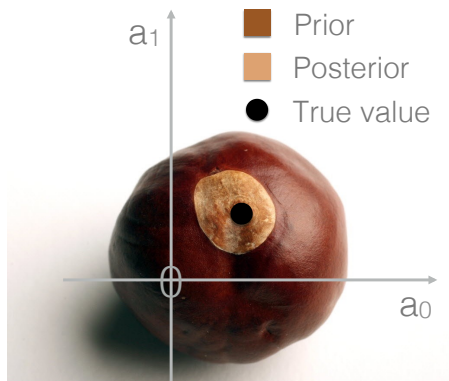


# Overview of Emulators for Nuclear Physics

Dick Furnstahl  
ISNET-9, May 2023



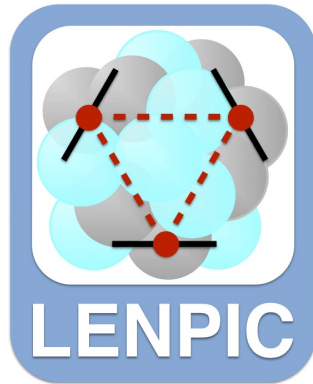
THE OHIO STATE UNIVERSITY



BUQEYE Collaboration

<https://buqeye.github.io/>

Python notebooks here!



<https://www.lenpic.org/>

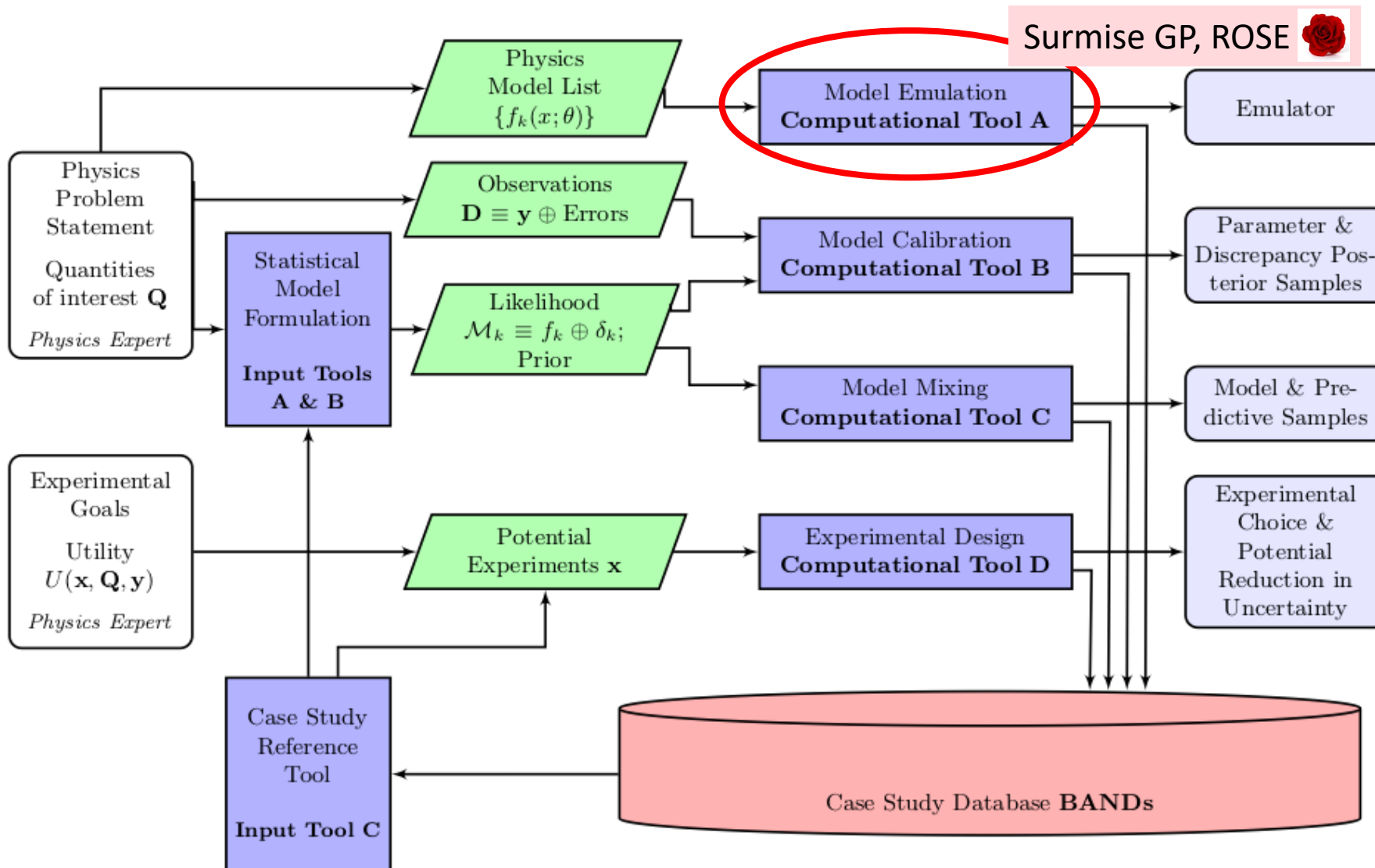
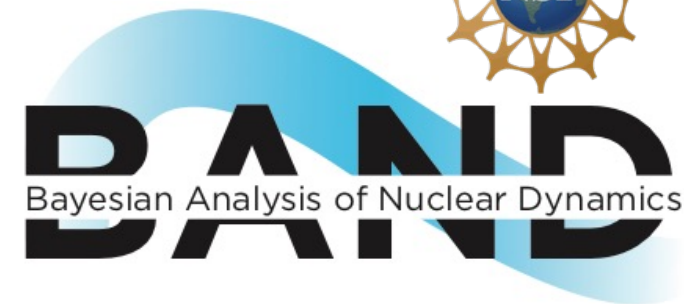
**NUCLEI**  
Nuclear Computational Low-Energy Initiative  
<https://nuclei.mps.ohio-state.edu/>

**BAND**  
Bayesian Analysis of Nuclear Dynamics  
<https://bandframework.github.io/>



# BAND (Bayesian Analysis of Nuclear Dynamics)

*Goal: Facilitate principled Uncertainty Quantification in Nuclear Physics*



**An NSF CSSI Framework  
(started 7/2020)**

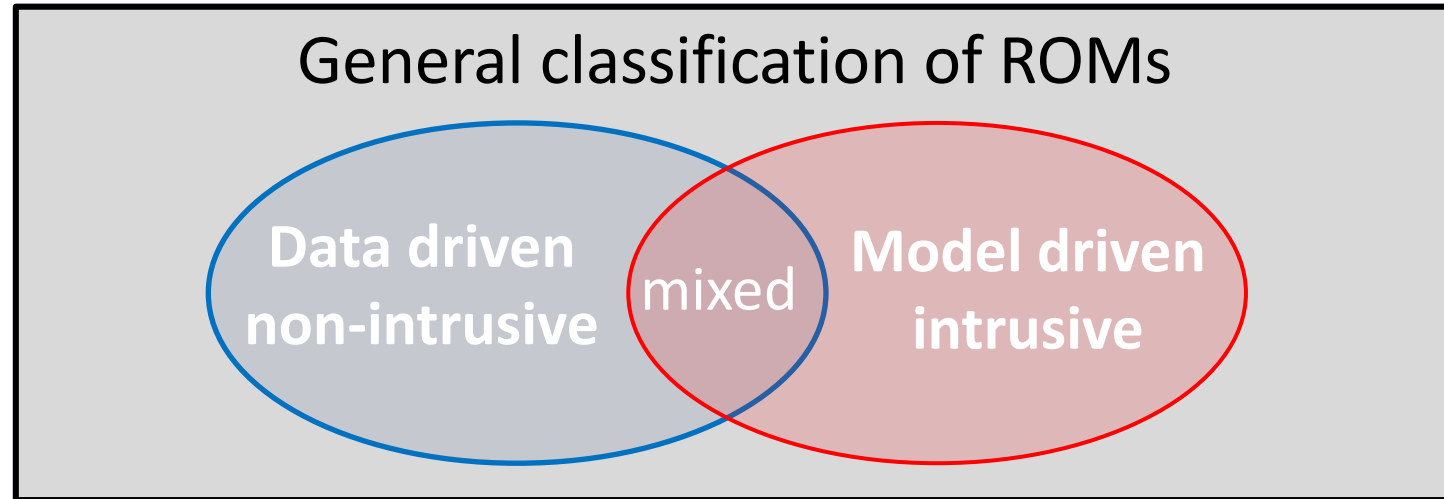
Look to  
<https://bandframework.github.io/> for papers,  
talks, and software!

# Model reduction methods → build nuclear emulators

**Need:** to vary parameters for design, control, optimization, UQ.

**Exploit:** much information in high-fidelity models is superfluous.

**Solution:** reduced-order model (ROM) → emulator (fast & accurate™).



**Data driven:** interpolate output of high-fidelity model w/o understanding → *non-intrusive*

Examples: Gaussian processes; dynamic mode decomposition; artificial neural network, also hybrid ML

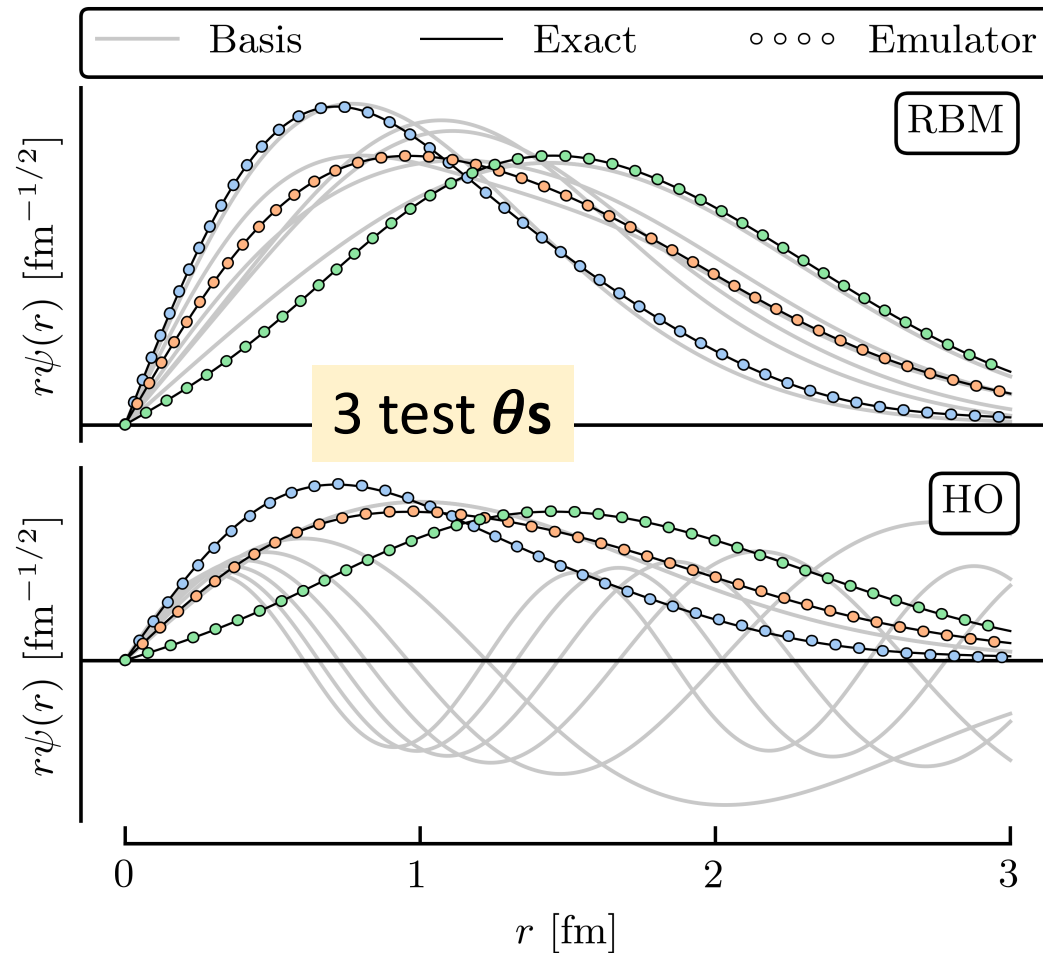
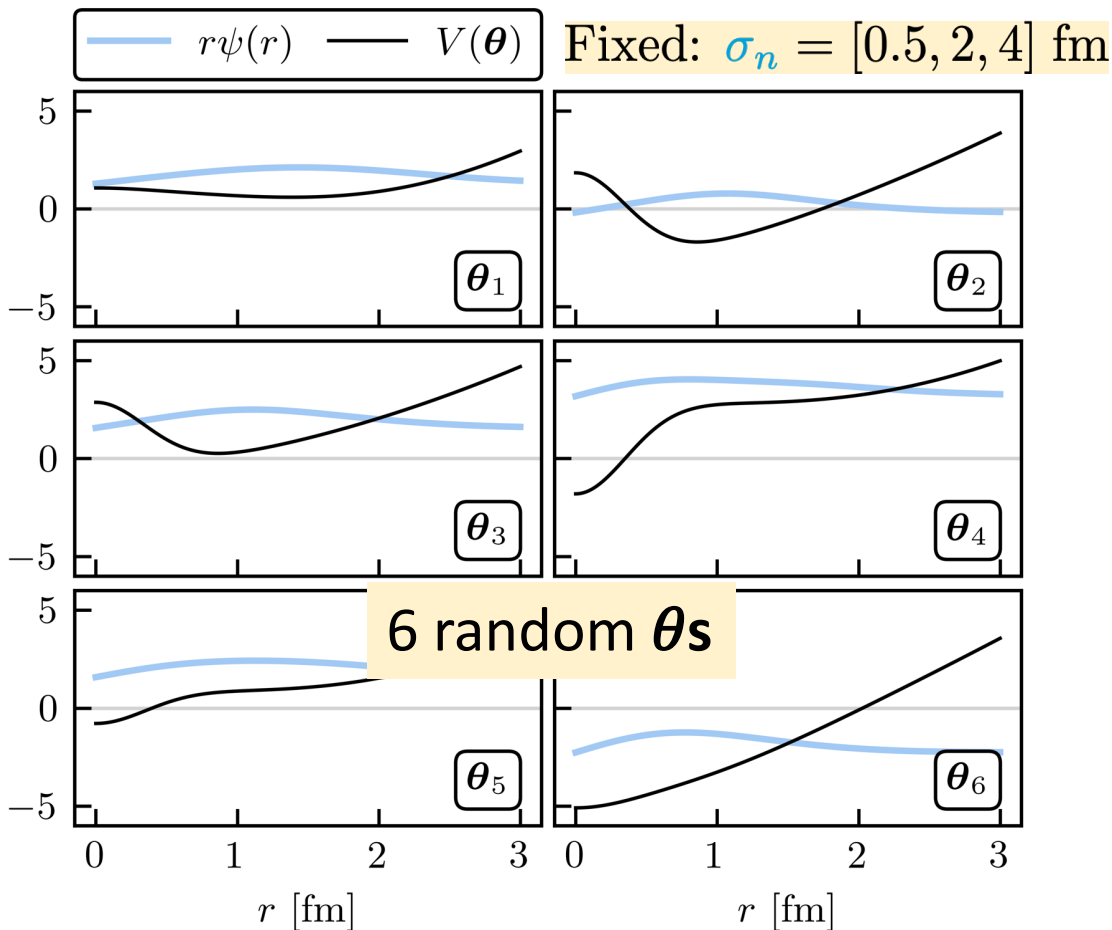
**Model driven:** derive reduced-order equations from high-fidelity equations → *intrusive*

Features: physics-based, respects underlying structure → can extrapolate; often uses projection

See [Melendez et al., 2022](#) for many references from the wide ROM literature; various types of emulators already successful in NP (e.g., refs. in [Drischler et al., 2022](#))

# Illustrative example: anharmonic oscillator [\[Try your own!\]](#)

Eigenvalue problem:  $H(\boldsymbol{\theta})|\psi\rangle = E|\psi\rangle$   $V(r; \boldsymbol{\theta}) = V_{\text{HO}}(r) + \sum_{n=1}^3 \theta^{(n)} e^{-r^2/\sigma_n^2}$  ← affine!

