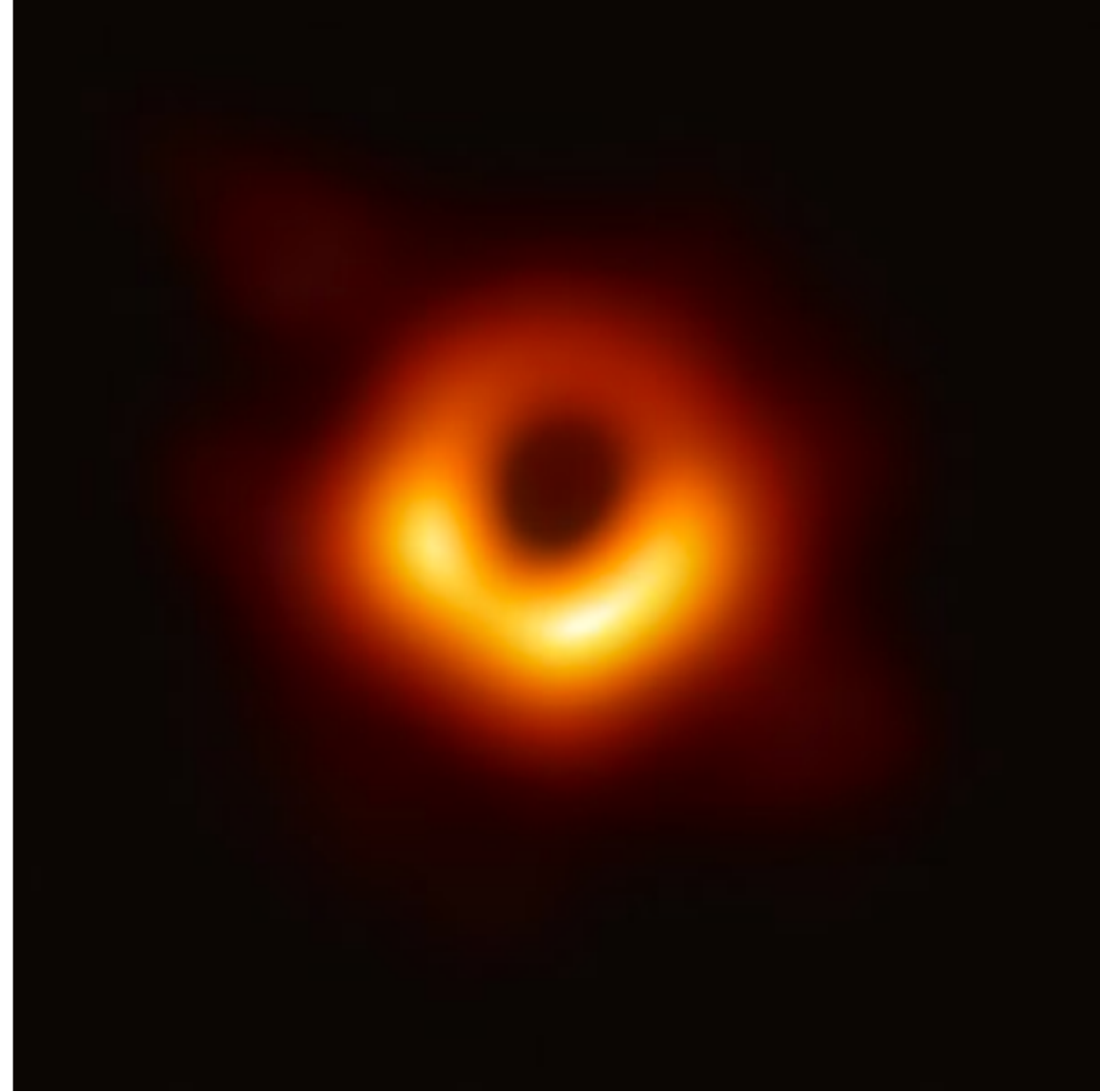


Some phenomenological consequences of neutrino emission from primordial black holes

Yuber F. Perez-Gonzalez



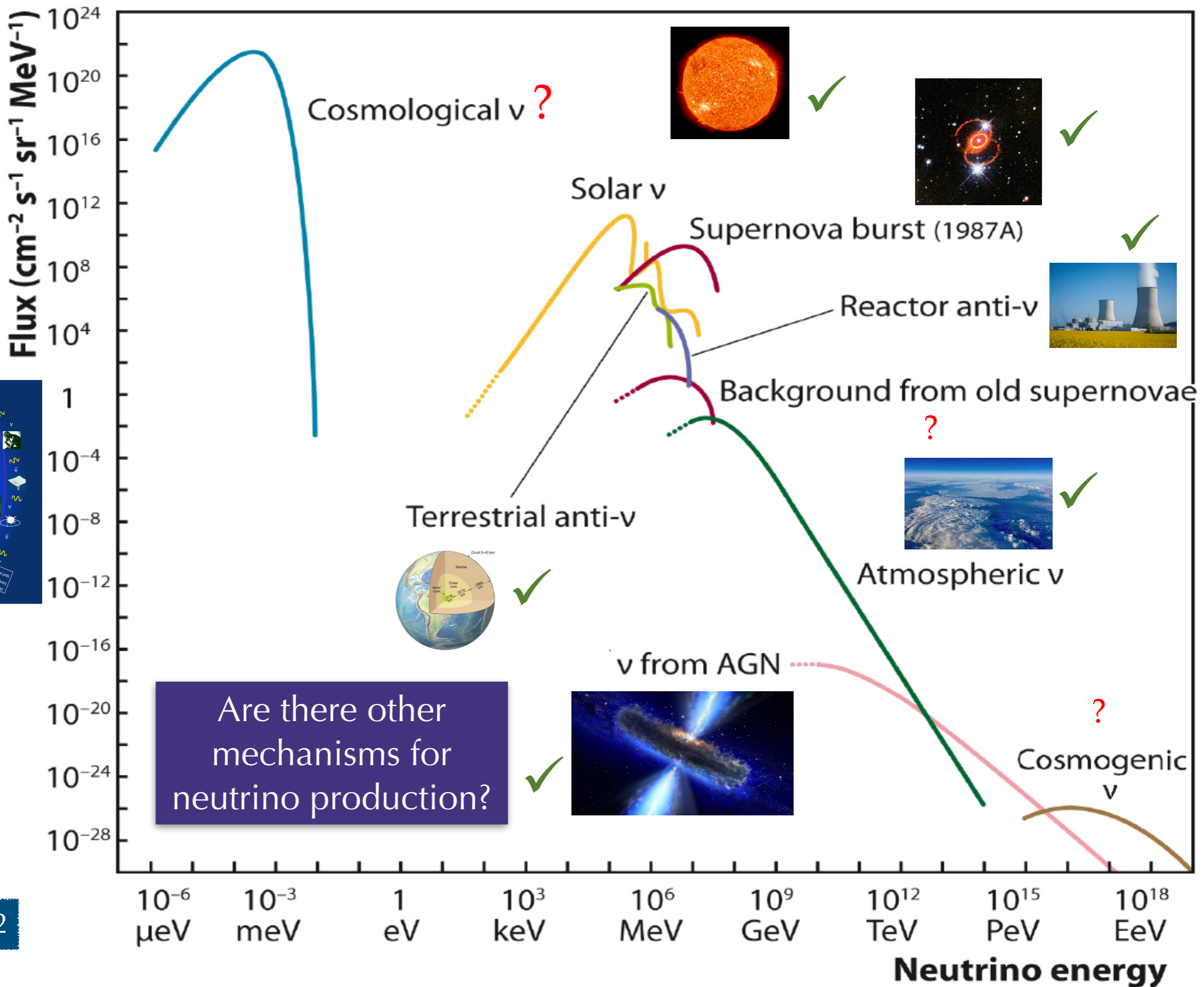
EHT Collaboration

Materia Oscura en Colombia

October 7, 2020



Northwestern



Katz et.al., 2012

Primordial Black Holes (PBH)

Astrophysical Black Holes

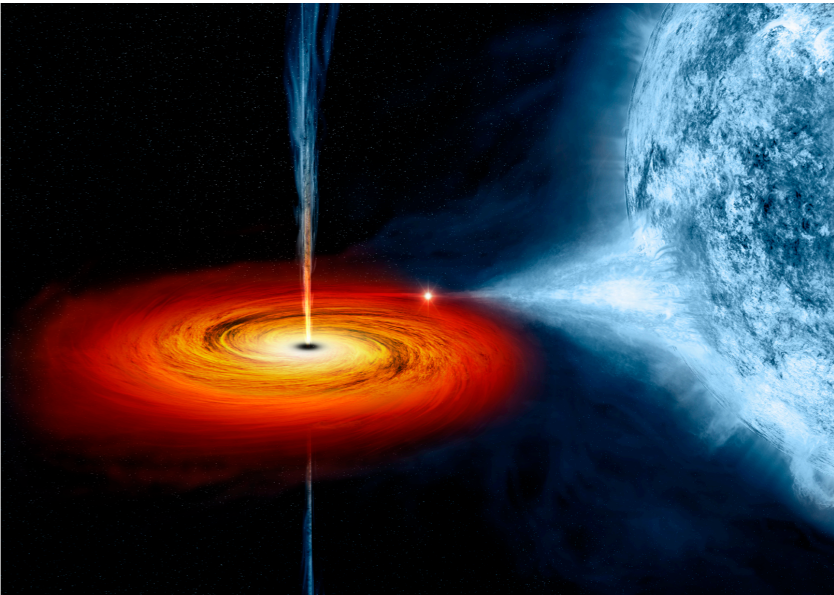
$$M \gtrsim 3M_{\odot}$$

Lighter Black Holes

Large densities

$$r_S = 2GM$$

$$M_i \sim \frac{t}{G} \sim 10^{15} \left(\frac{t}{10^{-23} \text{ s}} \right) \text{ g}$$



Primordial Black Holes

Formation

- ❖ Bubble collisions
- ❖ Pressure reduction
- ❖ Collapse of density fluctuations

