

Effective models for two-component scalar dark matter with Z_8 and Z_{10} symmetries

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Two-component scalar dark matter charged under Z_{2n} :

- Complex field $\rightarrow \phi_A$
- Real field $\rightarrow \phi_B$

Associated charges: $\omega_{2n}^A = \exp(i\pi A/n)$ and $\omega_{2n}^B = \exp(i\pi)$.

- Previous work [Yaguna and Zapata, 2021] on similar scenarios.
- Potentials for Z_8 and Z_{10} either lack new terms or are redundant.
- We extend Z_8 and Z_{10} scenarios to include effective operators of up to energy dimension 6.

Sym2Int

[Fonseca, 2017]

Operators

FeynRules

[Alloul et al., 2014]

Feynman Rules

micrOMEGAs

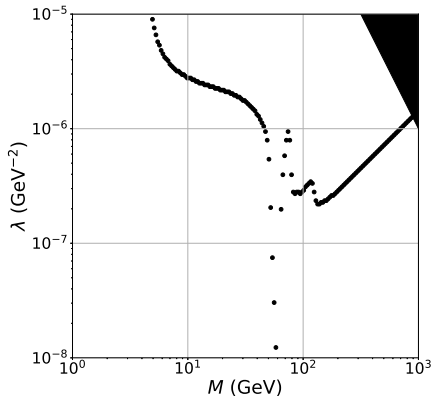
[Alguero et al., 2024]

Dark Matter
observables

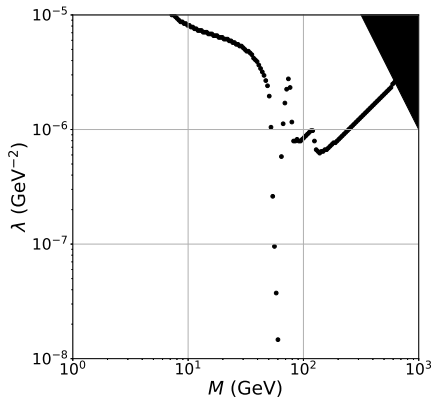
Single component case

$$\mathcal{L}_{DM} \supset M^2 |\phi_{A/B}|^2 / 2 + \lambda |H|^4 |\phi_{A/B}|^2$$

Real:

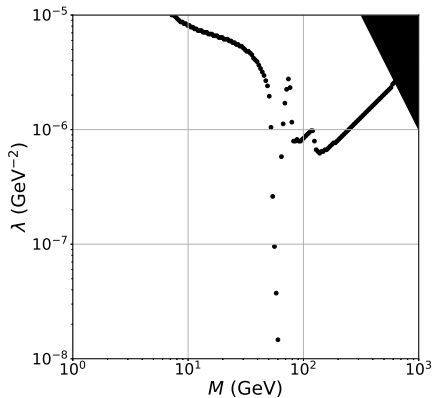
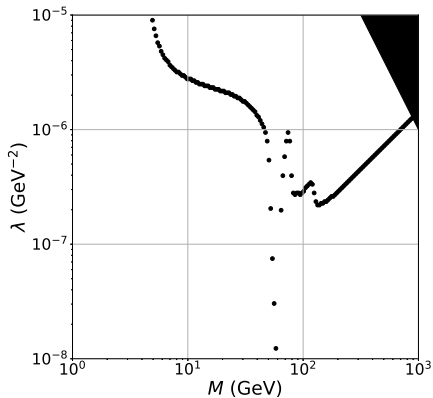


Complex:



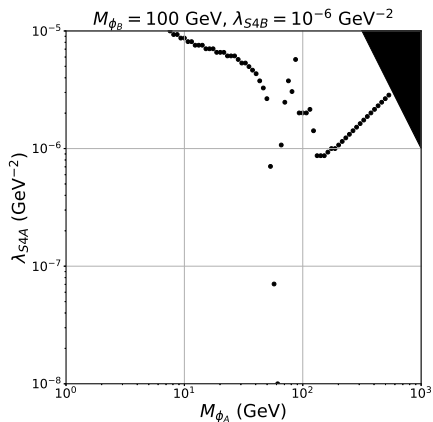
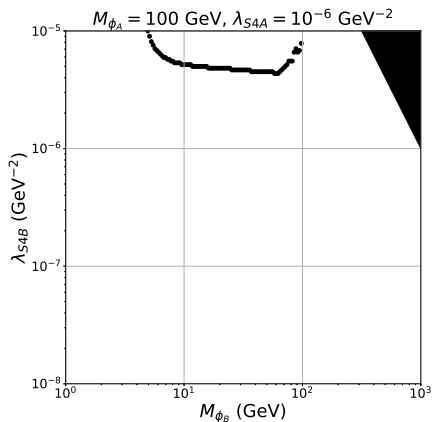
Single component case

Not completely new [Criado et al., 2021]



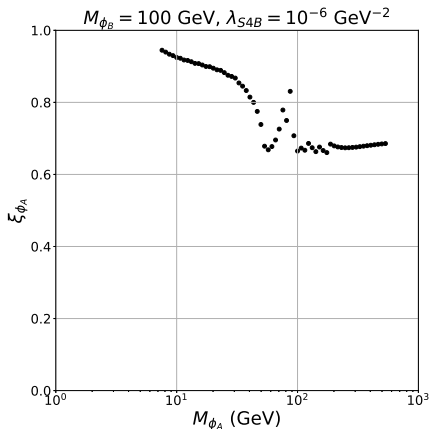
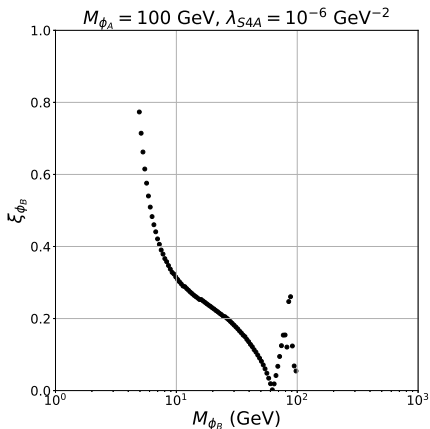
Two-component cases

$$M_{\phi_A}^2 |\phi_A|^2 + \frac{1}{2} M_{\phi_B}^2 \phi_B^2 + \lambda_{S4A} |H|^4 |\phi_A|^2 + \lambda_{S4B} |H|^4 \phi_B^2$$



Two-component cases

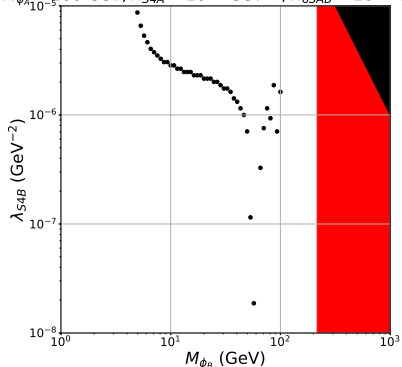
$$\xi_{\phi_i} = \Omega_{\phi_i} / \Omega_{DM}$$



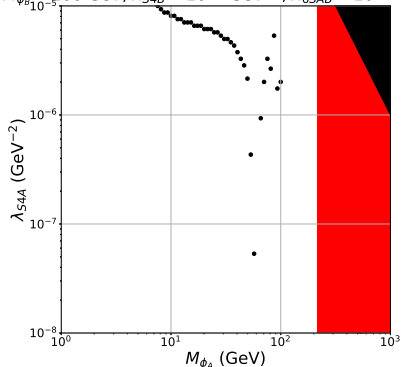
Two-component cases

$$M_{\phi_A}^2 |\phi_A|^2 + \frac{1}{2} M_{\phi_B}^2 \phi_B^2 + \lambda_{S4A} |H|^4 |\phi_A|^2 + \lambda_{S4B} |H|^4 \phi_B^2 + \lambda_{6SAB} |H|^2 |\phi_A|^2 \phi_B^2$$

