Solving the cosmological « lithium problem » with a sterile neutrino

A loophole to the standard theory of electromagnetic cascade

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Talk based on

PRL. 114 (2015) 9, 091101 PRD. 91 (2015) 10, 103007

In collaboration with Pasquale D. Serpico (LAPTh)

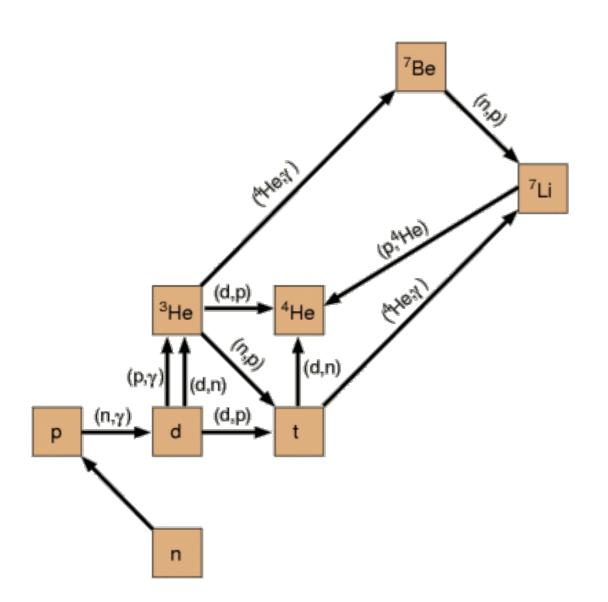


Texas Symposium, Geneva december 15, 2015



Big Bang Nucleosynthesis in a nutshell

- Happened 10 200 s after the BB when the Universe had T = [30,70] keV
- Main nucleus form ⁴He: $Y_p = 4n_{4He}/n_B \approx 0.25$, others $\mathcal{O}(10^{-5} 10^{-10})$



Only one free parameter:

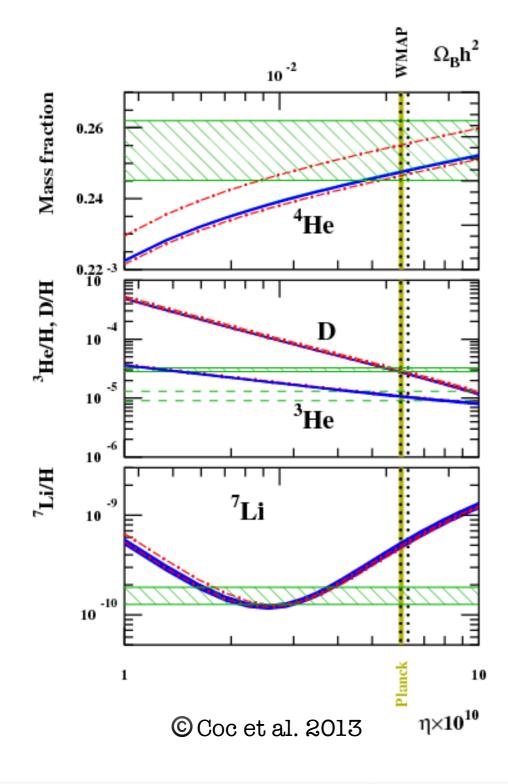
The photon-to-baryon ratio

$$\eta \equiv \frac{n_b}{n_\gamma} \sim 6 \times 10^{-10}$$

=> All abundances can be computed using numerical algorithm such as PArthENoPE

A typical reaction network, © Achim Weiss

Main results



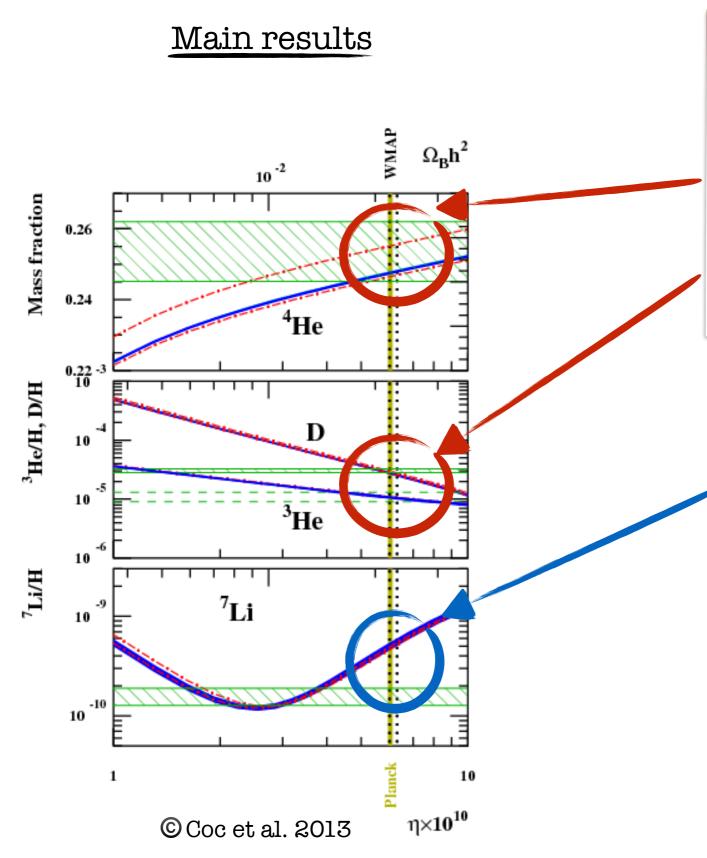
Main results WMAP $\Omega_{\rm B} h^2$ Mass fraction 0.26 0.24 ⁴He 0.22 -3 10 ³не/н, D/H ³He 10 7Li/H ⁷Li 10 -10) 10 $\eta{\times}10^{10}$ © Coc et al. 2013

