

Piercing of a boson star by a black hole

Zhen Zhong

Vitor Cardoso, Taishi Ikeda, Miguel Zilhão

ISCTE - University Institute of Lisbon

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Motivation

- Evidence for existence of matter of unknown nature and properties which interacts gravitationally
- We can expect “dark stars”, making up a significant fraction of astrophysical environments
- We will assume the existence of new fundamental, scalar degrees of freedom minimally coupled to gravity
- For complex scalar fields, boson stars (BSs) form rather generically as a consequence of gravitational collapse
- There are indications that such solutions describe well dark matter cores in haloes.
 - These models are often referred to as fuzzy dark matter models

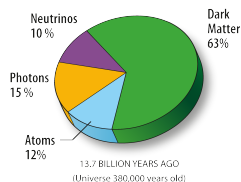
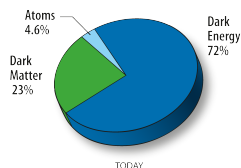


Figure: WMAP (2008)

Motivation

- Dark stars have so far gone undetected, Understanding the behavior of dark matter when moving perturbers drift by may be crucial for detections
- **Here we consider the motion of a small black hole as it pierces a large boson star structure**

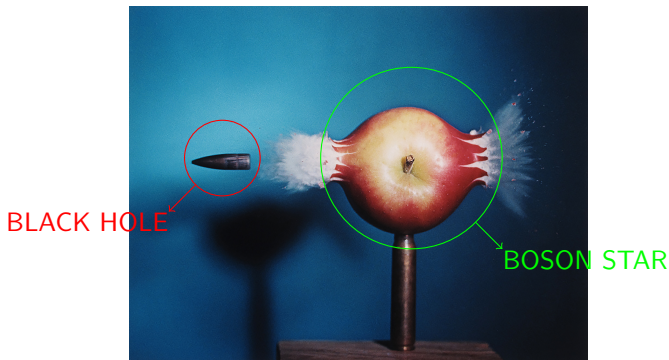


Figure: Harold Edgerton, Bullet Through Apple (1964)

Figure: The motion of a small black hole as it pierces a large boson star structure

Model

A **minimally coupled**, massive scalar field Φ described by the Einstein-Klein-Gordon action

$$S = \int d^4x \sqrt{-g} \left[\frac{\mathcal{R}}{16\pi} - (g^{ab} \nabla_a \Phi \nabla_b \Phi^* + \mu^2 \Phi \Phi^*) \right]$$

- BSSN
- Einstein Toolkit
 - Canuda



Figure: Einstein Toolkit

- Spherically-symmetric mini-boson star

$$\begin{aligned}\Phi(t, R) &= \phi(R)e^{i\omega t} \\ dl^2 &= \psi_{\text{BS}}(R)^4 (dx^2 + dy^2 + dz^2)\end{aligned}$$

- Boosted Schwarzschild black hole

$$dl^2 = \left(1 + \frac{M}{2R}\right)^4 (dx^2 + dy^2 + B_0(x, y, z)dz^2)$$

- Superposition of the conformal factors

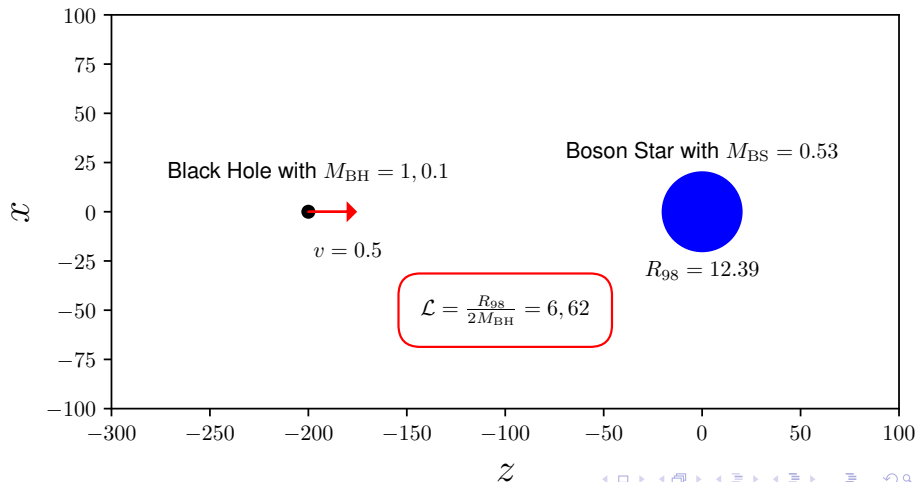
$$\begin{aligned}dl^2 &= \psi^4 (dx^2 + dy^2 + B(x, y, z)dz^2) \\ \psi &= 1 + \frac{M_{\text{BH}}}{2R} + \psi_{\text{BS}}\end{aligned}$$