$M_{\phi_A} = 100 \text{ GeV}, \lambda_{S4A} = 10^{-6} \text{ GeV}^{-2}, \lambda_{6SAB} = 10^{-5} \text{ GeV}^{-2}$



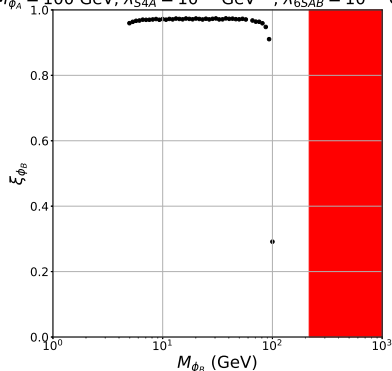
$M_{\phi_B} = 100 \text{ GeV}, \lambda_{S4B} = 10^{-6} \text{ GeV}^{-2}, \lambda_{6SAB} = 10^{-5} \text{ GeV}^{-2}$



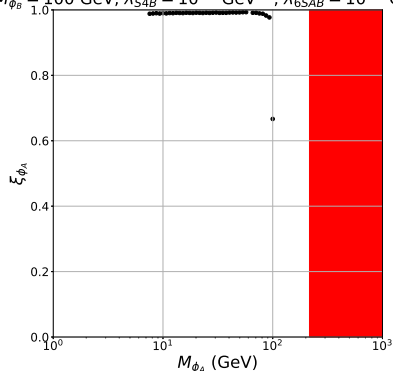
Two-component cases

$$M_{\phi_A}^2 |\phi_A|^2 + \frac{1}{2} M_{\phi_B}^2 \phi_B^2 + \lambda_{S4A} |H|^4 |\phi_A|^2 + \lambda_{S4B} |H|^4 \phi_B^2 + \lambda_{6SAB} |H|^2 |\phi_A|^2 \phi_B^2$$

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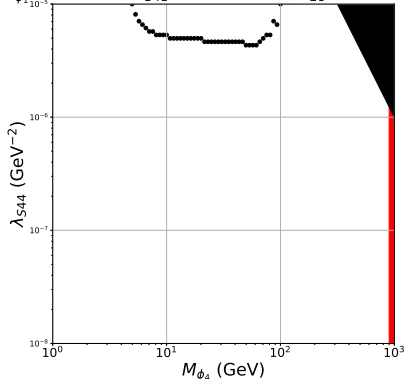
$M_{\phi_B} = 100 \text{ GeV}, \lambda_{S4B} = 10^{-6} \text{ GeV}^{-2}, \lambda_{6SAB} = 10^{-5} \text{ GeV}^{-2}$



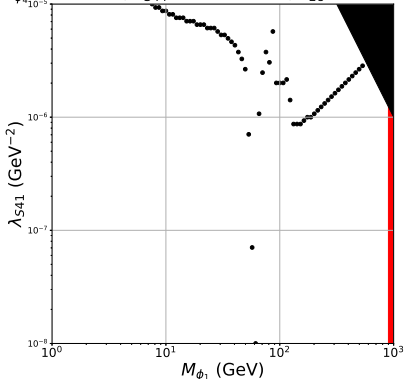
Two-component cases

$$Z_8 : M_{\phi_1}^2 |\phi_1|^2 + \frac{1}{2} M_{\phi_4}^2 \phi_4^2 + \lambda_{S41} |H|^4 |\phi_1|^2 + \lambda_{S44} |H|^4 \phi_4^2 + \lambda_{E8} \phi_1^4 \phi_4 + \text{h.c.}$$

$M_{\phi_1} = 100 \text{ GeV}, \lambda_{S41} = 10^{-6} \text{ GeV}^{-2}, \lambda_{E8} = 10^{-3} \text{ GeV}^{-1}$



