Coherent quantum states in resonant-mass gravitational

wave detectors

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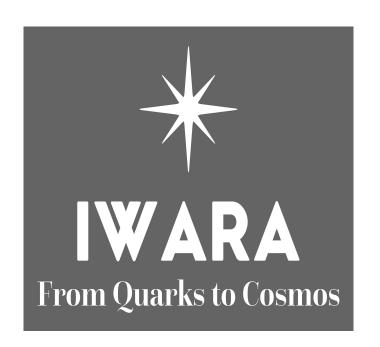
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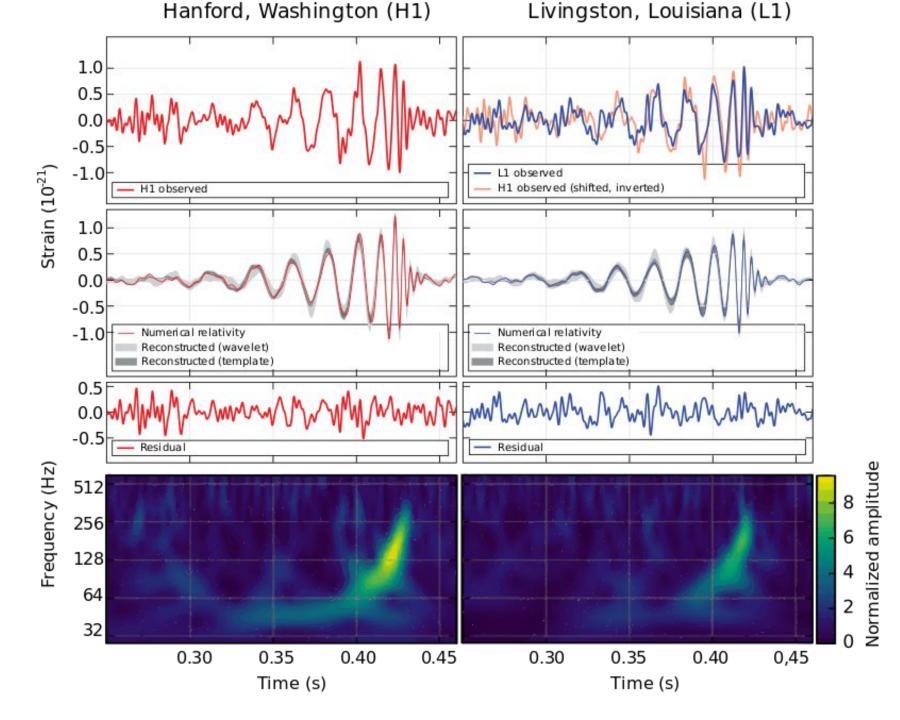
Abstract

Quantum coherent states is a new features in quantum mechanics. Resonant-mass gravitational waves detectors didn't succeed in detecting gravitational waves probably because of the operational frequency range chosen when the design of such detectors. But such detectors can be important in the future. This work describe the importance of quantum coherent states when the a resonant mass gravitational detects a energy from the gravitational wave very very close to the quantum limit. At this point quantum coherent states of vibration and quantum non-demolition measurements becomes essential.

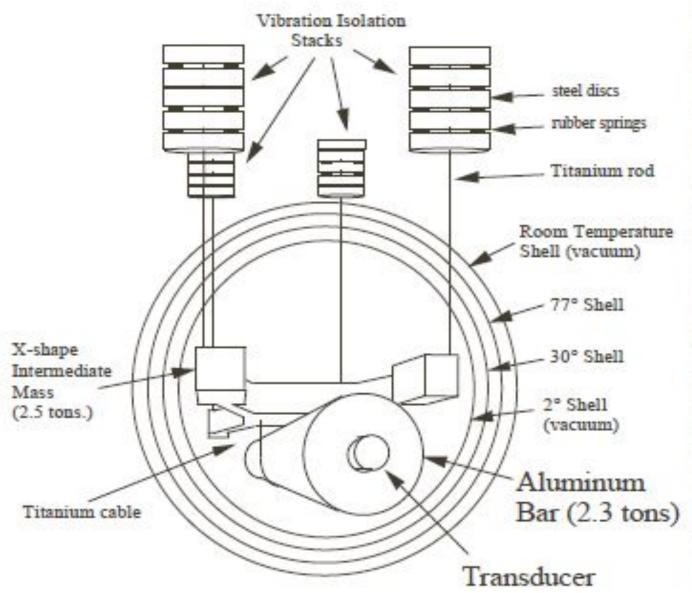
Outline

- Resonant mass gravitational wave detectors
- The GW detection frequency range
- IWARA 2020 presentation: quantum limit
- The classical solution
- Coherent quantum states
- The quantum solution