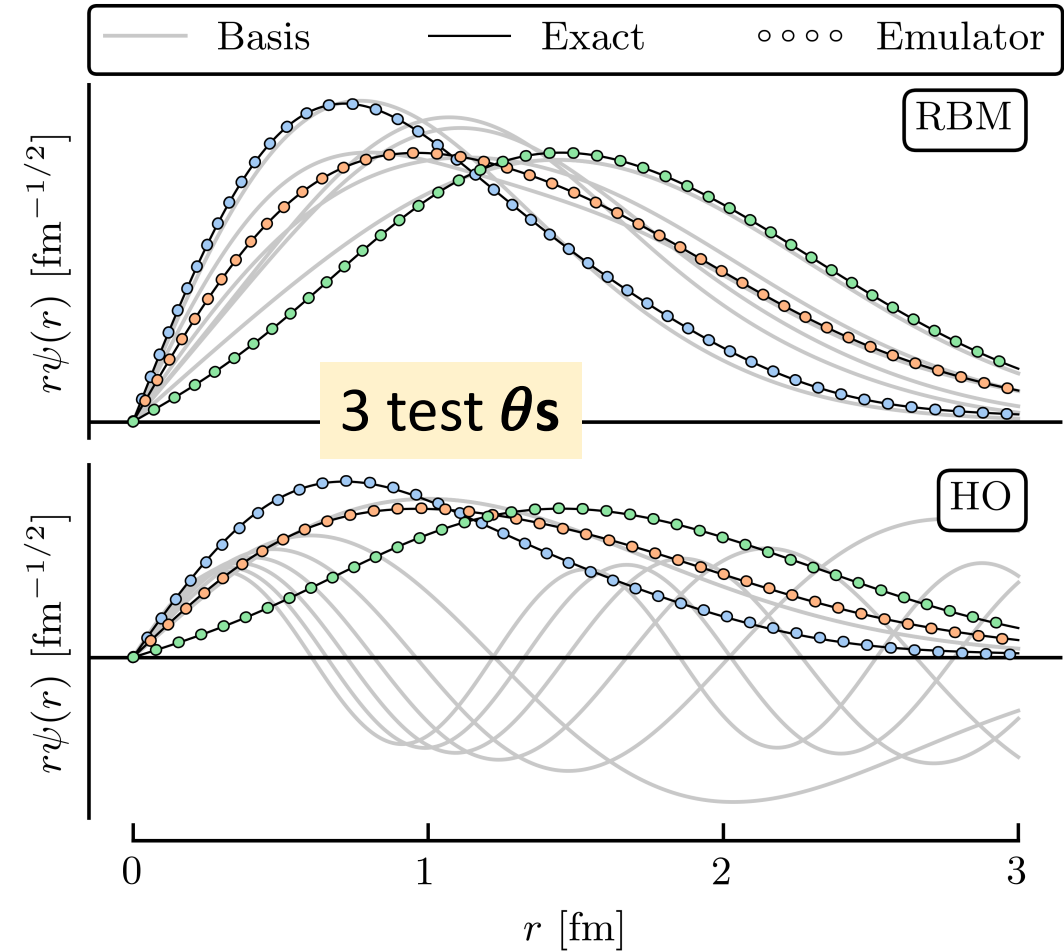
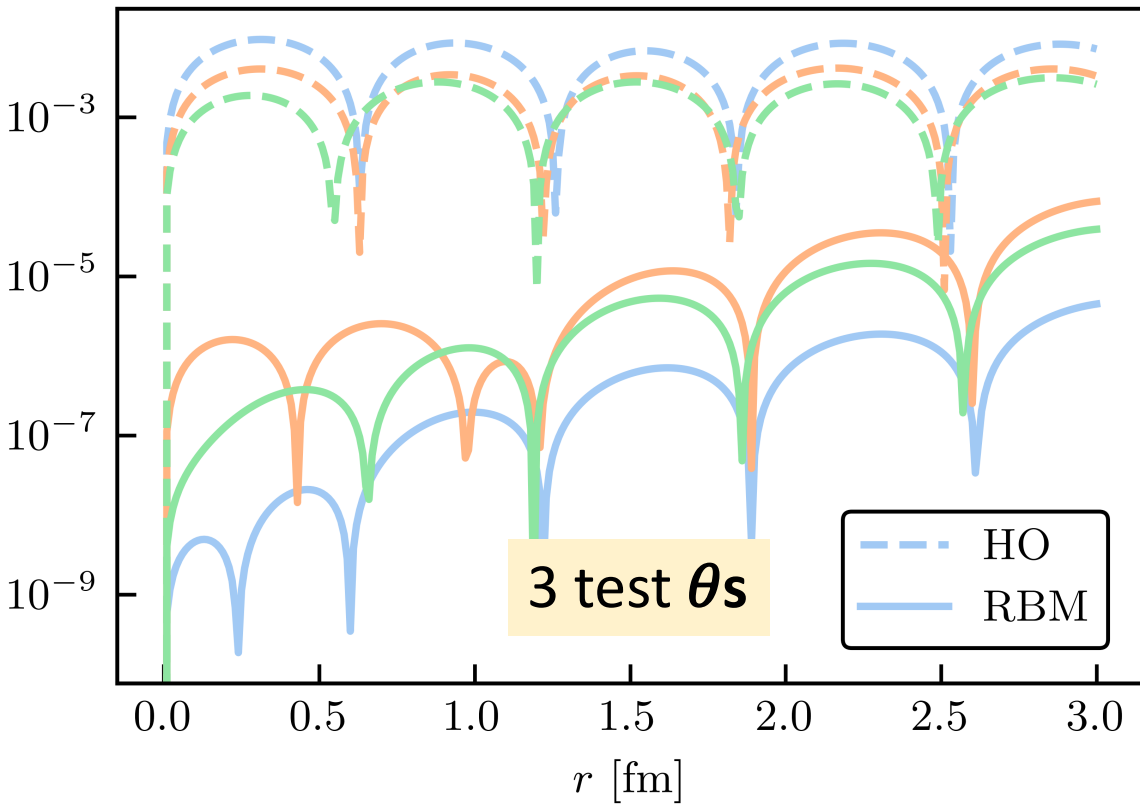


Variational emulator  $\rightarrow$  diagonalize the Hamiltonian  $H(\boldsymbol{\theta})$  in a *finite* basis:  $\sum_{i=1}^{n_b} \beta_i \psi_i$

# Illustrative example: anharmonic oscillator [\[Try your own!\]](#)

Eigenvalue problem:  $H(\boldsymbol{\theta})|\psi\rangle = E|\psi\rangle$   $V(r; \boldsymbol{\theta}) = V_{\text{HO}}(r) + \sum_{n=1}^3 \theta^{(n)} e^{-r^2/\sigma_n^2}$  ← affine!

Wave Function Absolute Residuals

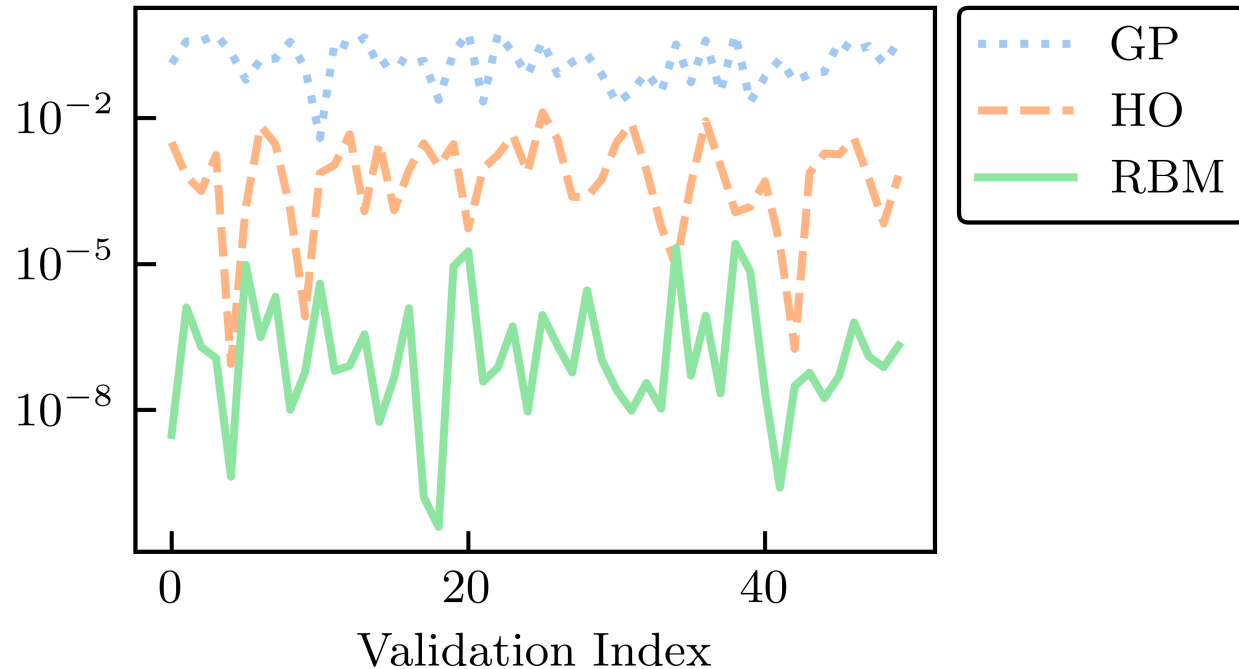


Variational emulator → diagonalize the Hamiltonian  $H(\boldsymbol{\theta})$  in a *finite* basis:  $\sum_{i=1}^{n_b} \beta_i \psi_i$

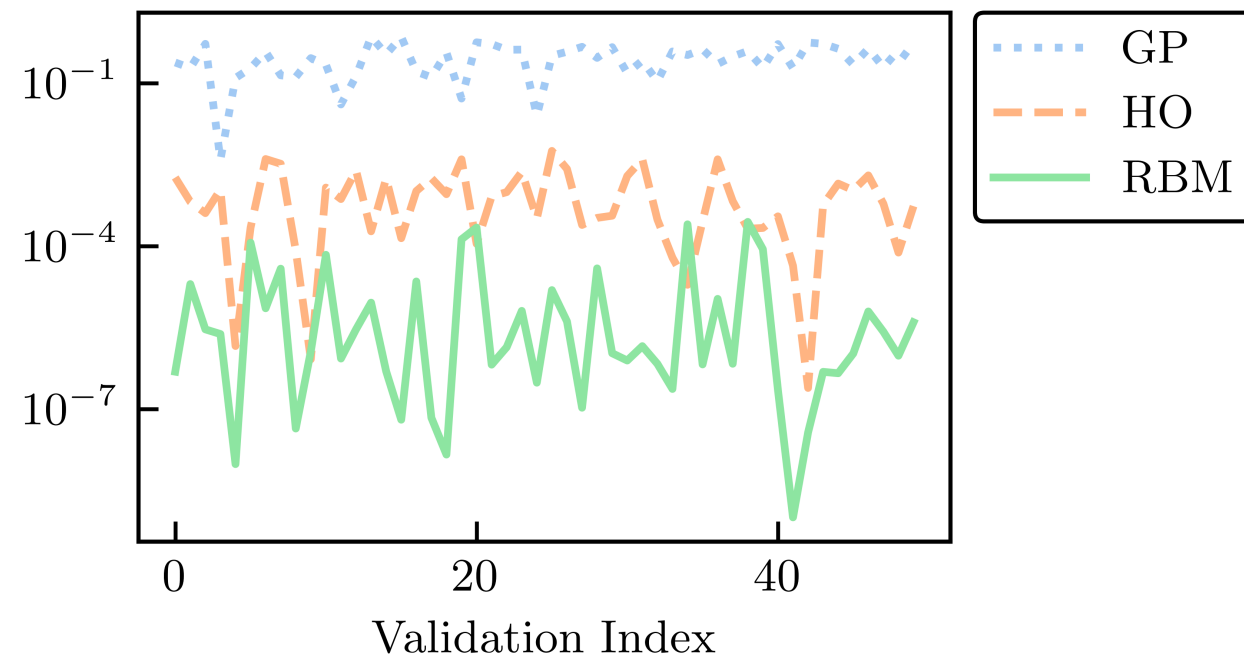
# Illustrative example: anharmonic oscillator [Try your own!]

$$V(r; \boldsymbol{\theta}) = V_{\text{HO}}(r) + \sum_{n=1}^3 \theta^{(n)} e^{-r^2/\sigma_n^2} \quad \leftarrow \text{affine!} \quad \text{Fixed: } \sigma_n = [0.5, 2, 4] \text{ fm}$$