Quantum effects are important

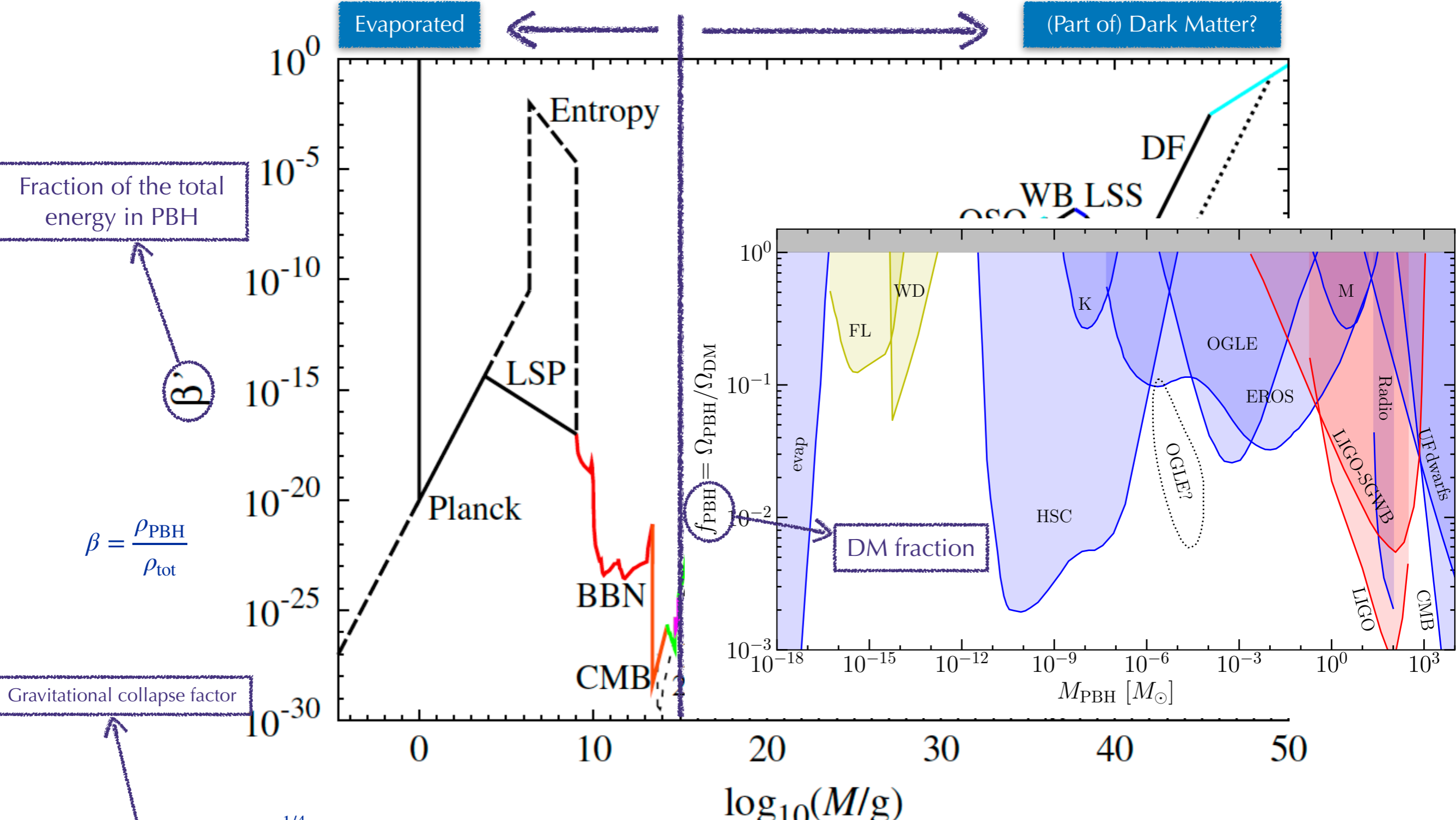
$$r_S \sim \lambda_C$$

Black Holes evaporate by thermal emission

Hawking, 1975

Carr et al, [2002.12778](#)

Primordial Black Holes (PBH)



Fraction of the total energy in PBH

β

$$\beta = \frac{\rho_{\text{PBH}}}{\rho_{\text{tot}}}$$

Gravitational collapse factor

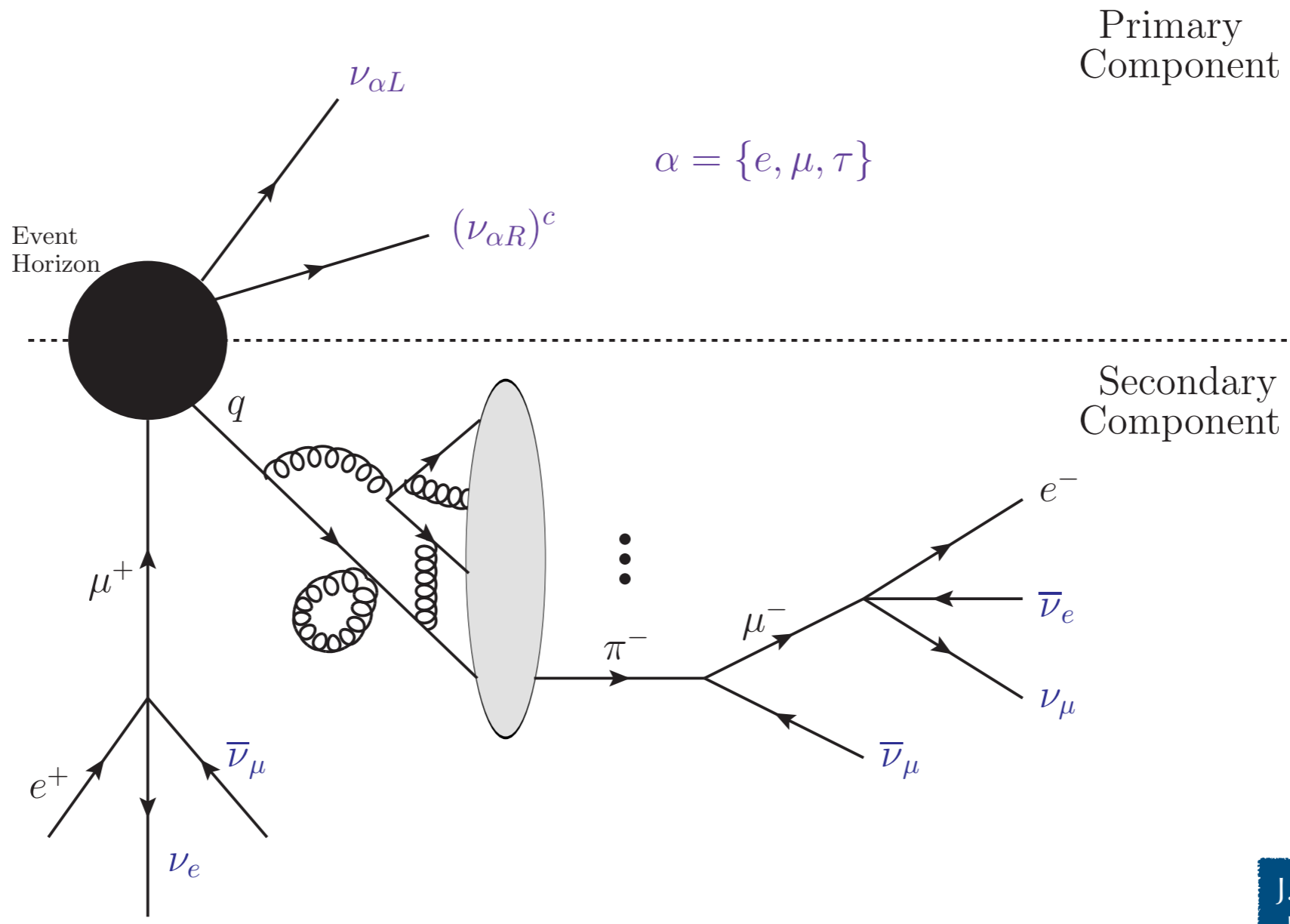
$$\beta' = \beta \gamma^{1/2} \left(\frac{g_{*f}}{106.75} \right)^{1/4}$$

Carr et al. [2002.12778](#)

B. Kavanagh
[10.5281/zenodo.3538999](https://zenodo.org/record/3538999)

Neutrino emission in the SM

B. Carr, 1976

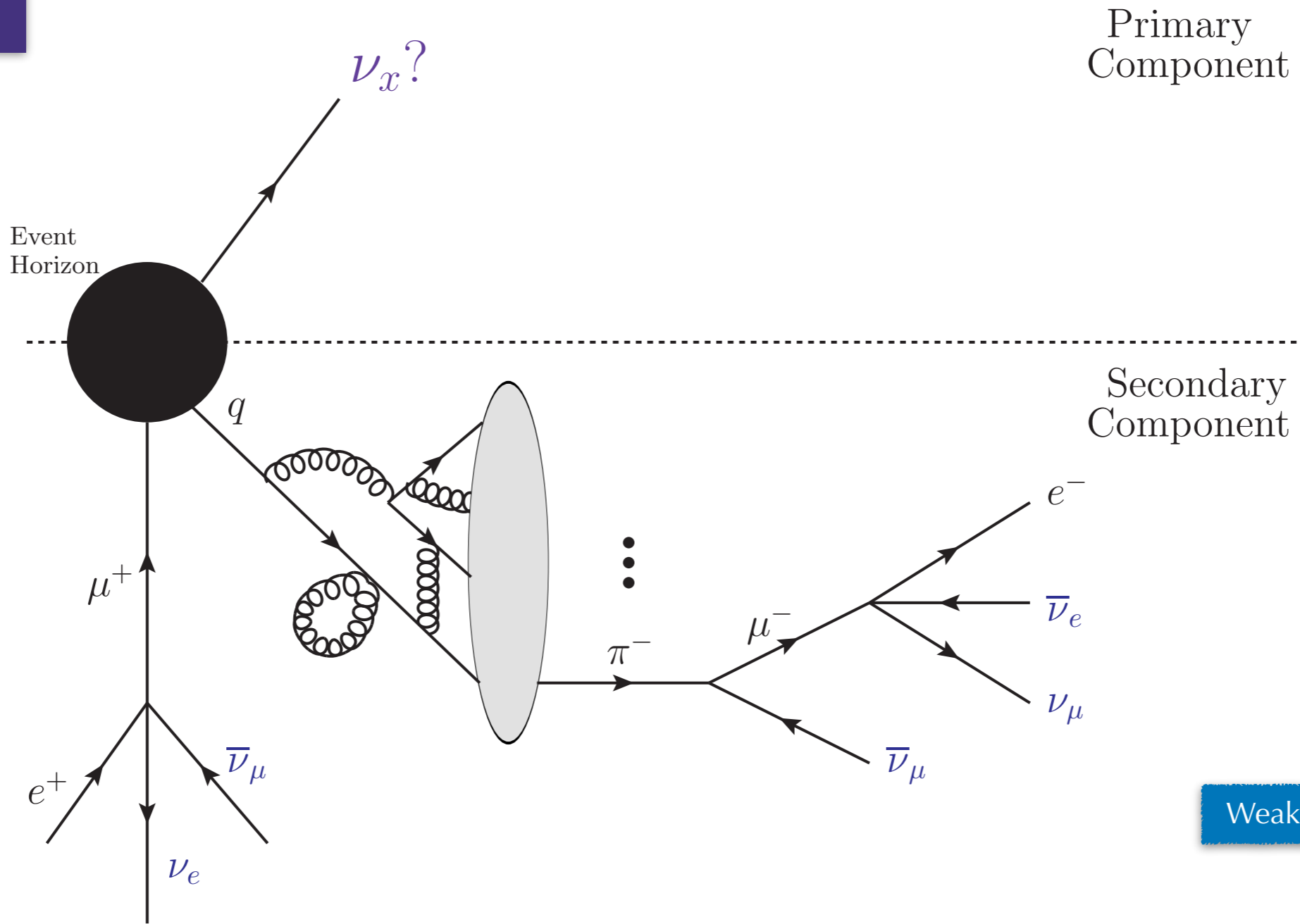


Constraints on the diffuse neutrino flux

J. MacGibbon, 1991
 F. Halzen, 9502268
 Bugaev, 0005295

What is the state of the emitted neutrino?

Neutrinos are massive



Primary Component

Secondary Component

Weak Interaction

What is the state of the emitted neutrino?

