

Extragalactic circuits, black holes,
and the ultimate energy transfers

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by

Philipp Kronberg

Dept of Physics

University of Toronto, Canada

Abstract

Energy, and electrical circuits in Space, and their connection to supermassive black holes

- A non-negligible fraction of Supermassive Black Hole's (SMBH) rest mass energy gets transported into extragalactic space -- by remarkable processes in jets which are not completely understood. The bulk of the energy flow from the SMBH (e.g. $\gtrsim 10^7 M_{\odot}$) appears to be electromagnetic, rather than *via* a particle beam flux. Also, remarkably, these jets contain current flows that remain largely intact over multi-kpc distances. Accretion disk models have independently calculated that a $\sim 10^8 M_{\odot}$ SMBH should generate $\mathcal{O} 10^{18} - 10^{19}$ Ampères in the vicinity of the SMBH.
- I describe the first and best yet observational estimate of the current flow along the axis of a jet that extends from the nucleus of the active elliptical galaxy in 3C303. This is $I \sim 10^{18}$ Ampères at a projected 40 kpc from the AGN. This points to the existence of cosmic scale electric circuits. The power flow is $P = I^2 Z$, watts, where $Z \sim 30$ Ohms, which is \mathcal{O} the *impedance* of free space $Z(\epsilon_0, \mu_0)$, (ϵ_0, μ_0) being the permittivity and magnetic permeability. These, in turn, uniquely determine c . The electrical potential drop ($\sim 10^{20}$ V) across the jet diameter (which is \approx a few times r_g of the SMBH) is, interestingly \approx that required to accelerate the Ultra High Energy Cosmic Rays (UHECR).
- Jets and high energy outflows have different progenitors, forms, sizes, luminosities, and ambient environments
- This talk focuses on
 1. electromagnetically dominated (Poynting flux) jets from supermassive BH's
 2. located in a rarified intergalactic environment -i.e. not in rich galaxy clusters

$$Z_0 = \frac{3}{c} \beta$$

BH (magnetic + CR) energy output ($\gtrsim 10^{60}$ ergs) is “captured” within a few Mpc,

compare with

η (photons), $\approx 10\%$ of $M_{\text{BH}}c^2$ (not captured) appears comparable to η (CR + B),

2147+816 giant radio galaxy

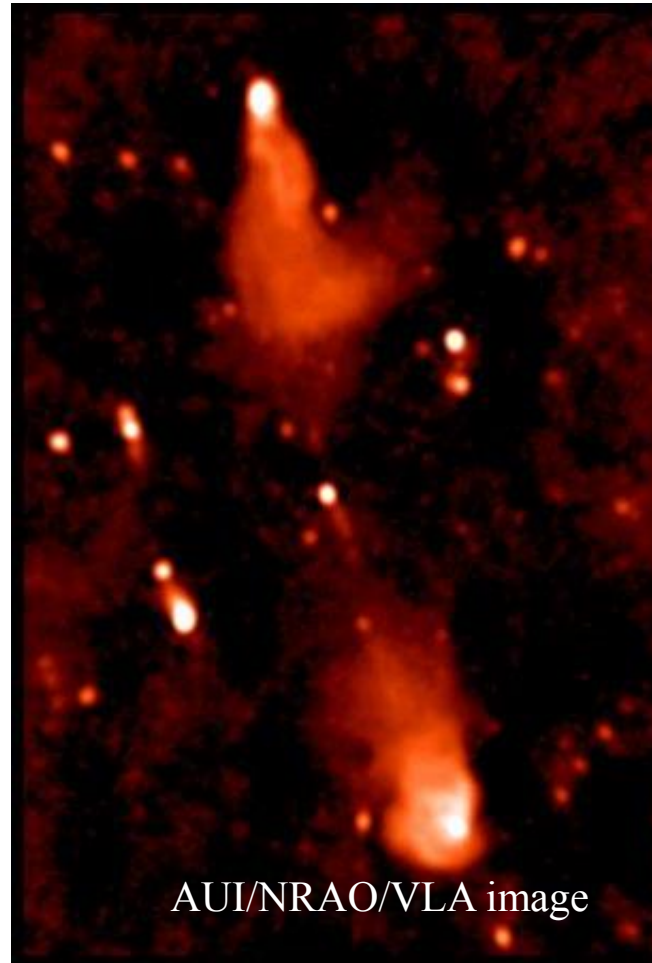
Analysis of ≈ 70 GRG images
Kronberg, Dufon, Li, Colgate
ApJ 2001

$z=0.146$

2.6 Mpc

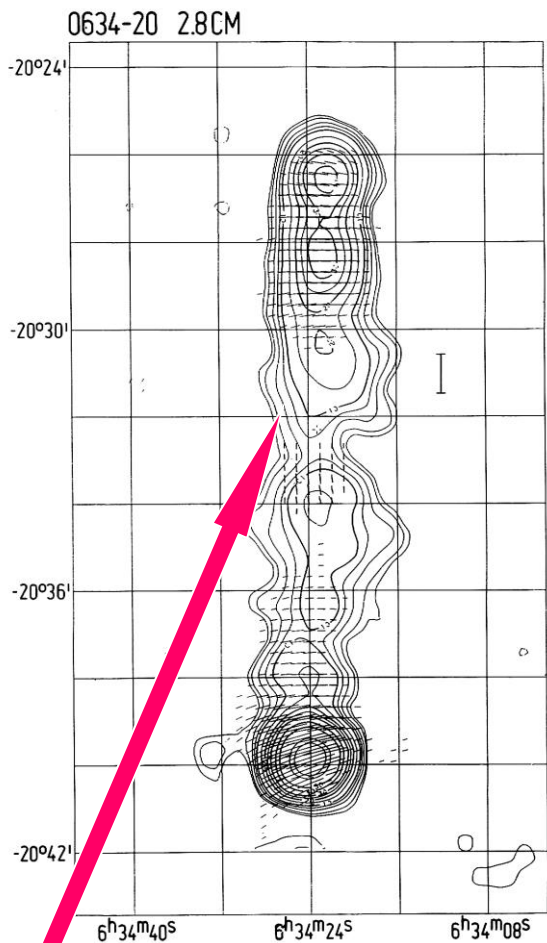
8 FR II-like GRG's, w. detailed,
multi- λ obs. & analysis
Kronberg, Colgate, Li, Dufon ApJL 2004

