

kinetic simulations of relativistic harmonic magnetic equilibria

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talking on
inflation (Thu)

content of this talk

- kinetic particle-in-cell simulations
- Zeltron code (B. Cerutti)
- harmonic magnetic equilibrium (“ABC fields”)
- fundamental unstable mode
- 2D periodic domain
- pair plasma
- focus on particle acceleration and evolution of current layers

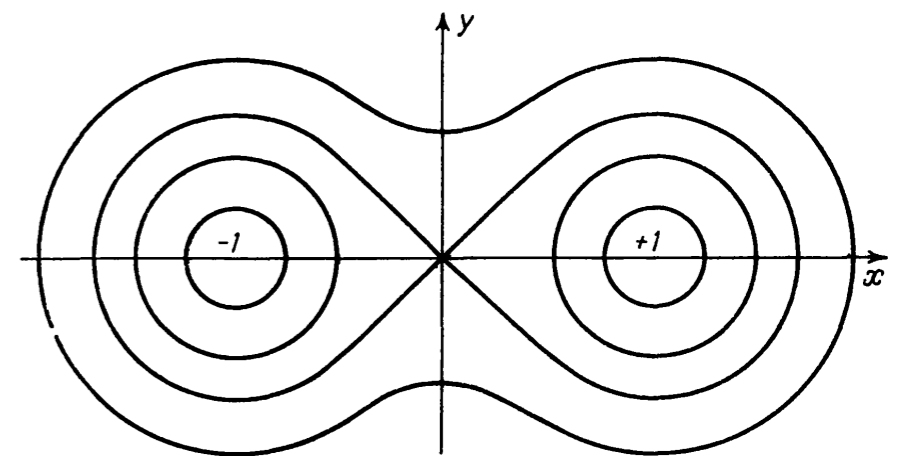
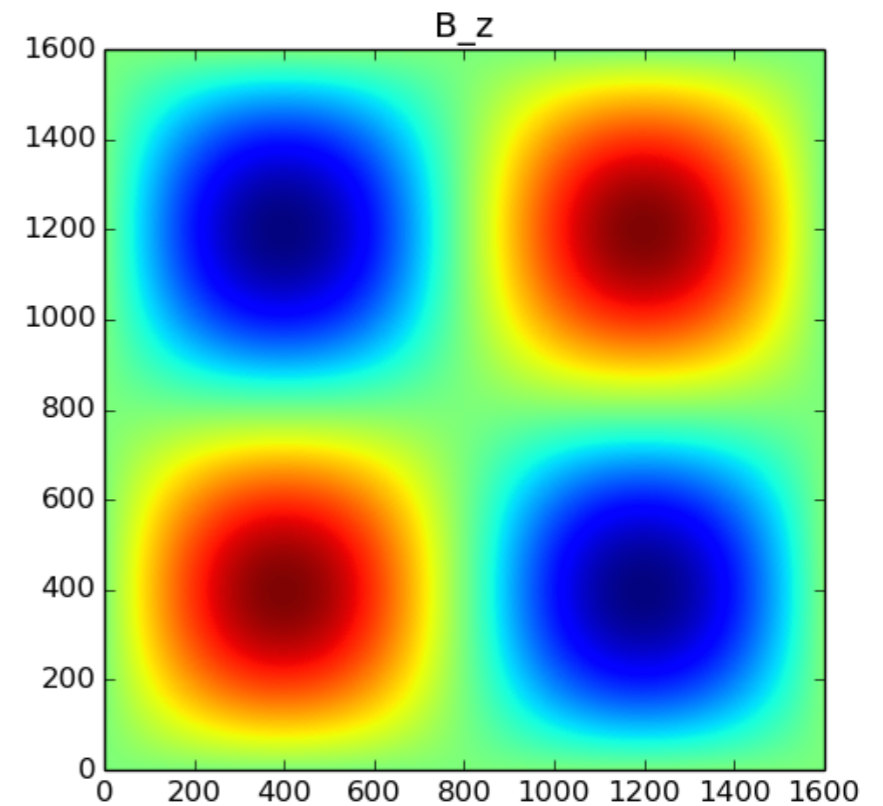
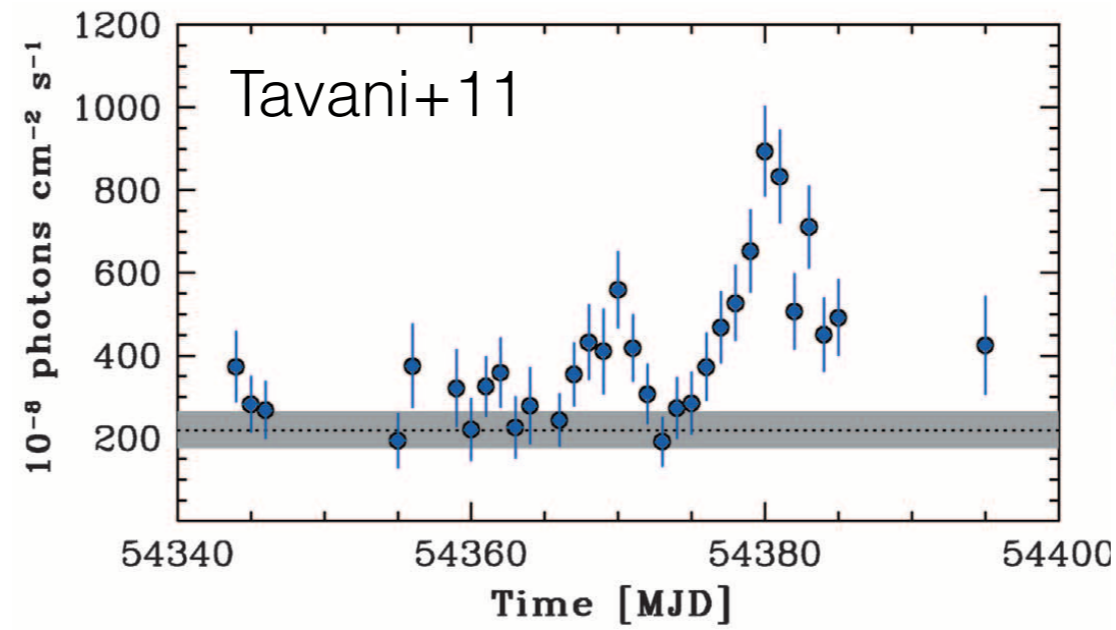
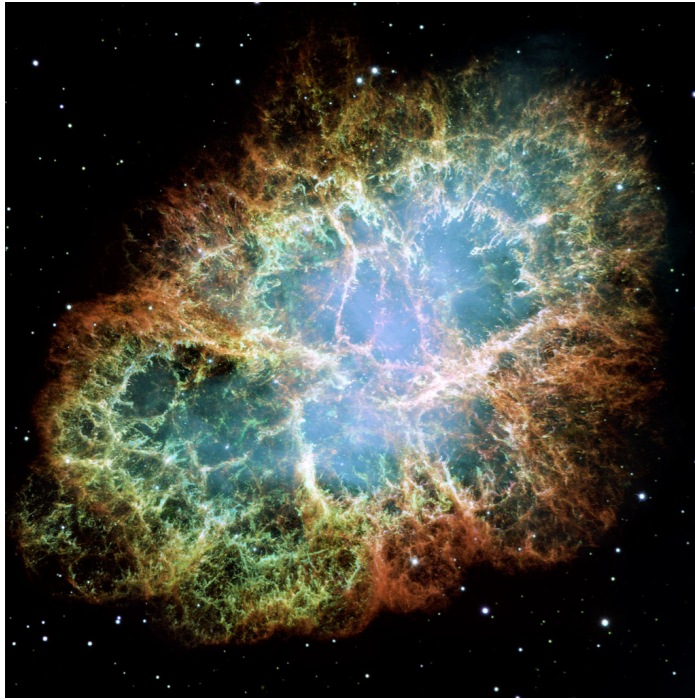


Fig. 1 Syrovatskii 66

similarities to Lyutikov,
Sironi & Komissarov

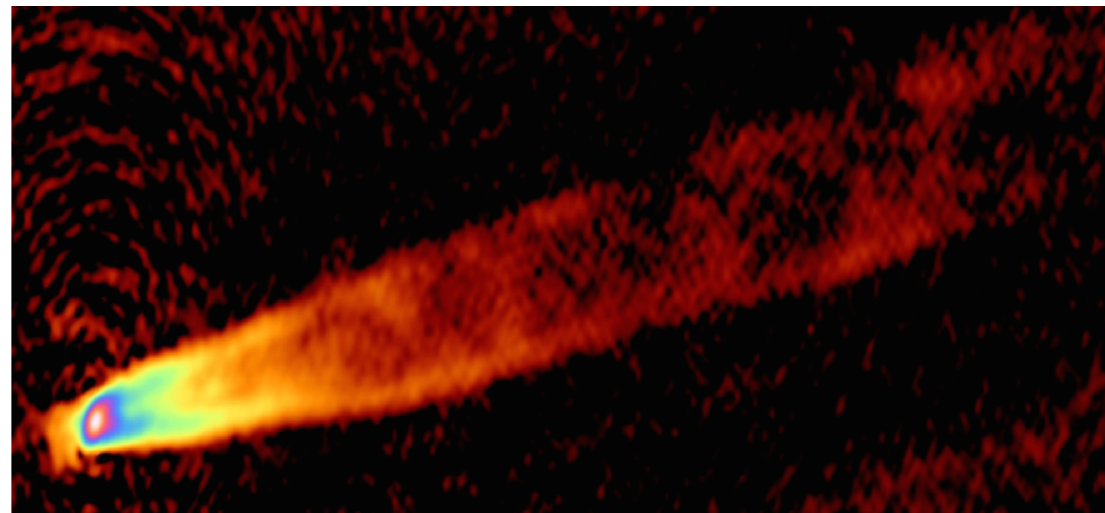
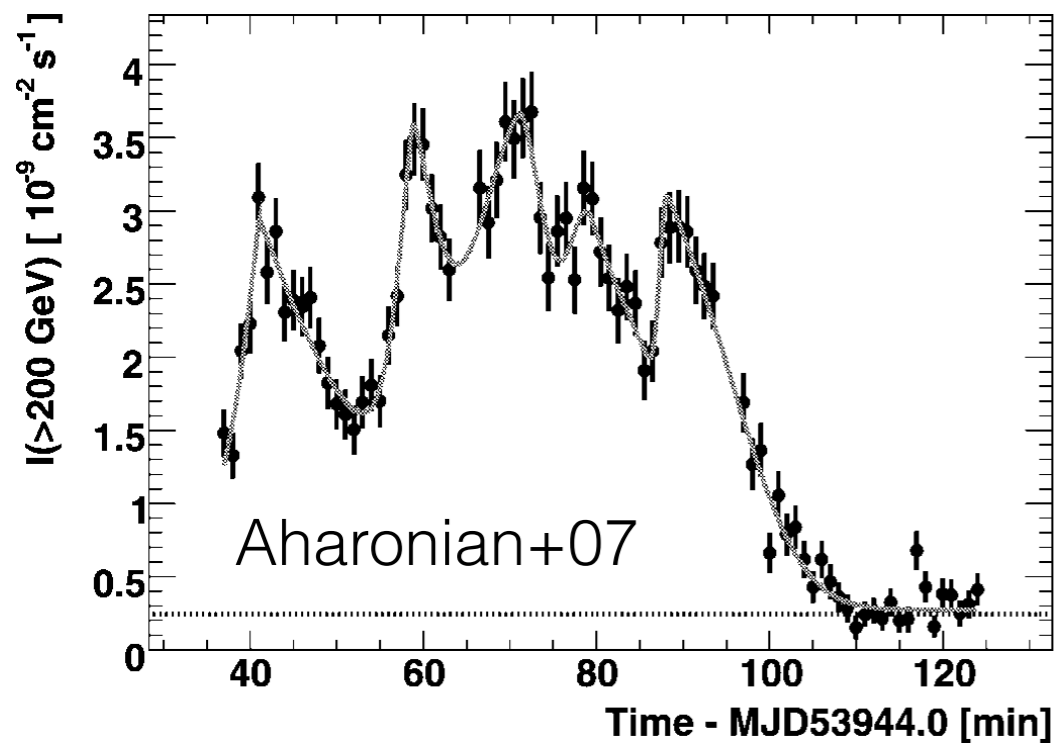
motivation