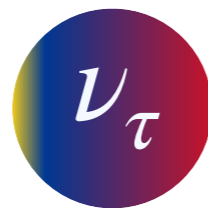
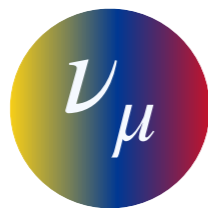
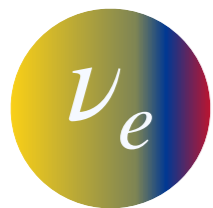
Weak interactions

$$n \rightarrow p^+ + e^- + \bar{\nu}_e$$

Interaction mediated by
a gauge boson

Associated with a
charged lepton

Flavor eigenstate



Hawking Effect

$$\langle 0_- | b_i^\dagger b_i | 0_- \rangle = \Gamma_{lm} \left[\exp(E_a/T_{\text{BH}} + 1) \right]^{-1}$$

Particle definition in a curved
spacetime is observer dependent

No associated
production of a charged
lepton

Mass eigenstate



Neutrino instantaneous spectrum

$$\frac{d^2 N_\nu}{dp dt} = \sum_{a=1,2,3} \frac{g_a^N \sigma_{\text{abs}}^\nu(M, p, m_a) p^2}{2\pi^2 \exp[E_a(p)/T] + 1} \frac{p}{E_a(p)}$$

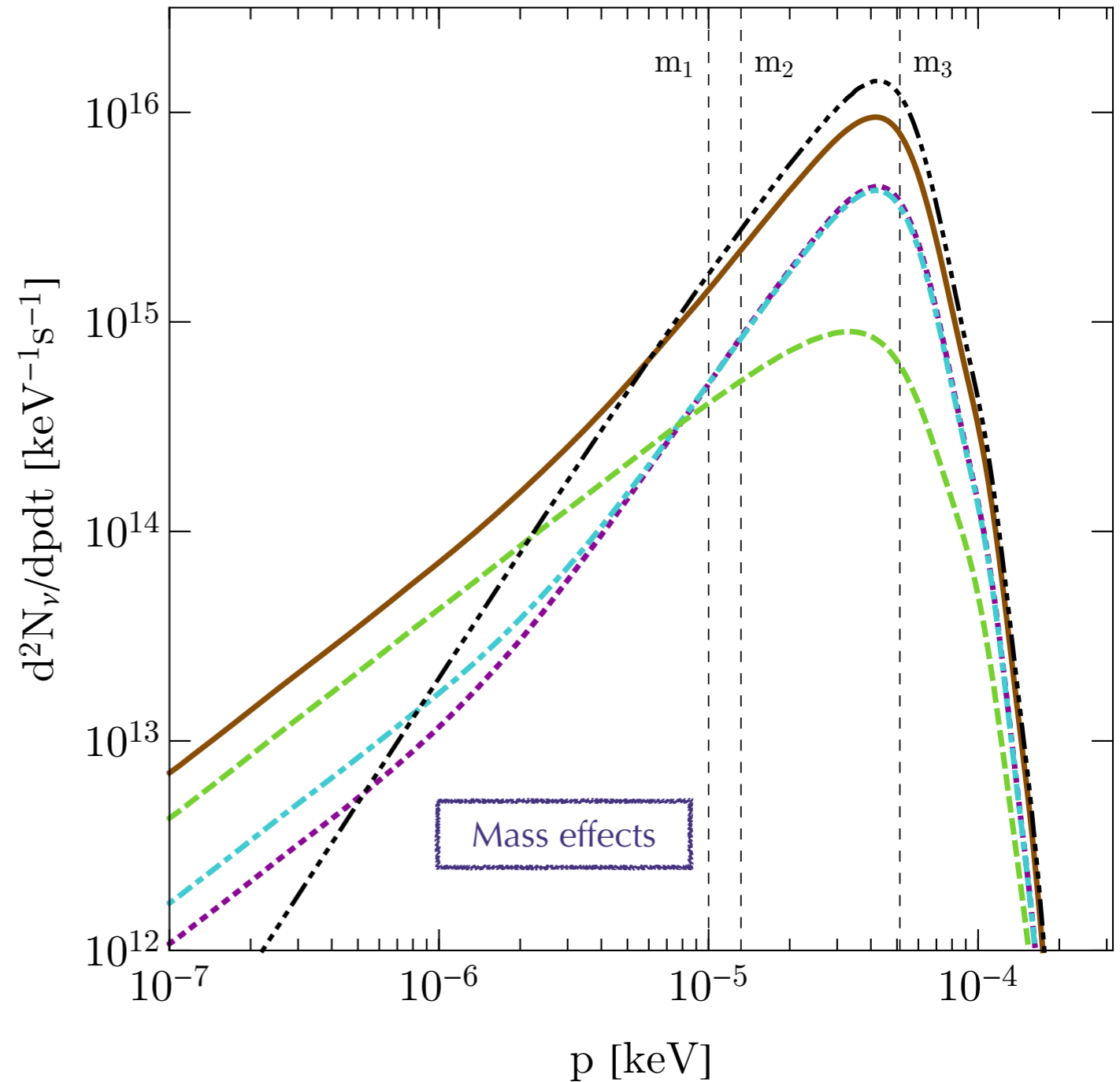
Degrees of freedom Absorption cross section

BH Temperature

$$T = \frac{1}{8\pi GM}$$

$$\approx 1.06 \left(\frac{10^{13} \text{ g}}{M} \right) \text{ GeV}$$

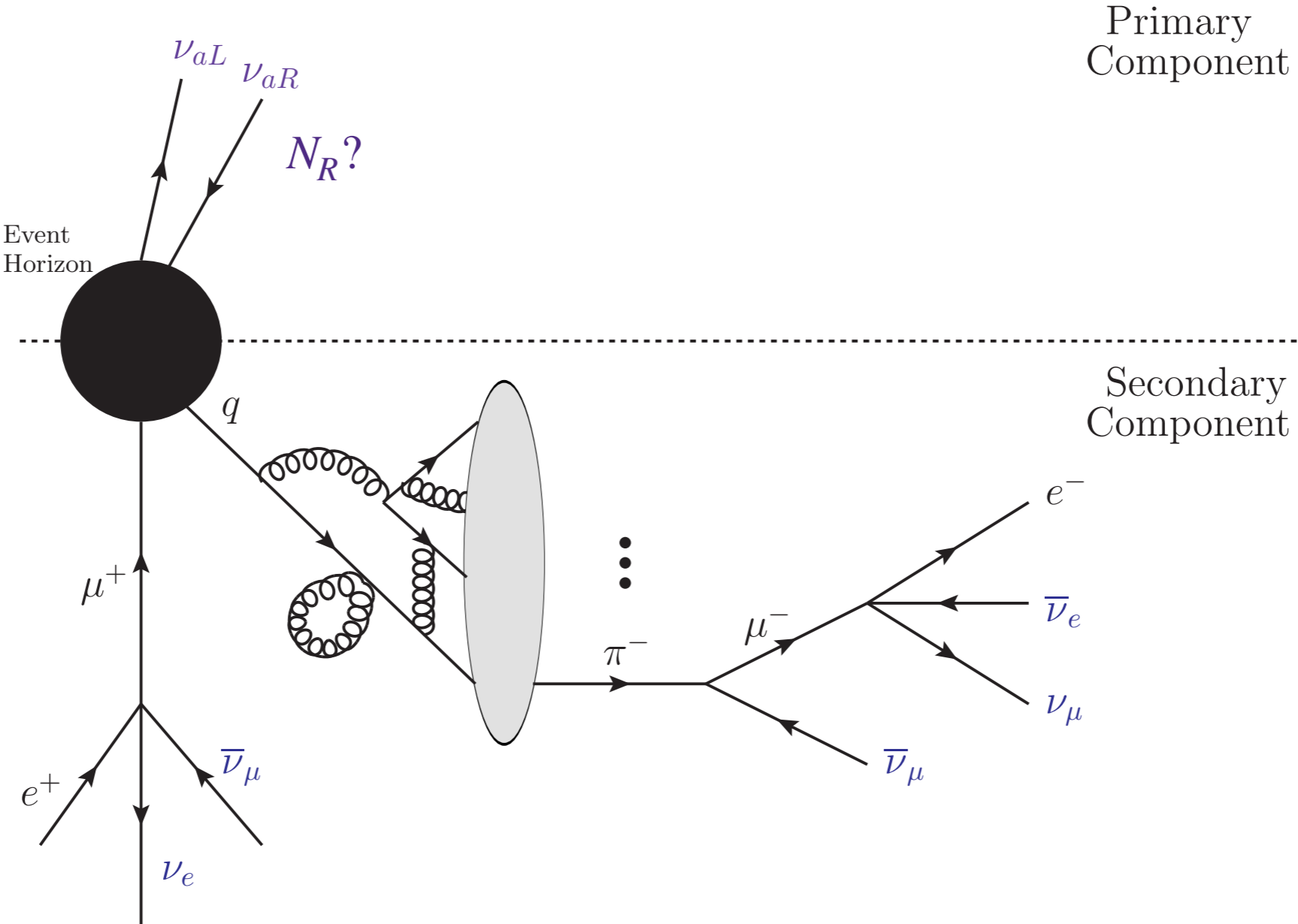
$M = 10^{24} \text{ g}$ ($T_{\text{BH}} \approx 0.01 \text{ eV}$)



$m_0 = 0.01 \text{ eV}$

Schwarzschild BH

Dirac vs Majorana



Majorana neutrinos

Dirac vs Majorana

Dirac neutrinos

$$\sigma_{\text{abs}}^{\nu}(+1/2) = \sigma_{\text{abs}}^{\nu}(-1/2)$$

Unruh, 1976

No helicity suppression

Production of light RH neutrinos!

Cecilia Lunardini, YFPG
JCAP08(2020)014
arXiv:1912:07864

Majorana neutrinos

Heavy RH neutrinos



PBH-driven Leptogenesis

Yamada and Iso, 1610.02586
Morrison et al, 1812.10606
Baldes et al, 2004.14773

Cecilia Lunardini, YFPG
and Jessica Turner,
in preparation

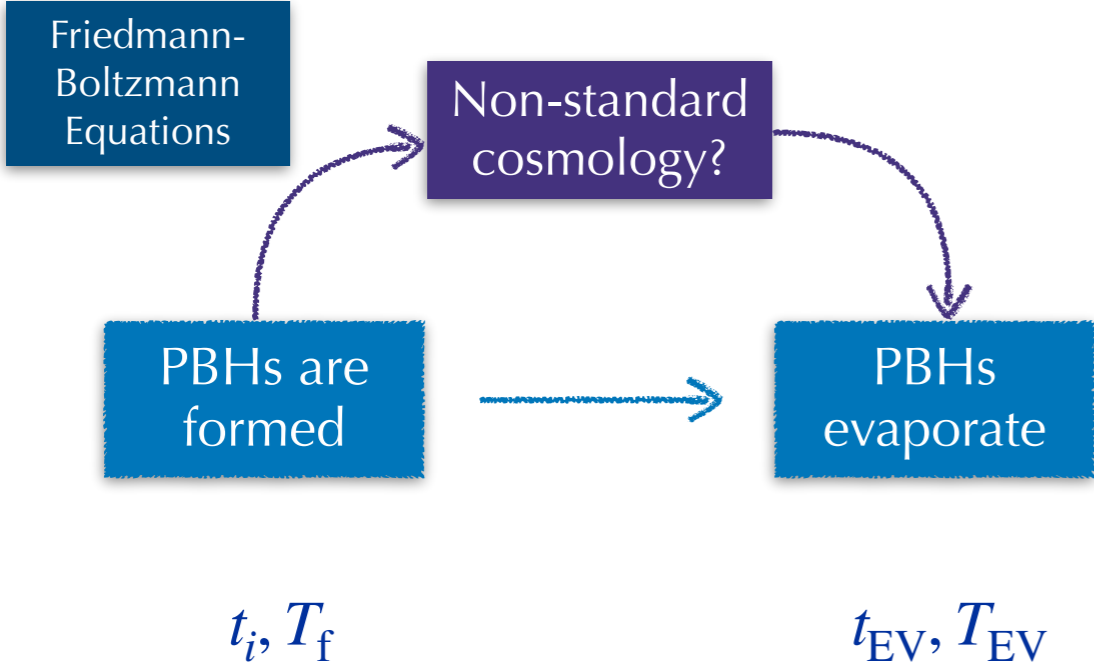
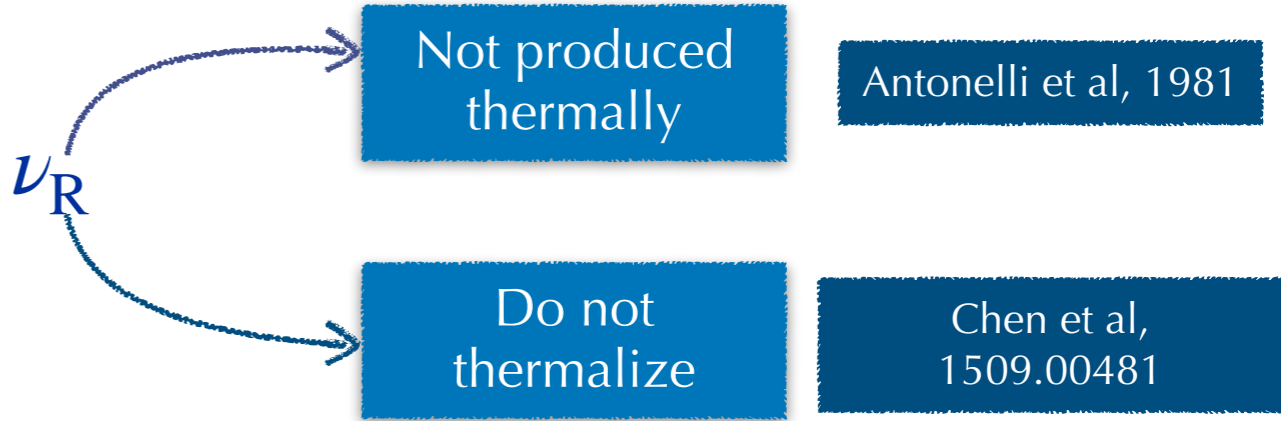
Dirac Neutrinos

