



Wydział
Fizyki

POLITECHNIKA WARSZAWSKA

From quantum theory of nuclear matter to hydrodynamics of neutron stars

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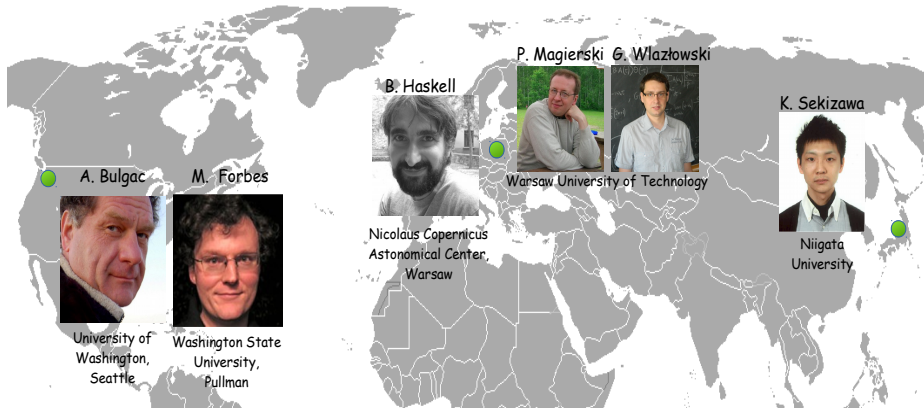
Piotr Magierski

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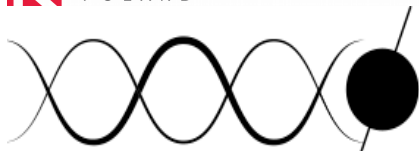
Nuclear theory group



Colaboration



NATIONAL SCIENCE CENTRE
POLAND



PHAROS
THE MULTI-MESSENGER
PHYSICS AND ASTROPHYSICS
OF NEUTRON STARS



Presentation plan

Open question: **How to perform transition from quantum to macroscopic regime preserving "the physics"?**

1 Introduction

- Star internal structure
- Description on different length scales

2 Towards mesoscopic model

- Methods
- Reduction of DoFs
- Vortex Filament Model

3 Neutron Star crust modelling

- Vortex self-interaction
- Phenomenological dissipation
- Applications

4 Future work objectives



Presentation plan

Correction to M. Antonelli: **Transition from superfluid coherence length to intervortex distance scale.**

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Plan of presentation

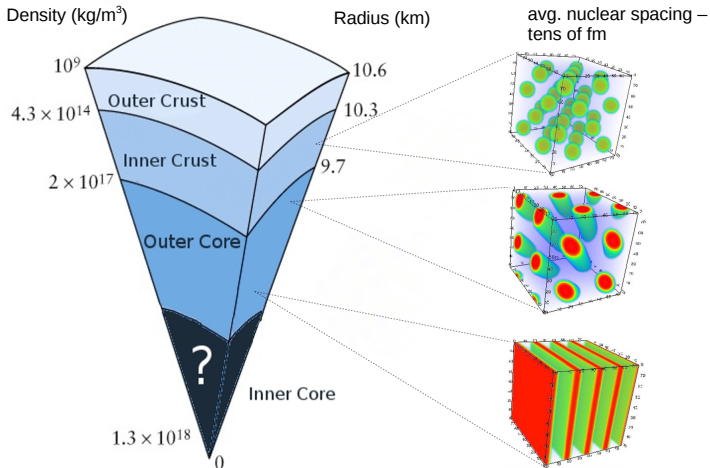
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Star internal structure

Macroscopic - hydrodynamics

Microscopic - quantum



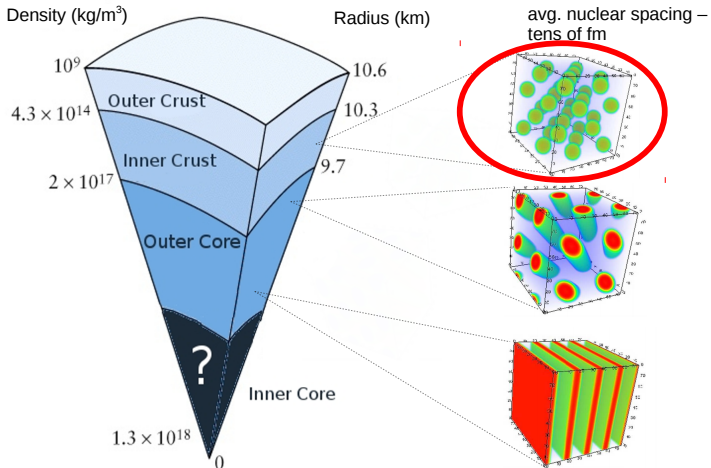
Source: <https://apatruno.wordpress.com/neutron-stars/>

M. Okamoto et al.,
Phys. Rev. C 88.025801 (2013)

Star internal structure

Macroscopic - hydrodynamics

Microscopic - quantum



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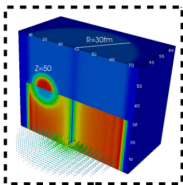


Description on different length scales



Microscopic

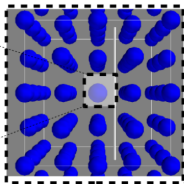
- Quantum many-body description
- Dynamics of neutrons and protons



