- compare to observations
- progenitor models
- large grids (M, Z)

Disadvantages:

- parametrized physics (e.g. convection)
- missing multi-D processes
- incapable of modelling turbulence

1D stellar models



What's missing?

- self-consistent physical descriptions of mass loss, **convection**, **rotation**, magnetic fields, opacity, binarity **and their interplay**

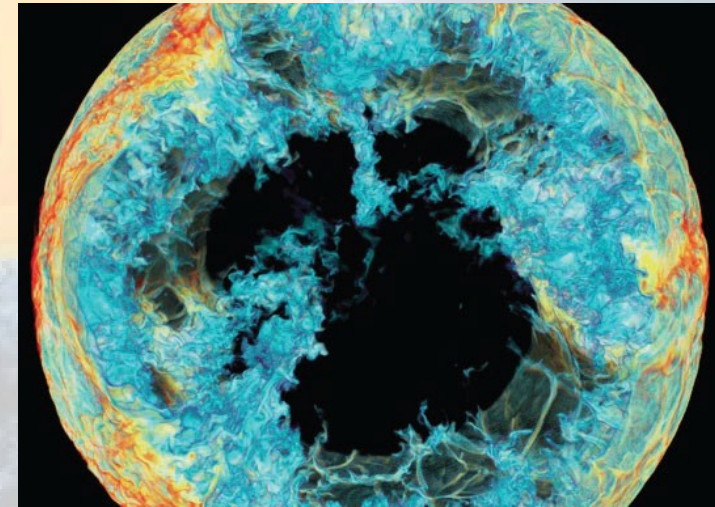
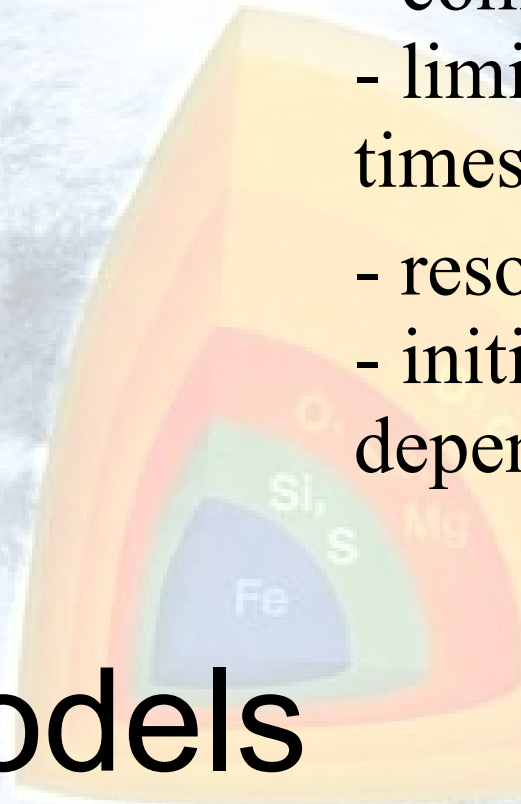
Advantages:

- model fluid instabilities (e.g. Rayleigh-Taylor)
- modeling 3D processes
- model diffusive and advective processes

Disadvantages:

- computational cost
- limited to dynamical timescales ($t_{\text{conv}} \sim 1\text{s} - \text{days}$)
- resolution dependent?
- initial condition dependent?

3D stellar models



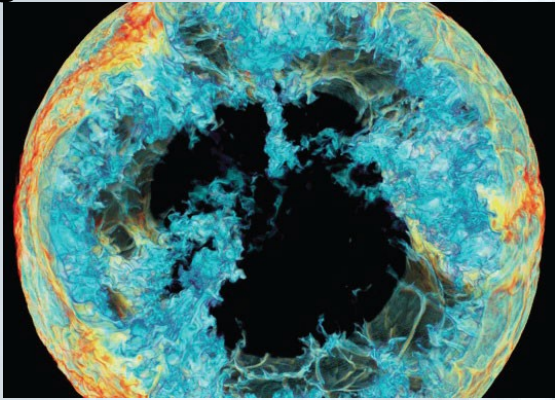
Herwig, Woodward et al 2013

What's missing?

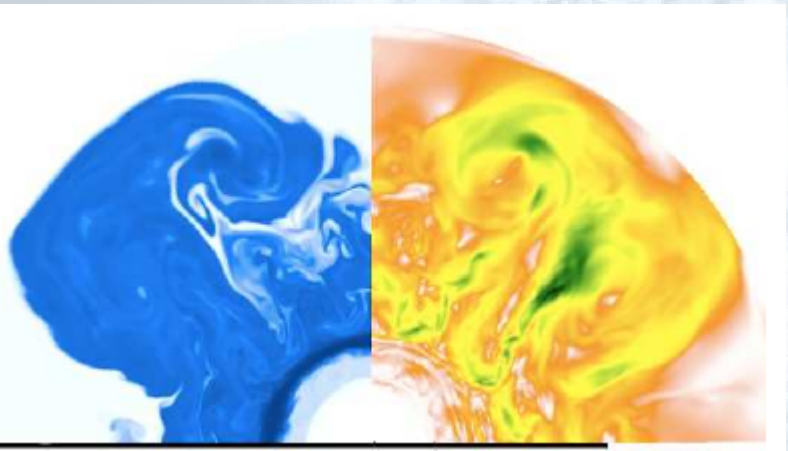
- full star or lifetime simulations
- Large scale (LES) and small scale (DNS) cannot be followed simultaneously

Way Forward: 1 to 3 to 1D link

Targetted 3D simulations

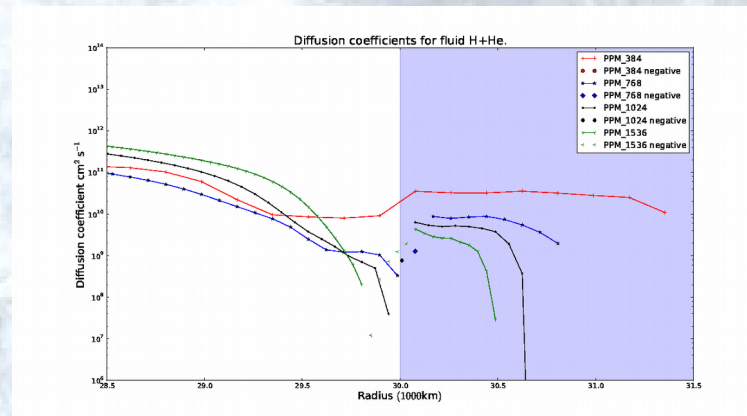
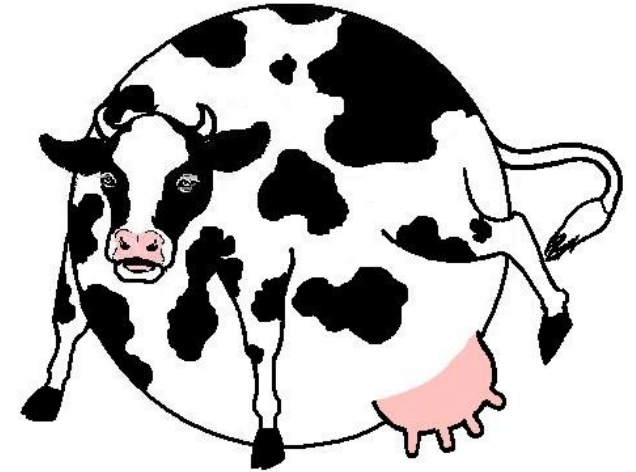


Herwig et al 06, Herwig, Woodward et al 2013



e.g. Arnett & Meakin 2011, ...
Mocak et al 2011,
Viallet et al 2013, ...

Uncertainties in 1D



Meakin et al 09 ; Bennett et al (thesis), Jones et al 16

→ Determine effective coefficient / improve theoretical prescriptions

2.3. Dynamical shear

The criterion for stability against dynamical shear instability is the Richardson criterion:

$$Ri = \frac{N^2}{(\partial U / \partial z)^2} > \frac{1}{4} = Ri_c, \quad (1)$$

Hirschi et al 2004

where U is the horizontal velocity, z the vertical coordinate and N^2 the Brunt-Väisälä frequency:

$$N^2 = \frac{g\delta}{H_P} [\nabla_{ad} - \nabla + \frac{\varphi}{\delta} \nabla_{\mu}] \quad (2)$$

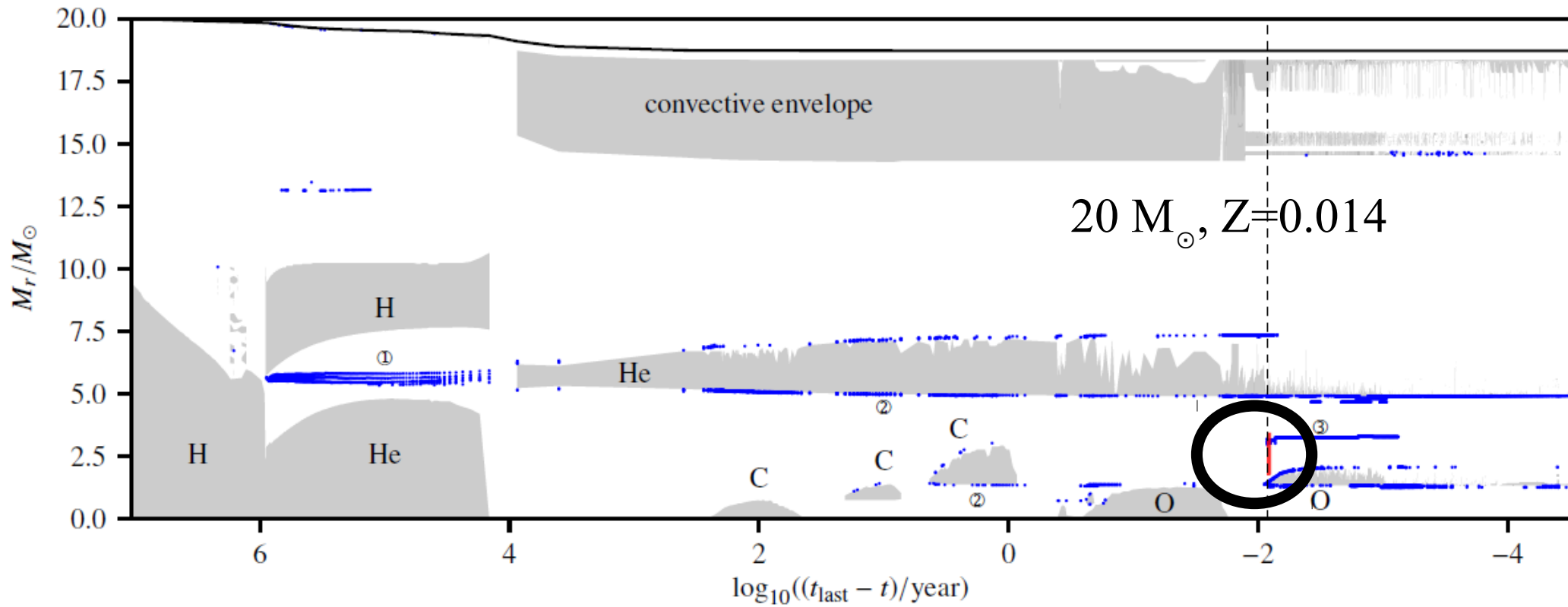
2.3.1. The recipe

The following dynamical shear coefficient is used, as suggested by J.-P. Zahn (priv. comm.):