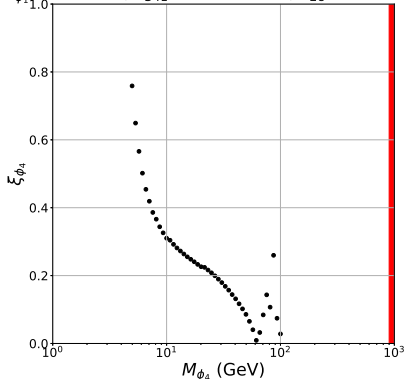
$M_{\phi_4} = 100 \text{ GeV}, \lambda_{S44} = 10^{-6} \text{ GeV}^{-2}, \lambda_{E8} = 10^{-3} \text{ GeV}^{-1}$



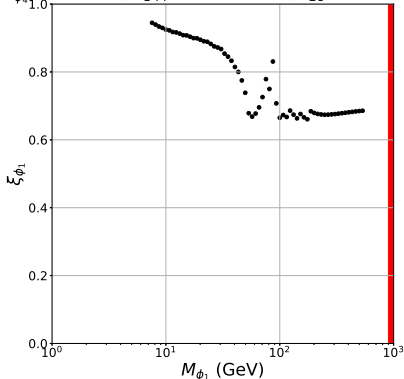
Two-component cases

$$Z_8 : M_{\phi_1}^2 |\phi_1|^2 + \frac{1}{2} M_{\phi_4}^2 \phi_4^2 + \lambda_{S41} |H|^4 |\phi_1|^2 + \lambda_{S44} |H|^4 \phi_4^2 + \lambda_{E8} \phi_1^4 \phi_4 + \text{h.c.}$$

$M_{\phi_1} = 100 \text{ GeV}, \lambda_{S41} = 10^{-6} \text{ GeV}^{-2}, \lambda_{E8} = 10^{-3} \text{ GeV}^{-1}$



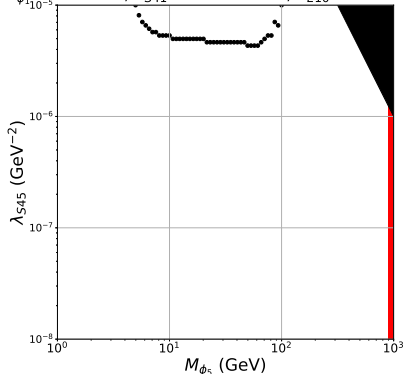
$M_{\phi_4} = 100 \text{ GeV}, \lambda_{S44} = 10^{-6} \text{ GeV}^{-2}, \lambda_{E8} = 10^{-3} \text{ GeV}^{-1}$



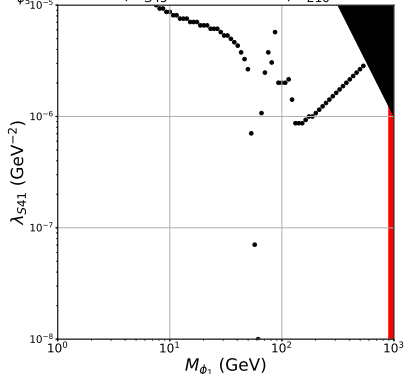
Two-component cases

$$Z_{10} : M_{\phi_1}^2 |\phi_1|^2 + \frac{1}{2} M_{\phi_5}^2 \phi_5^2 + \lambda_{S41} |H|^4 |\phi_1|^2 + \lambda_{S45} |H|^4 \phi_5^2 + \lambda_{E10} \phi_1^5 \phi_5 + \text{h.c.}$$

$M_{\phi_1} = 100 \text{ GeV}, \lambda_{S41} = 10^{-6} \text{ GeV}^{-2}, \lambda_{E10} = 10^{-6} \text{ GeV}^{-2}$



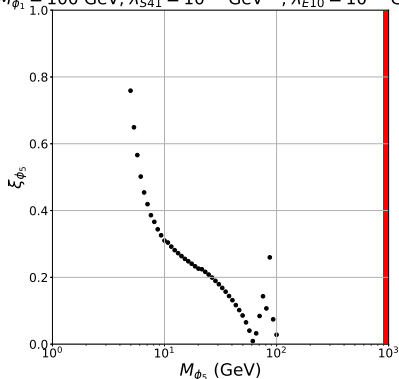
$M_{\phi_5} = 100 \text{ GeV}, \lambda_{S45} = 10^{-6} \text{ GeV}^{-2}, \lambda_{E10} = 10^{-6} \text{ GeV}^{-2}$



