

OBSERVING BLACK HOLES VIBRATIONS

VASCO GENNARI

DECEMBER 19-20, 2022

XV BLACK HOLES WORKSHOP - ISCTE

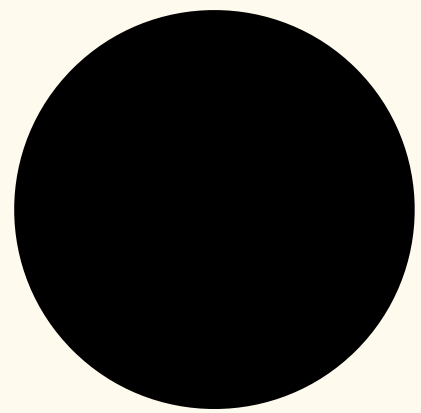


1. ringdown

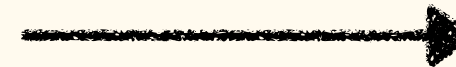
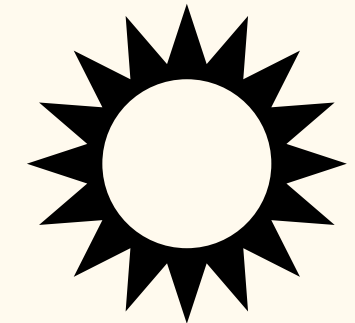
WHY STUDY RINGDOWN?

unique possibility of studying general relativity (GR)
in **strong field** and **extreme curvature** regimes

black hole



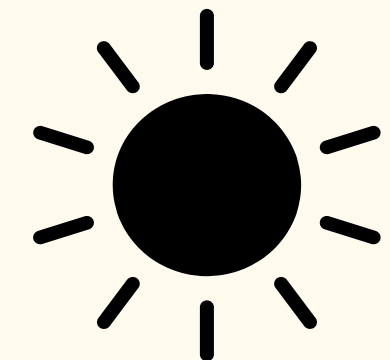
are we really observing
black holes?



is GR the correct theory of
gravity?



are there quantum effects
at the horizon?

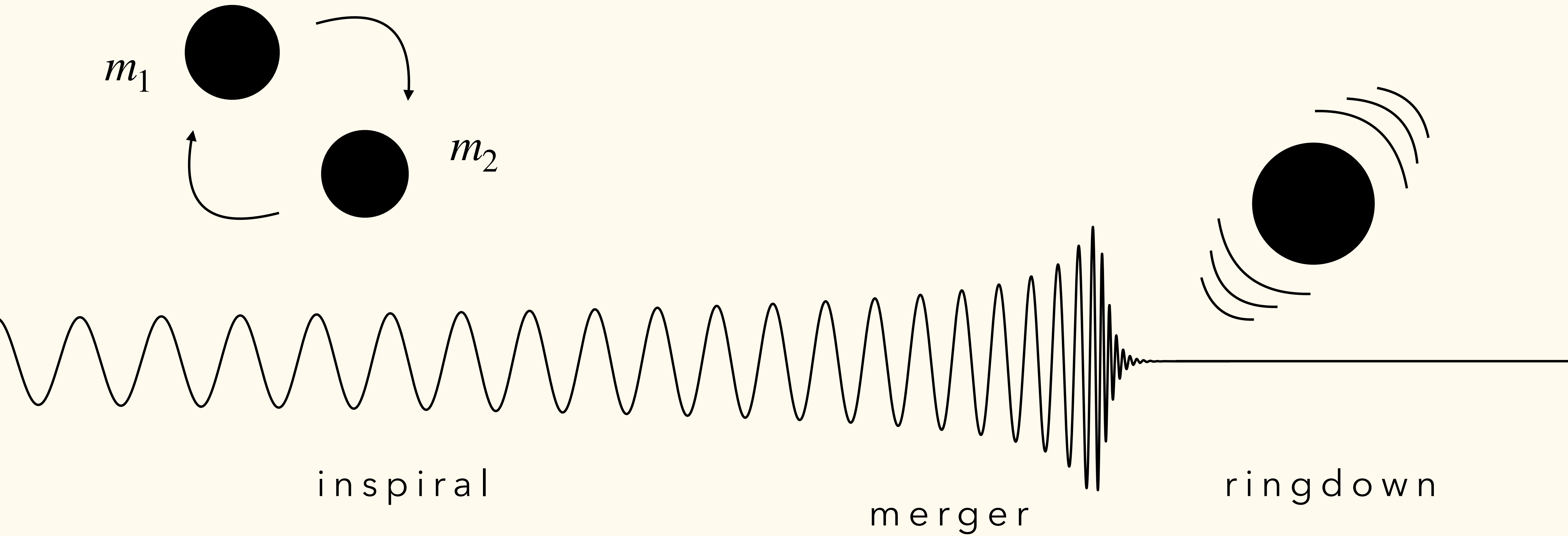


RINGDOWN BASICS

what is the **ringdown**?

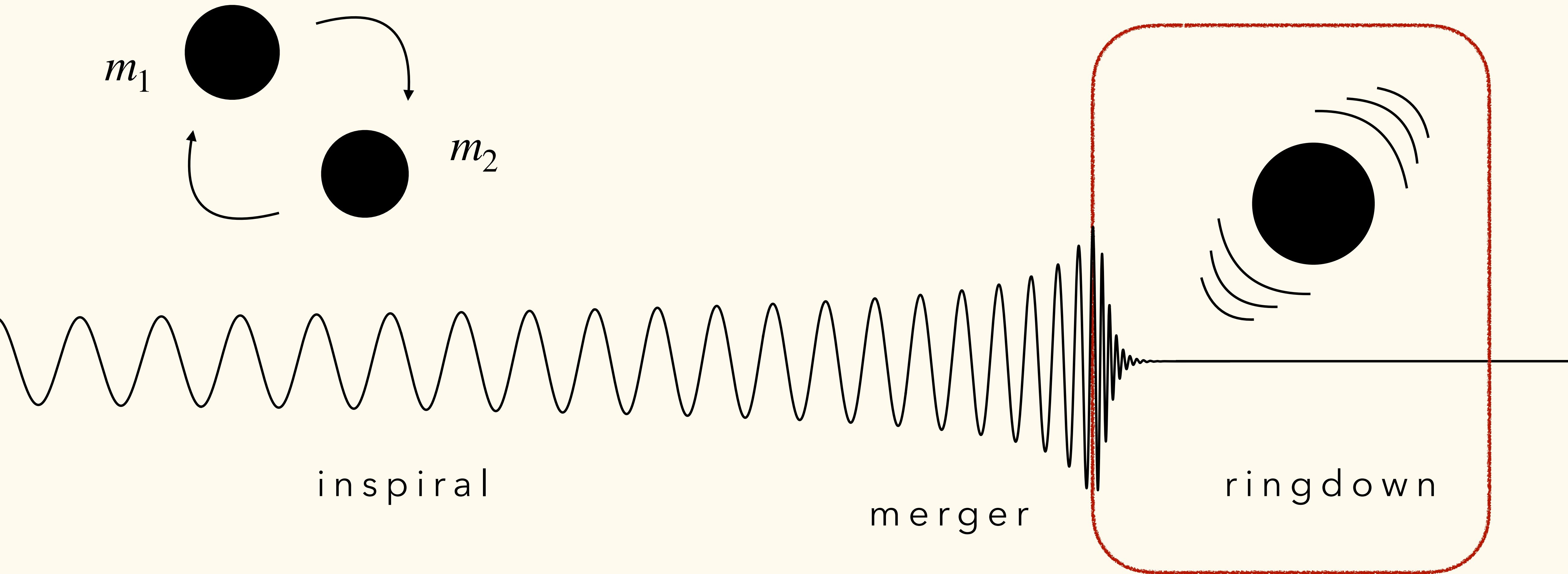
RINGDOWN BASICS

what is the **ringdown**?



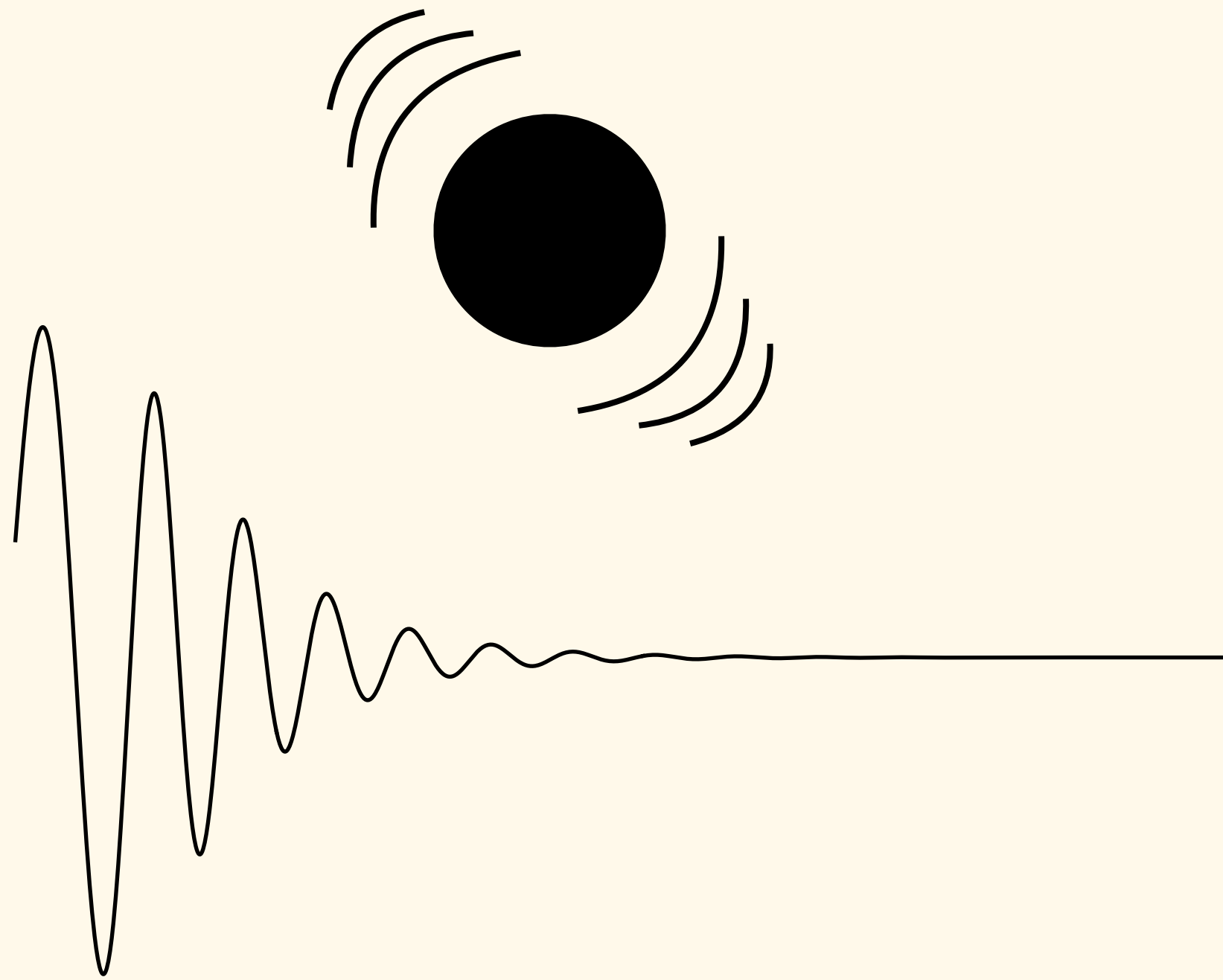
RINGDOWN BASICS

what is the **ringdown**?



