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# Electromagnetic probes of QGP

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In collaboration with **Taesoo Song**, Wolfgang Cassing, Pierre Moreau



*8th International Workshop on Astronomy and Relativistic  
Astrophysics 34th Winter Workshop on Nuclear Dynamics  
IWARA2018*

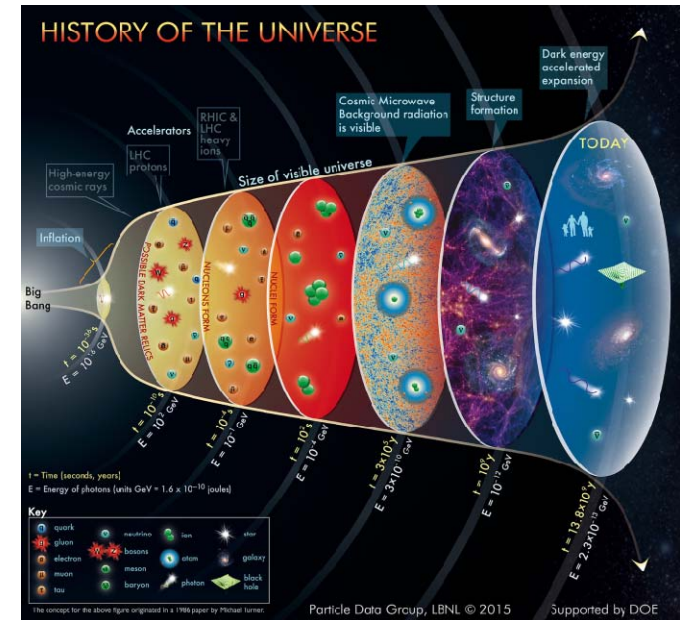
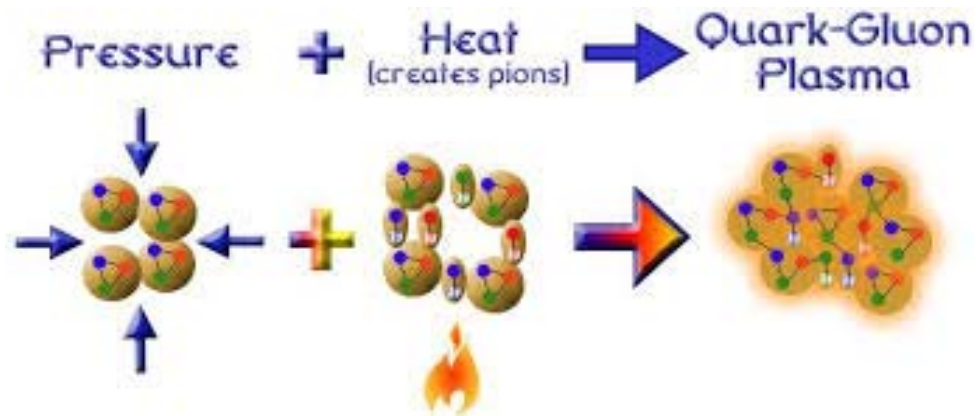
*Ollantaytambo, Sacred Valley, Peru,  
9 –15 September, 2018*



# Experiment: Heavy-ion collisions

## □ Heavy-ion collision experiment

→ ‚re-creation‘ of the Big Bang conditions in laboratory:  
matter at high **pressure** and **temperature**

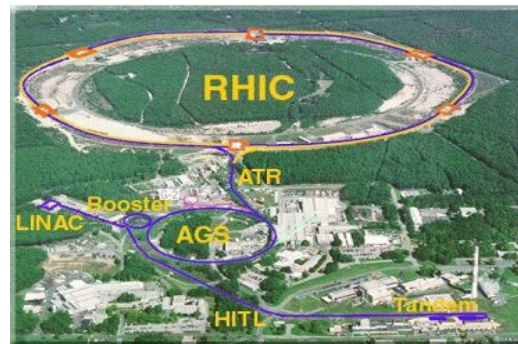


## □ Heavy-ion accelerators:

Large Hadron Collider -  
LHC (CERN):  
Pb+Pb up to 574 A TeV



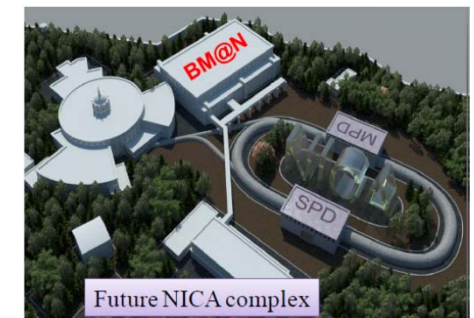
Relativistic-Heavy-Ion-Collider -  
RHIC (Brookhaven):  
Au+Au up to 21.3 A TeV



Facility for Antiproton and Ion  
Research – FAIR (Darmstadt)  
(Under construction)  
Au+Au up to 10 (30) A GeV



Nuclotron-based Ion Collider  
Facility – NICA (Dubna)  
(Under construction)  
Au+Au up to 60 A GeV





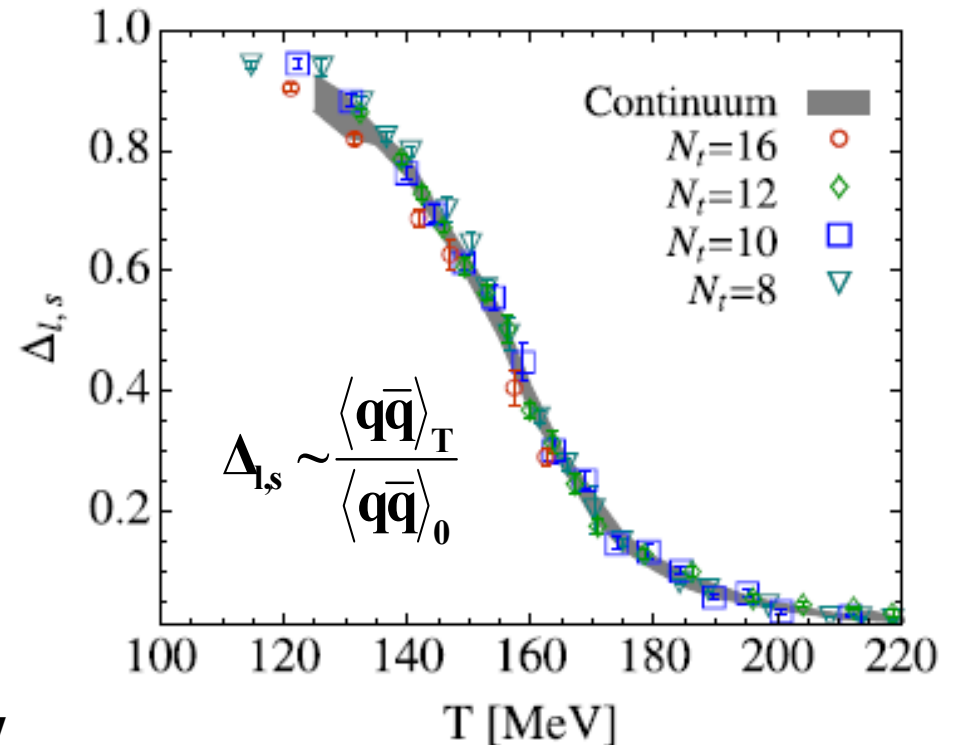
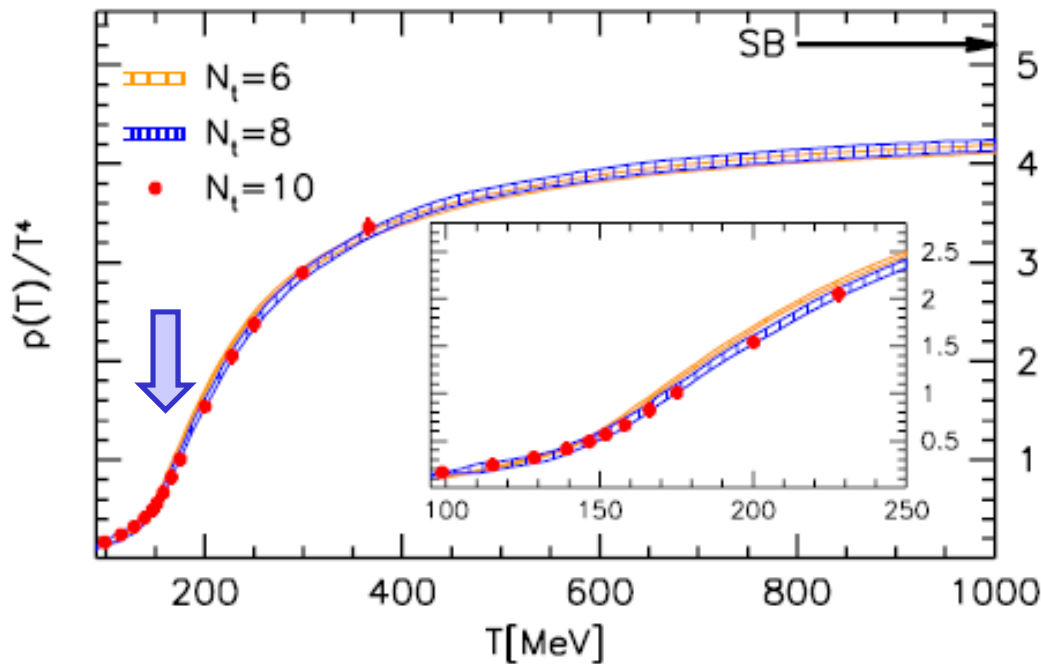
# Theory: Information from lattice QCD

I. deconfinement phase transition with increasing temperature



II. chiral symmetry restoration with increasing temperature

IQCD BMW collaboration:  $\mu_q=0$



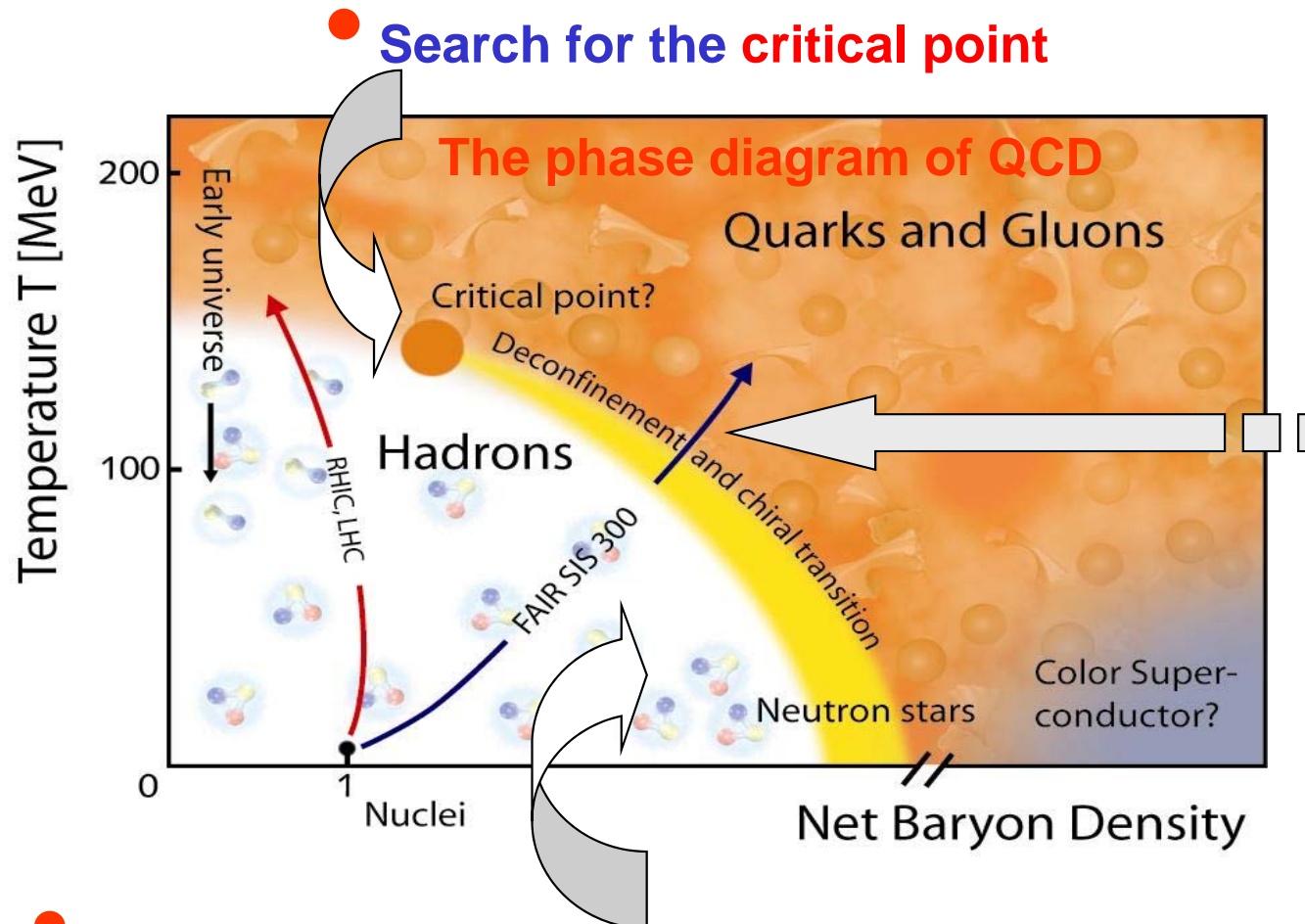
□ **Crossover:** hadron gas → QGP,  $T_C=156$  MeV

