

# Deviations to Tri-Bi-Maximal mixing in the limit of $\mu \leftrightarrow \tau$ symmetry

Diana Carolina Rivera Agudelo

in collaboration with A. Pérez-Lorenzana and Sergio Tostado

*3rd COMHEP  
USC, Cali, Colombia*

December 4, 2018



# Outline

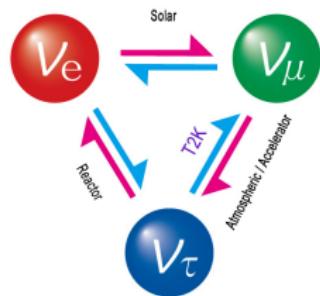
- 1 Introduction
- 2  $\mu - \tau$  symmetry in the mixing matrix and deviations from TBM pattern
- 3 TBM limit in the mass matrix
- 4 Results and conclusions

# Introduction: Neutrino mixing

Standard parametrization:  $c_{ij} \equiv \cos \theta_{ij}$ ,  $s_{ij} \equiv \sin \theta_{ij}$

$$U_{PMNS} = \begin{pmatrix} c_{12}c_{13} & s_{12}c_{13} & s_{13}e^{-i\delta_{CP}} \\ -s_{12}c_{23} - c_{12}s_{23}s_{13}e^{i\delta_{CP}} & c_{12}c_{23} - s_{12}s_{23}s_{13}e^{i\delta_{CP}} & s_{23}c_{13} \\ s_{12}s_{23} - c_{12}c_{23}s_{13}e^{i\delta_{CP}} & -c_{12}s_{23} - c_{23}s_{12}s_{13}e^{i\delta_{CP}} & c_{23}c_{13} \end{pmatrix} P_\nu.$$

- $P_\nu = \text{diag} \left[ 1, e^{-i\frac{\beta_1}{2}}, e^{-i\frac{\beta_2}{2}} \right]$
- Global analysis ( $1\sigma$ ):  
 $\theta_{12}/^\circ = 34.5^{+1.2}_{-1.0}$ ,  $\theta_{13}/^\circ = 8.45^{+0.16}_{-0.14}$ ,  
 $\theta_{23}/^\circ = 47.7^{+1.2}_{-1.7}$ ,  $\delta_{cp}/\pi = 1.32^{+0.21}_{-0.15}$
- Majorana nature  $\iff 0\nu\beta\beta$



Neutrino oscillation between three generations

Credits T2K coll.

[1]

<sup>1</sup>P.F. de Salas et al., PLB 782, 633 2018

# $\mu - \tau$ symmetry and TBM pattern

- If  $\theta_{13} = 0$ ,  $\theta_{23} = \pi/4$

$$U_{\mu-\tau} = \begin{pmatrix} c_{12} & s_{12} & 0 \\ \frac{-s_{12}}{\sqrt{2}} & \frac{c_{12}}{\sqrt{2}} & \frac{-1}{\sqrt{2}} \\ \frac{-s_{12}}{\sqrt{2}} & \frac{c_{12}}{\sqrt{2}} & \frac{1}{\sqrt{2}} \end{pmatrix},$$

- $c_{12}(s_{12})$  stands for  $\sin \theta_{12}(\cos \theta_{12})$
- When  $\sin^2 \theta_{12} = 1/3 \rightarrow$  Tri-Bi-Maximal<sup>2</sup> (TBM) mixing

## TBM main features

Null reactor angle, no CP violation, solar angle

$$\theta_{12} = 30^\circ$$

Disfavored by experiments

Maximal atmospheric angle  $\theta_{23} = 45^\circ$  is still allowed



Credits nobelprize.org

<sup>2</sup>T. Fukuyama and H. Nishiura, arXiv:hep-ph/9702253 [hep-ph].

# Deviations from TBM pattern

However, small deviations from  $\mu \leftrightarrow \tau$  symmetry can be adopted. Let us consider

- $U_{PMNS} = U_{TBM} U_{\text{Corr}}$ , where

## Correction matrix

$$U_{\text{Corr}} = U_{ij}(\phi, \sigma) \text{ Diag} \left( 1, e^{-i\frac{\alpha_1}{2}}, e^{-i\frac{\alpha_2}{2}} \right),$$

Two possibilities since  $U_{12}$  is trivial:

$$U_{13} = \begin{pmatrix} \cos \phi & 0 & \sin \phi e^{-i\sigma} \\ 0 & 1 & 0 \\ -\sin \phi e^{i\sigma} & 0 & \cos \phi \end{pmatrix}, \quad U_{23} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & \cos \phi & \sin \phi e^{-i\sigma} \\ 0 & -\sin \phi e^{i\sigma} & \cos \phi \end{pmatrix}$$

$\phi$  and  $\sigma$  encode the deviation from TBM mixing.  $\alpha_1$  and  $\alpha_2$  are aimed to account for the Majorana nature.