- Willis & Strom, 1978,80
- Kronberg, Wielebinski & Graham.1986,
- Mack *et al.* A&A 329, 431, 1998
- Schoenmakers *et al.* 1998,2000
- Subrahmanian *et al.* 1996
- Feretti *et al* 1999
- Lara *et al.* 2000
- Palma *et al.* 2000



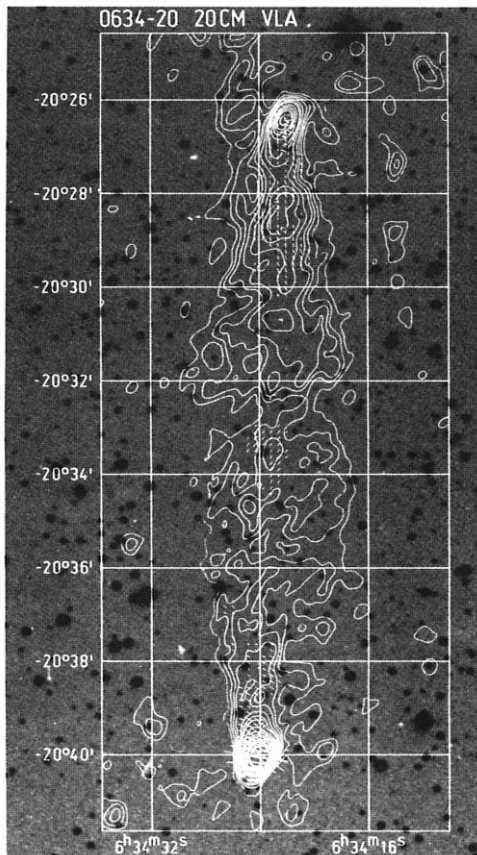
AUI/NRAO/VLA image

10 GHz



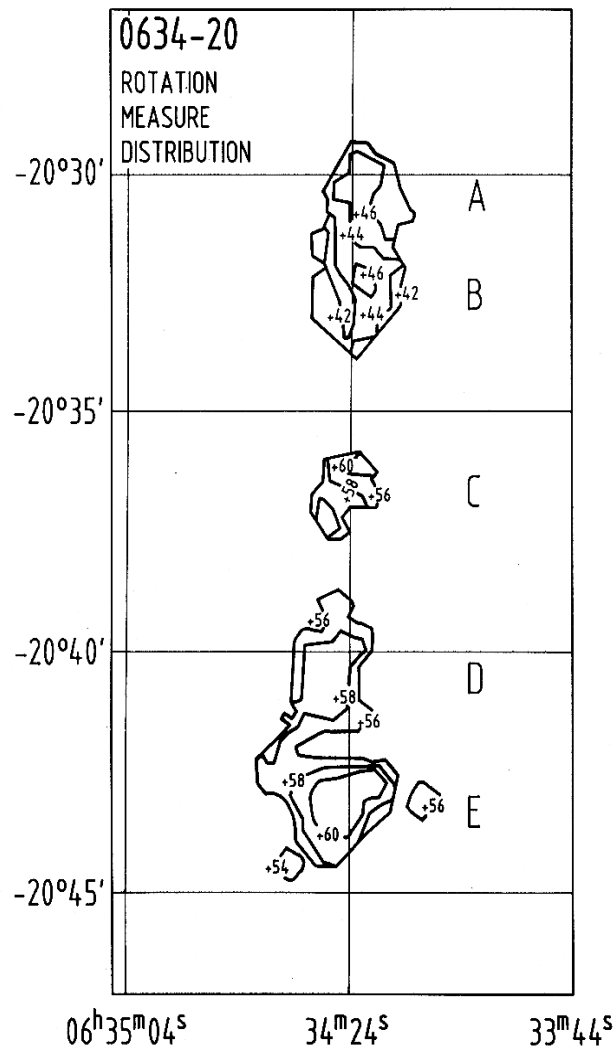
Effelsberg 100m.
Telescope 10.6 GHz

1.4 GHz



VLA 1.4GHz

Faraday RM(radians/m²)