□ **Scalar quark condensate  $\langle q\bar{q} \rangle$**  is viewed as an **order parameter** for the restoration of chiral symmetry:

$$\langle \bar{q}q \rangle = \begin{cases} \neq 0 & \text{chiral non-symmetric phase;} \\ = 0 & \text{chiral symmetric phase.} \end{cases}$$

→ both transitions occur at about the same temperature  $T_C$  for low chemical potentials

# The ,holy grail' of heavy-ion physics:



- Study of the **phase transition** from hadronic to partonic matter – **Quark-Gluon-Plasma**



- Study of the **in-medium** properties of hadrons at high baryon density and temperature

# Electromagnetic probes: photons and dileptons

Feinberg (76), Shuryak (78)

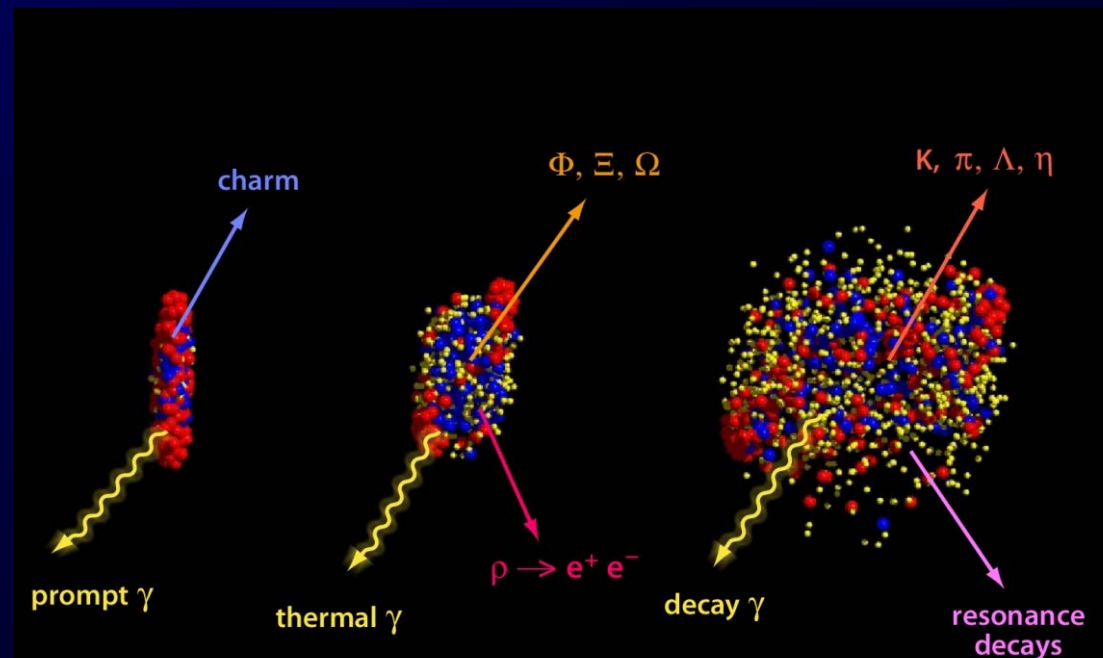
## ■ Advantages:

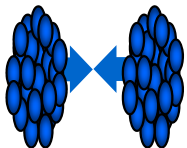
- ✓ dileptons and real photons are emitted from different stages of the reaction and not effected by final-state interactions
- ✓ provide undistorted information about their production channels
- ✓ promising signal of QGP – ,thermal' photons and dileptons

→ Requires **theoretical models** which describe the **dynamics** of heavy-ion collisions during the whole time evolution!

## □ Disadvantages:

- low emission rate
- production from hadronic corona
- many production sources which cannot be individually disentangled by experimental data





# Basic models for heavy-ion collisions

- Statistical models:

**basic assumption:** system is described by a (grand) canonical ensemble of non-interacting fermions and bosons in **thermal and chemical equilibrium**  
= **thermal hadron gas at freeze-out** with common  $T$  and  $\mu_B$

[ - : no dynamical information]

- Hydrodynamical models:

**basic assumption:** conservation laws + equation of state (EoS);  
assumption of **local thermal and chemical equilibrium**

- Interactions are ,hidden‘ in properties of the **fluid** described by **transport coefficients** (shear and bulk viscosity  $\eta$ ,  $\zeta$ , ..), which is **‘input‘** for the hydro models

[ - : simplified dynamics]

- Microscopic transport models:

**based on transport theory of relativistic quantum many-body systems**

- **Explicitly account for the interactions of all degrees of freedom** (hadrons and partons)  
in terms of cross sections and potentials

- Provide a unique dynamical description of **strongly interaction matter**

**in- and out-of equilibrium:**

- **In-equilibrium:** transport coefficients are calculated in a box – controlled by IQCD

- **Nonequilibrium dynamics** – controlled by HIC

**Actual solutions:** Monte Carlo simulations

[+ : full dynamics | - : very complicated]



# Parton-Hadron-String-Dynamics (PHSD)

PHSD is a non-equilibrium transport approach with

- explicit **phase transition** from hadronic to partonic degrees of freedom
- **IQCD EoS** for the partonic phase (‘crossover’ at low  $\mu_q$ )
- explicit **parton-parton interactions** - between quarks and gluons
- dynamical **hadronization**

□ **QGP phase is described by the Dynamical QuasiParticle Model (DQPM)** matched to reproduce lattice QCD

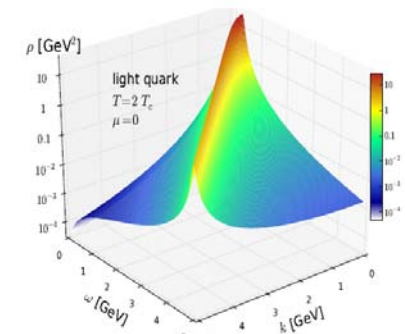
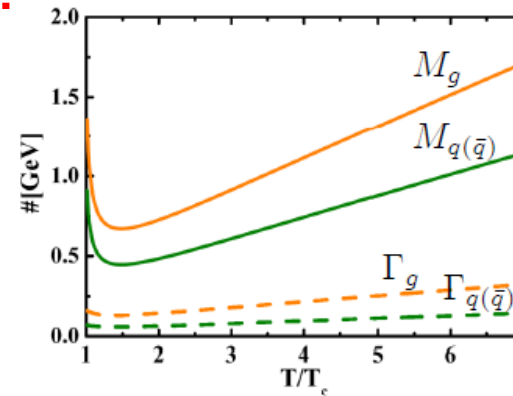
A. Peshier, W. Cassing, PRL 94 (2005) 172301;  
W. Cassing, NPA 791 (2007) 365; NPA 793 (2007)

▪ **strongly interacting quasi-particles:** massive quarks and gluons ( $g, q, q_{\text{bar}}$ ) with sizeable collisional widths in a self-generated **mean-field potential**

▪ **Spectral functions:**

$$\rho_i(\omega, T) = \frac{4\omega\Gamma_i(T)}{\left(\omega^2 - \vec{p}^2 - M_i^2(T)\right)^2 + 4\omega^2\Gamma_i^2(T)}$$

$(i = q, \bar{q}, g)$



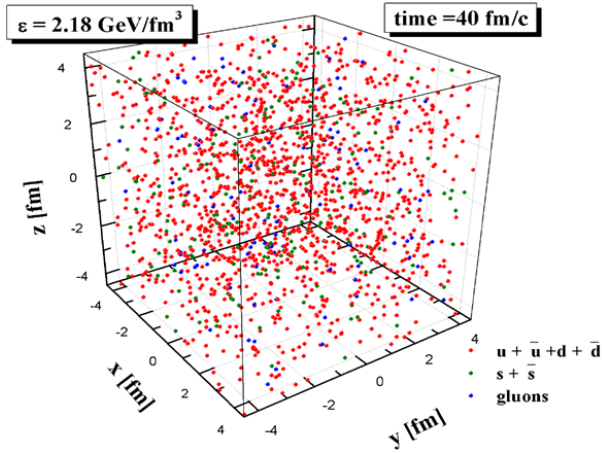
□ **Transport theory:** generalized off-shell transport equations based on the 1st order gradient expansion of Kadanoff-Baym equations (**applicable for strongly interacting systems!**)



# QGP in equilibrium: Transport properties at finite $(T, \mu_q)$ : $\eta/s$

Infinite hot/dense matter =

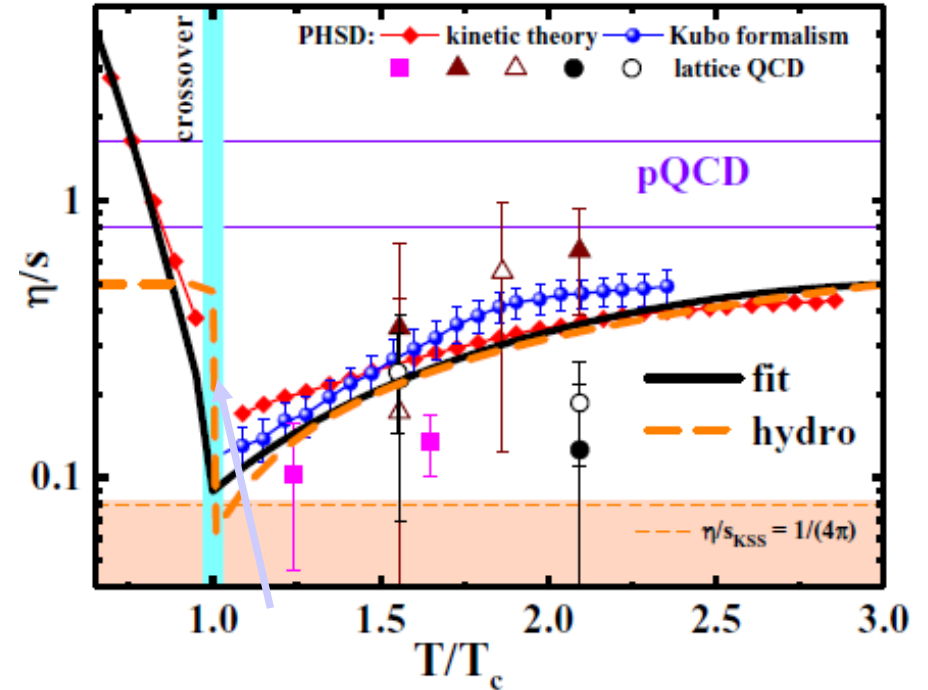
PHSD in a box:



## Shear viscosity $\eta/s$ at finite T

PHSD: V. Ozvenchuk et al., PRC 87 (2013) 064903

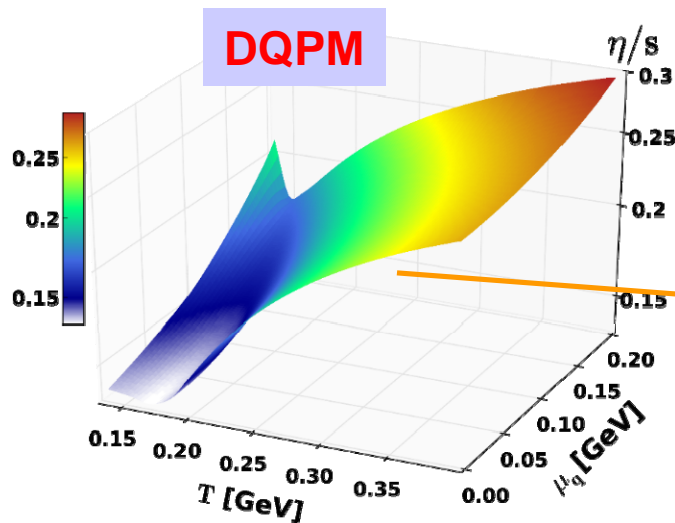
Hydro: Bayesian analysis, S. Bass et al., 1704.07671



## Shear viscosity $\eta/s$ at finite $(T, \mu_q)$

IQCD:

$$\frac{T_c(\mu_q)}{T_c(\mu_q = 0)} = \sqrt{1 - \alpha \mu_q^2} \approx 1 - \alpha/2 \mu_q^2 + \dots$$



**QGP in PHSD = strongly-interacting liquid-like system**

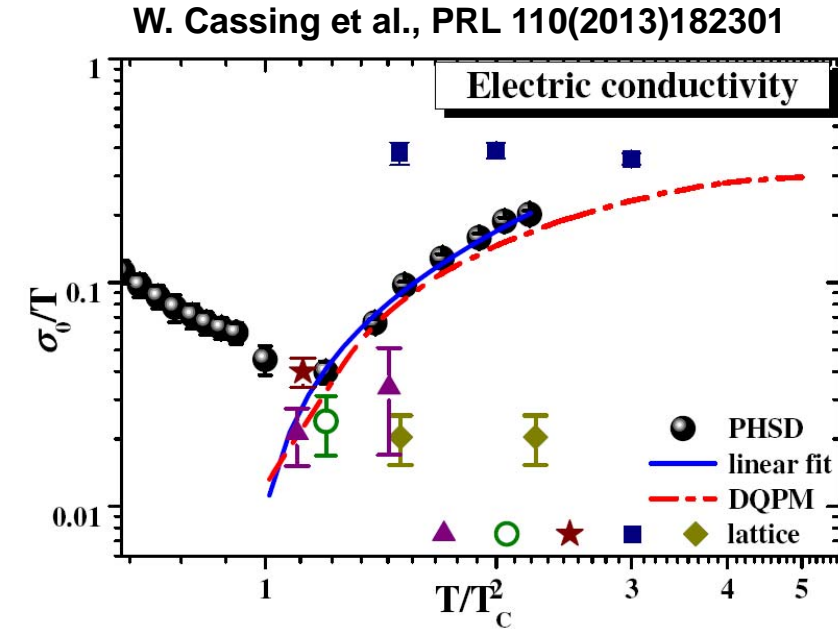
$\eta/s$ :  $\mu_q=0 \rightarrow$  finite  $\mu_q$ : smooth increase as a function of  $(T, \mu_q)$



# Transport properties at finite $(T, \mu_q)$ : $\sigma_e/T$

PHSD in a box:

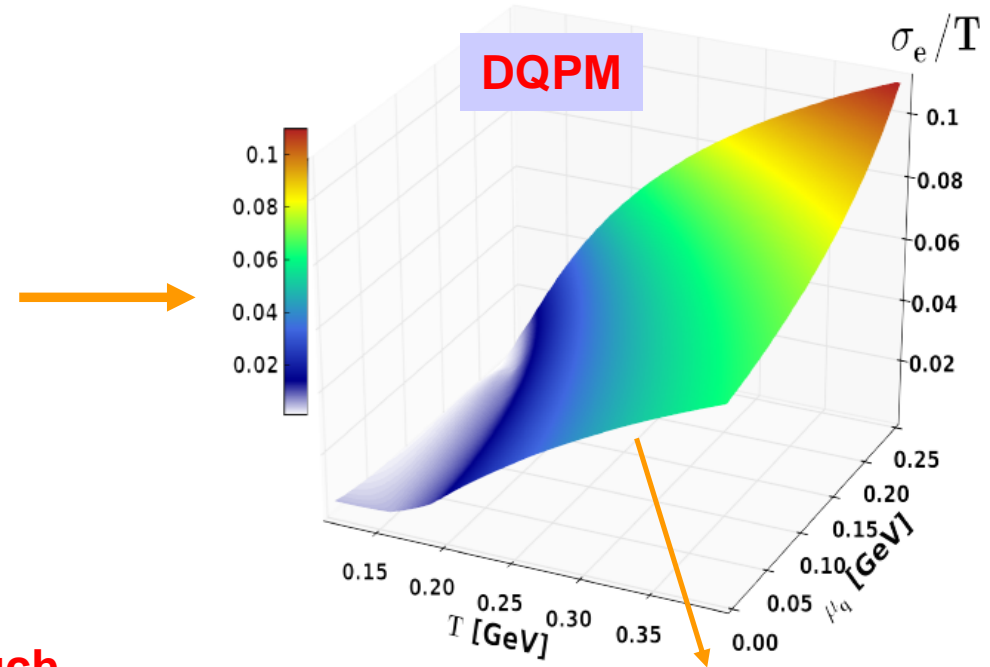
Electric conductivity  $\sigma_e/T$  at finite  $T$



the QCD matter even at  $T \sim T_c$  is a much better electric conductor than Cu or Ag (at room temperature) by a factor of 500 !

Electric conductivity  $\sigma_e/T$  at finite  $(T, \mu_q)$

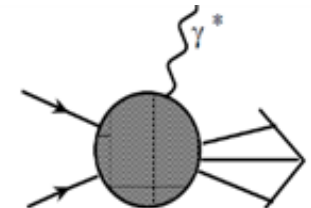
H. Berrehrah et al. , PRC93 (2016) 044914



$\sigma_e/T$  :  $\mu_q=0 \rightarrow$  finite  $\mu_q$ : smooth increase as a function of  $(T, \mu_q)$

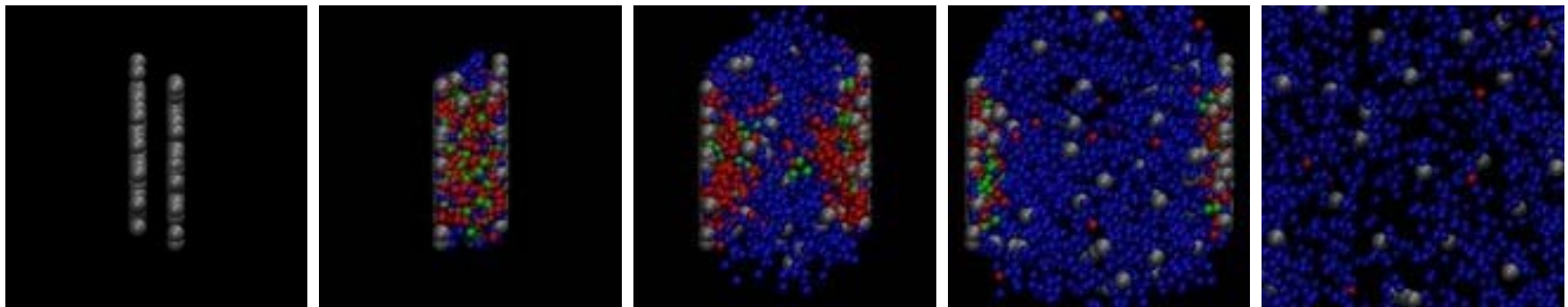
Photon emission: rates at  $q_0 \rightarrow 0$  are related to electric conductivity  $\sigma_0$

$$q_0 \left. \frac{dR}{d^4x d^3q} \right|_{q_0 \rightarrow 0} = \frac{T}{4\pi^3} \sigma_0$$



$\sigma_0 \rightarrow$  Probe of electromagnetic properties of the QGP

# „Bulk“ properties in Au+Au collisions

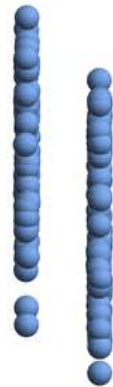







# Au+Au at 200 A GeV, b=2.2 fm

t = 0.1 fm/c



**Au + Au  $\sqrt{s_{NN}} = 200$  GeV**  
**b = 2.2 fm - Section view**



-  Baryons (394)
-  Antibaryons ( 0)
-  Mesons ( 0)
-  Quarks ( 0)
-  Gluons ( 0)






# Au+Au at 200 A GeV, b=2.2 fm

t = 1.63549 fm/c



**Au + Au  $\sqrt{s_{NN}} = 200$  GeV**  
**b = 2.2 fm - Section view**



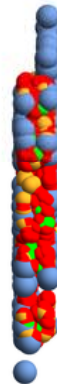
-  Baryons (394)
-  Antibaryons ( 0)
-  Mesons (1598)
-  Quarks (4383)
-  Gluons (344)






# Au+Au at 200 A GeV, b=2.2 fm

t = 2.06543 fm/c



**Au + Au  $\sqrt{s_{NN}} = 200$  GeV**  
**b = 2.2 fm – Section view**



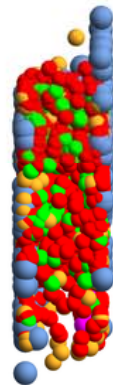
-  Baryons (396)
-  Antibaryons ( 2)
-  Mesons (1136)
-  Quarks (5066)
-  Gluons (516)






# Au+Au at 200 A GeV, b=2.2 fm

t = 3.20258 fm/c



**Au + Au  $\sqrt{s_{NN}} = 200$  GeV**  
**b = 2.2 fm – Section view**



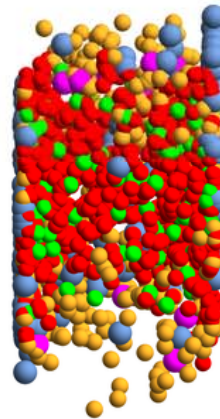
-  Baryons (413)
-  Antibaryons ( 13)
-  Mesons (1080)
-  Quarks (4708)
-  Gluons (761)






# Au+Au at 200 A GeV, b=2.2 fm

t = 5.56921 fm/c



**Au + Au  $\sqrt{s_{NN}} = 200$  GeV**  
**b = 2.2 fm – Section view**



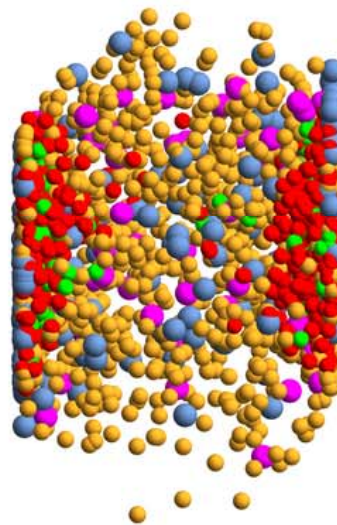
-  Baryons (472)
-  Antibaryons ( 70)
-  Mesons (1724)
-  Quarks (3843)
-  Gluons (652)