Initial Data: Boosted Black Hole + Boson Star

Run	M_{BH}	$R_{98}/(2M_{\text{BH}})$	v_0	M_{tot}
IB	1	6	0.5	1.68
IVB	0.1	62	0.5	0.65

$$dl^2 = \psi^4 (dx^2 + dy^2 + B(x, y, z)dz^2)$$

$$\psi = 1 + \frac{M_{\text{BH}}}{2R} + \psi_{\text{BS}}$$



Time Evolutions

Run	M_{BH}	$R_{98}/(2M_{\text{BH}})$	v_0	M_{tot}
IB	1	6	0.5	1.68
IVB	0.1	62	0.5	0.65

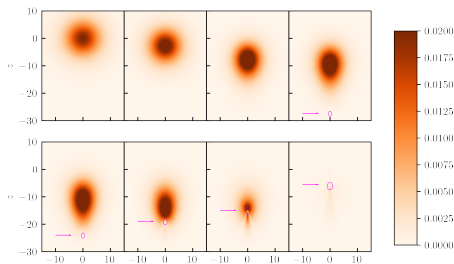


Figure: Run IB

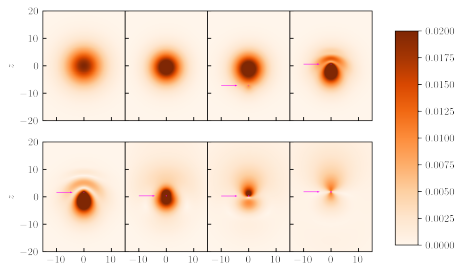


Figure: Run IVB

Accretion during collision

- even for case IVB: a BH which is 62 times smaller than the BS and moving at half the speed of light.
- after a relatively short timescale ~ 1000 , and after the BH has had time to interact with the BS, the final BH mass is at least 95% of the total initial energy.

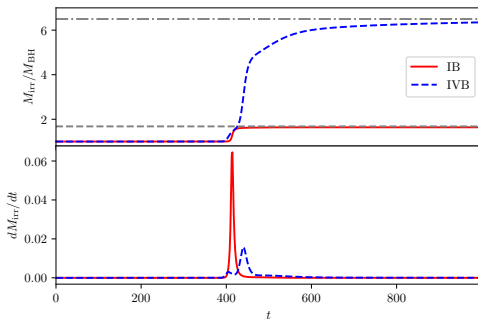


Figure: Accretion of scalar onto the BH, **BS is nearly totally accreted**

Black hole motion

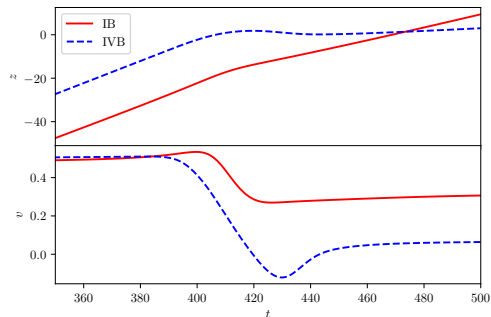


Figure: The location z and the speed v of the BH

- Interaction between BH and BS is clear, and translates into a deceleration starting at $t \sim 400$ lasting for time to cross bulk of BS
- Notice that the BH velocity is negative for a small amount of time, in simulation IVB: the BH is tidally captured by the boson star.

- $M_{\text{BH}} = 1$
- $\mathcal{L} \sim 6$
- $v_0 = 0.5$
- $M_{\text{tot}} = 1.68$

Figure: Movie for Run IB

- $M_{\text{BH}} = 0.1$
- $\mathcal{L} \sim 62$
- $v_0 = 0.5$
- $M_{\text{tot}} = 0.65$

Figure: Movie for Run IVB

- Challenging simulations of boosted BHs and BSs
- Even for length ratios ~ 62 the BS is entirely accreted by the BH
- At late times, a “gravitational atom” is formed: massive BH surrounded by a quasi-bound state of the scalar field (the BS remnant)
- One can expect that for even smaller BH masses the BH can pierce through without destroying the BS
 - In order to evolve such a small BH, one would need too large an amount of computational resources...
- Spinning and/or charged configurations may lead to new phenomena