RINGDOWN WAVEFORM

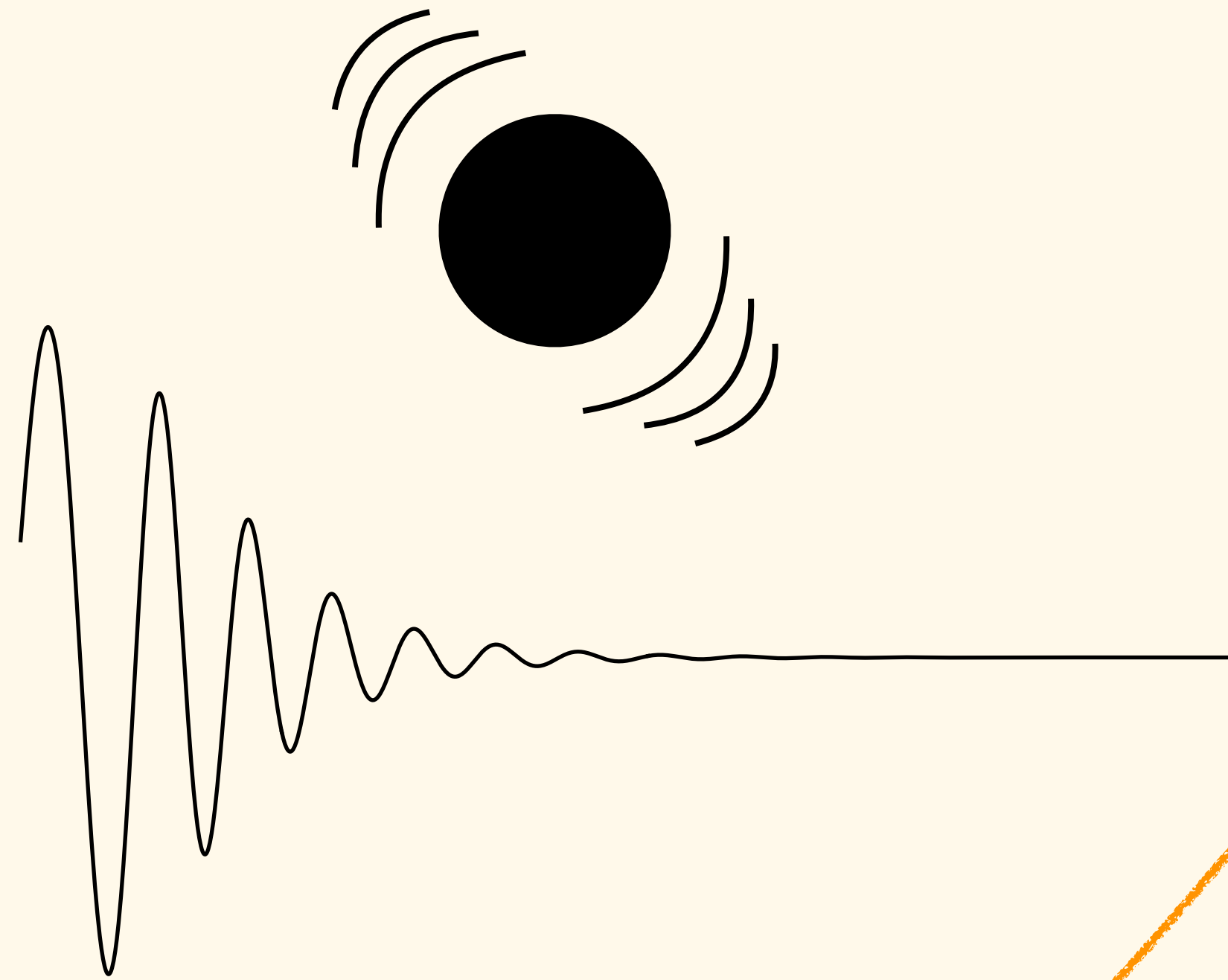
black hole (BH) linear perturbation theory predicts:



$$h = \sum_{lm} A_{lm} e^{-i\omega_{lm}t - t/\tau_{lm}} {}_2Y_{lm}$$

RINGDOWN WAVEFORM

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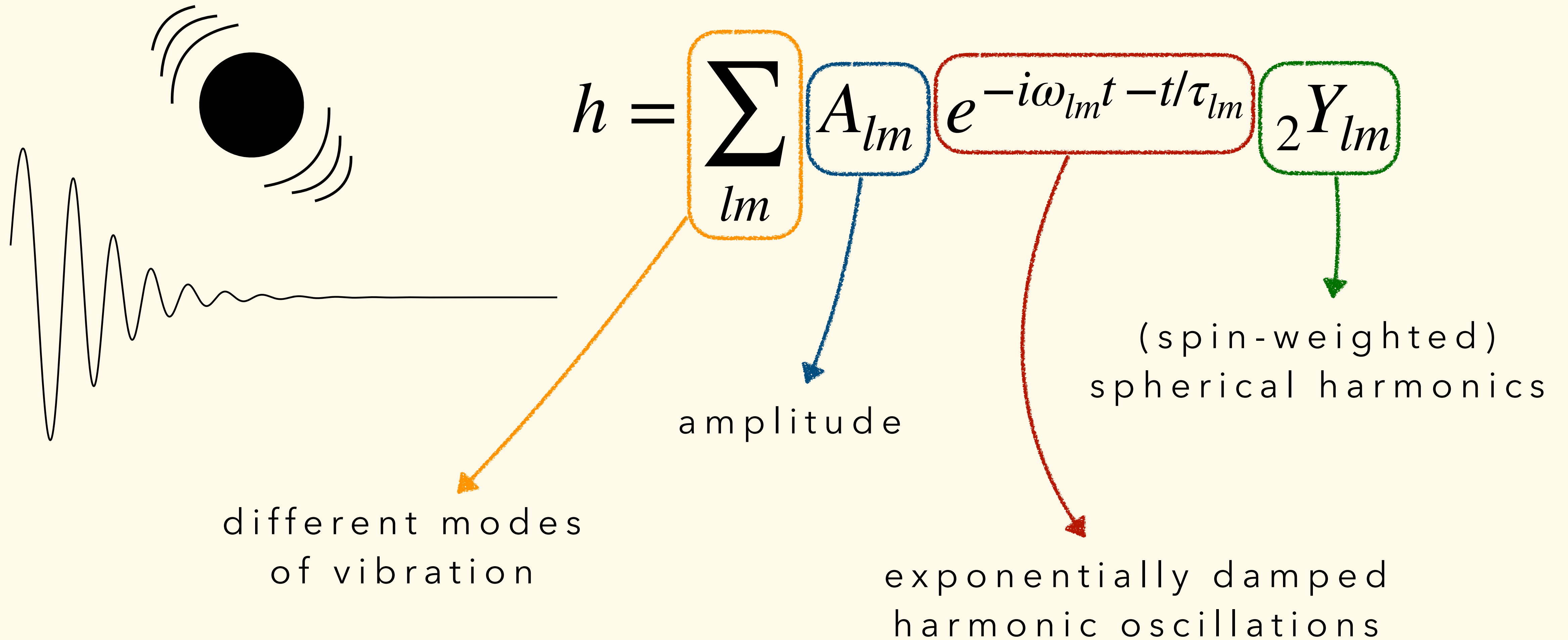


$$h = \sum_{lm} A_{lm} e^{-i\omega_{lm}t - t/\tau_{lm}} {}_2Y_{lm}$$

different modes
of vibration

RINGDOWN WAVEFORM

black hole (BH) linear perturbation theory predicts:



RINGDOWN WAVEFORM

$${}_2Y_{lm}$$

provide angular dependence
for the modes

→ inclination ι

$$h = \sum_{lm} A_{lm} e^{-i\omega_{lm}t - t/\tau_{lm}} {}_2Y_{lm}$$

(spin-weighted)
spherical harmonics

RINGDOWN WAVEFORM

$$h = \sum_{lm} A_{lm} e^{-i\omega_{lm}t - t/\tau_{lm}} {}_2Y_{lm}$$

${}_2Y_{lm}$

provide angular dependence
for the modes

↪ inclination ι

exponentially damped
harmonic oscillations

$$e^{-i\omega_{lm}t - t/\tau_{lm}}$$

ω_{lm} and τ_{lm} are known once M and χ are fixed

↪ quasinormal modes

RINGDOWN WAVEFORM

$$h = \sum_{lm} A_{lm} e^{-i\omega_{lm}t - t/\tau_{lm}} {}_2Y_{lm}$$

$${}_2Y_{lm}$$

provide angular dependence
for the modes

amplitude

inclination ι

$$e^{-i\omega_{lm}t - t/\tau_{lm}}$$

ω_{lm} and τ_{lm} are known once M and χ are fixed

quasinormal modes

$$A_{lm}$$

depend on the specific process that perturbs the BH

are not known analytically

2. higher modes

HIGHER MODES (HMS)

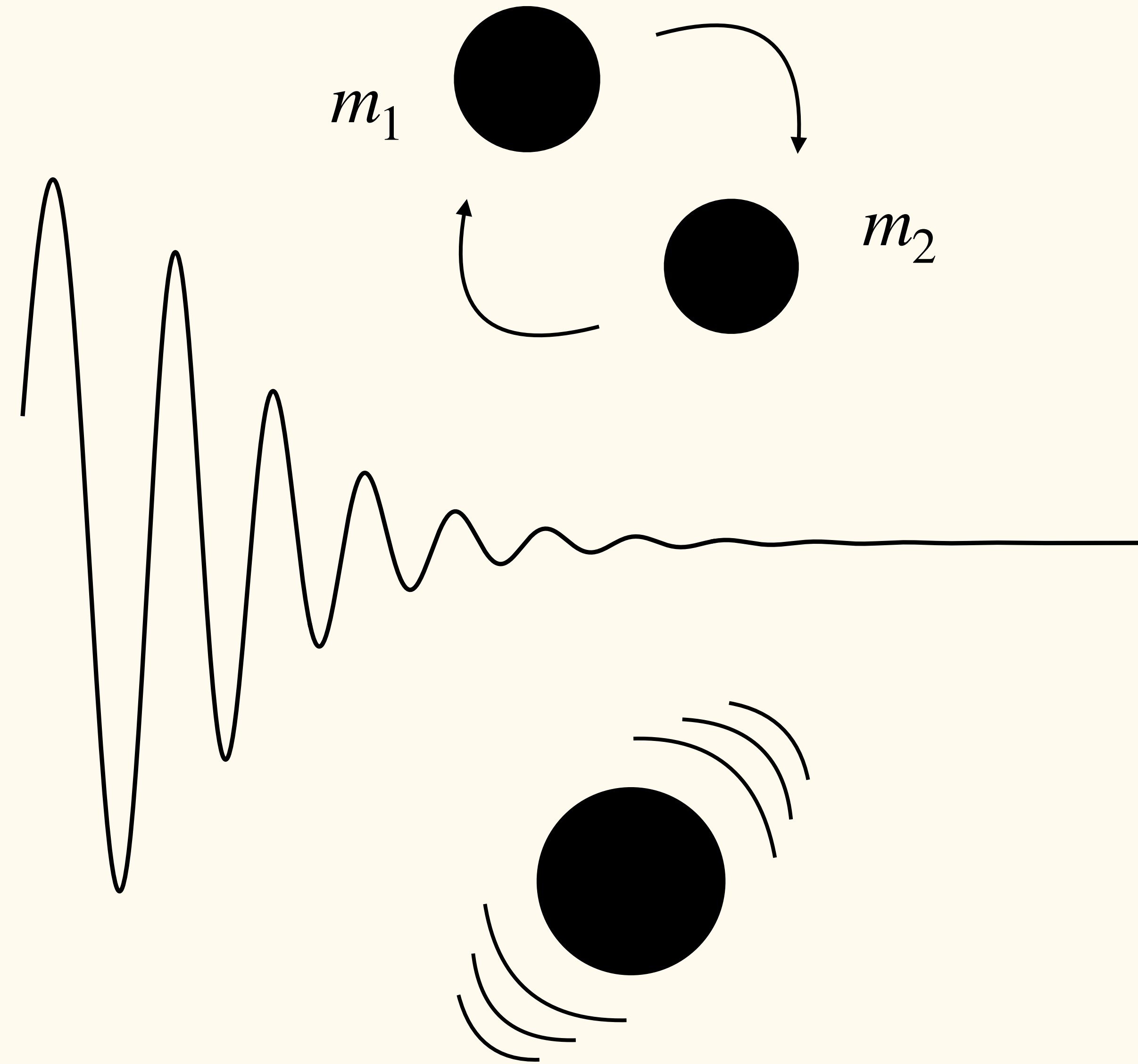
which modes are observable
in the ringdown?