# Au+Au at 200 A GeV, b=2.2 fm

t = 8.06922 fm/c



**Au + Au  $\sqrt{s_{NN}} = 200$  GeV**  
**b = 2.2 fm – Section view**



-  Baryons (559)
-  Antibaryons (139)
-  Mesons (2686)
-  Quarks (2628)
-  Gluons (442)

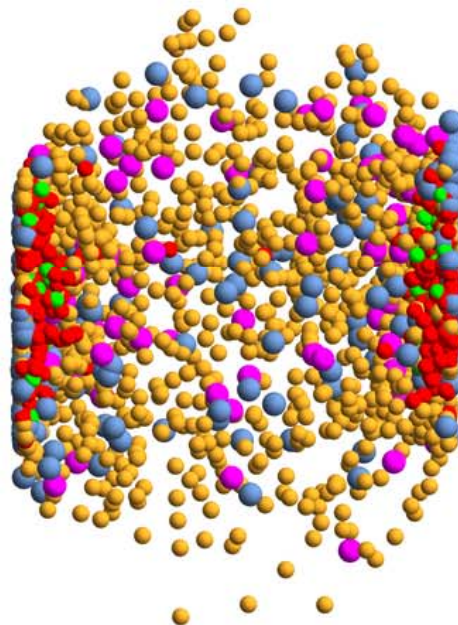







# Au+Au at 200 A GeV, b=2.2 fm

t = 10.5692 fm/c



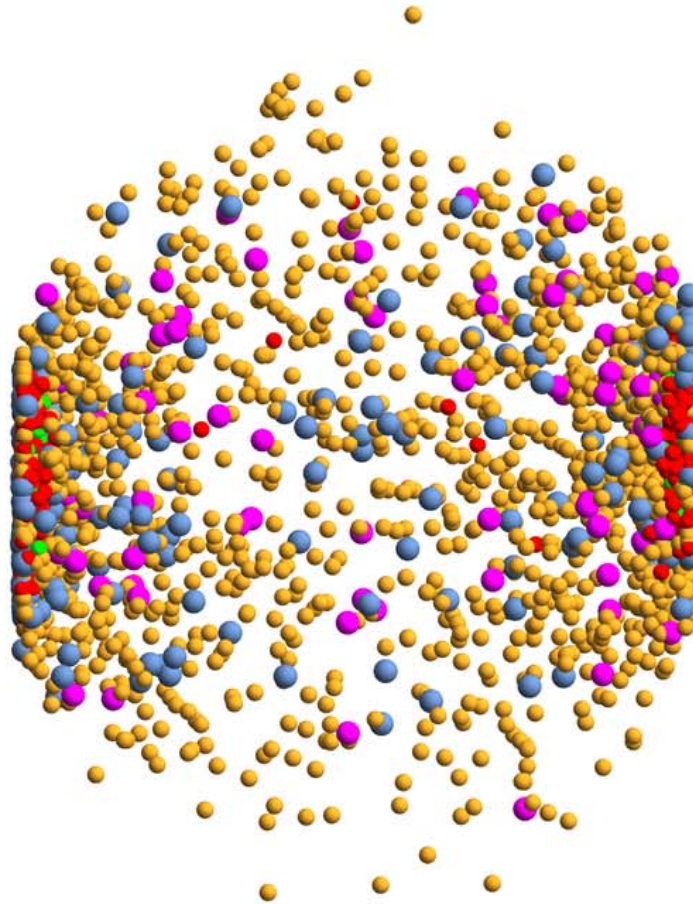
**Au + Au  $\sqrt{s_{NN}} = 200$  GeV**  
**b = 2.2 fm – Section view**







-  Baryons (604)
-  Antibaryons (187)
-  Mesons (3169)
-  Quarks (2076)
-  Gluons (319)

# Au+Au at 200 A GeV, b=2.2 fm

t = 15.5692 fm/c

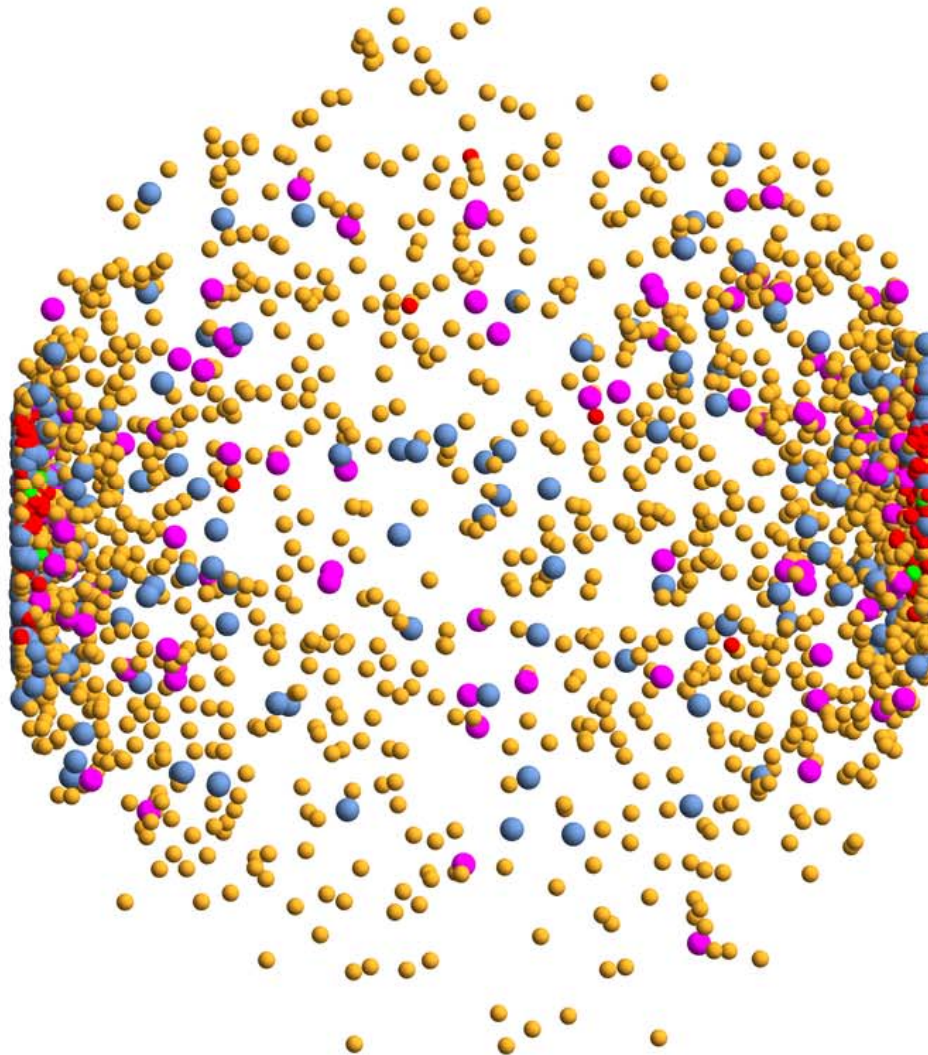


**Au + Au  $\sqrt{s_{NN}} = 200$  GeV**  
**b = 2.2 fm – Section view**






-  Baryons (662)
-  Antibaryons (229)
-  Mesons (3661)
-  Quarks (1499)
-  Gluons (175)

# Au+Au at 200 A GeV, b=2.2 fm

t = 20.5692 fm/c



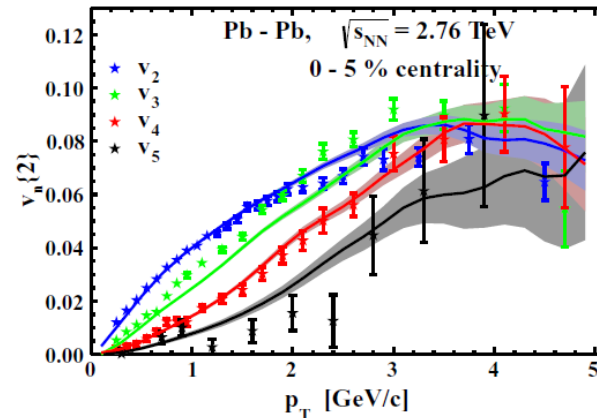
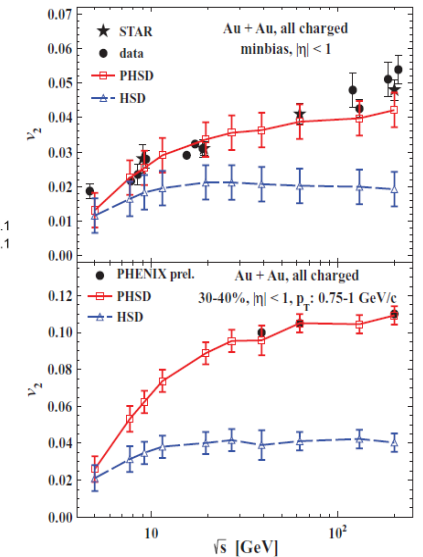
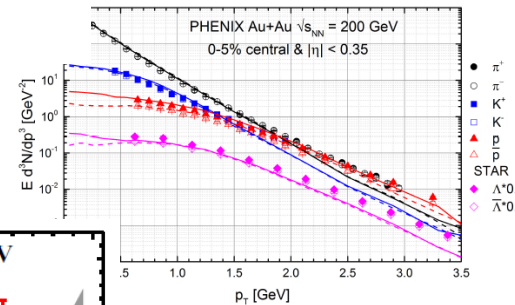
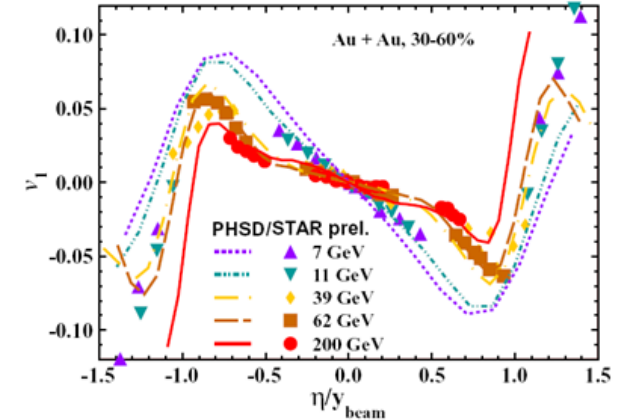
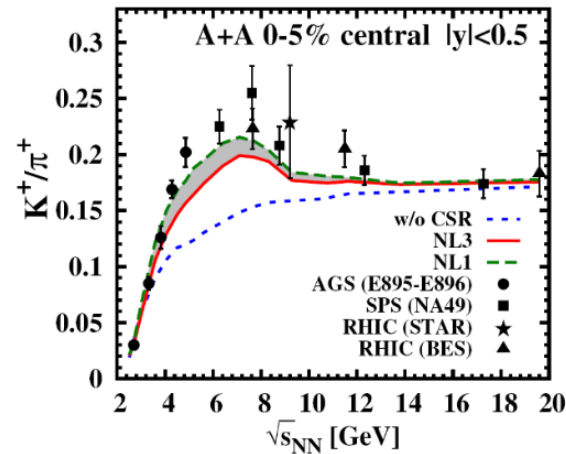
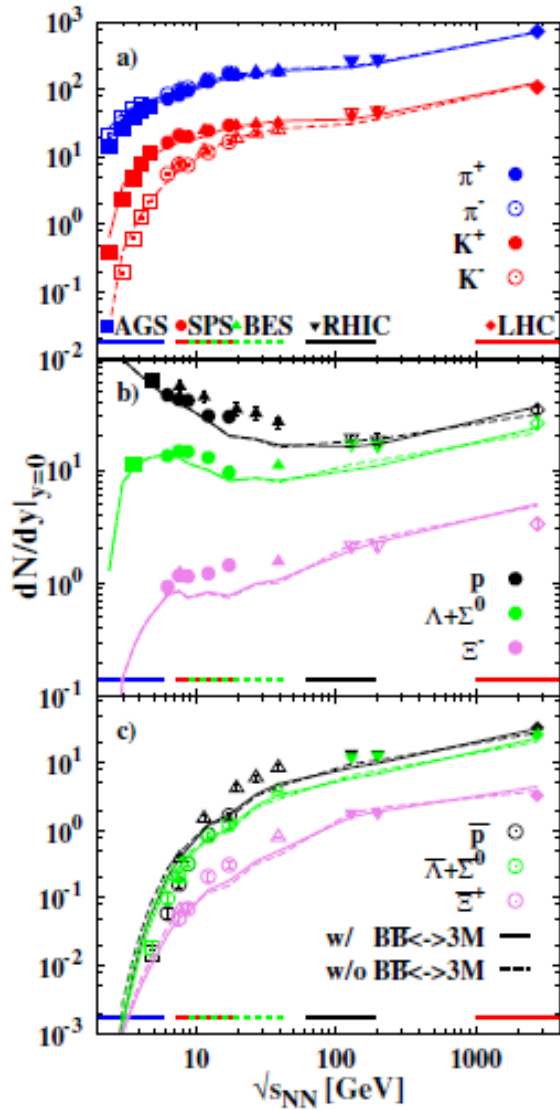
**Au + Au  $\sqrt{s_{NN}} = 200$  GeV**  
**b = 2.2 fm – Section view**