For 3 nuclei:

Strong observational constraints

$$Y_p > 0.2368$$

 $2.56 \times 10^{-5} < {}^{2}\text{H/H} < 3.48 \times 10^{-5}$
 ${}^{3}\text{He/H} < 1.5 \times 10^{-5}$



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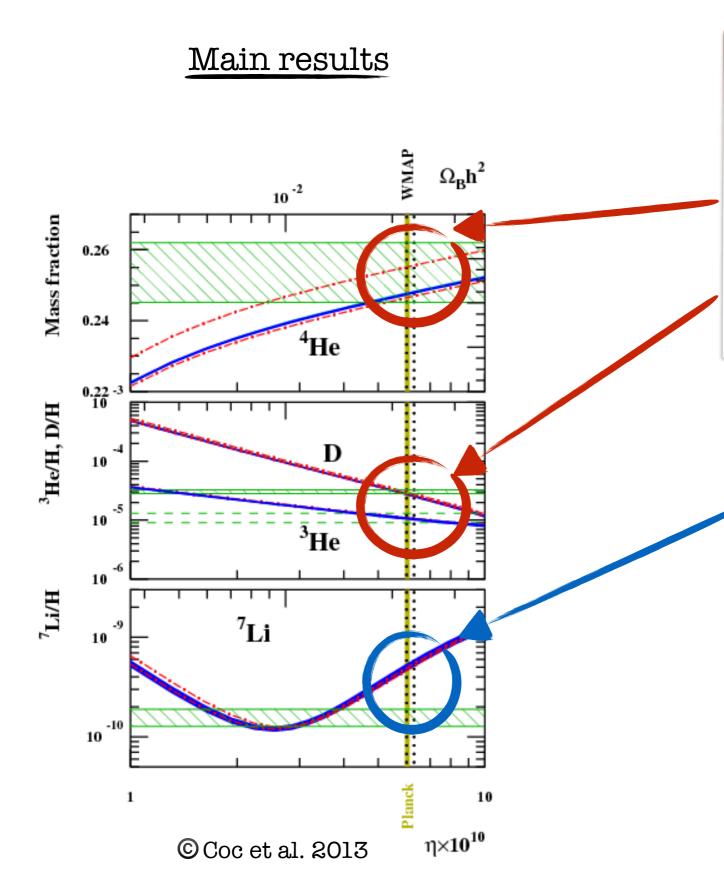
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The Lithium problem:

Overprediction of the ⁷Li abundance

$$Y_{\mathrm{Li}}^{\mathrm{theo}} \simeq 3 \times Y_{\mathrm{Li}}^{\mathrm{obs}}$$



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The Lithium problem:

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Lithium is indirectly produced!

$$^{3}\mathrm{He} + ^{4}\mathrm{He} \rightarrow ^{7}\mathrm{Be} + \gamma$$
 followed by

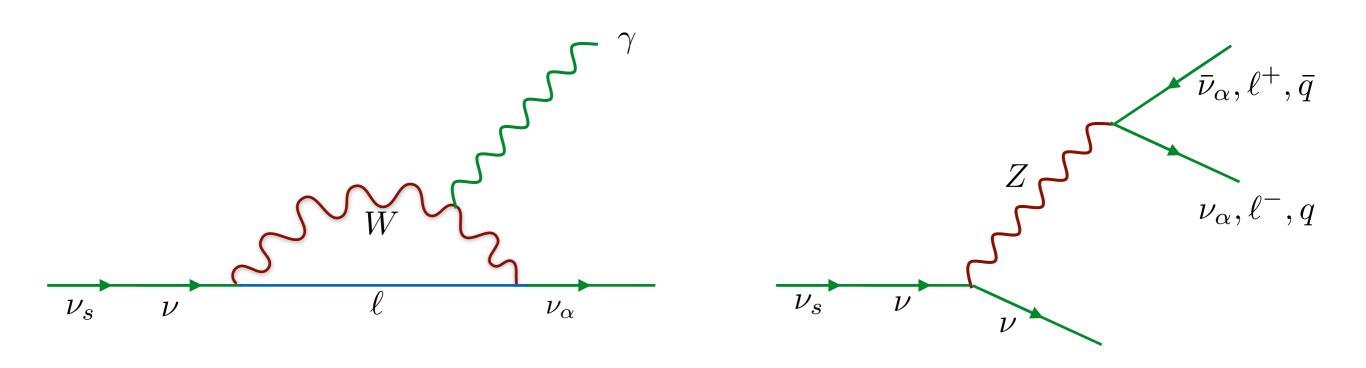
$$^{7}\mathrm{Be} + e^{-} \rightarrow ^{7}\mathrm{Li} + \nu_{e}, \ \tau_{\mathrm{Be}} \sim 53\mathrm{d}$$

One has to destroy the Beryllium!

see e.g. hep-ph/0308083 hep-ph/0008138 astro-ph/0410175

Sterile Neutrino and the BBN

- Modification of N_{eff} affects the expansion rate and the BBN outcome;
- If coupling $\nu_s \leftrightarrow \nu_e$, modification of the weak rates affects the n-p equilibrium which (mostly) sets the ⁴He abundance;
- Eventually, creation of a lepton asymmetry influencing BBN;
- Decay products directly interacting with nuclei can modify BBN yields.



