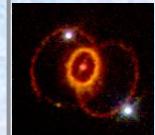


# Non-uniform (warm) matter and cluster distribution in core-collapse supernovae

Anthea F. Fantina ([anthea.fantina@ganil.fr](mailto:anthea.fantina@ganil.fr))

Collaborations : F. Gulminelli (LPC Caen, France)  
G. Grams, K. Yamasaki (Universidade Federal de Santa Catarina, Brasil)  
S. Giraud (GANIL)

Supernova Remnant 1987A in the Large Magellanic Cloud.  HUBBLESITE.org



# Open questions



- Treatment of non-uniform matter ?
  1. cluster formation & interactions
  2. SNA vs NSE

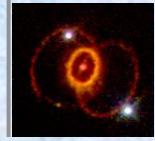
→ formalism to include cluster distribution in EoS  
& application in CCSNe
  
- Relevance of specific nuclear physics inputs / parameters ?

→ nuclear masses (experimental and reliable theoretical models)  
& CCSNe

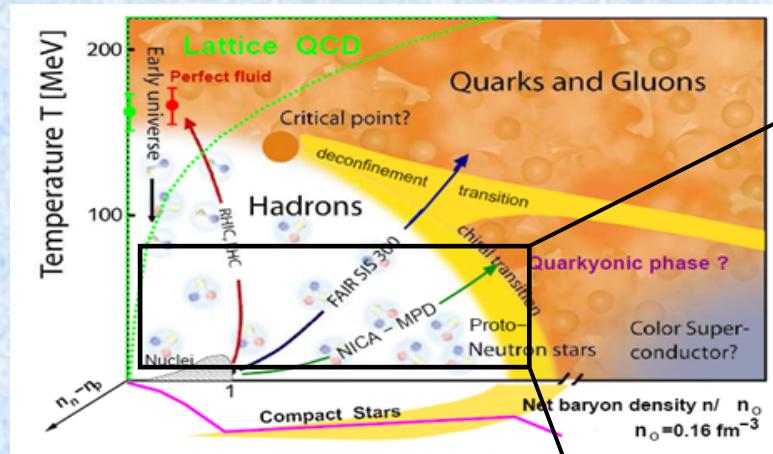


# Outline

- ❖ Introduction : EoS and treatment of non-uniform matter
- ❖ Nuclear distribution in CCSNe:
  - formalism for cluster distribution in a SNA EoS
  - impact of nuclear masses
- ❖ Conclusions



# Probing extreme conditions in SNe



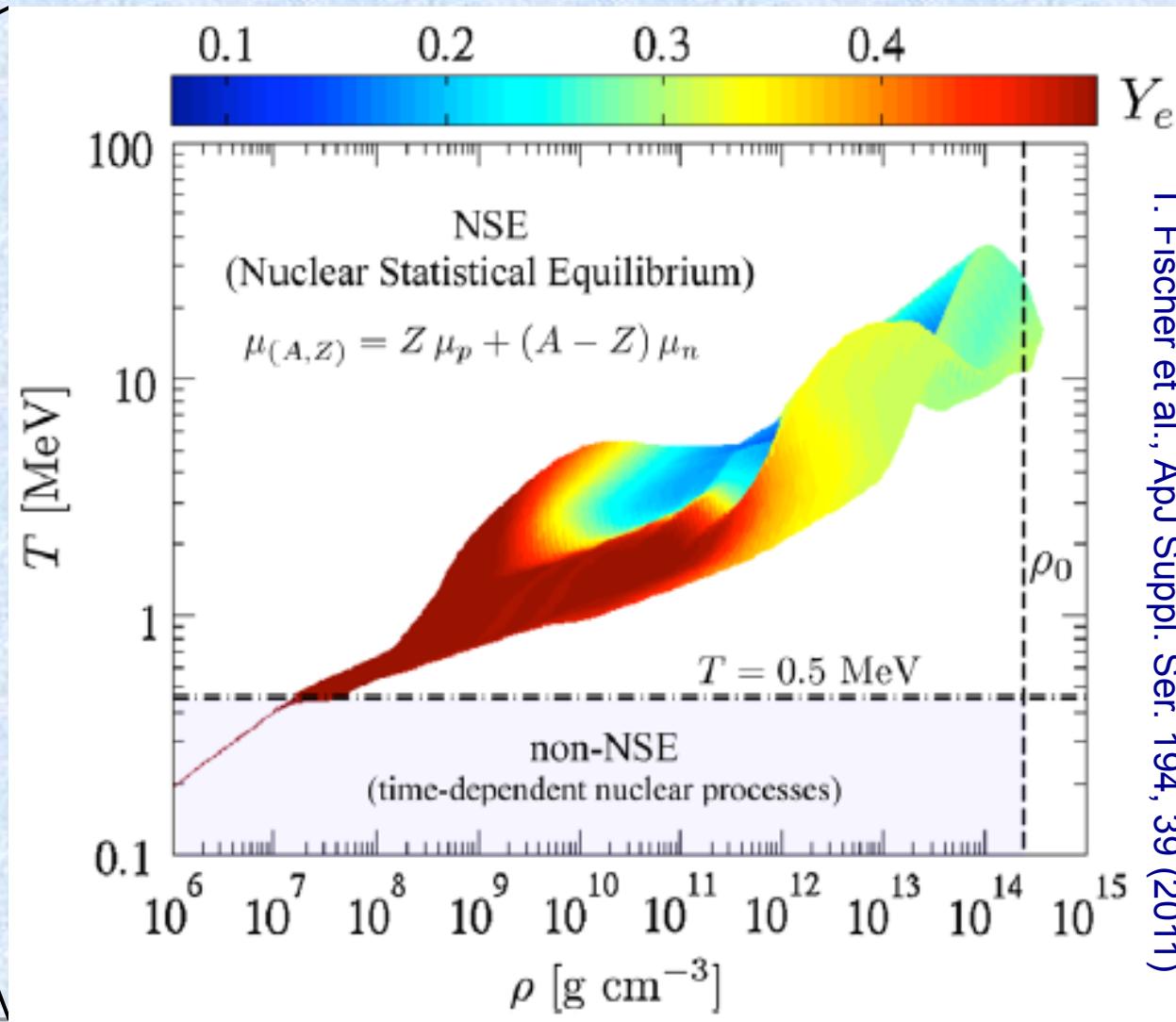
Condition in the core  
during collapse  
and NS formation :

$$\rho \in [10^5 - 10^{15}] \text{ g cm}^{-3}$$

$$T \in [0.1 - 100] \text{ MeV}$$

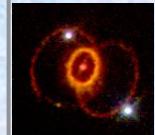
$$Y_e \in [0.05 - 0.5]$$

At present, best 3D hydro simulations  
do not reproduce satisfactory CCSN  
explosions !



see e.g., for a review, Oertel et al., Rev. Mod. Phys. 89, 015007 (2017)

Burgio & Fantina, Chap. 6 white book COST/NewCompStar, to be published



# Probing extreme conditions in CC

We focus on **infall phase**

- mostly at sub-saturation
- mostly  $T$  up to a few MeV



nuclei

Simulation in 1D, GR (Fantina, PhD thesis (2010))