Cecilia Lunardini,
YFPG
JCAP08(2020)014
arXiv:1912:07864

Constraints in the Dirac neutrino case

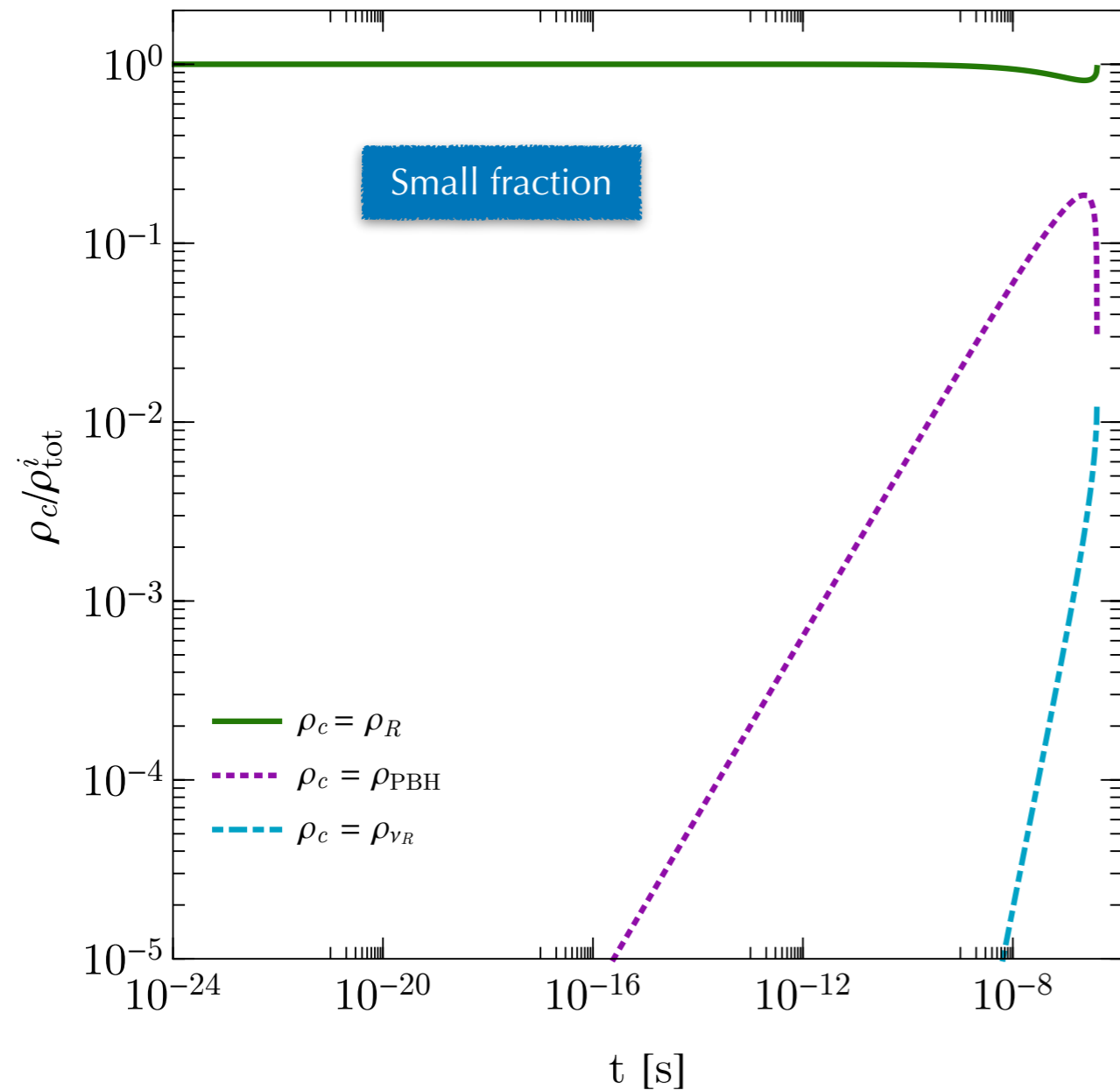
Let us consider the minimal extension

$$\mathcal{L}_Y = - Y_\nu^{ab} \bar{L}_L^a \widetilde{H} \nu_{bR}$$

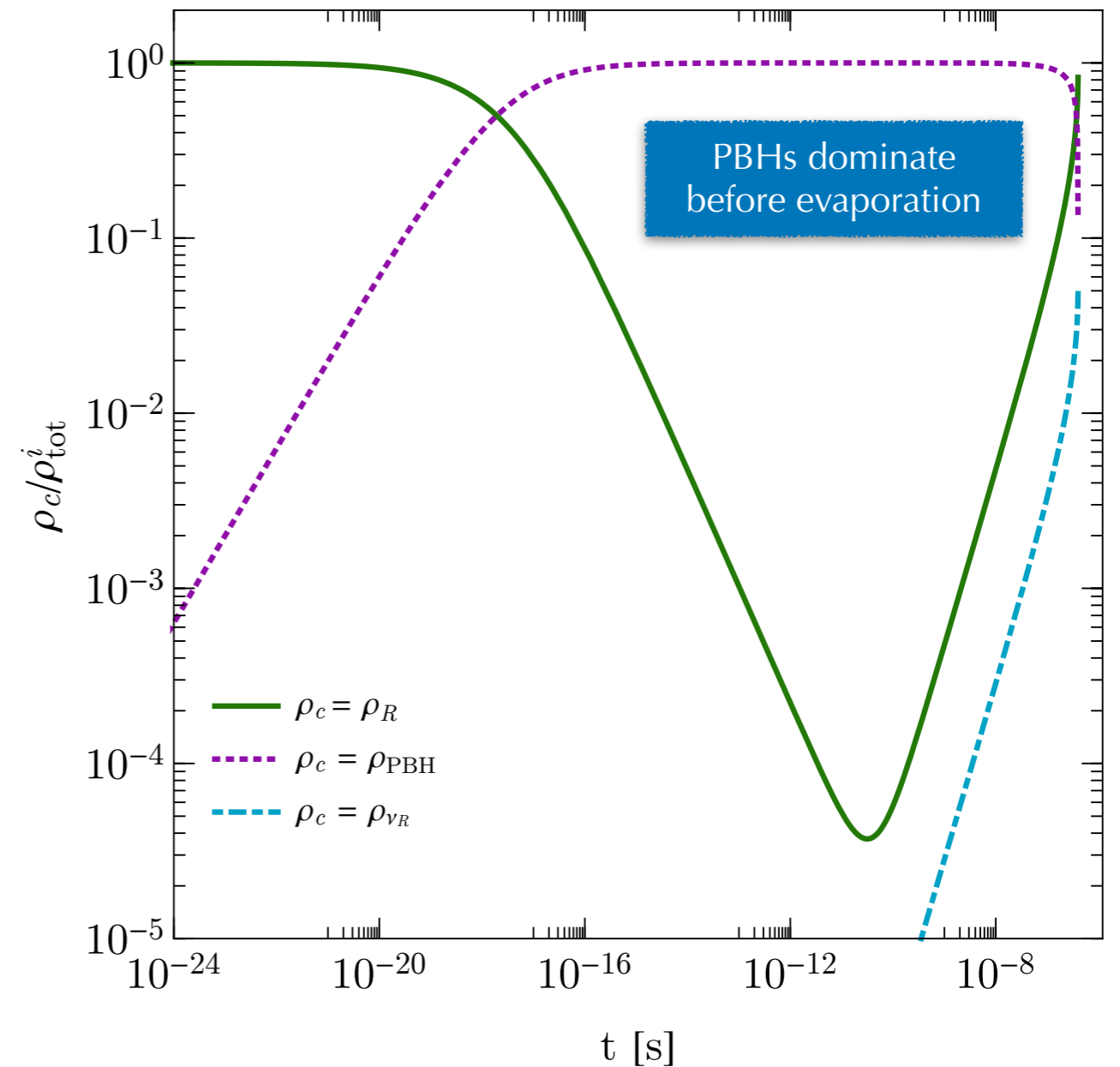


Constraints in the Dirac neutrino case

$$M_i = 10^7 \text{ g}, \beta' = 10^{-13}$$



$$M_i = 10^7 \text{ g}, \beta' = 10^{-7}$$



Initial fraction

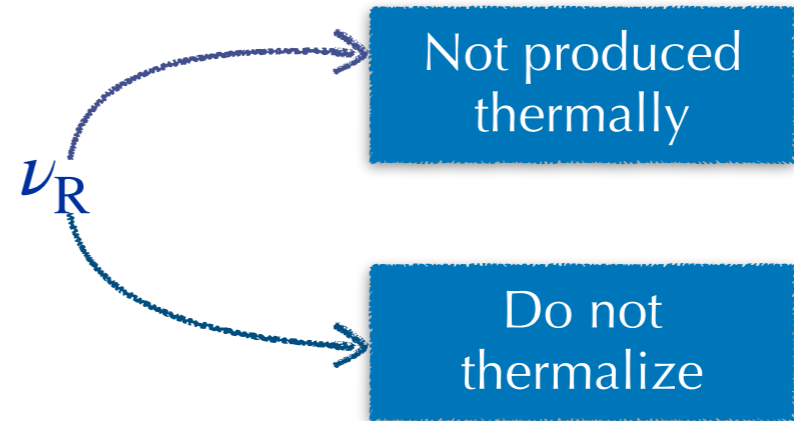
$$\beta' \gtrsim 2.5 \times 10^{-14} \left(\frac{g_*(T_f)}{106.75} \right)^{-\frac{1}{4}} \left(\frac{M_i}{10^8 \text{ g}} \right)^{-1} \left(\frac{\epsilon_D(M_i)}{15.35} \right)^{\frac{1}{2}}$$