An « old » problem:

Cyburt et al. Phys. Rev. D67 103521, 2003; Jedamzik, ,Phys. Rev. D74 103509, 2006; arXiv:0809.0631;arXiv:1403.5995...

- Big constraints from other nuclei
- Big constraints from entropy production and spectral distortions

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$$^{7}\mathrm{Be} + \gamma \rightarrow {}^{4}\mathrm{He} + {}^{3}\mathrm{He}$$

 $E_{\text{threshold}}(\text{Be}) = 1.58 \text{ MeV}$ $E_{\text{threshold}}(\text{De}) = 2.2 \text{ MeV}$

One « trick » : if $1.6 < E_0 < 2.2$ MeV it is possible to avoid all BBN constraints!

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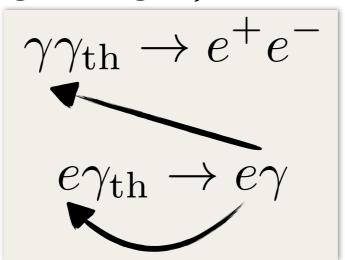
However, this was known to fail, why would it work now?

Electromagnetic Cascade in a nutshell

We want to describe electromagnetic energy injection in a plasma of photons (very few e+e-, nuclei):

what is the resulting metastable distribution of photons?

Basic processes are (at high energies)



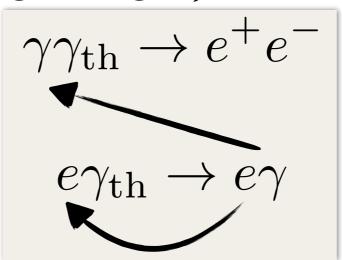
Particle multiplication and energy redistribution

Electromagnetic Cascade in a nutshell

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Particle multiplication and energy redistribution

The first process has a threshold, below it

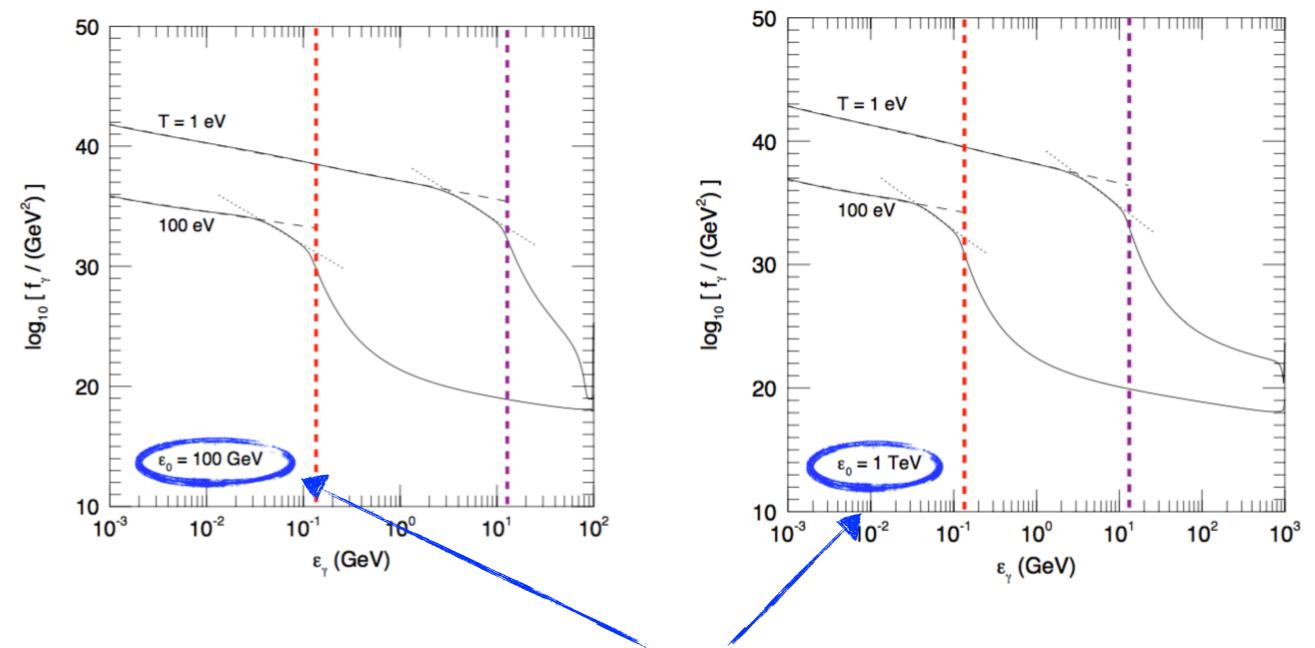
$$\gamma \gamma_{\rm th} \to \gamma \gamma$$

and eventually (very low rates)

$$\gamma N \to eN$$

$$\gamma e_{\rm th} \to \gamma e$$

Kawasaki & Moroi, ApJ 452,506 (1995) This has been shown to lead to a universal spectrum