$$\rho \approx [10^9 - \text{few } 10^{14}] \text{ g cm}^{-3}$$

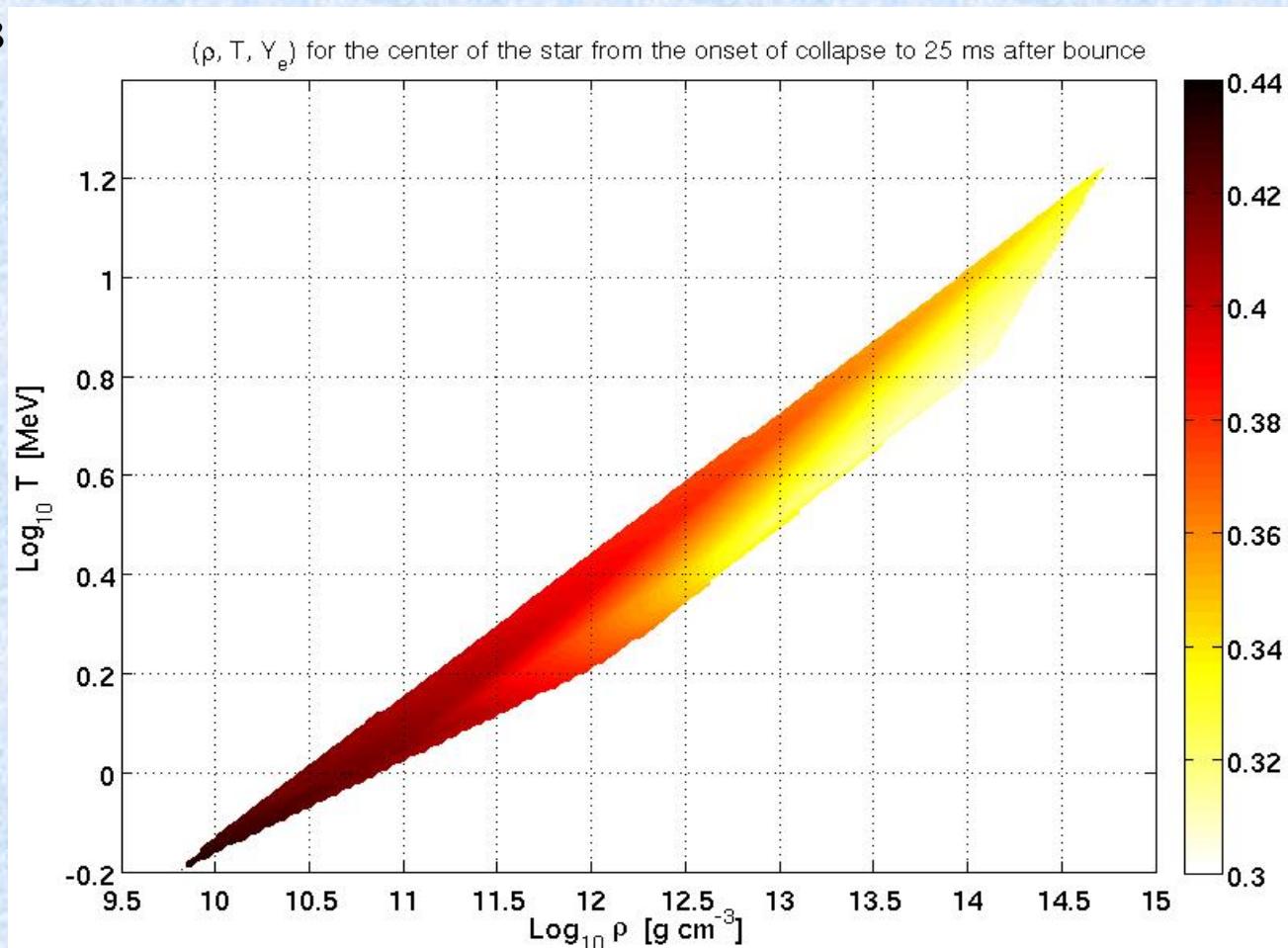
$$T \approx [0.6 - 20] \text{ MeV}$$

$$Y_e \approx [0.2 - 0.5]$$

For CC key nuclear data are:

- **EoS**  $\leftrightarrow$  Nuclear masses, level densities,...
- **EC**  $\leftrightarrow$  Q-val, GT response
- neutrino interactions

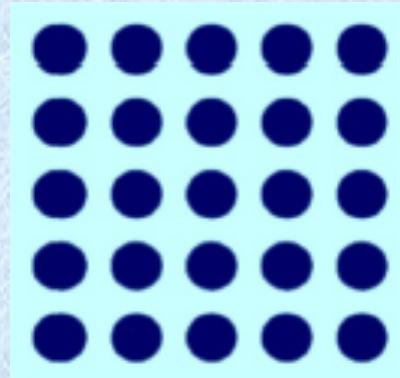
- ⇒ All possible nuclear species
- ⇒ What is the relevance of a specific nuclear physics input ?





# EoS in SN : SNA vs NSE

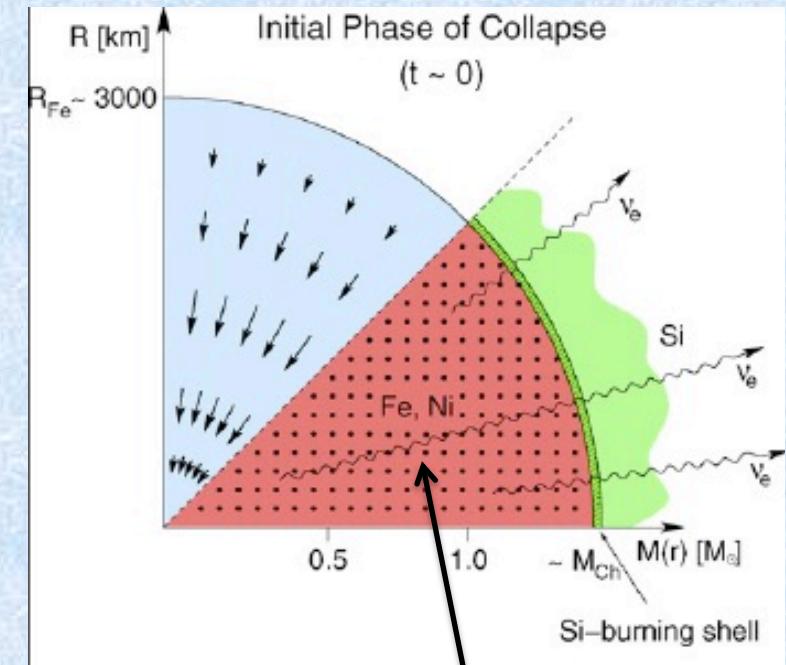
## Single Nucleus Approximation (SNA)



- One “heavy” ( $A, Z$ ) + gas  $n, p, e^-$  (+ alpha) for each thermo condition ( $n_B, T, Y_e$ )
- Used in most of EoSs in SN simulations
- ✓ OK for thermodynamic quantities
- ✓ Faster computationally

BUT:

- at finite  $T$  more microstates are populated
- reaction (EC) rates ?



Nuclei + gas of nucleons and leptons

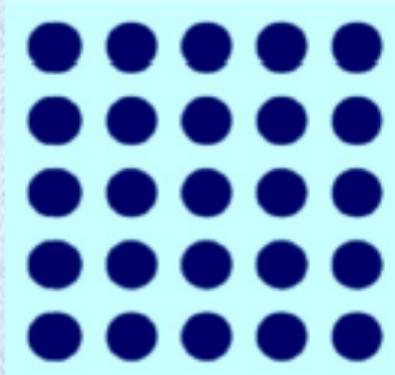
We focus on **infall phase**

- mostly at sub-saturation density
- mostly  $T$  up to a few MeV



# EoS in SN : SNA vs NSE

## Single Nucleus Approximation (SNA)

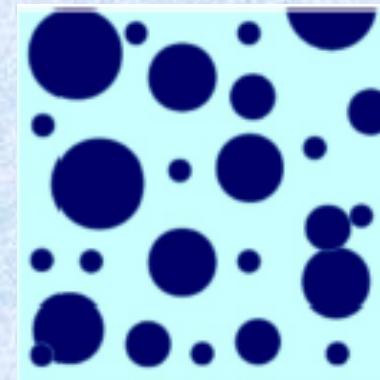


- One “heavy” ( $A, Z$ ) + gas  $n, p, e^-$  (+ alpha) for each thermo condition ( $n_B, T, Y_e$ )
- Used in most of EoSs in SN simulations
- ✓ OK for thermodynamic quantities
- ✓ Faster computationally

### BUT:

- at finite T more microstates are populated
- reaction (EC) rates ?