HIGHER MODES (HMS)

which modes are observable
in the ringdown?

for quasi-circular BHs with
masses $m_1 \simeq m_2$:

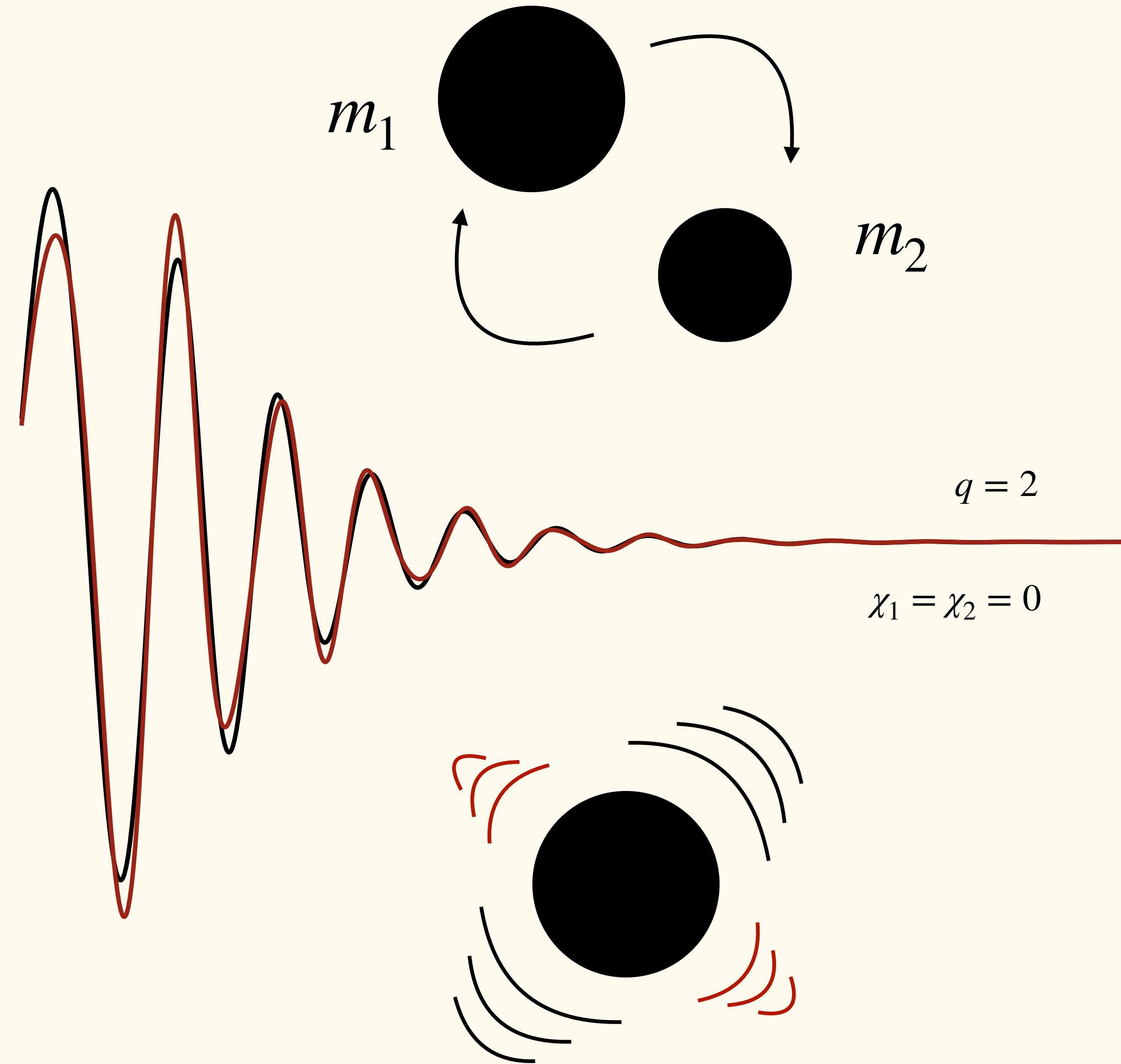
- dominant contribution
(2,2) **fundamental mode**
- subdominant contribution
(3,3), (2,1), (4,4) **higher modes**



HIGHER MODES EXCITATION

HMs can be excited by:

- increasing the mass ratio $q \equiv m_1/m_2$
- increasing the initial spins χ_1 and χ_2

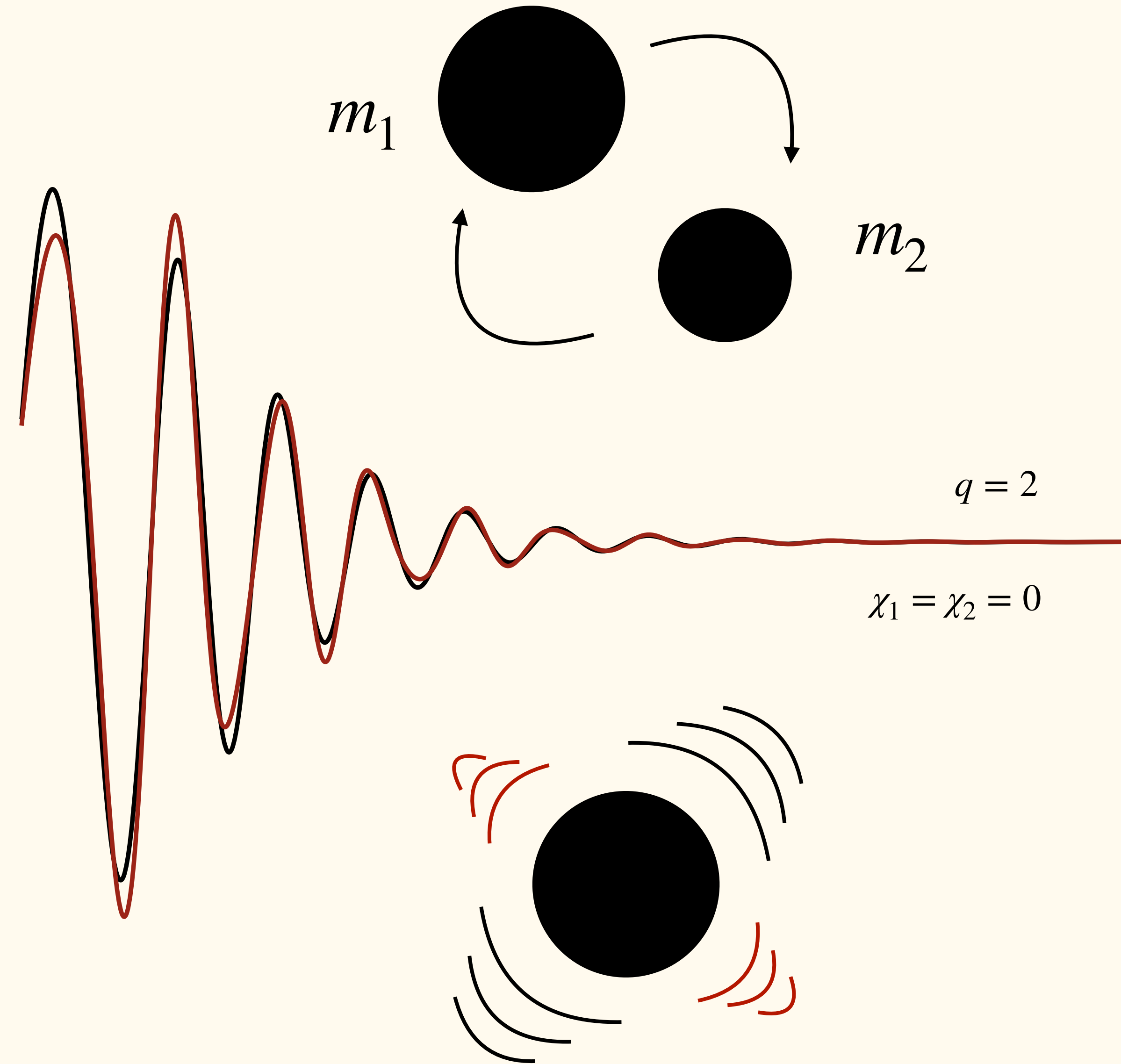


HIGHER MODES EXCITATION

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can we observe HMs with current detectors?



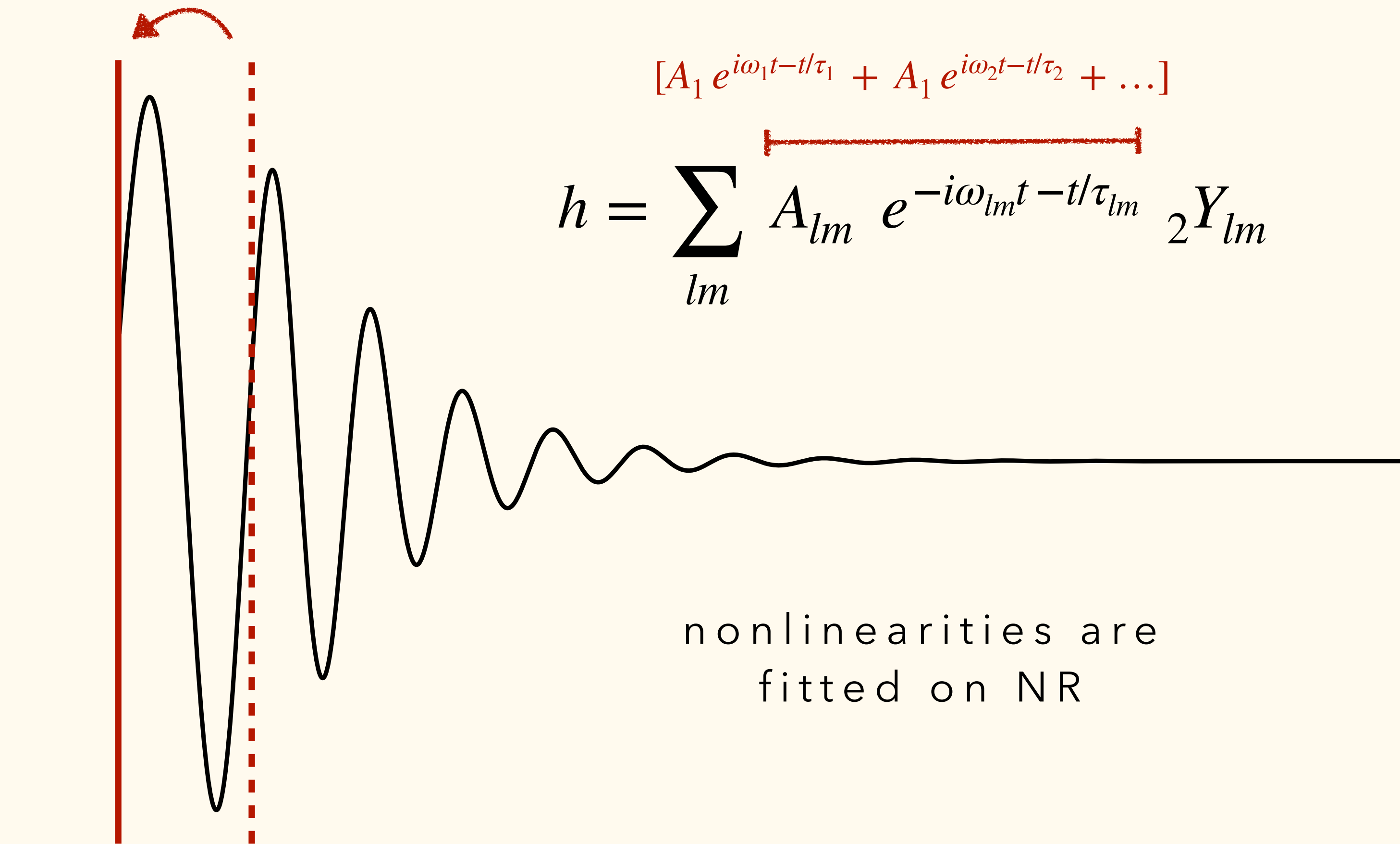
TEOBPM MODEL

Damour and Nagar (2014)
1406.0401

EOB model that includes
post-merger **nonlinearities**

advantages:

- fixes the starting time
- more data with high SNR
- includes the A_{lm}



$$h = \sum_{lm} A_{lm} e^{-i\omega_{lm} t - t/\tau_{lm}} {}_2Y_{lm}$$

$[A_1 e^{i\omega_1 t - t/\tau_1} + A_1 e^{i\omega_2 t - t/\tau_2} + \dots]$

nonlinearities are
fitted on NR

- unambiguous results
- more accurate results

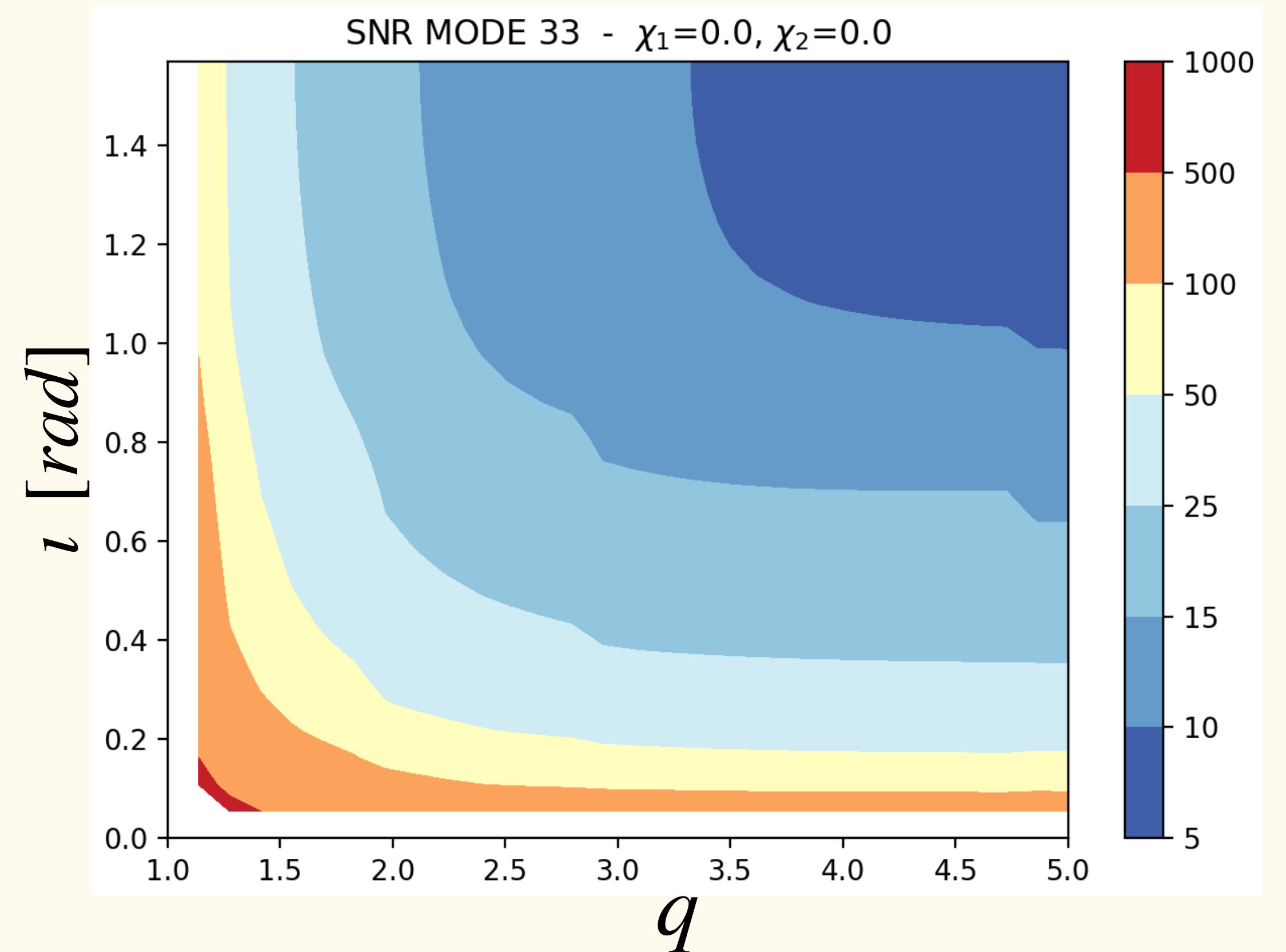
(~ 20% more SNR)

HMS DETECTABILITY

$$\ln B \simeq \frac{1}{2} (1 - FF^2) SNR^2$$

SNR needed to detect the (3,3) mode

(with $\ln B = 5$)