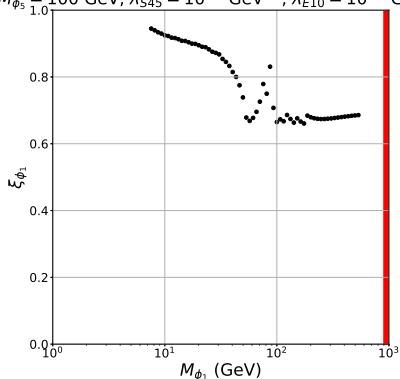
Two-component cases

$$Z_{10} : M_{\phi_1}^2 |\phi_1|^2 + \frac{1}{2} M_{\phi_5}^2 \phi_5^2 + \lambda_{S41} |H|^4 |\phi_1|^2 + \lambda_{S45} |H|^4 \phi_5^2 + \lambda_{E10} \phi_1^5 \phi_5 + \text{h.c.}$$

$M_{\phi_1} = 100 \text{ GeV}, \lambda_{S41} = 10^{-6} \text{ GeV}^{-2}, \lambda_{E10} = 10^{-6} \text{ GeV}^{-2}$








$M_{\phi_5} = 100 \text{ GeV}, \lambda_{S45} = 10^{-6} \text{ GeV}^{-2}, \lambda_{E10} = 10^{-6} \text{ GeV}^{-2}$



Conclusions

- Generated correct dark matter abundances for single component real and complex DM with effective operators.
- Constructed a two-component scenario combining both real and complex fields.
- Effective operators such as $|H|^4|\phi_{A/B}|^2$ and $|H|^2|\phi_A|^2\phi_B^2$ change the outlook for generating correct DM abundance.
- Operators $\phi_1^4\phi_4$ and $\phi_1^5\phi_5$ do not seem to meaningfully affect the DM abundances but more work is required.

References

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The End

BU: Cut-Off energy of the EFT

For an effective operator of dimension D its coupling c :

$$c = \frac{1}{\Lambda^{D-4}} \Rightarrow \Lambda = c^{-1/(D-4)}$$

With E the typical energy scale of a DM process we say that the EFT breaks down at $E \sim \Lambda$.

For effective operators that involve one field we take the typical energy of a DM process as $E \sim M_{DM}$:

$$M_{DM} < c^{-1/(D-4)}$$

For effective operator involving both DM fields we take the sum of the masses instead:

$$\sum M_{DM} < c^{-1/(D-4)}$$

In principle for Z_8 we can have $A = 1, 2, 3$:

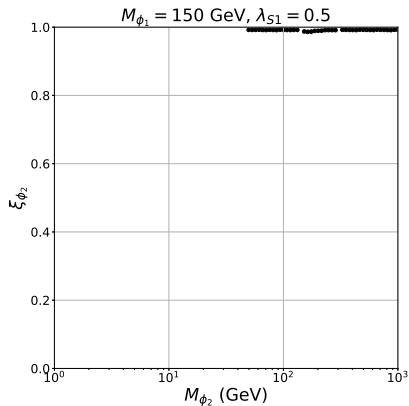
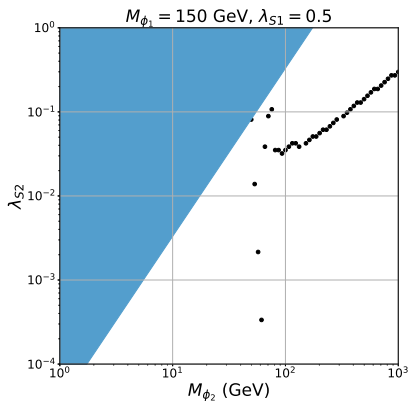
- $A = 2$ is analogous to $Z_4(12)$.
- $A = 1, 3$ allow the same operators up to dim 6 so we only show $Z_8(14)$ and drop the (14).