## Nuclear Statistical Equilibrium (NSE)

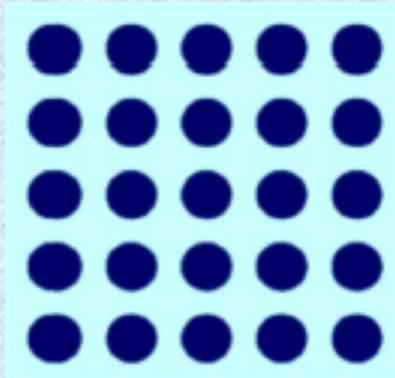


- Ensemble of nuclei ( $A_i, Z_i$ ) + gas  $n, p, e^-$  for each ( $n_B, T, Y_e$ )
  - Some NSE EoSs implemented in simulations
  - ✓ Distribution of nuclei
- BUT:
- implemented as tables → no flexibility
  - computationally expensive  
→ difficult to perform systematic calculations



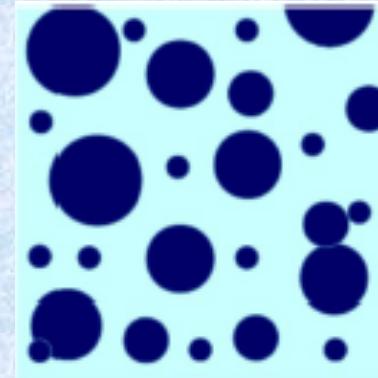
# EoS in SN : SNA vs NSE

## Single Nucleus Approximation (SNA)



- Compressible Liquid-drop models  
e.g. Baym et al. 1971, Lattimer & Swesty 1991, ...
- (Extended) Thomas-Fermi  
e.g. Shen et al. 1998, Onsi et al. 2008,  
Shen et al. 2011, ...
- Self-consistent mean field models  
e.g. Negele & Vautherin 1973, Baldo et al. 2007,  
Grill et al. 2011, Pais et al. 2014, ...

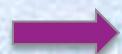
## Nuclear Statistical Equilibrium (NSE)

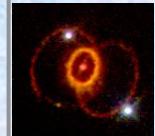


- (Extended) NSE  
e.g. Hillebrandt & Wolff 1985, Hempel &  
Schaffner-Bielich 2010, Raduta & Gulminelli  
2010, Gulminelli & Raduta 2014, Furusawa  
et al. 2013, 2017
- + in-medium effects:  
Virial EoS, models with in-medium mass  
shifts; e.g. Horowitz & Schwenk 2006, Ropke et  
al. 2011, 2013, 2015, Typel et al. 2010, 2014, ...

e.g., for a review, Oertel et al., Rev. Mod. Phys. 89, 015007 (2017) ; Burgio & Fantina, Chap. 6 white book COST/NewCompStar

N.B. either using non-relativistic functional or RMF models





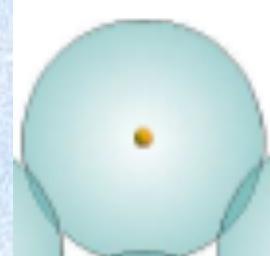
# Method in a nutshell (1)

**Start from an EoS model with SNA and add NSE perturbatively**

1. Start from the free energy density in a WS cell for an arbitrary EoS using SNA

$$F_{\text{WS}} = \underbrace{F_{\text{nuc}} + F_g}_{F_{\text{bar}}} + F_{\text{lept}} + F_{\gamma}$$

e.g.  $F_{\text{nuc}}^{(j)} = F_{\text{nuc}}^{(j)}(A, Z, n_g, V_C)$



the total free energy of the system being

$$F_{\text{tot}} = \sum_j N_j(k) F_{\text{WS}}^{(j)} \quad \langle N_j \rangle_k \longleftrightarrow \langle p_j \rangle_k \equiv p_j$$



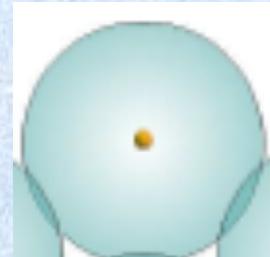
# Method in a nutshell (2)

**Start from an EoS model with SNA and add NSE perturbatively**

1. Start from the free energy density in a WS cell for an arbitrary EoS using SNA

$$F_{\text{WS}} = \underbrace{F_{\text{nuc}} + F_g}_{F_{\text{bar}}} + F_{\text{lept}} + F_{\gamma}$$

e.g.  $F_{\text{nuc}}^{(j)} = F_{\text{nuc}}^{(j)}(A, Z, n_g, V_C)$



the total free energy of the system being

$$F_{\text{tot}} = \sum_j N_j(k) F_{\text{WS}}^{(j)} \quad \langle N_j \rangle_k \longleftrightarrow \langle p_j \rangle_k \equiv p_j$$

2. Calculate the Gibbs free energy (Legendre transformation) working in a grand-canonical ensemble

$$G_{\text{nuc}}^{(j)} = F_{\text{nuc}}^{(j)} - \mu_n N^{(j)} - \mu_p Z^{(j)}$$

**N.B.** :  $F$  depends on the particular model employed

1. If  $F_{\text{nuc}}$  depends on the density (as in CLDM or EDF)  $\rightarrow$  *rearrangement term* must be added !
2. In-medium effects (e.g. excluded volume) have to be accounted for !
3. the gas is kept as in the original model  $\rightarrow$  only “non-uniform” part is changed



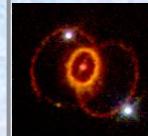
# Method in a nutshell (3)

Start from an EoS model with SNA and add NSE perturbatively

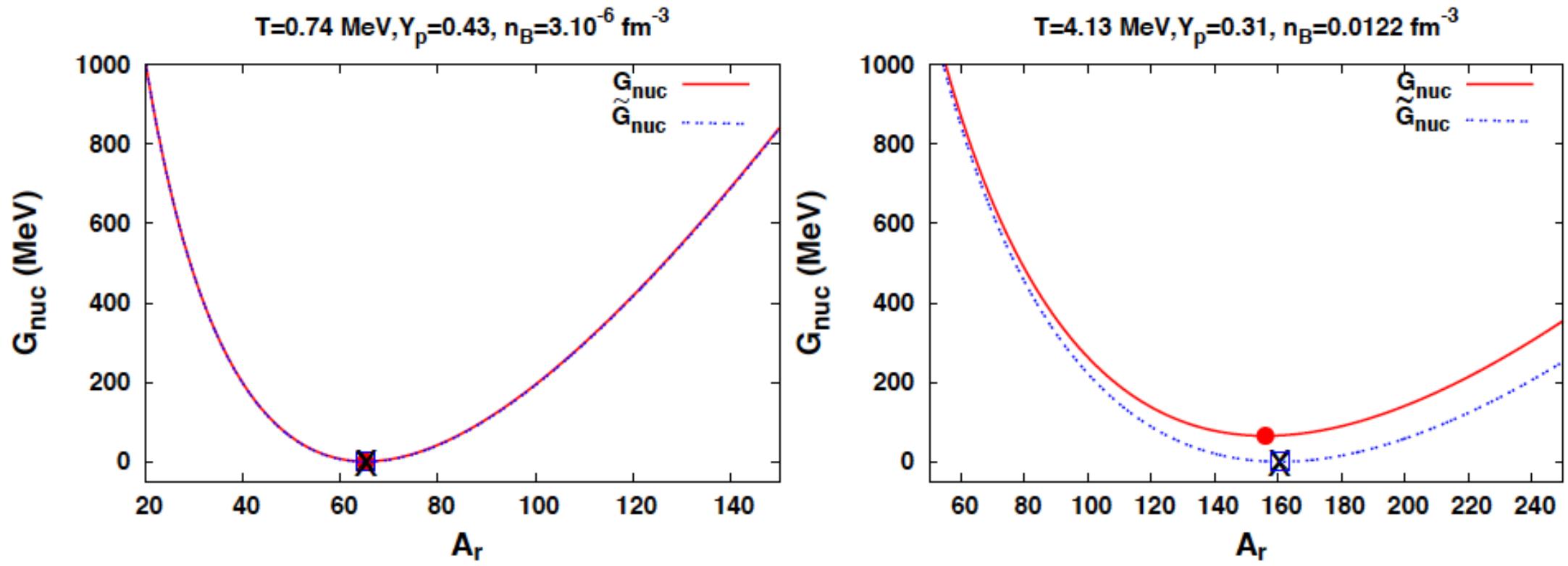
1. Start from the free energy density in a WS cell for an arbitrary EoS using SNA
2. Calculate the Gibbs free energy (Legendre transformation)  
working in a grand-canonical ensemble
3. Construct the probability associated to a nucleus ( $A, Z$ ) :

$$p_j = \frac{\langle N_j \rangle_k}{\sum_j \langle N_j \rangle_k} = \frac{\exp(-\tilde{G}_{\text{nuc}}^{(j)}/k_B T)}{\sum_j \exp(-\tilde{G}_{\text{nuc}}^{(j)}/k_B T)} \propto \exp(-\tilde{G}_{\text{nuc}}^{(j)}/k_B T) \rightarrow p(A, Z)$$