Method: TDDFT
DoF: neutrons and protons.
Scale: $\sim 10^{-13}$ m

Mesoscopic

- Semi-classical model
- **Distinguishable vortices and nuclei**



Method: Vortex Filament Model
DoF: impurities and vortices
Scale: $\sim 10^{-9}$ m

Macroscopic

- Relativistic hydrodynamics?
- Continuous vorticity



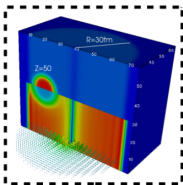
Method: Hydrodynamics
DoF: fluid elements
Scale: \sim size of star

Description on different length scales



Microscopic

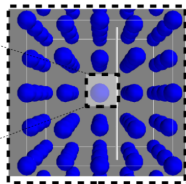
- Quantum many-body description
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Mesoscopic

- Semi-classical model
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Macroscopic

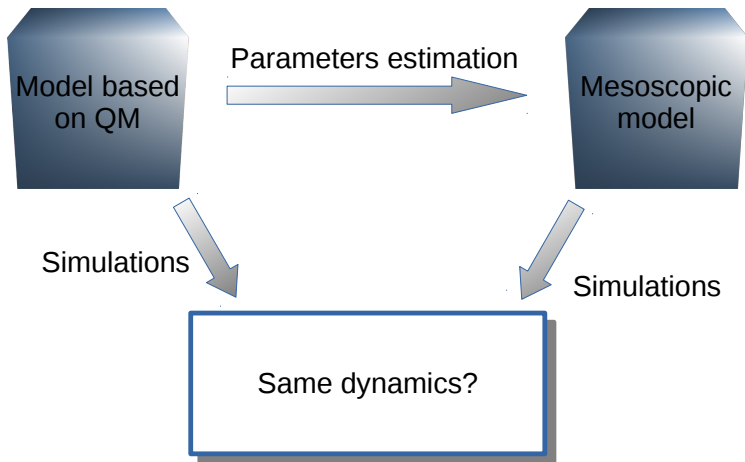
- Relativistic hydrodynamics?
- Continuous vorticity



Method: Hydrodynamics
DoF: fluid elements
Scale: \sim size of star

Collective effort required!

Our aim



Plan of presentation

- 1 Introduction
- 2 Towards mesoscopic model
- 3 Neutron Star crust modelling
- 4 Future work objectives



Methods - our proposal



Microscopic

- ASLDA framewrok
- Explicit s-wave superfluidity
- Equivalent of Schrödinger eq.

Mesoscopic

- Relativistic?
- Vortex Filament Model
- Explicit dynamics of nuclei and vortices

Macroscopic

- Relativistic
- Fluid of clusters
- Fluid of neutrons
- Continous vorticity

- At macroscopic level we want to describe fluid of nuclei and neutronic superfluid
- How to estimate friction between them?
- Pinning force between nuclei and vortices in neutron superfluid



A. Melatos, C. Peralta
Apj, 709, 77 (2010)

Methods



Microscopic

- ASLDA framewrok
- Explicit s-wave superfluidity
- Equivalent of Schrödinger eq.

Mesoscopic

- Relativistic?**
- Vortex Filament Model**
- Explicit dynamics of nuclei and vortices**

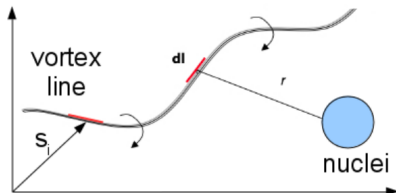
Macroscopic

- Relativistic
- Fluid of nuclei
- Fluid of neutrons
- Continous vorticity

$$\rho \mathbf{K} \times (\dot{\mathbf{s}} - \mathbf{v}_{ind}(\mathbf{s}) - \mathbf{v}_{ext}) + \mathbf{f}^{VN}(\mathbf{s}, \mathbf{r}) + \mathbf{f}^D(\dot{\mathbf{s}}, \dot{\mathbf{r}}) = 0$$

$$\mathbf{s} = \mathbf{s}(\xi)$$

where ξ is a parameter describing position of line element along curve





Methods

Microscopic

- ASLDA framewrok**
- Explicit s-wave superfluidity**
- Equivalent of Schrödinger eq.**

Mesoscopic

- Relativistic?
- Vortex Filament Model
- Explicit dynamics of nuclei and vortices

Macroscopic

- Relativistic
- Fluid of nuclei
- Fluid of neutrons
- Continous vorticity

Two sets HFB equations with density functional for protons and neutrons

$$i\hbar \frac{\partial}{\partial t} \begin{pmatrix} u_{n\uparrow}(\mathbf{r}, t) \\ u_{n\downarrow}(\mathbf{r}, t) \\ v_{n\uparrow}(\mathbf{r}, t) \\ v_{n\downarrow}(\mathbf{r}, t) \end{pmatrix} = \begin{pmatrix} h_{\uparrow,\uparrow}(\mathbf{r}, t) & h_{\uparrow,\downarrow}(\mathbf{r}, t) & 0 & \Delta(\mathbf{r}, t) \\ h_{\downarrow,\uparrow}(\mathbf{r}, t) & h_{\downarrow,\downarrow}(\mathbf{r}, t) & -\Delta(\mathbf{r}, t) & 0 \\ 0 & -\Delta^*(\mathbf{r}, t) & -h_{\uparrow,\uparrow}^*(\mathbf{r}, t) & -h_{\uparrow,\downarrow}^*(\mathbf{r}, t) \\ \Delta^*(\mathbf{r}, t) & 0 & -h_{\downarrow,\uparrow}^*(\mathbf{r}, t) & -h_{\downarrow,\downarrow}^*(\mathbf{r}, t) \end{pmatrix} \begin{pmatrix} u_{n\uparrow}(\mathbf{r}, t) \\ u_{n\downarrow}(\mathbf{r}, t) \\ v_{n\uparrow}(\mathbf{r}, t) \\ v_{n\downarrow}(\mathbf{r}, t) \end{pmatrix}$$