# Relations between experimental parameters and PMNS entries

Mixing angles:

$$\sin^2 \theta_{12} = \frac{|U_{e2}|^2}{1 - |U_{e3}|^2}, \quad \sin^2 \theta_{23} = \frac{|U_{\mu 3}|^2}{1 - |U_{e3}|^2}, \quad \sin^2 \theta_{13} = |U_{e3}|^2,$$

CP is obtained from Jarlskog invariant

$$\begin{aligned} J_{CP} &= \text{Im} [U_{e1} U_{\mu 2} U_{e2}^* U_{\mu 1}^*] \\ &= (1 - s^2 \theta_{13}) \sqrt{s^2 \theta_{13} s^2 \theta_{12} s^2 \theta_{23} (1 - s^2 \theta_{12}) (1 - s^2 \theta_{23})} \sin \delta_{CP}. \end{aligned}$$

## Case I: 1-3 Rotation

$$U_{PMNS} = U_{TBM} U_{13}(\phi, \sigma) \text{Diag} \left( 1, e^{-i\frac{\alpha_1}{2}}, e^{-i\frac{\alpha_2}{2}} \right)$$

$$\sin^2 \theta_{12} = \frac{1}{3 - 2 \sin^2 \phi},$$

$$\sin^2 \theta_{23} = \frac{1}{2} \left( 1 + \frac{\sqrt{3} \sin 2\phi \cos \sigma}{3 - 2 \sin^2 \phi} \right),$$

$$\sin^2 \theta_{13} = \frac{2}{3} \sin^2 \phi.$$

### Dirac CP

$$\sin \delta_{CP} = - \frac{(2 + \cos 2\phi) \sin \sigma}{[(2 + \cos 2\phi)^2 - 3 \sin^2 2\phi \cos^2 \sigma]^{1/2}}.$$

### Majorana phases

$$\beta_1 = \alpha_1, \quad \beta_2 = \alpha_2 + 2(\sigma - \delta_{CP})$$

## Case II: 2-3 Rotation

$$U_{PMNS} = U_{TBM} U_{23}(\phi, \sigma) \text{Diag} \left( 1, e^{-i\frac{\alpha_1}{2}}, e^{-i\frac{\alpha_2}{2}} \right)$$

$$\sin^2 \theta_{12} = 1 - \frac{2}{3 - \sin^2 \phi},$$

$$\sin^2 \theta_{23} = \frac{1}{2} \left( 1 - \frac{\sqrt{6} \sin 2\phi \cos \sigma}{3 - \sin^2 \phi} \right),$$

$$\sin^2 \theta_{13} = \frac{1}{3} \sin^2 \phi$$

### Dirac CP

$$\sin \delta_{CP} = - \frac{(5 + \cos 2\phi) \sin \sigma}{[(5 + \cos 2\phi)^2 - 24 \sin^2 2\phi \cos^2 \sigma]^{1/2}}$$

### Majorana phases

$$\beta_1 = \alpha_1, \quad \beta_2 = \alpha_2 + 2(\sigma - \delta_{CP})$$

## $\mu - \tau$ symmetric limit in the mass matrix

- Since  $M_\nu = U_\nu \text{diag}(m_1, m_2, m_3) U_\nu^T$ ,  $U_\nu = U_{TBM}$  implies