First
Gravitational
Wave
Detection



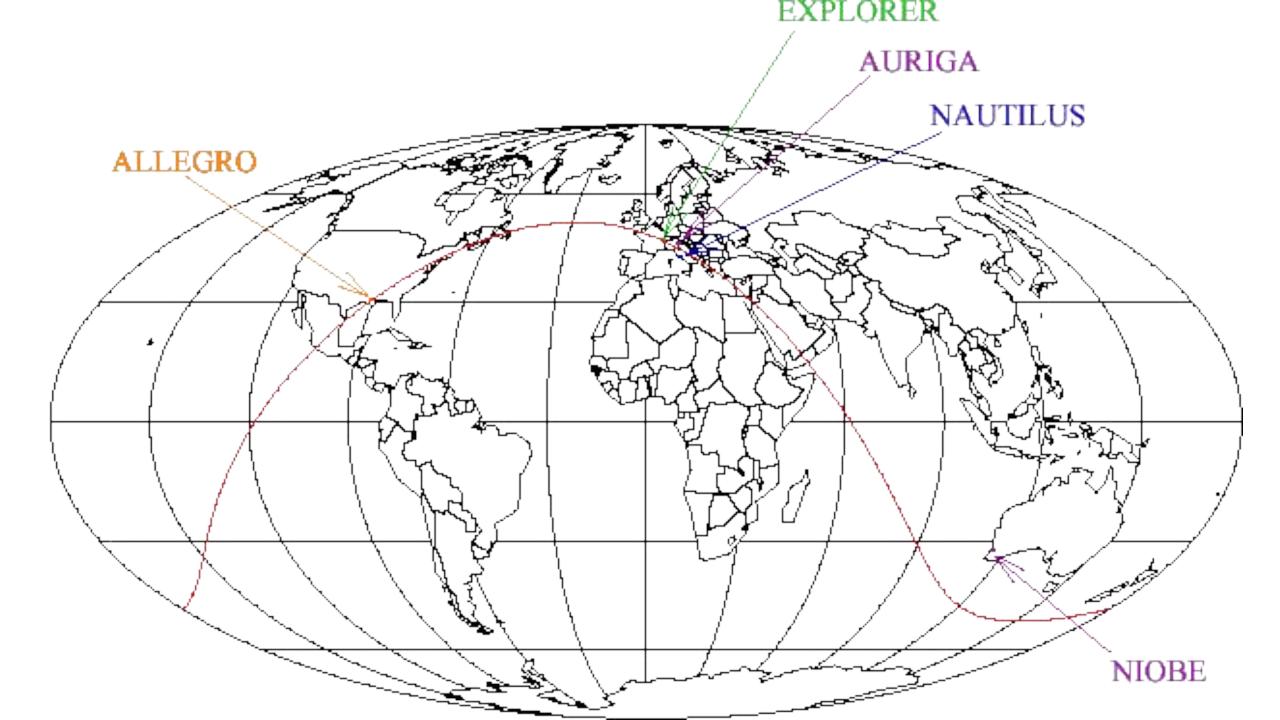
Resonant-mass GWD Easily calibrated.