For Z_{10} it is possible to have $A = 1, 2, 3, 4$:

- $A = 1, 3$ allow the same operators up to dim 6 so we only show $Z_{10}(15)$ and drop the (15).
- $A = 2, 4$ allow for a ϕ_1^5 term we have not considered yet.

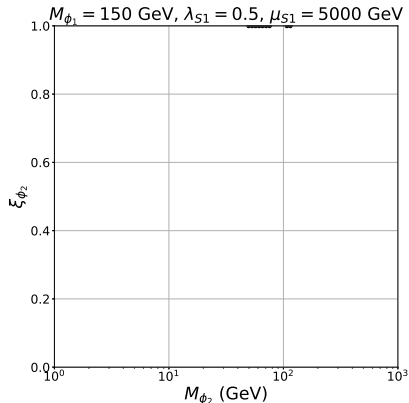
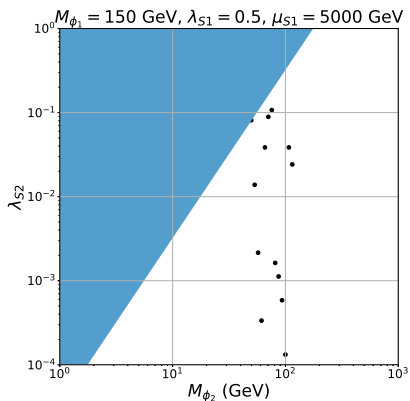
BU: Other Benchmark

$$\mu_A^2 |\phi_A|^2 + \frac{1}{2} \mu_B^2 \phi_B^2 + \lambda_{SA} |H|^2 |\phi_A|^2 + \lambda_{SB} |H|^2 \phi_B^2$$



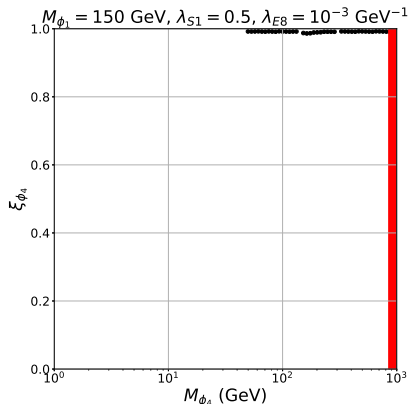
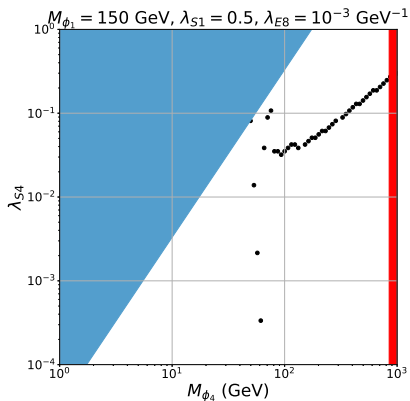
BU: Other Benchmark

$$Z_4 : \mu_1^2 |\phi_1|^2 + \frac{1}{2} \mu_2^2 \phi_2^2 + \lambda_{S1} |H|^2 |\phi_1|^2 + \lambda_{S2} |H|^2 \phi_2^2 + \mu_{S1} \phi_1^2 \phi_2 + h.c.$$



BU: Other Benchmark

$$Z_8 : \mu_1^2 |\phi_1|^2 + \frac{1}{2} \mu_4^2 \phi_4^2 + \lambda_{S1} |H|^2 |\phi_1|^2 + \lambda_{S4} |H|^2 \phi_4^2 + \lambda_{E8} \phi_1^4 \phi_4 + \text{h.c.}$$



BU: Other Benchmark

$$Z_{10} : \mu_1^2 |\phi_1|^2 + \frac{1}{2} \mu_5^2 \phi_5^2 + \lambda_{S1} |H|^2 |\phi_1|^2 + \lambda_{S5} |H|^2 \phi_5^2 + \lambda_{E10} \phi_1^5 \phi_5 + \text{h.c.}$$

