

The Effective Strength of Gravity and the Scale of Inflation

Subodh P. Patil

University of Geneva

XXVIII Texas Symposium, December 14 2015

Based on:

I. Antoniadis, S.P.Patil; arXiv:1410.8845

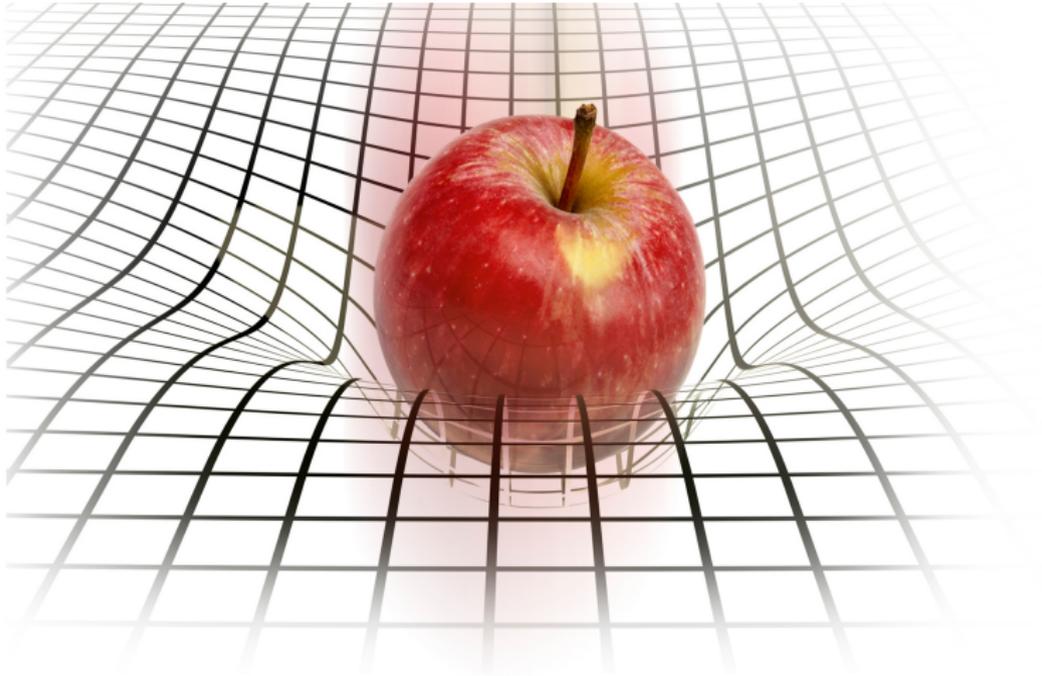
I. Antoniadis, S.P.Patil; arXiv:1510.06759

What is the strength of gravity?

The Effective Strength
of Gravity and the
Scale of Inflation

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The Effective Strength
of Gravity



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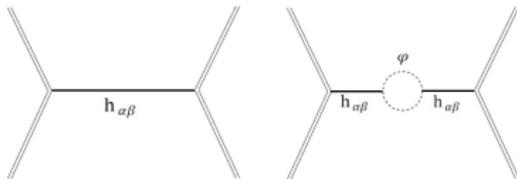
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- ▶ N.B. There are many subtleties and caveat emptors when dealing with running quantities in EFT of gravity.
(cf. [Anber, Donoghue, arXiv:1111.2875](https://arxiv.org/abs/1111.2875); [Bjerrum-Bohr et al, arXiv:1505.04974](https://arxiv.org/abs/1505.04974))

The Strong Coupling Scale

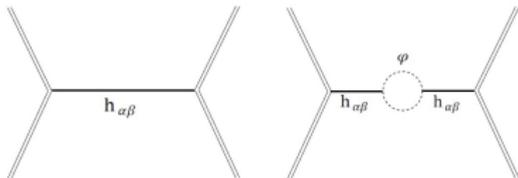
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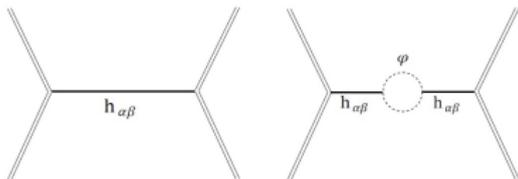
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- ▶ On a Minkowski background: [Dvali, Redi, arXiv:0710.4344](#)