A. Bulgac, Phys. Rev. A 76, 040502 (2007)



SLDA method capabilities

Two sets HFB equations with density functional for protons and neutrons

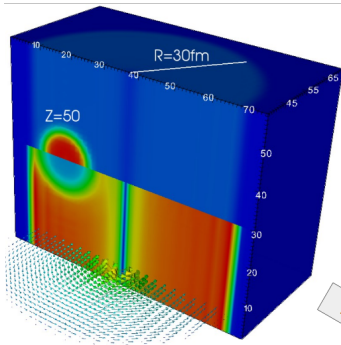
$$i\hbar \frac{\partial}{\partial t} \begin{pmatrix} u_{n\uparrow}(\mathbf{r}, t) \\ u_{n\downarrow}(\mathbf{r}, t) \\ v_{n\uparrow}(\mathbf{r}, t) \\ v_{n\downarrow}(\mathbf{r}, t) \end{pmatrix} = \begin{pmatrix} h_{\uparrow,\uparrow}(\mathbf{r}, t) & h_{\uparrow,\downarrow}(\mathbf{r}, t) & 0 & \Delta(\mathbf{r}, t) \\ h_{\downarrow,\uparrow}(\mathbf{r}, t) & h_{\downarrow,\downarrow}(\mathbf{r}, t) & -\Delta(\mathbf{r}, t) & 0 \\ 0 & -\Delta^*(\mathbf{r}, t) & -h_{\uparrow,\uparrow}^*(\mathbf{r}, t) & -h_{\uparrow,\downarrow}^*(\mathbf{r}, t) \\ \Delta^*(\mathbf{r}, t) & 0 & -h_{\uparrow,\downarrow}^*(\mathbf{r}, t) & -h_{\downarrow,\downarrow}^*(\mathbf{r}, t) \end{pmatrix} \begin{pmatrix} u_{n\uparrow}(\mathbf{r}, t) \\ u_{n\downarrow}(\mathbf{r}, t) \\ v_{n\uparrow}(\mathbf{r}, t) \\ v_{n\downarrow}(\mathbf{r}, t) \end{pmatrix}$$

Present computing capabilities:

(for nuclear systems)

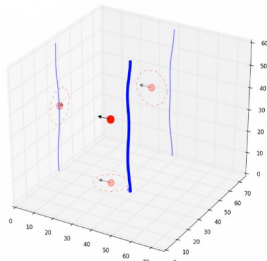
- ▶ full 3D (unconstrained) dynamics
- ▶ volumes up to 100^3 fm^3
- ▶ number of particles of order 10^4
- ▶ up to 10^6 time steps (for nuclear systems it gives trajectory length 10^{-19} sec)

Aim to construct mesoscopic model preserving low-energy physics



REDUCTION

When exact dynamics of quantum system is known a semi-classical model preserving low-energy excitations dynamics can be constructed.

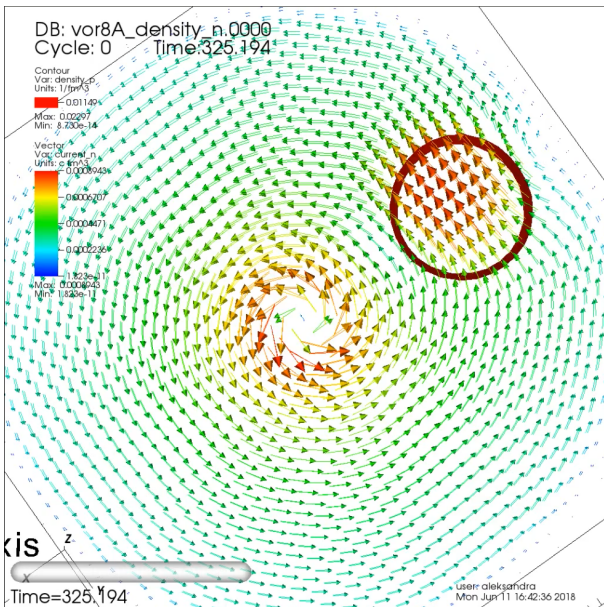


Exact dynamics of all fermionic DoFs
Computational cost limits size of system

$$i\hbar \frac{\partial}{\partial t} \begin{pmatrix} u_{n\uparrow}(\mathbf{r}, t) \\ u_{n\downarrow}(\mathbf{r}, t) \\ v_{n\uparrow}(\mathbf{r}, t) \\ v_{n\downarrow}(\mathbf{r}, t) \end{pmatrix} = \begin{pmatrix} h_{\uparrow,\uparrow}(\mathbf{r}, t) & h_{\uparrow,\downarrow}(\mathbf{r}, t) & 0 & \Delta(\mathbf{r}, t) \\ h_{\downarrow,\uparrow}(\mathbf{r}, t) & h_{\downarrow,\downarrow}(\mathbf{r}, t) & -\Delta(\mathbf{r}, t) & 0 \\ 0 & -\Delta^*(\mathbf{r}, t) & -h_{\uparrow,\uparrow}^*(\mathbf{r}, t) & -h_{\uparrow,\downarrow}^*(\mathbf{r}, t) \\ \Delta^*(\mathbf{r}, t) & 0 & -h_{\downarrow,\uparrow}^*(\mathbf{r}, t) & -h_{\downarrow,\downarrow}^*(\mathbf{r}, t) \end{pmatrix} \begin{pmatrix} u_{n\uparrow}(\mathbf{r}, t) \\ u_{n\downarrow}(\mathbf{r}, t) \\ v_{n\uparrow}(\mathbf{r}, t) \\ v_{n\downarrow}(\mathbf{r}, t) \end{pmatrix}$$