$$|m_{e\mu}| = |m_{e\tau}| \quad \text{and} \quad m_{\mu\mu} = m_{\tau\tau}$$

Deviations can be accommodated in a correction matrix *via*

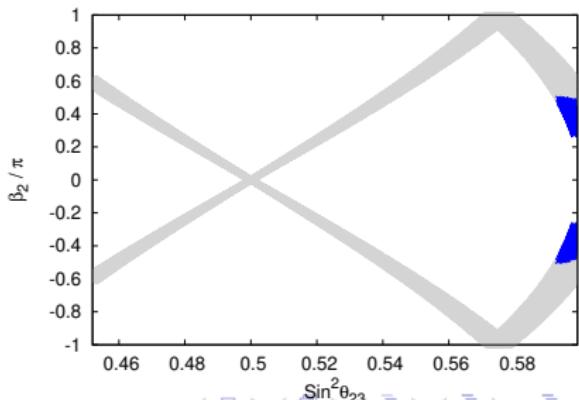
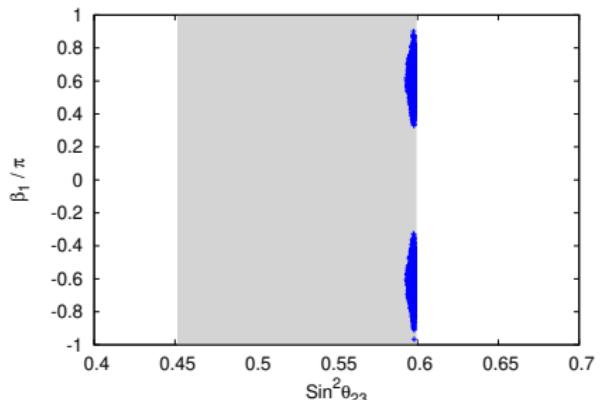
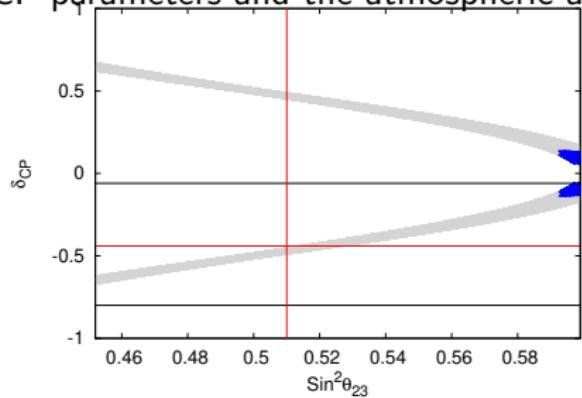
$$M_\nu = M_{\mu-\tau} + \delta M(\hat{\delta}, \hat{\epsilon}).$$

### Breaking parameters

$$\begin{aligned}\hat{\delta} &= \frac{\sum_i (U_{ei} U_{\tau i} - U_{ei} U_{\mu i}) m_i}{\sum_i U_{ei} U_{\mu i} m_i} , \\ \hat{\epsilon} &= \frac{\sum_i (U_{\tau i} U_{\tau i} - U_{\mu i} U_{\mu i}) m_i}{\sum_i U_{\mu i} U_{\mu i} m_i} .\end{aligned}$$

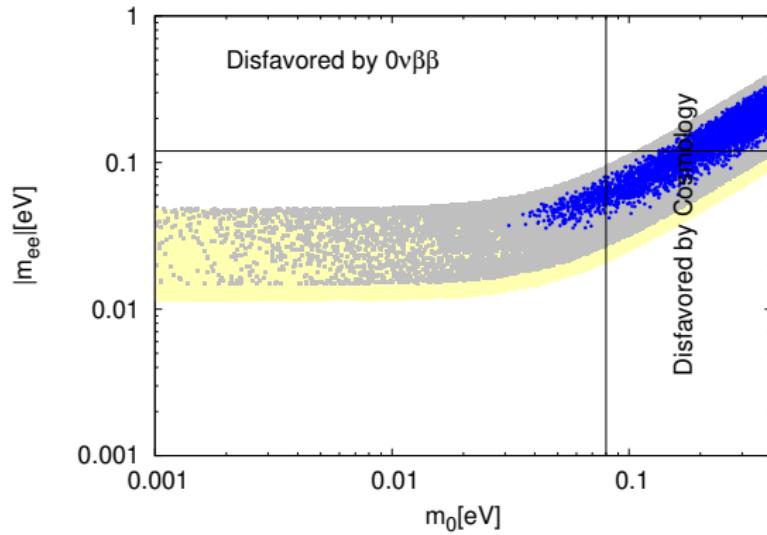
# Results Case I ( $U_{13}$ )

Correlation between CP parameters and the atmospheric angle. Case  $\alpha_2 = 0$ .