HMS DETECTABILITY

$$\ln B \approx \frac{1}{2} (1 - FF^2) SNR^2$$

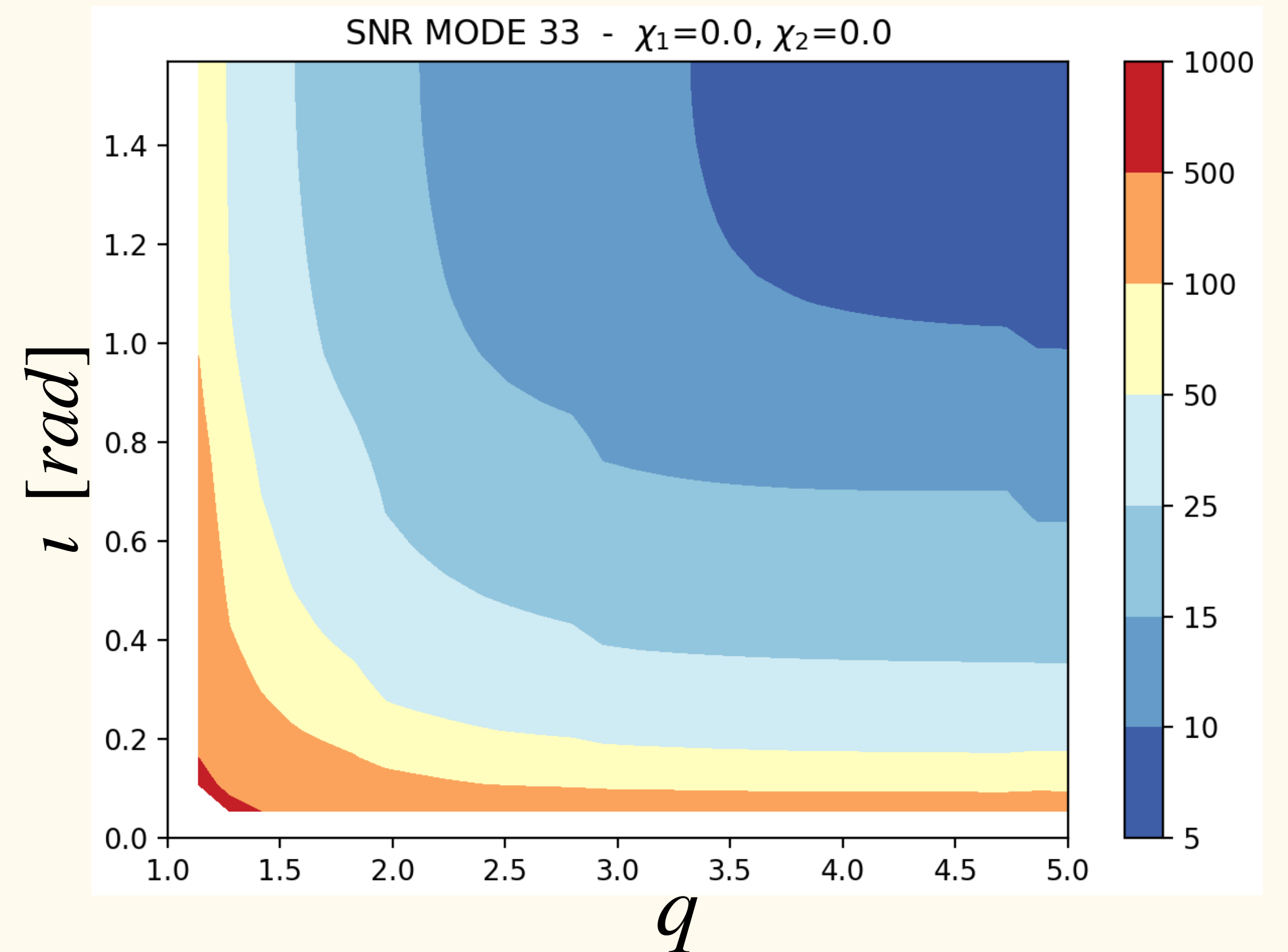
Bayes factor

fitting factor

signal-to-noise ratio

SNR needed to detect the (3,3) mode

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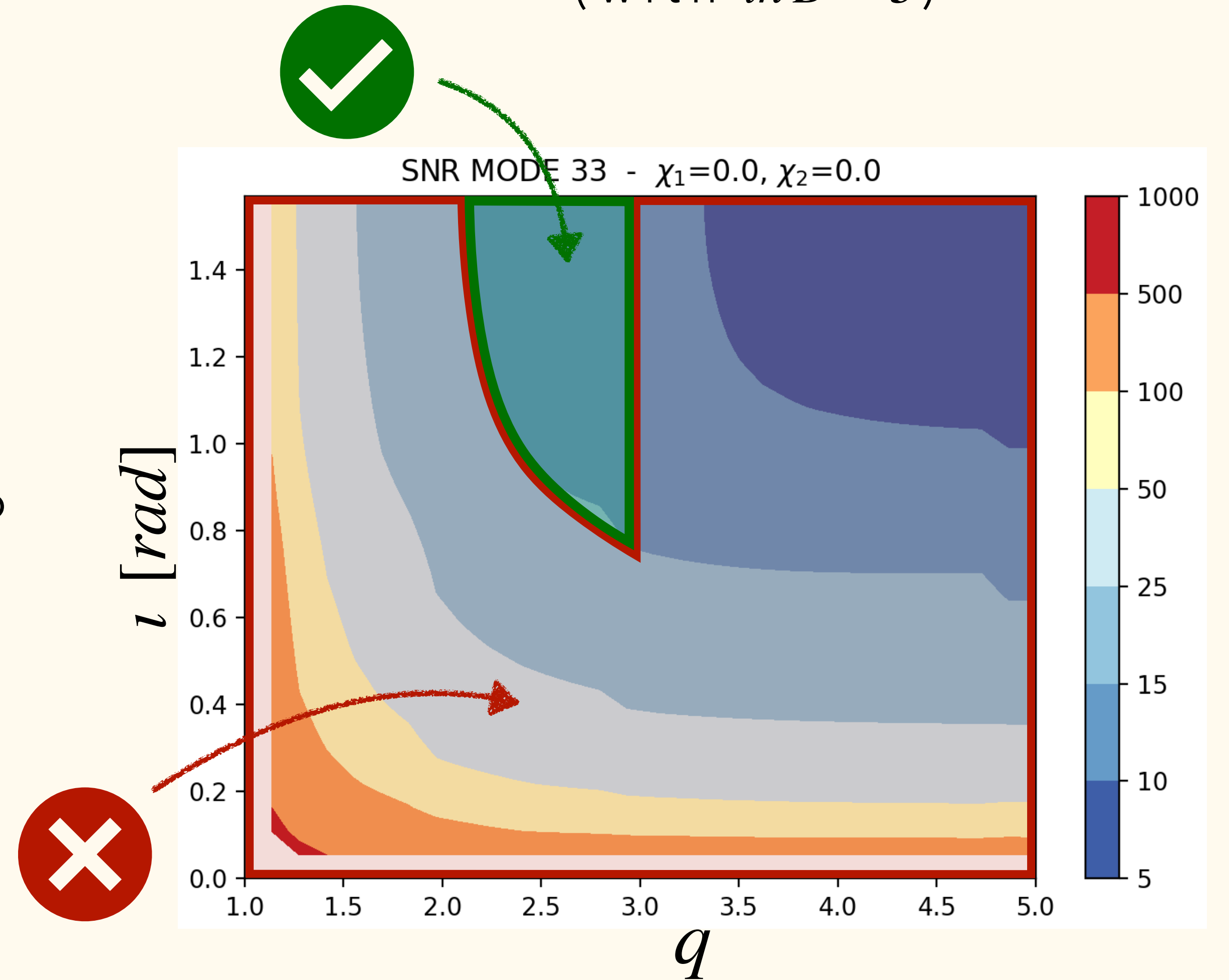
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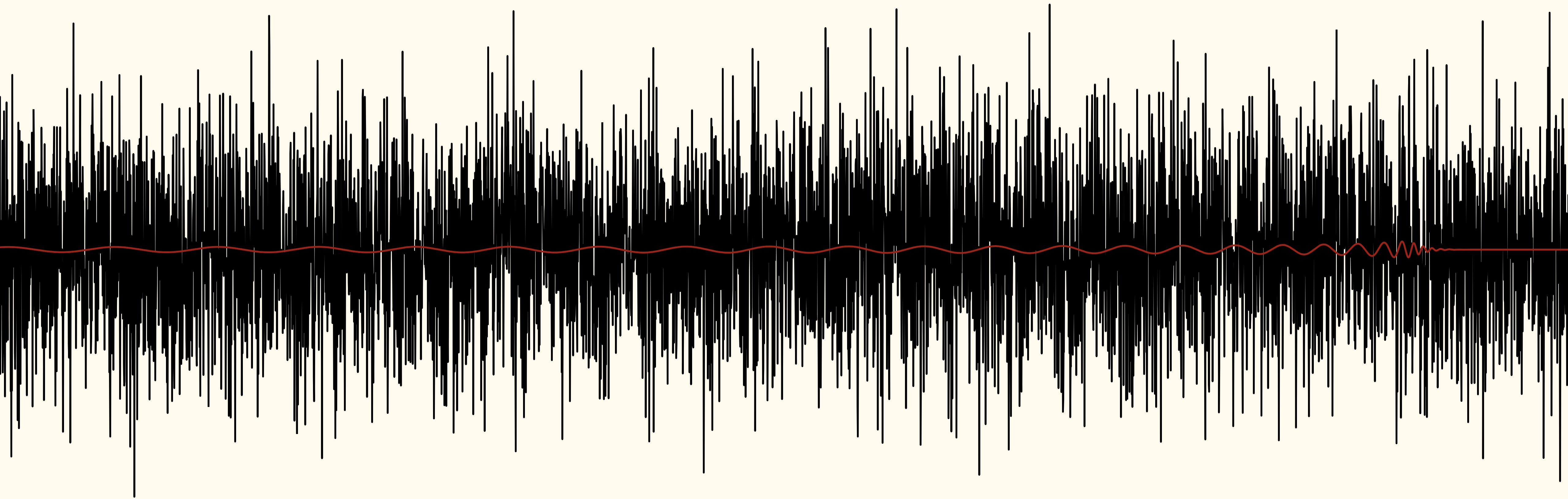
SNR needed to detect the (3,3) mode

(with $\ln B = 5$)