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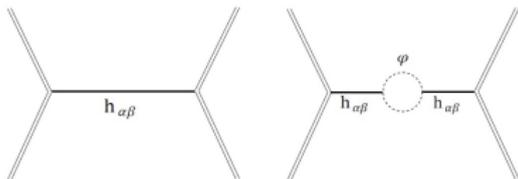
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- ▶ $c := N = \frac{4}{3} N_\phi + 8N_\psi + 16N_V$ Duff, Nucl. Phys. B 125, 334 (1977)

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- ▶ Comparison with the free propagator $1/(p^2 M_{pl}^2)$ implies that the perturbative expansion fails at $p = M_{**}$ where $M_{**} \sim \frac{M_{pl}}{\sqrt{N}}$.

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Generalized to curved backgrounds, places bounds on maximum allowed curvature.

- ▶ $S = \frac{M_{Pl}^2}{2} \int d^4x \sqrt{-g} R + \int d^4x \sqrt{-g} [c_1 R^2 + c_2 R^{\mu\nu} R_{\mu\nu}] + \dots$
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- ▶ ... but also get contributions from higher curvature terms s.t.

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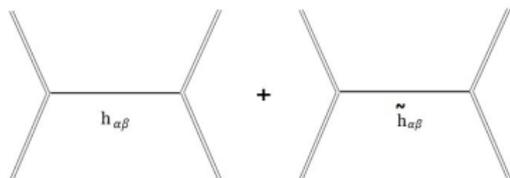
$$M_{pl}^2 \rightarrow M_{pl}^2 \left(1 + c \frac{H^2}{M_{pl}^2} + \dots \right)$$

- ▶ Expansion breaks down when $H^2 \sim M_{pl}^2/N$.
- ▶ Corollary: it is not possible to consistently *infer* a scale of inflation H greater than M_{pl}/\sqrt{N} .

The Strength of Gravity

The *strength* of gravity M_* is an independent quantity¹. Provided we are below M_{**} , any universally coupled species will also affect the *effective strength* of gravity depending on the process in question (equivalence principle, in general violated).

- ▶ KK gravitons do so universally for all conserved sources.

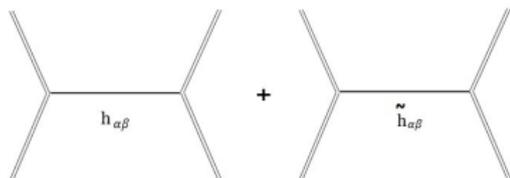


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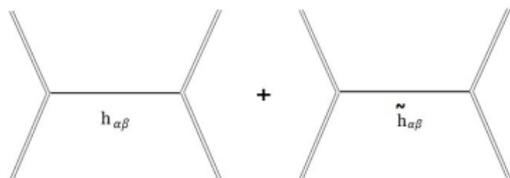
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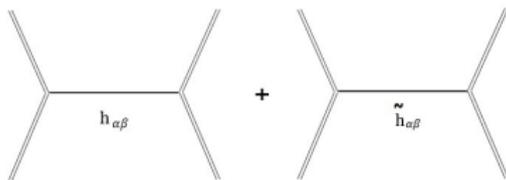
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- ▶ $M_*^2 = M_{pl}^2 / N_*$, N_* counts the number of species with masses below the momentum transfer of the process in question.