$$D = \frac{1}{3} v l = \frac{1}{3} \frac{v}{l} l^2 = \frac{1}{3} r \frac{d\Omega}{dr} \Delta r^2 = \frac{1}{3} r \Delta\Omega \Delta r \quad (5)$$

where r is the mean radius of the zone where the instability occurs, $\Delta\Omega$ is the variation of Ω over this zone and Δr is the extent of the zone. The zone is the reunion of

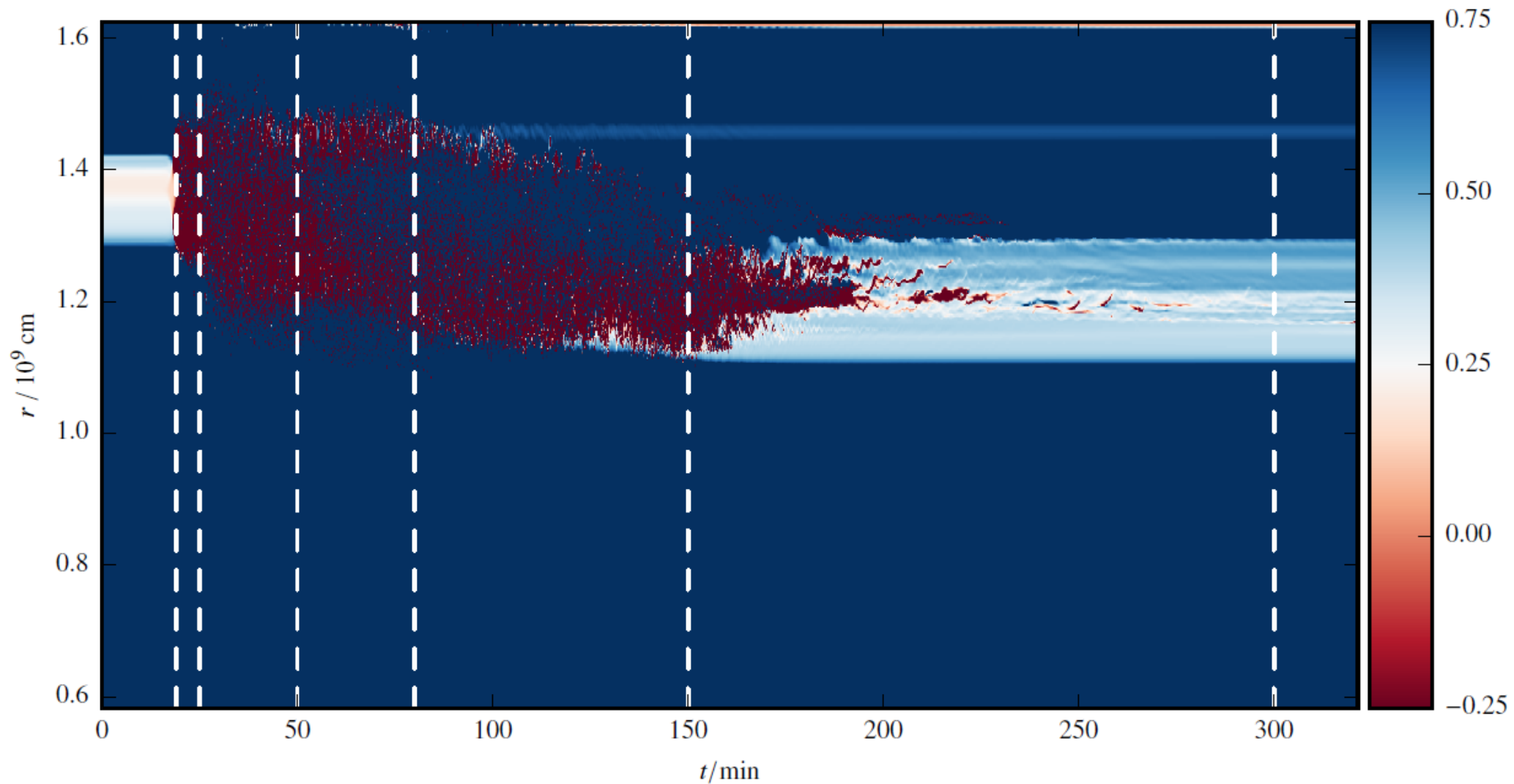
Dynamical Shear



Edelmann et al 2016, A&A

2D simulations with SLH

Dynamical Shear: Ri Time Evolution in 2D



- Threshold (Ri_c) confirmed.
- Mixing/transport rate is tricky to implement correctly in 1D codes

Convection: Current Implementation in 1D Codes

Multi-D processes:

Major contributor to turbulent mixing

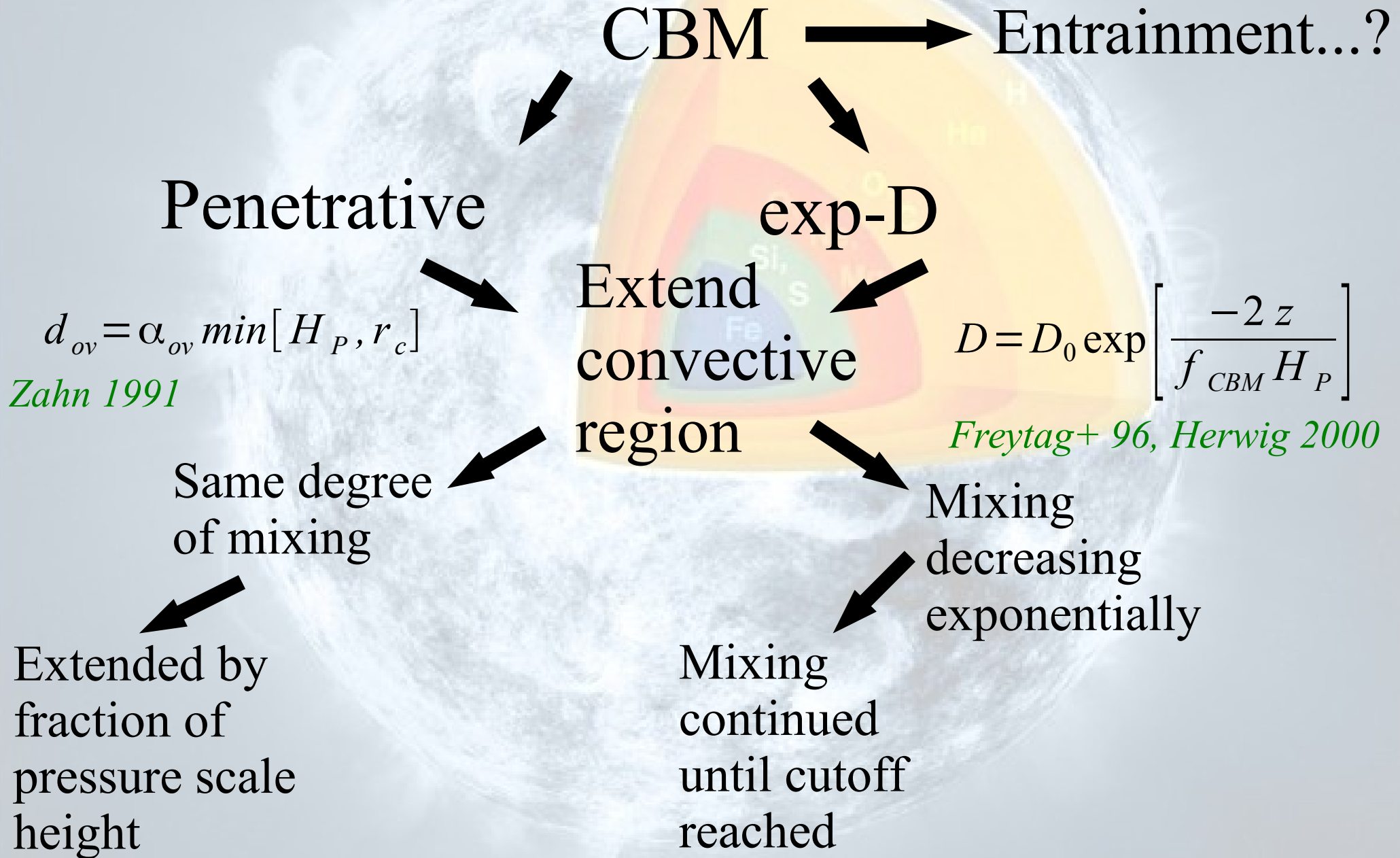
Turbulent entrainment at convective boundaries