Example of transition between micro- and meso- scopic regime



Quantum mechanics vs hydrodynamics

$$i\hbar\partial_t\Psi = \left[\hat{H}_{one-body} + \cancel{\hat{H}_{int}} + V_{ext}(x,t) \right] \Psi, \quad \Psi = \sqrt{n}e^{i\theta}$$

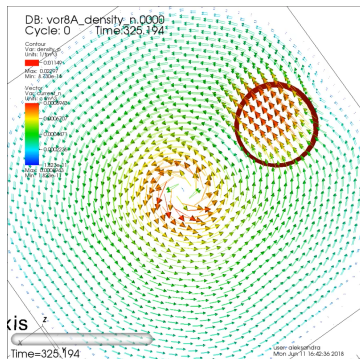
$$\downarrow$$

$$\frac{d\mathbf{v}}{dt} = \partial_t\mathbf{v} + \mathbf{v} \cdot \nabla\mathbf{v} = -\frac{1}{m}\nabla \left(-\frac{\hbar^2}{2m} \frac{\nabla^2\sqrt{n}}{\sqrt{n}} + V_{ext}(x,t) \right) \quad \partial_t n + \nabla(n\mathbf{v}) = 0$$

probability current

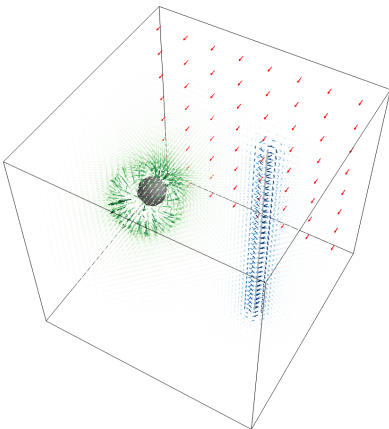
$$n\mathbf{v} = \mathbf{j} = \frac{\hbar}{2mi} [\Psi^*(\nabla\Psi) - \Psi(\nabla\Psi^*)]$$

Madelung, E. (1926). "Eine anschauliche Deutung der Gleichung von Schrödinger".
Naturwissenschaften. 14 (45): 1004–1004



Helmholz-Hodge decomposition (HHD)

Compressible and incompressible velocity fields



$$\mathbf{v} = \underbrace{\mathbf{v}_{ext}}_{\approx \text{const}} + \underbrace{\mathbf{v}_{ind}}_{\nabla \times \mathbf{A}} + \underbrace{\mathbf{v}_{nucl}}_{\nabla \phi_N} + \mathbf{v}_{\text{boundary}}$$

How to estimate $\nabla \phi_N$?



Vortex line motion

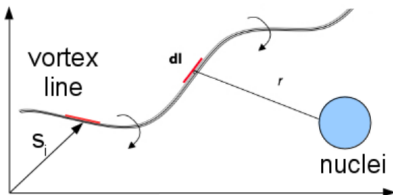
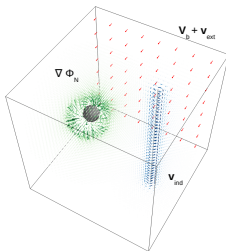
Description with fluid velocities

$$\dot{\mathbf{s}} = \mathbf{v}_{ind} + \mathbf{v}_b + \mathbf{v}_{ext} + \nabla\phi_N + \text{dissipative terms}$$

Description with forces acting on a vortex line

$$\rho_s \mathbf{K} \times \dot{\mathbf{s}} = \rho_s \mathbf{K} \times [\mathbf{v}_{ind} + \mathbf{v}_b + \mathbf{v}_{ext} + \nabla\phi_N] + \mathbf{f}^D$$

$$\rho_s \mathbf{K} \times \dot{\mathbf{s}} = \rho_s \mathbf{K} \times [\mathbf{v}_b + \mathbf{v}_{ext}] - \mathbf{f}^T - \mathbf{f}^{VN} - \mathbf{f}^D$$

