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#### Neutrino structure functions predictions for high energy and high precision neutrino experiments

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## **Motivations**

- Interaction of high energy neutrinos on hadron targets are an important probe to test QCD and to understand the parton properties of hadron structure
- Combinations of neutrino and anti-neutrino scattering data used to determine the structure functions
- The structure function F<sub>2</sub> is the singlet distribution

$$F_2^{\nu N} \propto xq^S = x\sum(q+\bar{q})$$

 Phenomenology using saturation models within the Color Dipole Approach successfully describes current small-x data <sup>1</sup>



<sup>1</sup>GAY DUCATI, M. B., MACHADO, M. M., MACHADO, M. V. T. – PLB 644 (2007) 340;

## Motivations (MINERvA)<sup>2</sup>

- Fermi National Accelerator Laboratory
- First neutrino experiment in the world to use a high-intensity beam to study neutrino reactions with five different nuclei
- First self-contained comparison of interactions in different elements.
- High precision measurements of neutrino interactions
- Information about proton's structure
- Water, helium, carbon, iron, lead, and plastic.





# Motivations (NuSOnG)

- Purpose of a new high-energy, ultra-high statistics neutrino scattering experiment
- Neutrino Scattering On Glass (NuSOnG)<sup>3</sup>
- Variety of issues including precision QCD measurements, extraction of structure functions, and the derived Parton Distribution Functions (PDFs)
- Tevatron-based neutrino beam to obtain DIS events
- Discerning tests of fundamental Standard Model parameters





<sup>3</sup>T. ADAMS et al, Int.J.Mod.Phys.A25:909-949,2010

## Neutrino-nucleon collision



$$s = 2ME$$
  $y = \frac{pq}{ME}$ 

 $x = \frac{Q^2}{2pq}$ 





- M is the nucleon mass
- E is the neutrino energy
- p and q are the nucleon and boson four-momenta
- Charged-current interactions and neutral current interactions



# Neutrino-nucleon cross section

$$\frac{\partial^2 \,\sigma_{(\nu,\,\bar{\nu})}^{\rm NC}}{\partial x \,\partial y} = \frac{G_F^2 \,m_N \,E_\nu}{\pi} \left(\frac{m_Z^2}{Q^2 + m_Z^2}\right)^2 \left[\frac{1 + (1 - y)^2}{2} F_2^{\rm NC}(x,Q^2) - \frac{y^2}{2} F_L^{\rm NC}(x,Q^2)\right]$$

G<sub>F</sub> is the Fermi constant

1.166.10<sup>-5</sup> GeV<sup>-2</sup>

- M<sub>i</sub> is the boson mass
- F<sub>2</sub>, F<sub>L</sub> and F<sub>3</sub> are the structure functions

 $Q^2_{min} = 1 GeV^2 \qquad \qquad y = Q^2/(xs)$ 

$$\frac{d\sigma^{\nu(N)}}{dxdy} = \frac{G_F^2 m_N E_{\nu}}{\pi} \left[ xq(x) + x\bar{q}(x)(1-y)^2 \right]$$

$$\frac{d\sigma^{\overline{v}(N)}}{dxdy} = \frac{G_F^2 m_N E_v}{\pi} \Big[ xq(x)(1-y)^2 + x\overline{q}(x) \Big]$$

$$\sigma_{(\nu,\bar{\nu})}^{\rm NC}(E_{\nu};\,A) = \int_{Q^2_{\rm min}}^s dQ^2 \int_{Q^2/s}^1 dx \,\frac{1}{xs} \frac{\partial^2 \,\sigma_{(\nu,\bar{\nu})}^{\rm NC}}{\partial x \,\partial y}$$



<sup>4</sup> ROBERTS, R. G., "The structure of the proton", Cambridge University Press (1993);

## Color dipole phenomenology



- Ψ's are the wave functions for electroweak bosons
- z is the momentum fraction of quark and (1-z) is the momentum fraction of the antiquark
- $\lambda_1$  and  $\lambda_2$  are the helicities of the quarks (1/2 or -1/2)
- r is the transversal size of the dipole
- $\sigma_{\text{dip}}$  is parametrized and fitted to the experiment .

$$\sigma_{L,R}(x,Q^2) = \int d^2 \mathbf{r} \int_0^1 dz \sum_{\lambda_1,\lambda_2} |\Psi_{L,R}^{\lambda_1,\lambda_2}(z,\mathbf{r},Q^2)|^2 \sigma_{dip}(x,\mathbf{r})$$