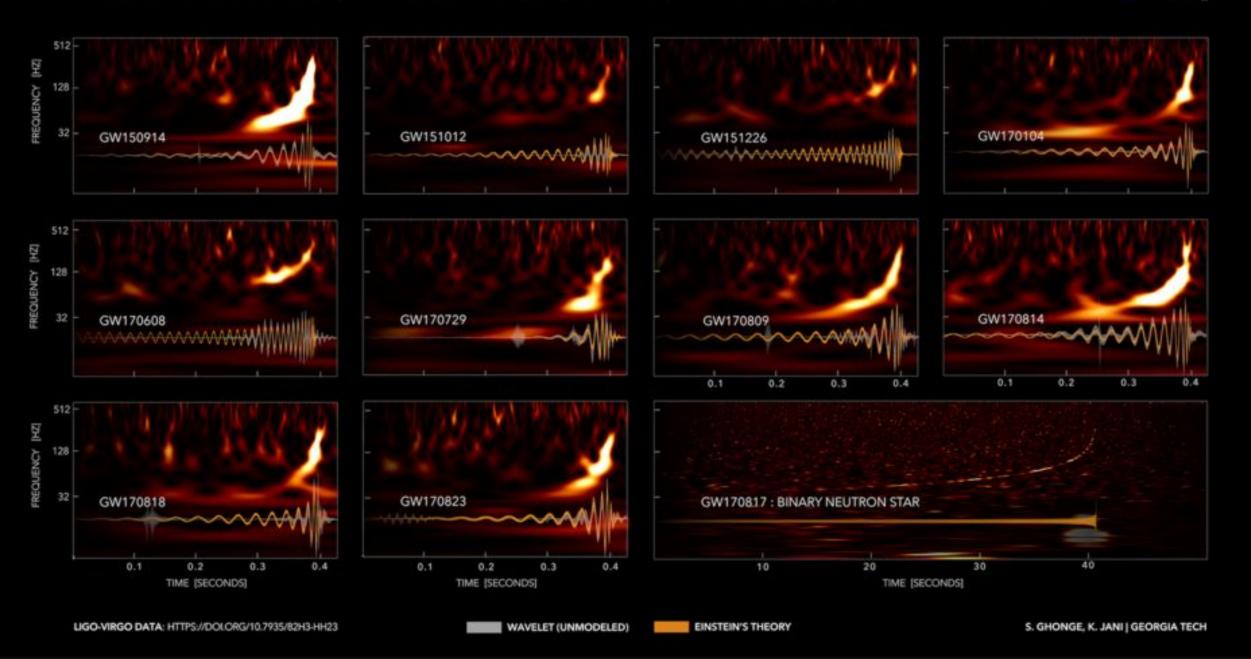


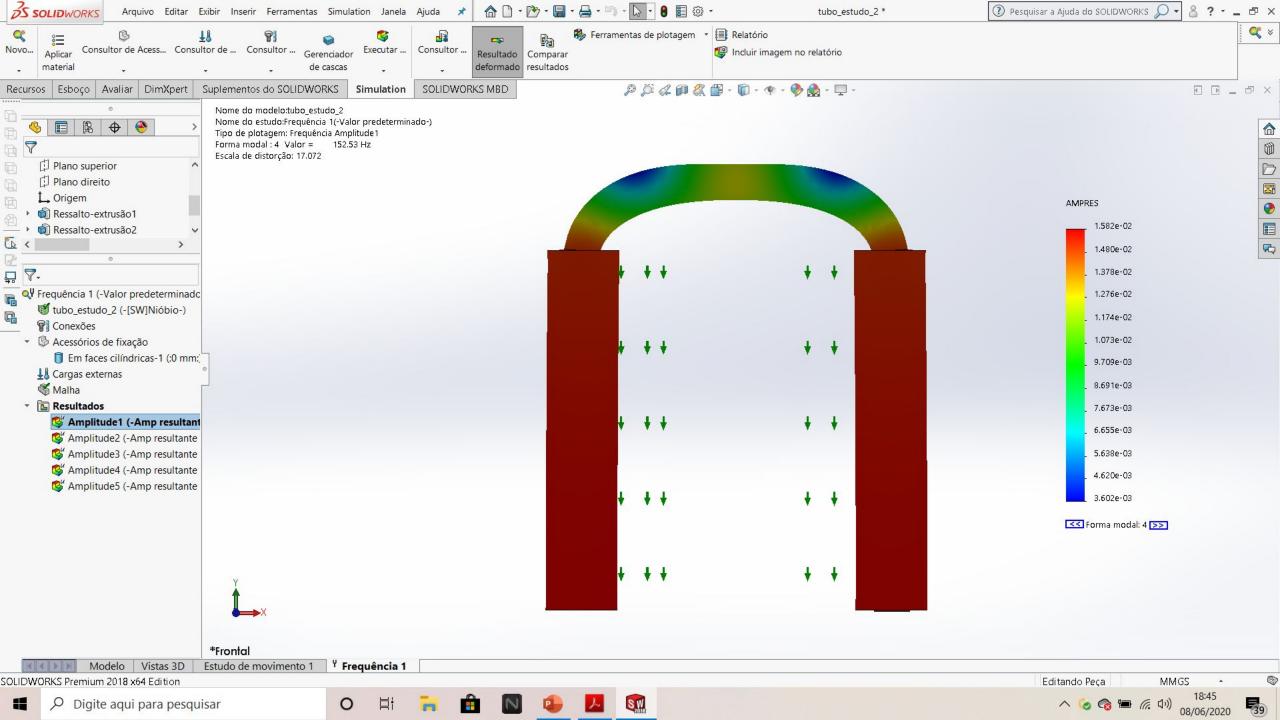


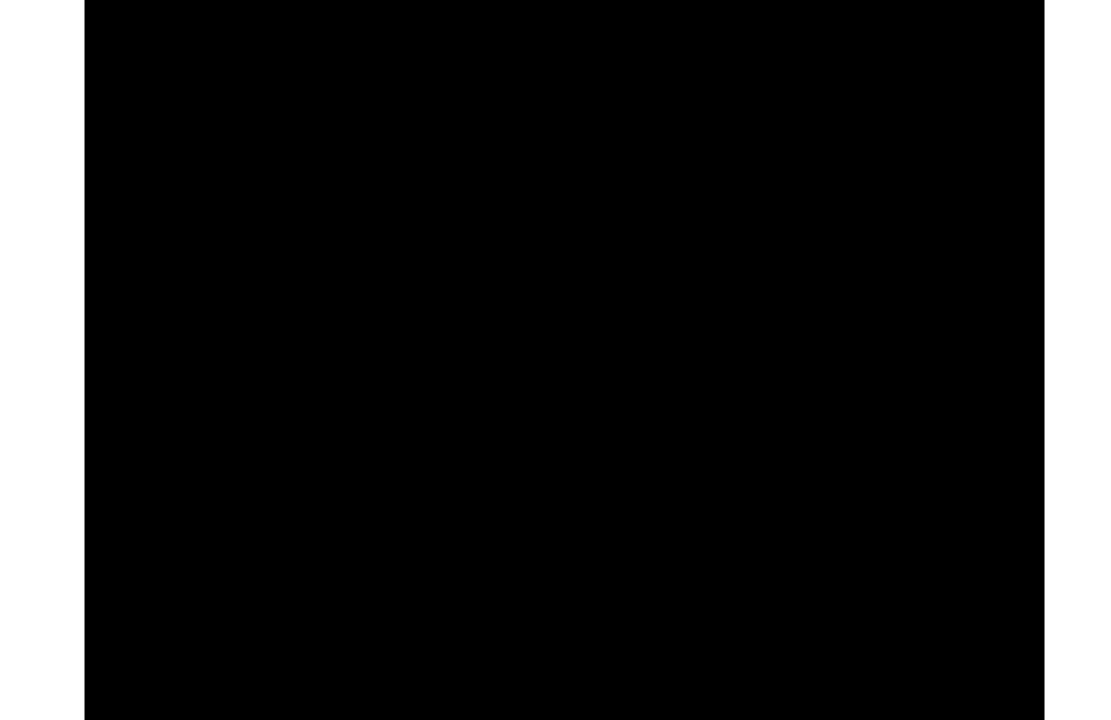