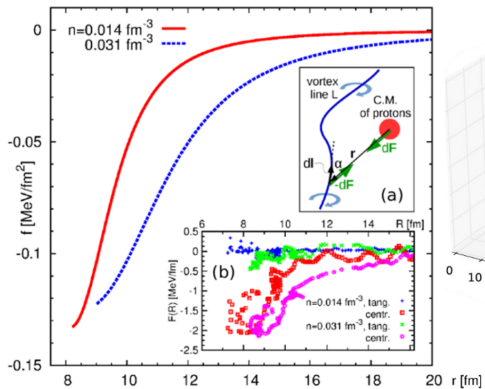


K. W. Schwarz, Phys. Rev. B 31, 5782 (1985)

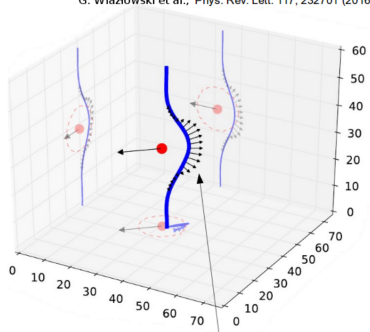




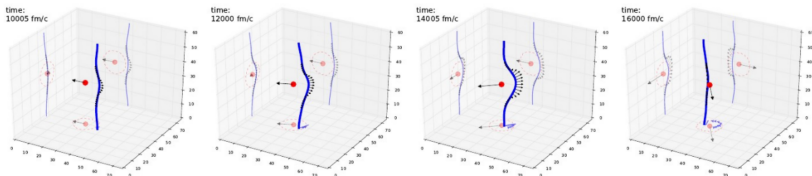
Vortex-nucleus interaction



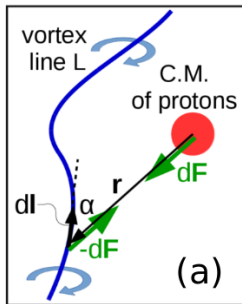
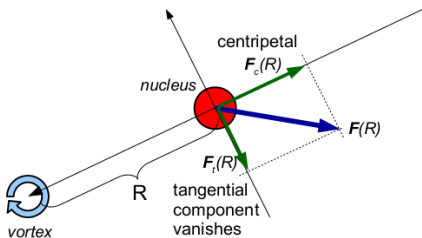
G. Wlazłowski et al., Phys. Rev. Lett. 117, 232701 (2016)



Nucleus is dragged by force adjusted to preserve uniform motion.
Interaction has to equilibrate system.

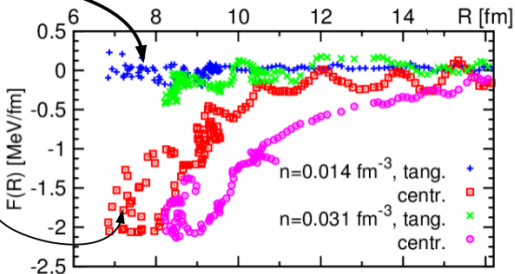


$$m_{nucl.} \frac{d}{dt} \mathbf{v}_{nucl.} = 0 = \mathbf{F}_{ext} + \hat{\mathbf{r}} \int_{\mathcal{L}} f(r) \sin \alpha dl$$



Force is central

At close separation the force is not a function of distance R only



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- 1 Introduction
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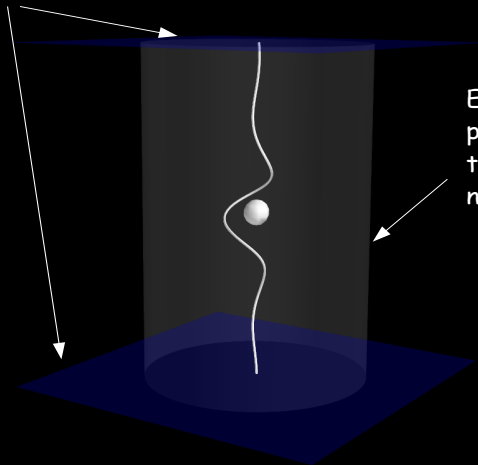


Setup of DFT simulations



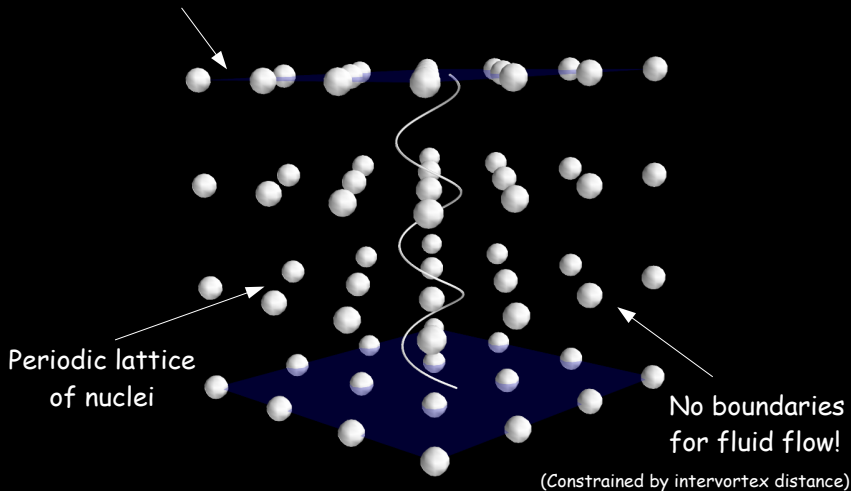
Periodic boundary conditions

External potential trapping neutron fluid



Setup of realistic system for NS crust simulation

Periodic boundary conditions - only horizontal direction



Self-induced velocity

$$\rho_s \mathbf{K} \times \dot{\mathbf{s}} = \rho_s \mathbf{K} \times \left[\mathbf{v}_b + \mathbf{v}_{ext} + \mathbf{v}_{ind} \right] - \mathbf{f}^{VN} - \mathbf{f}^D$$