Crab Nebula

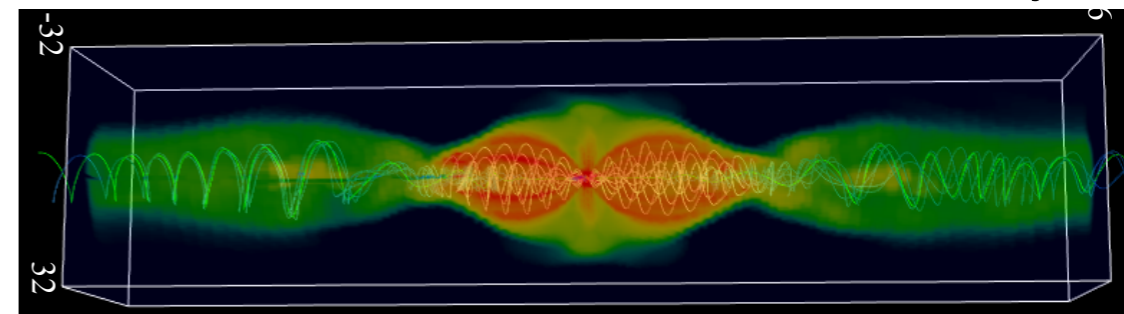


Porth+14

relativistic jets



Bromberg & Tchekhovskoy 15



magnetoluminescence

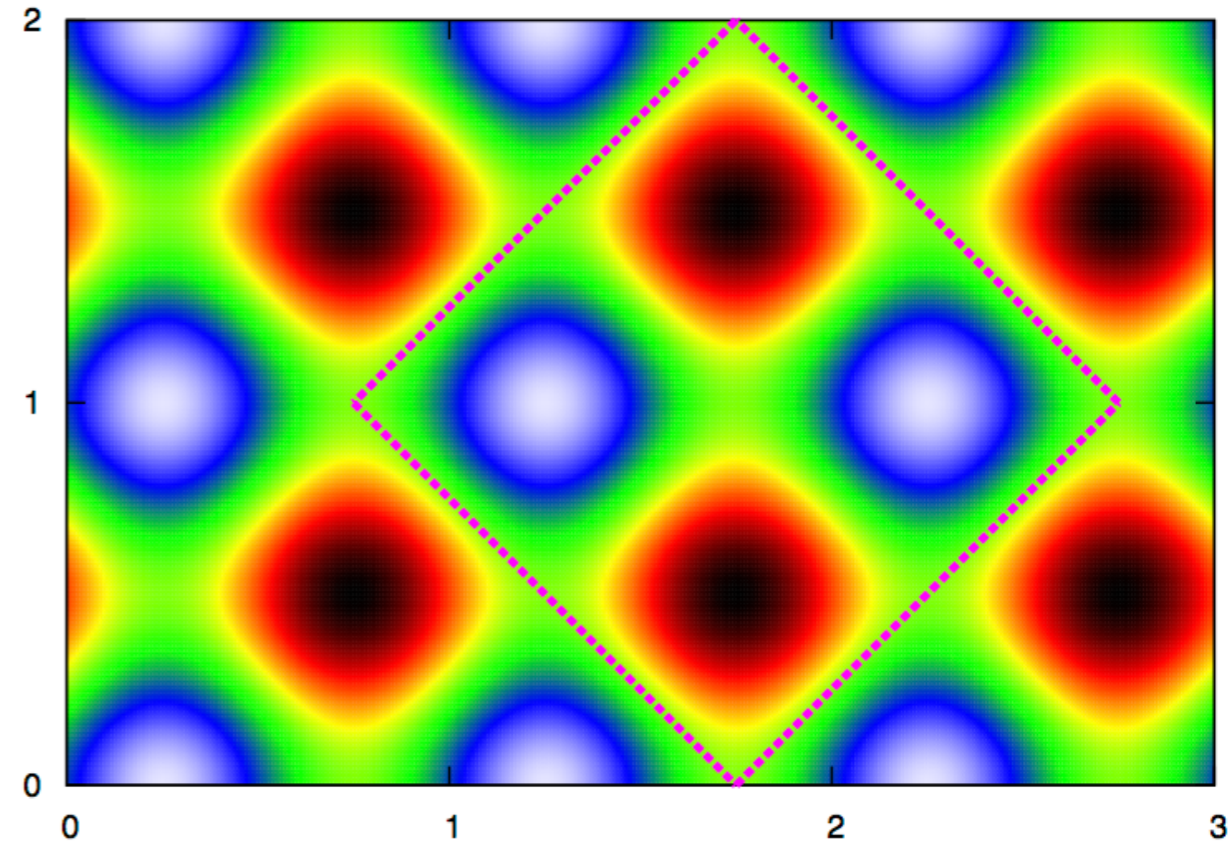
a process of extracting magnetic energy by means of dynamical instability (implosion) leading to transient current layers enabling efficient particle acceleration, and consequently a transient gamma-ray emission

Methods:

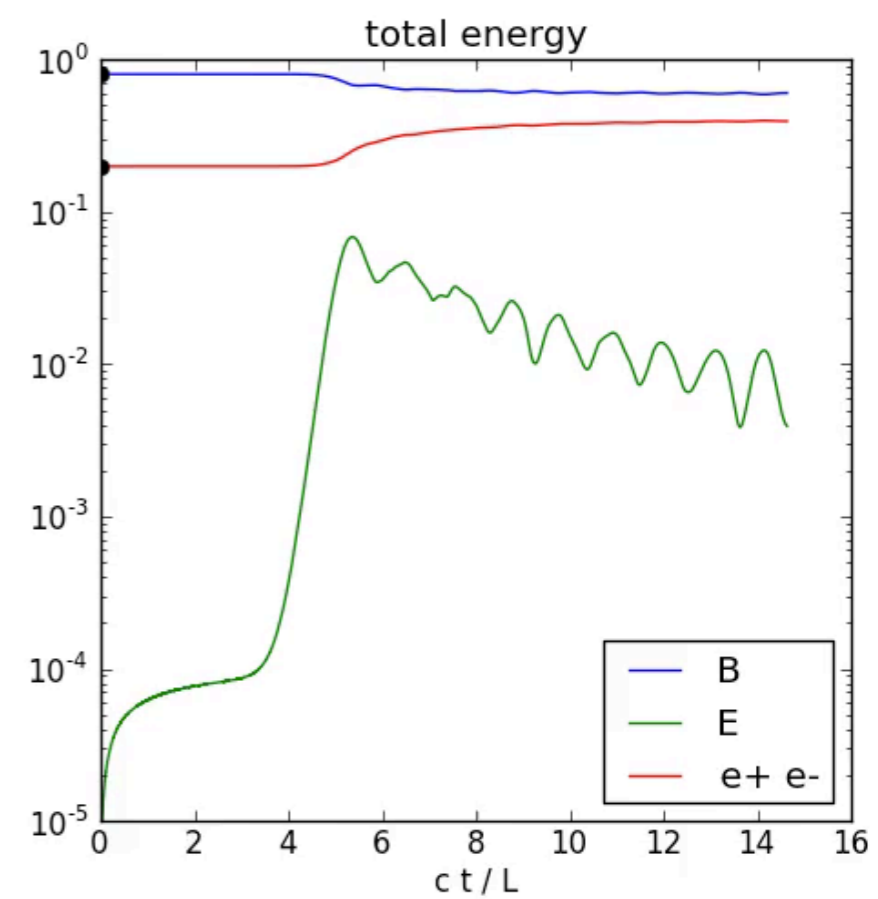
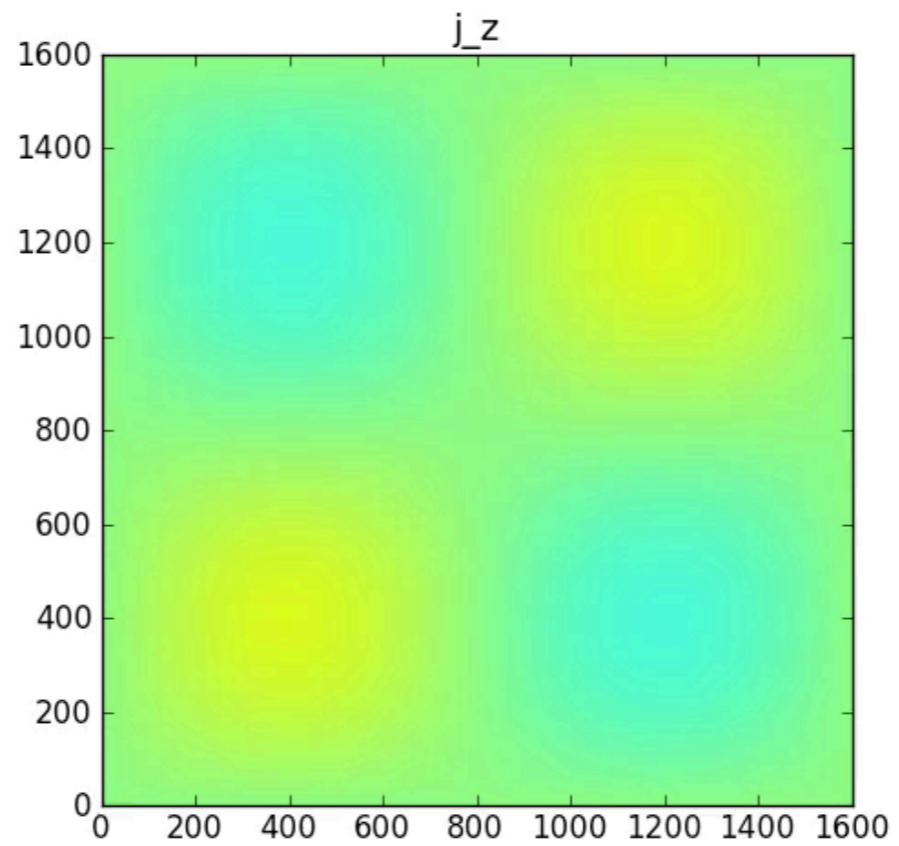
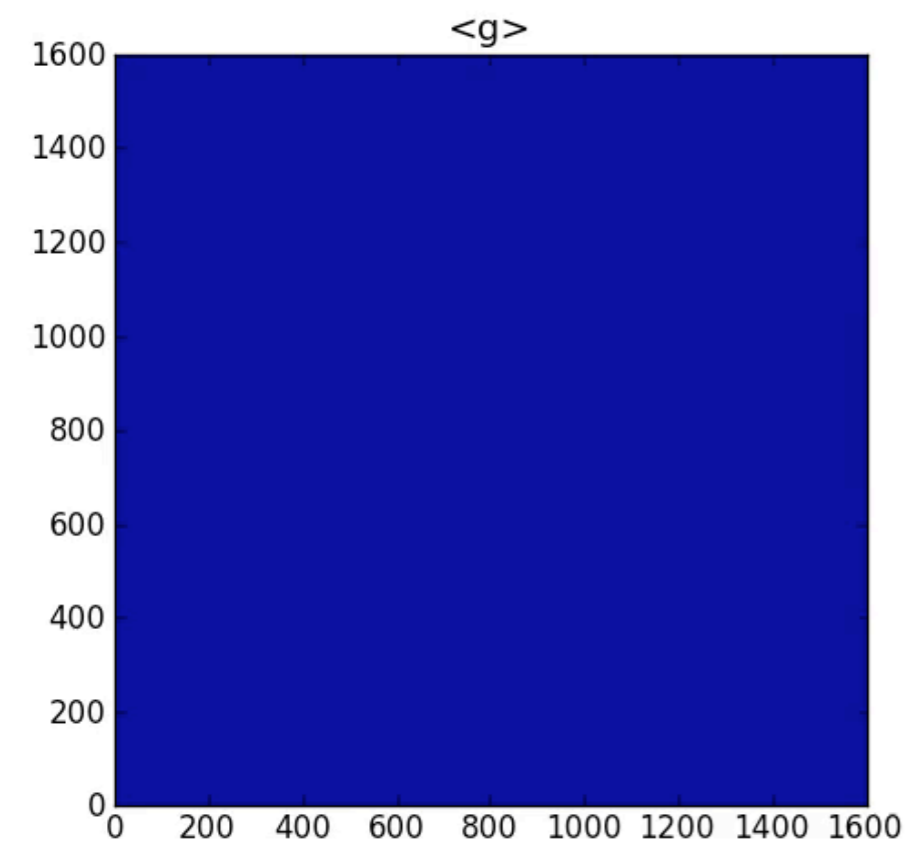
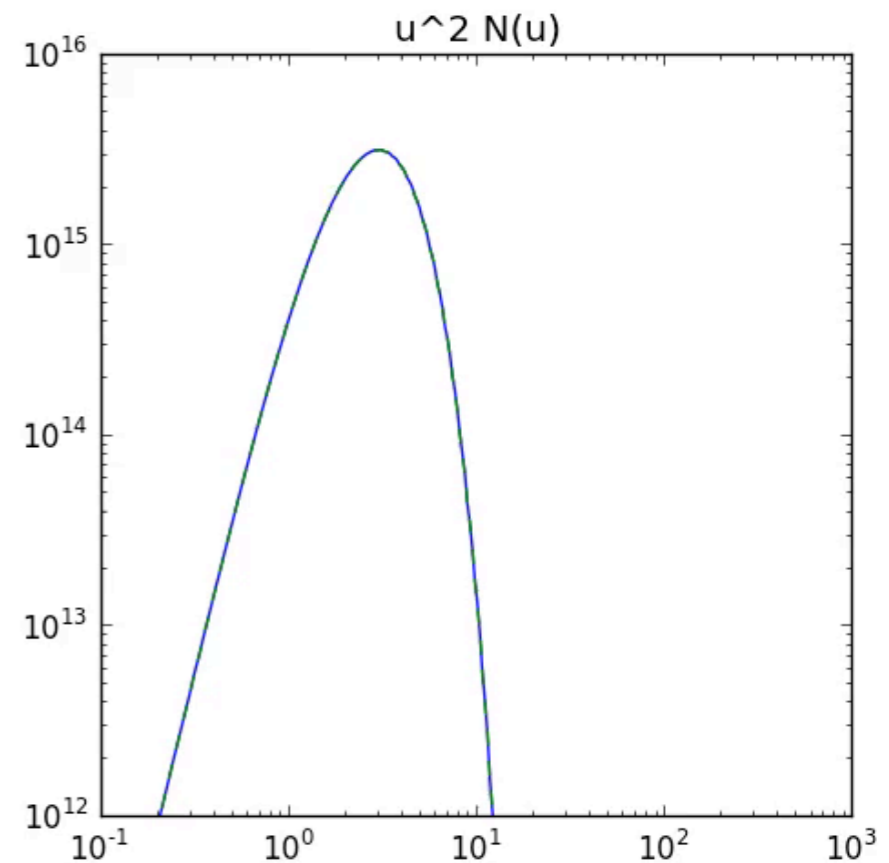
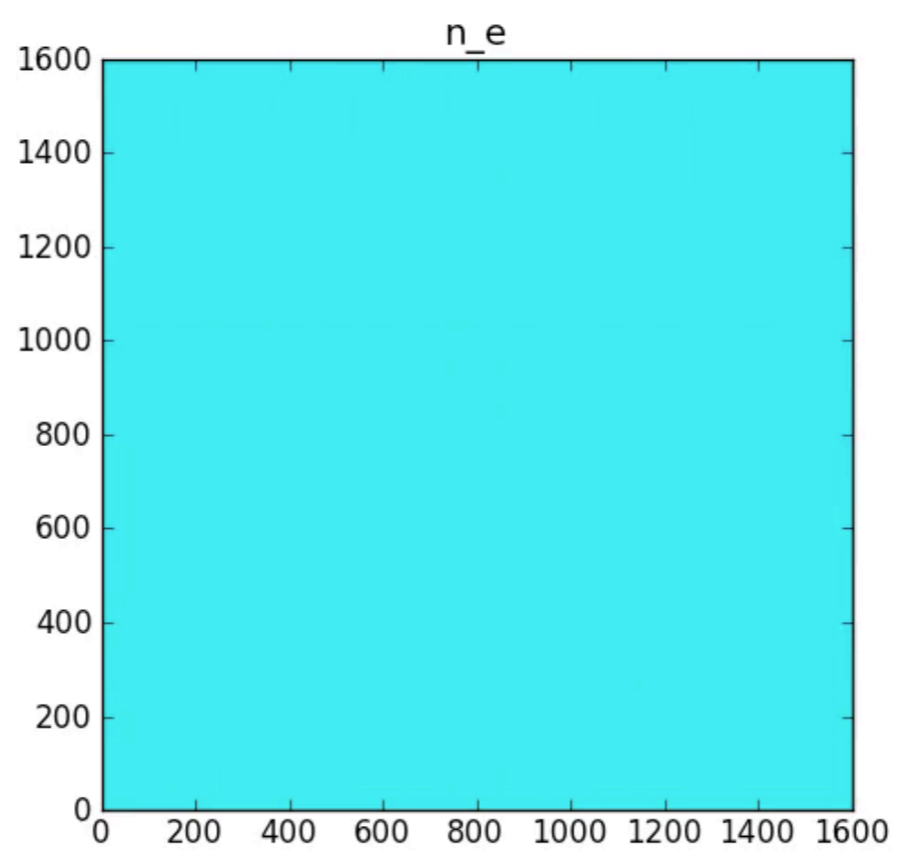
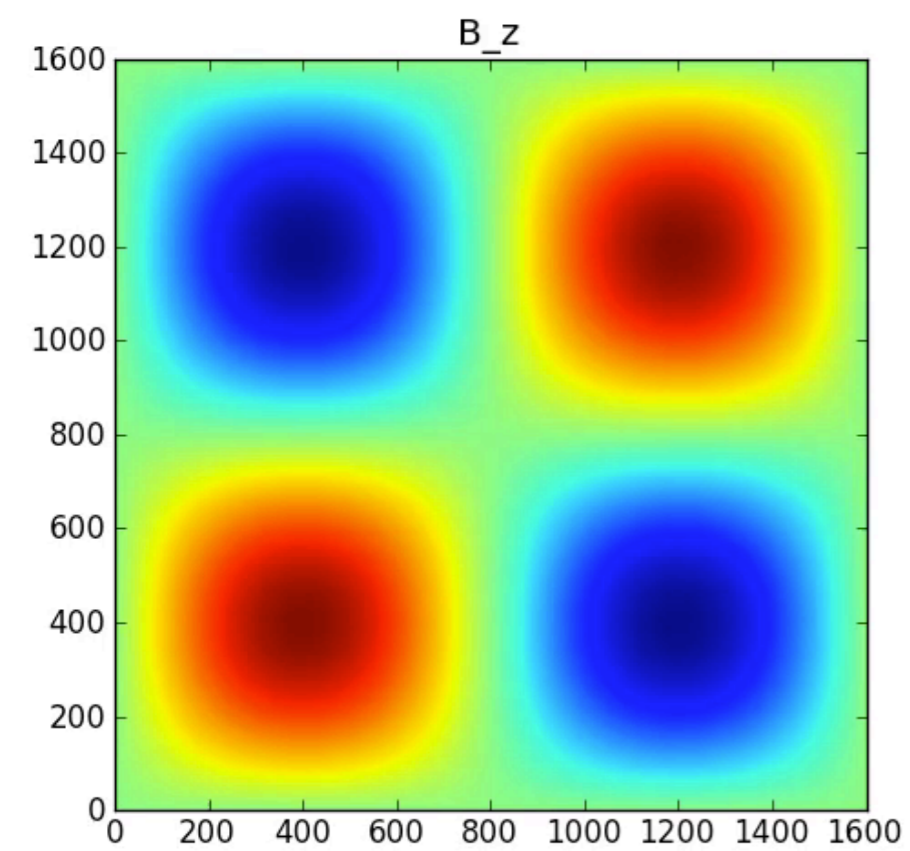
- analytical stability analysis (Y. Yuan)
- relativistic MHD (J. Zrake)
- relativistic force-free (W. East)
- particle-in-cell (this talk)
- radiative PIC (Y. Yuan)