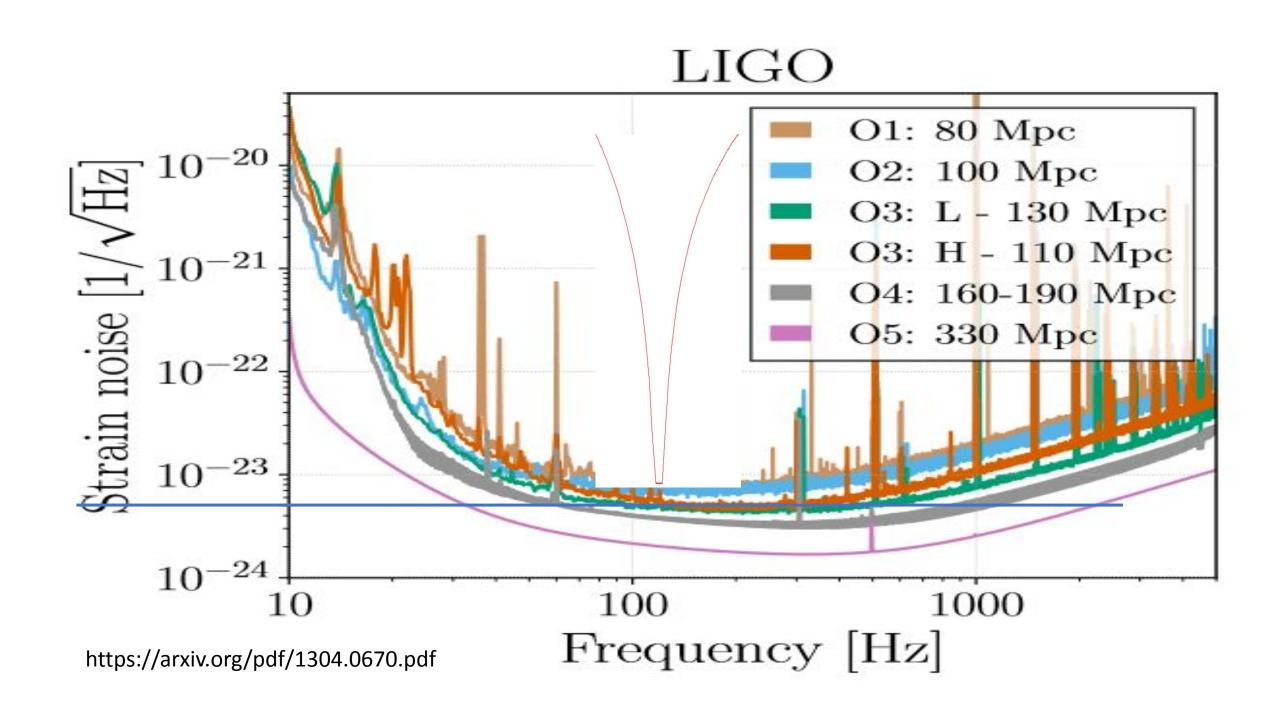
GRAVITATIONAL-WAVE TRANSIENT CATALOG-1