Ground-State Energy Residuals



Ground-State Radius Residuals

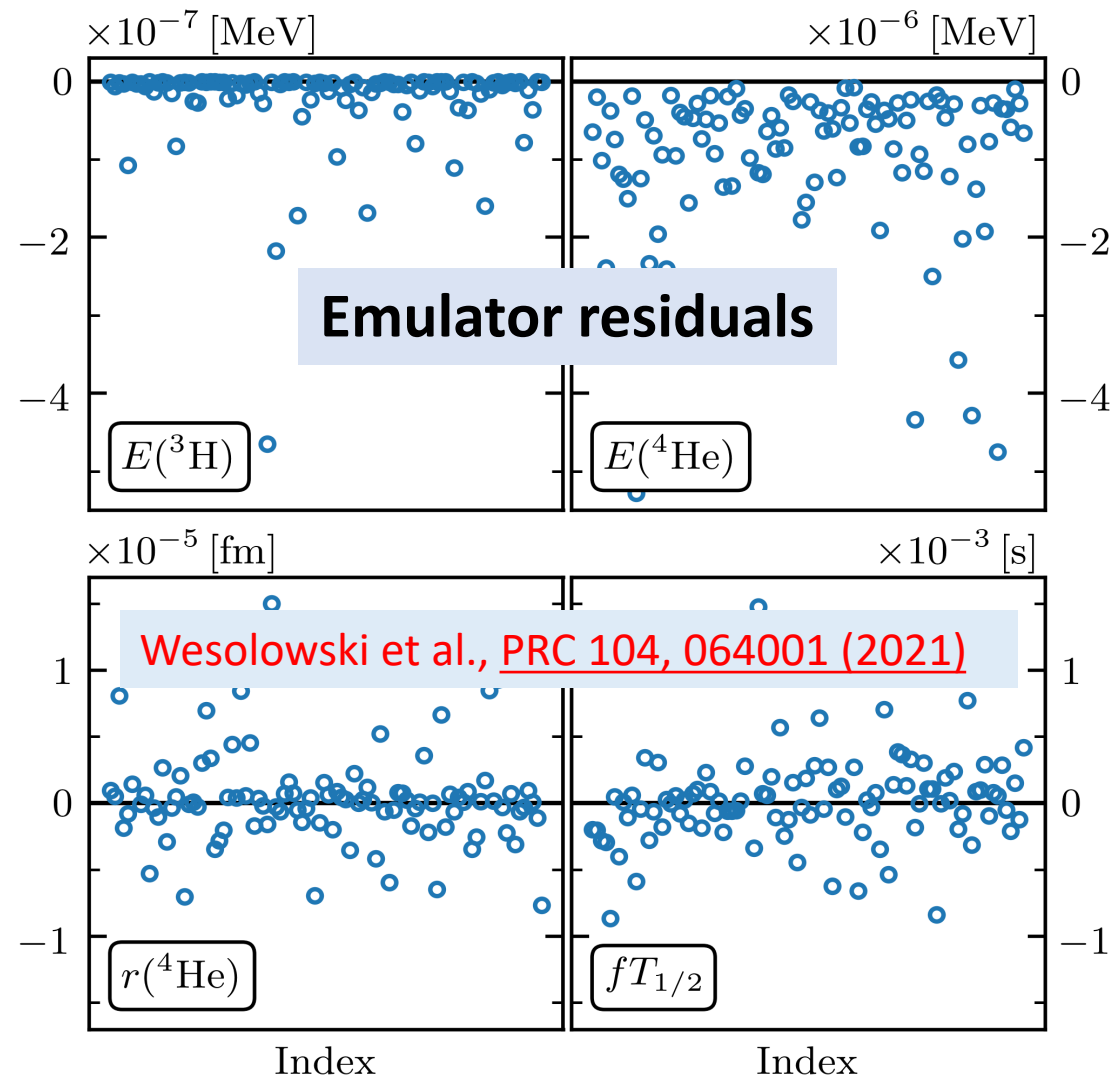
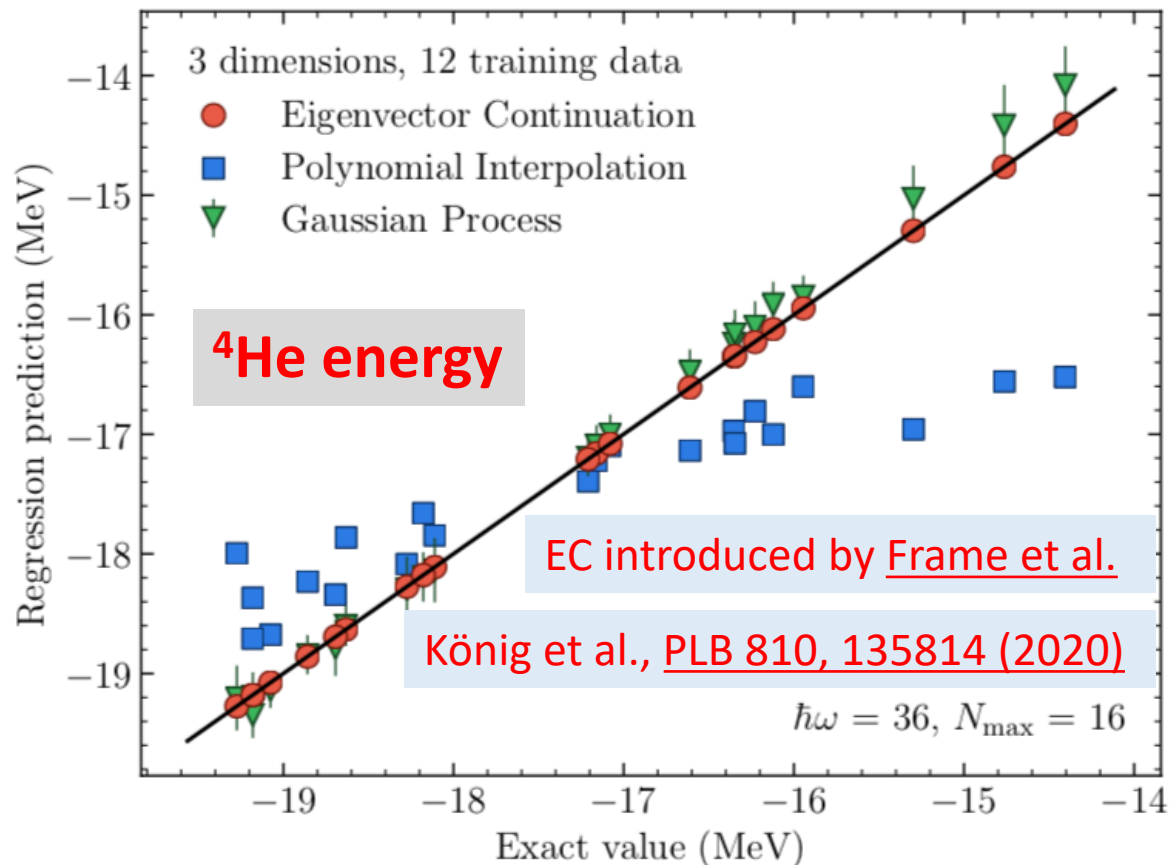


Summary: GP doesn't use the structure of the high-fidelity system to its advantage; HO emulator knows the problem to be solved is an eigenvalue problem; RBM (aka EC) training data are curves rather than scalars, takes advantage of system structure.

# Snapshot RBM emulators for nuclear observables

Ground-state eigenvectors from a selection of parameter sets is an extremely effective variational basis for other parameter sets.

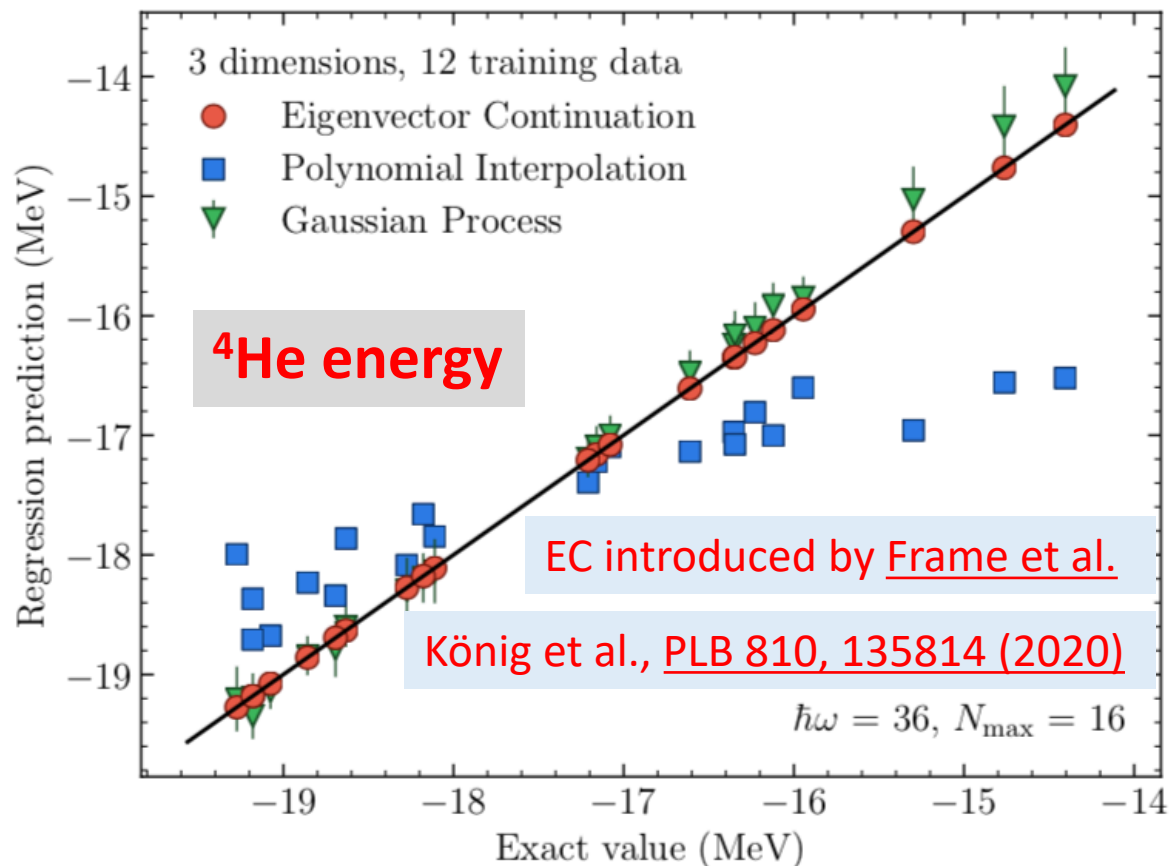
**Characteristics:** fast and accurate!



Emulator doesn't require specialized calculations!

# Snapshot RBM emulators for nuclear observables

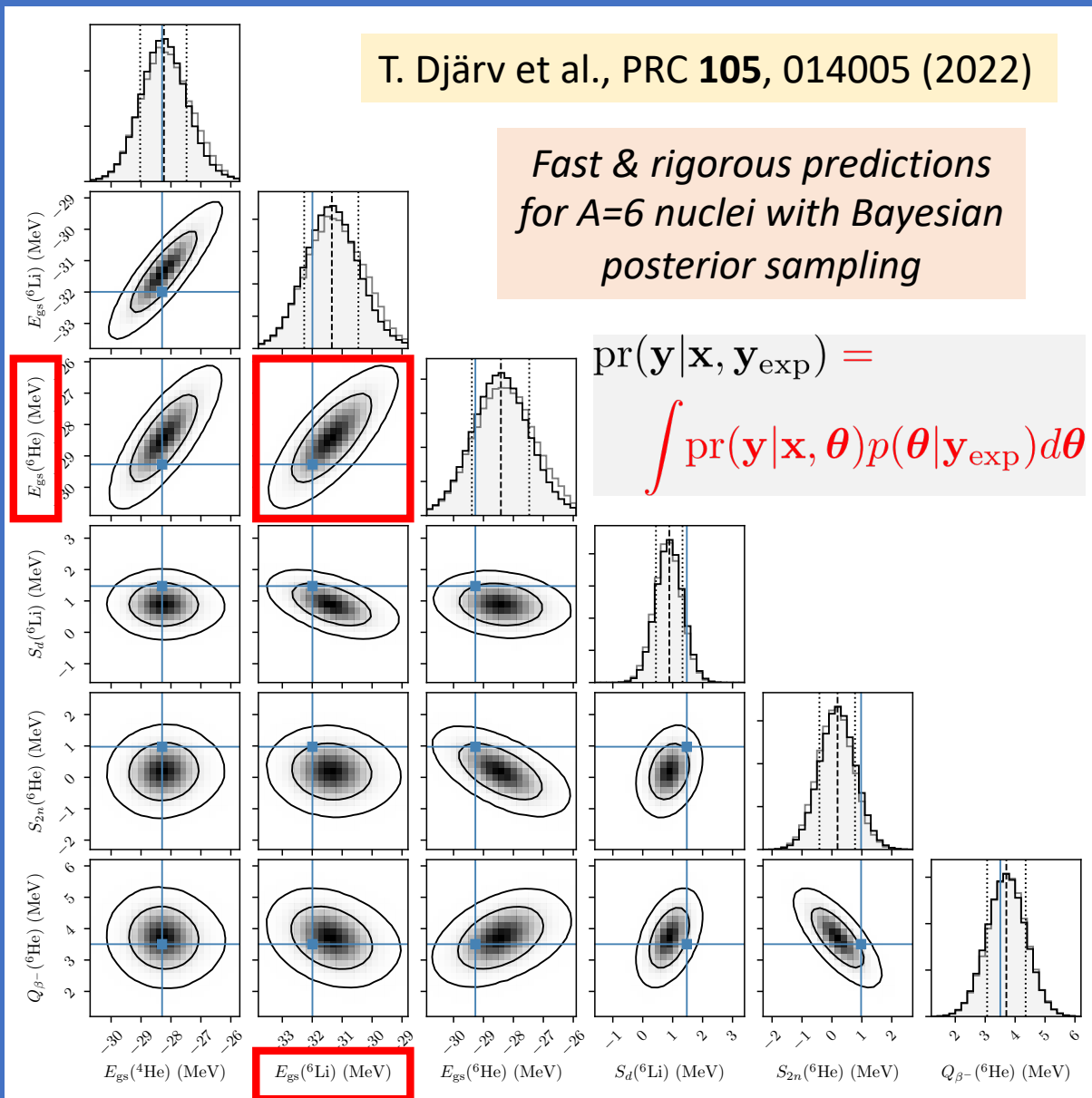
Ground-state eigenvectors from a selection of parameter sets is an extremely effective variational basis for other parameter sets.  
**Characteristics:** fast and accurate!



T. Djärv et al., PRC **105**, 014005 (2022)

*Fast & rigorous predictions for  $A=6$  nuclei with Bayesian posterior sampling*

$$\text{pr}(\mathbf{y}|\mathbf{x}, \mathbf{y}_{\text{exp}}) = \int \text{pr}(\mathbf{y}|\mathbf{x}, \boldsymbol{\theta}) p(\boldsymbol{\theta}|\mathbf{y}_{\text{exp}}) d\boldsymbol{\theta}$$





# Snapshot RBM emulators for nuclear observables

Ground-state eigenvectors from a selection of parameter sets is an extremely effective variational basis for other parameter sets.

**Characteristics:** fast and accurate!

**Already applied to many observables:**

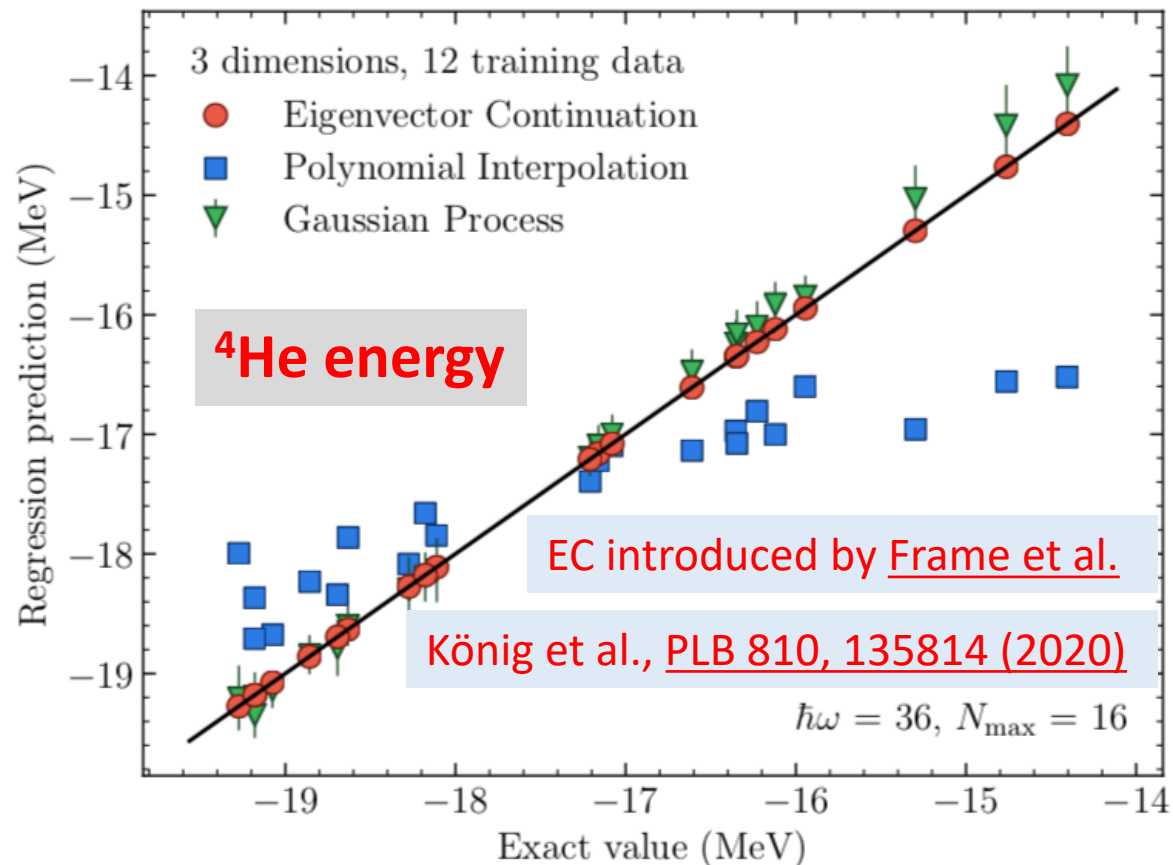
- Ground-state properties (energies, radii)
- Transition matrix elements
- Excited states
- Resonances