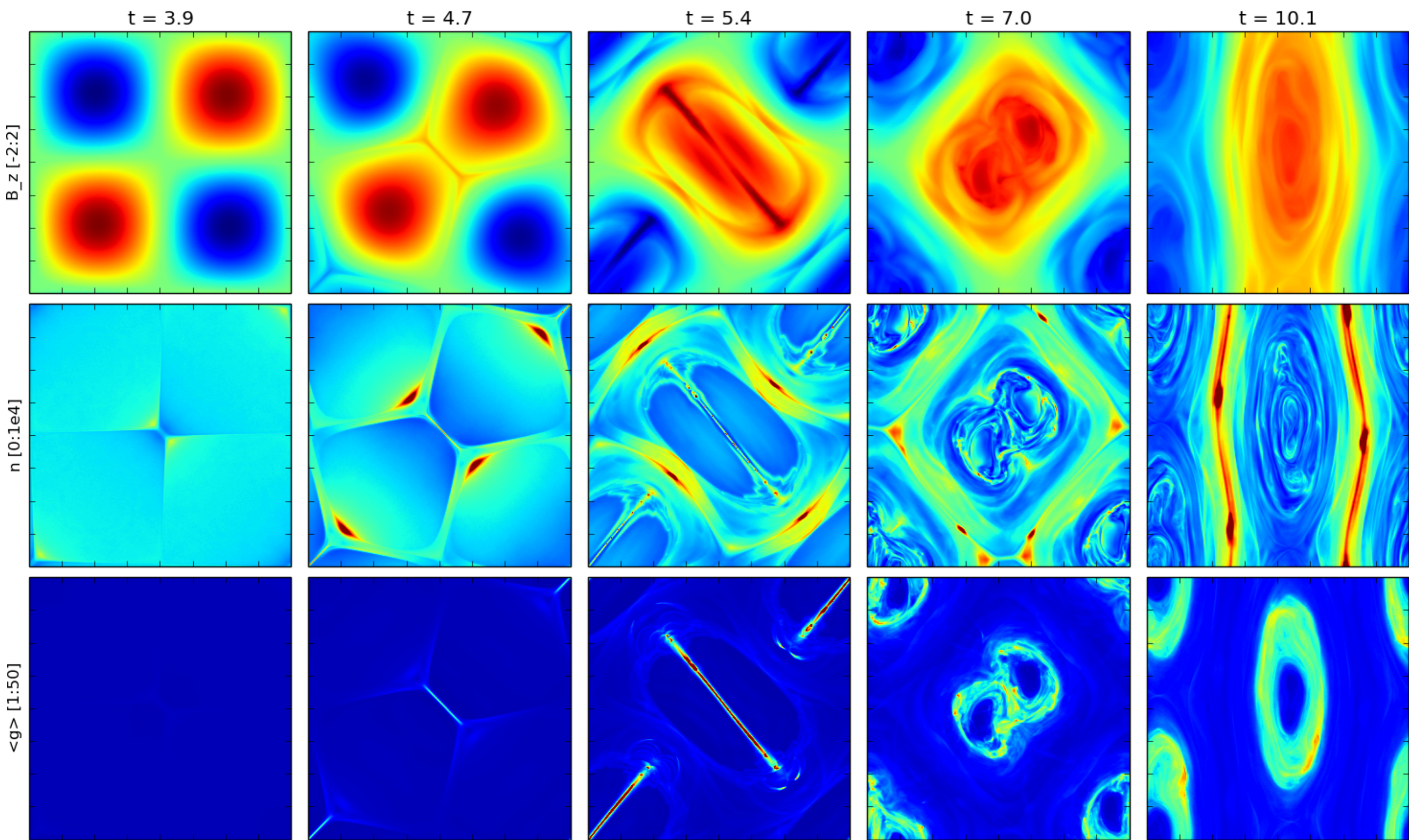
harmonic magnetic equilibria

- Beltrami condition:
 $\nabla \times B = \alpha B$
 $B = \alpha A, j = -(\alpha c/4\pi)B$
- ABC field:
 $B_x = B_3 \sin(\alpha z) + B_2 \cos(\alpha y)$
 $B_y = B_1 \sin(\alpha x) + B_3 \cos(\alpha z)$
 $B_z = B_2 \sin(\alpha y) + B_1 \cos(\alpha x)$
- 2D: $B_1 = B_2 = 1, B_3 = 0$
- fundamental unstable mode:
2 maxima and 2 minima of A_z
- no kinetic-scale initial structure



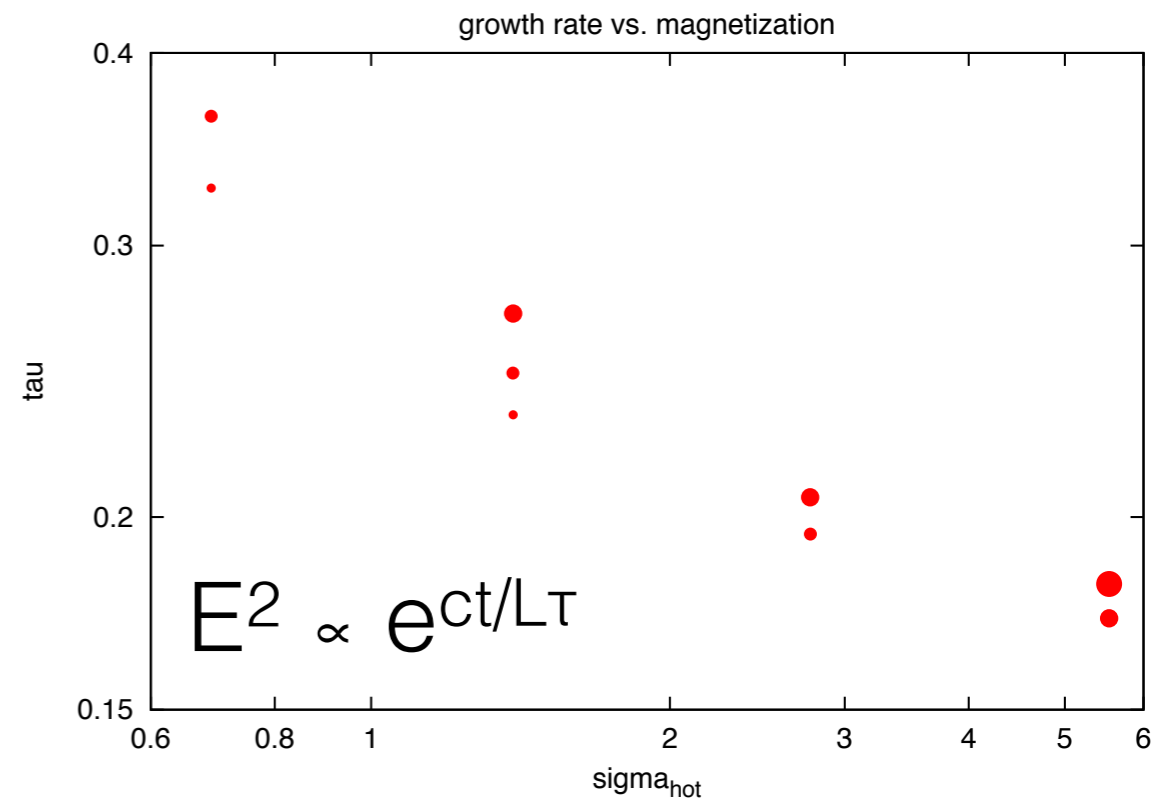
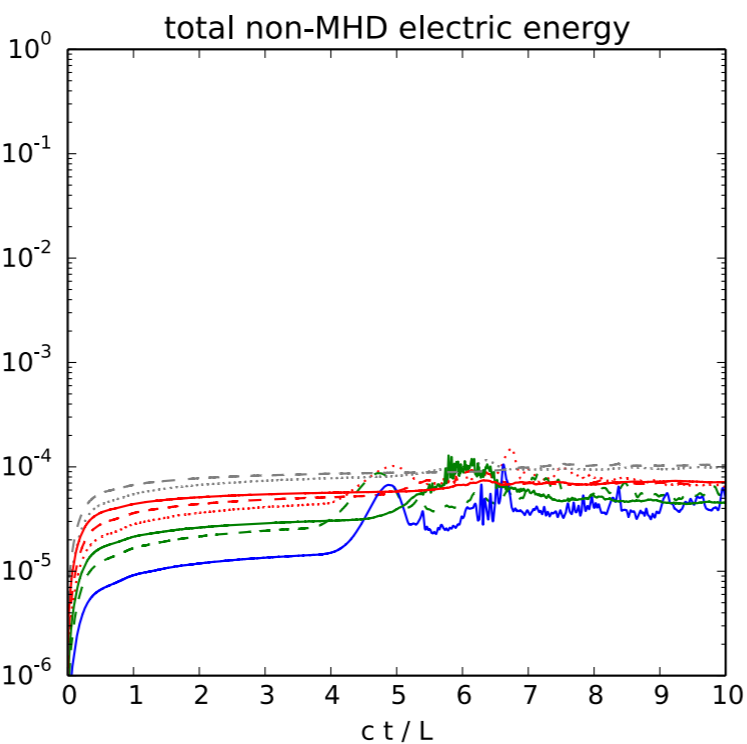
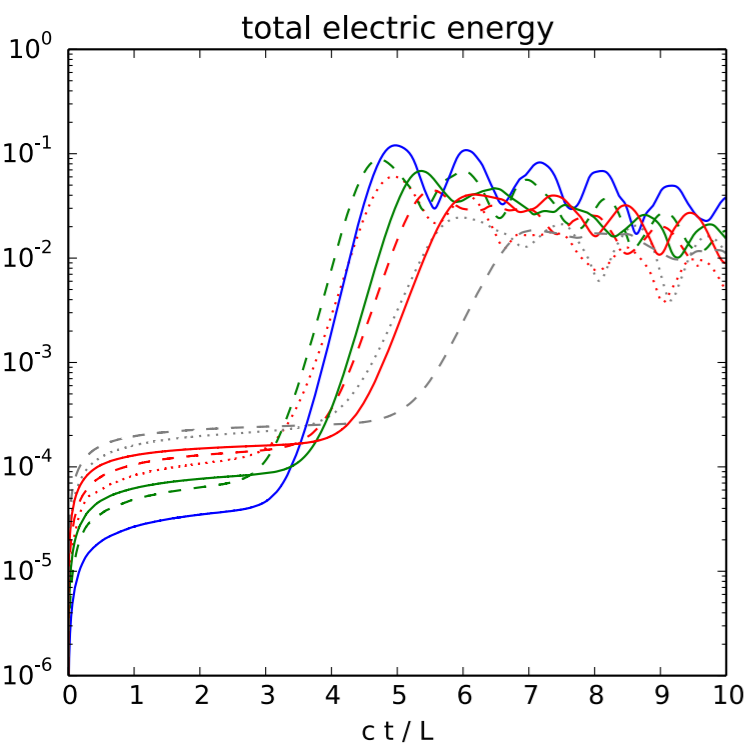
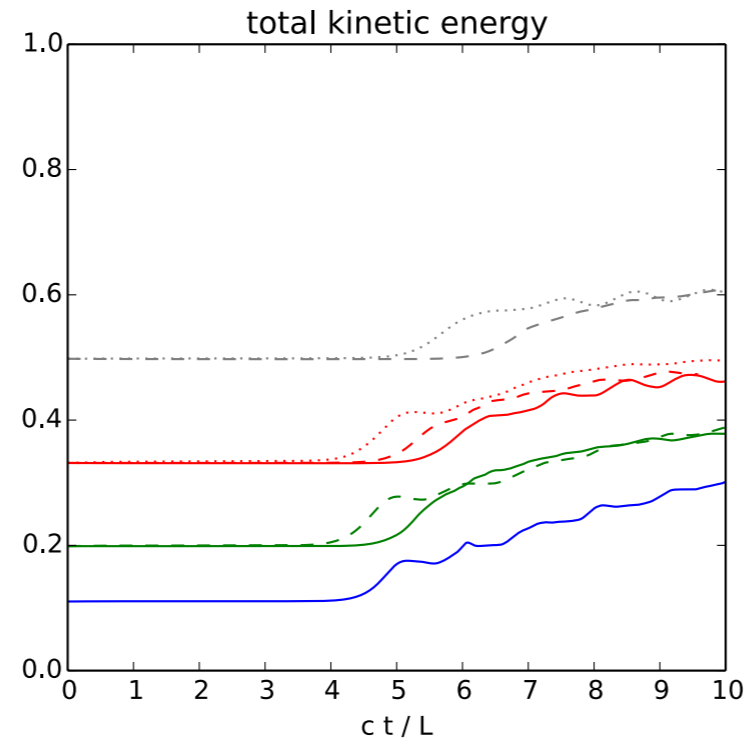
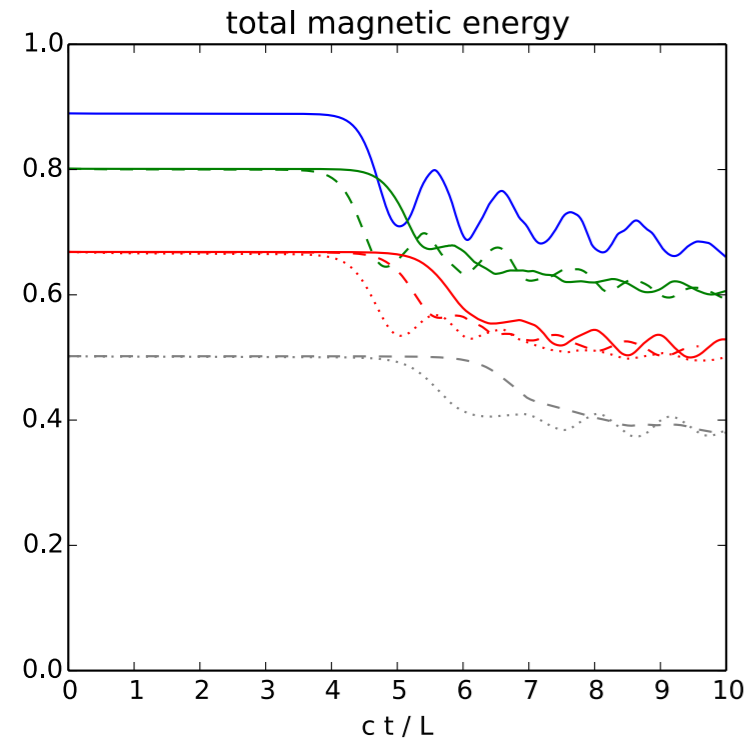
- current density from dipole moment a_1 in particle momentum distribution
- mean magnetization
 $\sigma \propto a_1(L/\rho_0)$