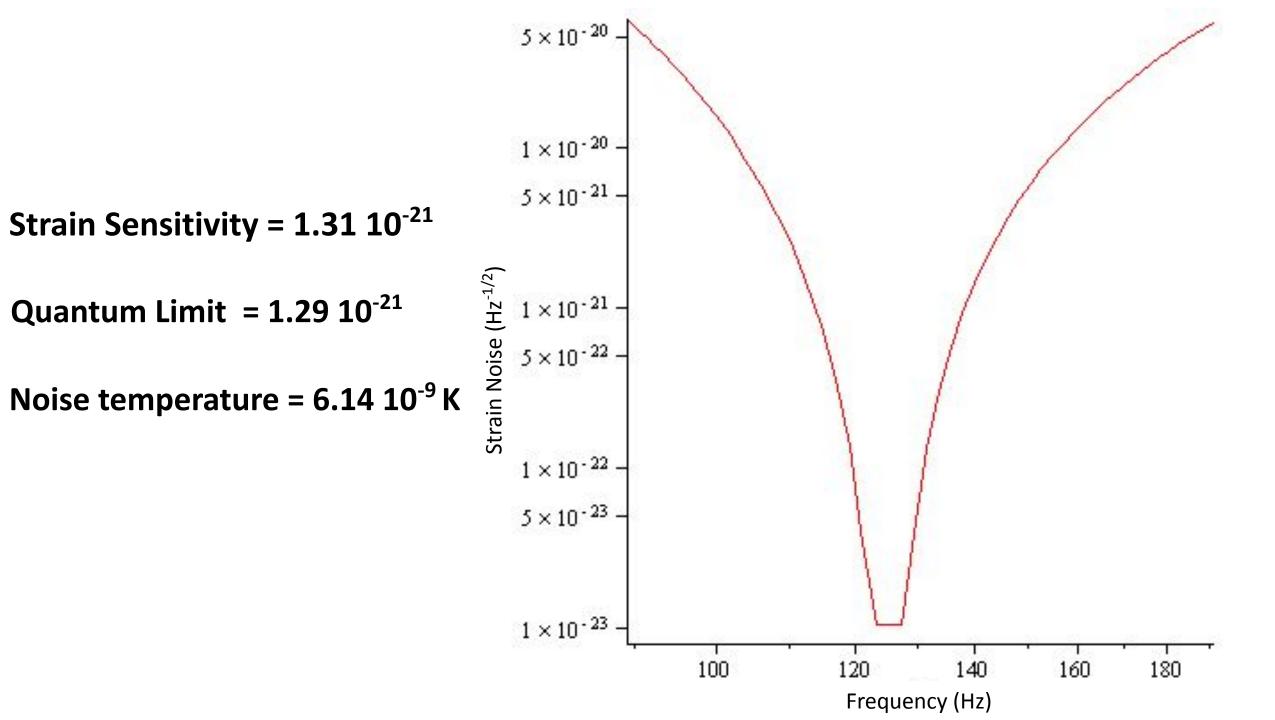




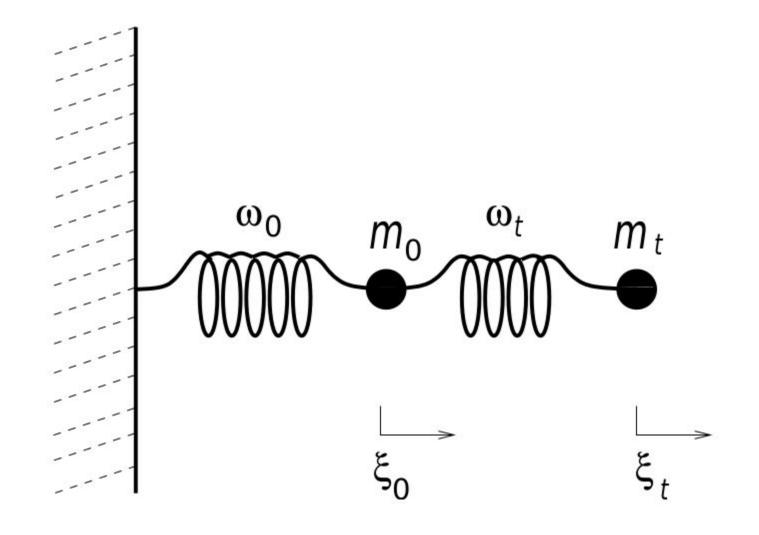


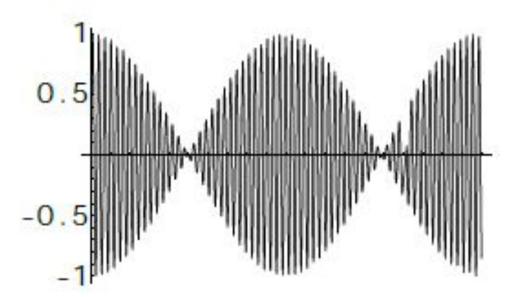


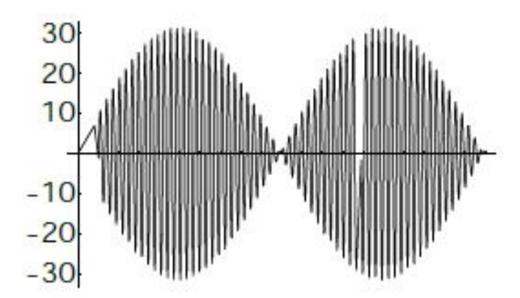




The force actuating in the mass is classical.