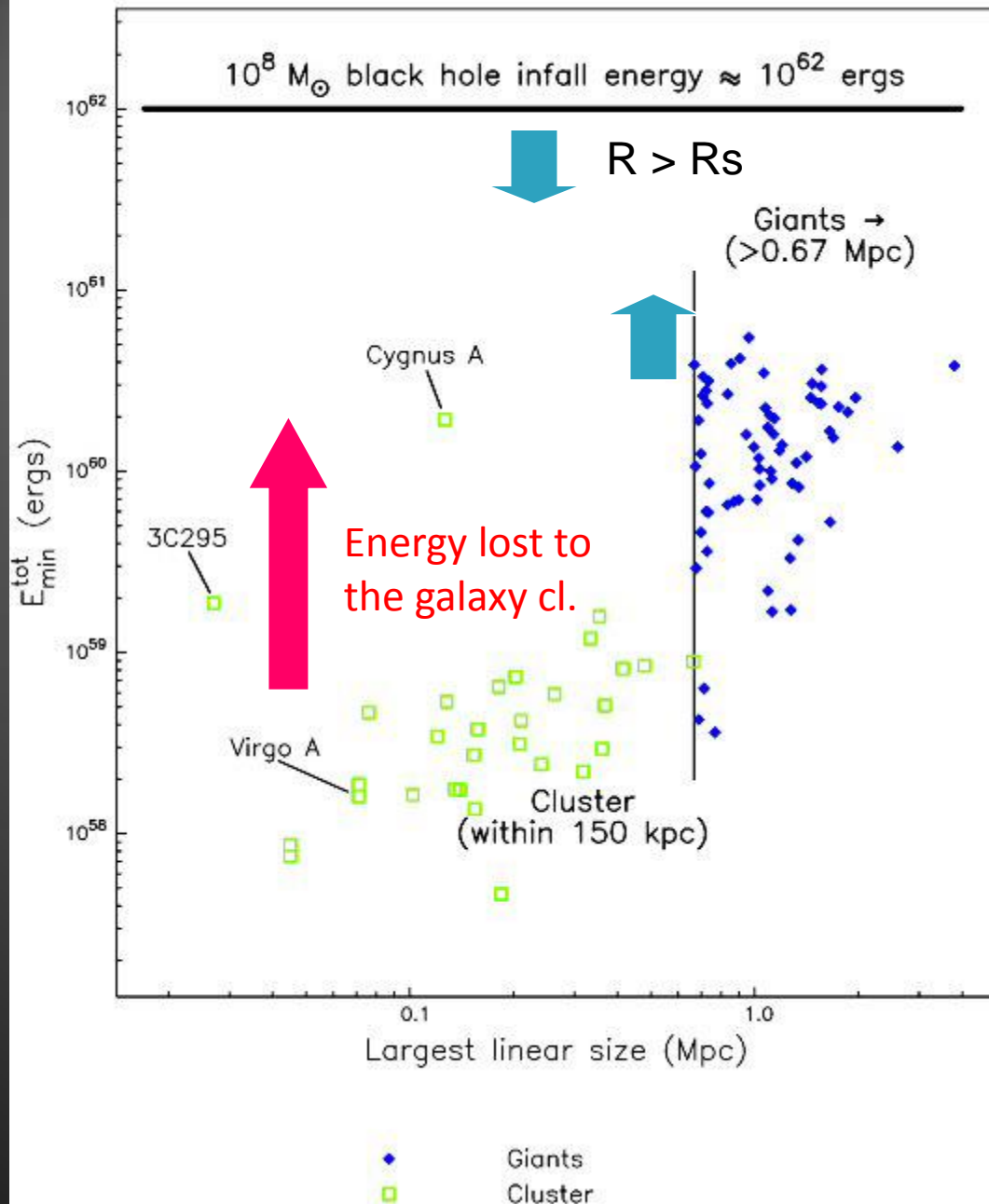


Kronberg, Wielebinski & Graham

A&A 169, 63, 1986

Freshly accelerated, and
"starved" of thermal plasma

$$E_{CR} \approx 10^{19} \left(\frac{B}{3\mu G} \right) \left(\frac{L}{1 \text{ Mpc}} \right) \text{ eV}$$



ENERGETICS:



$$=M_{\text{BH}}c^2$$



Mind the gap!!

Accumulated energy
 $(B^2/8\pi + \epsilon_{\text{CR}}) \times (\text{volume})$
 from "mature" BH-powered
 radio source lobes

Giant Radio Galaxies (GRG)
 capture the highest fraction
 of the BH energy
 released to the IGM

GRG's are the best BH energy
 calorimeters available

Overview of radio observational aspects

- Kpc-scale jets are likely electromagnetically dominated : ($P = I^2 Z$)
- Important measurables are power flow and current (I)
- Jets and lobes that proceed from a Supermassive Black Hole(SMBH) are good candidates for Hadron acceleration to the highest energies
- Why measurements of jet current are currently difficult at resolution scales of *e.g.* the NRAO JVLA .