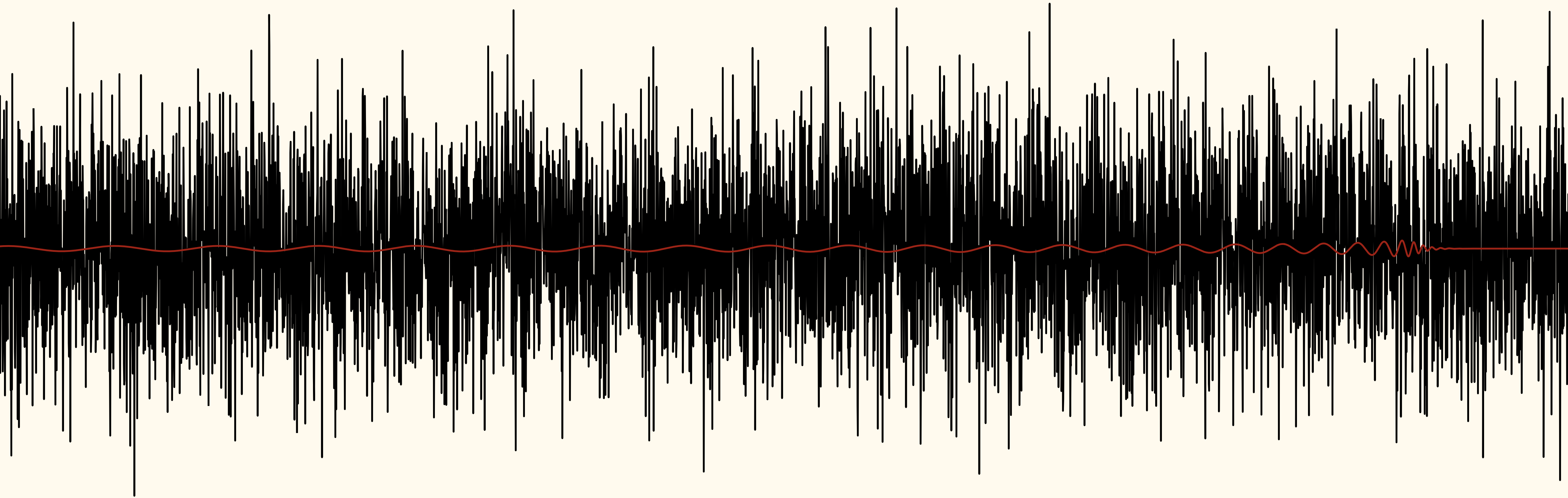


3. data analysis

INSTRUMENTAL NOISE



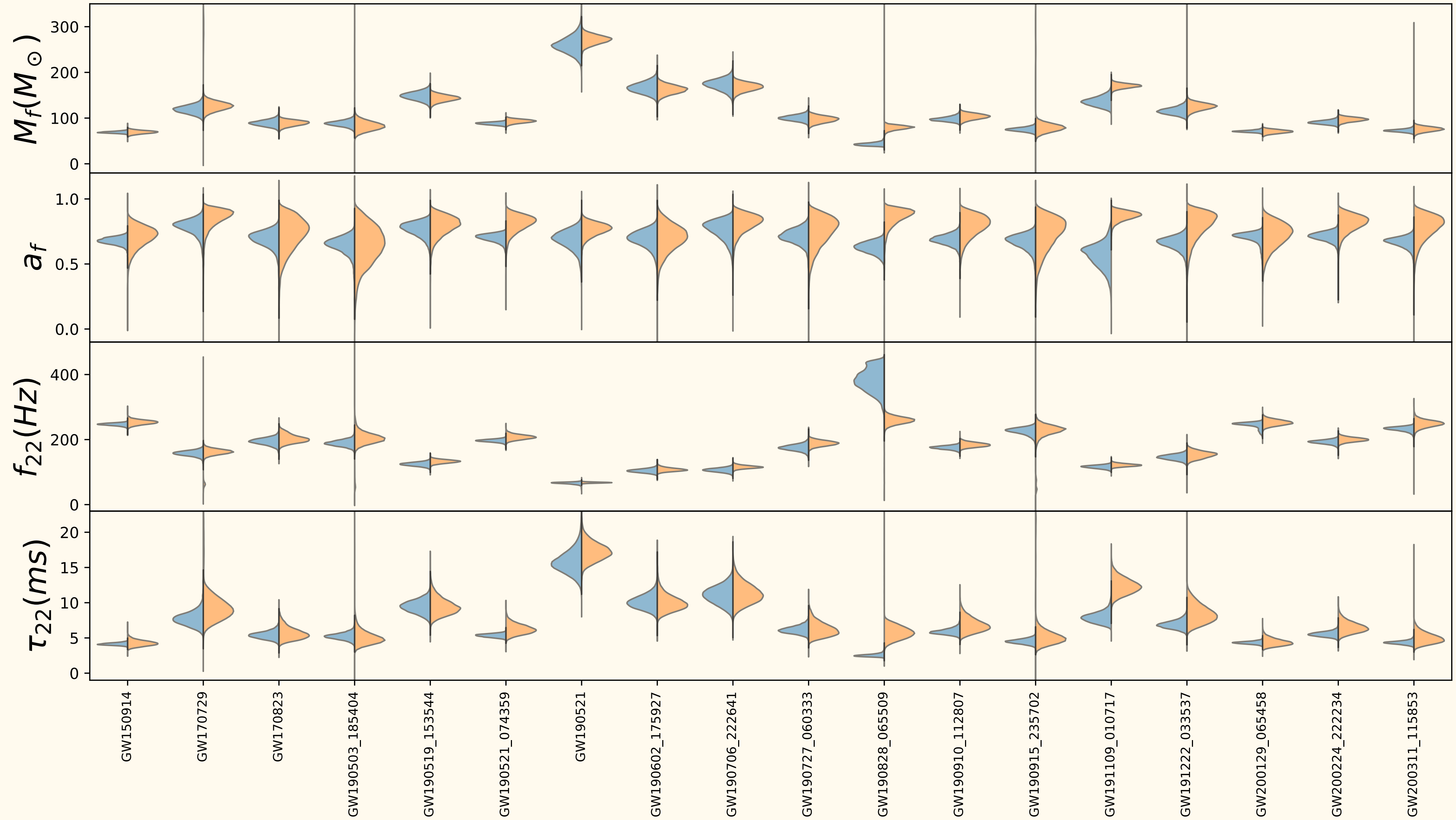
INSTRUMENTAL NOISE



we used TEOBPM for a time domain analysis of LIGO-Virgo data with pyRing

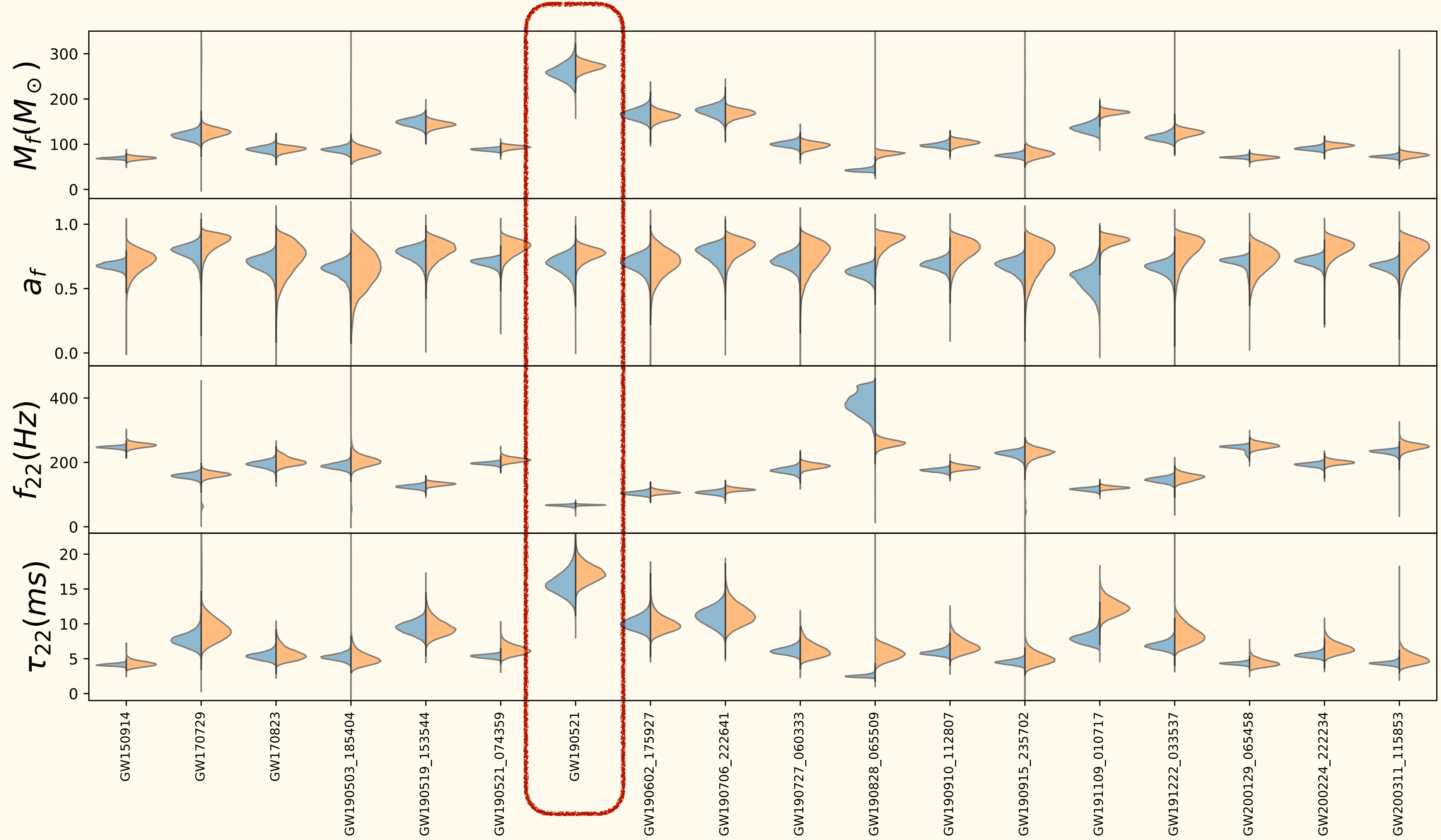
GWTC-3 ANALYSIS

LVK in IMR
TEOBPM in RD



GWTC-3 ANALYSIS

LVK in IMR
TEOBPM in RD



DID WE OBSERVED HMS?

observations of HMs in the literature:

RD-only analyses

event	reference	$\ln \mathcal{B}_{lm,22}$	coalescence	type
GW190521A	Capano et al. (2021)	3.8	IMBH	RD

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this work

no HMs in GW190521

$$\ln B_{33,22} = 0.13$$

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precession? eccentricity? other?
2009.01075 2106.05575

further analyses ongoing...

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this work

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 $\ln B_{33,22} = 0.13$

precession? eccentricity? other?
 2009.01075 2106.05575

further analyses ongoing...

IMR analyses

event	reference	$\ln \mathcal{B}_{lm,22}$	coalescence	type
GW170729	Chatziioannou et al. (2019)	1.6	BBH	IMR
GW190814A	Abbott et al. (2020c)	22.1	BH-(?)	IMR
GW190412A	Abbott et al. (2020a)	8.3	BBH	IMR

RD weakly measured
and not informative

4. tests of no-hair

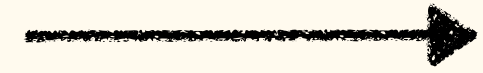
NO-HAIR THEOREM

why are higher modes important?

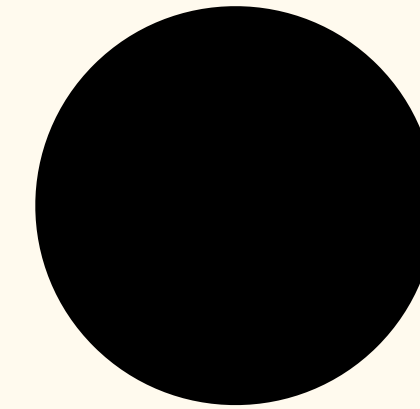
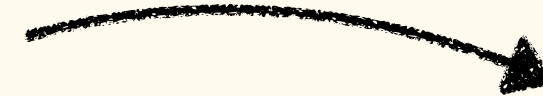
NO-HAIR THEOREM

why are higher modes important?

ω_{22}, τ_{22}

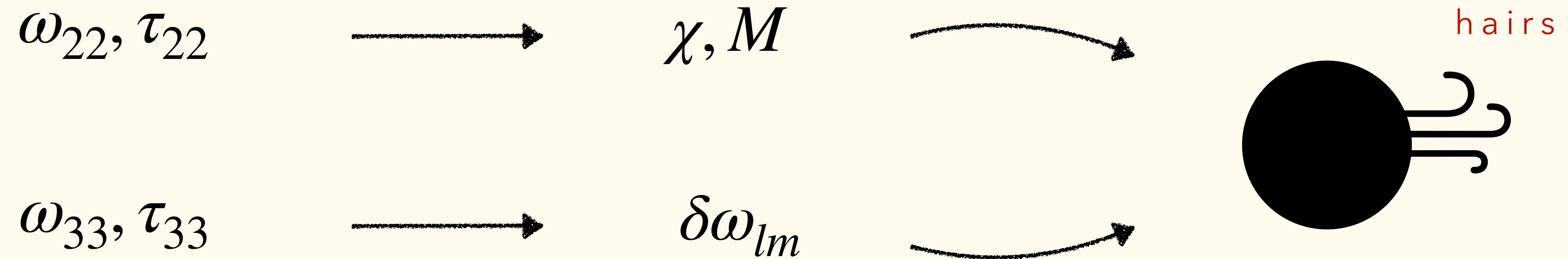


χ, M



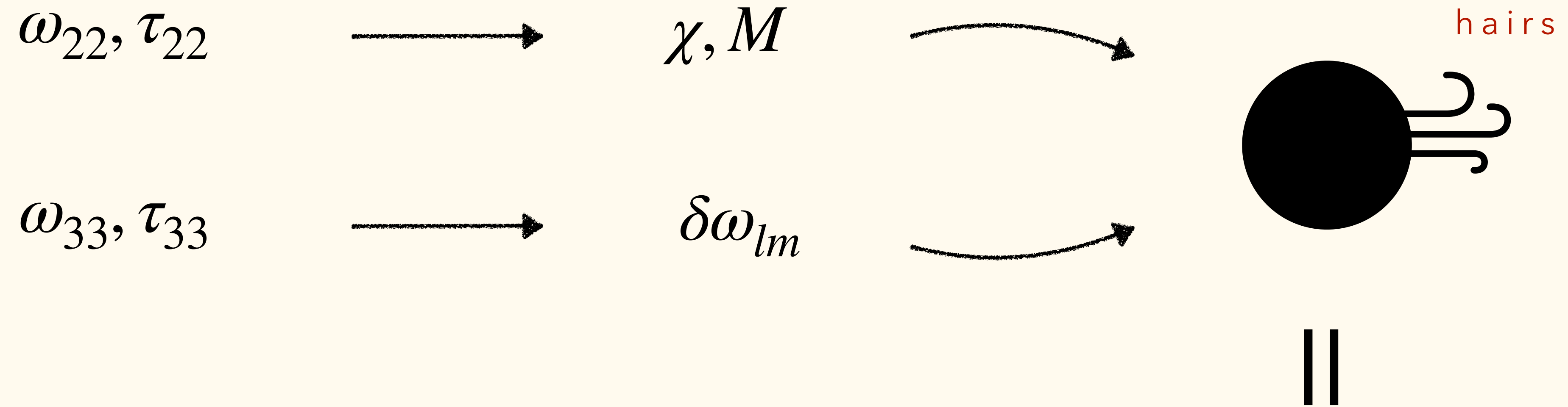
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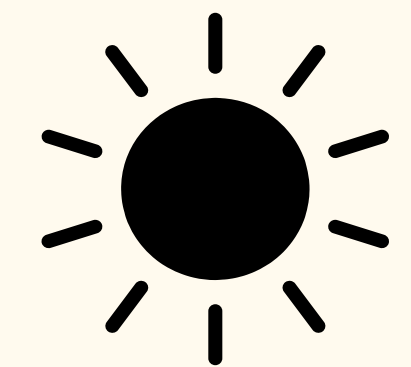
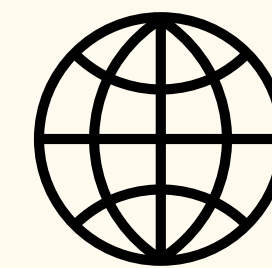
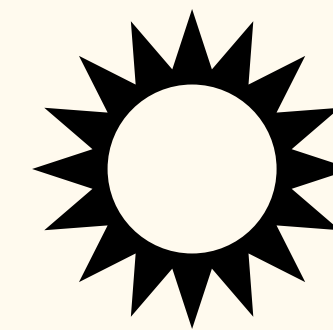


NO-HAIR THEOREM

why are higher modes important?