Hooper et al, 1905.01301

Constraints in the Dirac neutrino case

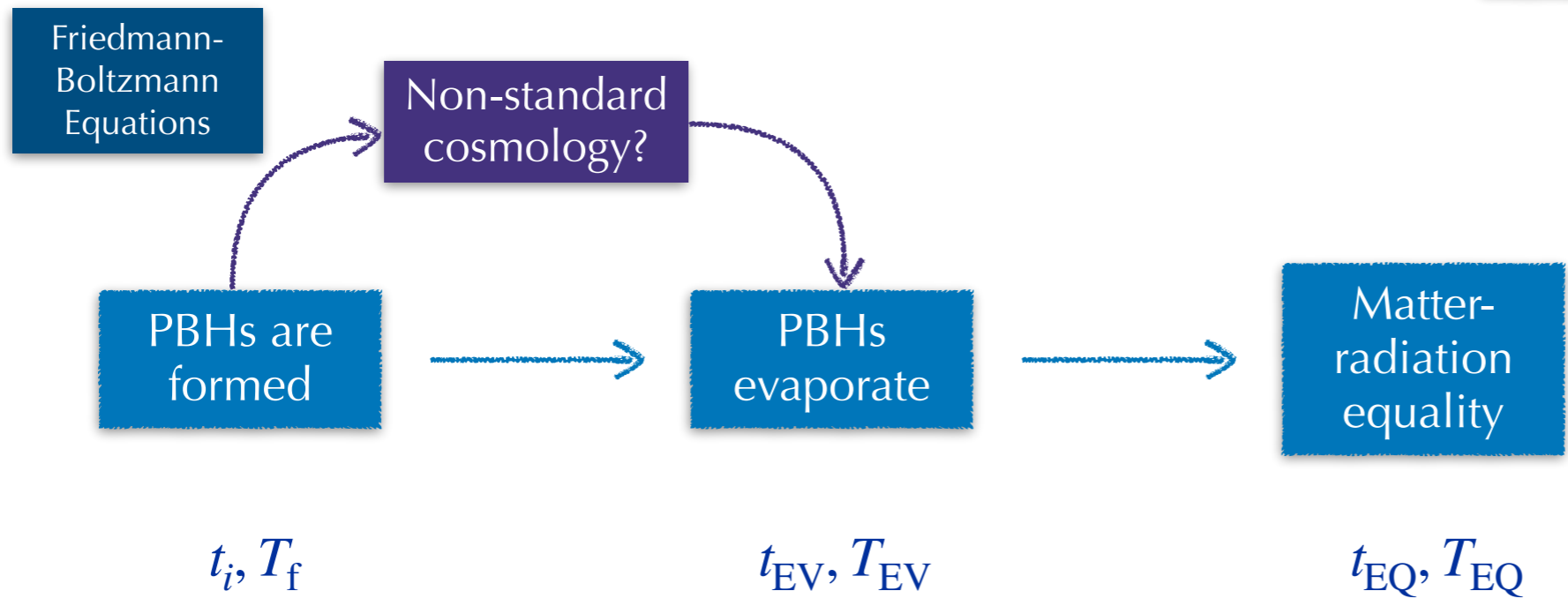
Let us consider the minimal extension

$$\mathcal{L}_Y = - Y_\nu^{ab} \bar{L}_L^a \widetilde{H} \nu_{bR}$$



Antonelli et al, 1981

Chen et al, 1509.00481

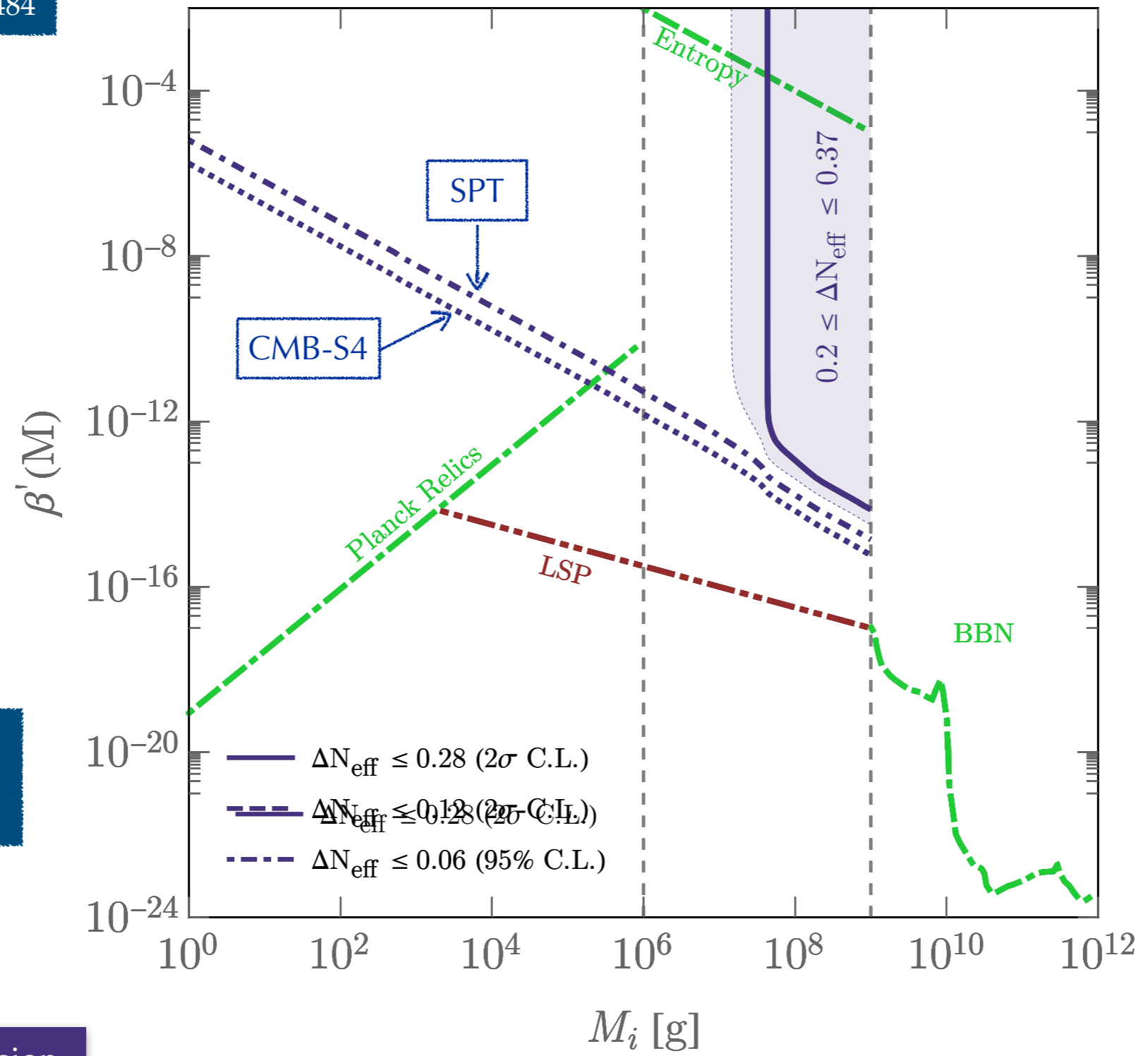


$$\Delta N_{\text{eff}} = \left\{ \frac{8}{7} \left(\frac{4}{11} \right)^{-\frac{4}{3}} + N_{\text{eff}}^{\text{SM}} \right\} \frac{\rho_{\nu_R}(T_{\text{EV}})}{\rho_R(T_{\text{EV}})} \left(\frac{g_*(T_{\text{EV}})}{g_*(T_{\text{EQ}})} \right) \left(\frac{g_{*S}(T_{\text{EQ}})}{g_{*S}(T_{\text{EV}})} \right)^{\frac{4}{3}}$$

Hooper et al, 1905.01301

Constraints in the Dirac neutrino case

Green and Liddle, 9903484



Zel'dovich et al, 1977

MacGibbon, 1987
Barrow et al, 1992
Carr et al, 1994

Carr et al, 0912.5297

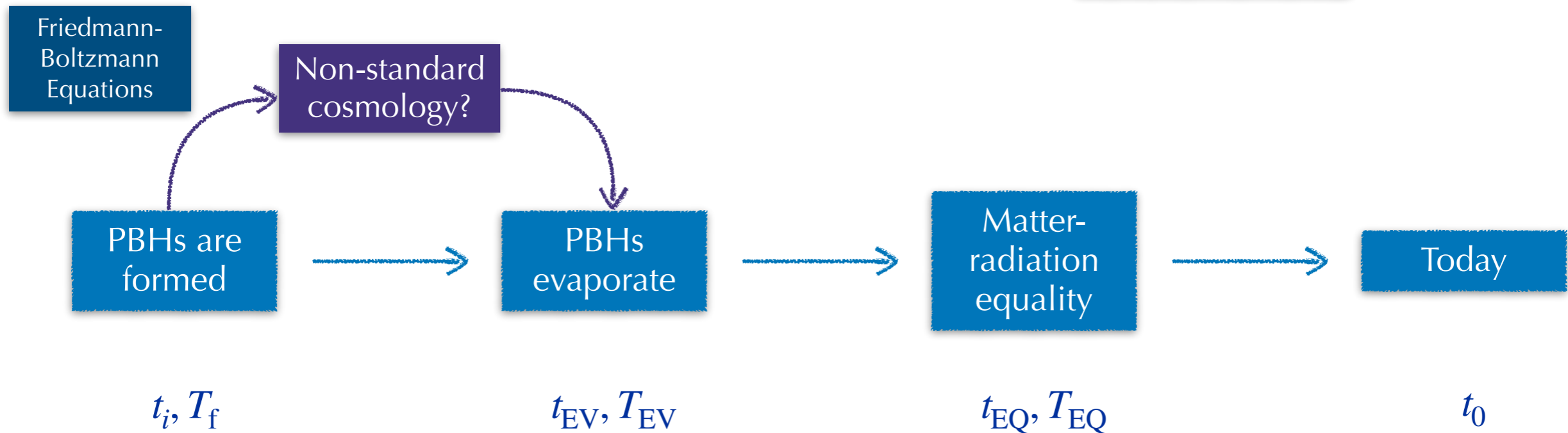
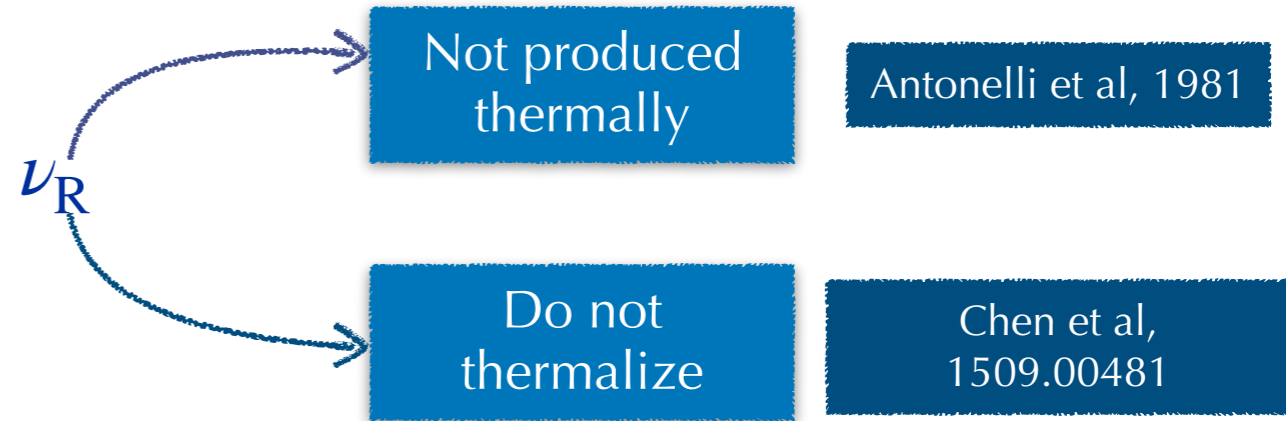
Alleviate the H_0 tension

What happens when $M \rightarrow M_{Pl}$?