**Adapted to special situations and methods**

- Pairing
- Coupled cluster approach; MBPT
- Systems in a finite box
- Subspace diagonalization on quantum computers

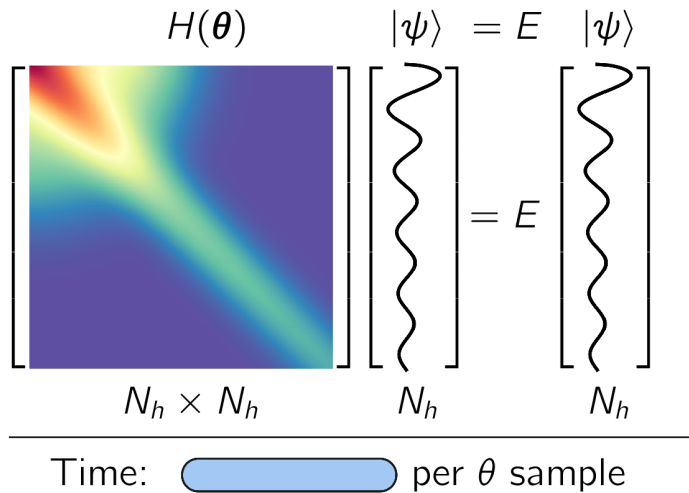
**Extended to non-eigenvalue problems**

- Reactions and scattering



# Constructing a reduced-basis model (aka emulator)

High-fidelity system



CPU time scales with the length of [blue bar]

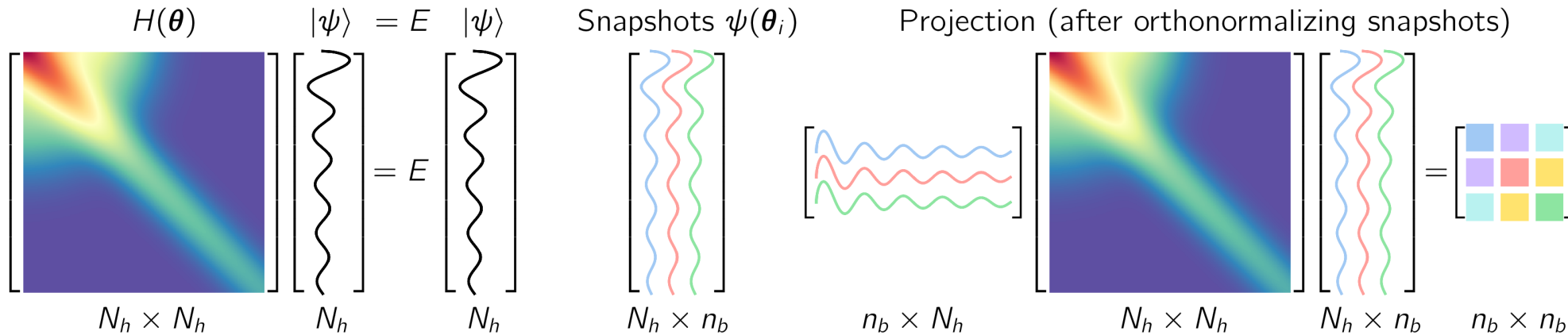
- [\*J. A. Melendez et al., J. Phys. G 49, 102001 \(2022\)\*](#)
- [\*E. Bonilla, P. Giuliani et al., Phys. Rev. C 106, 054322 \(2022\)\*](#)
- [\*P. Giuliani, K. Godbey et al., Front. Phys. 10, 1212 \(2022\)\*](#)
- [\*C. Drischler et al., Quarto + Front. Phys. 10, 1365 \(2022\)\*](#)

# Constructing a reduced-basis model (aka emulator)

High-fidelity system

Constructing a reduced-order model for bound states

Offline stage



Time:  per  $\theta$  sample

$n_b \times$  

$\sim$  

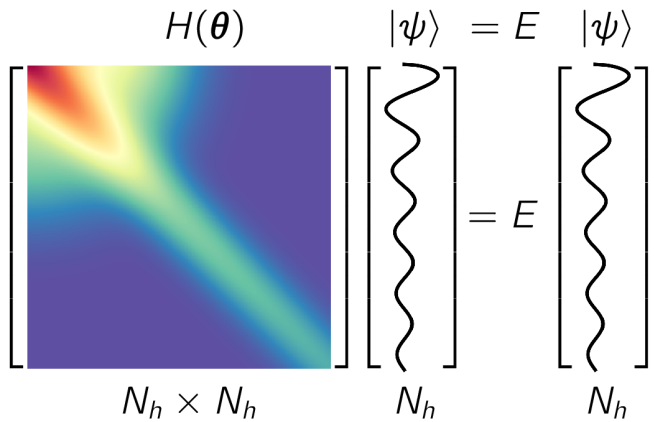
CPU time scales with the length of 

- Offline stage (pre-calculate):
  - Construct basis using snapshots from high-fidelity system (simulator)
  - Project high-fidelity system to small-space using snapshots

- [J. A. Melendez et al., J. Phys. G 49, 102001 \(2022\)](#)
- [E. Bonilla, P. Giuliani et al., Phys. Rev. C 106, 054322 \(2022\)](#)
- [P. Giuliani, K. Godbey et al., Front. Phys. 10, 1212 \(2022\)](#)
- [C. Drischler et al., Quarto + Front. Phys. 10, 1365 \(2022\)](#)

# Constructing a reduced-basis model (aka emulator)

High-fidelity system

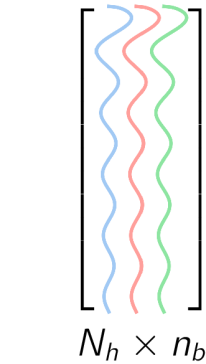


Time:  per  $\theta$  sample

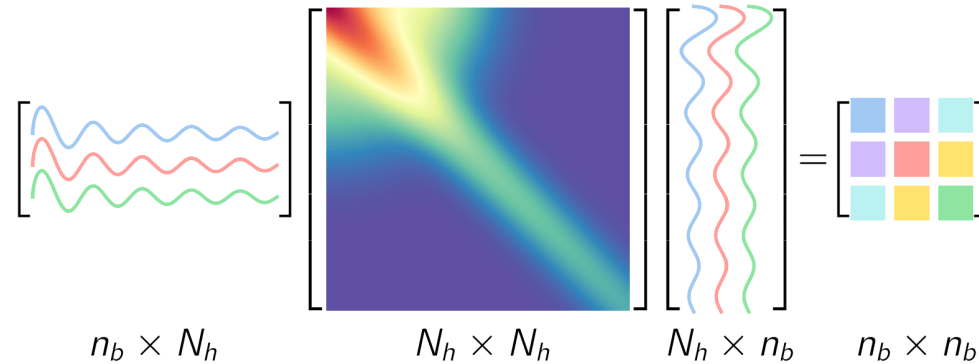
Constructing a reduced-order model for bound states

Offline stage

Snapshots  $\psi(\theta_i)$



Projection (after orthonormalizing snapshots)

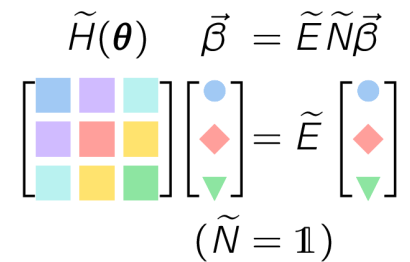


$n_b \times$  


$\sim$  

Online stage

Emulation ( $E \approx \tilde{E}$ )



All size- $n_b$  operations

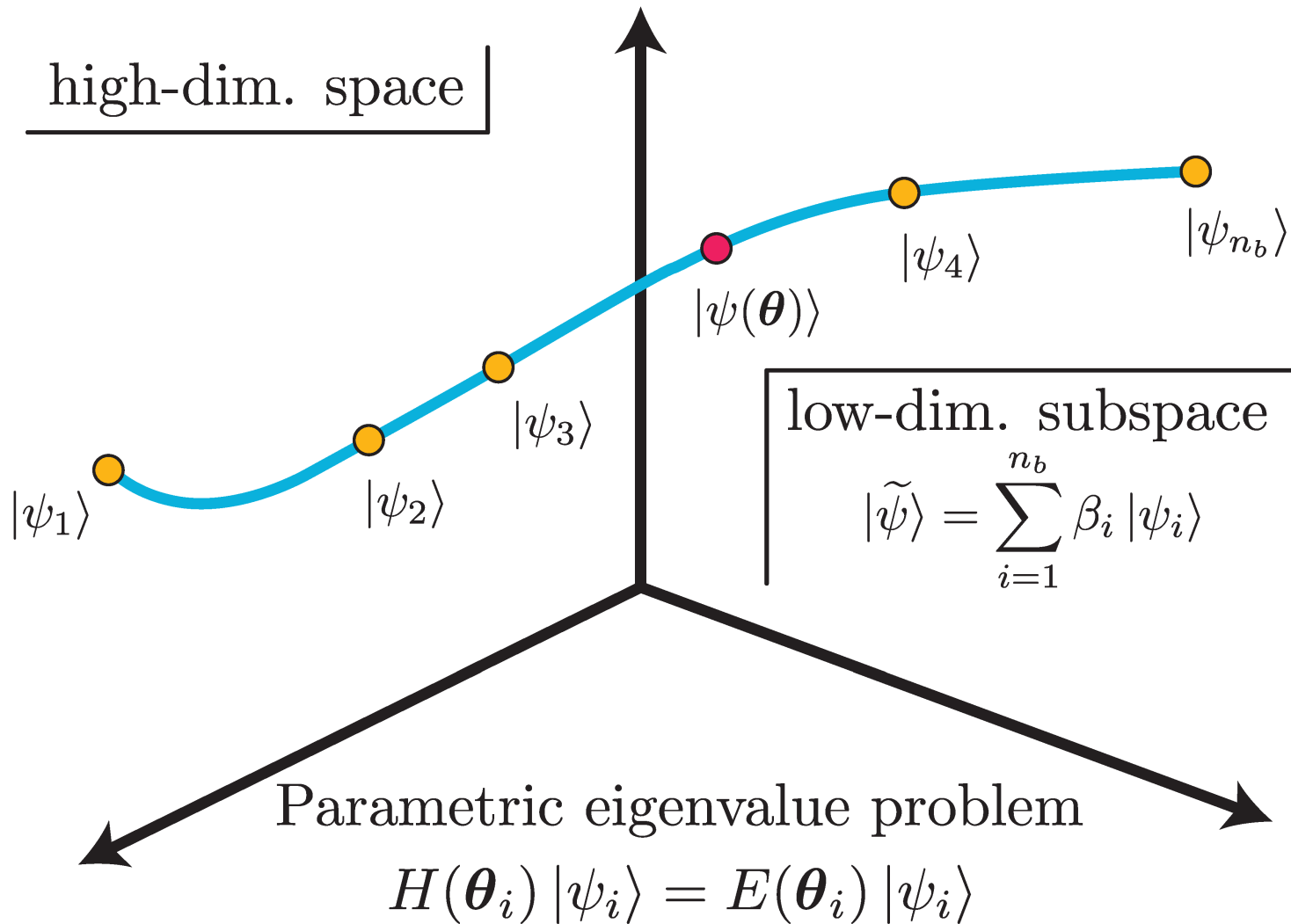
 per  $\theta$  sample

CPU time scales with the length of 

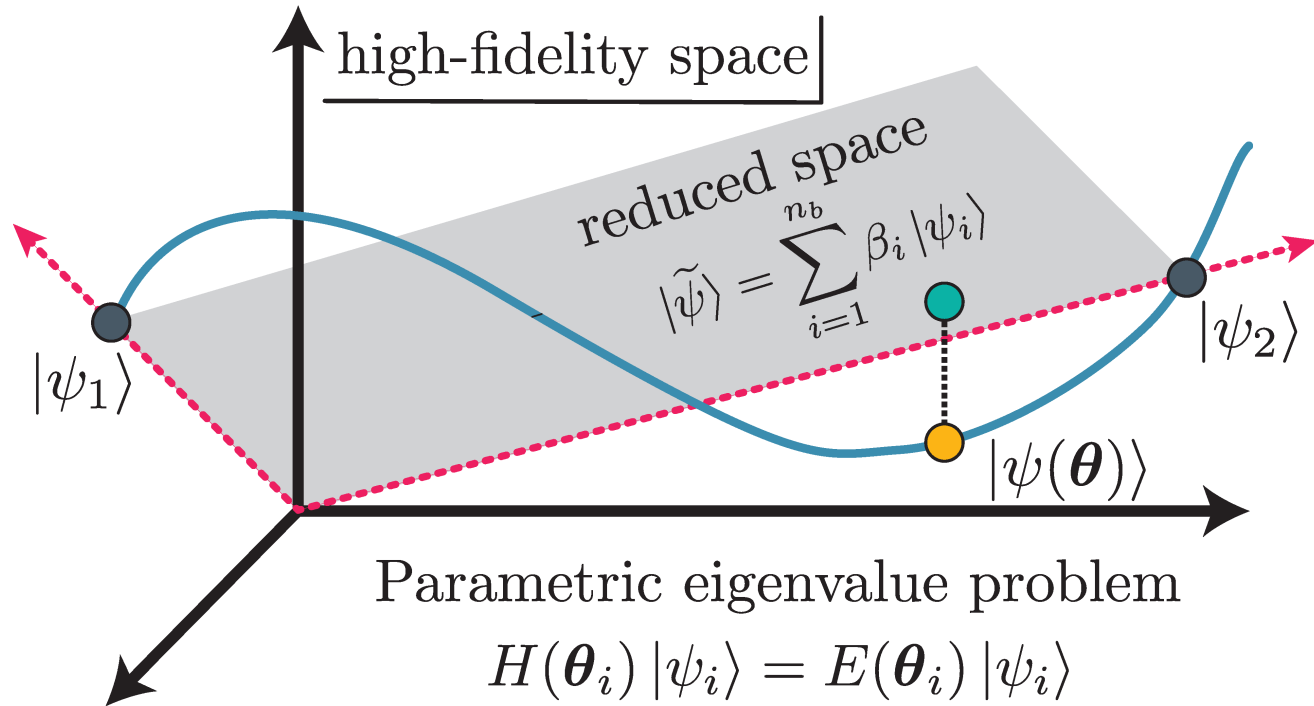
- Offline stage (pre-calculate size  $N_h$ ):
  - Construct basis using snapshots from high-fidelity system (simulator)
  - Project high-fidelity system to small-space using snapshots
- Online stage (emulation size  $n_b$  only):
  - Make many predictions fast & accurately (e.g., for Bayesian analysis)

- [J. A. Melendez et al., J. Phys. G 49, 102001 \(2022\)](#)
- [E. Bonilla, P. Giuliani et al., Phys. Rev. C 106, 054322 \(2022\)](#)
- [P. Giuliani, K. Godbey et al., Front. Phys. 10, 1212 \(2022\)](#)
- [C. Drischler et al., Quarto + Front. Phys. 10, 1365 \(2022\)](#)

