

Holographic aspects of four-dimensional asymptotically flat $N = 2$ BPS black holes

Gabriel Lopes Cardoso

XV Black Holes Workshop

with Pedro Aniceto and Suresh Nampuri, arXiv:2111.13190

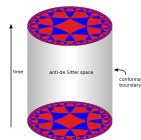
with A. Kidambi, S. Nampuri, V. Reys, M. Rosselló,
arXiv:2211.06873



BPS black holes

AdS/CFT correspondence: quantum gravity in an asymptotically AdS space-time encoded in CFT on its conformal boundary.

BPS black hole solutions in **four dimensions:**
asymptotically flat, near-horizon geometry $AdS_2 \times S^2$.



AdS_2/CFT_1 Sen, arXiv: 0805.0095

How is BPS black hole entropy encoded in CFT_1 ?

A conformal quantum mechanics model: DFF model.
de Alfaro, Fubini, Furlan, 1976

BPS black holes:

extremal, supersymmetric, **charged** (q_i, p^i)

4D, $N = 2$ supergravity theories reduced on S^2

Einstein-Maxwell-dilaton type actions with additional scalar fields Y^I .

Dimensional reduction on two-sphere S^2 :

$$ds_4^2 = ds_2^2 + v_2(t, r) \left(d\theta^2 + \sin^2 \theta d\varphi^2 \right) \quad , \quad v_2 > 0$$

$$ds_2^2 = h_{ij} dx^i dx^j \quad , \quad x^i = t, r \quad ,$$

$$F_{rt}^I = e^I(t, r) \rightarrow q^I \quad , \quad F_{\theta\varphi}^I = p^I \sin \theta \quad , \quad Y^I(t, r)$$

Rescaling $h_{ij} = \tilde{h}_{ij} / \sqrt{v_2}$, get 2D action in the form

$$L_2 = \frac{1}{2} \sqrt{-\tilde{h}} \left(-v_2 \tilde{R}_2 + U(v_2, q, p, Y) + v_2 (\partial Y)^2 + \dots \right)$$

Expanding $v_2(t, r) = \phi_0 + \epsilon \phi(t, r)$, $Y^I(t, r) = Y_0^I + \epsilon \mathcal{Y}^I(t, r)$,

$$L_2 = -\frac{1}{2} \sqrt{-\tilde{h}} \phi_0 \tilde{R}_2 + \frac{1}{2} \sqrt{-\tilde{h}} \phi \left(-\tilde{R}_2 + \frac{2}{v_1^{3/2}(q, p)} \right) + \mathcal{O}(\phi^2, (\mathcal{Y})^2, (\partial \mathcal{Y})^2)$$

JT gravity type action.

First-order perturbations

Constant curvature solution supported by **constant scalar fields** v_2 and Y^I : in Fefferman-Graham gauge

$$\begin{aligned} ds_2^2 &= h_{ij} dx^i dx^j = dr^2 + h_{tt}(t, r) dt^2 \\ \sqrt{-h_{tt}} &= \alpha(t) e^{r/\sqrt{v_1}} + \beta(t) e^{-r/\sqrt{v_1}} = \sqrt{-h_0} \\ v_2 &= v_1(q, p) \quad , \quad Y_0^I(q, p) \end{aligned}$$

$$R_2 = 2/v_1(q, p).$$

Near-horizon geometry of a **BPS black hole in 4D**:

Sen, arXiv: 0809.3304

black hole in $AdS_2 (\times S^2)$

$$ds_2^2 = v_1 [-(\rho^2 - 1) dt^2 + d\rho^2 / (\rho^2 - 1)] \quad , \quad \rho > 1 \quad \alpha = 1, \beta = -1$$

First-order perturbations:

$$\sqrt{-h} = \sqrt{-h_0} + \epsilon \sqrt{-h_1} \quad , \quad v_2 = v_1(q, p) + \epsilon \phi \quad , \quad Y^I = Y_0^I + \epsilon Y^I$$

First-order deformation

$$\begin{aligned}\sqrt{-h_1} &= -\frac{\sqrt{-h_0}}{v_1} \phi - 2 \partial_t \left(\frac{\partial_t \nu}{\alpha} \right) \\ \phi &= \nu(t) e^{r/\sqrt{v_1}} + \vartheta(t) e^{-r/\sqrt{v_1}} \\ \left(\square - \frac{2}{v_1} \right) \bar{y}' &= 0\end{aligned}$$

with

$$\vartheta = \frac{1}{\nu} \left(c_0 - \frac{v_1}{4} \left(\frac{\nu'}{\alpha} \right)^2 \right), \quad \beta = -\frac{\alpha}{\nu^2} \left(c_0 - \frac{v_1}{4} \left(\frac{\nu'}{\alpha} \right)^2 \right) - \frac{v_1}{2\nu} \left(\frac{\nu'}{\alpha} \right)'$$

Example: take a black hole in AdS_2 , deform by a constant $\nu = \nu_0$

$$\alpha = \frac{1}{\sqrt{v_1}}, \quad \beta = -\frac{\mu^2}{4\sqrt{v_1}} \implies c_0 = \frac{\mu^2}{4} \nu_0^2 > 0$$

Now $\nu(t) = \nu_0 + f(t)$, $\bar{y}' = 0$:

Lorentzian DFF / Euclidean 1D Liouville type action

Dynamics of $\nu(t)$ described by the renormalized on-shell action:

$$S_{ren} = \int_{\partial M} dt \left(\frac{c_0}{\nu} + \frac{(\nu')^2}{4\nu} \right), \quad \nu(t) = x^2(t)$$

$$S_{DFF} = \int_{\partial M} dt \left(\frac{c_0}{x^2(t)} + (x')^2 \right)$$

Conformal mechanics Lagrangian of DFF. $c_0 > 0$, 'wrong' sign.

Converting to global time, C, Kidambi, Nampuri, Reys, Rosselló, arXiv:2211.06873



$$S_{DFF} = \int dt \left[\frac{(\nu')^2}{4\nu} + c_0 \left(\frac{1}{\nu} + \nu \right) \right]$$

$$S_{Liouv} = \int dt \left[\frac{1}{2} (l')^2 + 2e^{-l} \right]$$

Holographic model for BPS BHs in $D = 4$.

Lin, Maldacena, Rozenberg, Shan, 2207.00408