-  Baryons (692)
-  Antibaryons (266)
-  Mesons (4022)
-  Quarks (1184)
-  Gluons ( 90)



# Non-equilibrium dynamics: description of A+A with PHSD

## PHSD: highlights



PRC 85 (2012) 011902; JPG42 (2015) 055106

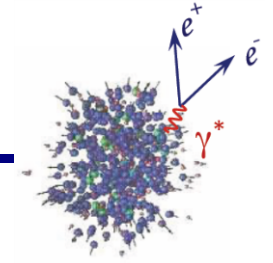
arXiv:1801.07557

PHSD provides a good description of 'bulk' observables ( $y$ -,  $p_T$ -distributions, flow coefficients  $v_n$ , ...) from SIS to LHC

# Dileptons as a probe of the QGP and in-medium effects



# Dilepton sources



from the QGP via partonic (q,qbar, g) interactions:



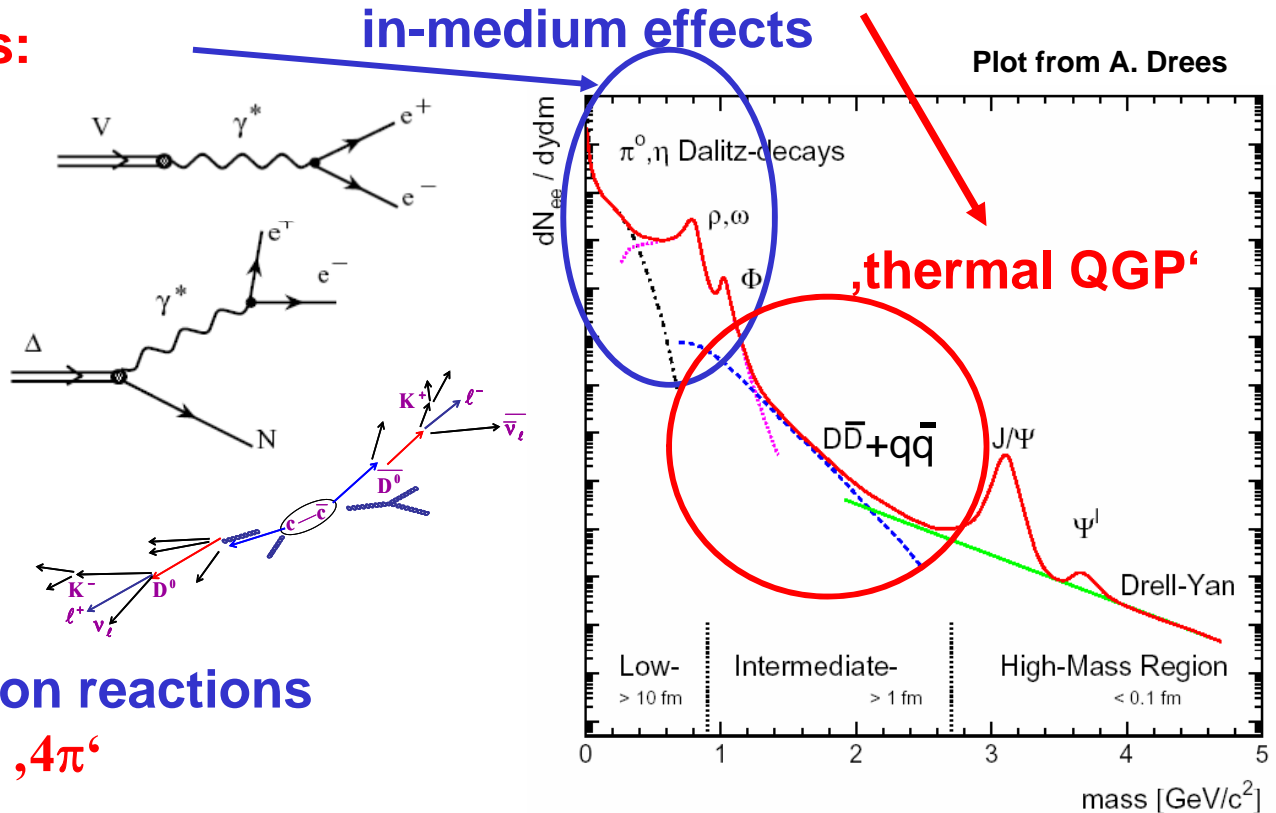
from hadronic sources:

- direct decay of vector mesons ( $\rho, \omega, \phi, J/\Psi, \Psi'$ )

- Dalitz decay of mesons and baryons ( $\pi^0, \eta, \Delta, \dots$ )

- correlated D+Dbar pairs

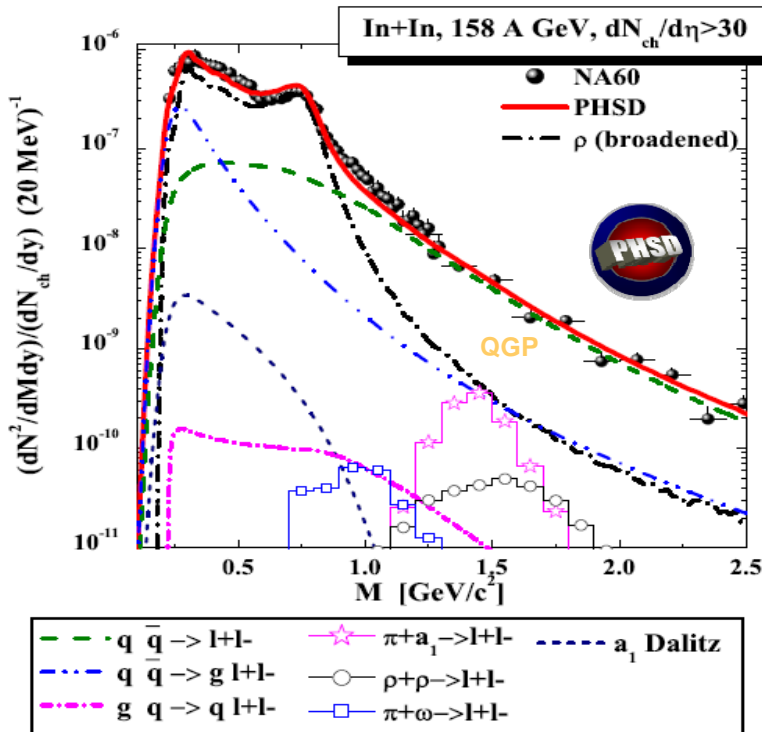
- radiation from multi-meson reactions ( $\pi+\pi, \pi+\rho, \pi+\omega, \rho+\rho, \pi+a_1$ ) -  $4\pi'$



**! Advantage of dileptons:**  
 additional „degree of freedom“ ( $M$ ) allows to disentangle various sources

# Lessons from SPS: NA60

## Dilepton invariant mass spectra:

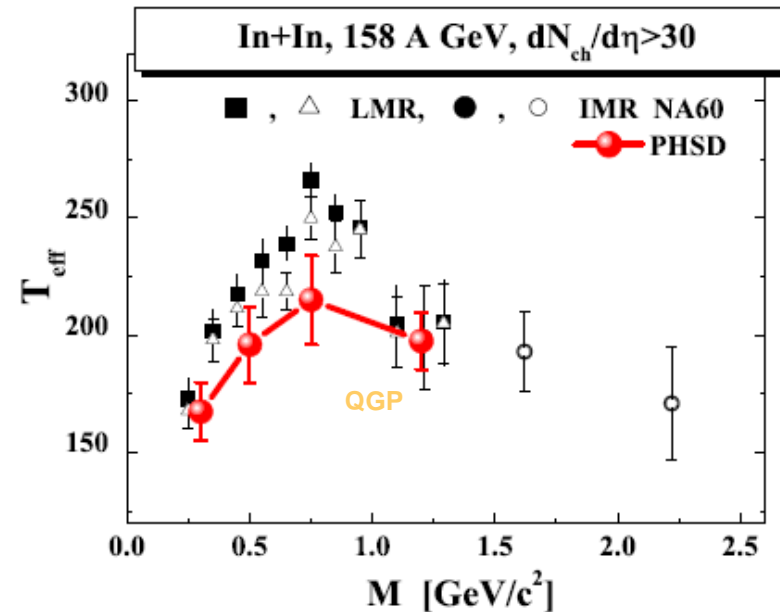


NA60: Eur. Phys. J. C 59 (2009) 607

PHSD: Linnyk et al, PRC 84 (2011) 054917

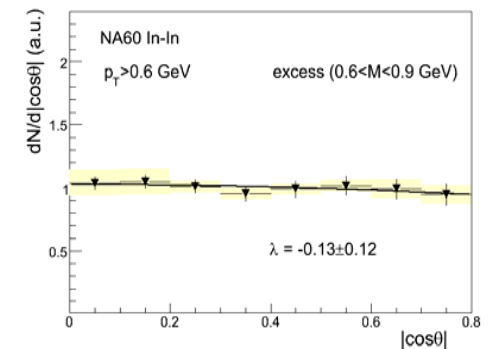
## Inverse slope parameter $T_{\text{eff}}$ :

spectrum from QGP is softer than from hadronic phase since the QGP emission occurs dominantly before the collective radial flow has developed



## Message from SPS: (based on NA60 and CERES data)

- 1) Low mass spectra - evidence for the **in-medium broadening of  $\rho$ -mesons**
- 2) Intermediate mass spectra above 1 GeV - dominated by **partonic radiation**
- 3) The rise and fall of  $T_{\text{eff}}$  – evidence for the thermal **QGP radiation**
- 4) **Isotropic angular distribution** – indication for a **thermal origin of dimuons**