Internal gravity waves

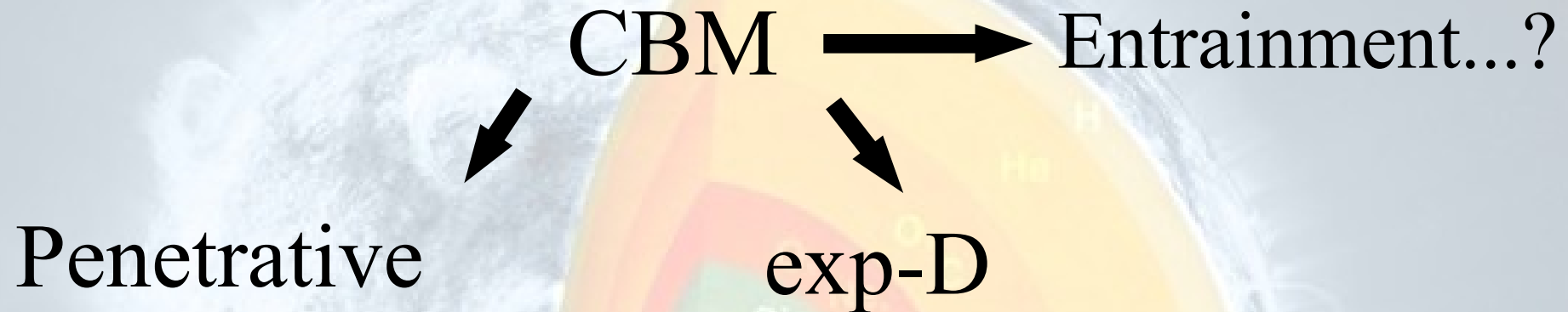
1D prescriptions:

- Energy transport in convective zone: **mixing length theory (MLT)** *Bohm-Vitense (1957,58)*, or updates, e.g. FST: *Canuto & Mazitelli (1991)*
- Boundary location: Schwarzschild criterion OR Ledoux (+semi-convection)
- Convective boundary mixing (CBM, also composition dependent)

Convective Boundary Mixing (CBM)



Convective Boundary Mixing (CBM)



- Inspired by theory & multi-D hydro simulations
- More simulations and **their interpretation in a theoretical framework** will help improve these prescriptions

Viallet, Meakin, Prat, Arnett 2015, Arnett et al 2015

3D Hydro Efforts/Priority List

* Convective boundary mixing during core hydrogen burning:

- +: many constraints (HRD, astero, ...)
- -: difficult to model due to important thermal/radiative effects
- -: long time-scale

* Silicon burning:

- +: important to determine impact on SNe of multi-D structure in progenitor (Couch et al 2015a,b, Mueller & Janka [aph1409.4783](#), Mueller et al [ArXiv1605.01393](#))
- +: possible shell mergers occurring after core Si-burning (e.g. Tur et al 2009ApJ702.1068; Sukhbold & Woosley 2014ApJ783.105) strongly affect core compactness
- +: radiative effects small/negl.
- -: $\sim 10^9$ CPU hours needed for full silicon burning phase will be ok soon;
- -: might be affected by convective shell history

* AGB thermal pulses/H-ingestion:

- +: already doable (e.g. Herwig et al 2014ApJ729.3, 2011ApJ727.89, Mocak et al 2010A&A520.114, Woodward et al 2015)
- +: thermal/radiative effects not dominant
- ?: applicable to other phases?

* Oxygen shell: (Meakin & Arnett 2007ApJ667.448/665.448, Viallet et al 2013ApJ769.1, Jones et al [ArXiv1605.03766](#))

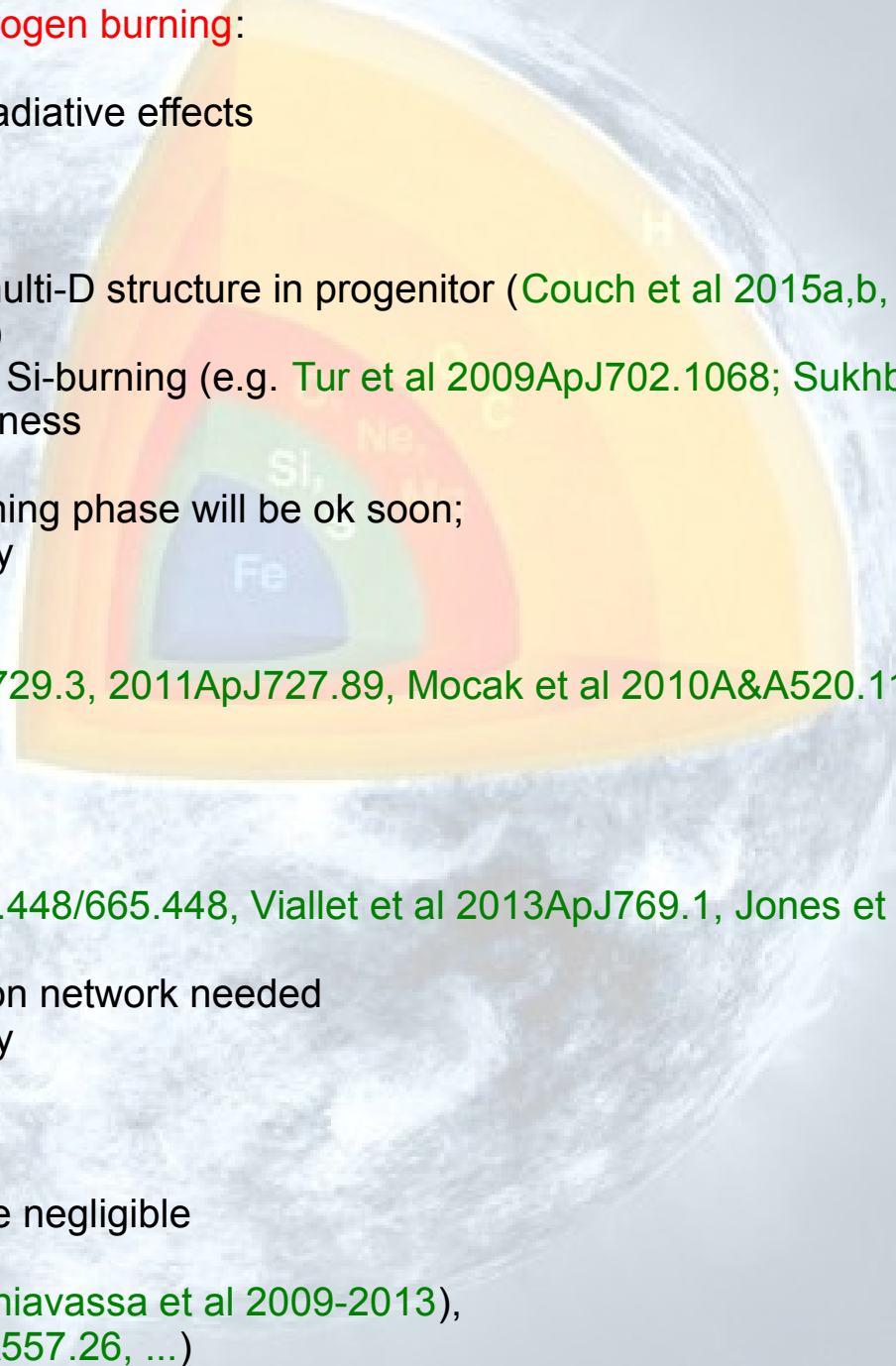
- +: similar to silicon burning but smaller reaction network needed
- -: might be affected by convective shell history

* Carbon shell: (PhD A. Cristini)

- +: not affected by prior shell history
- +: first stage for which thermal effects become negligible

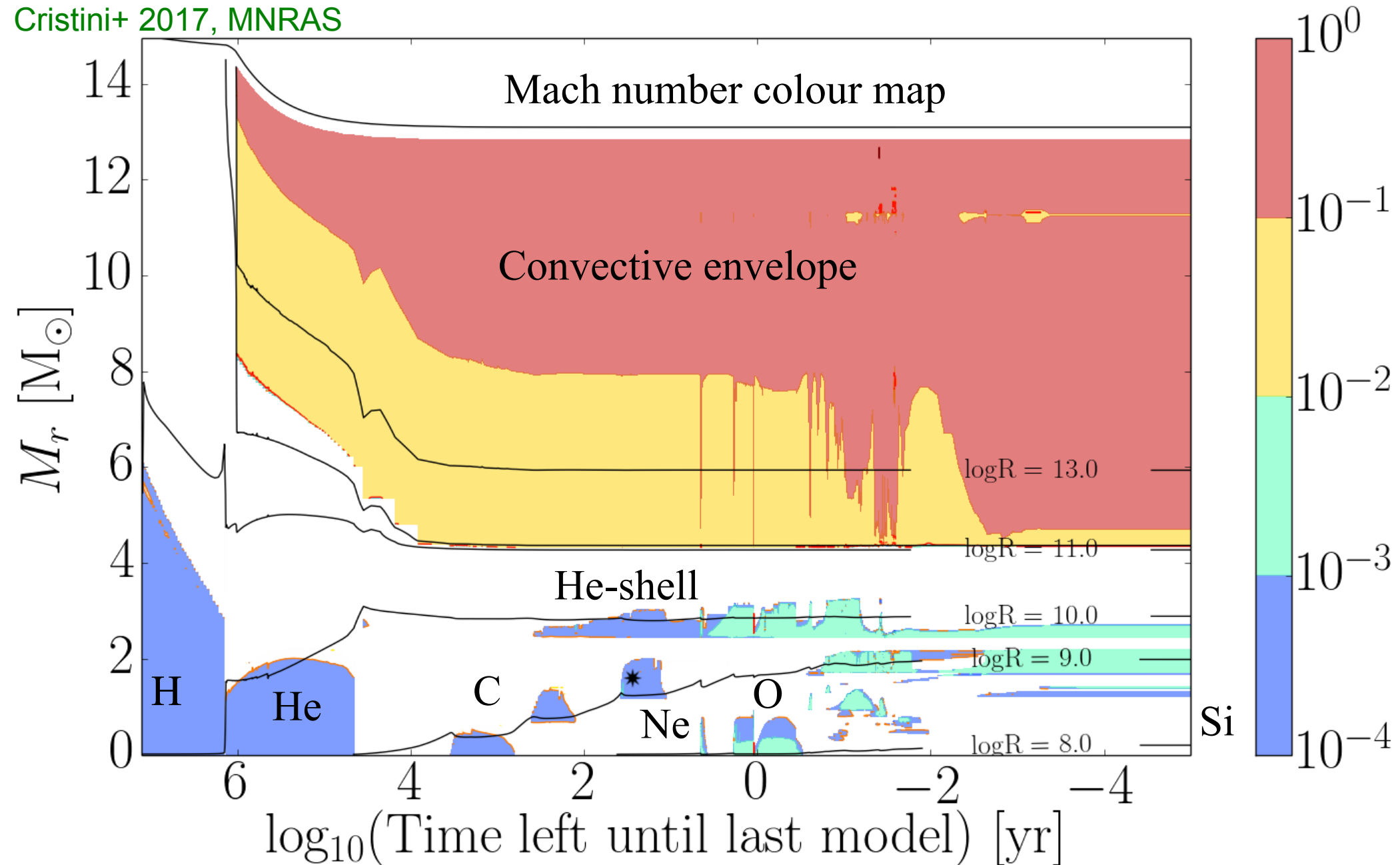
* Envelope of RSG (e.g. Viallet et al. 2013, Chiavassa et al 2009-2013),

* Solar-type stars (e.g. Magic et al. 2013A&A557.26, ...)