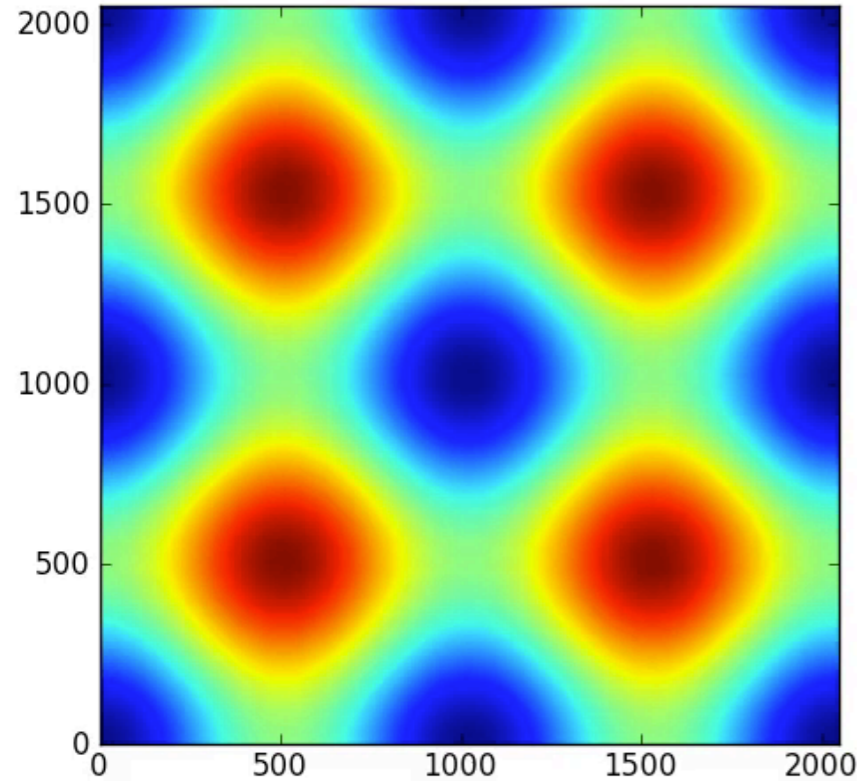
total energy

- linear instability seen in total electric energy
- non-ideal electric energy appears insignificant
- relative magnetic dissipation efficiency is constant

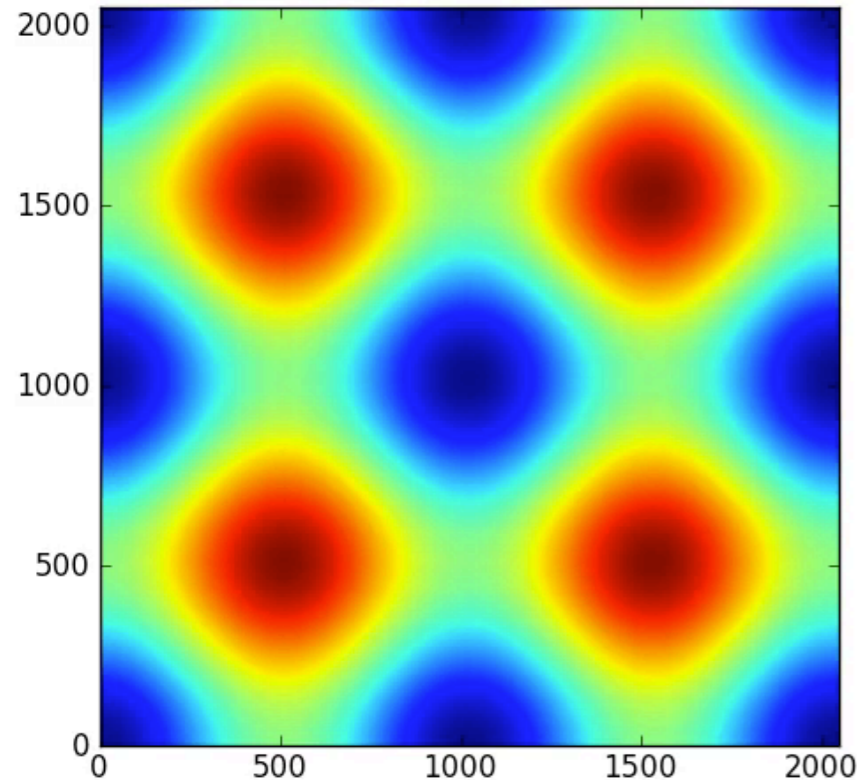


magnetic helicity

$B^2, t=0$

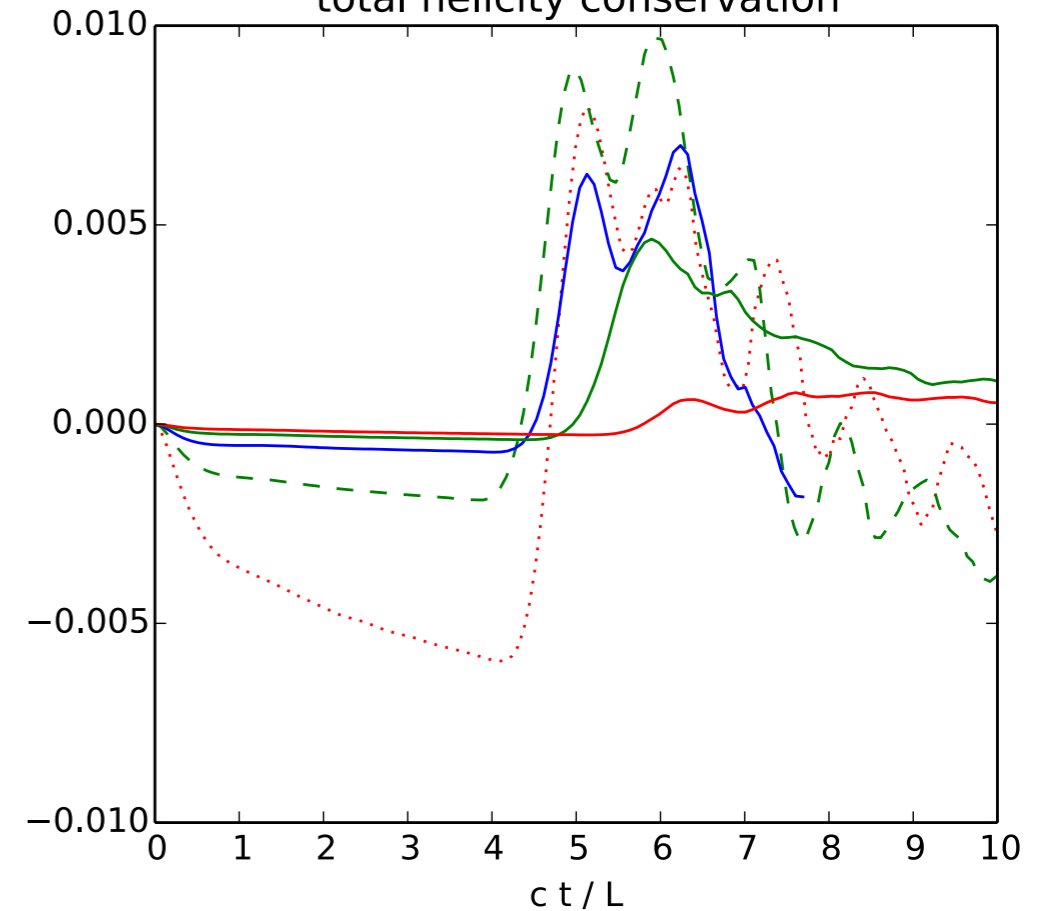


$A \cdot B, t=0$

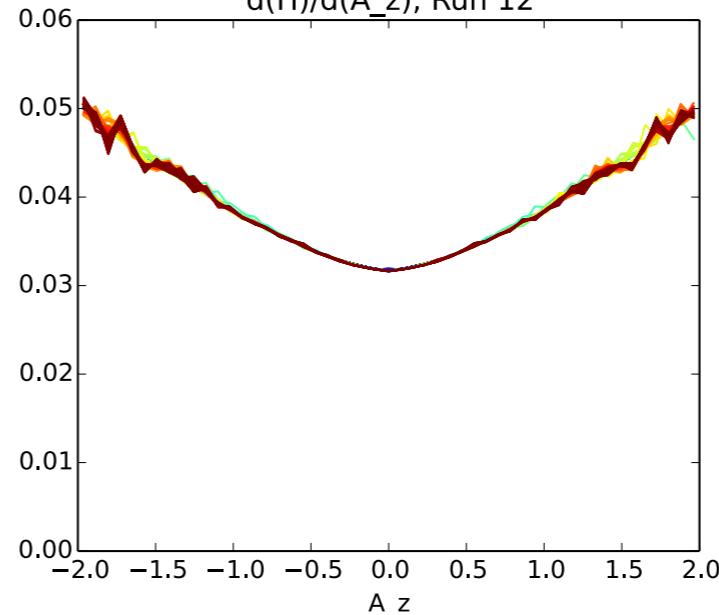


- $H = \int (A \cdot B) dV$
- conserved quantity in ideal MHD
- conservation in PIC simulations better than 1%, especially for small dipole moment a_1