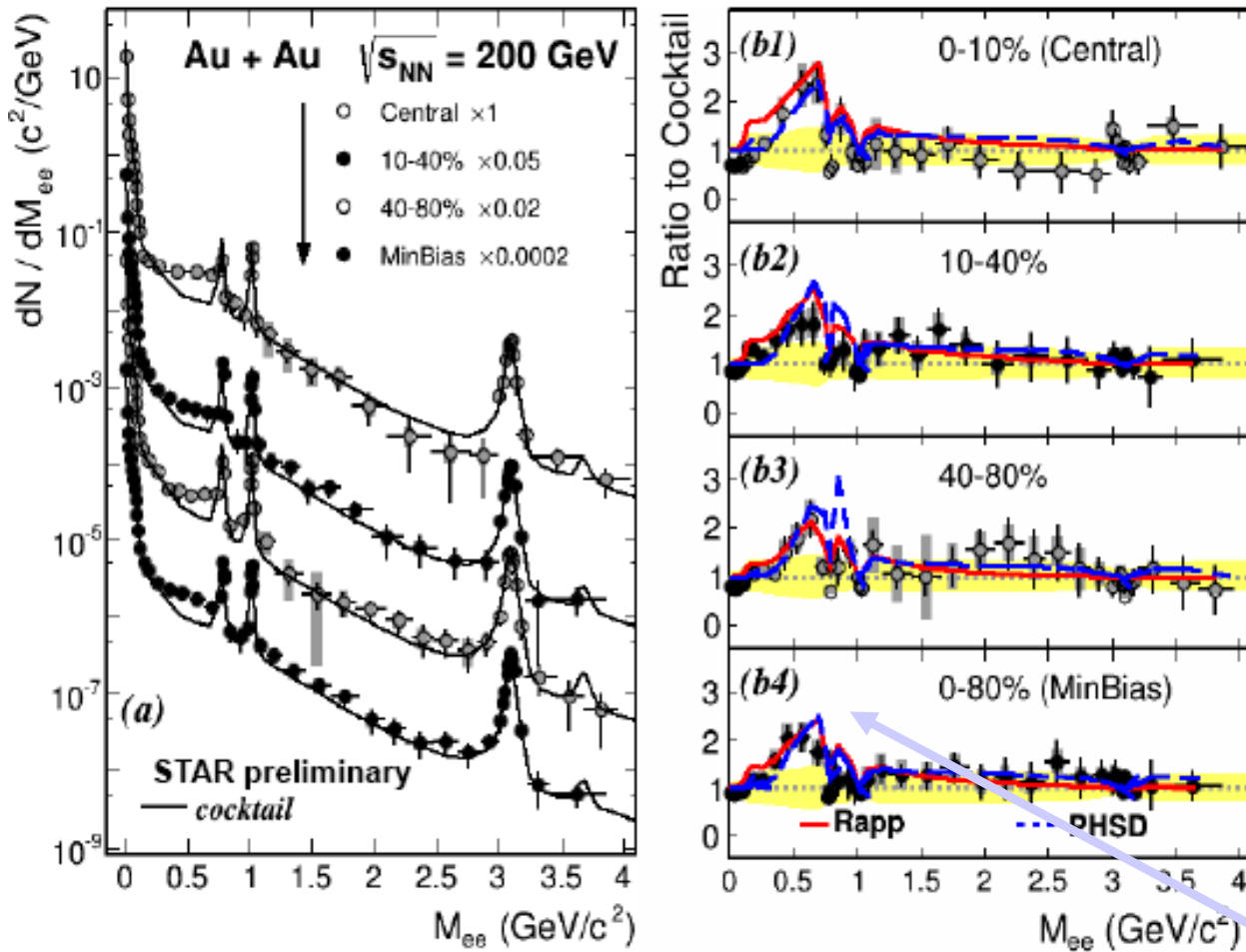


PRL 102 (2009) 222301

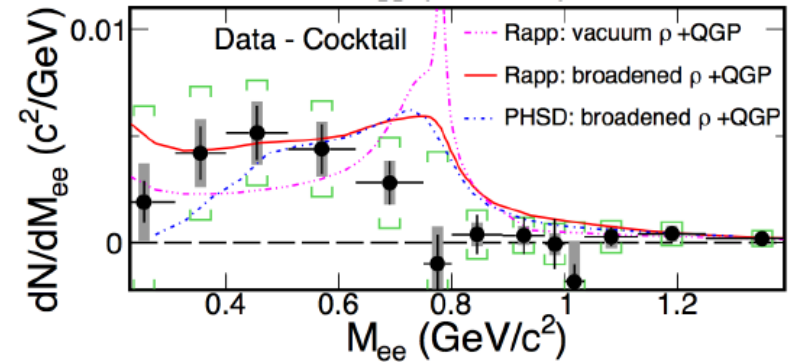
# Dileptons at RHIC: STAR data vs model predictions

PRC 92 (2015) 024912

## Centrality dependence of dilepton yield



## Excess in low mass region, min. bias



Models:

- Fireball model – R. Rapp
- PHSD

Low masses:

collisional broadening of  $\rho$

Intermediate masses:

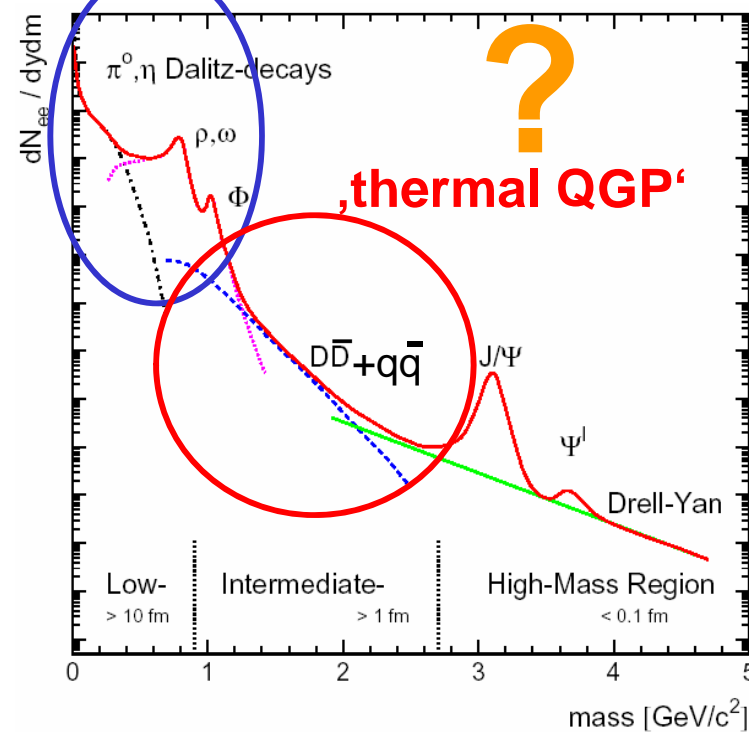
QGP dominant

**Message:** STAR data are described by models within a collisional broadening scenario for the vector meson spectral function + QGP