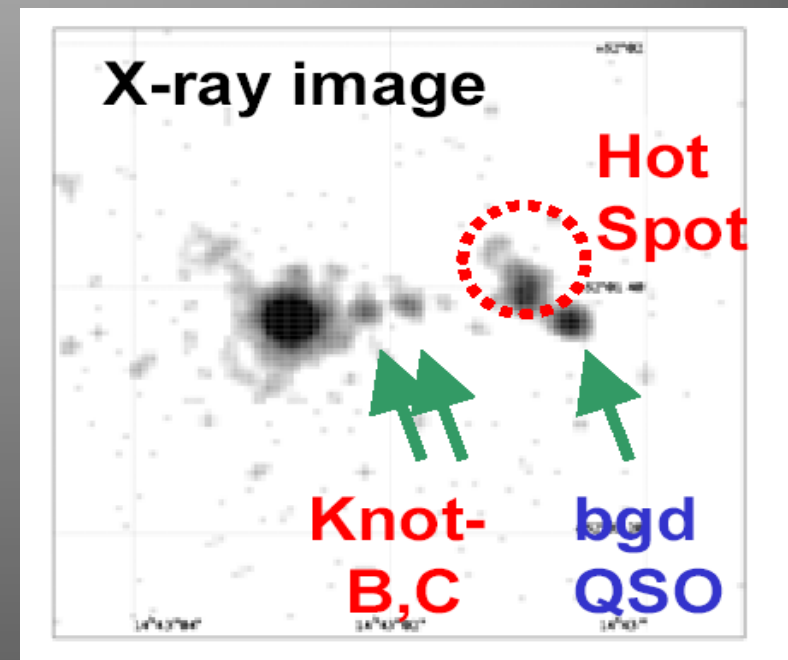
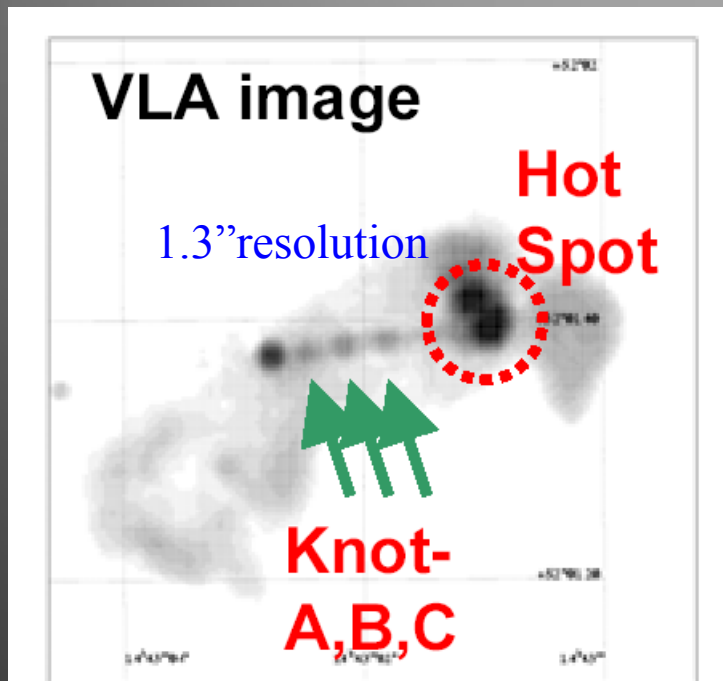
Knots and Hotspots of **3C303** ($z=0.141$, $D \sim 600$ Mpc)

Radio (VLA) and **X-Ray** (CHANDRA)

P. Kronberg, Can.J. Phys **64**, 449, 1986

P. Leahy & R. Perley, Astr. J. **102**, 537, 1991

*J. Kataoka, P. Edwards,
M. Georganopoulos, F. Takahara,
& S. Wagner A&A* **399**, 91, 2003