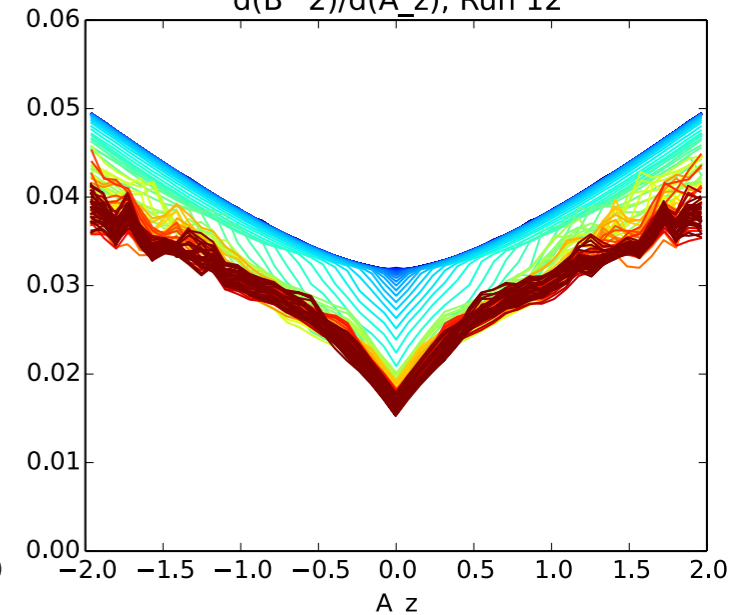
total helicity conservation



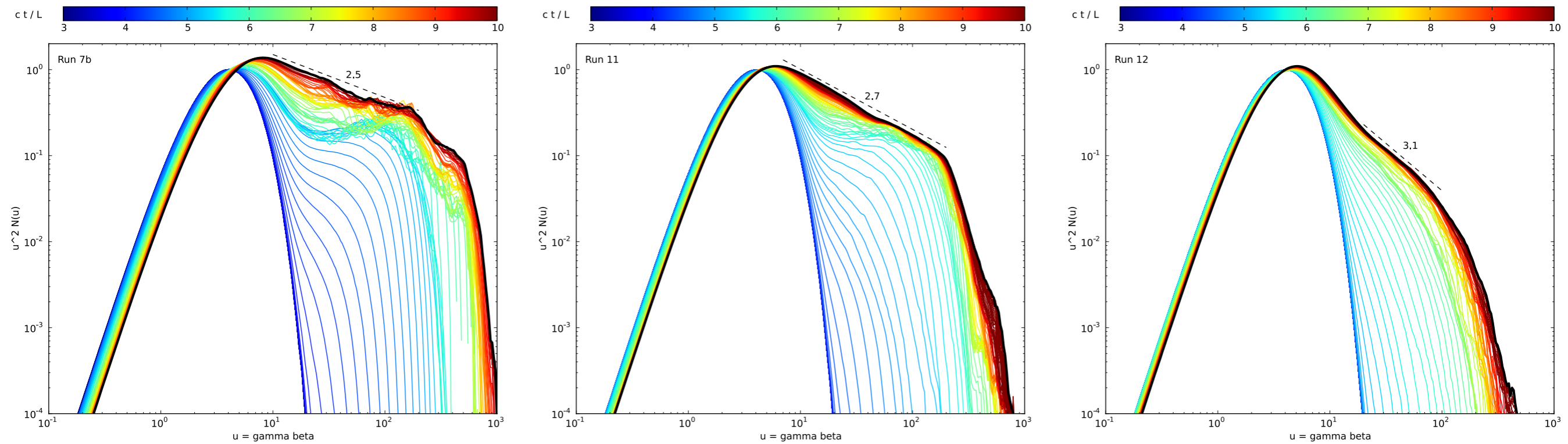
$d(H)/d(A_z), \text{Run 12}$



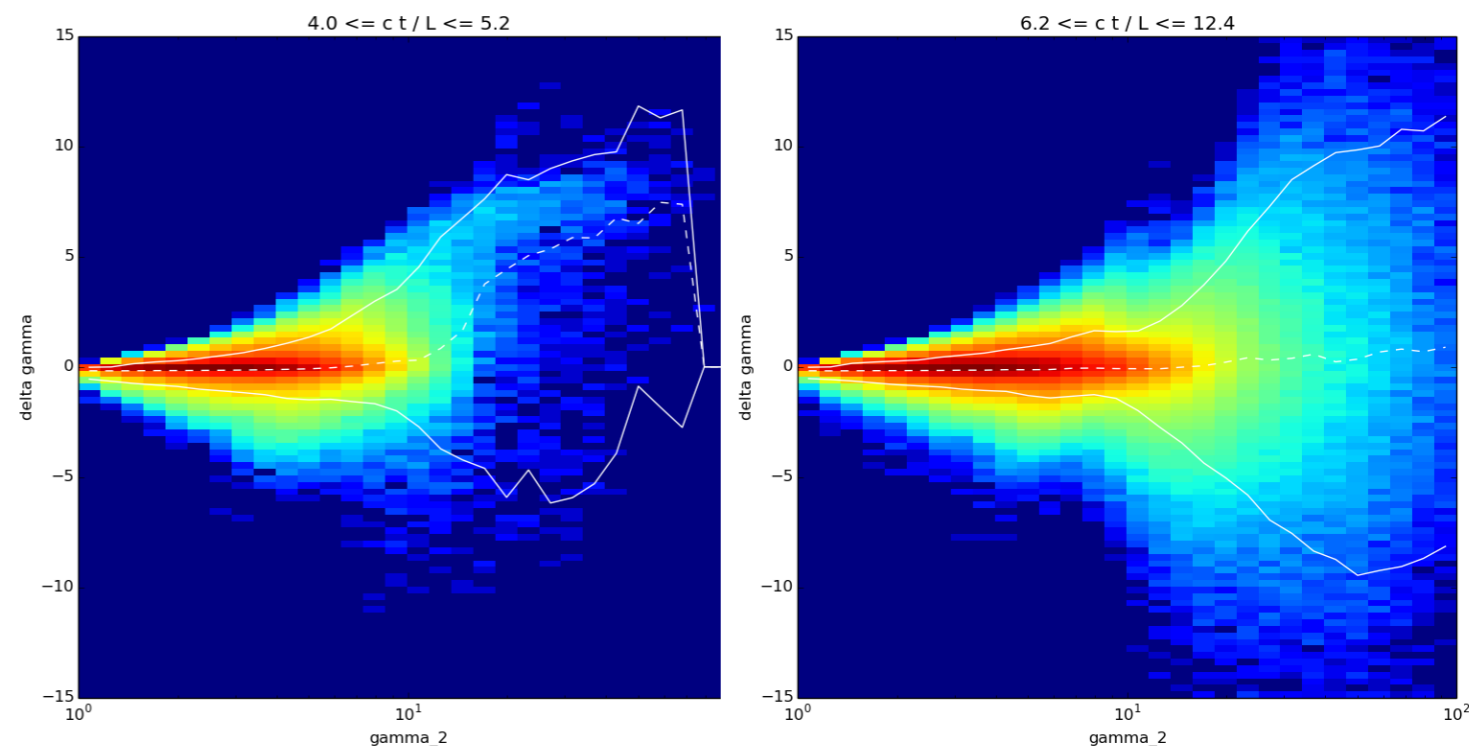
$d(B^2)/d(A_z), \text{Run 12}$



particle energy distribution

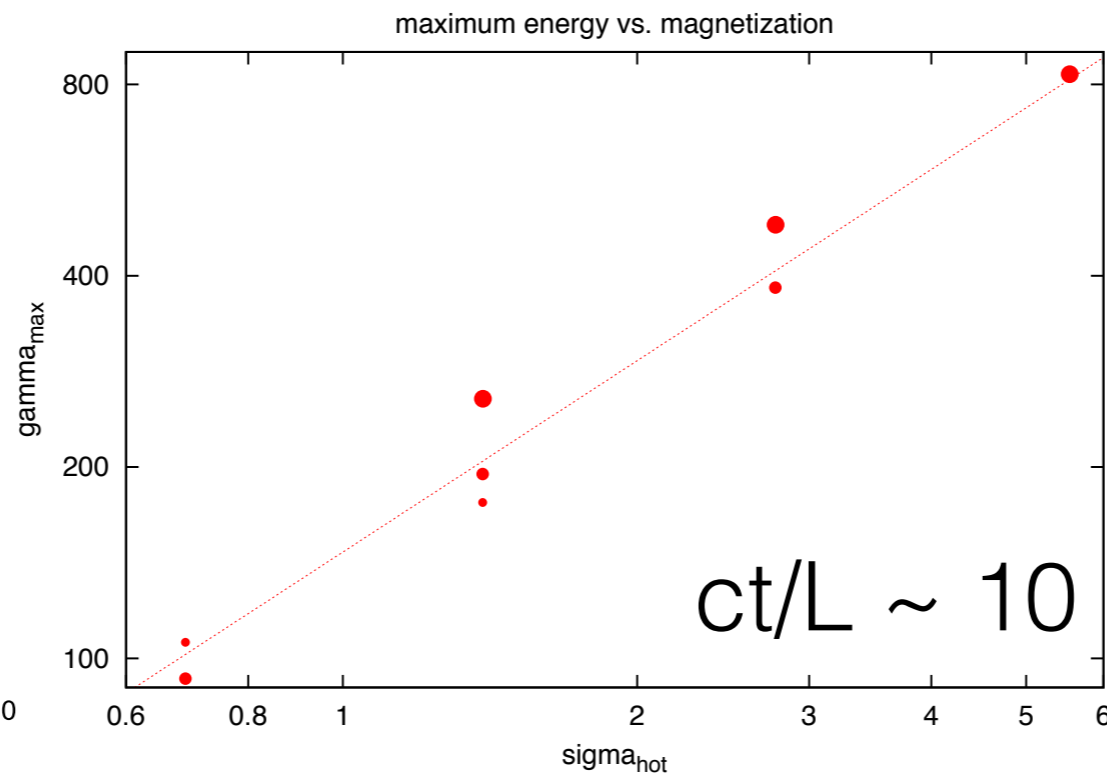
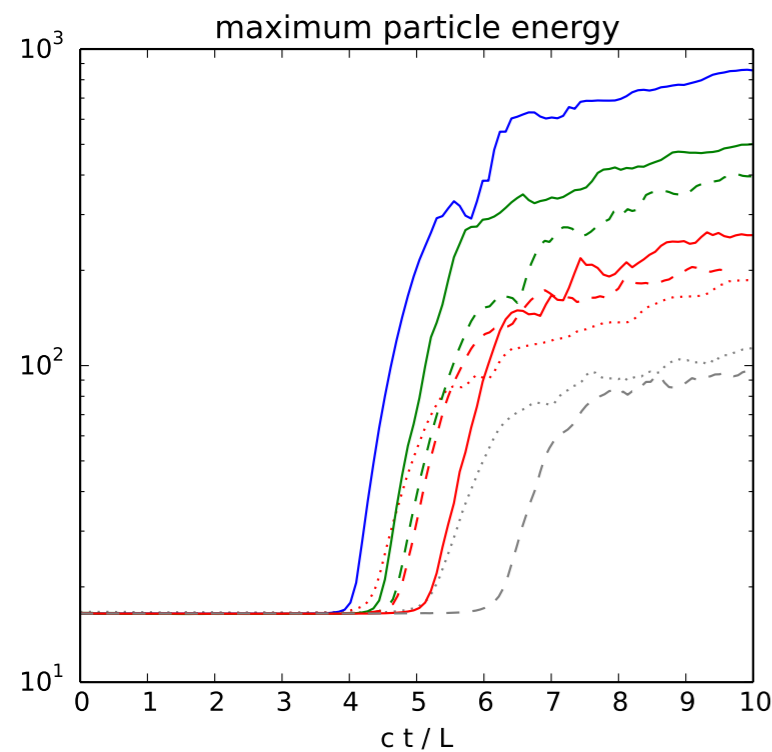
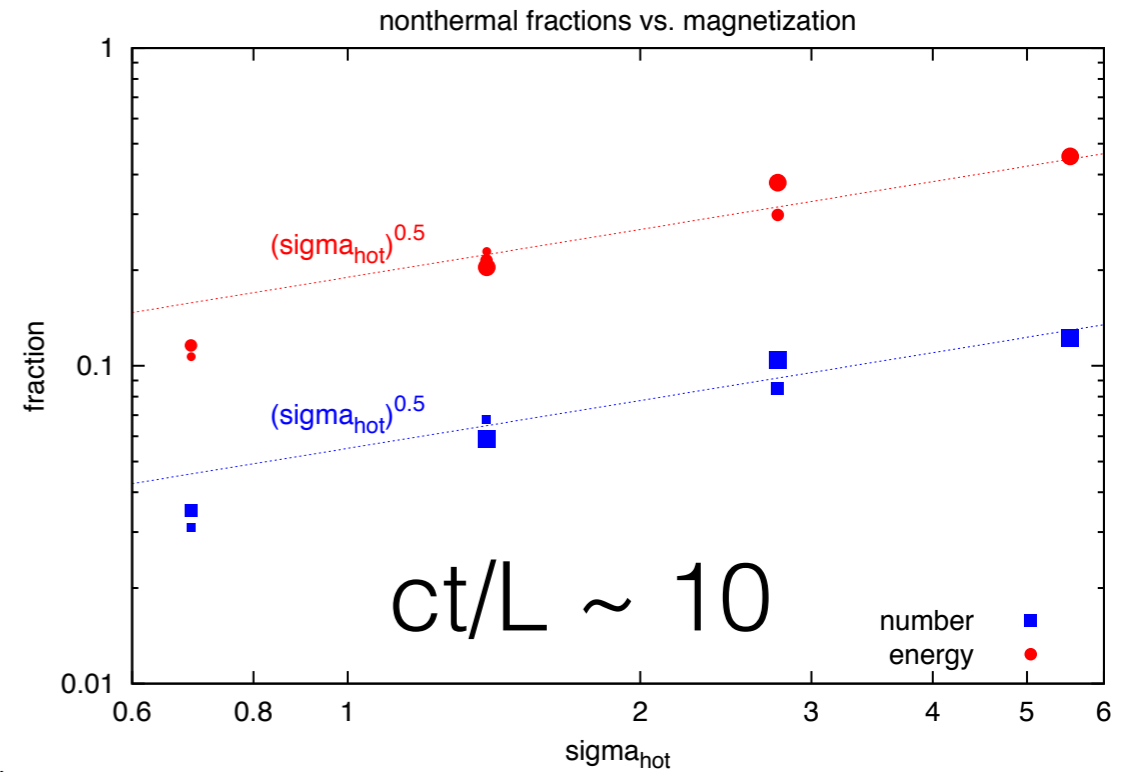