Where to Start?

Cristini+ 2017, MNRAS



Convection takes place during most burning stages

Plan

Monthly Notices

of the

ROYAL ASTRONOMICAL SOCIETY



MNRAS **471**, 279–300 (2017)
Advance Access publication 2017 June 20

doi:10.1093/mnras/stx1535

3D hydrodynamic simulations of carbon burning in massive stars

A. Cristini,^{1★} C. Meakin,^{2,3} R. Hirschi,^{1,4★} D. Arnett,^{3★} C. Georgy,^{1,5}
M. Viallet⁶ and I. Walkington¹

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MNRAS **484**, 4645–4664 (2019)
Advance Access publication 2019 February 1

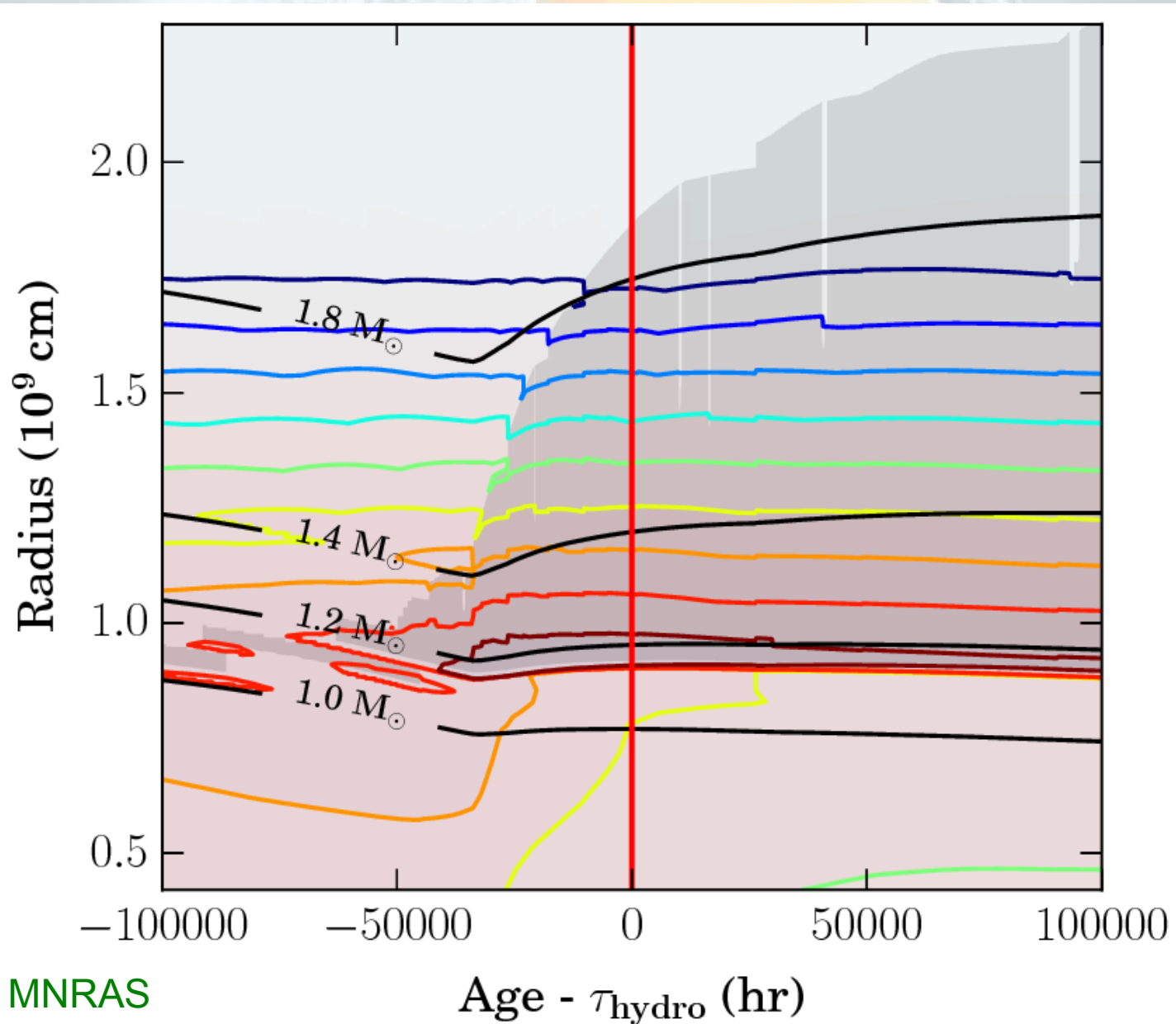
doi:10.1093/mnras/stz312

Dependence of convective boundary mixing on boundary properties and turbulence strength