#### **Neutrino Structure functions**

$$\begin{split} F_{T,L}^{\text{CC,NC}}(x,Q^2) &= \frac{Q^2}{4\pi^2} \int d^2 r \int_0^1 dz \, |\psi_{T,L}^{W^{\pm},Z^0}|^2 \, \sigma_{dip} \\ |\psi_T^{W^{\pm}}| &= \frac{4N_c}{(2\pi)^2} \left\{ [(1-z)^2 m_q + z^2 m_{\bar{q}}^2] K_0^2(\epsilon r) + [z^2 + (1-z)^2] \epsilon^2 K_1^2(\epsilon r) \right\} \\ |\psi_L^{W^{\pm}}| &= \frac{4N_c}{(2\pi)^2 Q^2} \left\{ [(z(1-z)Q^2 + \epsilon^2)^2 + m_q^2 m_{\bar{q}}^2] K_0^2(\epsilon r) + [\frac{a_-^2 + a_+^2}{2}] \epsilon^2 K_1^2(\epsilon r) \right\} \\ |\bar{\psi}_L^T(r,z,Q^2)|^2 &= \frac{3}{2\pi^2} (L_u^2 + L_d^2 + R_u^2 + R_d^2) [z^2 + (1-z)^2] \bar{Q}^2 K_1^2(\bar{Q}r) \\ |\bar{\psi}_L^Z(r,z,Q^2)|^2 &= \frac{6}{\pi^2} (L_u^2 + L_d^2 + R_u^2 + R_d^2) z^2 (1-z)^2 Q^2 K_0^2(\bar{Q}r) \\ |\bar{\psi}_L^Z(r,z,Q^2)|^2 &= \frac{6}{\pi^2} (L_u^2 + L_d^2 + R_u^2 + R_d^2) z^2 (1-z)^2 Q^2 K_0^2(\bar{Q}r) \\ \text{Chiral coupling} \begin{cases} L_u &= 1 - \frac{4}{3} sin^2 \theta_W & L_d = -1 + \frac{2}{3} sin^2 \theta_W \\ R_u &= -\frac{4}{3} sin^2 \theta_W & R_d = \frac{2}{3} sin^2 \theta_W \end{cases}$$



### **Quark distribution**

 Gluon emits a quark-antiquark pair changing the quark distribution in the nucleon

- These quarks are called sea quarks
- Quark content is given by the sum of valence quarks and sea quarks

$$u(x) = u_v(x) + u_s(x),$$
  

$$d(x) = d_v(x) + d_s(x)$$
  

$$u_s(x) = \bar{u}_s(x) = d_s(x) = \bar{d}_s(x) = s_s(x) = \bar{s}_s(x)$$
  

$$2xF_1 = F_2 = 2x \sum_i (q_i + \bar{q}_i)$$
  

$$xF_3 = 2x \sum_i (q_i - \bar{q}_i)$$

aytambo, Sacred Valley, Peru, 9 - 15 Septem

## **Dipole cross section**

• Golec-Biernat-Wusthoff (GBW) <sup>5</sup>

$$\sigma_{dip}(x,r^2) = \sigma_0 \left[ 1 - \exp\left(-\frac{r^2 Q_{sat}^2}{4}\right) \right]$$

$$\bar{x} = x[1 + (4m_f^2/Q^2)]$$

- $\sigma_0 = 23 \text{ mb},$
- λ ~ 0.288,
- $x_0 \sim 3.10^{-4} \,\mathrm{m}$ ,
- m<sub>f</sub> = 0.14 GeV

$$Q_{sat}^2(x) = \left(\frac{x_0}{\bar{x}}\right)^\lambda GeV^2$$

<sup>5</sup> GOLEC-BIERNAT, K; WUSTHOFF, M. PRD 60, 1140231 (1998);



## Neutrino-nuclei interaction <sup>6</sup>

 Dipole cross section for bosons transversally or longitudinally polarized are extended for nuclei using Glauber-Gribov formalism

$$\begin{split} \sigma_{L,R}(x,Q^2) = & < \Psi_{\lambda} \mid \sigma^A(x,r) \mid \Psi_{\lambda} > \\ & = \int d^2 r \int_0^1 dz \sum_{\lambda_1,\lambda_2} \left| \Psi_{L,R}^{\lambda_1,\lambda_2}(z,r,Q^2) \right|^2 \sigma^A_{dip.}(x,r) \\ \sigma_{dip}^{\text{nucleus}}(x,r;A) = 2 \int d^2 b \left\{ 1 - \exp\left[-\frac{1}{2} T_A(b) \, \sigma_{dip}^{\text{nucleon}}(x,r)\right] \right\} \end{split}$$

Nuclear profile function T<sub>A</sub> (b)

$$T_A(b) = \int_{-\infty}^{+\infty} dz n(\sqrt{z^2 + b^2})$$

 b is the impact parameter and n(r) is the nuclear matter density normalized as

$$\int d^3 r n(r) = A$$

<sup>6</sup> WATT, G. KOWALSKI, H. PRD 78 (2008) 014016

## xF1 Structure functions (x fixed GBW model)



#### xF<sub>3</sub> Structure functions (x fixed)



### ΔxF<sub>3</sub> Structure functions (x fixed)





#### Conclusions

 Analysis of small-x neutral current neutrino-nucleus is performed within the color dipole formalism

- Two new experimental with nuclei collisions (MINERvA and NuSOnG)
- Structure functions  $xF_1$ ,  $\Delta xF_3$  and  $xF_3$  are investigated
- Predictions for Iron and Glass nuclei collisions
- Further investigations are requested