$$\mathbf{v}_{ind}(\mathbf{r}, t) = \underbrace{\frac{\kappa}{4\pi} \int_{\mathcal{L}} dz \frac{(\mathbf{s}(z, t) - \mathbf{r}) \times \hat{\mathbf{t}}(z, t)}{|\mathbf{s}(z, t) - \mathbf{r}|^3}}_{\text{over vortex line}}$$

is divergent when $\mathbf{r} \rightarrow \mathbf{s}(z)$

Local Induction Approximation (LIA) - a canonical regularization

$$\mathbf{v}_{ind}(\mathbf{s}, t) = \mathbf{v}_{LIA}(\mathbf{s}, t) + \frac{\kappa}{4\pi} \underbrace{\int_{\mathcal{L}} dz \frac{(\mathbf{s}(z, t) - \mathbf{s}) \times \hat{\mathbf{t}}(z, t)}{|\mathbf{s}(z, t) - \mathbf{s}|^3}}_{\text{omitting divergent point}} \approx \mathbf{v}_{LIA}(\mathbf{r}, t)$$

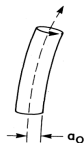
$$\mathbf{v}_{LIA}(\mathbf{s}(z), t) = \frac{\kappa}{4\pi R(z)} \ln \left(\frac{R(z)}{a_0} \right) \hat{\mathbf{b}}$$

a_0 - vortex core radius

$R(\mathbf{r})$ - radius of curvature

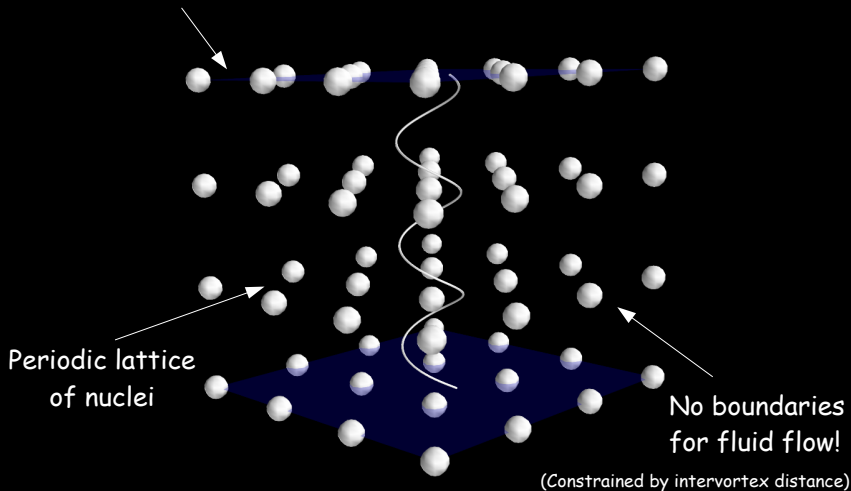


K. W. Schwarz,
Phys. Rev. B 31, 5782
(1985)



Setup of realistic system for NS crust simulation

Periodic boundary conditions - only horizontal direction





Self-induced velocity

$$\rho_s \mathbf{K} \times \dot{\mathbf{s}} = \rho_s \mathbf{K} \times \left[\mathbf{v}_b + \mathbf{v}_{ext} + \mathbf{v}_{ind} \right] - \mathbf{f}^{VN} - \mathbf{f}^D$$

$\boldsymbol{\omega}[\mathbf{s}] = \nabla \times \mathbf{v}[\mathbf{s}]$ - vorticity field, \mathbf{s} - vortex line element position

$$\partial_t \boldsymbol{\omega}[\mathbf{s}] = \nabla \times [\mathbf{v}[\mathbf{s}] \times \boldsymbol{\omega}[\mathbf{s}]] \quad \Longrightarrow \quad \nabla^2 \mathbf{A} = -\boldsymbol{\omega}[\mathbf{s}]$$

Effect of different domain

$$\mathbf{v}_{ind} = \nabla \times \mathbf{A}$$

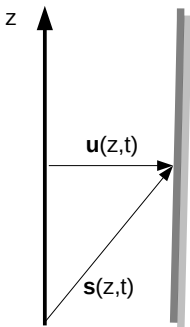
$$\mathbf{A} = \sum_{n=-\infty}^{+\infty} e^{ik_n(z-z')} \int_0^L \frac{\mathbf{K}(z')}{2\pi} K_0 \left(|k_n| \sqrt{(x-s_x(z'))^2 + (y-s_y(z'))^2} \right) dz'$$

Also divergent when $\mathbf{r} \rightarrow \mathbf{s}(z)$

Vortex self-interaction

$$\rho_s \mathbf{K} \times \dot{\mathbf{s}} = \rho_s \mathbf{K} \times [\mathbf{v}_b + \mathbf{v}_{ext}] - \mathbf{f}^T - \mathbf{f}^{VN} - \mathbf{f}^D$$

$$\mathbf{v}_{ind}(\mathbf{r}, t) = \frac{\kappa}{4\pi} \underbrace{\int_{\mathcal{L}} dz}_{\text{over vortex line}} \frac{(\mathbf{s}(z, t) - \mathbf{r}) \times \hat{\mathbf{t}}(z, t)}{|\mathbf{s}(z, t) - \mathbf{r}|^3} \text{ is divergent when } \mathbf{r} \rightarrow \mathbf{s}(z)$$