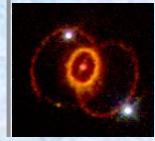
- ✓ method is general
- ✓ masses are input in the model → can be changed
- ✓ we apply it to the Lattimer & Swesty (LS) EoS (based on CLDM) widely used in SN simulations  
(Lattimer & Swesty, Nucl. Phys. A 535, 331 (1991))



# Application to the LS EoS (1)

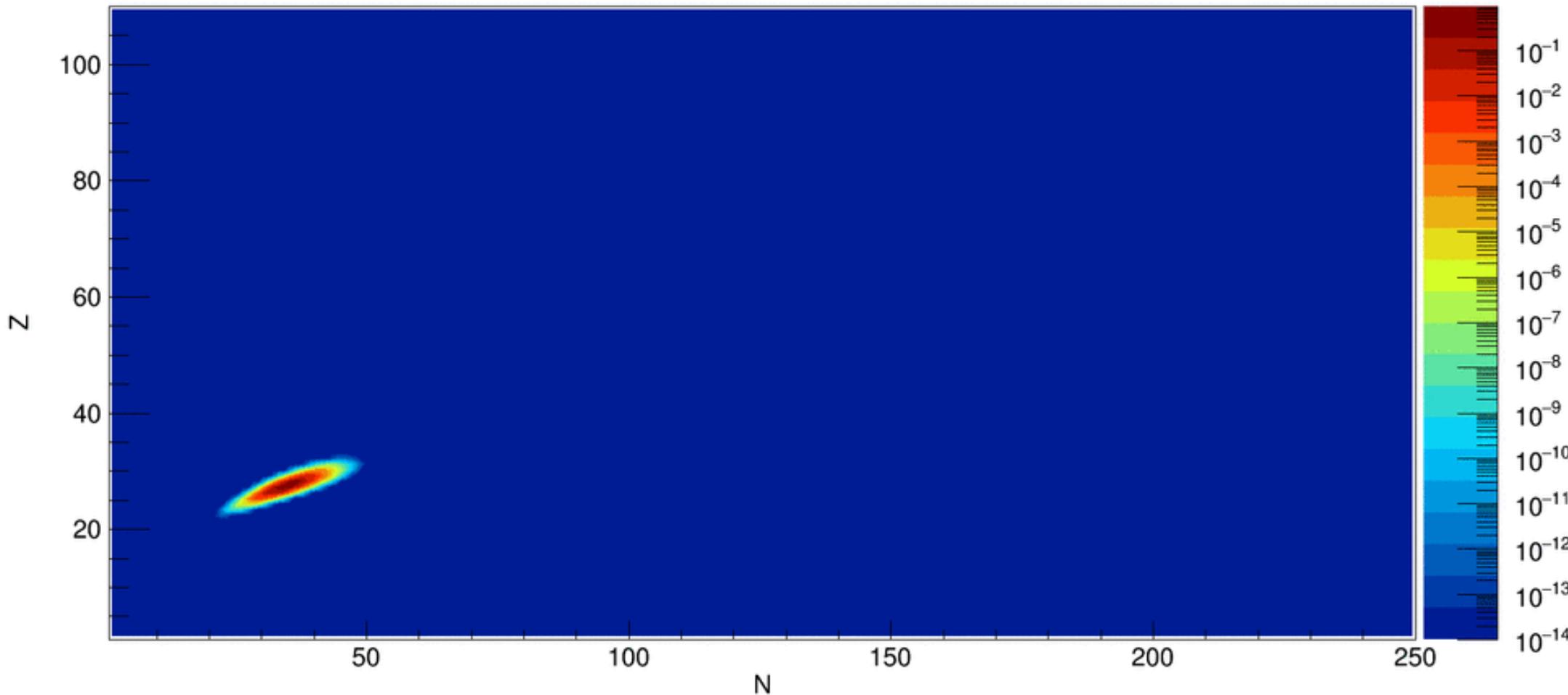


→ only including the rearrangement term allows to correctly reproduce the SNA result insuring thermodynamic consistency



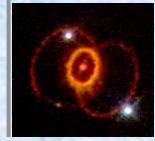
# Application to the LS EoS (2)

$T=0.665582$   $Y_p=0.442719$   $n_b=0.000002$

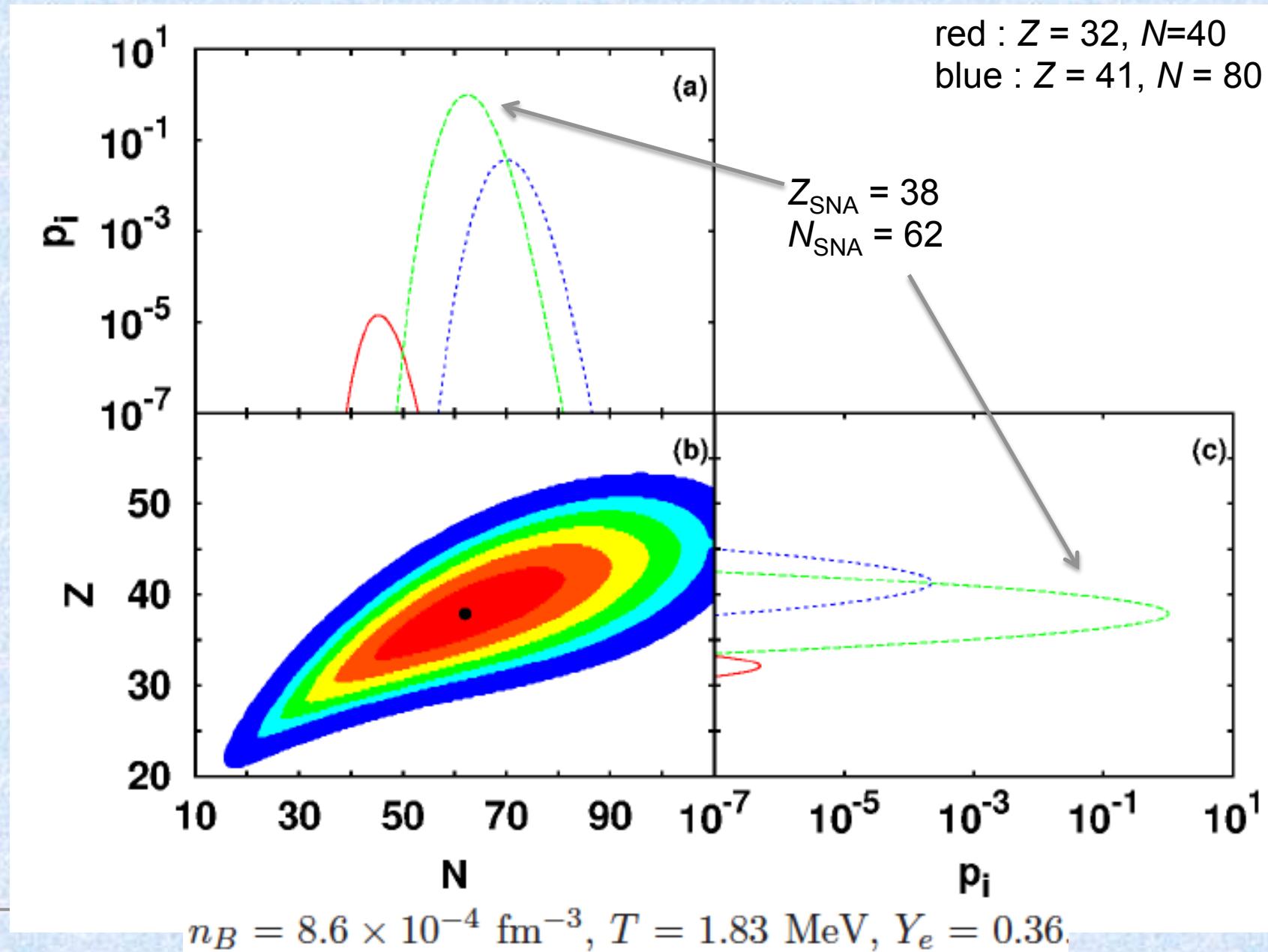


Thermo conditions fixed from a CCSN trajectory – 40 solar mass progenitor (A. Fantina, PhD thesis 2010)