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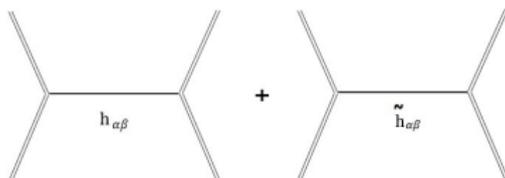
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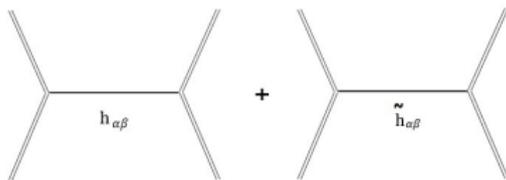
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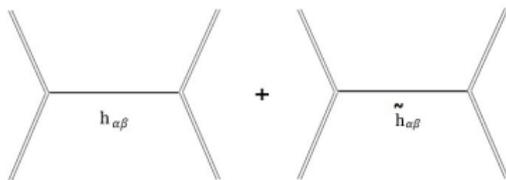
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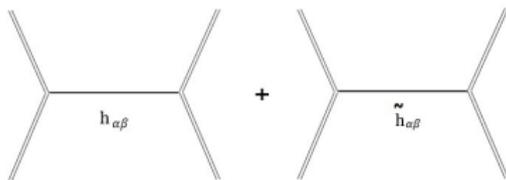
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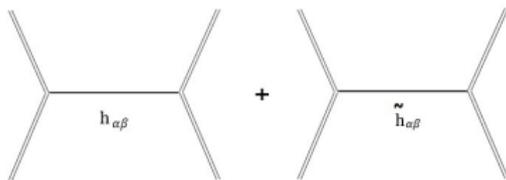
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- ▶ Non-minimal couplings do the same. $M_* = M_{pl}/N_*$, N_* a (process dependent) weighted index.

The Strength of Gravity and the Scale of Inflation

It is widely understood that any detection of primordial tensor modes in the CMB determines the scale of inflation. Or does it?

- ▶ In foliation where inflaton fluctuations are gauged away (comoving/ unitary gauge):

$$h_{ij}(t, x) = a^2(t)e^{2\mathcal{R}(t,x)}\hat{h}_{ij}; \quad \hat{h}_{ij} := \exp[\gamma_{ij}], \quad \partial_i\gamma_{ij} = \gamma_{ii} = 0$$

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- ▶ But what is M_{*s} ?

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- ▶ Consider a non-minimally coupled species with a mass between 10^{-2}eV and the putative scale of inflation:

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- ▶ $g_{\mu\nu} = \left(1 - \frac{2\xi\eta^2}{M_{pl}^2} \right)^{-1} \tilde{g}_{\mu\nu} := F^{-1}(\eta)\tilde{g}_{\mu\nu}$

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- ▶ Naively: $\mathcal{L}_{\text{eff}} \supset \frac{M_{pl}^2}{2} \left(1 - \frac{2\xi\eta^2}{M_{pl}^2}\right) R$

- ▶ $g_{\mu\nu} = \left(1 - \frac{2\xi\eta^2}{M_{pl}^2}\right)^{-1} \tilde{g}_{\mu\nu} := F^{-1}(\eta)\tilde{g}_{\mu\nu}$

- ▶ $\mathcal{L}_{\text{eff}} \supset \frac{M_{pl}^2}{2} \tilde{R} + F^{-2}(\eta)\mathcal{L}_m [F^{-1}(\eta)\tilde{g}_{\mu\nu}, \psi, A_{\mu}, \phi, \eta] + \dots$

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- ▶ $M_{*s}^2 = M_{\text{pl}}^2 \frac{F^2(\eta_*)}{F^2(\eta_0)} \approx \frac{M_{\text{pl}}^2}{1+N_*}$; where $\tilde{N}_* := \sum_i g_i$ where for example, $g_i = 2\Delta(\xi_i\eta_i^2)/M_{\text{pl}}^2$.

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$$V_*^{1/4} = \frac{M_{pl}}{\sqrt{N_*}} \left(\frac{3\pi^2 \mathcal{A} r_*}{2 \cdot 10^{10}} \right)^{1/4}; \quad N_* \sim (1 + N_{KK}^T)(1 + N_s)$$

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- ▶ * N.B. Higuchi bound: massive spin 2 fields inconsistent on dS backgrounds s.t. $m^2 < 2H^2$. However inflation is not dS, KK evidently gravitons are not described by the Fierz-Pauli Lagrangian *in progress*.

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- ▶ For putatively low scale inflationary models, extra species can bring down the scale of inflation even further. Collider constraints?
- ▶ Given $M_* \leq M_{pl}$ and that $M_{**} = M_{pl}/N$, N being *total* number of species, can infer the absolute bound $N \leq \frac{9.15}{r_*} \times 10^7 \left(\frac{M_{pl}^2}{M_{*T}^2} \right)$