A. Cristini ,^{1★} R. Hirschi,^{1,2★} C. Meakin,^{3,4} D. Arnett,³ C. Georgy^{1,5} and
I. Walkington¹

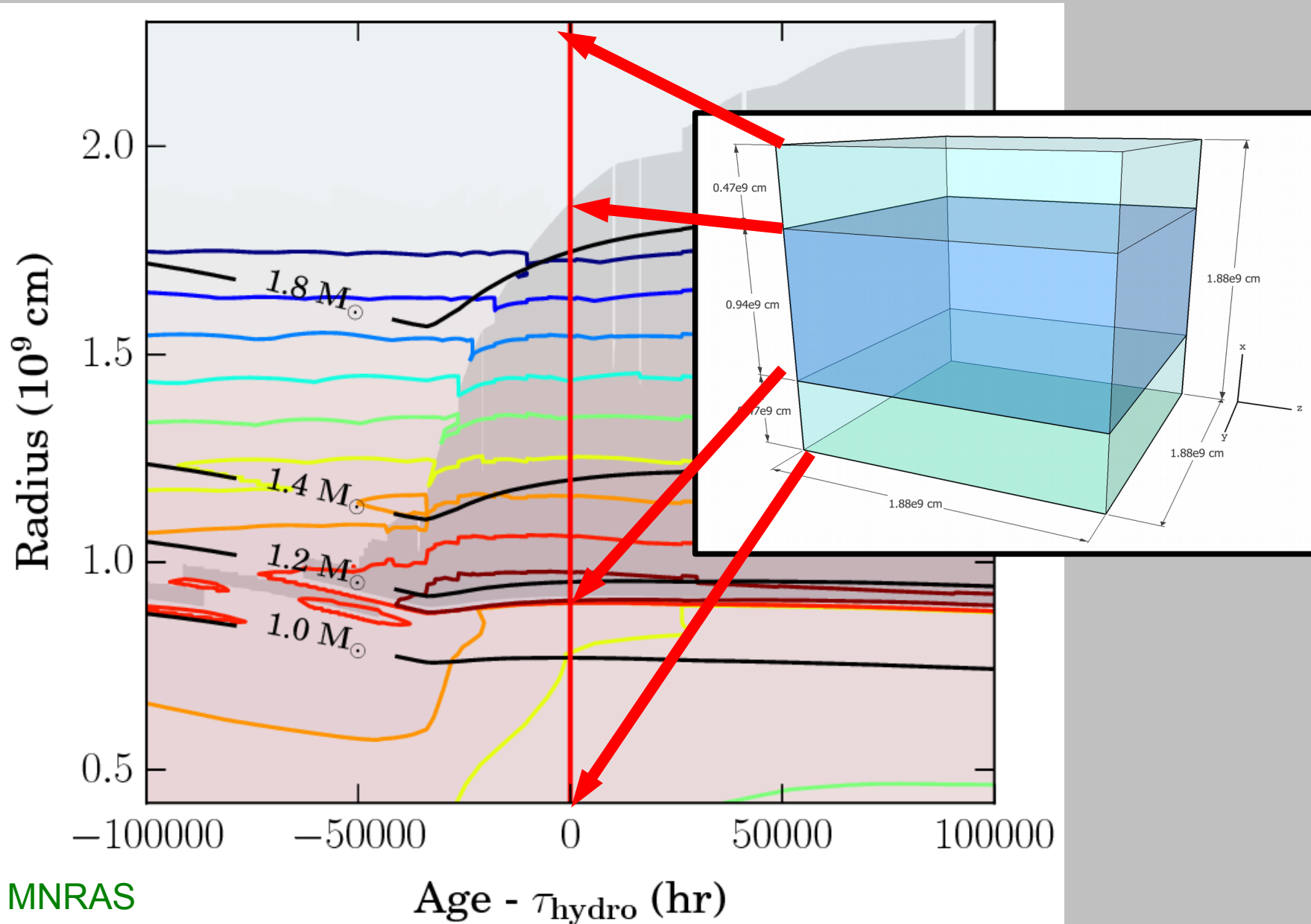
From 1D to 3D

C-shell in $15 M_{\odot}$, $Z=0.014$ 1D stellar evolution model

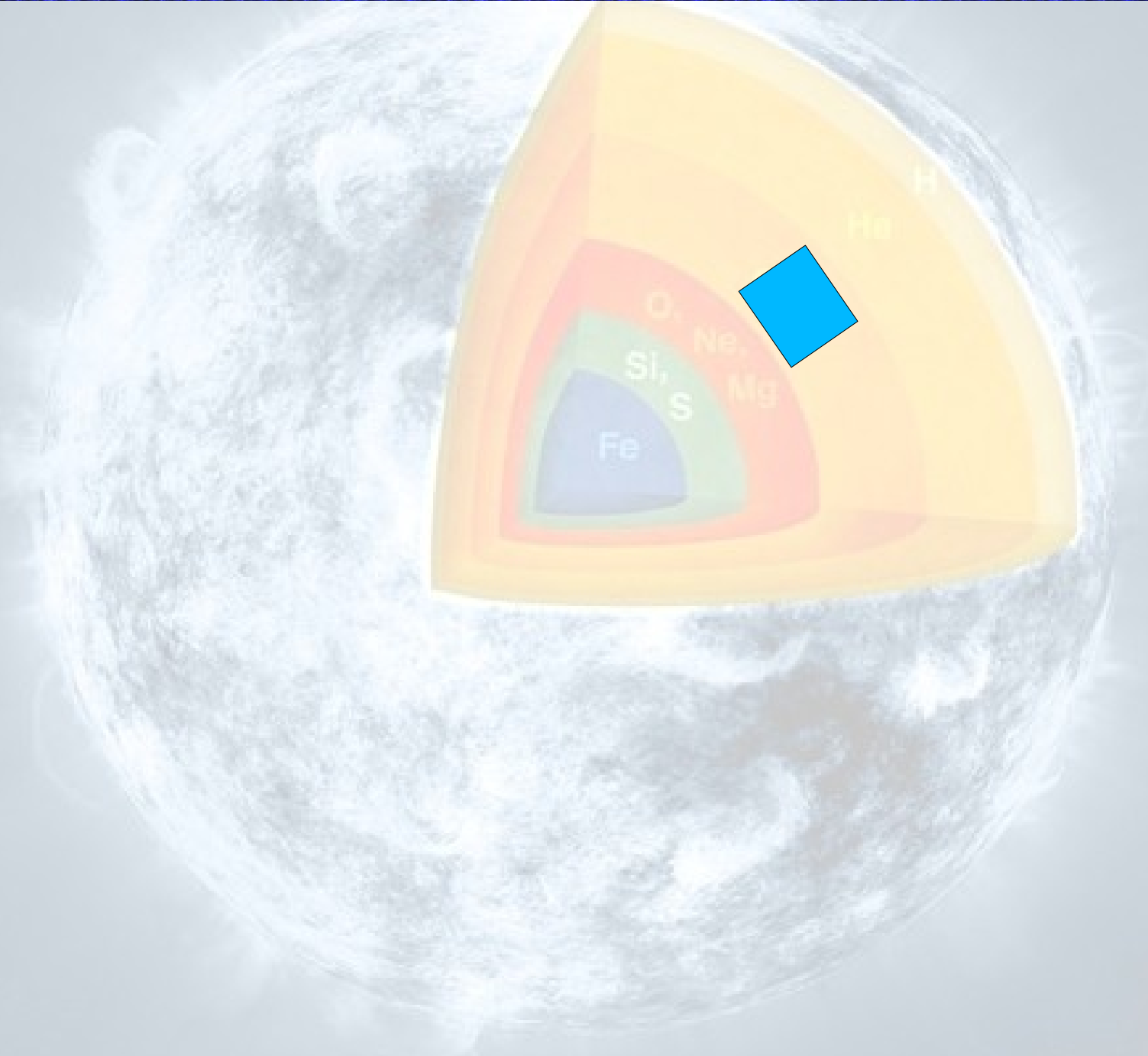


From 1D to 3D

C-shell in $15 M_{\odot}$, $Z=0.014$ 1D stellar evolution model



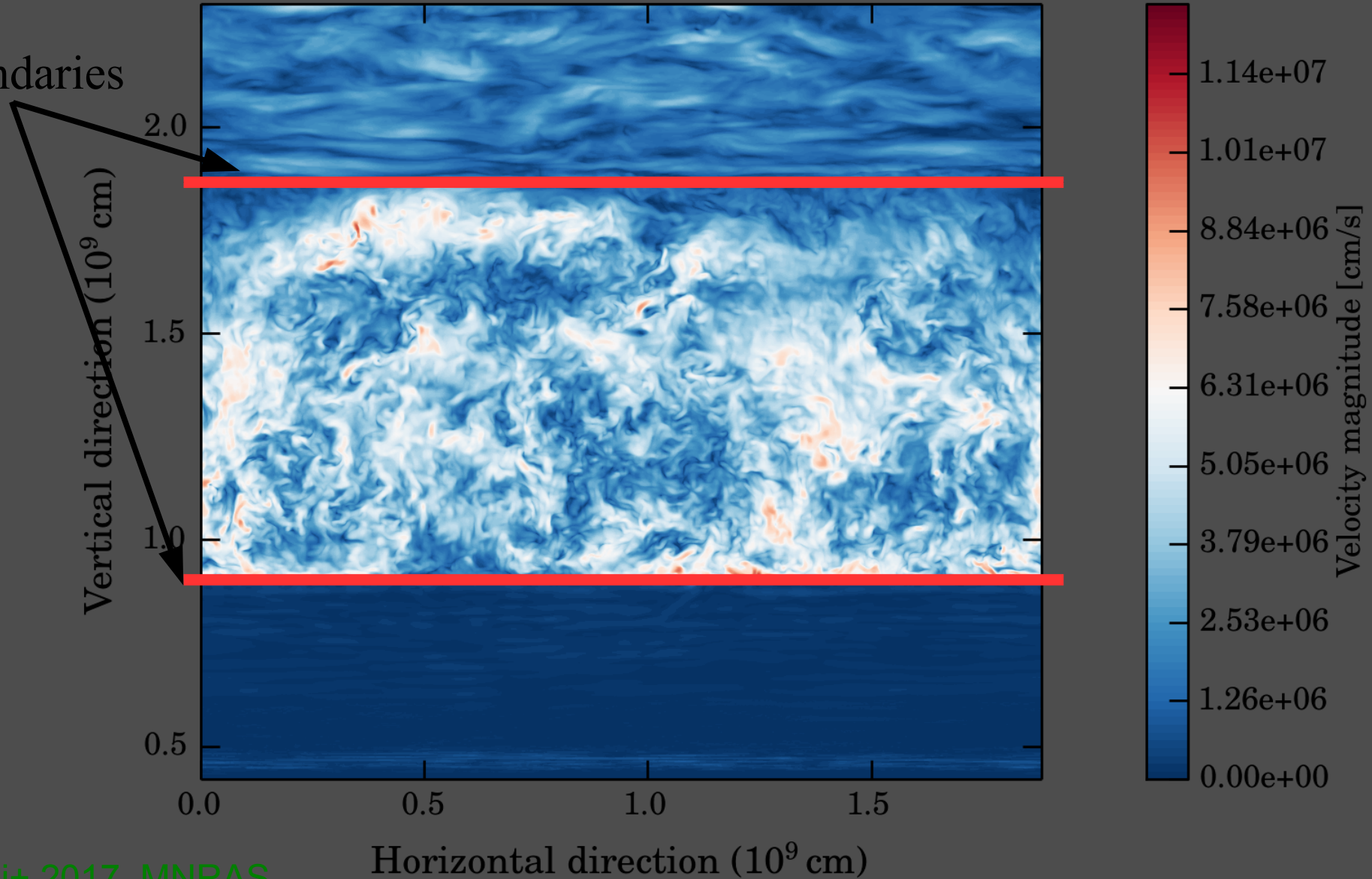
3D simulations of Carbon burning



3D C-shell Simulations

Snapshot from 1024^3 resolution run: Gas Velocity $\|v\|$

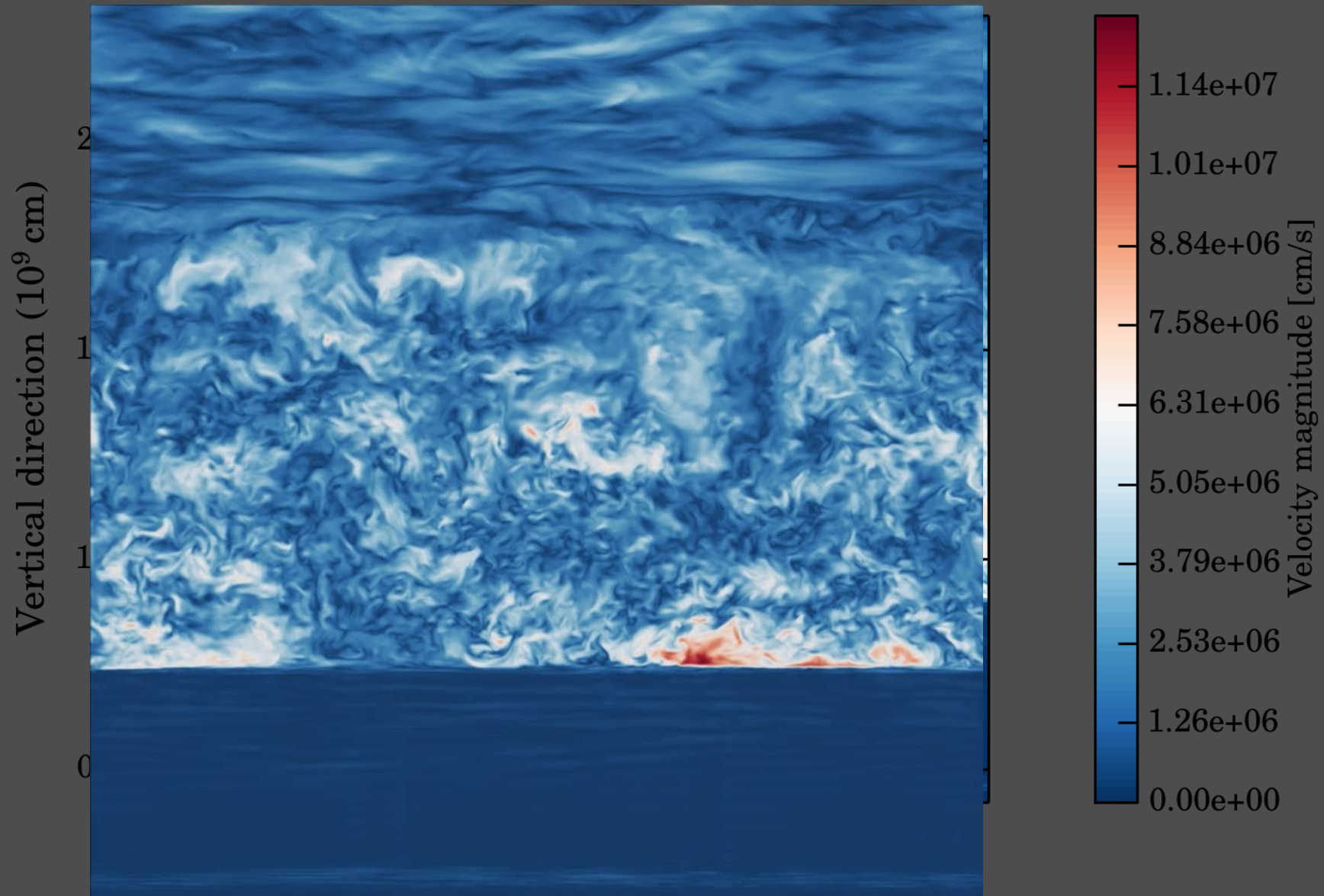