# Schematic picture of projection-based emulators



# Schematic picture of projection-based emulators



- Two high-fidelity snapshots ( $\theta_1, \theta_2$ )
- They span the ROM subspace (grey)
- High-fidelity trajectory is in blue.
- Subspace projection shown for  $|\psi(\theta)\rangle$

**Variational  $\rightarrow$  stationary functional**

$$\mathcal{E}[\psi] = \langle \psi | H(\theta) | \psi \rangle - E(\theta) (\langle \psi | \psi \rangle - 1)$$

Use trial  $|\tilde{\psi}\rangle = \sum_{i=1}^{n_b} \beta_i |\psi_i\rangle$  and  $\langle \delta\tilde{\psi} |$

Solve generalized eigenvalue problem:

$$\tilde{H}(\theta) \vec{\beta}(\theta) = \tilde{E}(\theta) \tilde{N} \vec{\beta}(\theta)$$

$$[\tilde{H}(\theta)]_{ij} = \langle \psi_i | H(\theta) | \psi_j \rangle, [\tilde{N}(\theta)]_{ij} = \langle \psi_i | \psi_j \rangle$$

**Galerkin projection  $\rightarrow$  use weak form**

$$\langle \zeta | H(\theta) - E(\theta) | \psi \rangle = 0, \forall \langle \zeta |$$

Reduce dimension:  $|\psi\rangle \rightarrow |\tilde{\psi}\rangle = \sum_{i=1}^{n_b} \beta_i |\psi_i\rangle$

Limit orthogonality:  $\langle \zeta_i | H(\theta) - \tilde{E}(\theta) | \tilde{\psi} \rangle = 0$

Choose  $\langle \zeta_i | = \langle \psi_i |$  (Ritz)  $\equiv$  variational

More general:  $\langle \zeta_i | \neq \langle \psi_i |$  (Petrov-Galerkin)

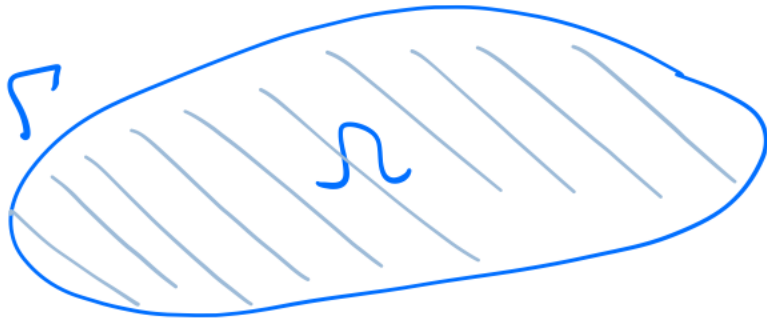
# Variational vs. Galerkin for differential equations

Projection-based emulator for solution  $\psi$  to

$$D(\psi; \boldsymbol{\theta}) = 0 \text{ in } \Omega; \quad B(\psi; \boldsymbol{\theta}) = 0 \text{ on } \Gamma$$

where D and B are operators. Example:

$$[-\nabla^2 \psi = g(\boldsymbol{\theta})]_{\Omega} \quad \left[ \frac{\partial \psi}{\partial n} = f(\boldsymbol{\theta}) \right]_{\Gamma}$$



If affine  $g(\boldsymbol{\theta})$ ,  $f(\boldsymbol{\theta}) \rightarrow$  calculate high-fidelity offline.  
If nonlinear or nonaffine  $\rightarrow$  hyper-reduction, etc.

See [Drischler et al., \(2022\)](#) for details and references

**Variational  $\rightarrow$  stationary functional**

$$S[\psi] = \int_{\Omega} d\Omega F[\psi] + \int_{\Gamma} d\Gamma G[\psi]$$

Use trial  $|\tilde{\psi}\rangle = \sum_{i=1}^{n_b} \beta_i |\psi_i\rangle$  and  $\langle \delta\tilde{\psi}|$

Solve linear algebra problem for  $\vec{\beta}_*$ :

$$\delta S = A\vec{\beta}_* + \vec{b} = 0$$

**Galerkin projection  $\rightarrow$  use weak form**

$$\int_{\Omega} d\Omega \zeta D(\psi) + \int_{\Gamma} d\Gamma \bar{\zeta} B(\psi) = 0$$

Reduce dimension:  $|\psi\rangle \rightarrow |\tilde{\psi}\rangle = \sum_{i=1}^{n_b} \beta_i |\psi_i\rangle$

Test bases:  $|\zeta\rangle = \sum_{i=1}^{n_b} \delta\beta_i |\zeta_i\rangle$ ,  $|\zeta\rangle \rightarrow |\bar{\zeta}\rangle$

$$\Rightarrow \delta\beta_i \left[ \int_{\Omega} d\Omega \zeta_i D(\tilde{\psi}) + \int_{\Gamma} d\Gamma \bar{\zeta}_i B(\tilde{\psi}) \right] = 0$$

# Variational vs. Galerkin emulators via concrete example

E.g., Poisson equation with Neumann BCs  $\rightarrow [-\nabla^2\psi = g(\boldsymbol{\theta})]_{\Omega}$  with  $[\frac{\partial\psi}{\partial n} = f(\boldsymbol{\theta})]_{\Gamma}$

Emulator  $\rightarrow \psi(\boldsymbol{\theta}) \approx \tilde{\psi}(\boldsymbol{\theta}) = \sum_{i=1}^{n_b} (\vec{\beta}_*)_i \psi_i = X \vec{\beta}_*$ ,  $X \equiv [\psi_1 \ \psi_2 \ \cdots \ \psi_{n_b}]$  find optimal  $\vec{\beta}_*$  online

## Variational (Ritz)

$$S[\psi] = \int_{\Omega} d\Omega \left( \frac{1}{2} \nabla\psi \cdot \nabla\psi - g\psi \right) - \int_{\Gamma} d\Gamma f\psi$$

$$\Rightarrow \delta S = \int_{\Omega} d\Omega \delta\psi (-\nabla^2\psi - g) + \int_{\Gamma} d\Gamma \delta\psi \left( \frac{\partial\psi}{\partial n} - f \right)$$

So  $\delta S = 0$  gives the Poisson eq. and BCs. Emulate  $\psi(\boldsymbol{\theta})$ :

$$S[\tilde{\psi}] \rightarrow \delta S[\tilde{\psi}] = \sum_{i=1}^{n_b} \frac{\partial S}{\partial \beta_i} \delta\beta_i = 0 \rightarrow n_b \text{ equations for } \vec{\beta}_*$$

If linear (as here)  $\rightarrow$

$$\tilde{A}\vec{\beta}_* = \vec{g} + \vec{f}, \quad \tilde{A}_{ij} = \int_{\Omega} \nabla\psi_i \cdot \nabla\psi_j,$$

$$g_i = \int_{\Omega} g(\boldsymbol{\theta})\psi_i, \quad f_i = \int_{\Gamma} f(\boldsymbol{\theta})\psi_i$$

If affine  $g(\boldsymbol{\theta})$ ,  $f(\boldsymbol{\theta}) \rightarrow$  calculate high-fidelity offline.

## Ritz-Galerkin

Weak formulation with *test function*  $\zeta$

$$\int_{\Omega} d\Omega \zeta (-\nabla^2\psi - g) + \int_{\Gamma} d\Gamma \zeta \left( \frac{\partial\psi}{\partial n} - f \right) = 0$$

$$\Rightarrow \int_{\Omega} d\Omega (\nabla\zeta \cdot \nabla\psi - g\zeta) - \int_{\Gamma} d\Gamma f\zeta = 0$$

Assert holds for  $\psi \rightarrow \tilde{\psi} = X\vec{\beta}$  and  $\zeta = \sum_{i=1}^{n_b} \delta\beta_i \psi_i$

$$\delta\beta_i \left[ \int_{\Omega} d\Omega \underbrace{(\nabla\psi_i \cdot \nabla\psi_j \beta_j)}_{\tilde{A}_{ij}} - \underbrace{g\psi_i}_{g_i} \right] - \int_{\Gamma} d\Gamma \underbrace{f\psi_i}_{f_i} = 0$$

$\rightarrow$  same result as variational, but Galerkin is more general. If  $\zeta_i \neq \psi_i$ , then *Petrov-Galerkin*.

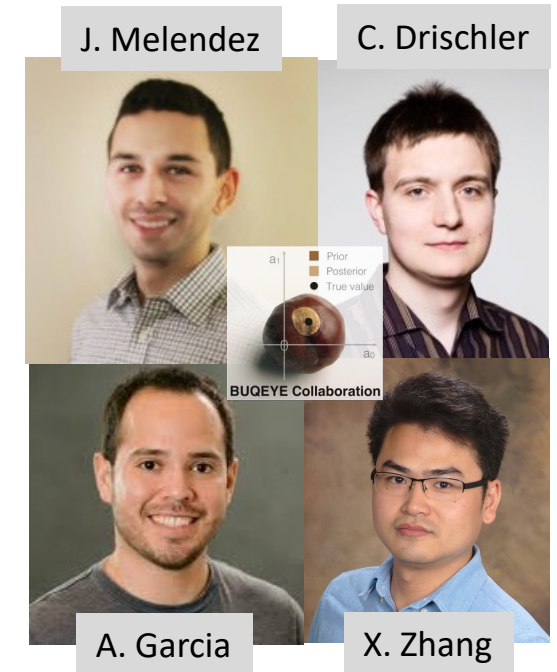


# RBM implementation freedom: examples from scattering

Quantum mechanical two-body scattering problem can be formulated in multiple ways: Schrödinger equation in coordinate or momentum space; variational methods; ...

Variational Principle		Galerkin Projection Information			
Name	Functional for $K$	Strong Form	Trial Basis	Test Basis	Constrained?
Kohn ( $\lambda$ )	$\tilde{K}_E + \langle \tilde{\psi}   H - E   \tilde{\psi} \rangle$	$H   \psi \rangle = E   \psi \rangle$	$ \psi_i\rangle$	$\langle \psi_i  $	Yes
Kohn (No $\lambda$ )	$\langle \tilde{\chi}   H - E   \tilde{\chi} \rangle + \langle \phi   V   \tilde{\chi} \rangle$ + $\langle \phi   H - E   \phi \rangle + \langle \tilde{\chi}   V   \phi \rangle$	$[E - H]   \chi \rangle = V   \phi \rangle$	$ \chi_i\rangle$	$\langle \chi_i  $	No
Schwinger	$\langle \tilde{\psi}   V   \phi \rangle + \langle \phi   V   \tilde{\psi} \rangle$ - $\langle \tilde{\psi}   V - VG_0V   \tilde{\psi} \rangle$	$ \psi\rangle =  \phi\rangle + G_0V  \psi\rangle$	$ \psi_i\rangle$	$\langle \psi_i  $	No
Newton	$V + VG_0\tilde{K} + \tilde{K}G_0V$ - $\tilde{K}G_0\tilde{K} + \tilde{K}G_0VG_0\tilde{K}$	$K = V + VG_0K$	$K_i$	$K_i$	No

See [Drischler et al., \(2022\)](#) for details and references



# RBM implementation freedom: examples from scattering

Quantum mechanical two-body scattering problem can be formulated in multiple ways:  
Schrödinger equation in coordinate or momentum space; variational methods; ...

Variational Principle		Galerkin Projection Information			
Name	Functional for $K$	Strong Form	Trial Basis	Test Basis	Constrained?
Kohn ( $\lambda$ )	$\tilde{K}_E + \langle \tilde{\psi}   H - E   \tilde{\psi} \rangle$	$H  \psi\rangle = E  \psi\rangle$	$ \psi_i\rangle$	$\langle \psi_i  $	Yes
Kohn (No $\lambda$ )	$\langle \tilde{\chi}   H - E   \tilde{\chi} \rangle + \langle \phi   V   \tilde{\chi} \rangle$ + $\langle \phi   H - E   \phi \rangle + \langle \tilde{\chi}   V   \phi \rangle$	$[E - H]  \chi\rangle = V  \phi\rangle$	$ \chi_i\rangle$	$\langle \chi_i  $	No
Schwinger	$\langle \tilde{\psi}   V   \phi \rangle + \langle \phi   V   \tilde{\psi} \rangle$ - $\langle \tilde{\psi}   V - VG_0V   \tilde{\psi} \rangle$	$ \psi\rangle =  \phi\rangle + G_0V  \psi\rangle$	$ \psi_i\rangle$	$\langle \psi_i  $	No
Newton	$V + VG_0\tilde{K} + \tilde{K}G_0V$ - $\tilde{K}G_0\tilde{K} + \tilde{K}G_0VG_0\tilde{K}$	$K = V + VG_0K$	$K_i$	$K_i$	No

See [Drischler et al., \(2022\)](#)  
for details and references

Every variational way  
for scattering has a  
Galerkin counterpart!

Non-variational, also,  
e.g., “origin” emulator  
 $(r\psi)(0) = 0$ ,  $(r\psi)'(0) = 1$   
(see later talks)