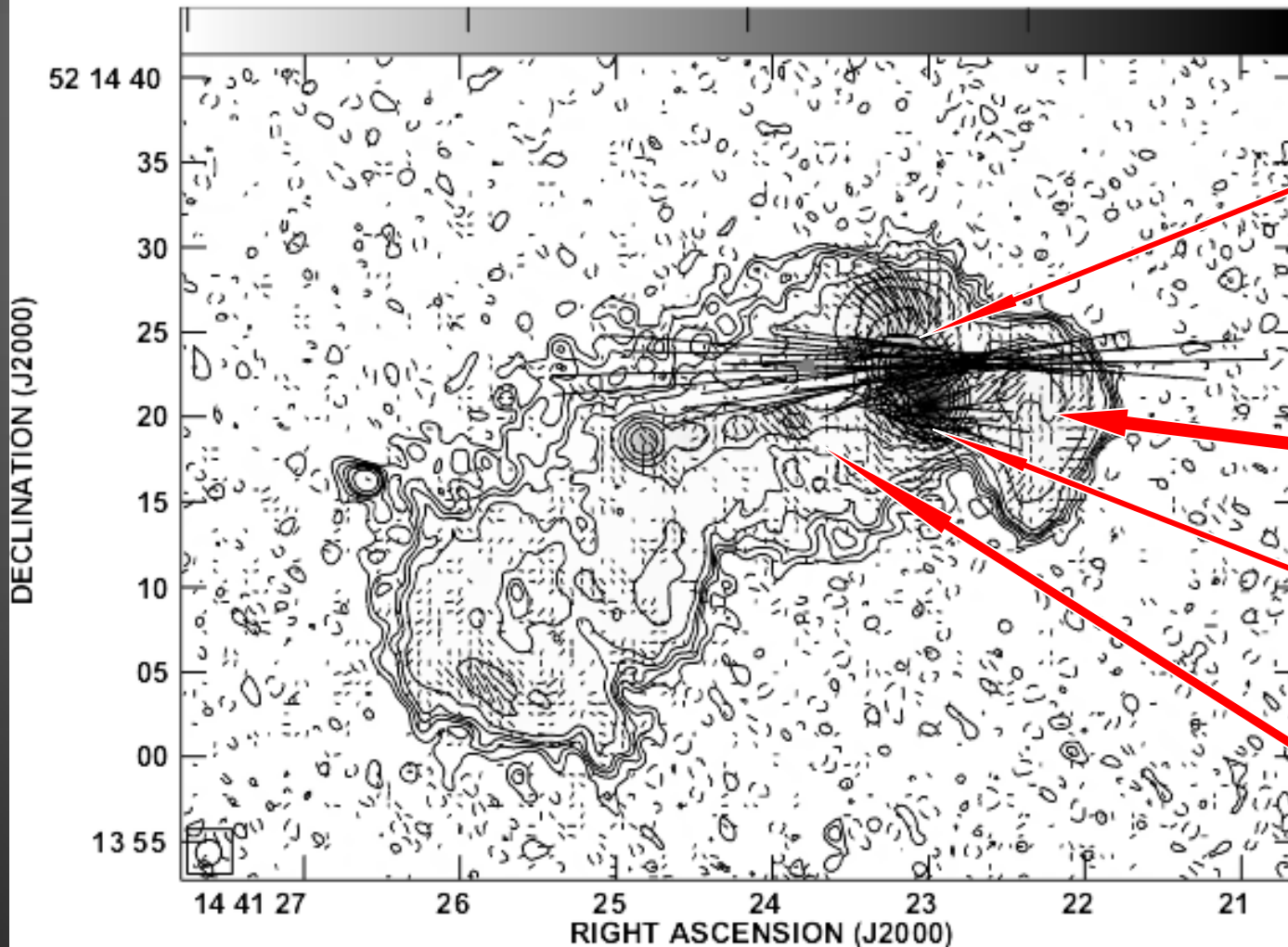


3C303 1.4GHz, 1.4'' resolution

PLot file version 5 created 02-MAY-2011 18:51:05

ALL: 3C303 IPOL 1406.750 MHZ 3C303L.ICL3.1

0 100 200 300



Grey scale flux range= -1.3 395.7 MilliJY/BEAM

Cont peak flux = 3.9566E-01 JY/BEAM

Levs = 5.000E-04 * (-1, 1, 2, 3, 4, 6, 12, 24, 48, 100, 200, 300)

Pol line 1 arcsec = 2.5000E-03 JY/BEAM

3 spheroid "islands"
Each has high
Mag field ordering
& current signatures

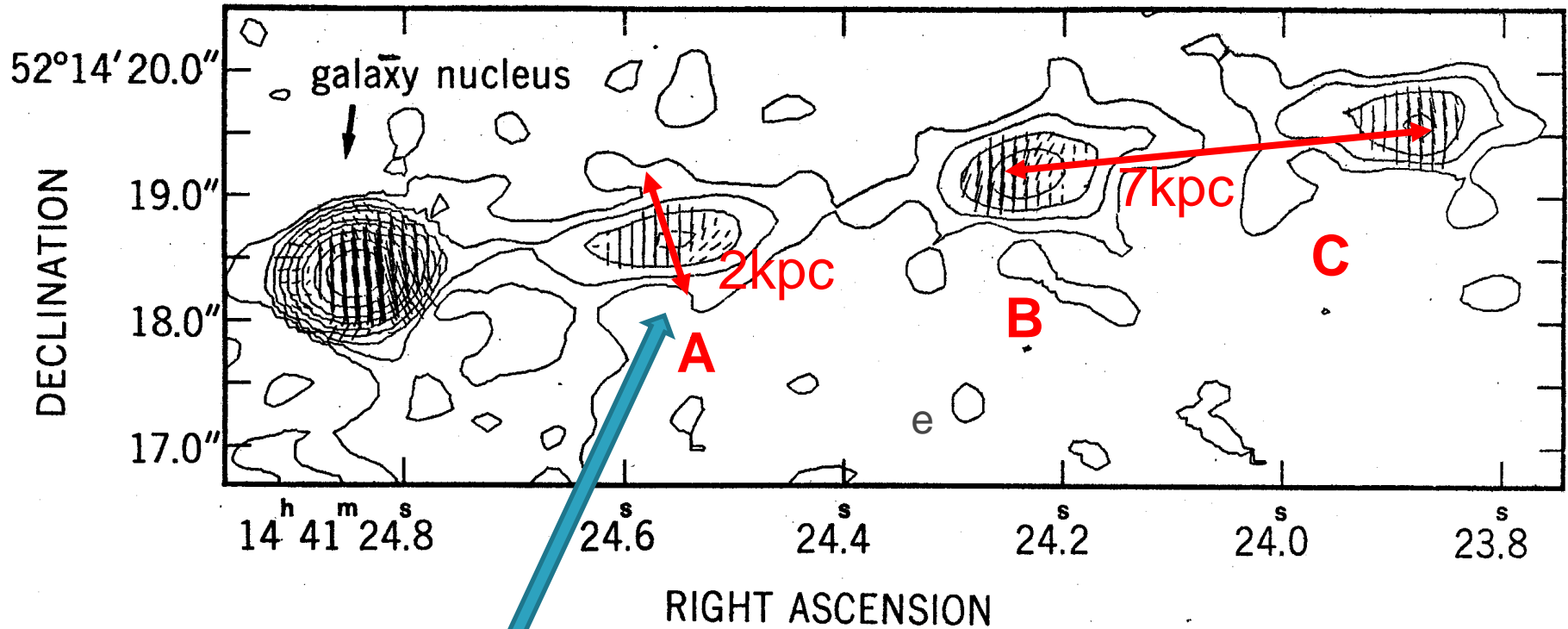
jet continues
to here

jet disruption
point

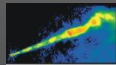
Knot "E3" has a
measured ∇RM
vector

4.9GHz VLA image at 0.3'' resolution

3C303 4866 MHz 0.35'' angular resolution



Compare scales!



The M87 jet on the physical scale of one 3C303 "knot"

M87 Knot cocoons are ~ 12,000 times (vol.) **smaller** than those in 3C303!
SMBH-powered jets are very scale-independent systems!