- Shape independent of the energy / temperature of the bath:
 Only dictates the overall normalisation;
- Threshold due to pair production.

$$E_{\rm cutoff}(1~{\rm keV}) \sim 12~{\rm MeV}$$
 $E_{\rm cutoff}(10~{\rm eV}) \sim 1.2~{\rm GeV}$

All cases simulated inject energy such that $E_{\gamma} \gg E_{\rm cutoff}$ => « Theoritical prejudice »!

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After « standard » BBN : $E_{\rm threshold}({\rm Be})$ = 1.58 MeV < $E_{\rm cutoff}$

 $\label{eq:energy_energy} If \ E_{\rm threshold} < E_0 < E_{\rm cutoff} \\ \ results \ in \ the \ literature \ are \ wrong \ !$

$$\gamma \gamma_{\rm th} \to \gamma \gamma, \ \gamma e_{\rm th}^{\pm} \to \gamma e^{\pm}, \ \gamma N \to N e^{\pm}$$

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Relevant Boltzmann equation writes:

$$\frac{\partial f_{\gamma}(E_{\gamma})}{\partial t} = -\Gamma_{\gamma}(E_{\gamma}, T(t)) f_{\gamma}(E_{\gamma}, T(t)) + \mathcal{S}(E_{\gamma}, t)$$

whose stationary solution is

$$f_{\gamma}^{\mathrm{S}}(E_{\gamma}) = \frac{\mathcal{S}(E_{\gamma}, t)}{\Gamma_{\gamma}(E_{\gamma}, t)}$$

Hubble rate much smaller than all particle physics interaction rate, thus neglected

where for a decaying particle

$$S(E_{\gamma}, t) = \frac{n_{\gamma}^{0} \zeta_{X} (1 + z(t))^{3} e^{-t/\tau_{X}}}{E_{0} \tau_{X}} p_{\gamma}(E_{\gamma}, t)$$

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exact at the end-point, then iterate

$$S(E_{\gamma}, t) \to S(E_{\gamma}, t) + \int_{E_{\gamma}}^{\infty} dx K_{\gamma}(E_{\gamma}, x, t) f_{\gamma}(x, t)$$

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Finally compute nuclei abundances:

$$\frac{dY_A}{dt} = \sum_T Y_T \int_0^\infty dE_\gamma f_\gamma(E_\gamma, t) \sigma_{\gamma+T\to A}(E_\gamma) - Y_A \sum_P \int_0^\infty dE_\gamma f_\gamma(E_\gamma, t) \sigma_{\gamma+A\to P}(E_\gamma)$$

$$Y_A \equiv n_A/n_b$$

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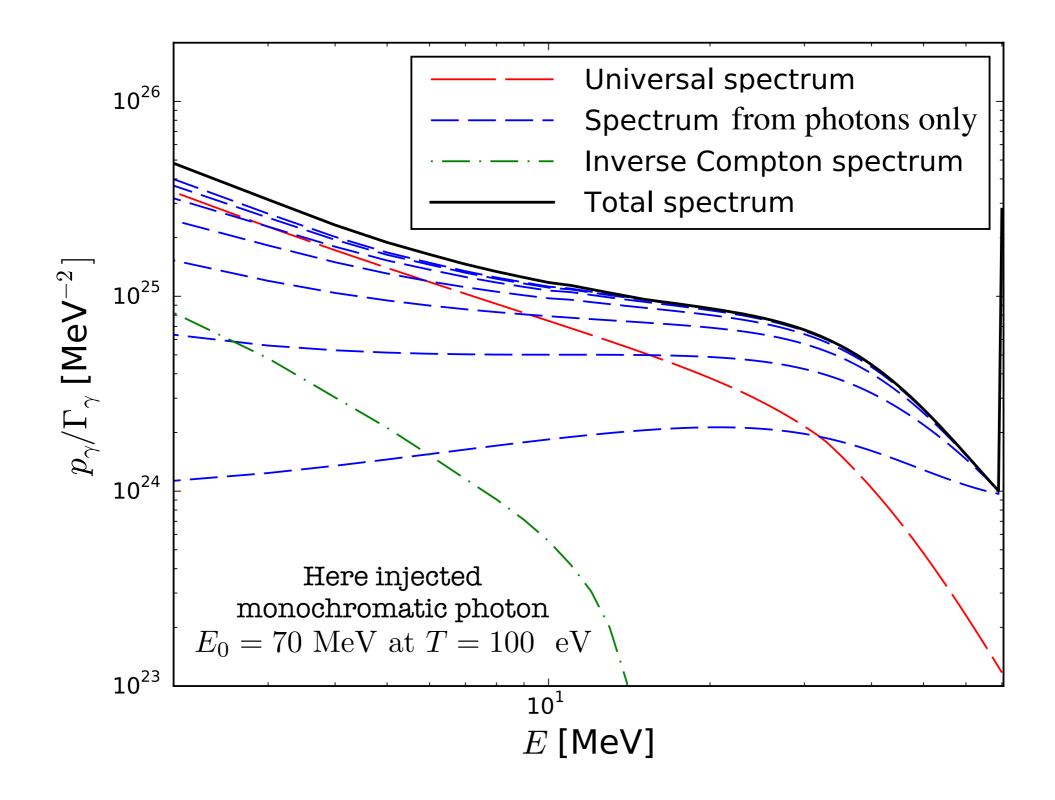
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Production from photodissociation of heavier nuclei