# What is the best energy range to observe thermal dileptons from QGP ?

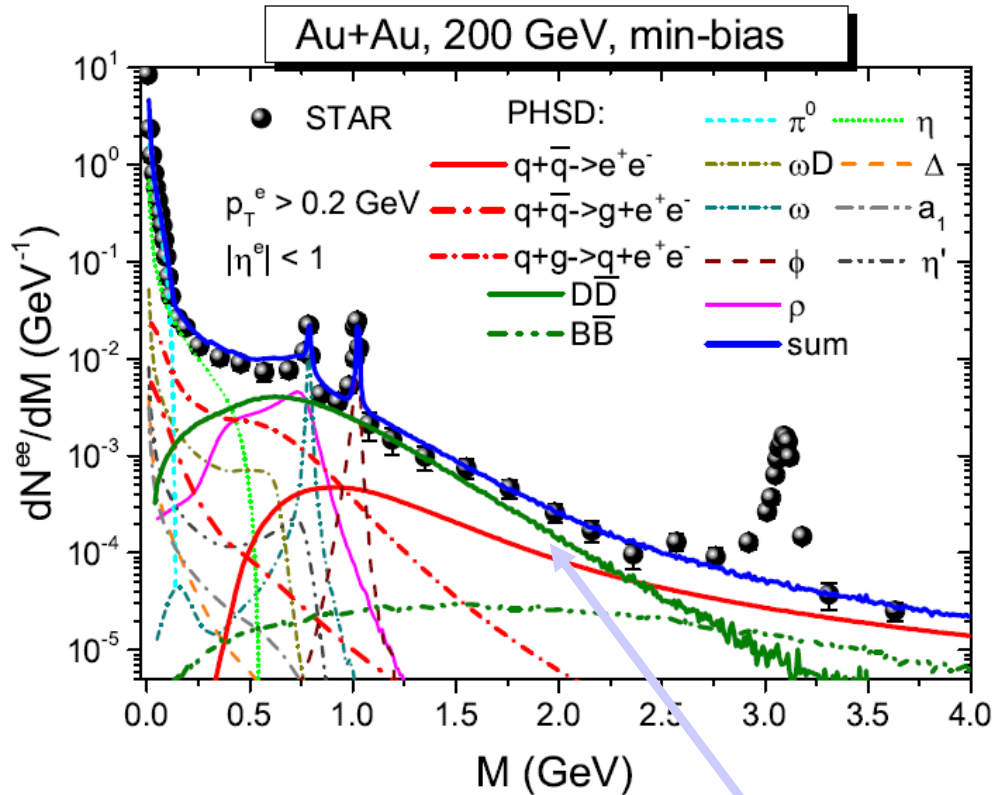
in-medium effects



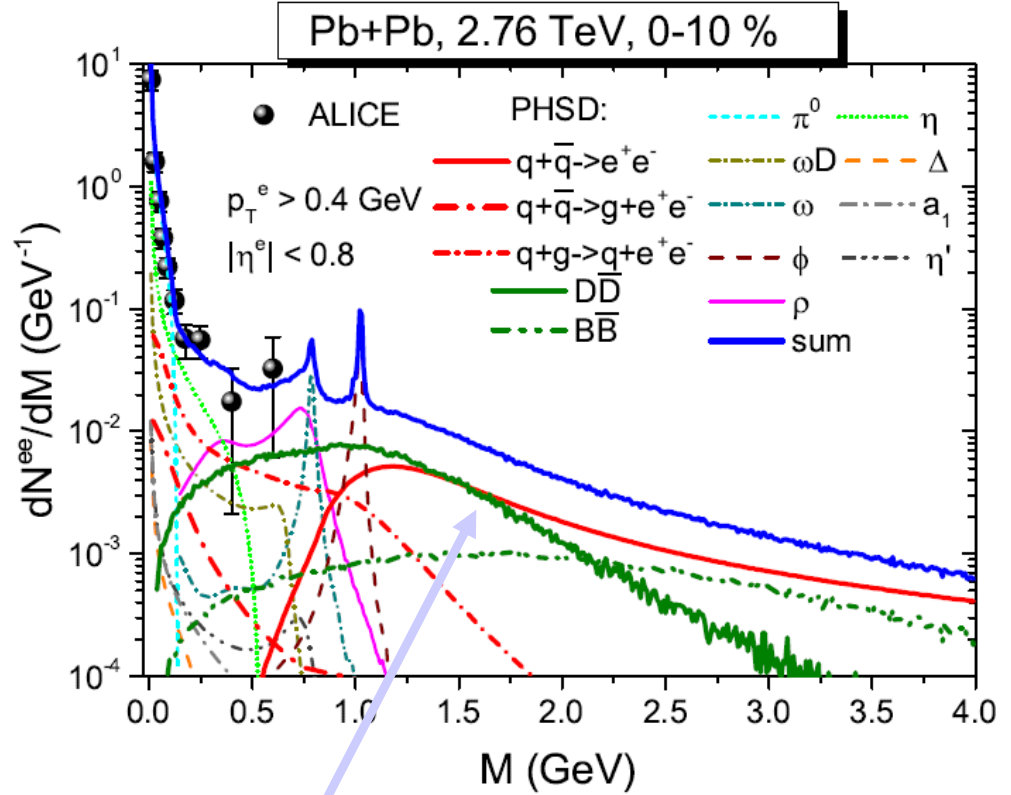


# Dileptons at RHIC and LHC

**RHIC**



**LHC**



**Message:**

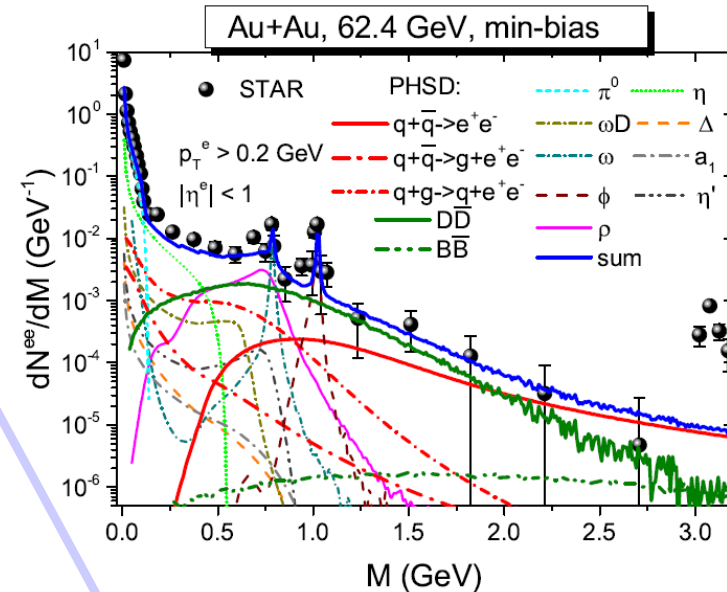
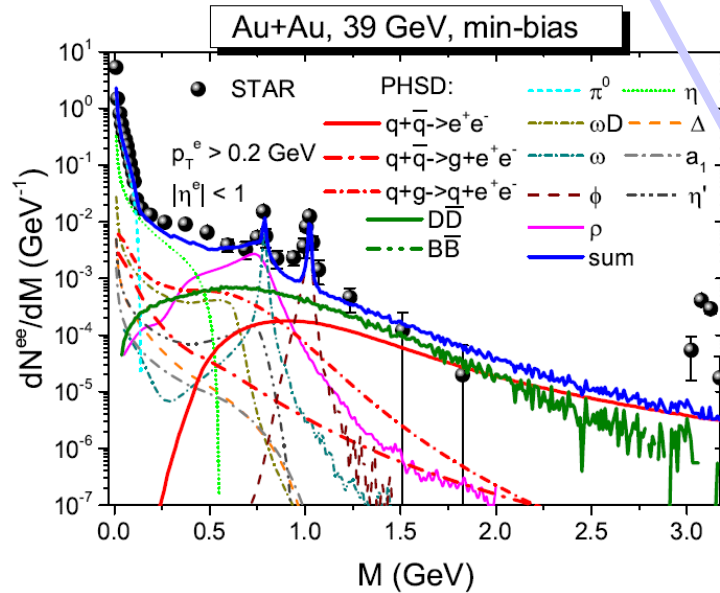
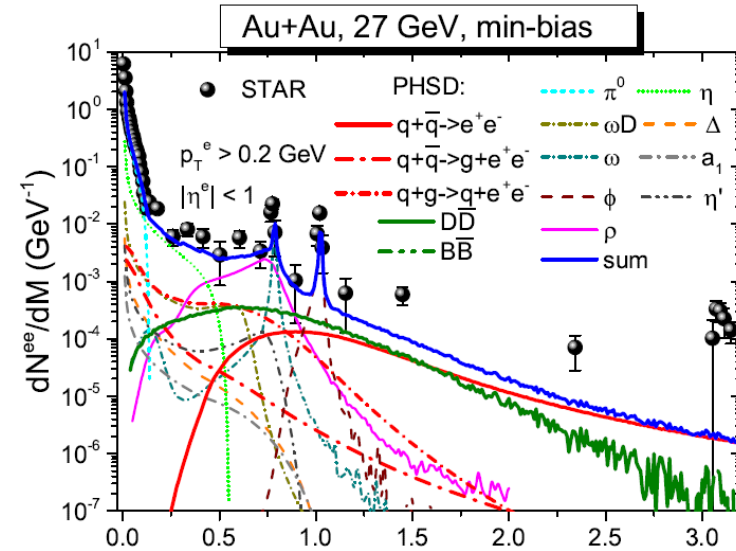
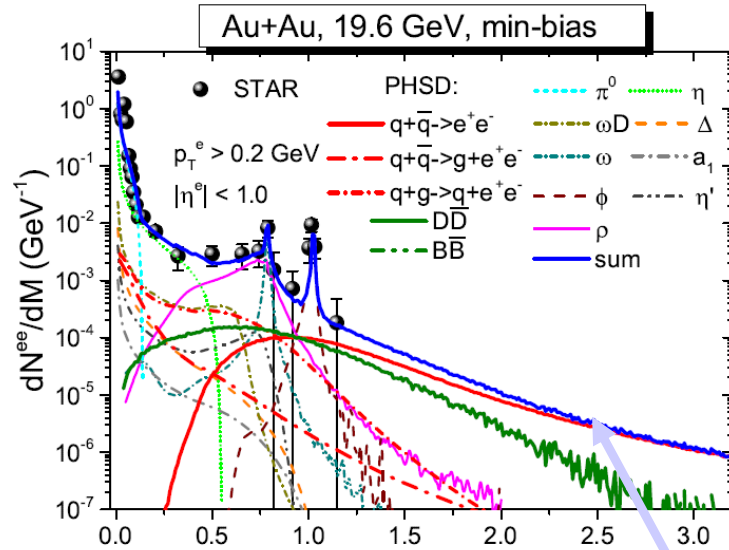
STAR data at 200 GeV and the ALICE data at 2.76 TeV are described by PHSD within

- 1) a **collisional broadening** scenario for the **vector meson** spectral functions + **QGP** + **correlated charm**
- 2) **Charm contribution** is dominant for  $1.2 < M < 2.5$  GeV



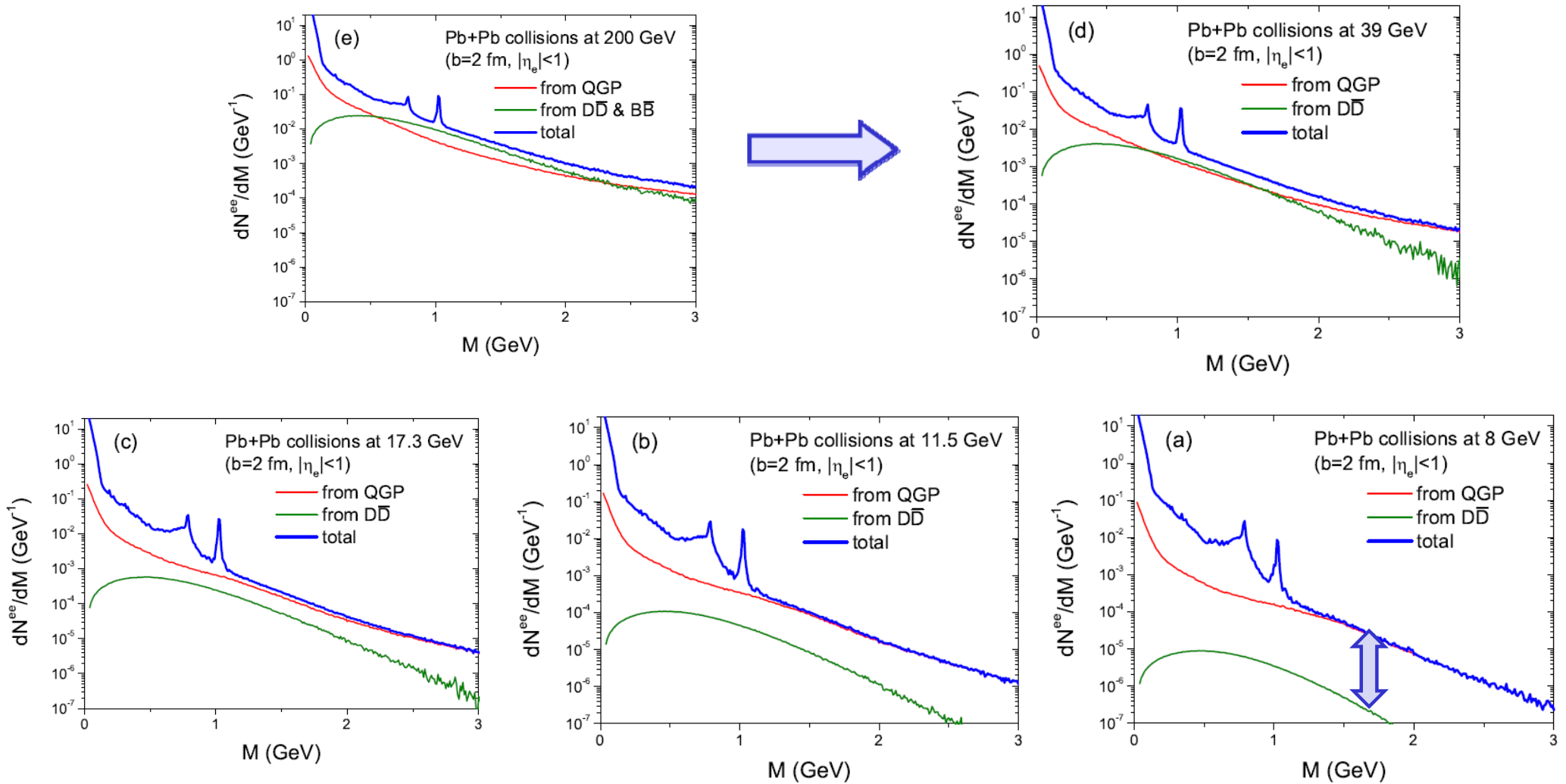
# Dileptons from RHIC BES: STAR

T. Song, W.Cassing, P.Moreau and E.Bratkovskaya, Phys. Rev. C 97 (2018) 064907



**QGP and charm are dominant contributions for intermediate masses at BES RHIC**  
**→ measurements of charm at BES RHIC are needed to control charm production !**

# Dileptons at FAIR/NICA energies: predictions



Relative contribution of **QGP** versus **charm** increases with **decreasing energy**!