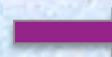
A. F. Fantina Movie by S. Giraud (PhD, GANIL)

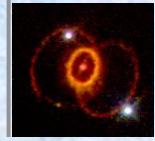


# Application to the LS EoS (3)



Grams, Giraud, Fantina, Gulminelli, PRC 97, 035807 (2018)





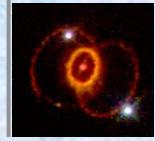
# Application to the LS EoS (4)

- ✓ Nuclear distribution centred on the most probable which coincides with the SNA → consistency
- ✓ Nuclei bigger and distribution larger as  $n_B, T$  increase
- ✓ but : No shell or pairing effects  
→ original LS model is simplified !  
(compressible liquid drop with simplified density-dependent nucleon-nucleon interaction)



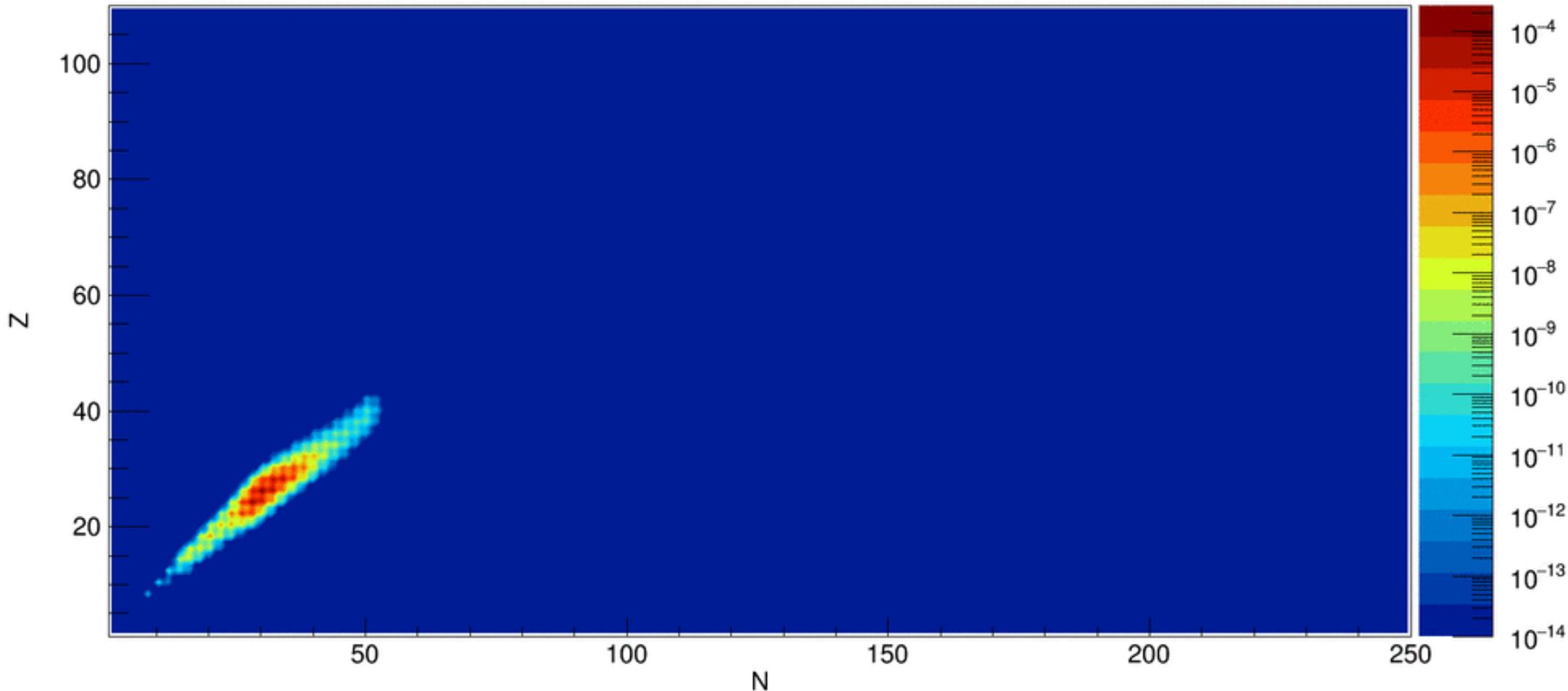
need of experimental masses or theoretical mass models !

- **More microscopic realistic mass models** → we use HFB-24 ([Goriely et al. 2013, PRC88, 024308](#))
  - HFB calculations with 16-parameter generalised Skyrme NN interaction with realistic pairing force
  - fit the 2353 experimental masses in AME2012 with  $Z, N > 8$  with a rms of about 0.5 MeV



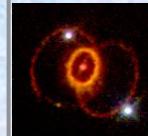
# Application to the LS EoS (5)

$T=0.665582$   $Y_p=0.442719$   $nb=0.000002$



Thermo conditions fixed from a CCSN trajectory – 40 solar mass progenitor (A. Fantina, PhD thesis 2010)

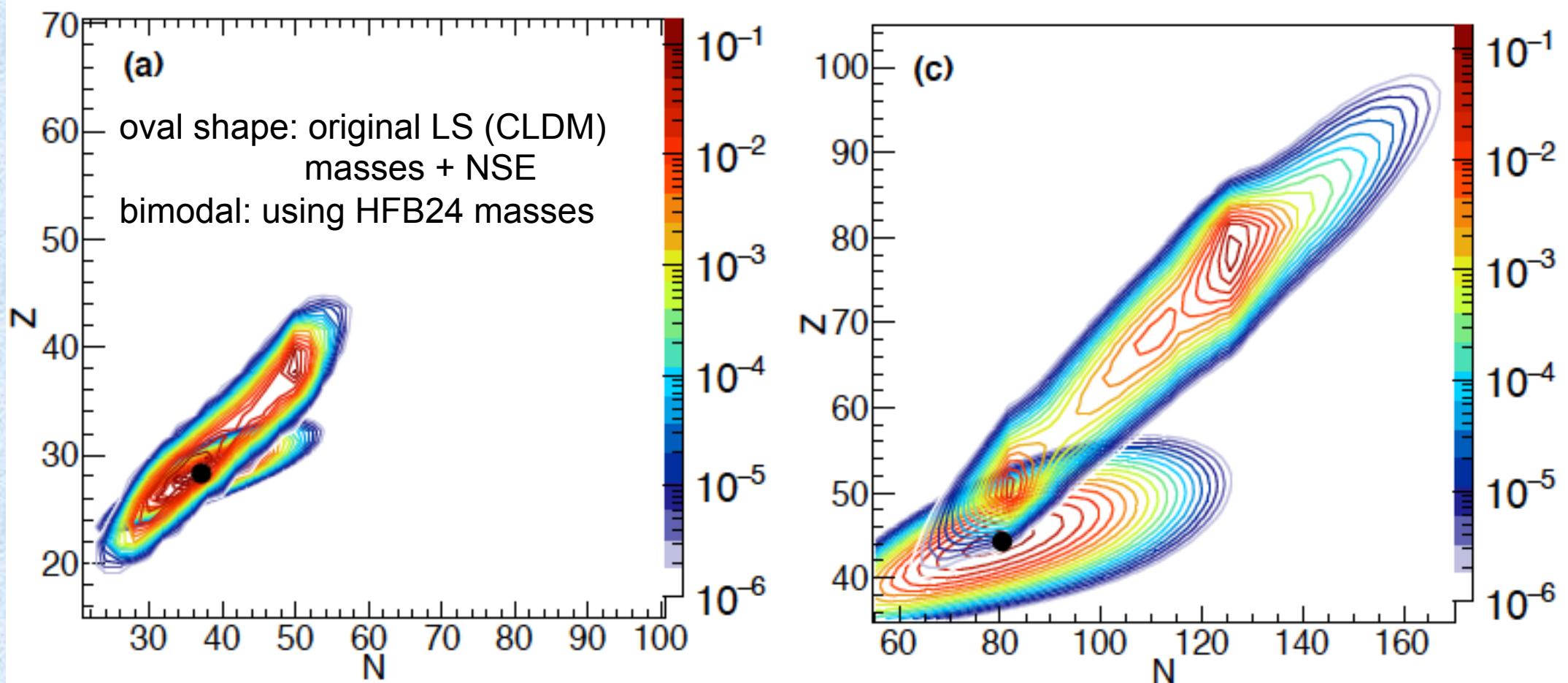
A. F. Fantina Movie by S. Giraud (PhD, GANIL)