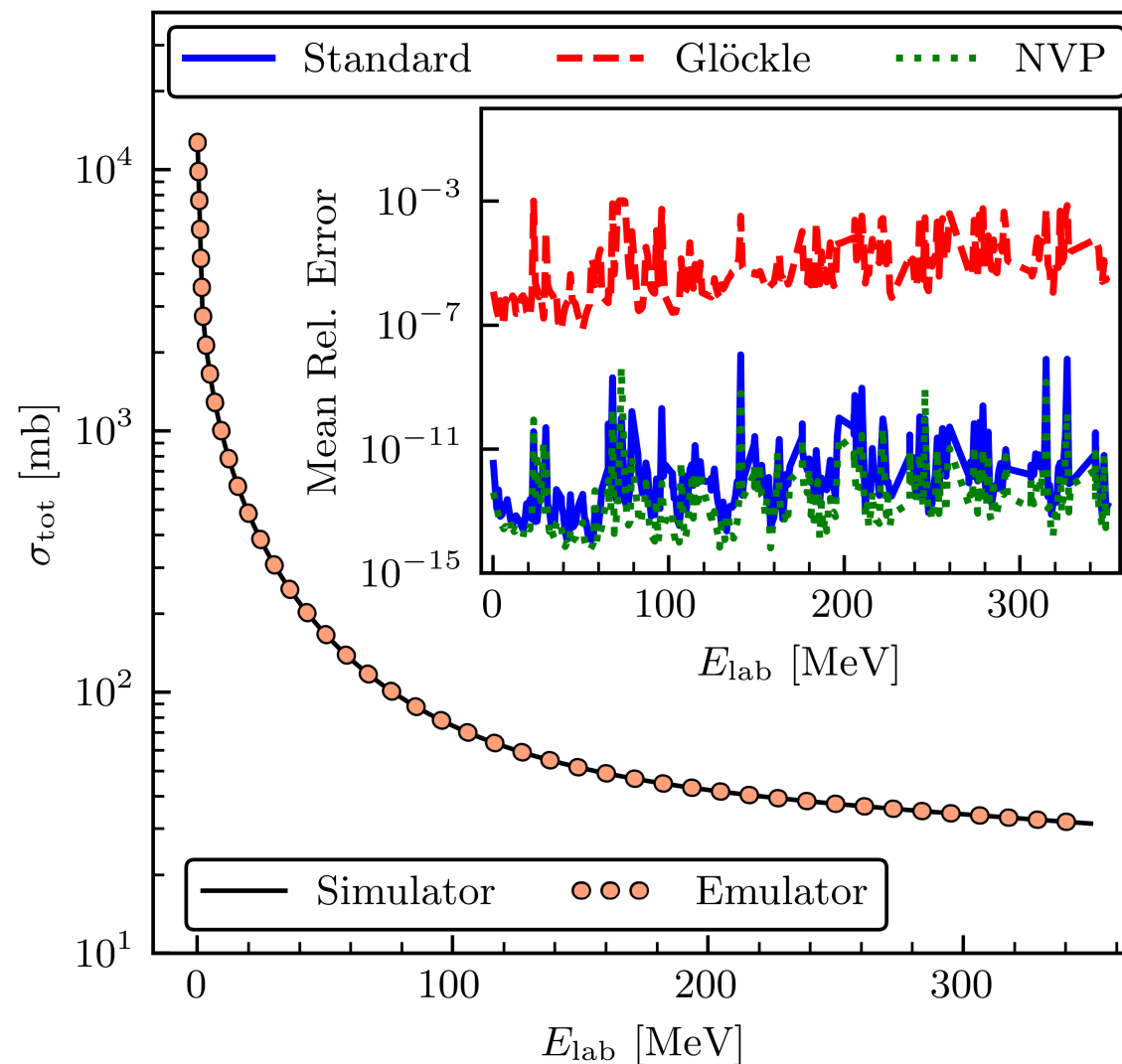
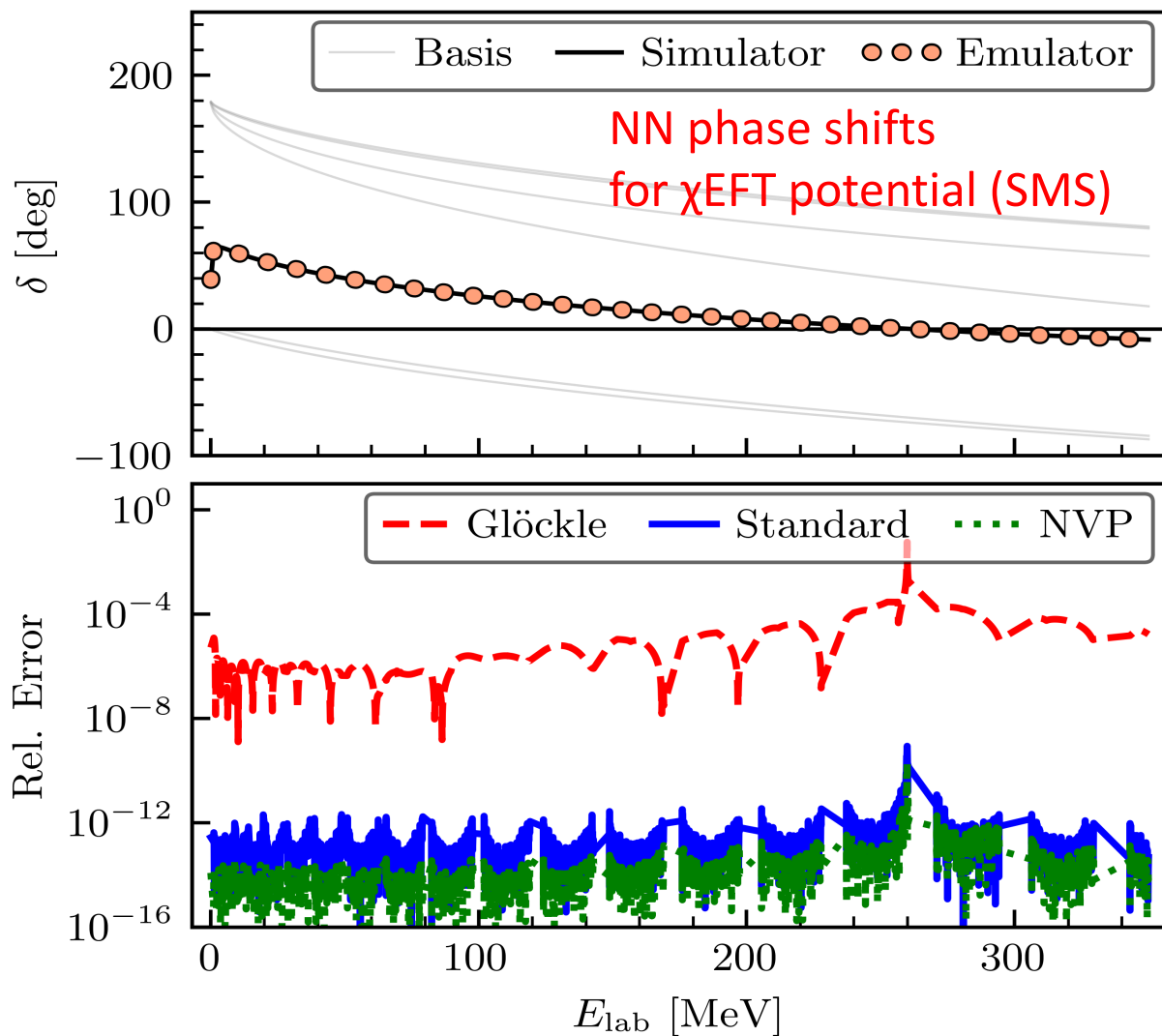
## What is the best way to implement a 3-body scattering emulator?

- E.g, for Bayesian  $\chi$ EFT LEC estimation or nuclear reactions.
- X. Zhang, rjf, [PRC \(2022\)](#) gave proof of principle (bosons) using KVP.

# RBM emulators for NN scattering in chiral EFT (affine!)

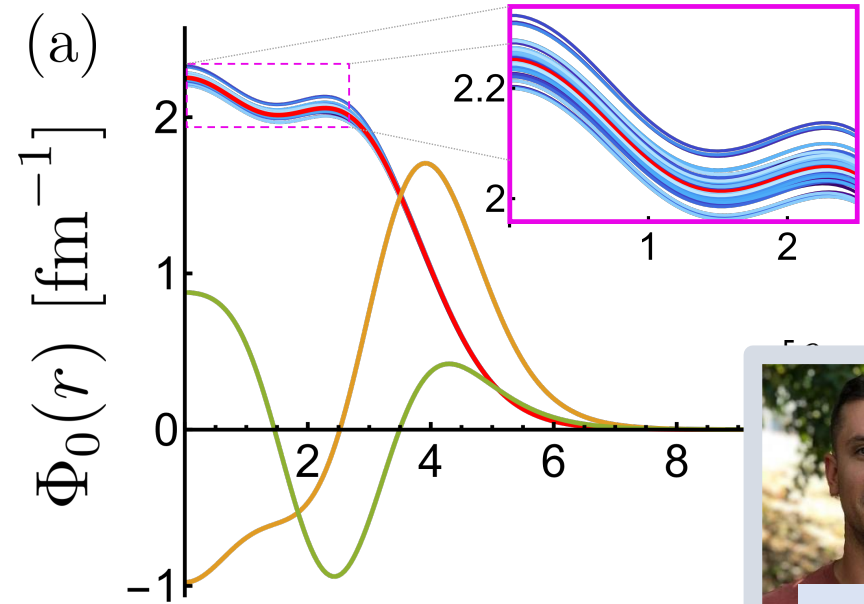
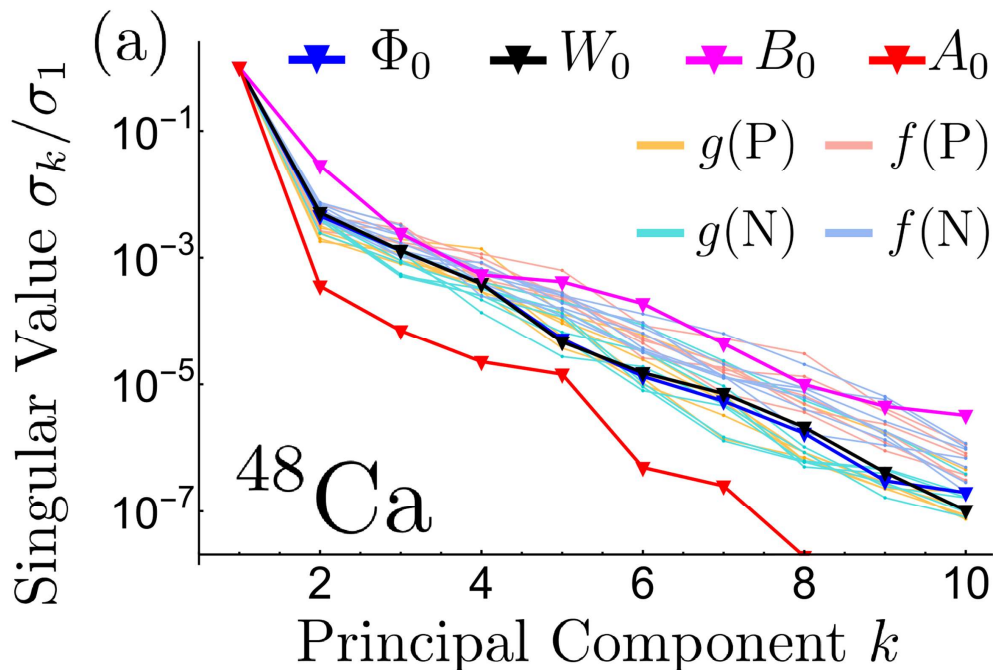
Compare NVP to two implementations of KVP

A. Garcia et al., [PRC 107 \(2023\)](#)



# RBM emulators for EDFs and non-affine

- Energy density functionals (EDFs) present new challenges.
- P. Giuliani et al., "[Bayes goes fast ...](#)" (also "[Training and Projecting](#)")  
→ apply Galerkin RBM to EDFs (covariant mean field, Skyrme)
- Efficient basis to evaluate functional for many parameter sets.
- → Fast and accurate emulation, ideal for Bayesian inference!



Also today: RBM for non-linear, non-affine problems. BAND: ROSE software (e.g., for opt. potl.)

# Summary of key RBM elements

Vast range of problems have been attacked with MOR in science and engineering, including heat transfer, fluid dynamics, electronic DFT, ...  $\rightarrow$  coupled ode's and pde's (incl. time-dependent and nonlinear); eigenvalue problems; and more!

**There's likely something out there in the MOR literature analogous to what you do!**

Large speed-ups from **offline-online paradigm** if heavy compute resources are offline.

$\rightarrow$  move size- $\psi$  operations offline so that emulation varying  $\theta$  online is efficient.

Key: exploit *affine* parameter dependence in operators, e.g.,  $H(\theta) = \sum_n h_n(\theta) H_n$

For non-linear systems and non-affine parameters, use *hyper-reduction* methods.

**Projection-based:** (i) choose *low-dimensional rep. of  $\psi$*  and (ii) write in integral form.

For (i):  $\tilde{\psi}(\theta) \equiv \sum_{i=1}^{N_b} \beta_i \psi_i = X \vec{\beta}$ ,  $X \equiv [\psi_1 \psi_2 \cdots \psi_{N_b}]$  with  $X$  found offline.

**Snapshot approaches:** construct  $X$  from high-fidelity solutions  $\psi_i = \psi(\theta_i)$  at set  $\{\theta_i\}$ .

# Research avenues for emulator applications in NP (I)

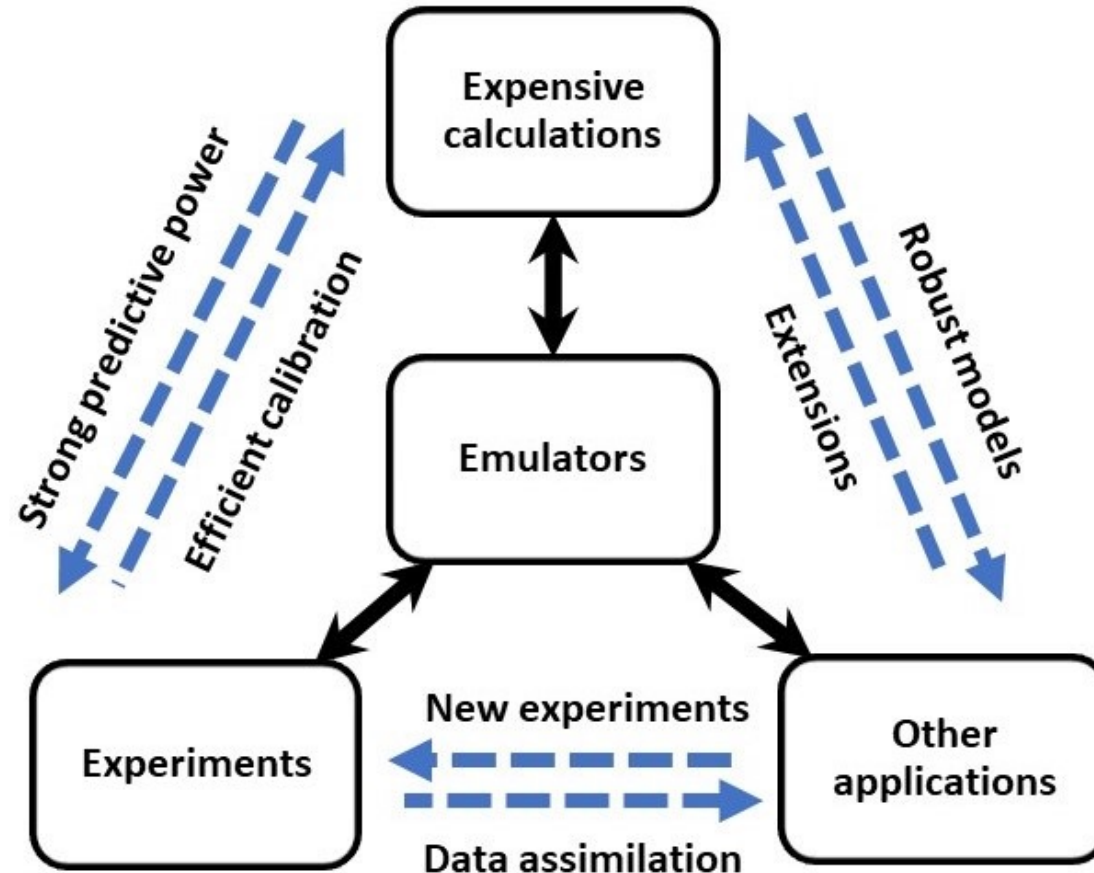
- Emulator uncertainties need to be robustly quantified; this should be facilitated by the extensive literature on uncertainties in the RBM.
- What are best practices for efficient implementation of NP emulators? Can we exploit MOR software libraries from other fields, such as pyMOR?
- Galerkin and variational emulators for bound-state and scattering calculations are equivalent for properly chosen test and trial basis. But [Petrov]-Galerkin emulators are more general; applications to nonlinear problems in NP can be fruitful but face challenges, e.g., hyperreduction methods need approximations that worsen accuracy and need UQ.
- Technical aspects to explore further, e.g., greedy (or active-learning) and SVD-based algorithms for choosing training points more effectively.

# Research avenues for emulator applications in NP (II)

- Further applications to reactions: long-range Coulomb interactions and optical potentials beyond two-body systems; emulators for time-dependent DFT; emulators for nuclear dynamics at much higher energy scales (JLAB/EIC).
- Emulators for extrapolation far from support of training ([Frame et al., 2018](#)); emulators as resummation tool to increase convergence radius of series expansions ([Demol et al., 2020](#)); emulators to extrapolate finite-box simulations of quantum systems ([Yapa and König, 2022](#)); emulation in the complex energy plane for general quantum continuum states ([Zhang, 2022](#)).
- Exploring synergy between projection-based and machine learning methods is a new direction for MOR (e.g., POD-DL-ROM by [Fresca and Manzoni, 2022](#)).
- Can we exploit in emulator applications use of field theory and RG methods for analyzing deep neural networks (e.g., *Why is AI hard and Physics simple?* by [Roberts \(2021\)](#))?