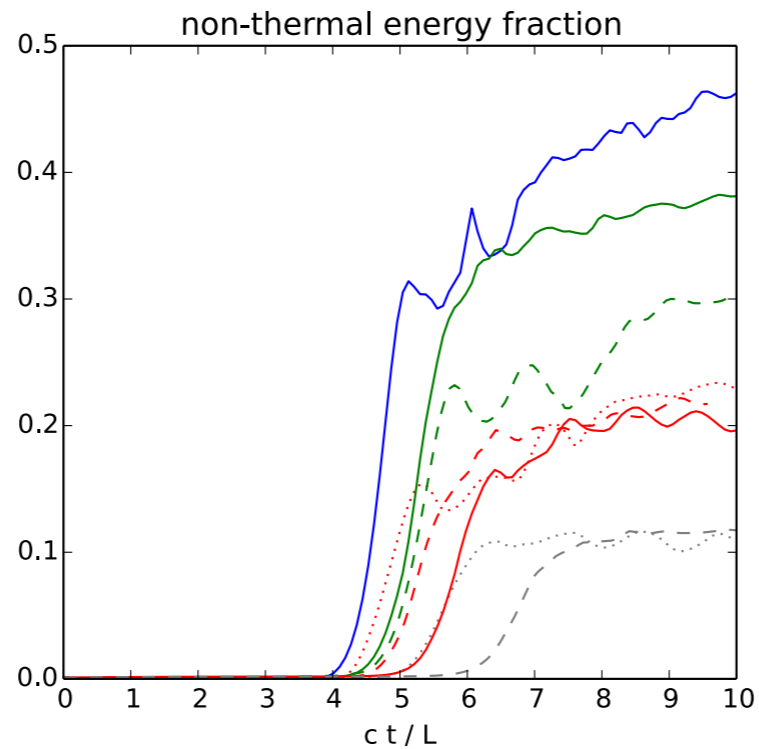
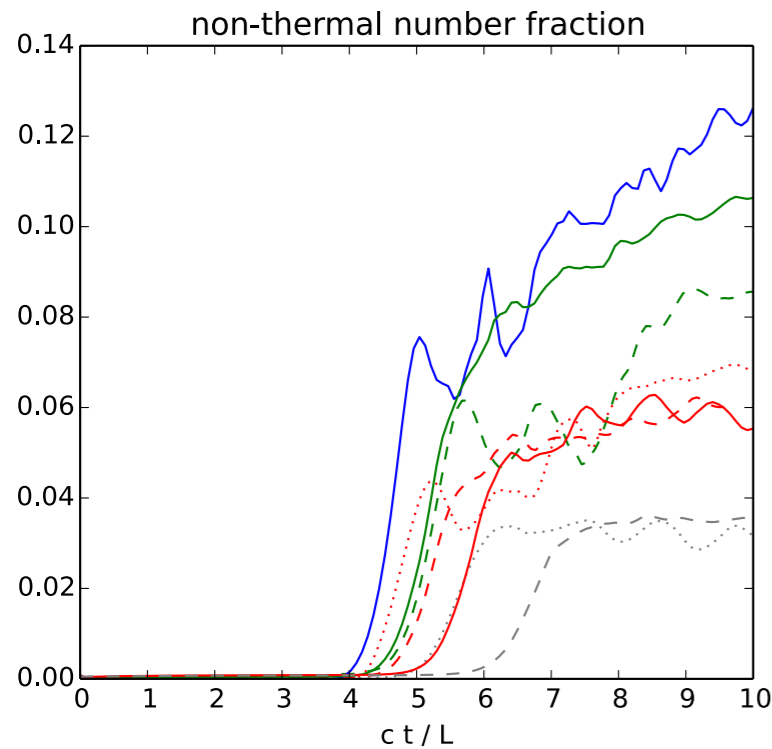


- steady direct acceleration in the linear phase
- stochastic acceleration in the non-linear phase produces a power-law



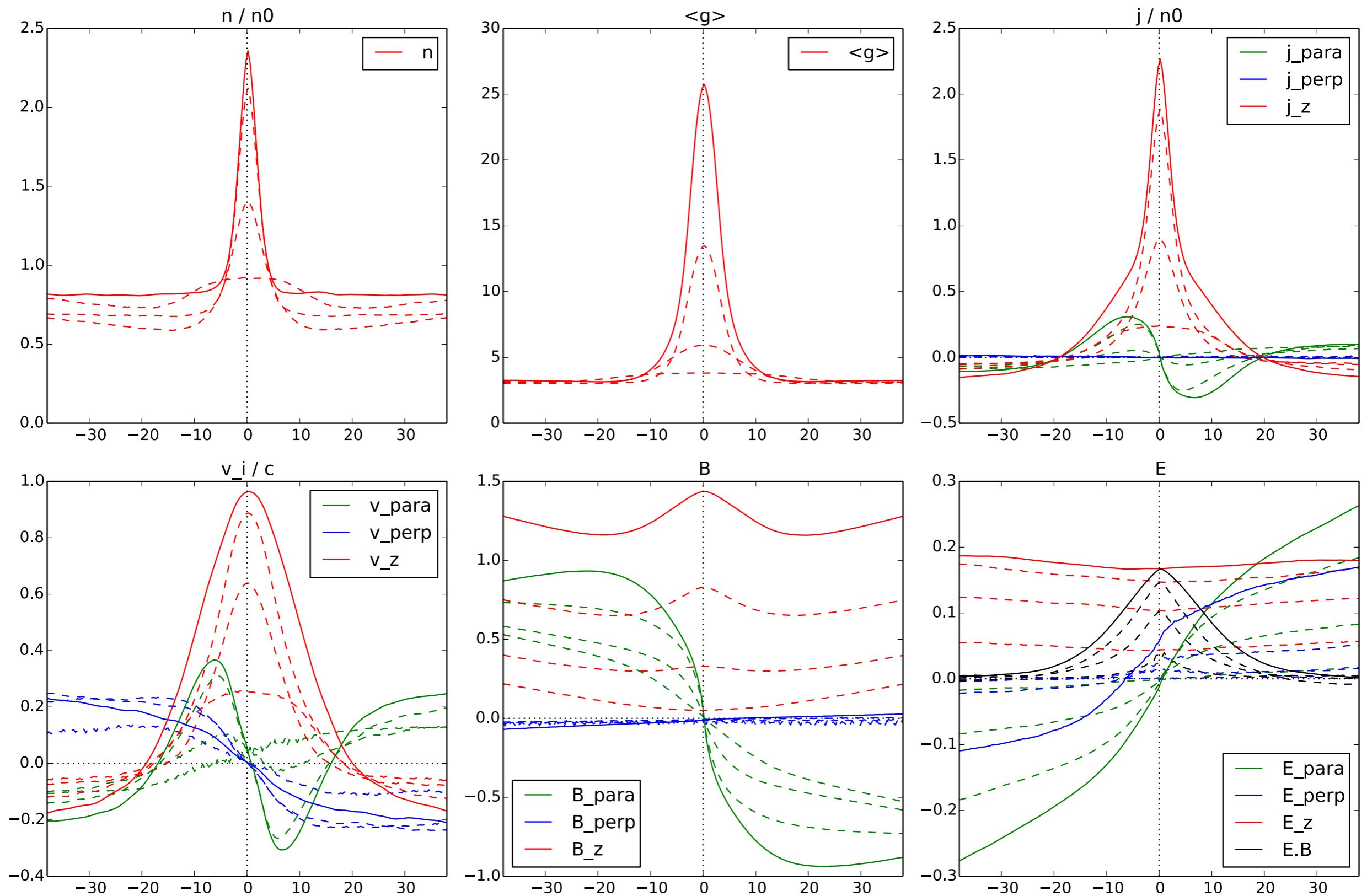
high-energy bump

- particle number and energy fraction beyond the Maxwellian component
- both fractions systematically increase with the magnetization