Destruction from its photodissociation

 $Y_A \equiv n_A/n_b$

Typical results for a given energy and a given temperature of the thermal bath

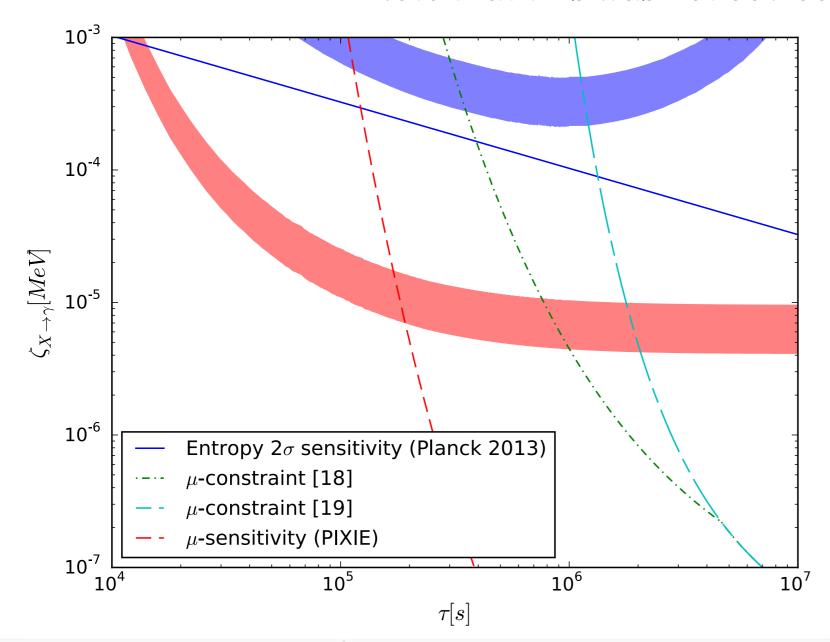


<u>Proof of principle solution</u>: monochromatic photon injection

In our case, it is possible to solve the lithium problem, while fulfilling other constraints.