# Role of emulators: new workflows for NP applications

From [Xilin Zhang](#), rjf, *Fast emulation of quantum three-body scattering*, Phys. Rev. C **105**, 064004 (2022).



How can ISNET facilitate these new workflows based on shared emulators?

If you can create fast & accurate™ emulators for observables, you can do calculations without specialized expertise and expensive resources!



# Thank you!

## Coming attractions:

2023: **Workshop on Eigenvector continuation method in nuclear structure and reaction theory**, May 30-June 2, at CEA, France

2023: **FRIB-TA Summer School on Practical Uncertainty Quantification and Emulator Development in Nuclear Physics**, June 26-28, at FRIB.

## Jupyter and Quora *books* for nuclear applications:

[Learning from Data \(OSU course Physics 8820\)](#)

[BUQEYE Guide to Projection-Based Emulators in Nuclear Physics](#)

[Reduced Basis Methods in Nuclear Physics](#)

Extra slides

# ANNs and GPs meet effective theories and RG

- Recent developments\* merge field theory and renormalization group (RG) insights and methods to describe ANNs (e.g., *Why is AI hard and Physics simple?* by [Roberts \(2021\)](#)).

- Principle of *sparsity* plus effective theory approach (cf. Ising Model for counting):

$$2^{\mathcal{O}(N)} \xrightarrow[\text{locality}]{k} \mathcal{O}(N^k) \xrightarrow[\text{locality}]{\text{spatial}} \mathcal{O}(N) \xrightarrow[\text{invariance}]{\text{translational}} \mathcal{O}(1)$$

- Exploit large width limit of ANNs, in which they become GPs (via generalized central limit theorem). Finite width expansion in depth / width of network; RG flow to criticality.
- Effective [field] theory and RG approaches are natural for (nuclear) many-body theory! The perturbative approach to leading non-trivial order is like Ginzburg-Landau form.
- **Can we apply insights to emulators and forge connections with reduced basis methods?**

[\*For up-to-date references, see *Structures of neural network effective theories* by Banta et al., [arXiv:2305.02334](#).]

# Lexicon for Model Order Reduction (MOR)

Term	Definition or usage
<i>High fidelity</i>	Highly accurate, usually for costly calculation [Full-Order Model (FOM)]
<i>Reduced-order model</i>	General name for an emulator resulting from applying MOR techniques.
<i>Intrusive</i>	Non-intrusive treats FOM as black box; intrusive requires coding.
<i>Offline-online paradigm</i>	Heavy compute done once (offline); cheap to vary parameters (online).
<i>Affine</i>	Parameter dependence factors from operators, e.g., $H(\boldsymbol{\theta}) = \sum_n h_n(\boldsymbol{\theta})H_n$
<i>Snapshots</i>	High-fidelity calculations at a set of parameters and/or times.
<i>Proper Orthogonal Decomposition (POD)</i>	Generically the term POD is used for PCA-type reduction via SVD. In snapshot context, PCA is applied to reduce/orthogonalize snapshot basis.
<i>Greedy algorithm</i>	Serially find snapshot locations $\boldsymbol{\theta}_i$ at largest expected error (fast approx.).
<i>Reduced basis methods</i>	Or RBMs. Implement snapshot-based projection methods.
<i>Hyper-reduction methods</i>	Approximations to non-linearity or non-affineness (e.g., EIM).

# Parametric MOR emulator workflow

Bird's eye view but still for projection-based PMOR only (i.e., not an exhaustive set!)

## (1) Sampling across range of parameters $\theta$ for $N_{\text{sample}}$ candidate snapshots $\rightarrow \{\theta_i\}$

- E.g., space-filling design (like latin hypercube) or center near emulated values.
- Want  $N_b \leq N_{\text{sample}}$  snapshots; locate wisely based on basis construction method.

## (2) Generating a basis $X$ from the snapshots to create. Multiple options, including:

- *Proper Orthogonal Decomposition* (POD) [cf. PCA]  $\rightarrow$  extract most important basis vectors. Compute all  $N_{\text{sample}}$  snapshots  $\psi(\theta_i)$  but keep  $N_b$  based on SVD.
- *Greedy algorithm* is an iterative approach: next location  $\theta_i$  from *fast* estimated emulator error at  $N_{\text{sample}}$  values and choose value with largest expected error.
- For time-dependent case, sample also in time or frequency. Many options here!

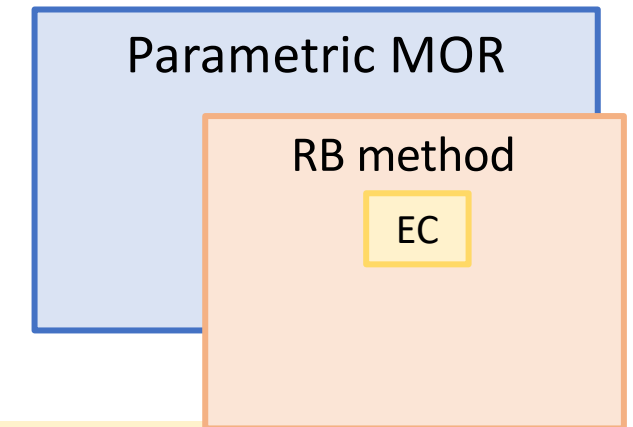
## (3) Construct the reduced system. Single basis $X$ or multiple bases across $\theta$

- Linear system and affine operators  $\rightarrow$  projecting to single basis works well.
- If non-linear or non-affine  $\rightarrow$  *hyper-reduction* approaches: e.g., empirical interpolation method EIM or DEIM, which finds an affine (separable) expansion.

# Some model reduction methods in context

*Reduced Basis* method (1980) widely used to emulate PDEs in reduced-order approach. Specific choices in MOR framework:

- Parameter set chosen using greedy algorithm (or POD)
- Single basis  $X$  constructed from snapshots
- RB model built from global basis projection



*Eigenvector continuation* (EC) is a particular implementation of the RB method

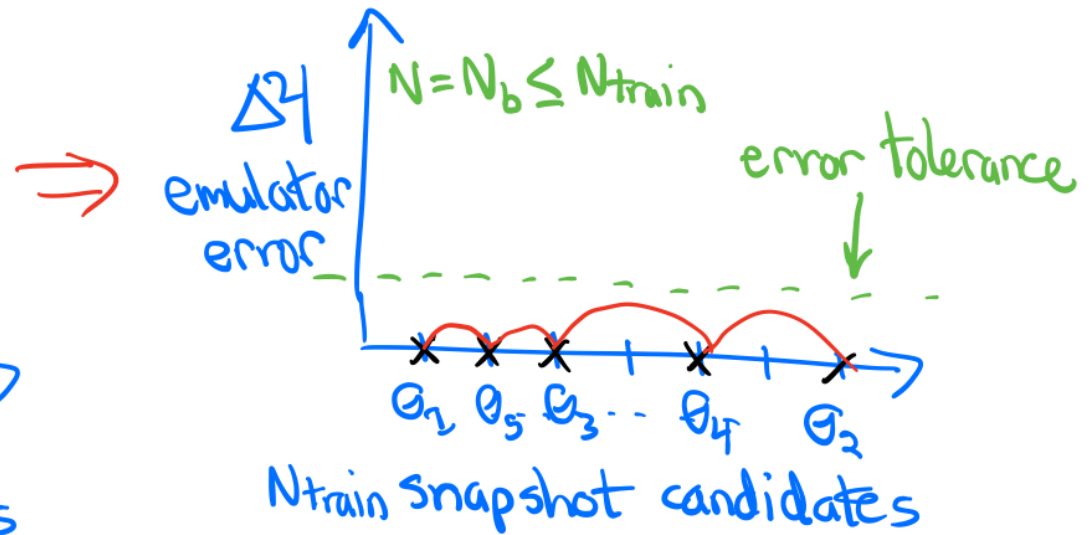
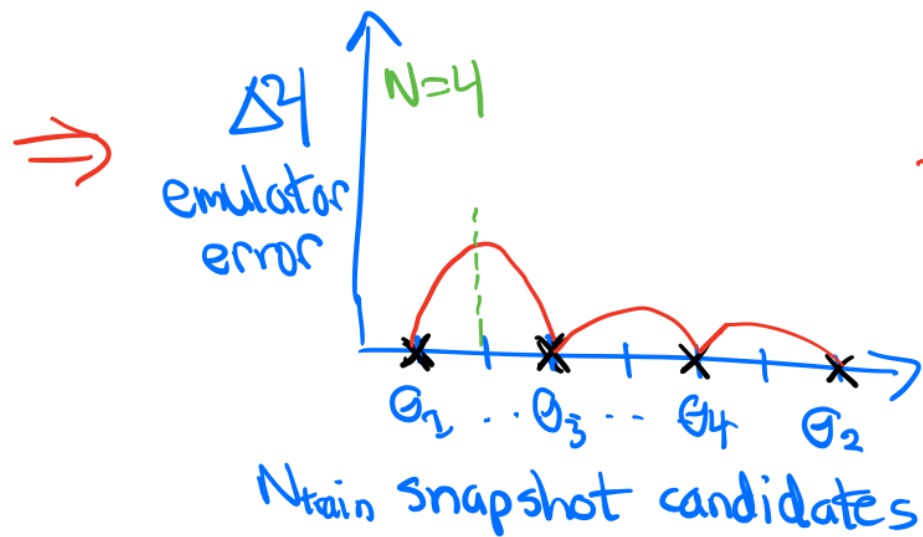
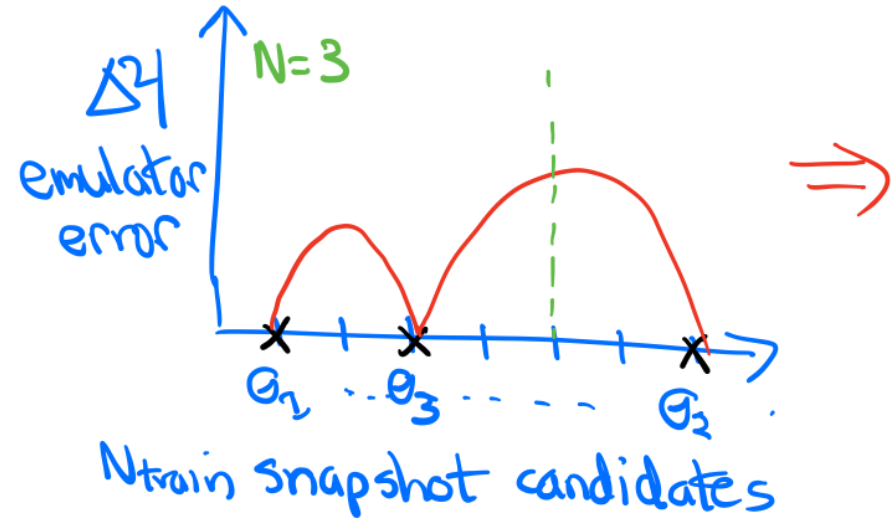
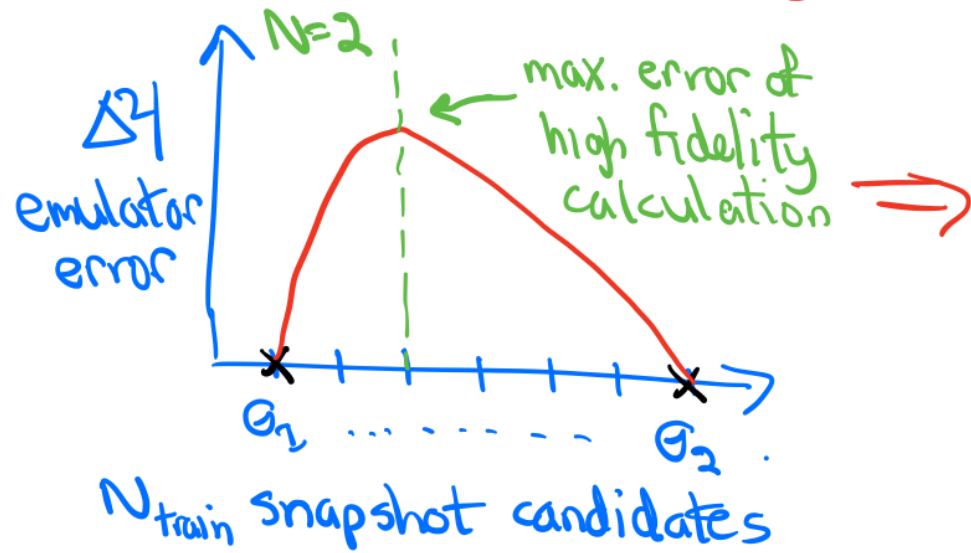
→ parametric reduced-order model for an eigenvalue problem (lots of prior art)

- Global basis constructed with snapshot-based POD approach
- “Active learning” by Sarkar and Lee adds greedy sampling algorithm for next  $\theta_i$

**Summary:** general features of *good* reduced-order emulators

- System dependent → works best when QOI lies in low-D manifold and operations on  $\psi$  can be avoided during online phase
- Relative smoothness of parameter dependence
- Affine parameter dependence (or effective hyper-reduction or other approach)

# CARTOONS FOR GREEDY ALGORITHM



Stop once desired error tolerance reached (or given # steps)

# Empirical interpolation method for nonaffine/nonlinear

Key: avoid costly order  $\psi$  (i.e., FOM) evaluations  $\rightarrow$  approximation strategy.

- Some cases: use low-order terms of a Taylor series expansion.
- More general: selective sampling of nonlinear terms with interpolation.
- Includes empirical interpolation method (EIM) and discrete variant DEIM.

## EIM basics [adapted from Hesthaven (2016)]

- Ingredients are ( $Q$  is an integer):
  - $Q$  interpolation points  $\mathbf{x}_1, \dots, \mathbf{x}_Q$
  - $Q$  parameter points  $\boldsymbol{\theta}_1, \dots, \boldsymbol{\theta}_Q$  ( $\boldsymbol{\theta} \equiv \boldsymbol{\mu}$ )
  - $Q$  basis functions  $h_1, \dots, h_Q$
- The function  $g$  is nonaffine in  $\mathbf{x}$  and  $\boldsymbol{\theta}$
- Interpolation is  $I_Q[g_{\boldsymbol{\theta}}](x) = \sum_{q=1}^Q c_q(\boldsymbol{\theta}) h_q(x)$  where  $I_Q[g_{\boldsymbol{\theta}}](x_j) = g_{\boldsymbol{\theta}}(x_j)$   $j = 1, \dots, Q$  is found by solving 
$$\sum_{q=1}^Q c_q(\boldsymbol{\theta}) h_j(x_j) = g_{\boldsymbol{\theta}}(x_j) \quad j = 1, \dots, Q$$
- The  $h_j$  are found as linear combinations of snapshots  $g_{\boldsymbol{\theta}_1}, \dots, g_{\boldsymbol{\theta}_Q}$  (see box at right).

### Algorithm: Empirical Interpolation Method

**Input:** A family of functions  $g_{\boldsymbol{\mu}} : \Omega \rightarrow \mathbb{R}$ , parametrized by a parameter  $\boldsymbol{\mu} \in \mathbb{P}_{\text{EIM}}$  and a target error tolerance  $\text{tol}$ .