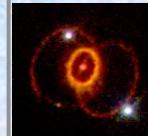
# Application to the LS EoS (6)

$n_B = 3.86 \times 10^{-6} \text{ fm}^{-3}$ ,  $T = 0.79 \text{ MeV}$ ,  $Y_e = 0.43$

$n_B = 3.01 \times 10^{-3} \text{ fm}^{-3}$ ,  $T = 2.68 \text{ MeV}$ ,  $Y_e = 0.33$

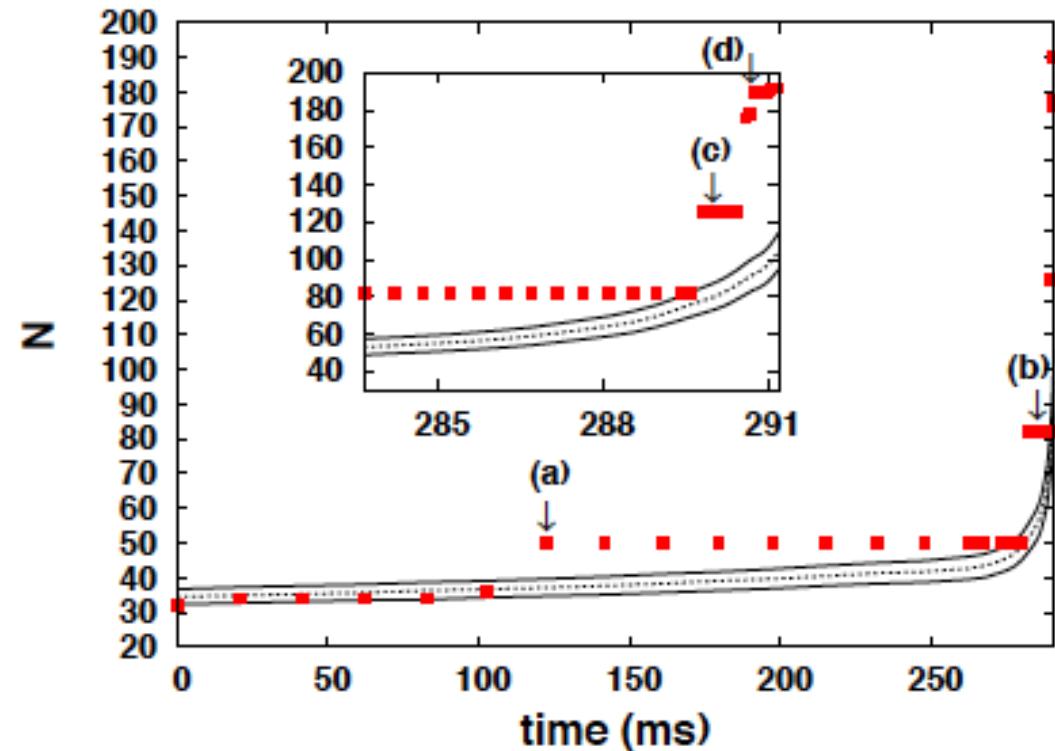
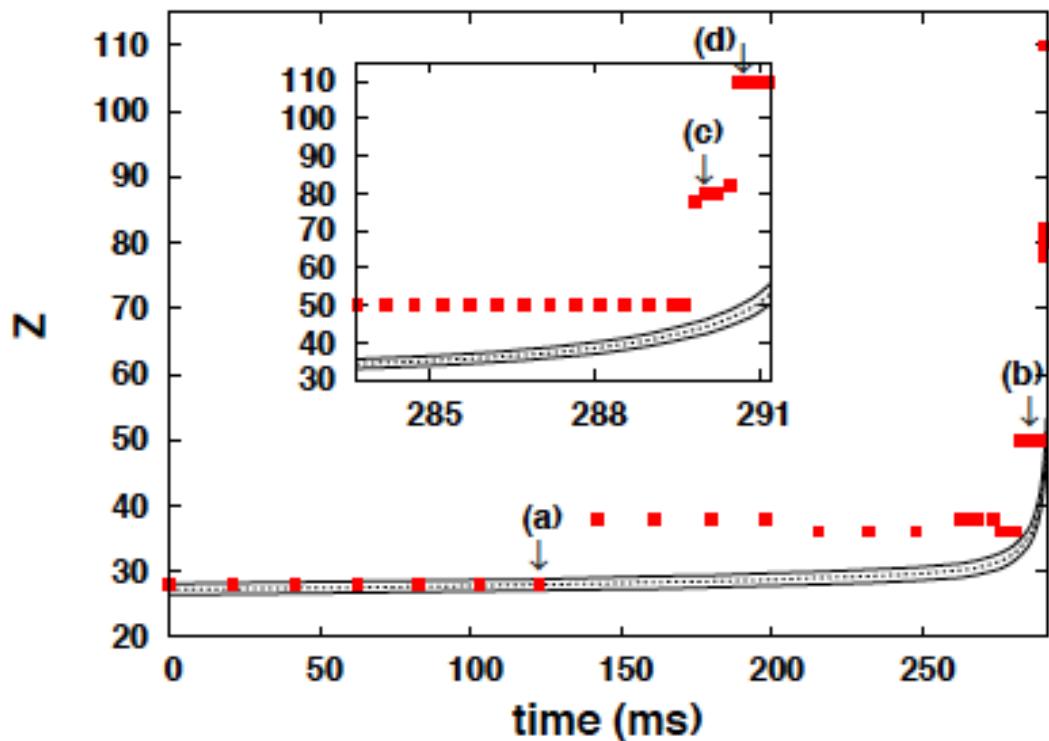


- **Bimodal distribution around magic nuclei**
- **in late stage of CC predictions for the distribution differ considerably**  
( $G$  and  $\mu_i$  should be calculated with same model)



# Application to the LS EoS (7)

Solid lines : original Lattimer & Swesty EoS model (compressible liquid drop model, no shell effects)  
Dots : original masses replaced by HFB-24 mass model (shell and pairing included)



→ Magic nuclei dominate trajectory when realistic mass model employed  
→ might impact electron-capture rates thus the collapse dynamics