Constraints in the Dirac neutrino case

Let us consider the minimal extension

$$\mathcal{L}_Y = - Y_\nu^{ab} \bar{L}_L^a \widetilde{H} \nu_{bR}$$

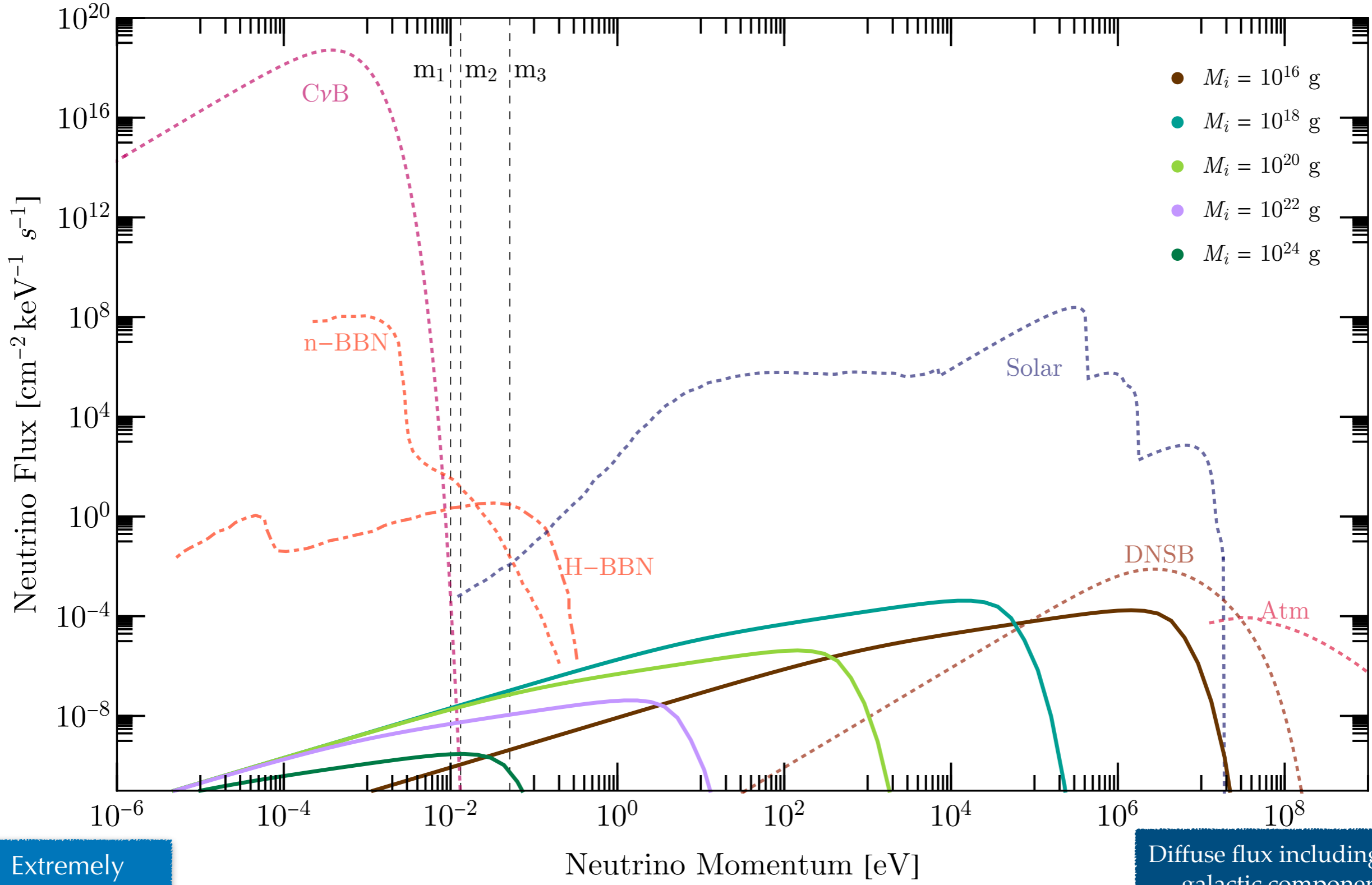


$$\Delta N_{\text{eff}} = \left\{ \frac{8}{7} \left(\frac{4}{11} \right)^{-\frac{4}{3}} + N_{\text{eff}}^{\text{SM}} \right\} \frac{\rho_{\nu_R}(T_{\text{EV}})}{\rho_R(T_{\text{EV}})} \left(\frac{g_*(T_{\text{EV}})}{g_*(T_{\text{EQ}})} \right) \left(\frac{g_{*S}(T_{\text{EQ}})}{g_{*S}(T_{\text{EV}})} \right)^{\frac{4}{3}}$$

Hooper et al, 1905.01301

Diffuse flux from non-evaporating PBHs

Total Majorana neutrino flux, $m_0 = 0.01$ eV, $M_i > M_*$

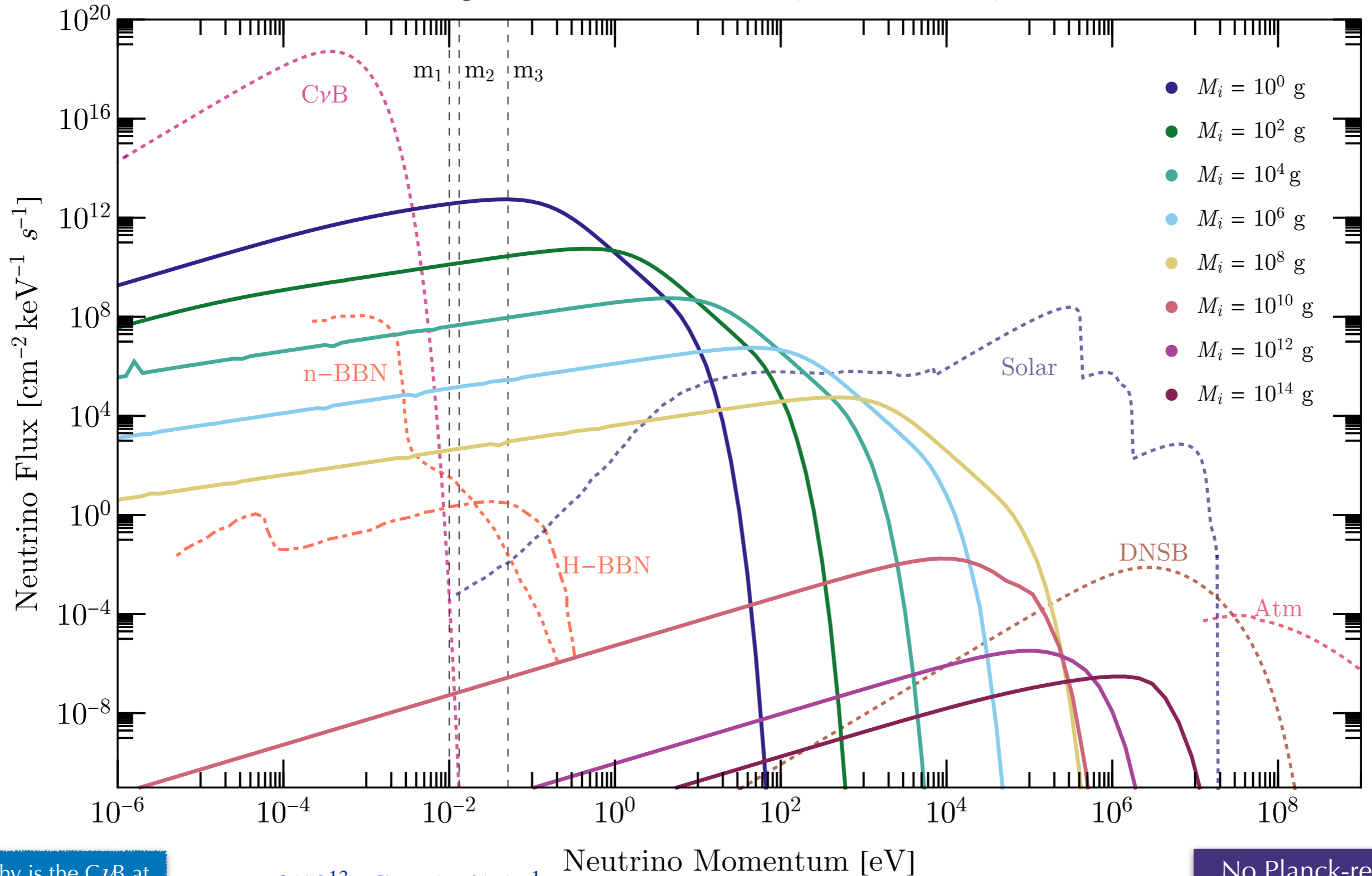


Extremely suppressed flux

Diffuse flux including the galactic component
arXiv: 1912.01014

Diffuse flux of RH neutrinos from PBHs

Total right-helical neutrino flux, $m_0 = 0.01$ eV, $M_i \leq M_*$



Why is the CνB at lower momenta?

$$E_\nu \sim \mathcal{O}(10^{13}) \text{ GeV } (M/1\text{g})^{-1}$$

No Planck-relic constraint

Detection?

Helicity
suppression

$$m_0 = 0.01 \text{ eV}$$

$$\frac{m_\nu}{E_\nu} \sim 10^{-1} \longrightarrow M = 1 \text{ g}$$

PTOLEMY?



$$\Gamma_{C\nu B}^D \sim 40 \text{ [kg - year]}^{-1}$$

$$\Gamma_{\text{PBH}}^\nu \sim 10^{-2} \text{ [kg - year]}^{-1}$$

PBH RH
flux is still
suppressed

Are there other possible
ways to try to detect this
RH neutrino flux?