Plasma Diagnostics of the 3C303 jet

Lapenta & Kronberg ApJ 625, 37-50, 2005

(1) <(Total energy flow rate)> = $E_{\min}^T / \tau \approx \underline{2.8 \times 10^{43}} \tau_7^{-1} \text{ erg/s}$

(2) Jet's total photon luminosity radio \rightarrow X-ray = $\underline{1.7 \times 10^{42}} \text{ erg s}^{-1}$

\rightarrow Radiative dissipation from the jet
is $\approx 10\%$ of the energy flow along jet!

(3) Measure knots' synchrotron luminosity & size (D_{knot}) $\rightarrow \underline{B_{\text{int}}^{\text{knot}} = 10^{-3} \text{ G}}$

(4) Measure Faraday rotation isolated in the knots, $\underline{RM} \propto n_{\text{th}} \times \underline{B_{\text{int}}^{\text{knot}}} \times \underline{D_{\text{knot}}}$

gives n_{th} in knots for 3C303) $\rightarrow \underline{n_{\text{th}} \approx 1.4 \times 10^{-5} \text{ cm}^{-3}}$ (a low, extragalactic - level density!)

(3) & (4) \rightarrow estimate of V_A within knots : $V_A^{\text{knot}} \propto B_{\text{int}}^{\text{knot}} n^{-0.5}$

RESULT: $\underline{V_A^{\text{knot}} \approx c}$. i.e. in the relativistic range -- V_A^{rel}

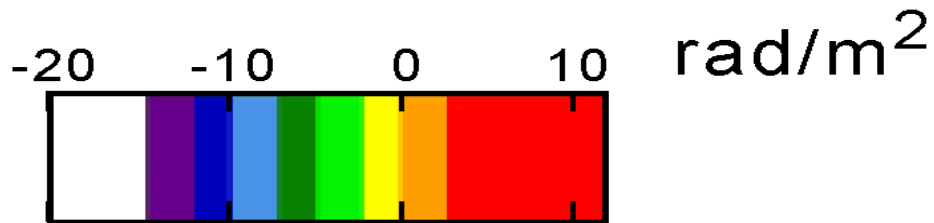
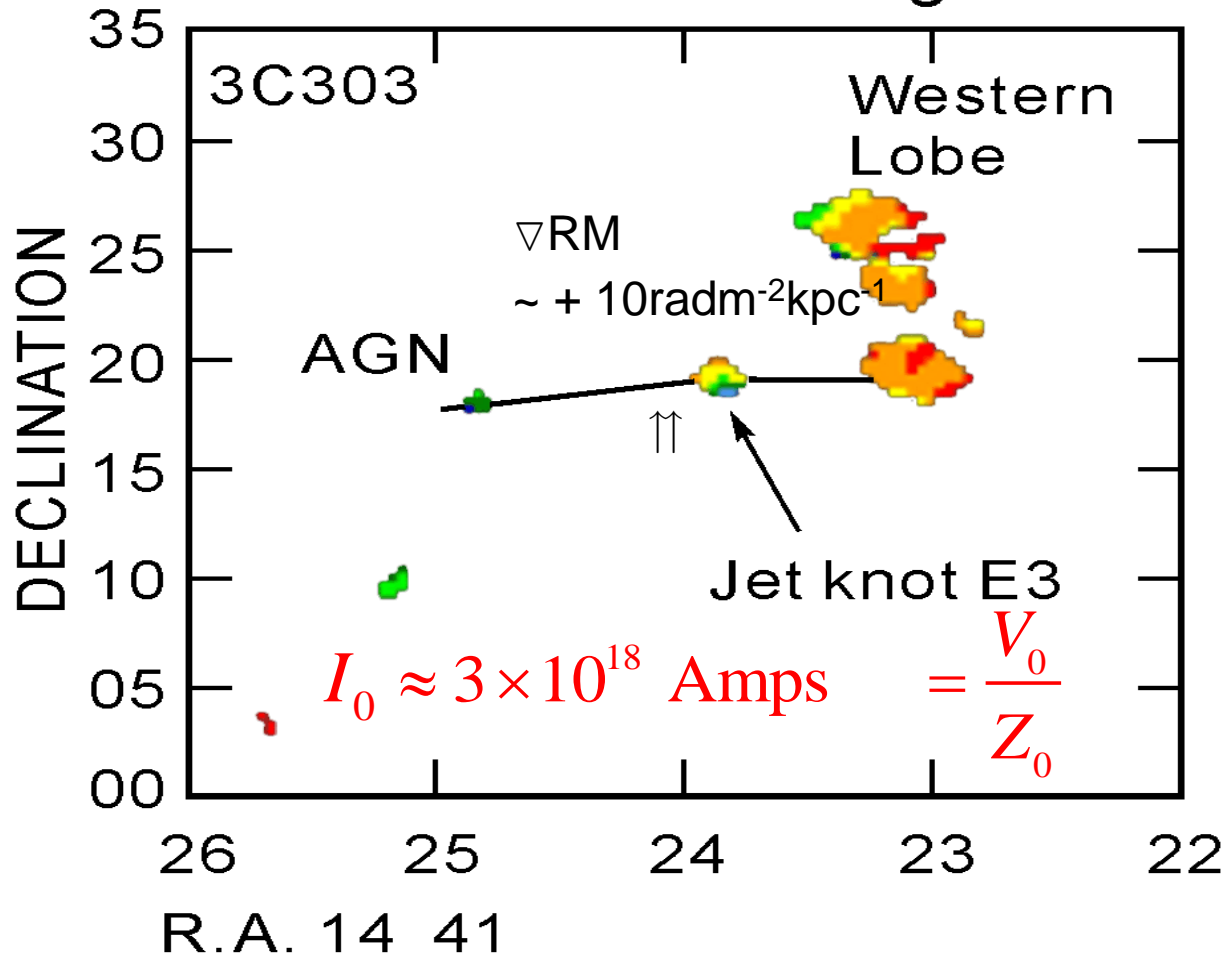
How to estimate the jet current? --
what measurements are required :

