

# A direct measurement of tomographic lensing power spectra from CFHTLenS

(in press at MNRAS; arXiv1509.04071)

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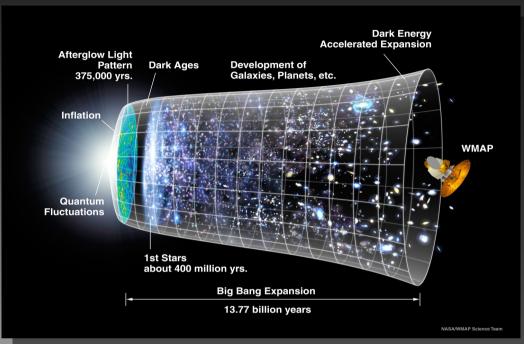
Collaborators: M. Viola, W. Valkenburg, B. Joachimi, H. Hoekstra, K. Kuijken

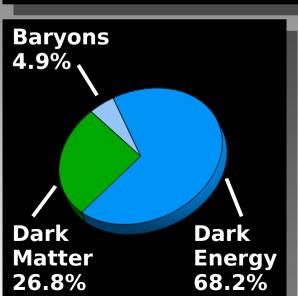
**Texas Symposium 2015, Geneva** 

**17 December 2015** 

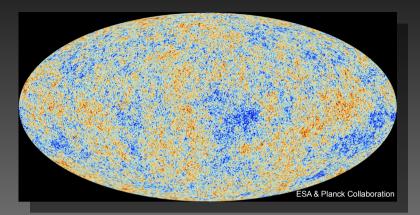
# I. Introduction

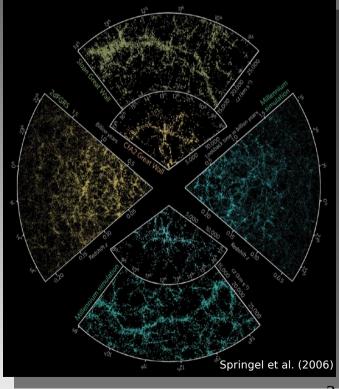
### Cosmological model



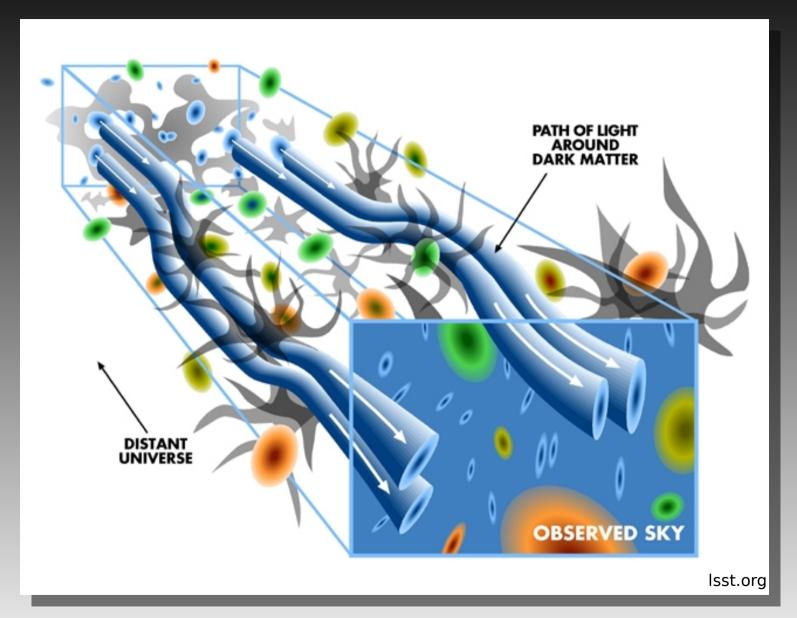






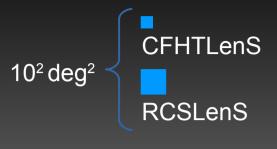


## Lensing statistic:



shear-shear (a.k.a. cosmic shear)

### Weak lensing: Future



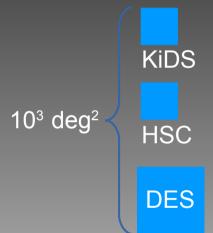
17 gals/arcmin²

6 gals/arcmin<sup>2</sup>

- $z_{\rm m} = 0.75$
- $z_{m} = 0.60$

#### completed

ongoing



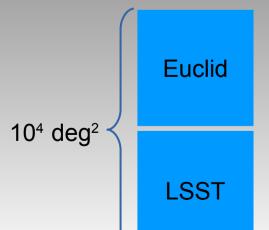
9 gals/arcmin²

 $z_{\rm m} = 0.70$ 

20 gals/arcmin²

• 8 gals/arcmin<sup>2</sup>

- $z_{\rm m} = 1.00$
- $z_{\rm m} = 0.65$



- 30 gals/arcmin²
- $z_{\rm m} = 0.90$

- 31 gals/arcmin<sup>2</sup>
- $Z_{\rm m} = 1.00$

> 2020

### Weak Lensing: Challenges

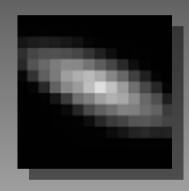
1) Accurate photometric redshifts

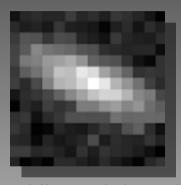
#### 2) Shape noise:











Bridle et al. (2009)



"The bigger (deeper) the survey the smaller the uncertainties!"

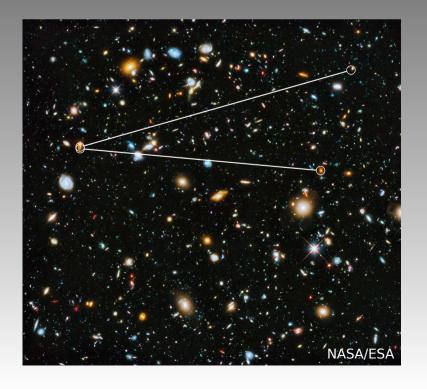
3) Blending (!)

### II. Cosmic Shear

### Lensing of LSS:

#### **Theory:**

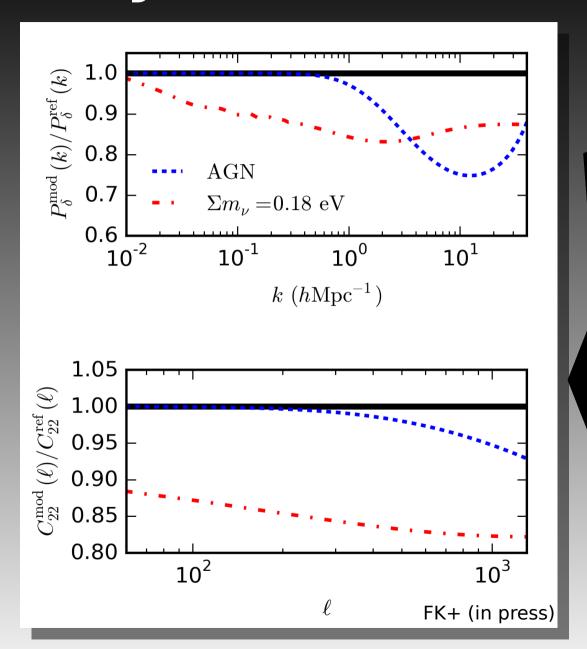
$$C_{\mu\nu}^{\rm EE}(\ell) = \frac{9\Omega_{\rm m}^2 H_0^4}{4c^4} \int_0^{\chi_{\rm H}} \mathrm{d}\chi \underbrace{\frac{g_{\mu}(\chi)g_{\nu}(\chi)}{a^2(\chi)}} P_{\delta}\left(k = \frac{\ell}{f_{\rm K}(\chi)};\chi\right)$$
 "geometry" "physics"



#### measurements:

**correlation functions**  $\Longrightarrow$  power spectra

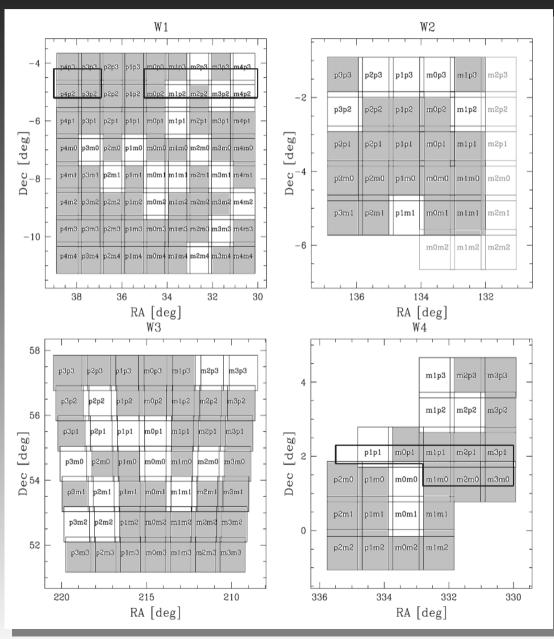
### **Baryons & neutrinos:**



integration over lensing kernel

- AGN feedback from OWLS after Harnois-Déraps et al. (2015)
- 3 degenerate, massive neutrinos with  $\Sigma m_v = 0.18$  eV

### The CFHTLenS case:



 $\sim$ 154 deg<sup>2</sup> ( $\sim$ 115 deg<sup>2</sup>)