Energy flow can be rewritten as

$$E_V = \sum_n \omega(k_n) \hat{\mathbf{u}}(k_n) \cdot \hat{\mathbf{u}}^*(k_n)$$

Leading to

$$\mathbf{f}^T = \mathcal{F}^{-1} \left[\frac{\delta E_V(k_n)}{\delta \hat{\mathbf{u}}^*(k_n)} \right]$$

$$\mathbf{f}^T = \mathcal{F}^{-1} [\omega(k_n) \hat{\mathbf{u}}(k_n)]$$

A. Fetter, Phys. Rev. 162.143 (1967)
 sir W. Thompson, Phil. Mag. 10, 155 (1880)



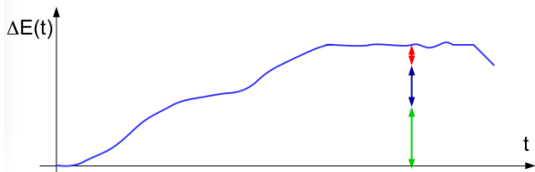
Dissipation force

$$\rho_s \mathbf{K} \times \dot{\mathbf{s}} = \rho_s \mathbf{K} \times [\mathbf{v}_b + \mathbf{v}_{ext}] - \mathbf{f}^T - \mathbf{f}^{VN} - \mathbf{f}^D$$

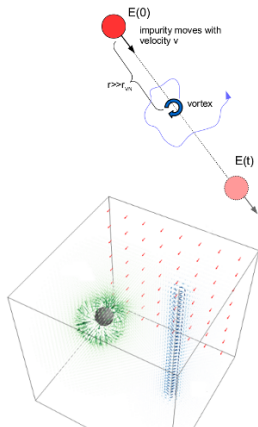
- Dissipation through vortex-phonon interactions $\mathbf{f}_1^D = -\eta_1 (\dot{\mathbf{u}} - \mathbf{v}_{ext})$
- Dissipation through nuclei deformation $\mathbf{f}_2^D = -\eta_2 (\dot{\mathbf{u}} - \dot{\mathbf{r}}_N)$

Estimation methodology

drag impurity with various velocities...
look at excitation energy...



$$\Delta E(t) = \Delta E_{\text{vortex}}(t) + \Delta E_{\text{compressible}}(t) + \Delta E_{\text{internal}}(t)$$



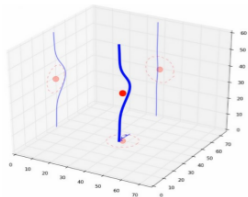


Dissipation force

$$\rho_s \mathbf{K} \times \dot{\mathbf{s}} = \rho_s \mathbf{K} \times [\mathbf{v}_b + \mathbf{v}_{ext}] - \mathbf{f}^T - \mathbf{f}^{VN} - \mathbf{f}^D$$

- Dissipation through vortex-phonon interactions $\mathbf{f}_1^D = -\eta_1 (\dot{\mathbf{u}} - \mathbf{v}_{ext})$
- Dissipation through nuclei deformation $\mathbf{f}_2^D = -\eta_2 (\dot{\mathbf{u}} - \dot{\mathbf{r}}_N)$

Estimation methodology



Reflected by the vortex line deformation in VFM

Use "ansatz" here for the dissipation force, and try to reproduce $W_d(t)$...

$$W_d(t) = \int_{\mathcal{L}} \left[\int_0^t \mathbf{f}^D(t') \cdot \dot{\mathbf{s}}_l(t') dt' \right] dl$$

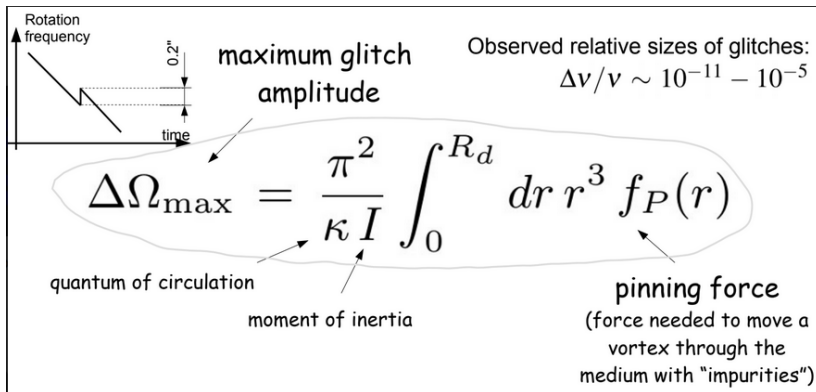
$W_d(t)$ = Dissipated energy
(from point of view of VFM)

Known from simulations...

$$\Delta E(t) = \Delta E_{\text{vortex}}(t) + \Delta E_{\text{compressible}}(t) + \Delta E_{\text{internal}}(t)$$



Glitch size estimation

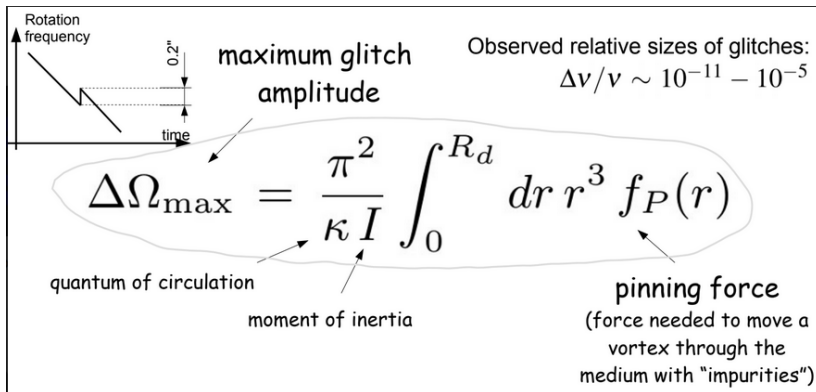


M. Antonelli, A. Montoli, P. M. Pizzochero, MNRAS 475, 5403 (2018)

Test for crustal vs flux tubes pinning!