tests of the **no-hair theorem**



SUMMARY

- RD provides unique access to strong field and extreme curvature regimes
- the RD waveform is a superposition of quasinormal modes
- the excitation of HMs strongly depends on mass ratio and inclination
- we used a RD model that includes post-merger nonlinearities
- we developed an analytical procedure to predict the detectability of HMs
- we analysed the RD signals in GWTC-3 using TEOBPM
- we found results consistent with LVK and no HMs in GW190521
- importance of our results for future tests of the no-hair theorem

backup slides

GW SIGNAL

$$h_+ - ih_\times = \sum_{lm} A_{lm} e^{-i\omega_{lm}t - t/\tau_{lm}} {}_2Y_{lm}$$

$$h_{ij} = h_+ e_{ij}^+ + h_\times e_{ij}^\times$$

$$h \equiv F_+ h_+ + F_\times h_\times$$

ringdown waveform

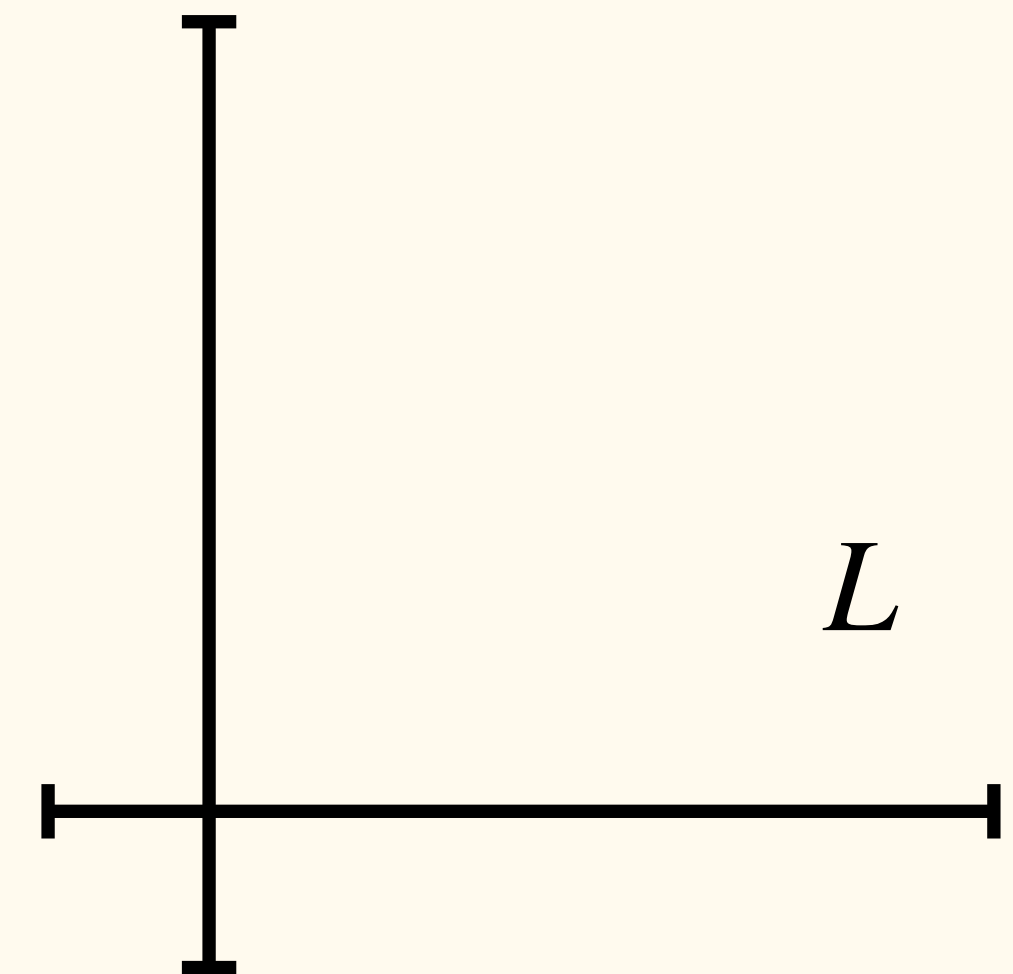
polarization tensors

antenna functions

differential length of
the **interferometer**

$$h = \frac{\Delta L}{L}$$

GW interaction is encoded in the
amplitude and **phase** of the laser output



BAYESIAN ANALYSIS

- parameter estimation

from Bayes theorem, we can find the probability density of the parameters

$$\frac{p(\mathbf{d} | \boldsymbol{\theta}, H) p(\boldsymbol{\theta} | H)}{p(\mathbf{d} | H)} = p(\boldsymbol{\theta} | \mathbf{d}, H)$$

data

parameters

posterior distribution

- model selection

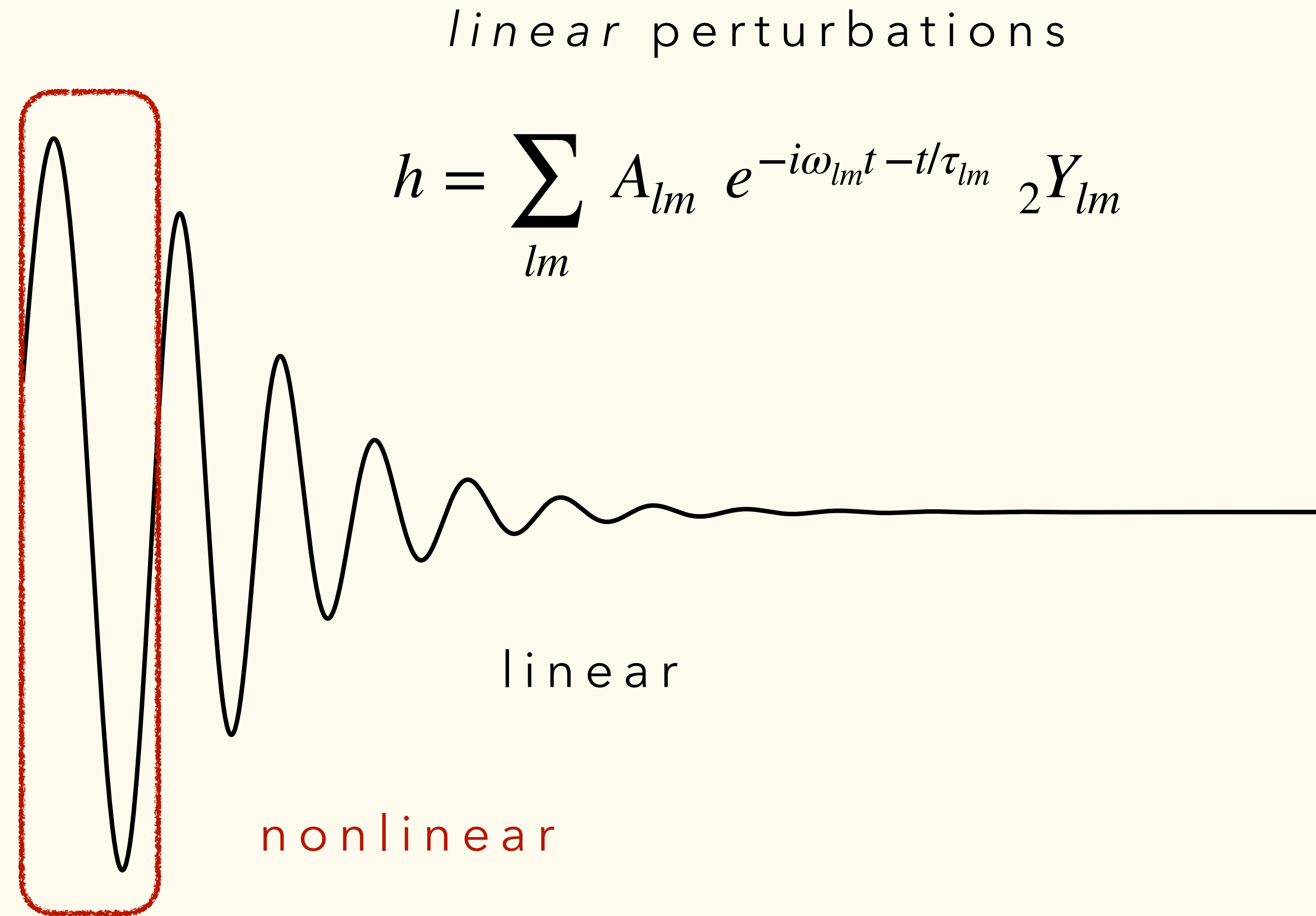
the Bayes factor tells us which model better describes the data

$$\frac{p(\mathbf{d} | H_1)}{p(\mathbf{d} | H_2)} \equiv B_{1,2}$$

RD STARTING TIME

when the RD starts?

- too late, surely linear but lose all the signal
- too early, linear model to nonlinear data