1D boundaries



3D C-shell Simulations: $|v|$ movie

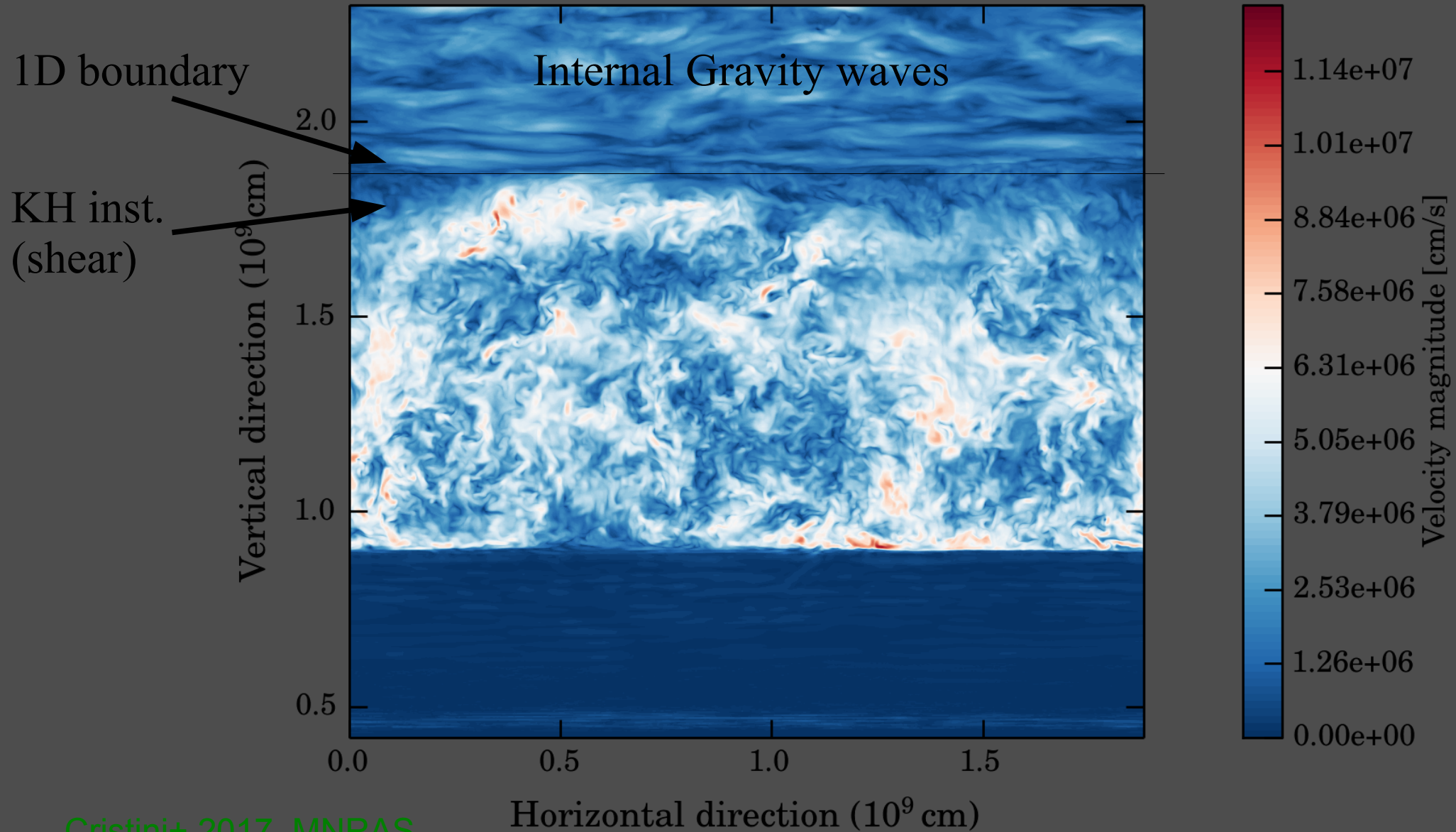
Cristini+ 2017, MNRAS

Gas Velocity $\|v\|$

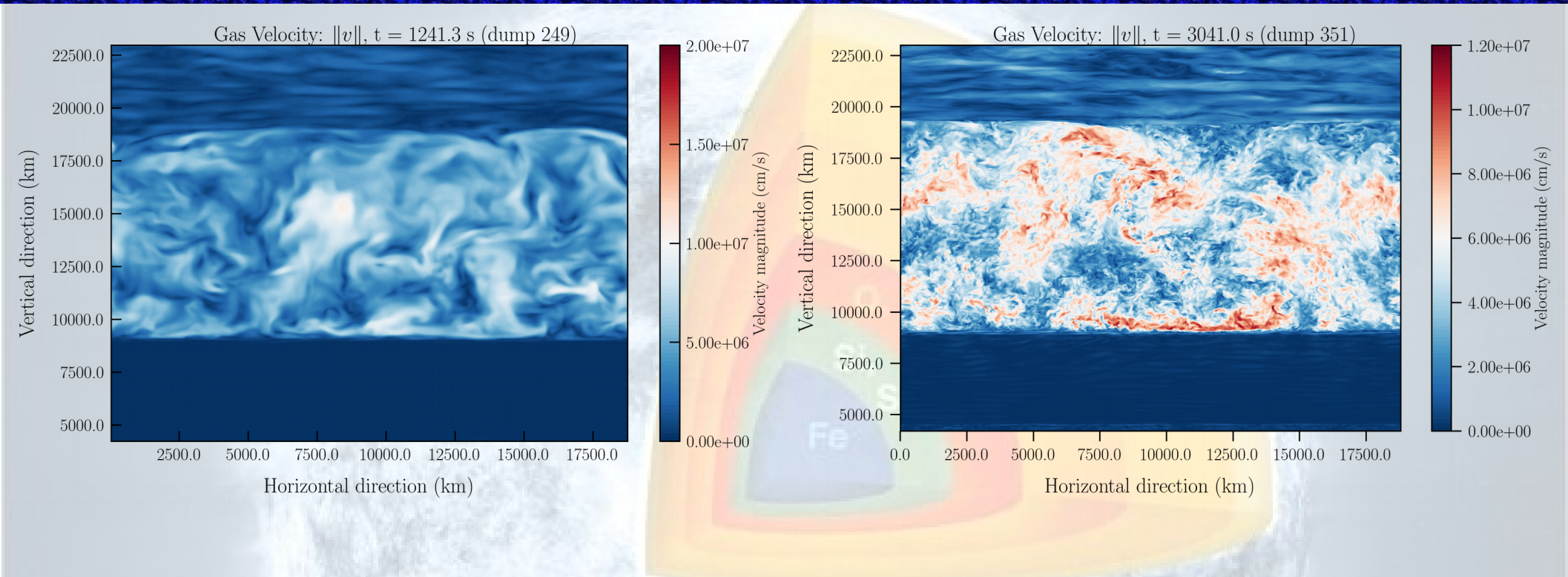


3D C-shell Simulations

Snapshot from 1024^3 resolution run: Gas Velocity $\|\mathbf{v}\|$



Resolution & Luminosity Study



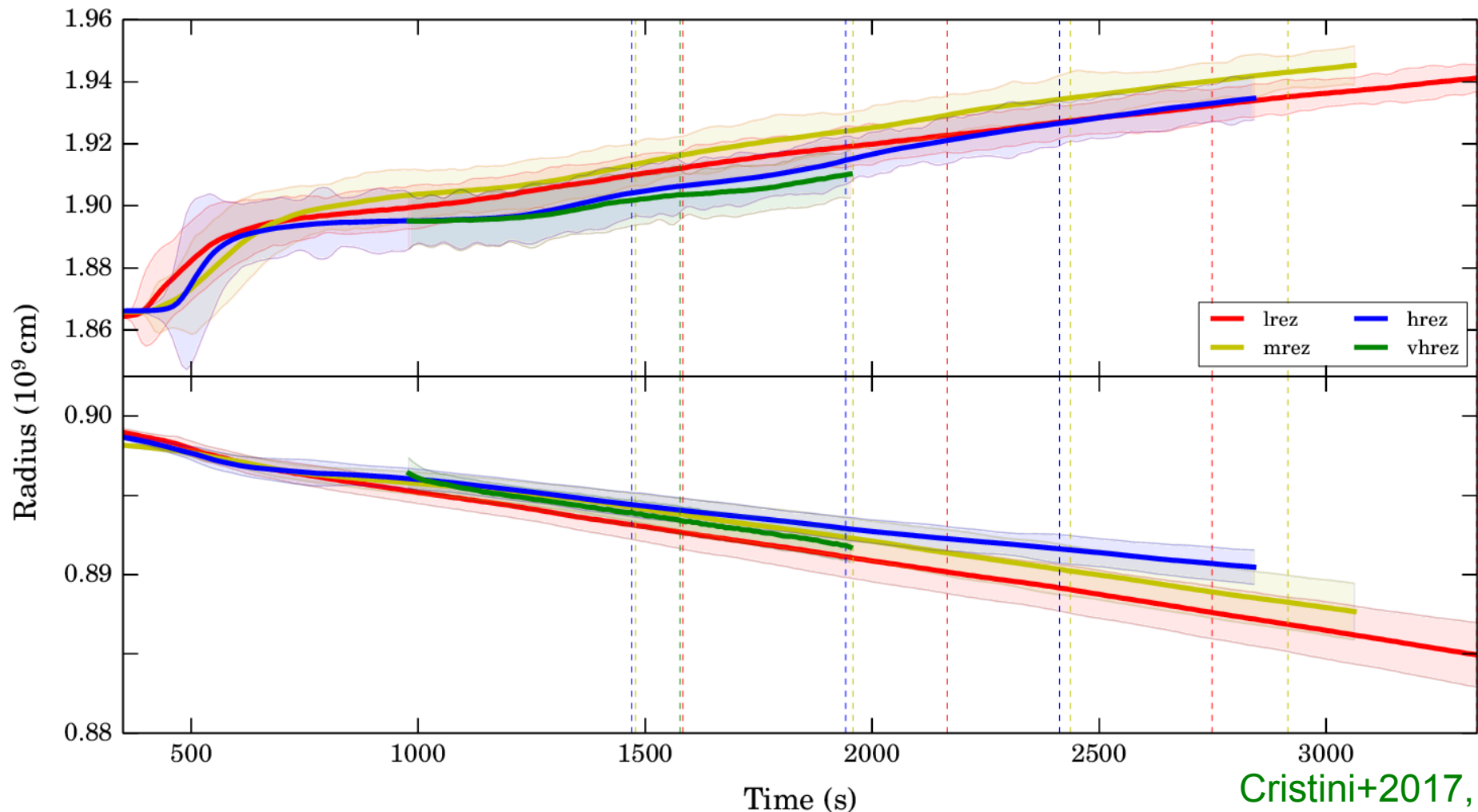
Luminosity		128	256	512	1024
Resolution					
	1			eps1	
	33			eps33	
	100			eps100	
	333			eps333	
	1000	lrez	mrez	hrez/eps1k	vhrez
	3333			eps3k	
	10000			eps10k	
	33333			eps33k	