Glitch size estimation

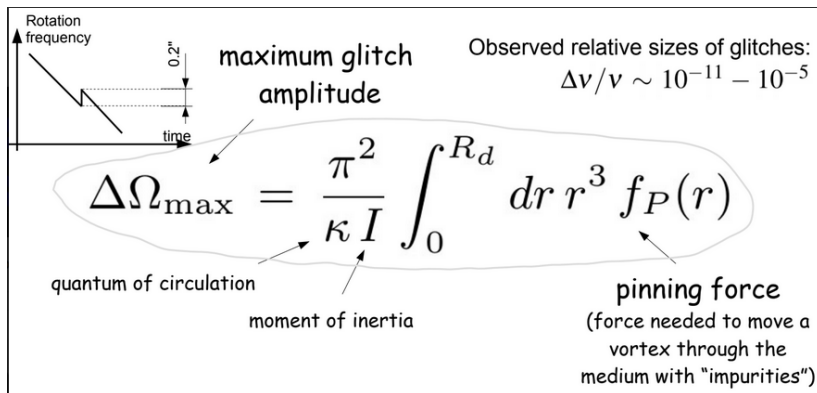


M. Antonelli, A. Montoli, P. M. Pizzochero, MNRAS 475, 5403 (2018)

Could provide test for microscopic models?



Glitch size estimation



M. Antonelli, A. Montoli, P. M. Pizzochero, MNRAS 475, 5403 (2018)

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Future work objectives



- Development of nuclear SLDA
 - 1 More accurate density functionals - BSk
 - 2 Nuclei lattice - structure and dynamics
- Vortex Filament Model validation
 - 1 Comparison between LIA and small deflections approach
 - 2 Extraction for Kelvin Waves dispersion relation
 - 3 Inclusion of nuclei dynamics
 - 4 Extraction of dissipation coefficients
- Large-scale VFM simulations
 - 1 Extraction of pinning force
 - 2 Connection to relativistic hydrodynamics of NS
- Rigorous testing of density functionals?

High performance computing

Rank	Site	System	Cores	Rmax (TFlop/s)	Rpeak (TFlop/s)	Power (kW)
1	DOE/SC/Dak Ridge National Laboratory United States	Summit - IBM Power System AC922, IBM POWER9 Z2C 3.070Hz, NVIDIA Volta GV100, Dual-rail Mellanox EDR Infiniband IBM	2,282,544	122,300.0	187,459.3	8,806
2	National Supercomputing Center in Wuxi China	Sunway TaihuLight - Sunway MPP, Sunway SW26010 260C 1.450Hz, Sunway NRCPC	10,649,600	93,014.6	125,435.9	15,371
3	DOE/NSA/LNL United States	Sierra - IBM Power System S922LC, IBM POWER9 Z2C 3.1GHz, NVIDIA Volta GV100, Dual-rail Mellanox EDR Infiniband IBM	1,572,480	71,610.0	119,193.6	
4	National Super Computer Center in Guangzhou China	Tianhe-2A - TH-IV9-FEP Cluster, Intel Xeon E5-2692v2 12C 2.20Hz, TH Express-2, Matrix-2000 NUOT	4,981,760	61,444.5	100,478.7	18,482
5	National Institute of Advanced Industrial Science and Technology (AIST) Japan	AI Bridging Cloud Infrastructure (ABCI) - PRIMERGY CX2550 M4, Xeon Gold 6148 20C 2.40Hz, NVIDIA Tesla V100 SXM2, Infiniband EDR Fujitsu	391,680	19,880.0	32,576.6	1,649
6	Swiss National Supercomputing Centre (CSCS) Switzerland	Piz Daint - Cray XC30, Xeon ES-2690v3 12C 2.60Hz, Aries interconnect , NVIDIA Tesla P100 Cray Inc.	361,760	19,590.0	25,326.3	2,272
7	DOE/SC/Dak Ridge National Laboratory United States	Titan - Cray XK7, Opteron 6274 14C 2.200GHz, Cray Gemini interconnect, NVIDIA K20x Cray Inc.	560,440	17,590.0	27,112.5	8,209
19	GSIC Center, Tokyo Institute of Technology Japan	TSUBAME3.0 - SGI ICE XA, IP139-SXM2, Xeon ES-2680v4 14C 2.40Hz, Intel Omni-Path, NVIDIA Tesla P100 SXM2 HPE	135,820	8,125.0	12,127.1	792

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TOP 500
The List.

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Thank you!

Effects of boundaries

$$\rho_s \mathbf{K} \times \dot{\mathbf{s}} = \rho_s \mathbf{K} \times \left[\mathbf{v}_b + \mathbf{v}_{ext} \right] - \mathbf{f}^T - \mathbf{f}^{VN} - \mathbf{f}^D$$

Flow past boundaries have to be suppressed $\implies \mathbf{v}^\infty \cdot \hat{\mathbf{n}} = \mathbf{v}_b \cdot \hat{\mathbf{n}}$