 $n_{gal} = 17 \text{ gals/arcmin}^2$ 

two redshift slices:

 $z_1$ :  $0.50 < z \le 0.85$ 

 $z_2$ : 0.85 <  $z \le 1.30$ 

minimize intrinsic alignments

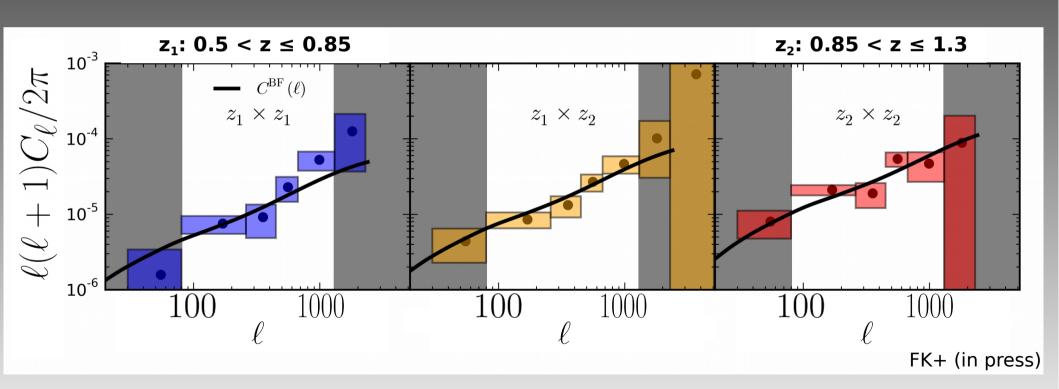
!!! PUBLIC data !!!

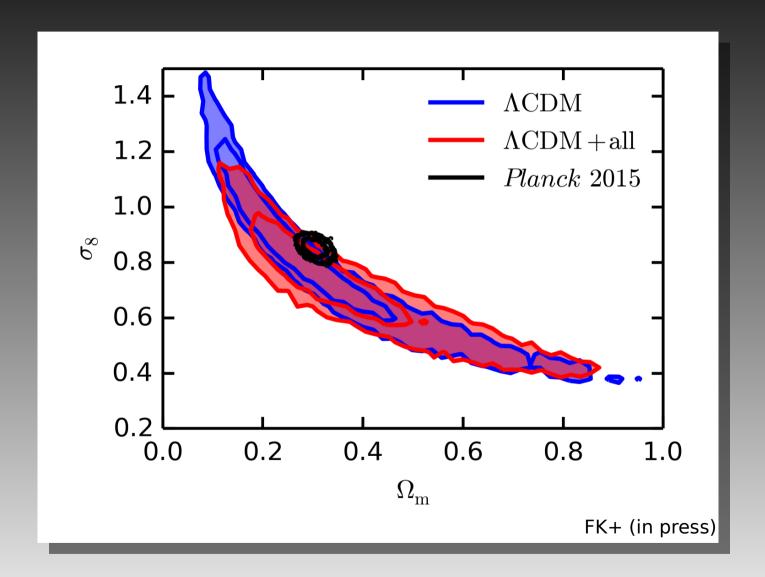
Erben et al. (2012)