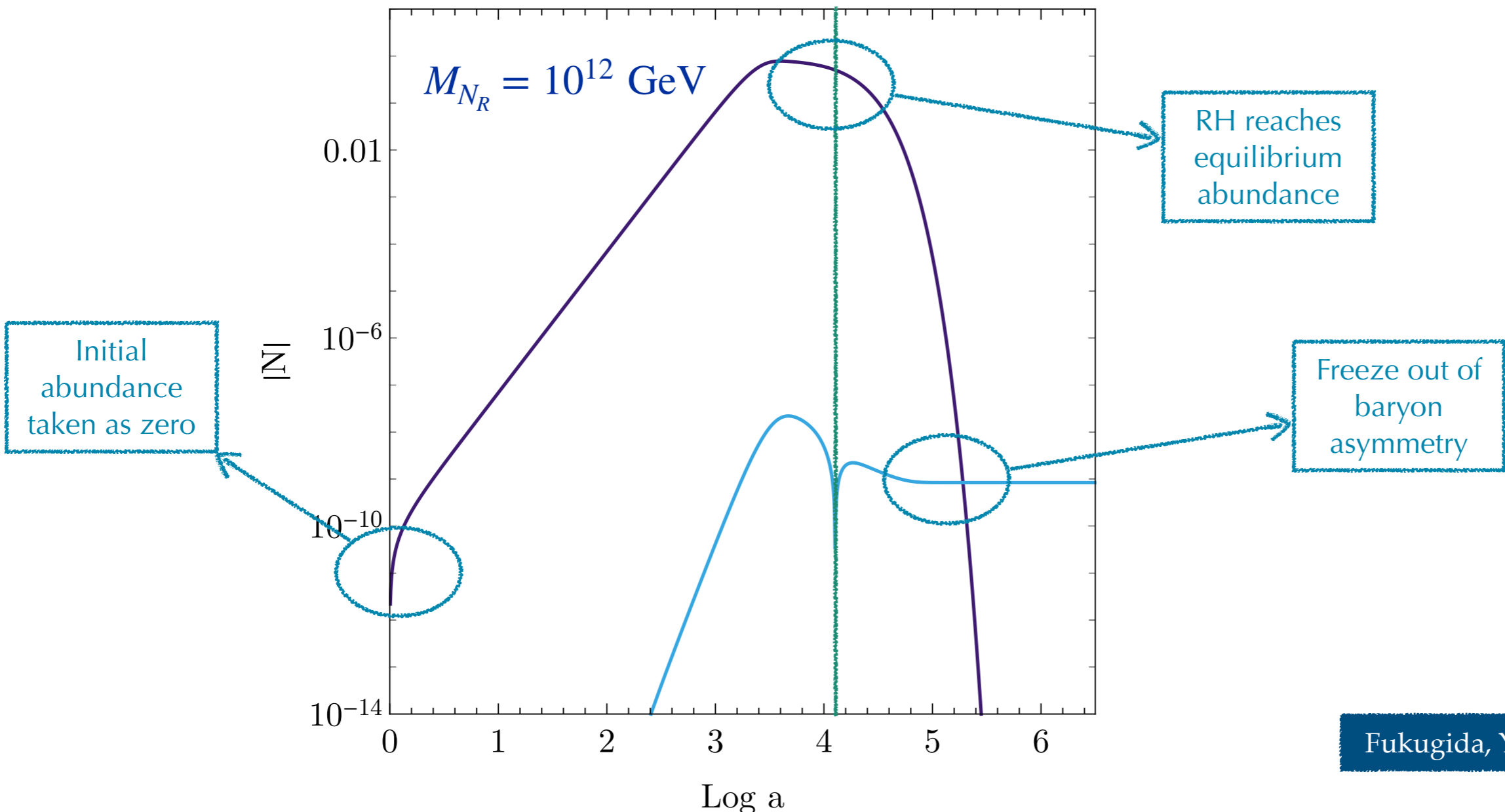
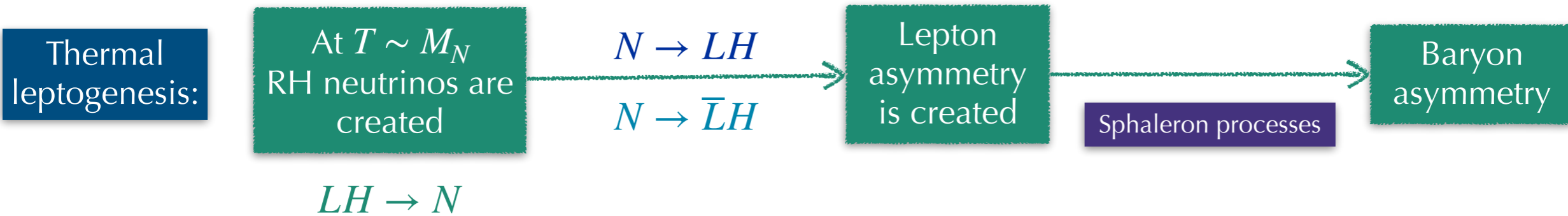


Majorana Neutrinos

Preliminary results

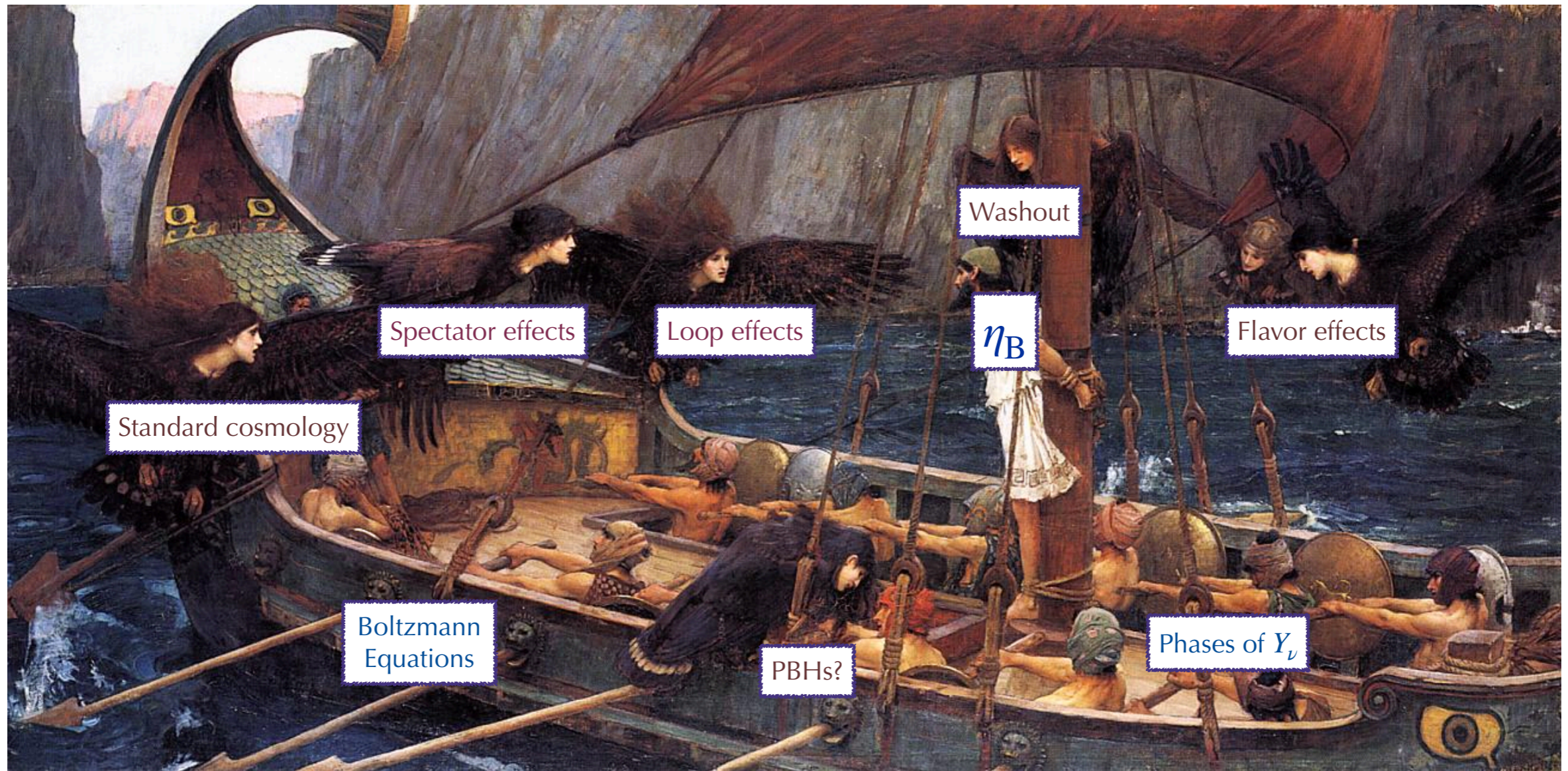
YFPG and Jessica Turner
2010.XXXXX

Leptogenesis in a nutshell



Fukugida, Yanagida, '86

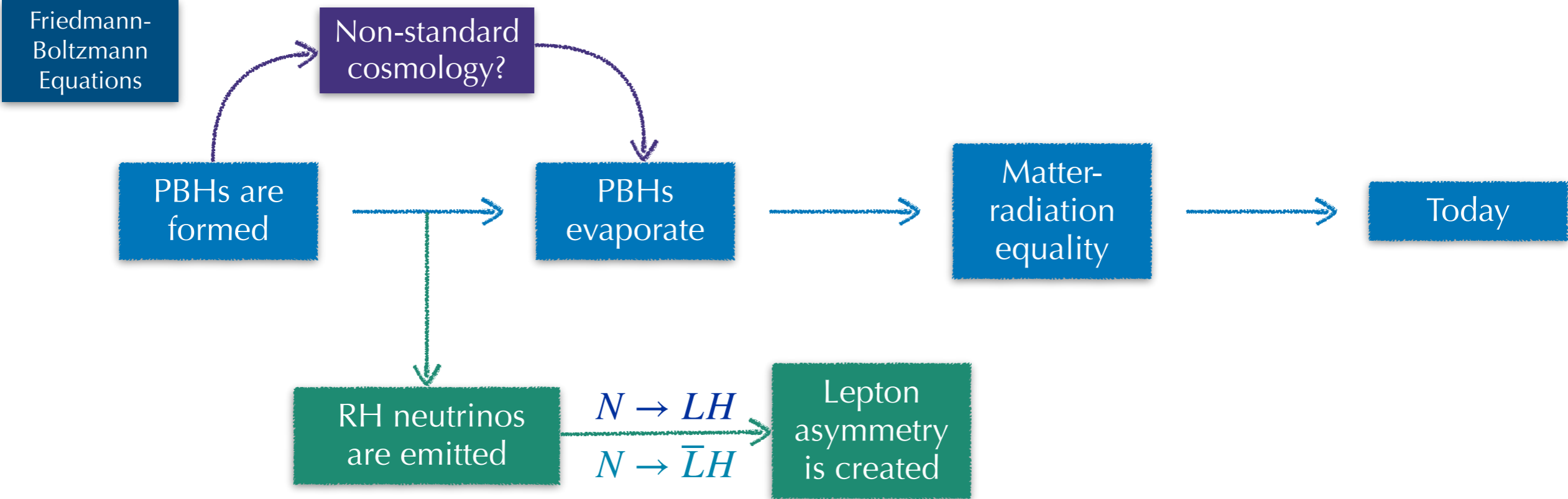
Universal LeptogeneSiS Equation Solver (ULYSSES)



A Graneli, K Moffat, YFPG,
H Schulz and Jessica Turner,
arXiv: [2007.09150](https://arxiv.org/abs/2007.09150)

- ❖ Leptogenesis via decays and resonant leptogenesis
- ❖ Easy parallelization
- ❖ Rapid evaluation
- ❖ Multidimensional scan of the parameter space

PBH-driven Leptogenesis



Three scenarios

- A. PBH evaporate before RH are thermally produced
- B. Evaporation happens at more or less at the same time
- C. PBH create a RH density when the plasma would not be able to do.