1. Need sub-arcsec resolution, + adequate sensitivity in Stokes IQU at ≥ 2 frequencies,
2. Faraday RM distribution over the jet $\frac{\Delta\chi}{\Delta\lambda^2}(x, y)$ at a common angular resolution
2. High resolution X-ray image (\sim keV) gives estimated T and ρ of gas surrounding the jet
4. Need **nearby sky** RM's to establish the zero-level of RM
i.e. subtract $\langle \text{RM}_{\text{backgnd sources}} \rangle$ from the RM's in the jet image
(normally only feasible outside a galaxy cluster)

P.P. Kronberg Can J. Phys **64**, 449, 1986 (original results)

P.P. Kronberg, R.V.E. Lovelace, G. Lapenta, & S.A. Colgate, ApJL **741**, L15, 2011

DEC. 52 14 RM Image



*Kronberg, Lovelace, Lapenta
& Colgate ApJLett 741, L15, 2011*

Foreground RM
Correction uncertainty

Transmission line analogy:

Line Voltage:

$$V_0 = -\frac{1}{2}r_0 \int_0^{\bar{r}_2} d\bar{r} E_r(\bar{r}) = \frac{r_0}{3^{1/4}} \frac{B_0}{\sqrt{\mathcal{R}}}$$

$$= 3.4 \times 10^{20} \text{ beta Volts}$$

Axial current:

$$\mathcal{R} \equiv r_0/r_g \geq 1.$$

$$I_0 = -\frac{1}{2}cr_2 B_\phi(r_2) = \frac{V_0}{Z_0}$$

$= 3.8 \times 10^{18} \text{ A}$

This leads to a straightforward electrical circuit analogue for BH energy transfer into ``empty'' space

P.P. Kronberg, R.V.E. Lovelace, G. Lapenta & S.A. Colgate ApJL 741, L15 2011

R.V.E. Lovelace & P.P. Kronberg, MNRAS 430, 2828-2835, 2013

- *Low thermal density around knots suggests dominance of a Poynting flux*
- $P \sim 10^{37}$ watts of directed e.m. power, and $I = 3.3 \times 10^{18}$ ampères of axial current, directed (in this case) away from the BH (Sign of ∇RM gives the I direction).

Poynting jet's electrical properties: (current, impedance, voltage):

$$I_0 = cr_2 B_{\phi(r_2)} = \frac{V_0}{Z_0} \approx 3 \times 10^{18} \text{ Amps (MKS)}$$

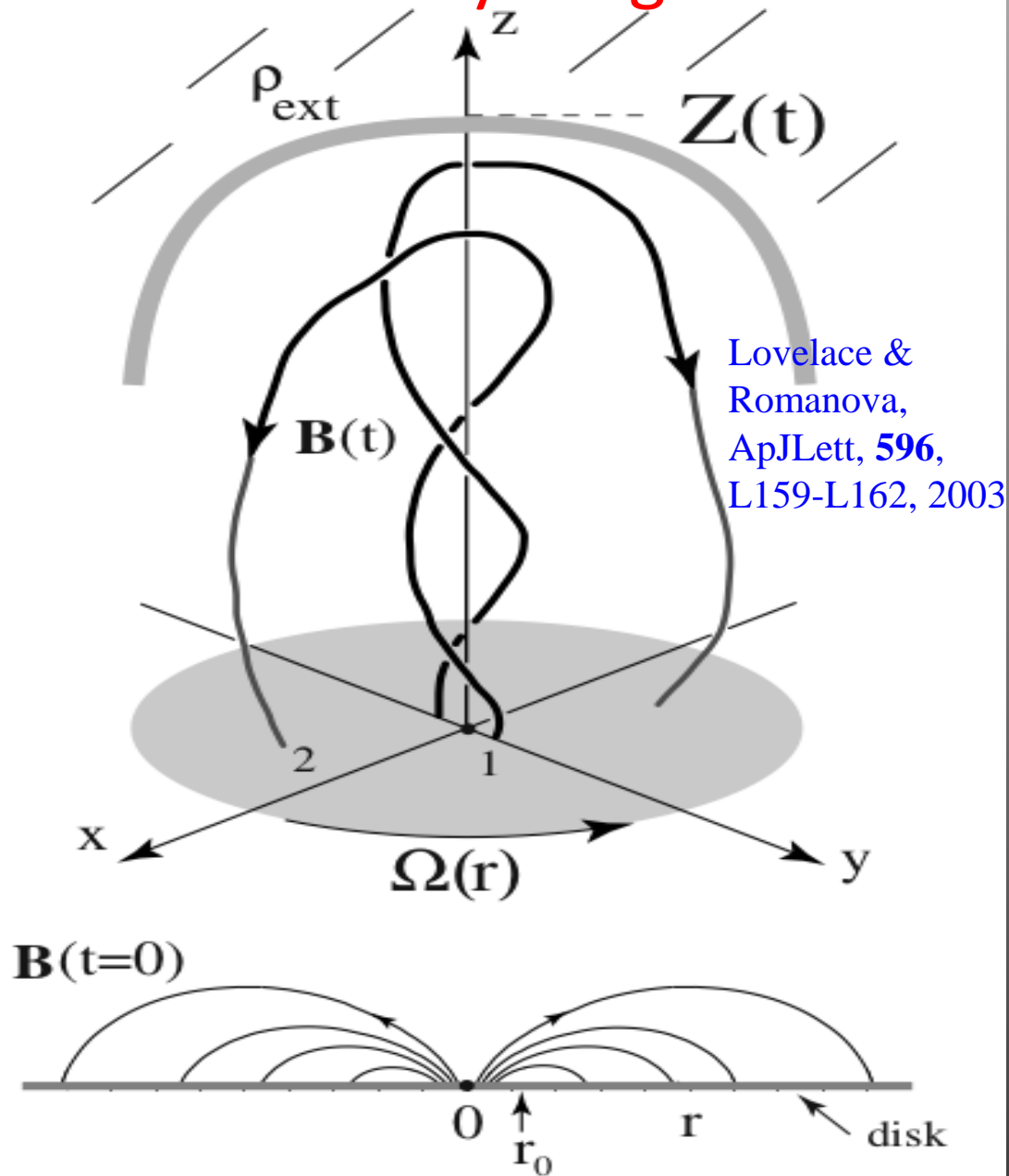
$$Z_0 = \frac{3}{c} \beta \text{ (cgs)} = 90\beta \text{ Ohms (MKS)}$$