**Output:** A set of  $Q$  basis functions  $\{h_q\}_{q=1}^Q$  and interpolation points  $\{x_q\}_{q=1}^Q$ .

Set  $q = 1$ . Do while  $\text{err} < \text{tol}$ :

1. Pick the sample point

$$\boldsymbol{\mu}_q = \arg \sup_{\boldsymbol{\mu} \in \mathbb{P}_{\text{EIM}}} \|g_{\boldsymbol{\mu}} - \mathcal{I}_{q-1}[g_{\boldsymbol{\mu}}]\|_{\mathcal{X}_{\Omega}},$$

and the corresponding interpolation point

$$x_q = \arg \sup_{x \in \Omega} |g_{\boldsymbol{\mu}_q}(x) - \mathcal{I}_{q-1}[g_{\boldsymbol{\mu}_q}](x)|. \quad (5.5)$$

2. Define the next basis function as the scaled error function

$$h_q = \frac{g_{\boldsymbol{\mu}_q} - \mathcal{I}_{q-1}[g_{\boldsymbol{\mu}_q}]}{g_{\boldsymbol{\mu}_q}(x_q) - \mathcal{I}_{q-1}[g_{\boldsymbol{\mu}_q}](x_q)}. \quad (5.6)$$

3. Define the error

$$\text{err} = \|\text{err}_p\|_{L^{\infty}(\mathbb{P}_{\text{EIM}})} \quad \text{with} \quad \text{err}_p(\boldsymbol{\mu}) = \|g_{\boldsymbol{\mu}} - \mathcal{I}_{q-1}[g_{\boldsymbol{\mu}}]\|_{\mathcal{X}_{\Omega}},$$

and set  $q := q + 1$ .



# Eigenvector continuation (EC) for scattering

$$\hat{H}(\boldsymbol{\theta}) = \hat{T} + \hat{V}(\boldsymbol{\theta}) = \hat{T} + \sum_a \theta^{(a)} \mathcal{O}^{(a)} \quad \text{with LECs } \boldsymbol{\theta} = \{\theta^{(a)}\} \quad \begin{array}{l} \text{Affine dependence} \\ \text{(here chiral)} \end{array}$$

**K matrix:**  $k_\ell(E) = \tan \delta_\ell(E)$  [cf.  $s_\ell(E) = e^{2i\delta_\ell(E)}$ ] Take  $\ell = 0$  here,  $p \equiv \sqrt{2\mu E}$

$$\text{Kohn: } \delta \left[ \frac{[k_0(E)]_{\text{trial}}}{p} - \frac{2\mu}{\hbar^2} \langle \psi_{\text{trial}} | \hat{H}(\boldsymbol{\theta}) - E | \psi_{\text{trial}} \rangle \right] = 0 \quad \text{with } |\psi_{\text{trial}}\rangle \xrightarrow{r \rightarrow \infty} \frac{1}{p} \sin(pr) + \frac{k_0(E)}{p} \cos(pr)$$

# Eigenvector continuation (EC) for scattering

$$\hat{H}(\boldsymbol{\theta}) = \hat{T} + \hat{V}(\boldsymbol{\theta}) = \hat{T} + \sum_a \theta^{(a)} \mathcal{O}^{(a)} \quad \text{with LECs } \boldsymbol{\theta} = \{\theta^{(a)}\} \quad \begin{array}{l} \text{Could be chiral EFT or} \\ \text{AV18 or ...} \end{array}$$

**K matrix:**  $k_\ell(E) = \tan \delta_\ell(E)$  [cf.  $s_\ell(E) = e^{2i\delta_\ell(E)}$ ] Take  $\ell = 0$  here,  $p \equiv \sqrt{2\mu E}$

**Kohn:**  $\delta \left[ \frac{[k_0(E)]_{\text{trial}}}{p} - \frac{2\mu}{\hbar^2} \langle \psi_{\text{trial}} | \hat{H}(\boldsymbol{\theta}) - E | \psi_{\text{trial}} \rangle \right] = 0$  with  $|\psi_{\text{trial}}\rangle \xrightarrow{r \rightarrow \infty} \frac{1}{p} \sin(pr) + \frac{k_0(E)}{p} \cos(pr)$

**EC:**  $|\psi_{\text{trial}}\rangle = \sum_{i=1}^N c_i |\psi_E(\boldsymbol{\theta}_i)\rangle \implies c_i = \sum_j (\Delta \tilde{U})_{ij}^{-1} ([k_0/p]_j - \lambda)$  and  $\lambda = \frac{\sum_{ij} (\Delta \tilde{U})_{ij}^{-1} ([k_0/p]_j - 1)}{\sum_{ij} (\Delta \tilde{U})_{ij}^{-1}}$

with  $\Delta \tilde{U}_{ij}(E) \equiv \frac{2\mu}{\hbar^2} \langle \psi_E(\boldsymbol{\theta}_i) | 2\hat{V}(\boldsymbol{\theta}) - \hat{V}(\boldsymbol{\theta}_i) - \hat{V}(\boldsymbol{\theta}_j) | \psi_E(\boldsymbol{\theta}_j) \rangle$  ← Coulomb cancels!

- Stationary functional for  $k_\ell(E)$  but not an upper (or lower bound) → still works!
- Use nugget regularization to deal with ill-conditioning and/or mix boundary conditions
- EC works for local or non-local potentials, r-space or k-space, complex potentials, 3-body
- More recent: also works for complex  $E$  and extrapolating in  $E$  (Xilin Zhang)

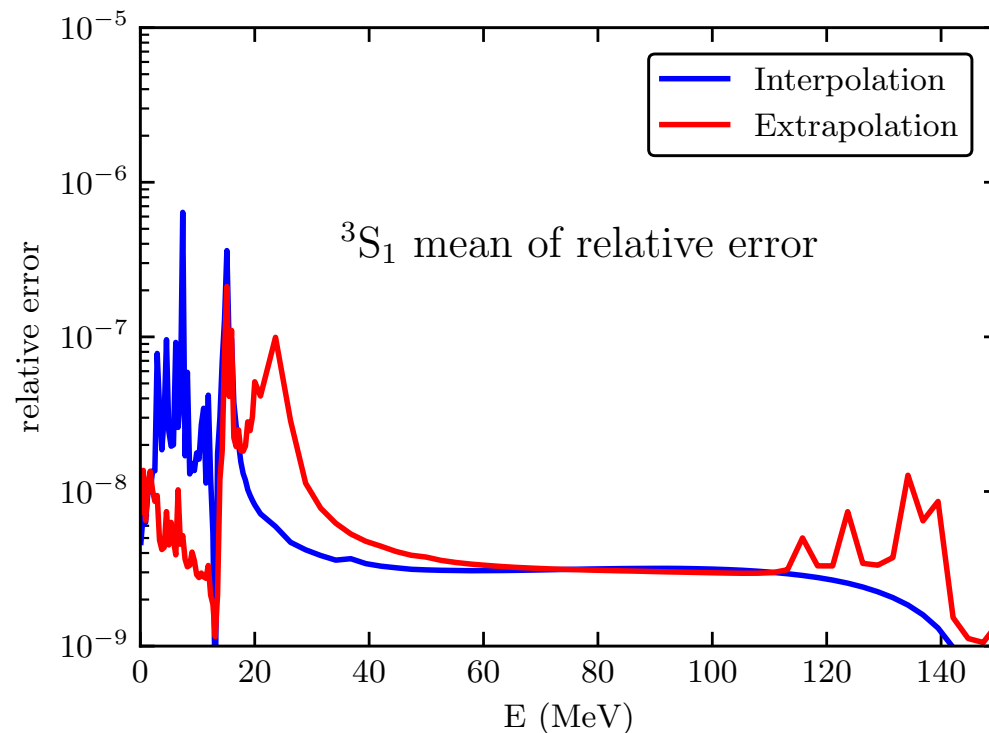
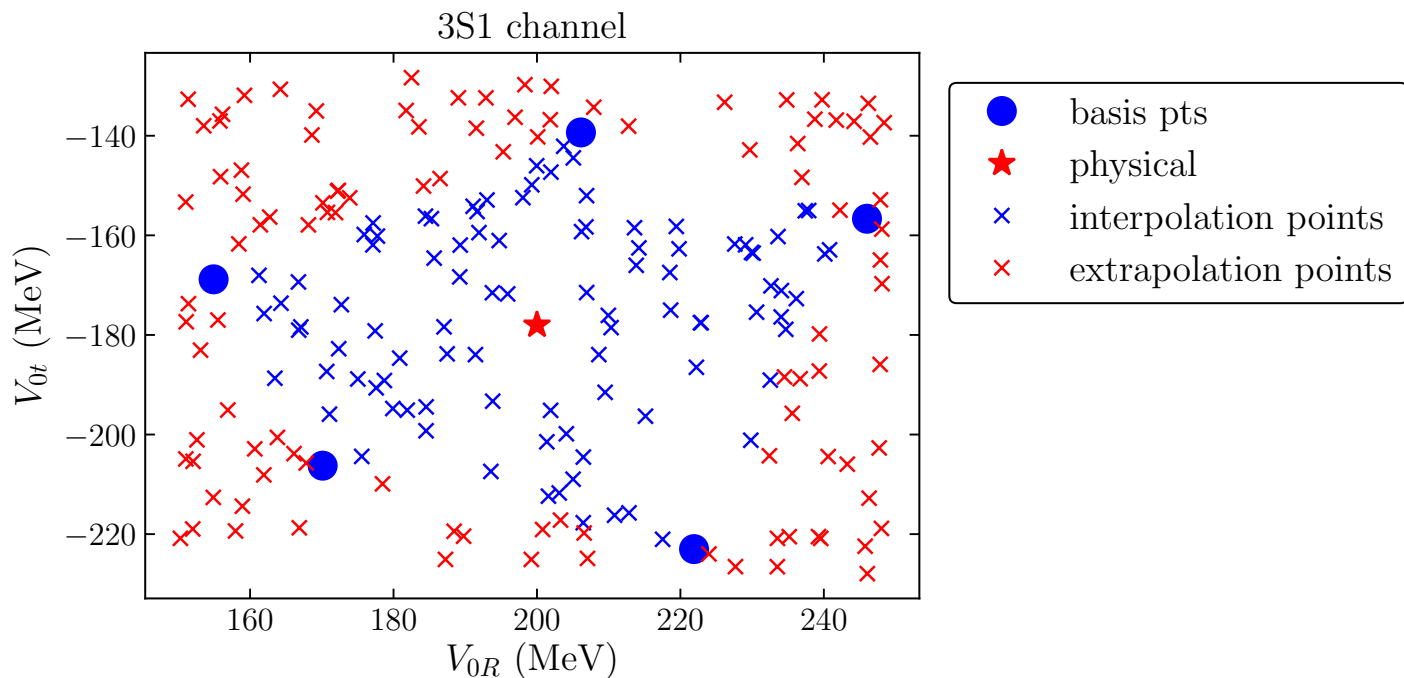
# Testing eigenvector continuation (EC) for scattering

Many different model problems tested: square well, + Coulomb, Yamaguchi potential, ...

→ one example: Minnesota potential in  $^3S_1$  channel (other plots available with notebooks)

$$V_{3S_1}(r) = V_{0R} e^{-\kappa_R r^2} + V_{0t} e^{-\kappa_t r^2} \quad \text{with } \kappa_R = 1.487 \text{ fm}^{-2} \quad \kappa_t = 0.639 \text{ fm}^{-2} \quad (\text{fixed})$$

$$\theta = \{V_{0R}, V_{0t}\} \xrightarrow{\text{“physical”}} \{200 \text{ MeV}, -178 \text{ MeV}\}$$



Better: choose basis points by “greedy algorithm”

# Emulating the Lippmann-Schwinger (LS) equation

LS equation:

Sets of parameters:

K-matrix formulation:

$$K(\vec{a}) = V(\vec{a}) + V(\vec{a}) G_0(E_q) K(\vec{a}) \rightarrow \{\vec{a}_i\} \rightarrow K_\ell(E_q) = -\tan \delta_\ell(E_q)$$
$$E_q = q^2 / 2\mu$$

Newton variational principle (NVP):

$$\tilde{K}(\vec{\beta}) = \sum_{i=1}^{n_t} \beta_i K_i \rightarrow \mathcal{K}[\tilde{K}] = V + V G_0 \tilde{K} + \tilde{K} G_0 V - \tilde{K} G_0 \tilde{K} + \tilde{K} G_0 V G_0 \tilde{K}$$
$$\mathcal{K}[K_{\text{exact}} + \delta K] = K_{\text{exact}} + (\delta K)^2$$

Implementation:

$$\langle \phi' | \mathcal{K}(\vec{a}, \vec{\beta}) | \phi \rangle = \langle \phi' | V(\vec{a}) | \phi \rangle + \vec{\beta}^T \vec{m}(\vec{a}) - \frac{1}{2} \vec{\beta}^T M(\vec{a}) \vec{\beta}$$

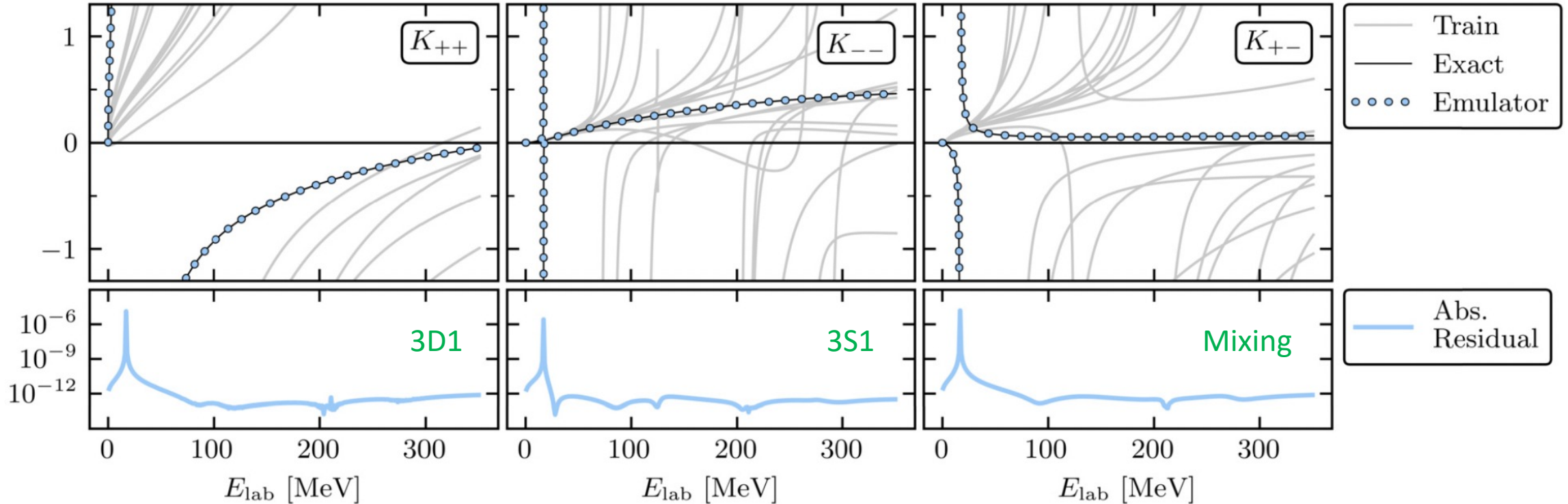
$$\frac{d\mathcal{K}}{d\vec{\beta}} = 0 \rightarrow \langle \phi' | \mathcal{K}(\vec{a}, \vec{\beta}) | \phi \rangle \approx \langle \phi' | V(\vec{a}) | \phi \rangle + \frac{1}{2} \vec{m}^T M^{-1}(\vec{a}) \vec{m}$$

*J. A. Melendez et al., Phys. Lett. B 821, 136608 (2021)*

# NVP emulation: SMS chiral potential

- Emulation of 3S1-3D1 coupled channel
- Basis size of 12 at  $N^4LO+$

*Dealing with anomalies/singularities:  
C. Drischler et al.,  
arXiv: 2108.08269 (2021)*



*J. A. Melendez et al., Phys.  
Lett. B 821, 136608 (2021)*