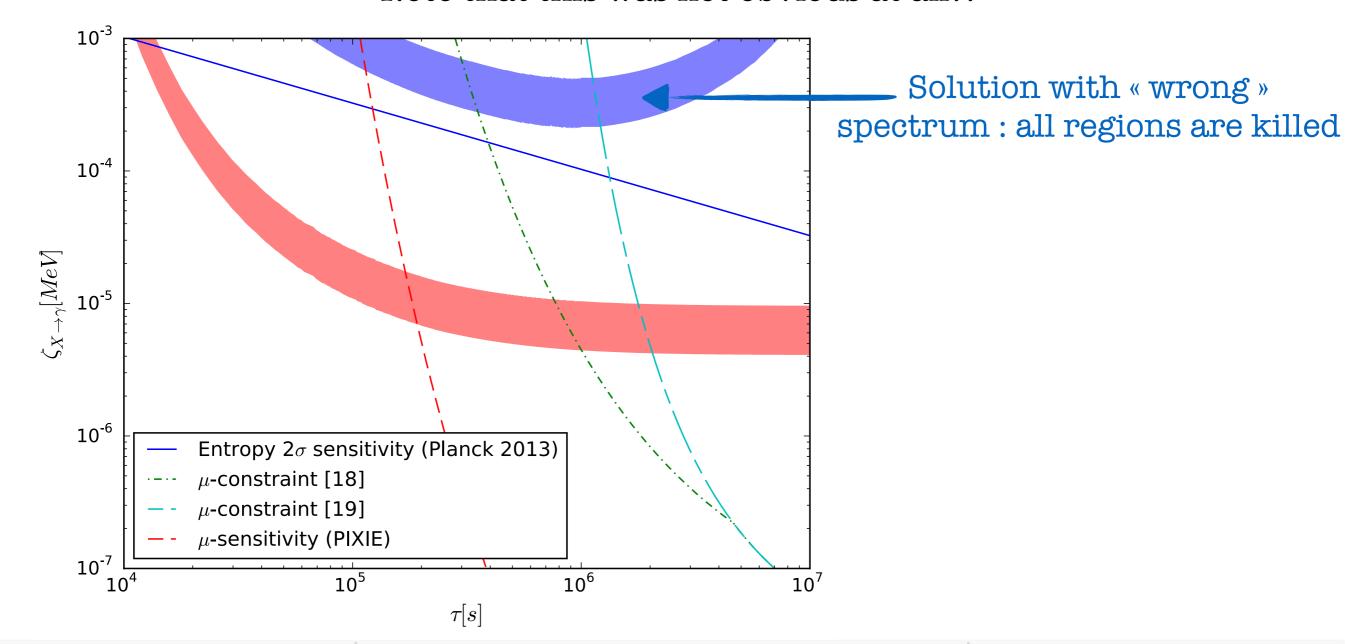
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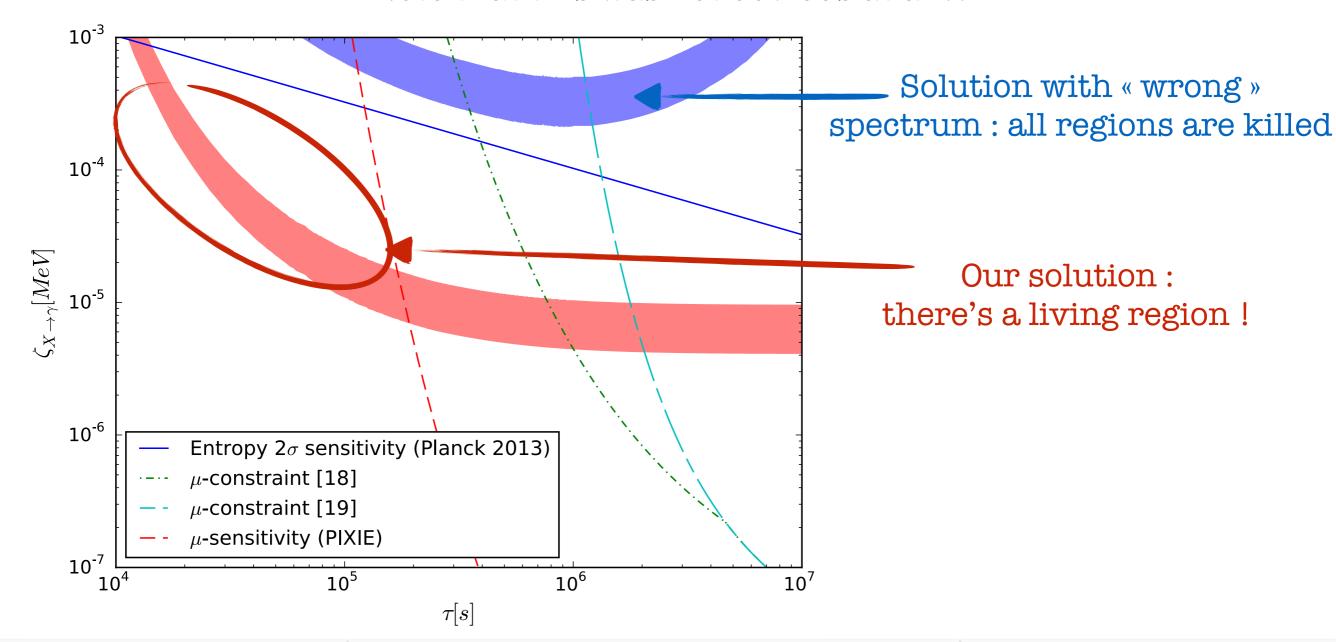
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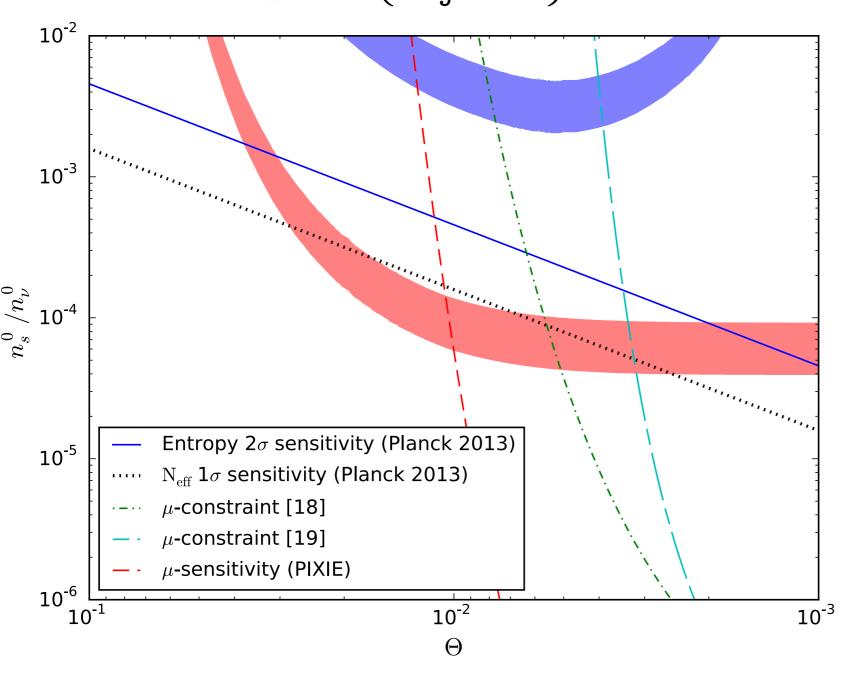