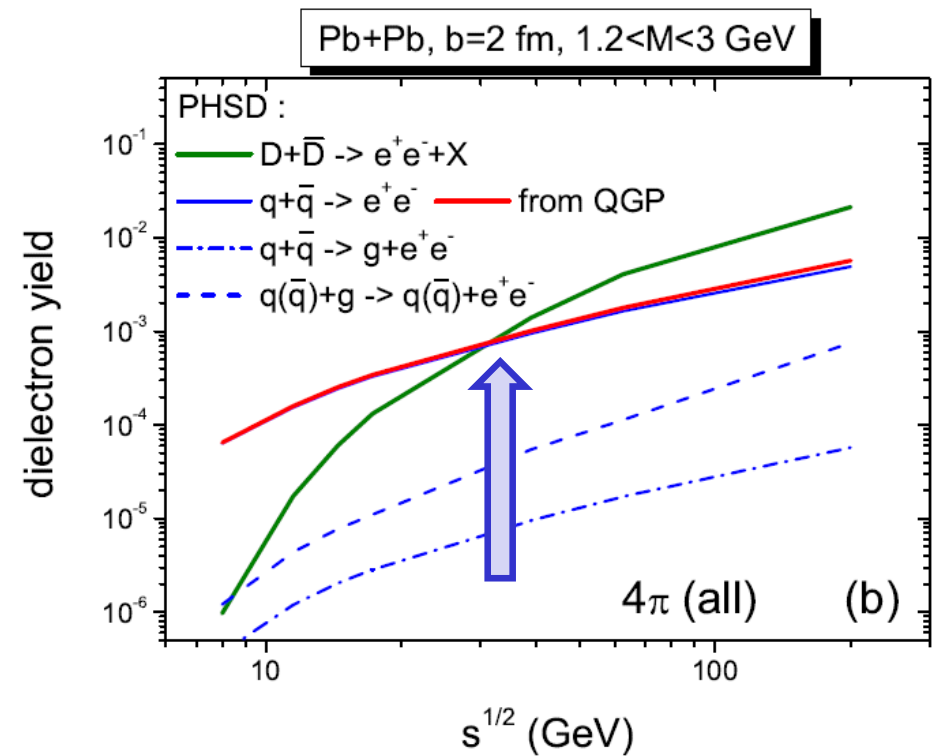
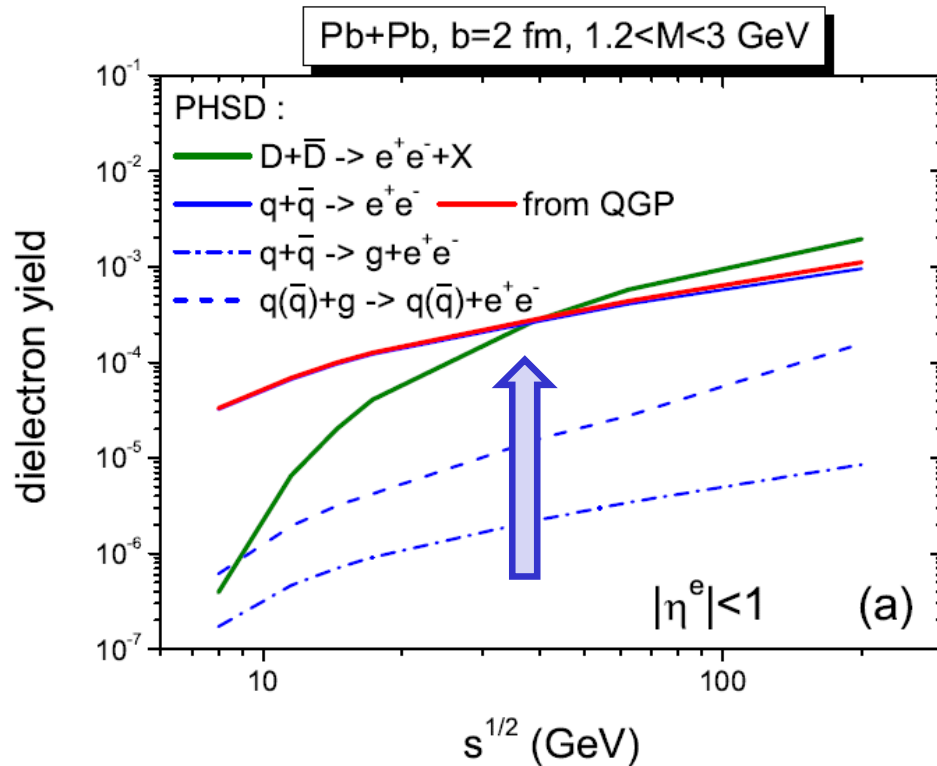


# Dileptons: QGP vs charm

Excitation function of dilepton multiplicity integrated for  $1.2 < M < 3 \text{ GeV}$

mid-rapidity

all rapidities ( $4\pi$ )

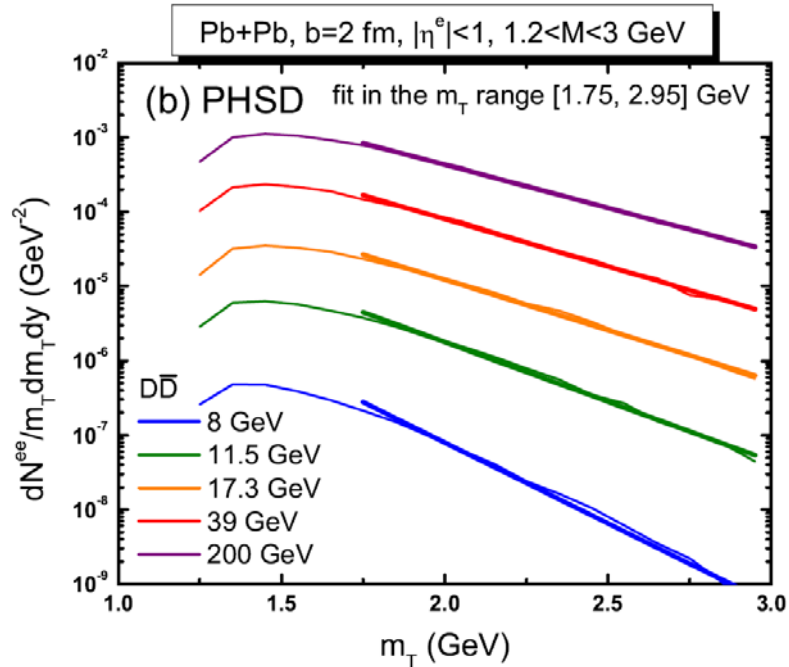
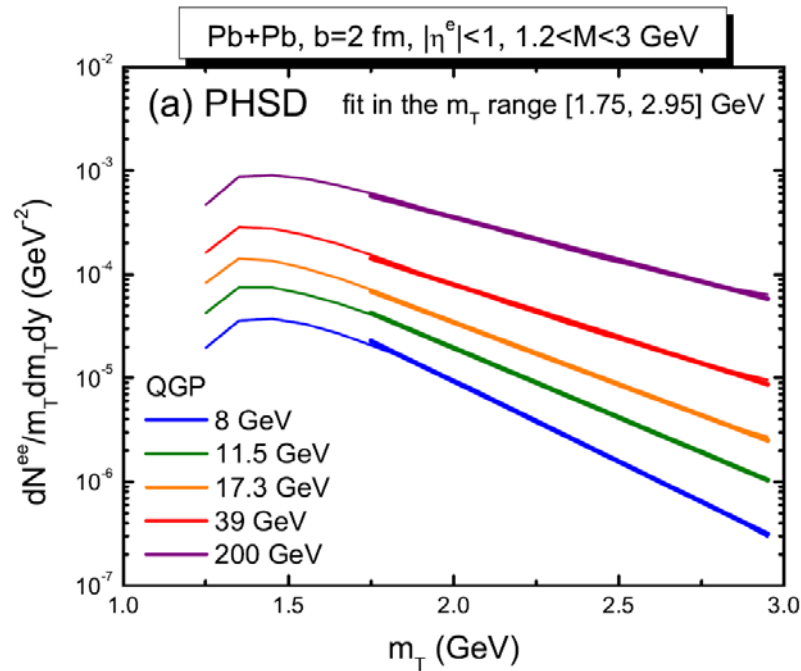


**QGP contribution overshines charm with decreasing energy!**

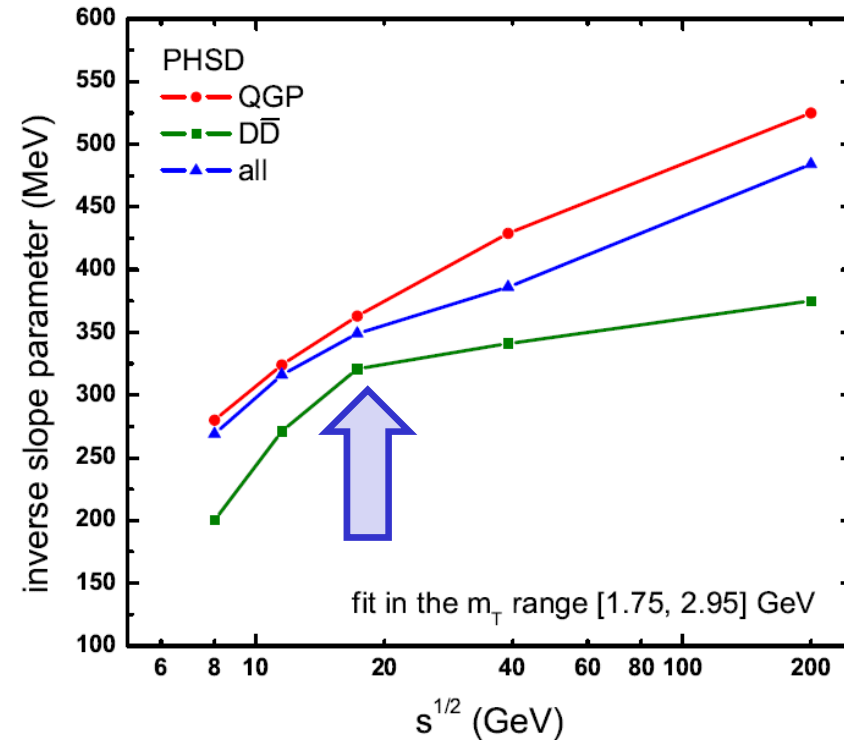
**→ Good perspectives for FAIR/NICA and BES RHIC!**



# Dilepton transverse mass spectra



The **inverse slope parameter** in the mass range [1.75, 2.95]



- **Inverse slope parameter: QGP contribution is **harder** than that from D-Dbar**
- **The excitation function of the total inverse slope parameter shows **characteristic changes at  $s^{1/2} > 20$  GeV****

# Messages from the dilepton study

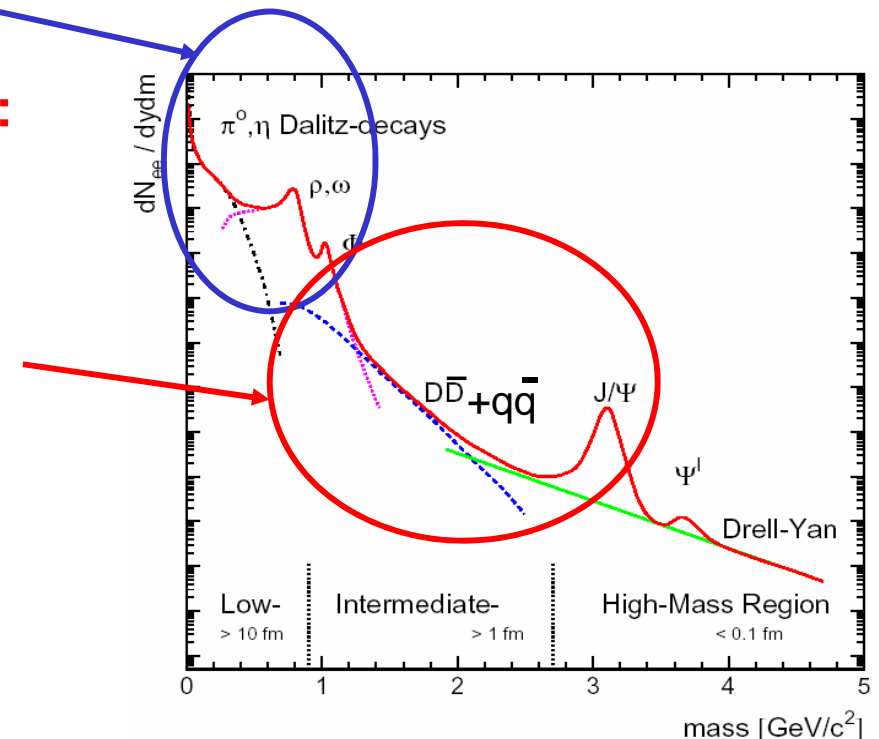


## □ Low dilepton masses:

- Dilepton spectra show **sizeable changes due to the in-medium effects – modification of the properties of vector mesons** (as collisional broadening) – which are observed experimentally
- **In-medium effects** can be observed at **all energies from SIS to LHC;** **excess increasing with decreasing energy** due to a longer  $\rho$ -propagation in the high baryon-density phase

## □ Intermediate dilepton masses $M > 1.2$ GeV :

- Dominant sources : **QGP** ( $q\bar{q}$ ), correlated charm  $D/D\bar{c}$
  - Fraction of QGP **grows** with increasing energy; however, the relative contribution of QGP to dileptons from charm pairs increases with decreasing energy
- Good perspectives for **FAIR/NICA**



Review: O. Linnyk et al., Prog. Part. Nucl. Phys. 89 (2016) 50

T. Song, W. Cassing, P. Moreau and E. Bratkovskaya, Phys. Rev. C 97 (2018) 064907

**Thank you for your attention !**



**Thanks to the Organizers !**