difficult to choose the starting time

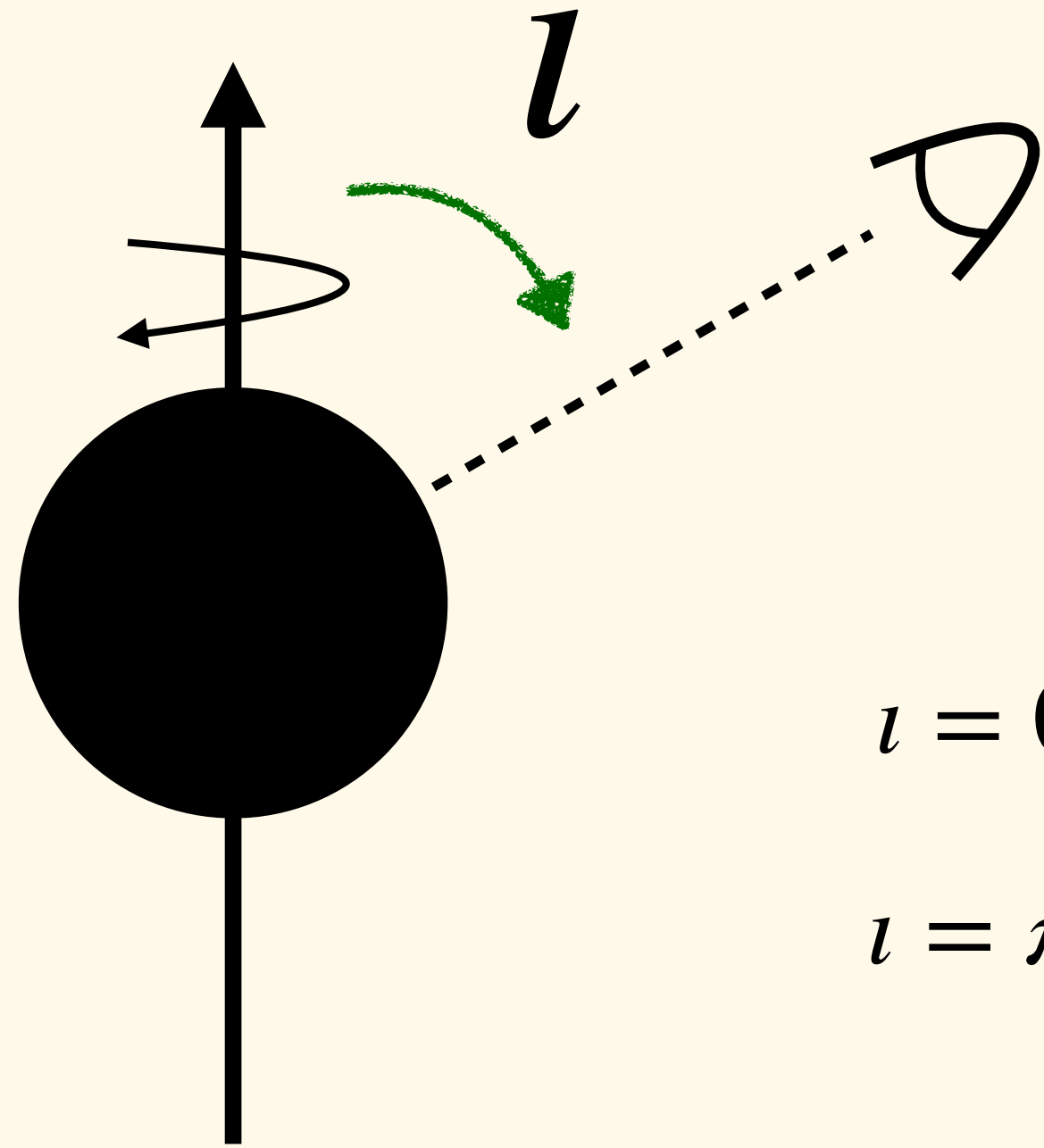


results depend on the starting time

SPHERICAL HARMONICS

$${}_2Y_{lm}(\iota, \varphi)$$

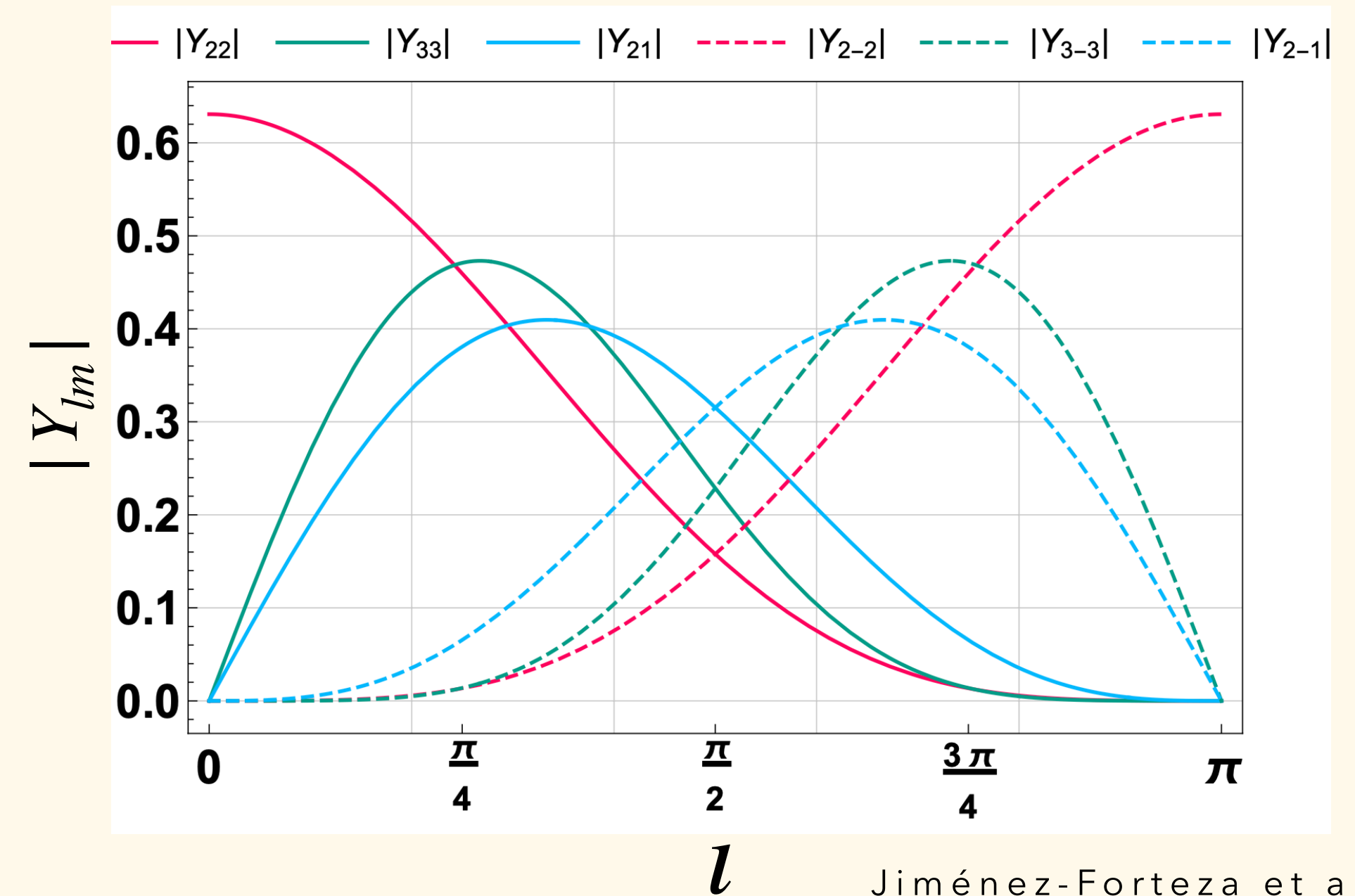
ι is the **inclination**



$\iota = 0$ face-off
 $\iota = \pi/2$ edge-on

$$h = \sum_{lm} A_{lm} e^{-i\omega_{lm}t - t/\tau_{lm}} \boxed{{}_2Y_{lm}}$$

(spin-weighted)
spherical harmonics



Jiménez-Forteza et al. (2020)
2005.03260

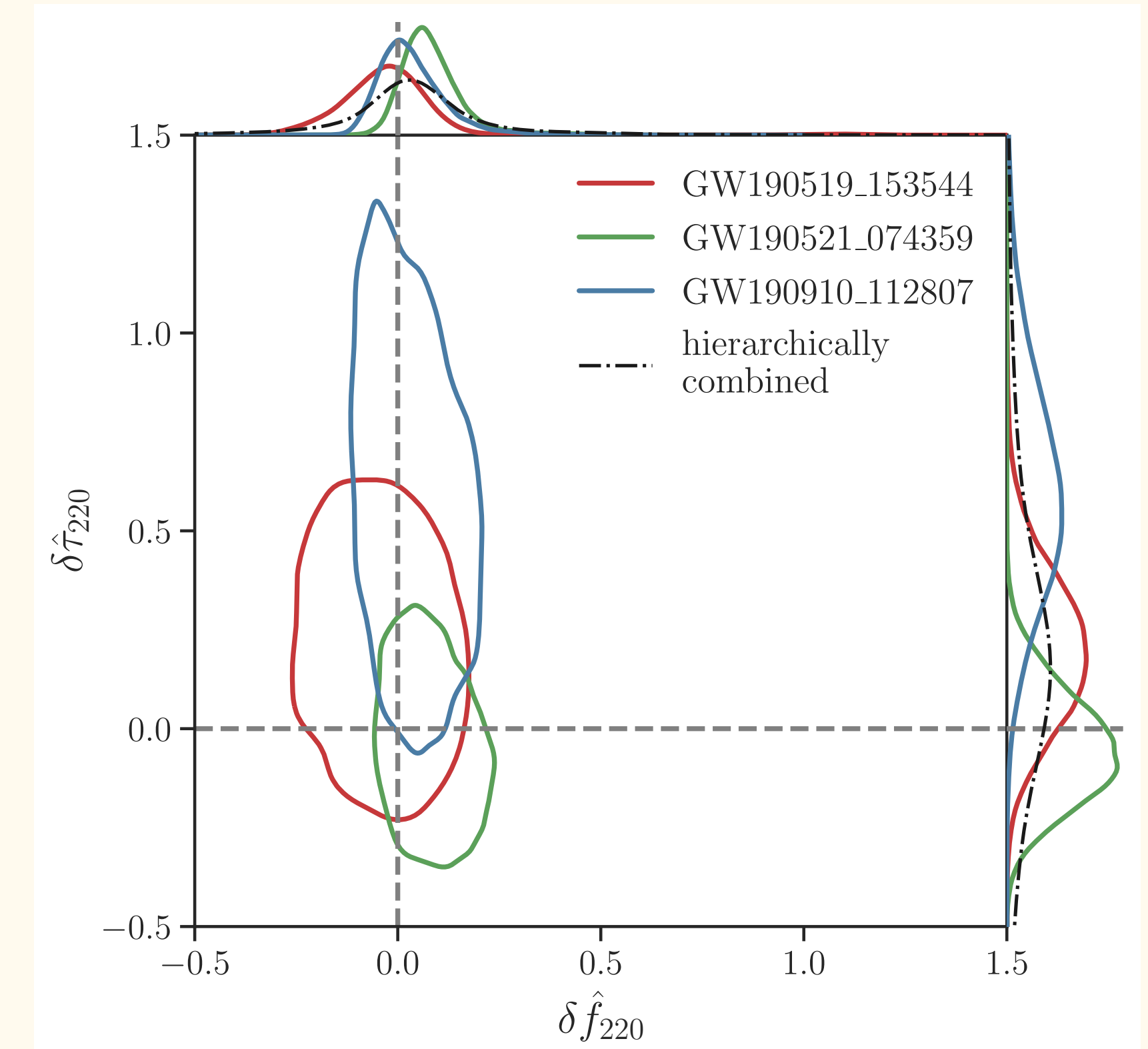
TESTS OF NO-HAIR

consider **fractional deviations** from GR

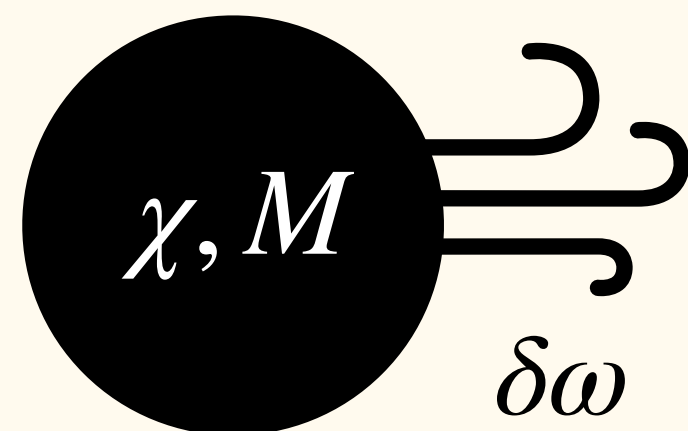
$$\omega_{lm} = \omega_{lm}^{GR} (1 + \delta\omega_{lm})$$

$$\tau_{lm} = \tau_{lm}^{GR} (1 + \delta\tau_{lm})$$

if the posteriors on $\delta\omega_{lm}$ and $\delta\tau_{lm}$ support zero, then GR is correct



Abbott et al. (2020)
2010.14529

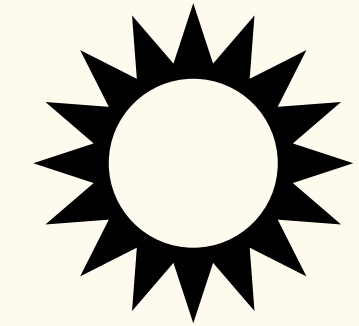


- $\delta\omega_{lm}, \delta\tau_{lm}$ are generic additional degrees of freedom
- possibility to map $\delta\omega_{lm}, \delta\tau_{lm}$ to specific theories

ALTERNATIVE SCENARIOS

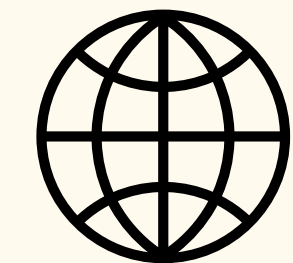
- are we really observing black holes?

exotic compact objects (boson star, gravastar, fuzzball, ...)



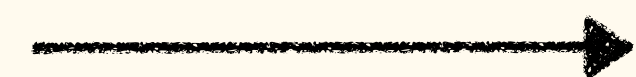
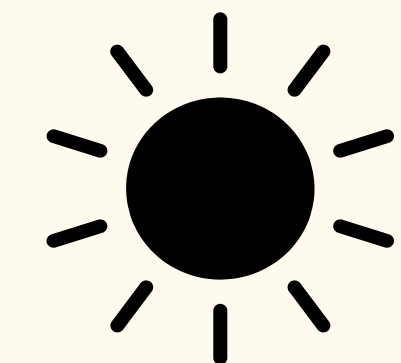
- is GR the correct theory of gravity?

modified theories of gravity (EdGb, dCS, EFT, ...)



- are there quantum effects at the horizon?

area quantisation (Bekenstein-Mukhanov conjecture),
BH entropy (Bekenstein-Hod bound)



systematics? unmodeled properties? *environmental effects?*