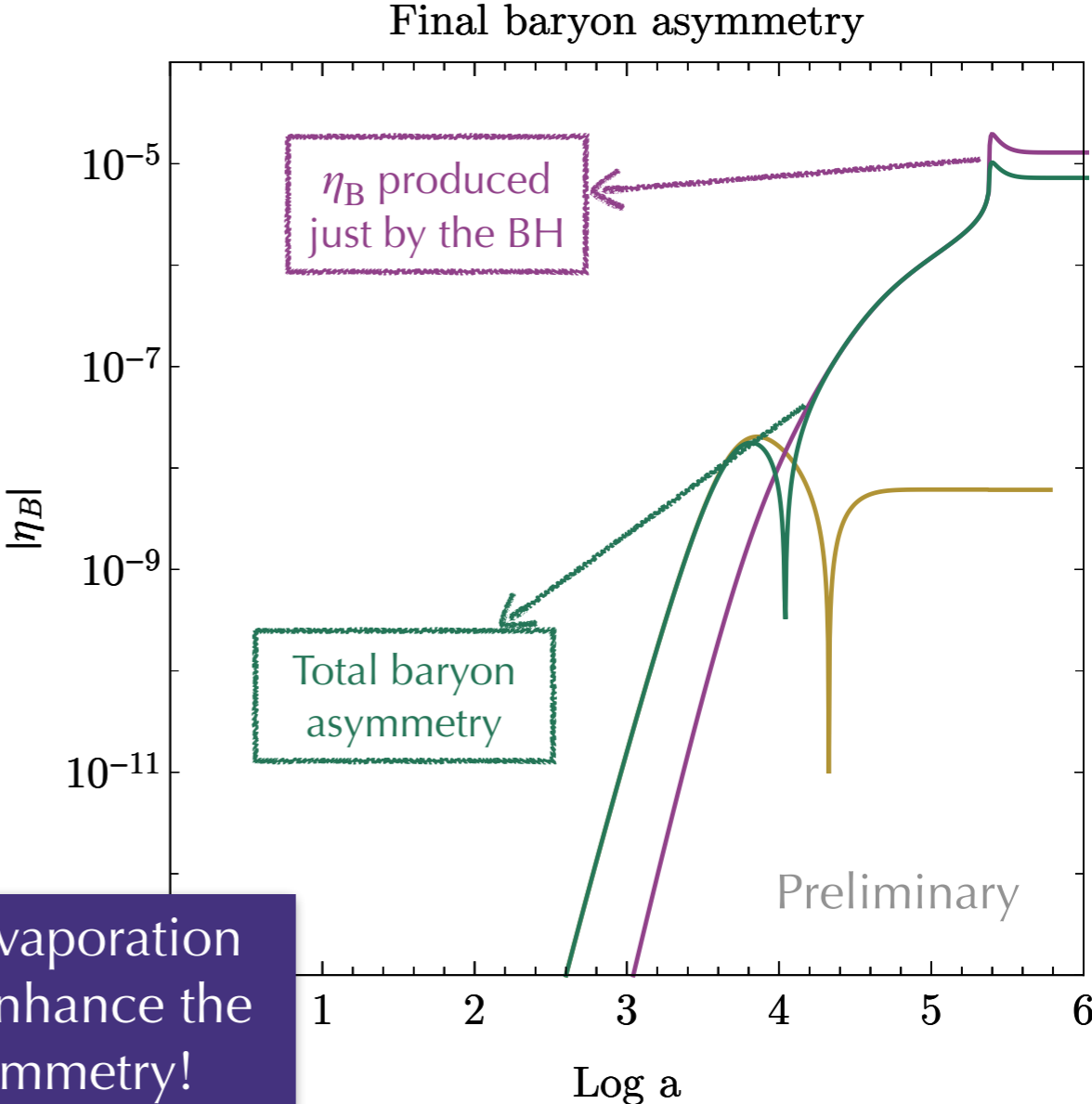
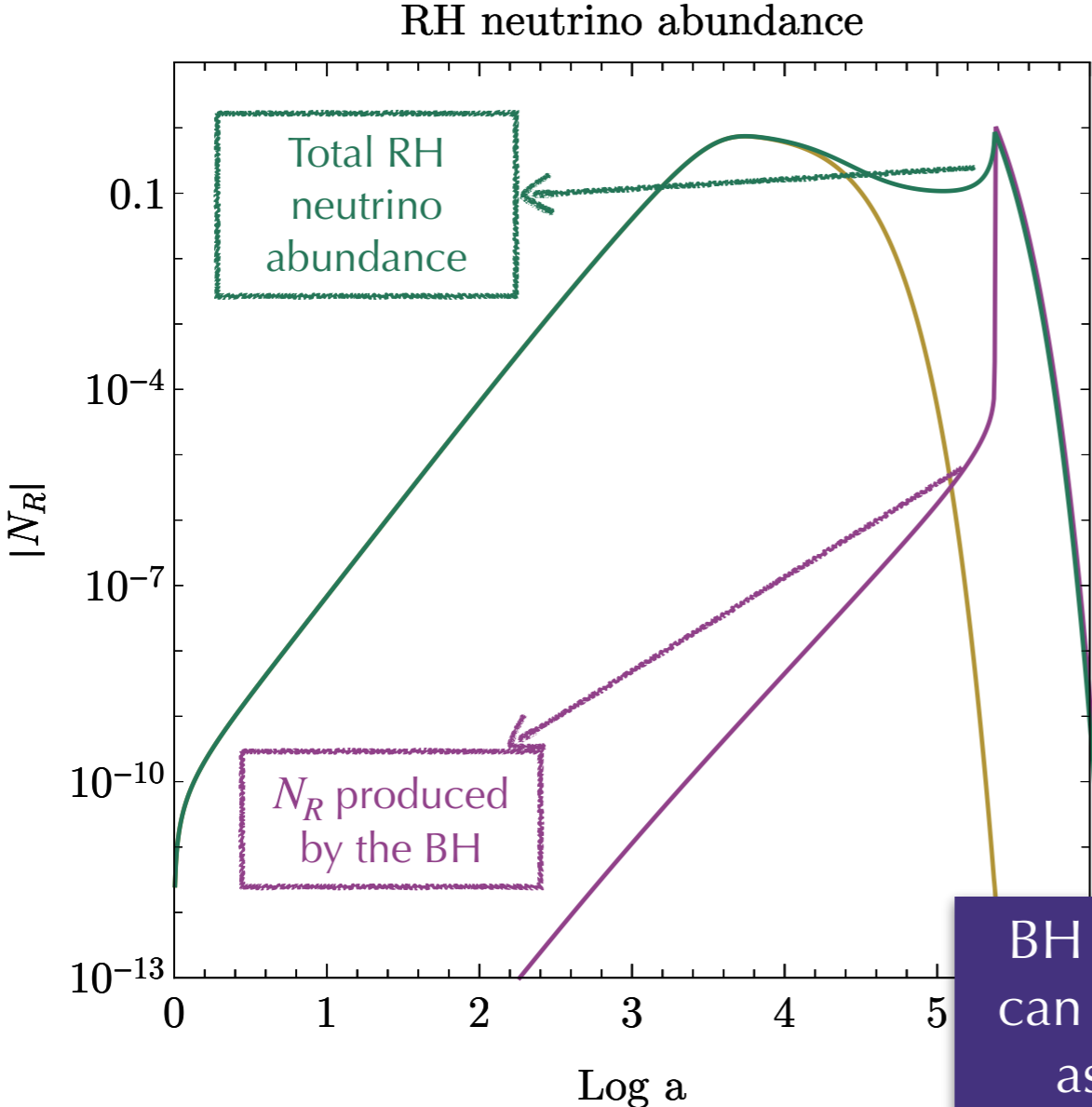
PBH-driven Leptogenesis

- A. PBH evaporate before RH are thermally produced
- B. Evaporation happens at more or less at the same time

The thermal plasma would equilibrate the over abundance created by the PBHs

$$M_i = 0.3 \text{ g}$$

$$\beta'_i = 10^{-3}$$



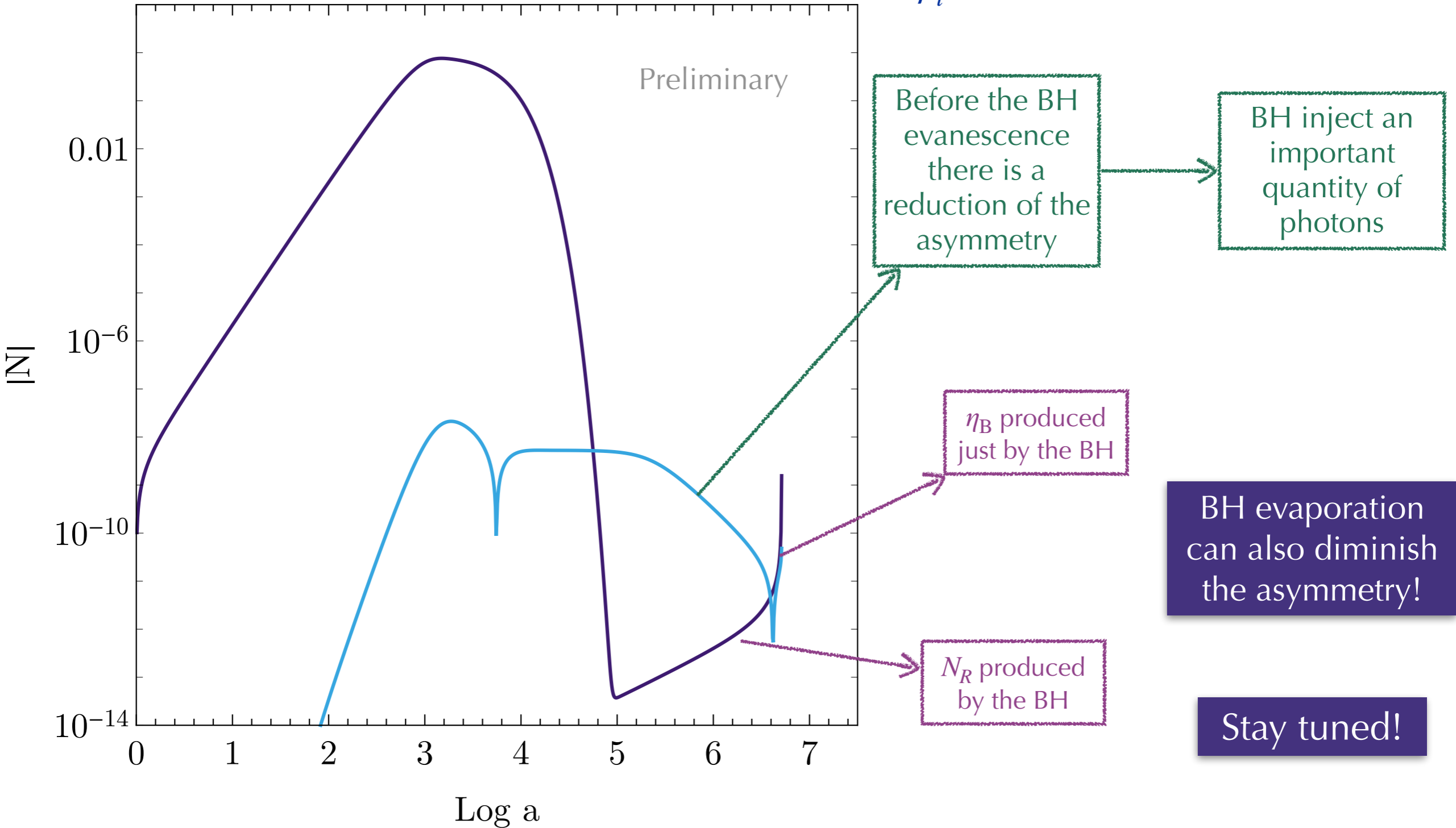
BH evaporation can enhance the asymmetry!

PBH-driven Leptogenesis

C. PBH create a RH density when the plasma would not be able to do.

$$M_i = 3 \text{ g}$$

$$\beta'_i = 10^{-3}$$



Conclusions

- The PBH evaporation depends on whether neutrinos are Dirac or Majorana particles
- In the Dirac scenario, there is not a helicity suppression of the emission of right-handed neutrinos
- We derived a constraint on the initial PBH fraction given the measurement of N_{eff} by Planck
- For certain values, it is possible to ease the Hubble measurement tension
- The diffuse flux of RH neutrinos can be large, but more careful analysis on its possible detection should be performed
- Preliminary results show that black hole evaporation can enhance or diminish the baryon asymmetry in the leptogenesis scenario.
- What about spinning black holes?