$$H = \frac{p_0^2}{2m_0} + \frac{p_t^2}{2m_t} + \frac{1}{2}k_0\xi_0^2 + \frac{1}{2}\kappa(\xi_t - \xi_0)^2$$

$$\Xi_{0} = \left(\frac{m_{0}}{m_{t}}\right)^{1/4} \xi_{0}; \quad \Xi_{t} = \left(\frac{m_{t}}{m_{0}}\right)^{1/4} \xi_{t}$$

$$P_{0} = \left(\frac{m_{t}}{m_{0}}\right)^{1/4} p_{0}; \quad P_{t} = \left(\frac{m_{0}}{m_{t}}\right)^{1/4} p_{t}$$

$$\mu \equiv (m_0 m_t)^{1/2}$$

$$H = \frac{P_0^2}{2\mu} + \frac{P_t^2}{2\mu} + \frac{1}{2}\mu\omega^2\Xi_0^2 + \frac{1}{2}\left[\left(\frac{m_0}{m_t}\right)^{1/4}\Xi_t - \left(\frac{m_t}{m_0}\right)^{1/4}\Xi_0\right]$$

$$\begin{bmatrix} \xi_{+} \\ \xi_{-} \end{bmatrix} = \begin{bmatrix} \cos \alpha & \sin \alpha \\ -\sin \alpha & \cos \alpha \end{bmatrix} \begin{bmatrix} \Xi_{0} \\ \Xi_{t} \end{bmatrix} \\
\begin{bmatrix} p_{+} \\ p_{-} \end{bmatrix} = \begin{bmatrix} \cos \alpha & \sin \alpha \\ -\sin \alpha & \cos \alpha \end{bmatrix} \begin{bmatrix} P_{0} \\ P_{t} \end{bmatrix}$$

$$H = \frac{p_+^2}{2\mu} + \frac{1}{2}\mu\omega_+^2\xi_+^2 + \frac{p_-^2}{2\mu} + \frac{1}{2}\mu\omega_-^2\xi_-^2$$

$$\omega_{+} = \left\{ \omega_0^2 \cos^2 \alpha + \omega_t^2 \sin^2 \alpha + \frac{\kappa}{\mu} \left[\left(\frac{m_0}{m_t} \right)^{1/4} \sin \alpha - \left(\frac{m_t}{m_0} \right)^{1/4} \cos \alpha \right]^2 \right\}^{1/2}$$

And

$$\omega_{-} = \left\{ \omega_0^2 \sin^2 \alpha + \omega_t^2 \cos^2 \alpha + \frac{\kappa}{\mu} \left[\left(\frac{m_0}{m_t} \right)^{1/4} \cos \alpha - \left(\frac{m_t}{m_0} \right)^{1/4} \sin \alpha \right]^2 \right\}^{1/2}$$

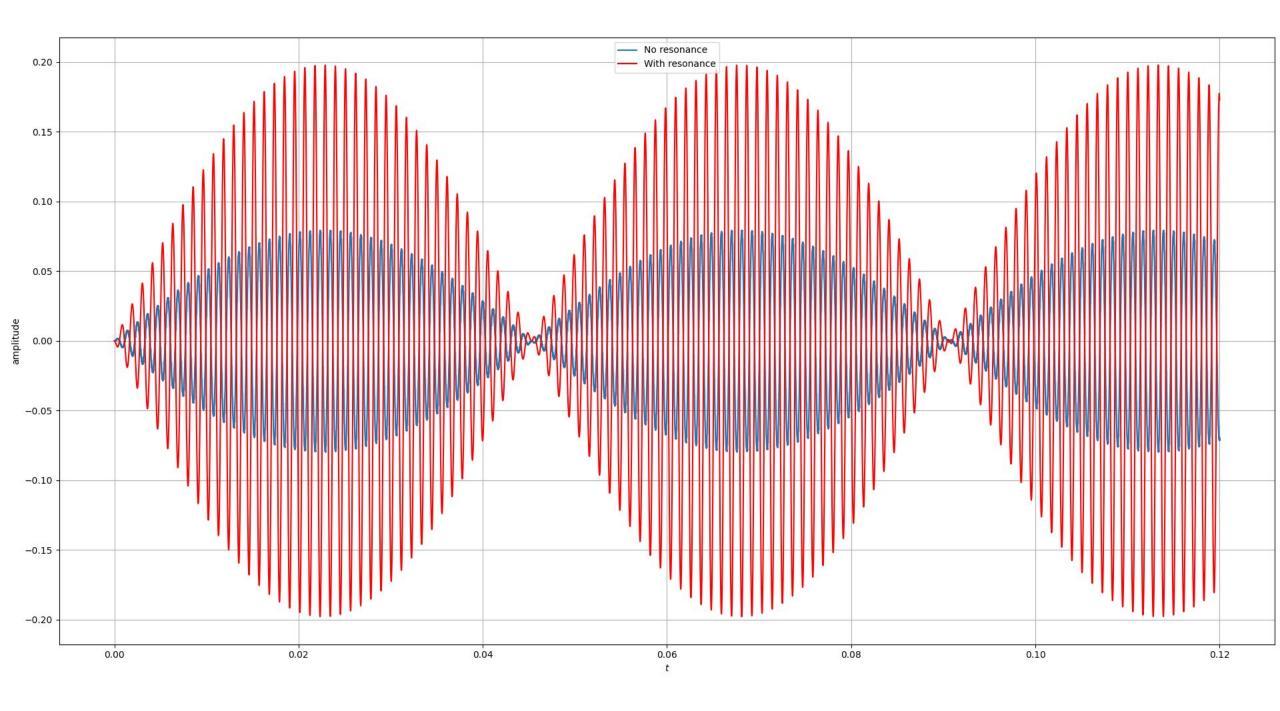
$$\hat{H} = \hat{a}_+^{\dagger} \hat{a}_+ + \hat{a}_-^{\dagger} \hat{a}_-$$