Cristini+2017, MNRAS

Resolution study

Luminosity study

Boundary Entrainment



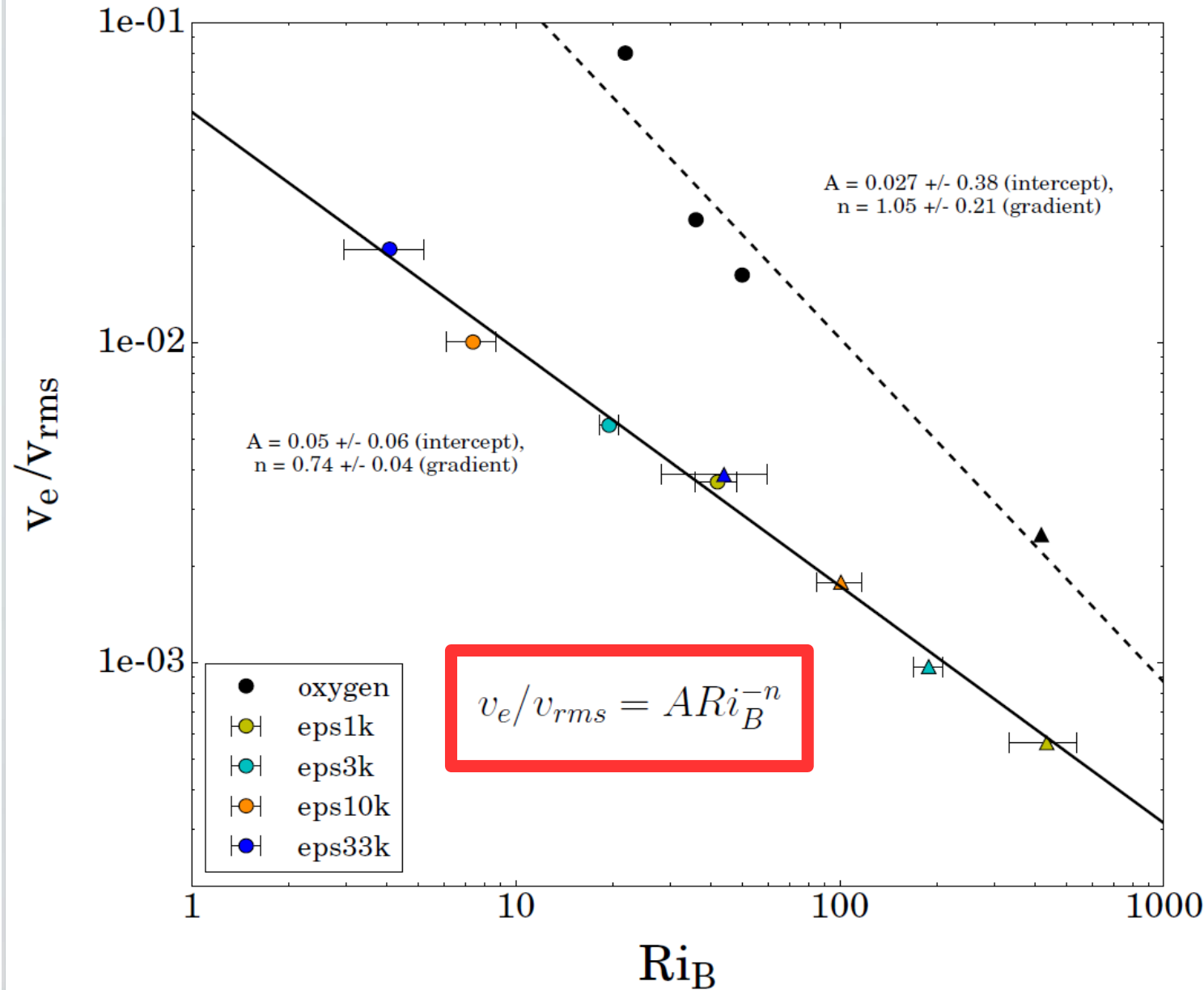
Cristini+2017, MNRAS

Top: $u_e \sim 20,000$ cm/s; Bottom: $u_e \sim 3,000$ cm/s. Rescaled for ϵ_{burn} boosting (1/1000)
→ In 1 year, top: $\Delta R \sim 6 \times 10^8$ cm, bottom: $\Delta R \sim 10^8$ cm: large but reasonable

Consistent with oxygen-shell results and entrainment law.

Entrainment Law

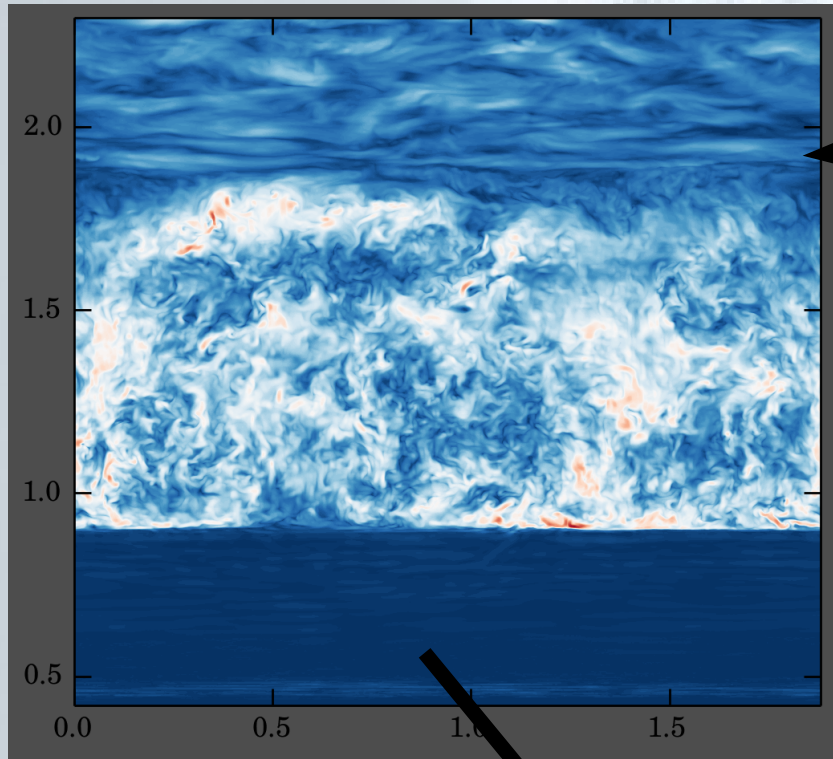
Cristini et al 2019 (see also Garcia & Mellado, 2014, Deardor 1980, Chemel, Staquet and Chollet 2010, Fernando, 1991, Stevens and Lenschow, 2001, Jonker+ 2013)



$$Ri_B = \frac{\Delta B \times l}{v_{rms}^2}$$
$$Ri_B = \frac{\text{stabilising potential}}{\text{turbulent kinetic energy}}$$