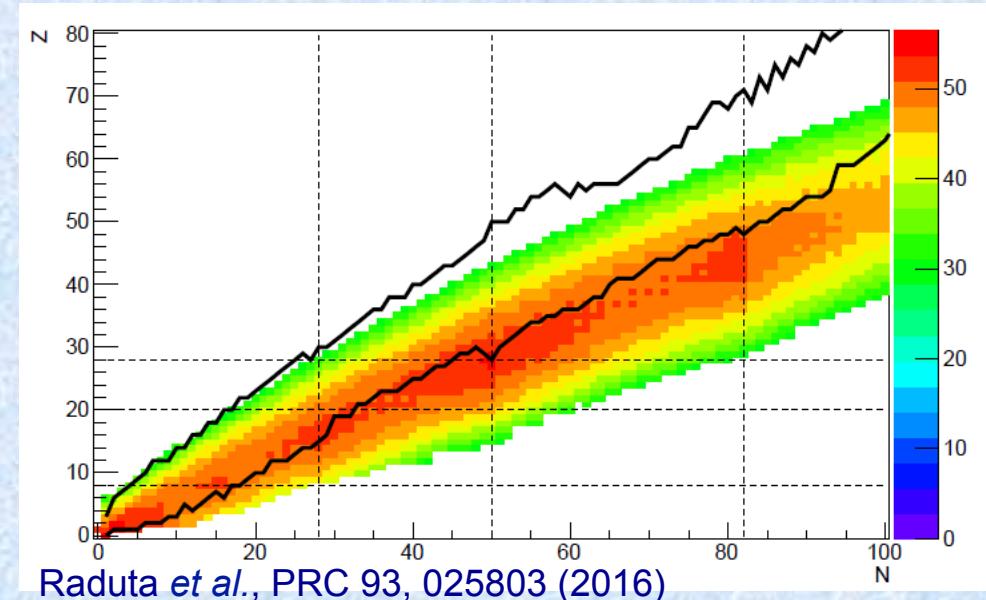
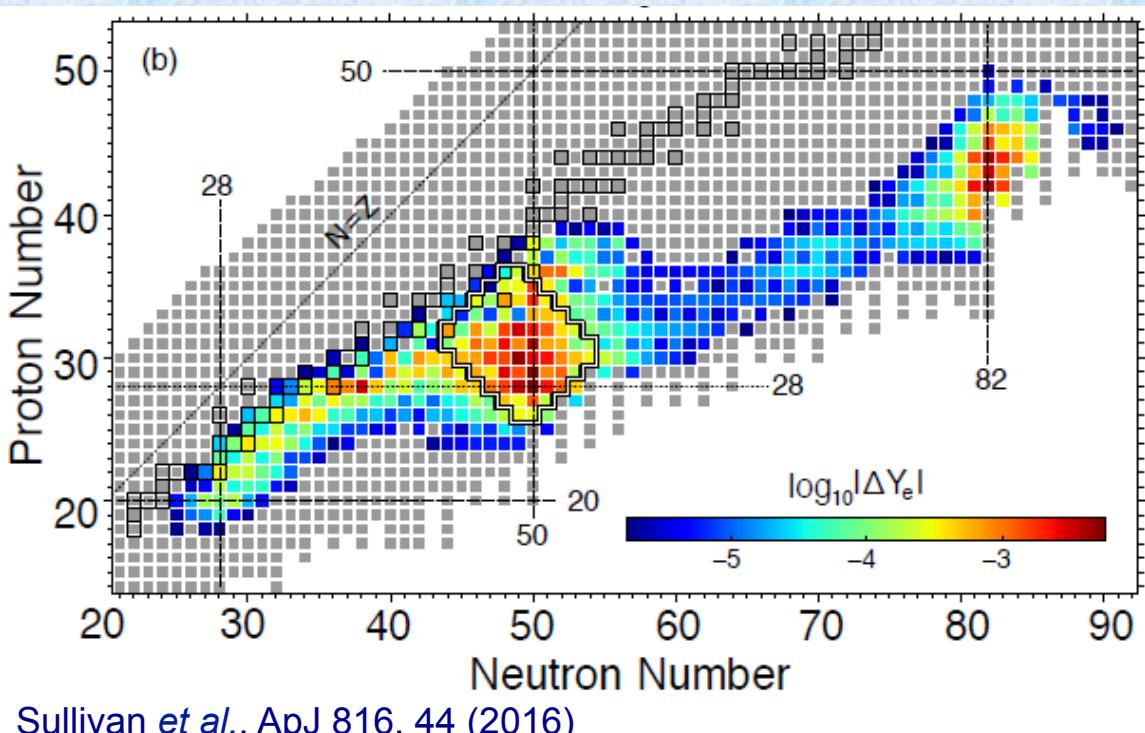


# Mass measurements & SN (1)

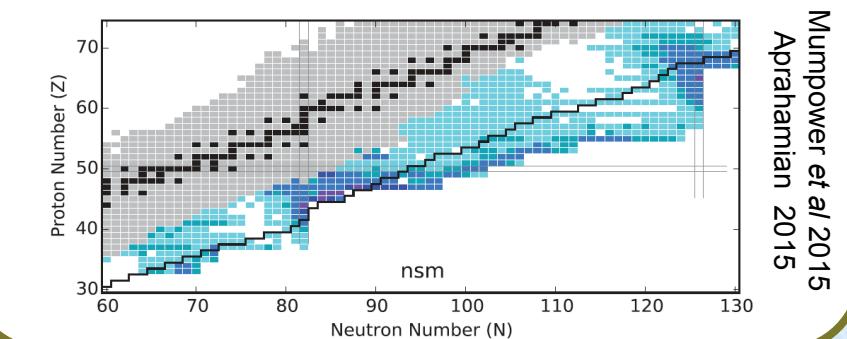
❖ Very precise mass measurements (within  $\sim 100$  KeV) are necessary for the computation of :

- ✓ Q value in EC rates
- ✓ EoS – composition of SN matter

❖ Which nuclei matter ?