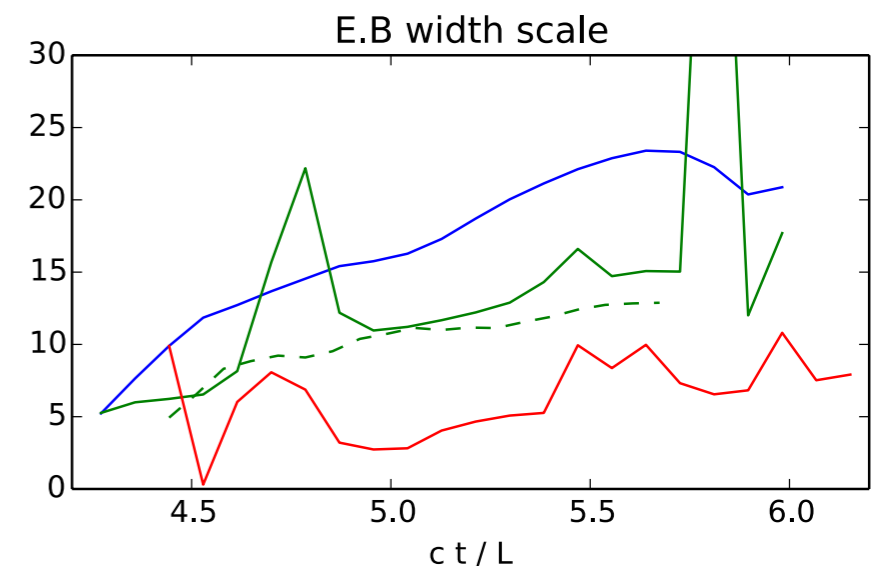
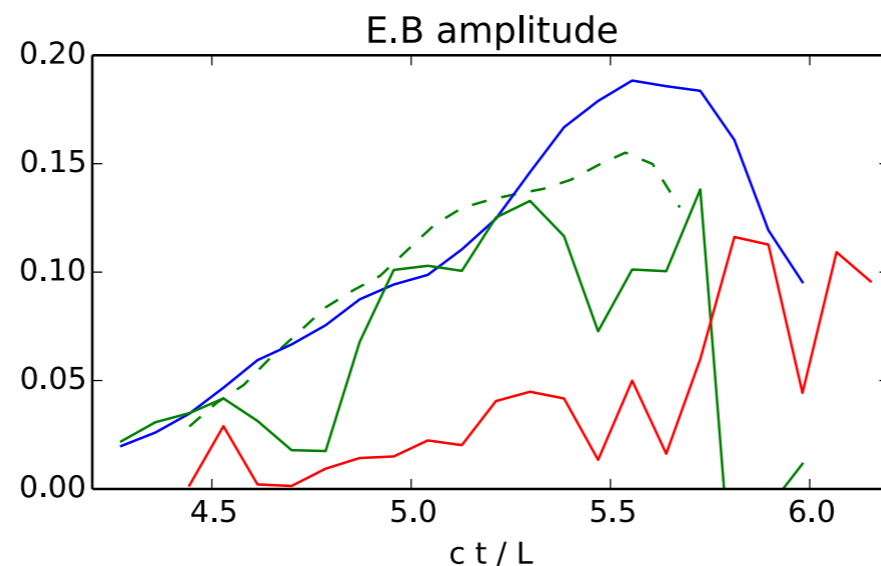
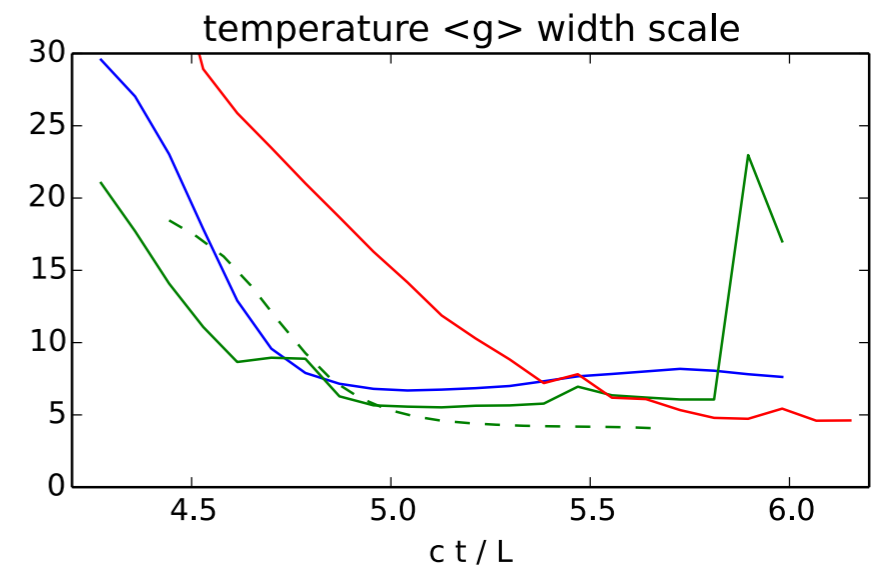
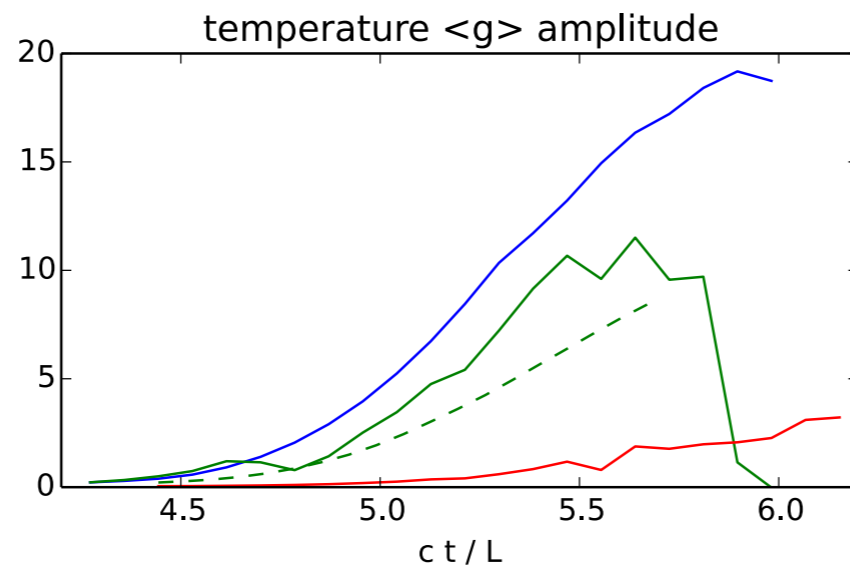
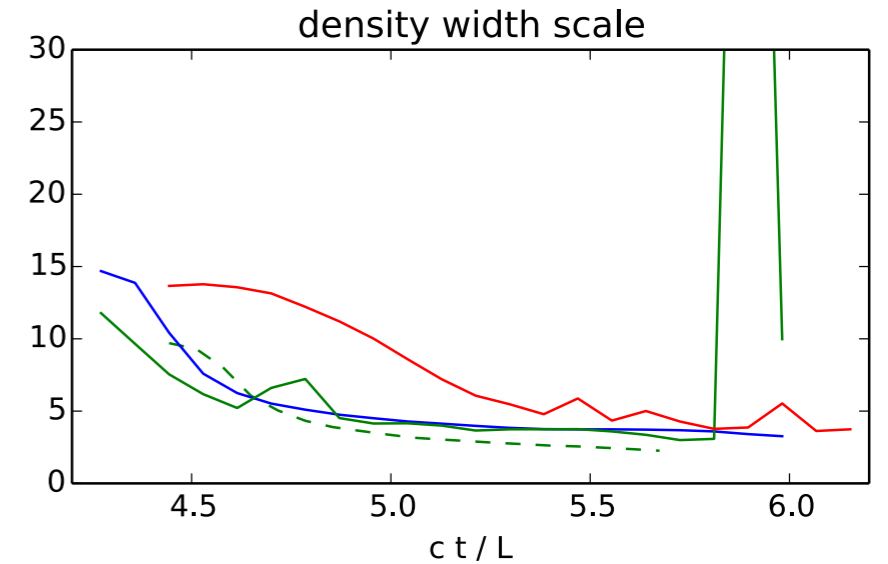
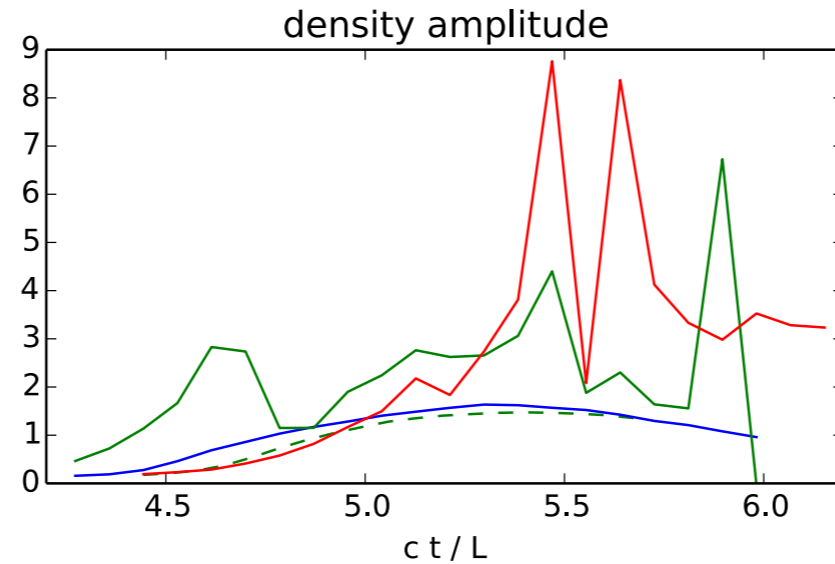
- maximum particle energy measured at $u^2 N(u) = 10^{-3}$

structure of current layers



evolution of current layers

- density width scale consistent with the skin-depth
- E.B width scale consistent with the gyro radius
- E.B volume increasing with the magnetization



summary

- simulations of unstable harmonic magnetic equilibria provide a generic model of efficient particle acceleration in magnetic dissipation
- the current layers forming in the linear instability phase are evolving on dynamical time scale, and very different from the Harris equilibria
- the efficiency of magnetic dissipation is governed by the conservation of magnetic helicity
- the efficiency of particle acceleration depends on magnetization, as it regulates the volume of non-ideal electric fields