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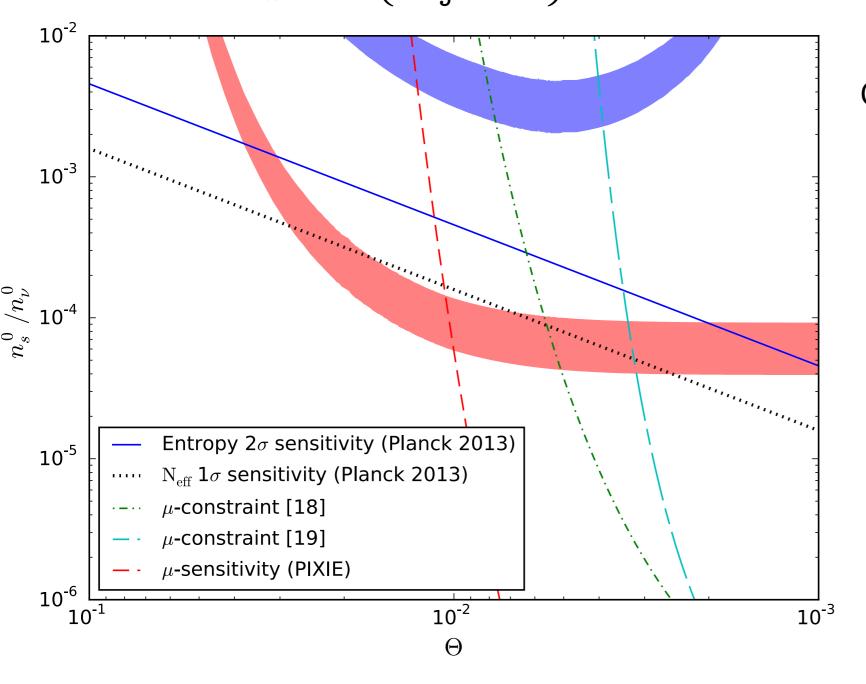
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Try with a « real » model that was known to fail when using universal spectrum : the Sterile (majorana) Neutrino



H. Ishida, M. Kusakabe and H.Okada, PRD 90, 8, 083519 (2014) Try with a « real » model that was known to fail when using universal spectrum : the Sterile (majorana) Neutrino



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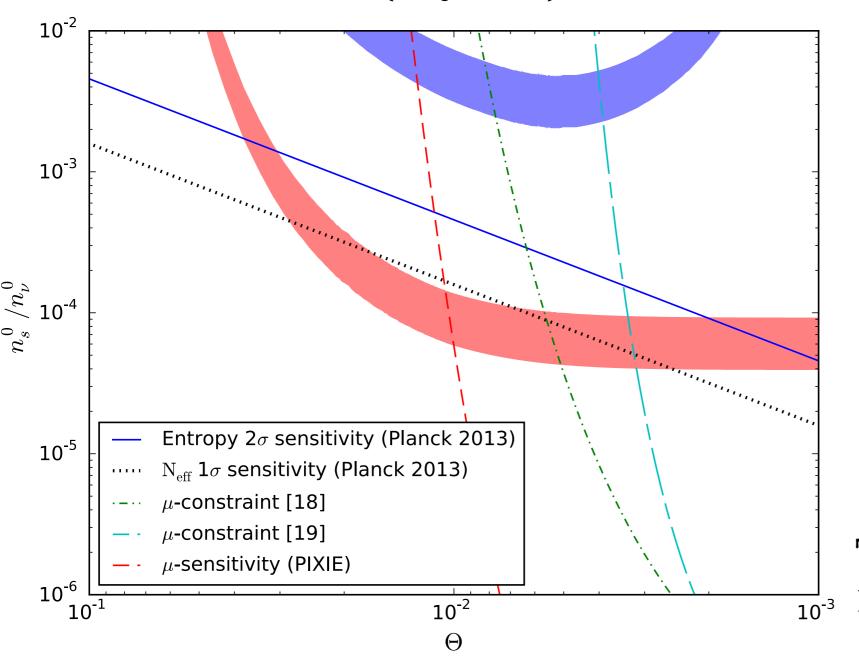
Convert the variables

$$au o \Theta$$

mixing angle

$$\zeta \to n_s^0/n_\nu^0$$

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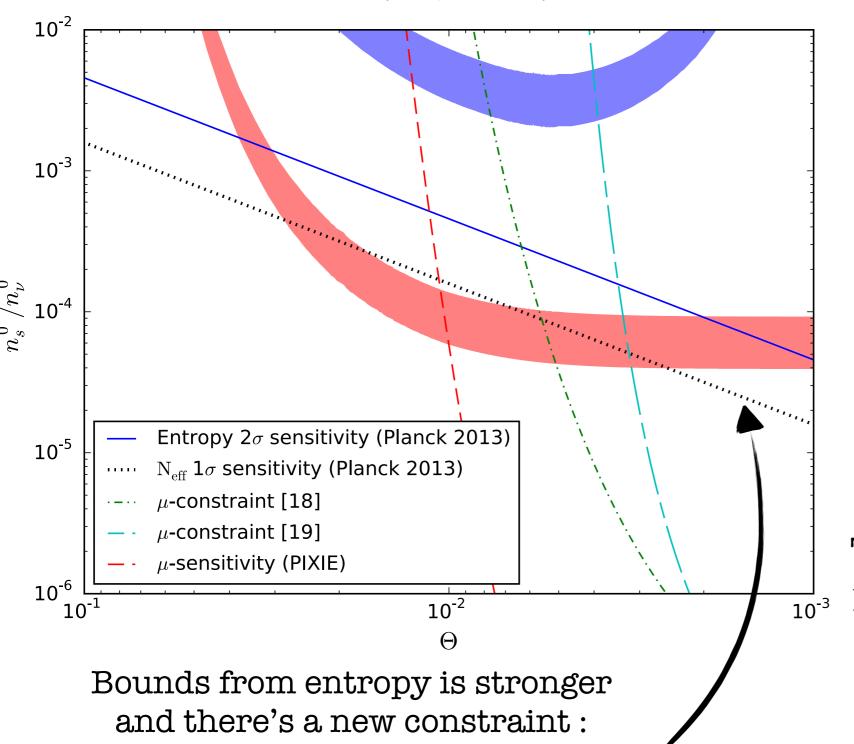
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To avoid constraints from cosmology and labs mixing required to be mostly ν_{μ} or ν_{τ}