**NOTE :** the **N=82** closed shell nuclei are **relevant for both core collapse and for r-process**.



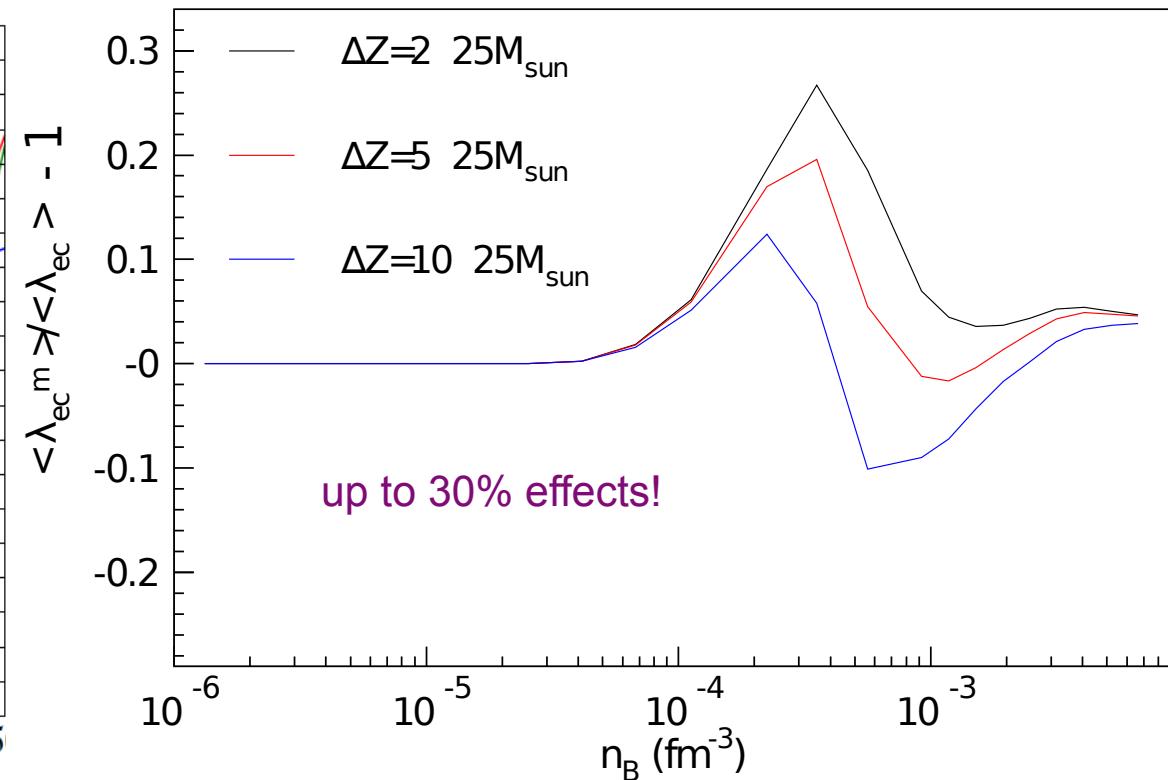
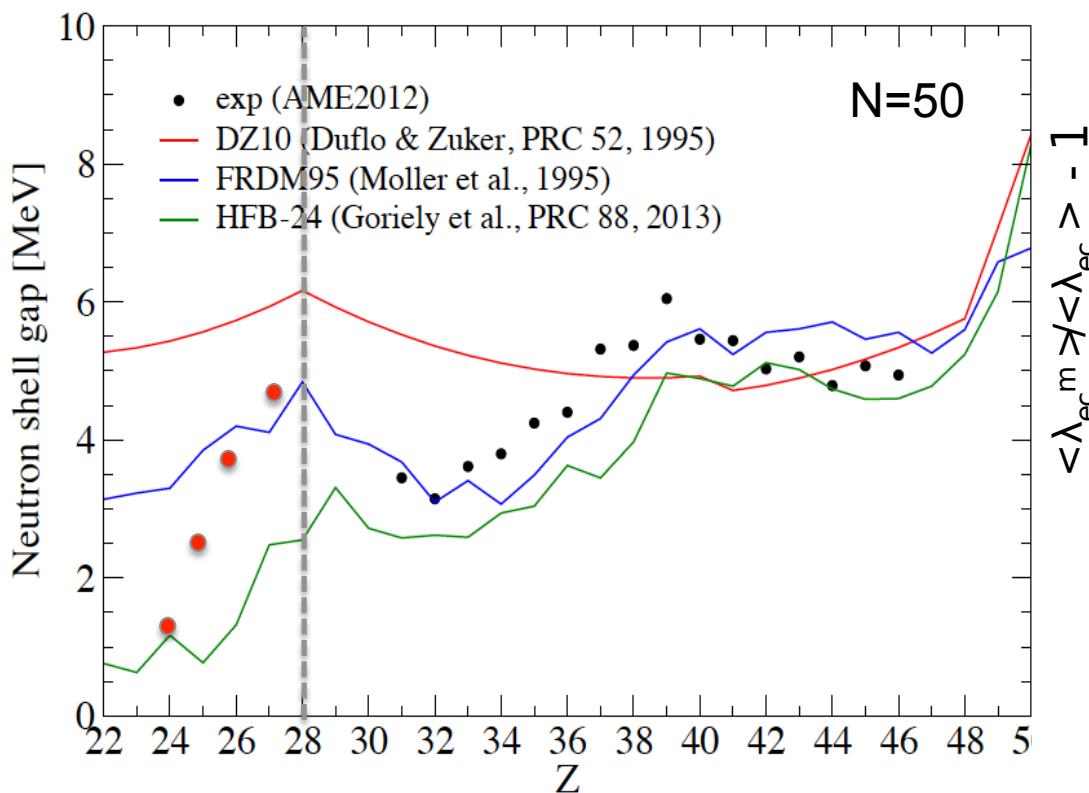


# Mass measurements & SN (2)

- ❖ Exotic nuclei around N=50 ( $^{78}\text{Ni}$ ) and N=82 ( $^{128}\text{Pd}$ ) dominate because predicted to be magic.

→ What about more exotic nuclei ?!

e.g.  $N = 50$  gap far from stability → model predictions differ !



Raduta, Gulminelli, Oertel, PRC 93, 025803 (2016),  
using DZ10 with quenching function for exotic nuclei

→ in turn, impact on CCSN dynamics, neutrino luminosity, ...



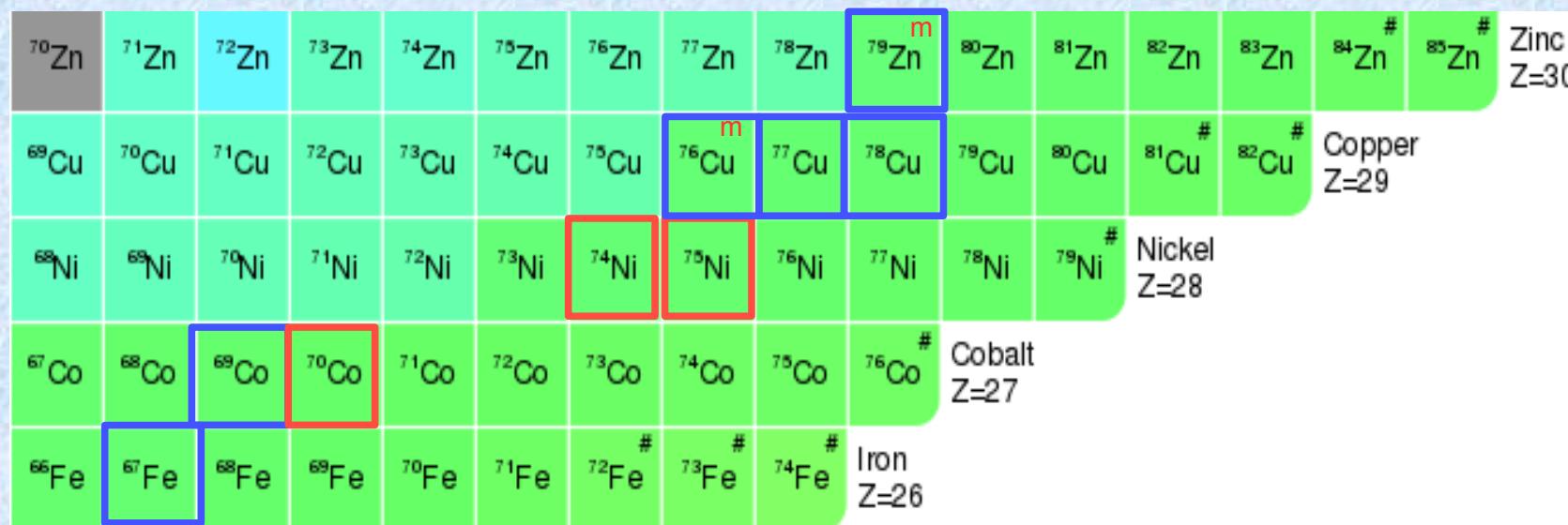
# Mass measurements & SN (3)

→ experimental proposal @JYFL : mass measurements around  $^{78}\text{Ni}$  (run 11/2017, under analysis)

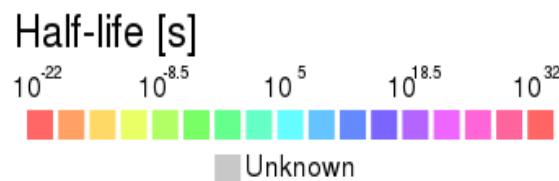
Spokespersons : B. Bastin, A. Kankainen (exp) ; A. F. Fantina, F. Gulminelli (theo)

Collaboration : GANIL, JYFL, LPC Caen, IFIN-HH, LUTh, IPNL, CENBG, IPNO, KU Leuven  
(PhD thesis of S. Giraud, 2016-2019)

IGISOL technique + penning traps & MR-ToF-MS



Courtesy of B. Bastin, S. Giraud (GANIL)



- This work**
  - Improved the precision
  - New masses
  - m New isomer masses



# Conclusions & Outlook

- ❖ **Formalism to include cluster distribution in a SNA EoS**
- ❖ **CCSN** dynamics sensitive to microphysics (EoS & EC rates)
  - **Masses** of n-rich nuclei ( $^{78}\text{Ni}$ ,  $^{128}\text{Pd} \rightarrow N = 50, 82$ )
  - **Magic nuclei dominate trajectory when realistic mass model employed**
    - might impact electron-capture rates thus the collapse dynamics



- ✓ need of (microscopic & reliable) theoretical model when no data
- ✓ need of experimental data to calibrate the models

- Implementation of new formalism in CCSN simulation (in progress)
- Check impact on EC (in progress)
- Impact of new results of **I220** experiment → test impact of new masses

Obrigada!