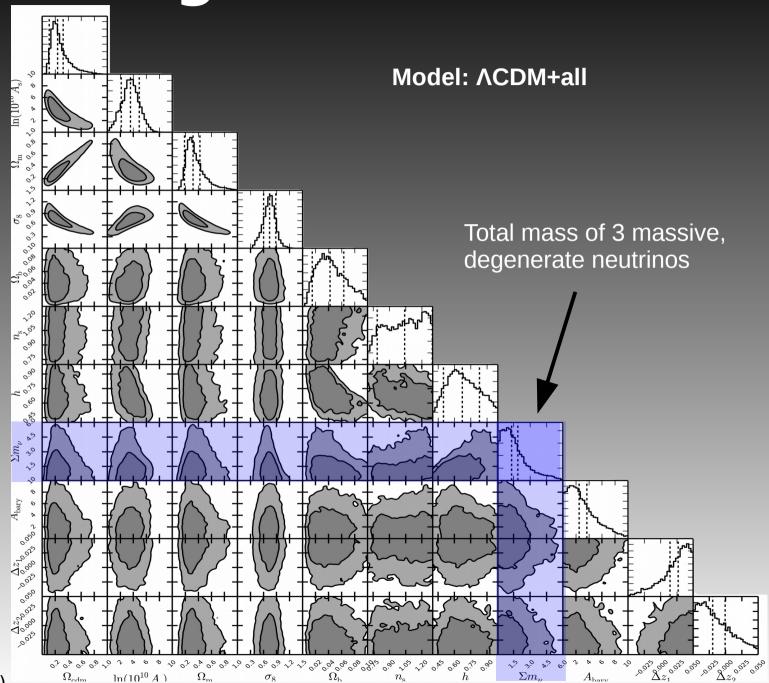
### Results: Multipole Space

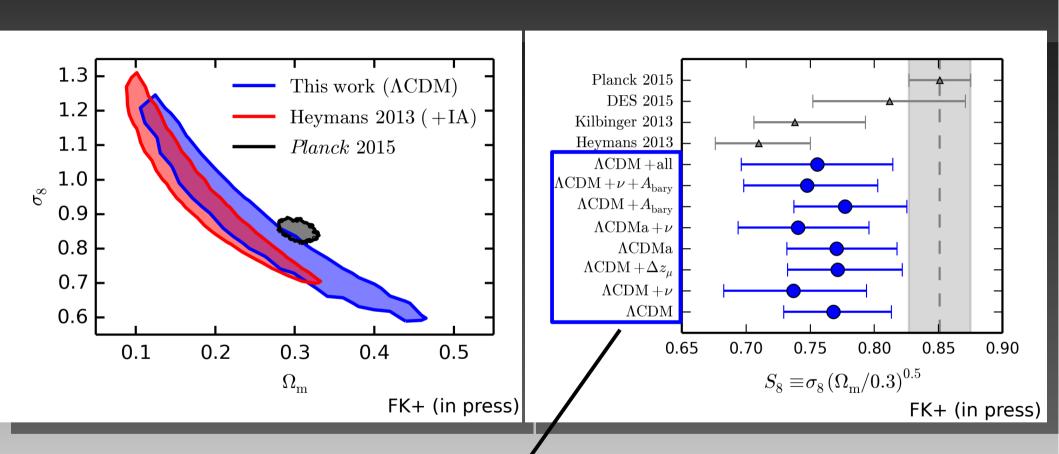
quadratic estimator method (Hu & White 2001) expanded to include photometric redshift bins

WL power spectra from CFHTLenS (W1, W2, W3 & W4 combined with inverse variance weights)









Which model describes the data the best?

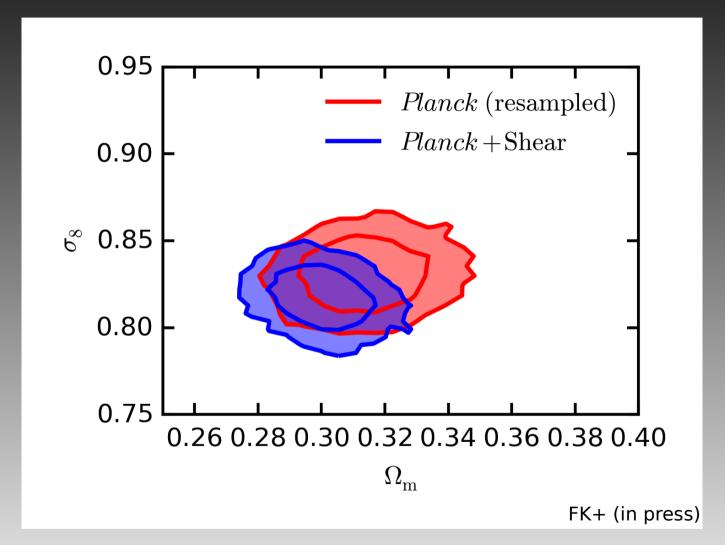
### **Evidences:**

likelihood analysis performed with *Monte Python* (Audren et al. 2012) and *Multinest* (Feroz et al. 2008, 2009, 2013)



Model	$\ln \mathcal{Z}$	$2 \ln K \ (K \equiv \mathcal{Z}_i / \mathcal{Z}_{\Lambda \mathrm{CDM}})$
$\Lambda \mathrm{CDM}$	$-40.96 \pm 0.06$	0
$\Lambda CDMa$	$-41.07 \pm 0.06$	-0.22
$\Lambda \text{CDM} + \nu$	$-41.63 \pm 0.07$	-1.34
$\Lambda CDMa + \nu$	$-41.83 \pm 0.07$	-1.74
$\Lambda$ CDM + $A_{\text{bary}}$	$-41.66 \pm 0.06$	-1.40
$\Lambda \text{CDM} + \nu + A_{\text{bary}}$	$-42.48 \pm 0.07$	-3.04
$\Lambda CDM + \Delta z_{\mu}$	$-40.75 \pm 0.07$	0.42
$\Lambda \text{CDM} + \text{all}$	$-42.19 \pm 0.07$	-2.46

FK+ (in press)





# III. Conclusions

A direct extraction of the lensing power spectrum is a "cleaner" way to compare data with theory.

The power spectrum results show overall consistency with previous results based on correlation-functions.

Ongoing and future lensing surveys have the potential to constrain distinct features in multipole space such as left by massive neutrinos or baryon feedback with high precision.