Typical branching ratio

$$1:0.1:0.01 \text{ in } 3\nu:\nu e^+e^-:\nu\gamma$$

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variation of N_{eff} (planck sensitivity)

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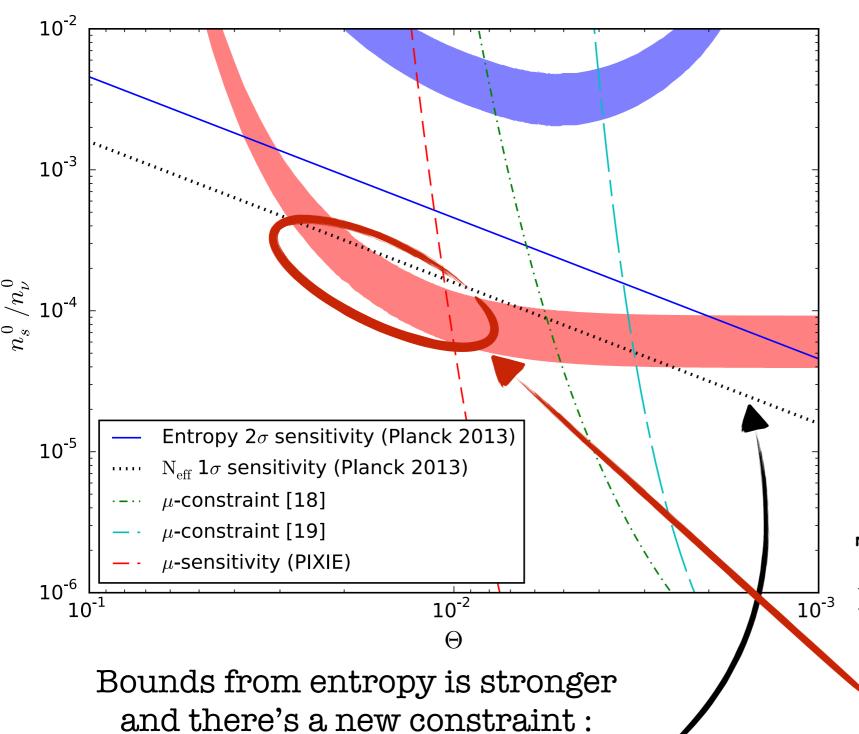
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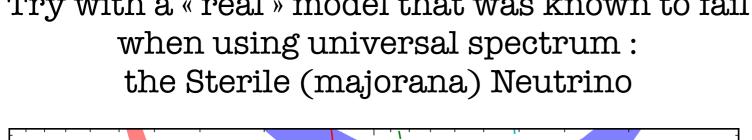
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It works!

10⁻²

10⁻³

Try with a « real » model that was known to fail when using universal spectrum:



H. Ishida, M. Kusakabe and H.Okada, PRD 90, 8, 083519 (2014)

Convert the variables

$$\tau \to \Theta$$

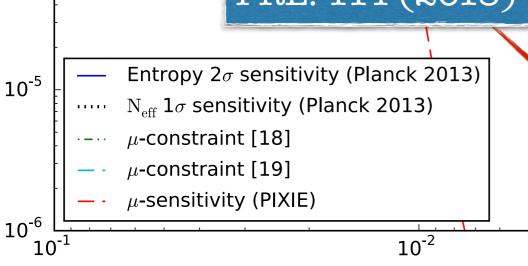
mixing angle

$$\zeta \to n_s^0/n_\nu^0$$

normalise to active neutrino density

More details:

PRL. 114 (2015) 9, 091101 [arXiv:1502.01250]



aints from cosmology and labs mixing required to be mostly ν_{μ} or ν_{τ}

Typical branching ratio

$$1:0.1:0.01 \text{ in } 3\nu:\nu e^+e^-:\nu\gamma$$

Bounds from entropy is stronger and there's a new constraint: variation of N_{eff} (planck sensitivity)

It works!

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- We have shown that the universality hypothesis breaks down.
 The resulting spectrum can be very different from the universal one.

- We have shown how it might ease particle physics (electromagnetic) solution to the lithium problem, as illustrated with the sterile neutrino model.
- The same phenomenon also has important consequences for BBN bounds: they are more stringent and non-universal.