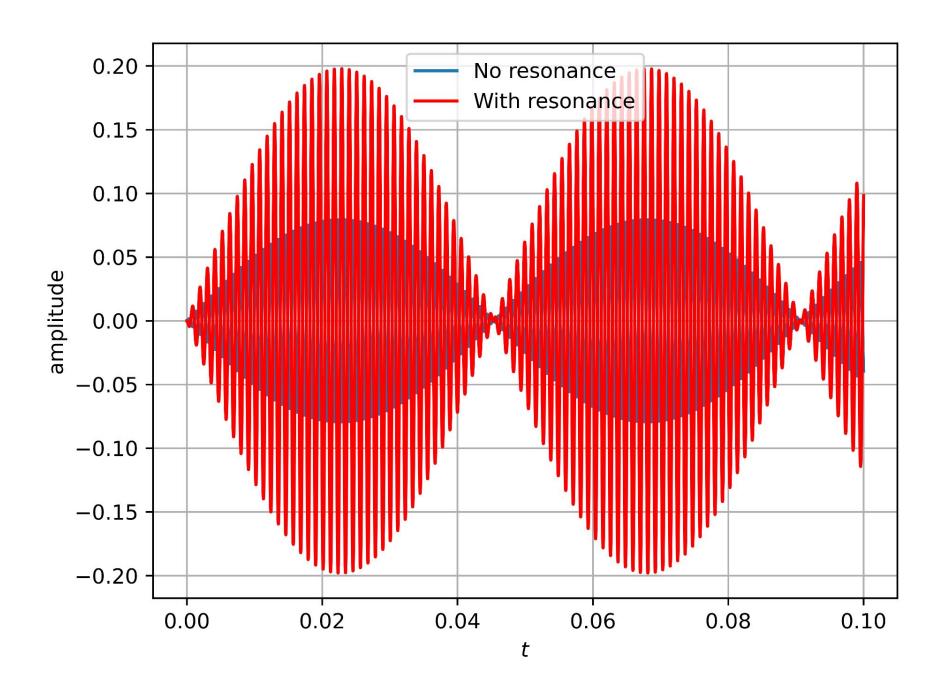
$$\hat{H} = \hat{N}_+ + \hat{N}_-$$

$$\Psi(\xi_0, \xi_t, 0) = \left(\frac{m_0 m_t \omega_0 \omega_t}{\pi^2 \hbar^2}\right)^{1/4} exp\left(-\frac{m_0 \omega_0 \xi_0^2 + m_t \omega_t (\xi_t - A)^2}{2 \hbar}\right)$$

$$\langle y_0 \rangle = \left(\frac{m_t}{m_0}\right)^{\frac{1}{2}} A \sin \alpha \cos \alpha [\cos \omega_+ t - \cos \omega_- t]$$

.





Main conclusion

The solution for the averaged position is the same as the classical one.

Next step is to study the measurement process during the detection.

On Thursday poster with a complementary work will be presented.

Thank you for your kind attention.