Way Forward: 1 to 3 to 1D link

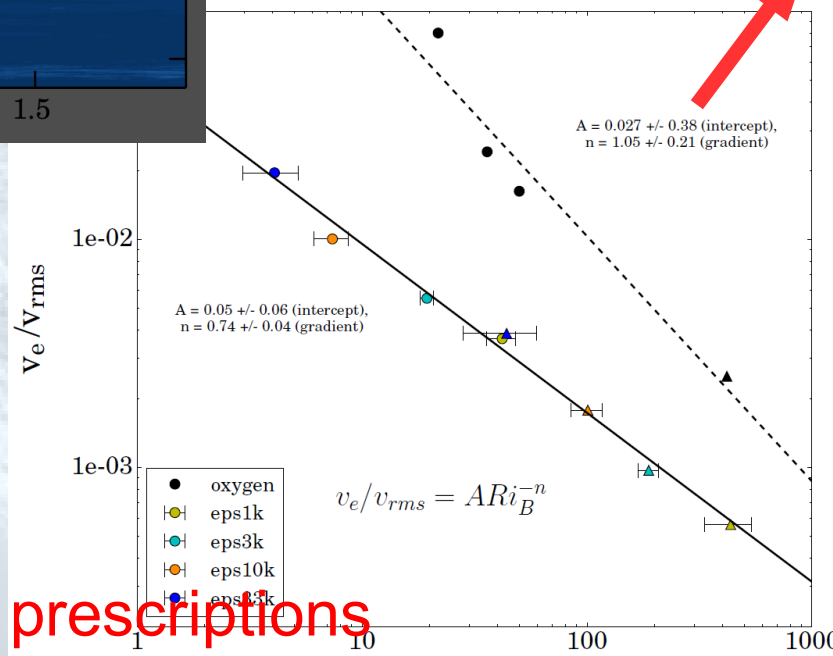
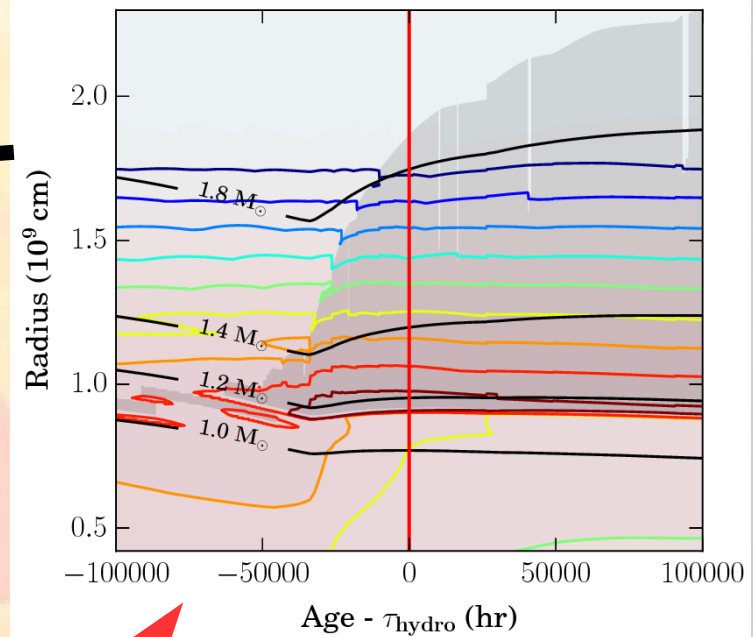
Targeted 3D simulations



Cristini+2017



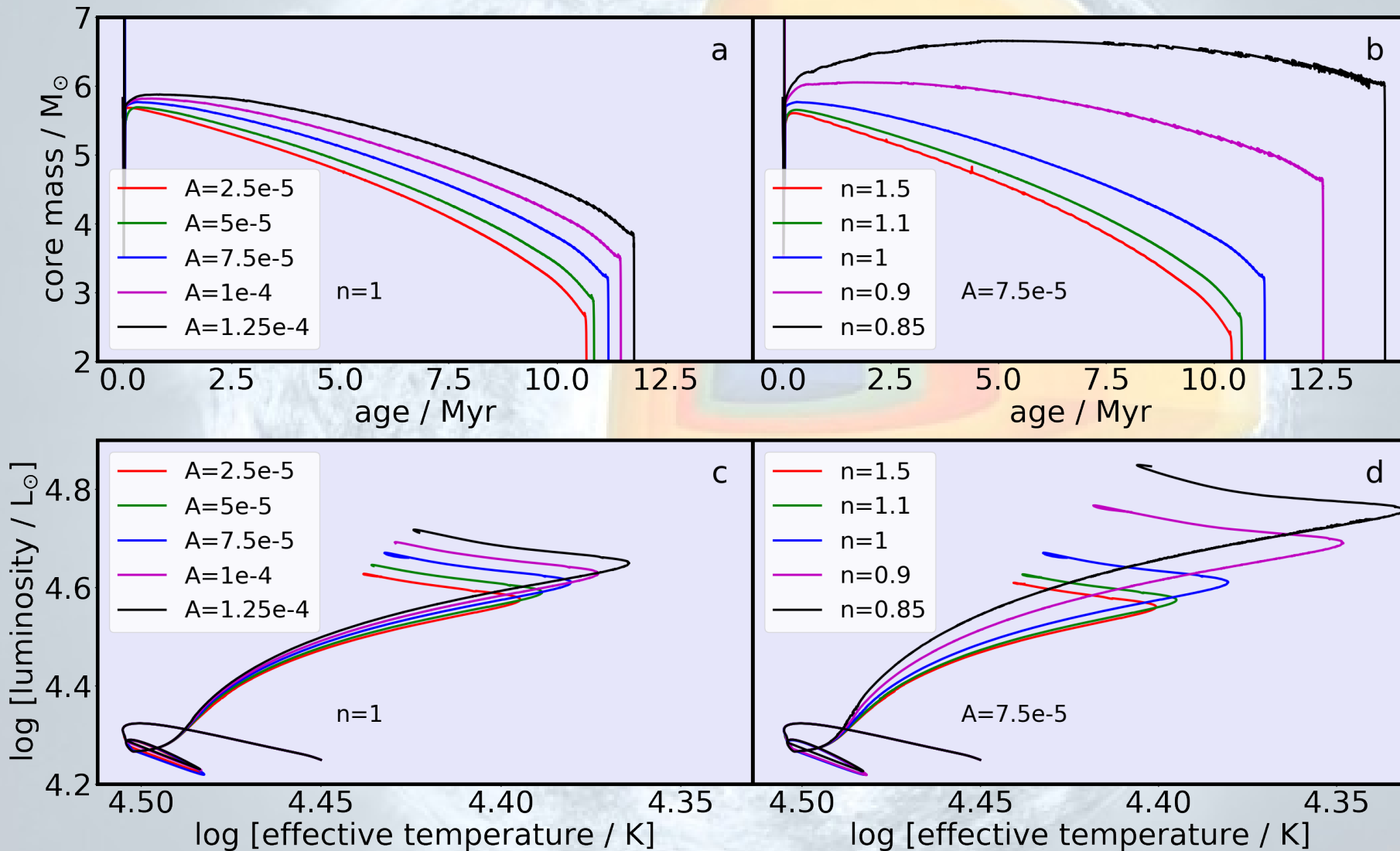
Uncertainties in 1D



→ Improve theoretical prescriptions

Entrainment in 1D model

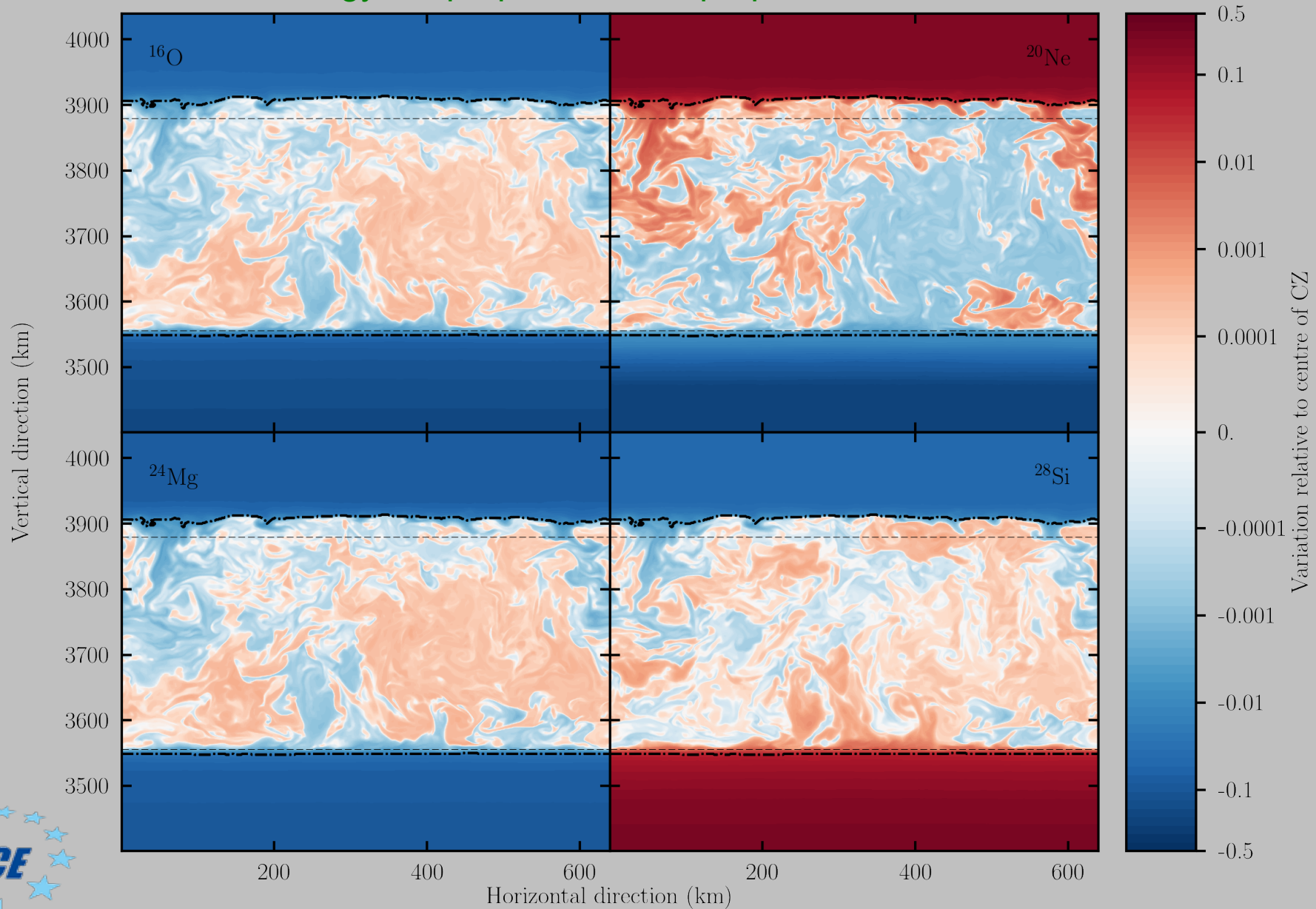
Scott+ in prep



Ne-burning & He-burning Underway

Variation relative to centre of CZ, $t = 673.8$ s (dump 450)

Georgy+ in prep ; Horst+ in prep



Conclusions & Outlook

- 1 - 3 - 1 D modelling of stellar interiors:
 - 2D simulations of dynamical shear
 - 3D C-shell simulations follow entrainment law
 - Entrainment law being tested in 1D code
 - & **new CBM prescriptions under development!**
 - Exciting times ahead:
 - - Great observational constraints
 - - Complex physics explored in unprecedented details
 - - CPU-GPU supercomputers
- 