$$V_0 = \frac{r_0 B_0}{3^{1/4} \sqrt{R}} = 2.7 \times 10^{20} \text{ Volts (MKS)}$$

$\beta = \frac{U}{c} \lesssim 1$, where r_1, r_2 are the inner & outer transmission line radii (Lovelace & Ruchi, 1983)

Concept of Poynting flux-
dominated energy flow
from a BH accretion disc

Relativistic Poynting Flux Jets



Pre 1900 - problem of telegraphic signal (energy) propagation over very long distances -- telegrapher's equations:

Time and space-dependent perturbations of a Poynting-flux jet are described by the Telegrapher's equations,

$$\frac{\partial \Delta V}{\partial t} = -\frac{1}{C} \frac{\partial \Delta I}{\partial z}, \quad \frac{\partial \Delta I}{\partial t} = -\frac{1}{L} \frac{\partial \Delta V}{\partial z}, \quad (20)$$

where $(\Delta V, \Delta I)$ represent deviations from the equilibrium values (V_0, I_0) . The equations can be combined to give the wave equations

O. Heaviside, Electromagnetic Theory 1893

Van Nostrand NY,

$$\left(\frac{\partial^2}{\partial t^2} - u_\varphi^2 \frac{\partial^2}{\partial z^2} \right) (\Delta V, \Delta I) = 0, \quad (21)$$

where

$$u_\varphi = \frac{1}{\sqrt{LC}} = c \left[1 + \frac{1}{2\Gamma^2} + \frac{1}{2\Gamma^2} \ln \left(\frac{r_3}{r_2} \right) \right]^{-1/2}, \quad (22)$$

Magnetic Insulation: It breaks down (less common):

when $|\mathbf{E}| \geq |\mathbf{B}|$

Example: where B approaches zero on some surface. At this point, normal Resistive MHD does not apply.

Specific examples :

1. In a magnetic reconnection layer,
or (relevant here)
2. If an electromagnetic jet encounters a “load” with an inductive component, the reflected ΔI and ΔV are no longer exactly in phase. This can create $|\mathbf{E}| \geq |\mathbf{B}|$, for some distance and time period Δt , ---
e.g. $\sim 1000\text{yr}$ in the 3C303 jet, --
--- creating conditions for coherent particle acceleration

Why are jet currents so difficult to measure?

Jet-associated ΔRM (& ∇RM) are intrinsically small.

They also need to be isolated from other source-intrinsic, and ambient (e.g. cluster ICM) emission

*example: ΔRM only $\sim 10 \text{ rad m}^{-2}$ across the 3C303 jet. Even for $I \sim 10^{18} \text{ A}$ at $z=0.14$, & 10 rad m^{-2} , the *angle of rotation* is only 25° at 1 GHz.*

This detectability problem gets worse at higher z

e.g. for 3C9 at $z=2.012$, $\Delta RM = 10 \text{ rad m}^{-2}$ becomes $\frac{10}{(1+2.012)^2} = 1.1 \text{ rad m}^{-2}$ - undetectable!

LOFAR frequencies down to $\sim 200 \text{ MHz}$ can detect RM's at low n_{IG} and $B_{IG} \leq 10^{-6} \text{ G}$

Future magnetoplasma probes of jets will require:

1. Much higher resolution ($\ll 1''$) transverse to the jet “spine”.
2. higher brightness sensitivity.
3. frequency ranges $\lesssim 1$ GHz for Faraday RM imaging.

Low freq. telescopes ($\lesssim 1$ GHz) will achieve these goals

recent references:

R.V.E. Lovelace & P.P. Kronberg, *MNRAS* **430**, 2828, 2013

P.P. Kronberg & R.V.E. Lovelace, *EPJ Web of Conf.* **99**,13005, 2015

Summary:

1. We have estimated the circuit parameters for a supra-galaxy--scale jet.
2. Applied it to observations of the jet in radio galaxy 3C 303.
3. We have developed a simple transmission line model for extragalactic Poynting flux jets.
4. Introduced the concept of magnetic insulation, relevant to VHE particle acceleration, and applied it for a jet-ambient plasma $\beta \approx 10^{-5}$,

Philipp Kronberg, Texas Symposium, Geneva 2015

END