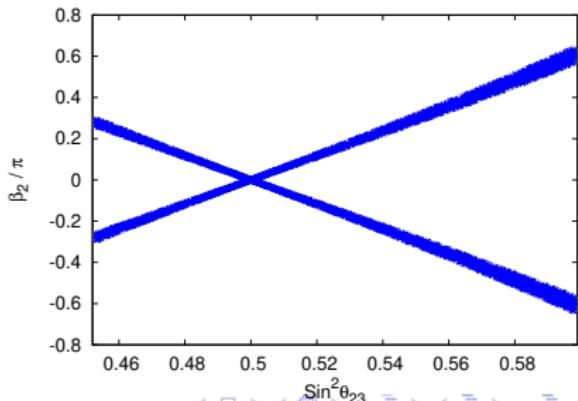
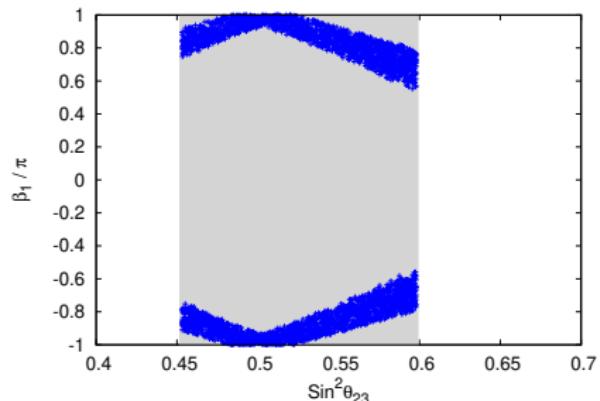
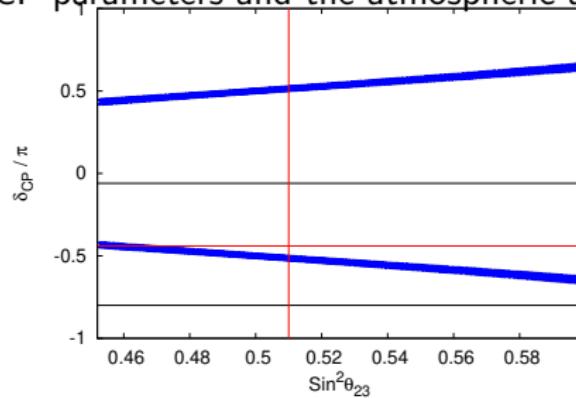
# Results Case I ( $U_{13}$ )

Neutrinoless double beta decay. Case  $\alpha_2 = 0$ .



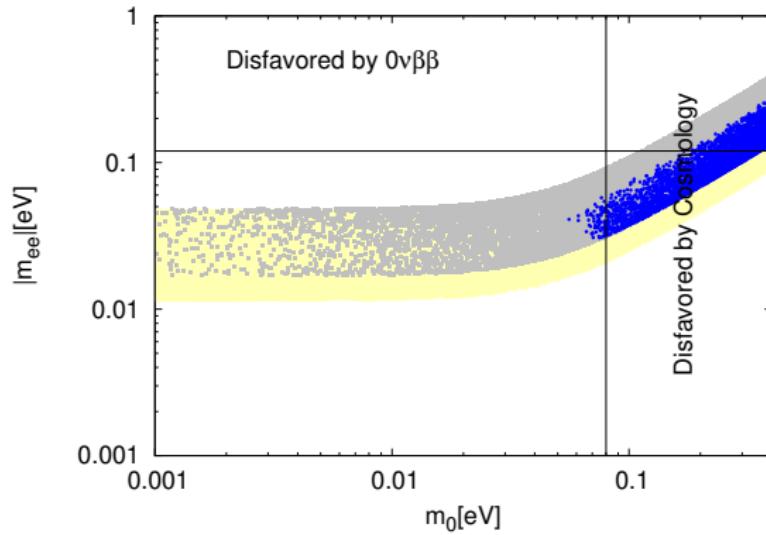
## Results Case II ( $U_{23}$ )

Correlation between CP parameters and the atmospheric angle. Case  $\alpha_2 = 0$ .



## Results Case II ( $U_{23}$ )

Neutrinoless double beta decay. Case  $\alpha_2 = 0$ .



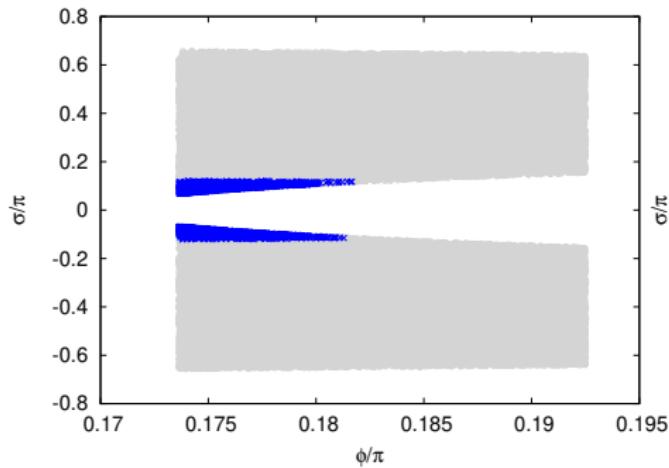
# Concluding Remarks

- TBM pattern is actually disfavored by experiments
- Corrections to TBM matrix can predict current mixing angles
- Small  $\mu - \tau$  symmetry breaking in the mass matrix could help bounding Majorana phases via a precise determination of the atmospheric angle
- Small breaking favors degenerate hierarchy and can be tested in future experiments

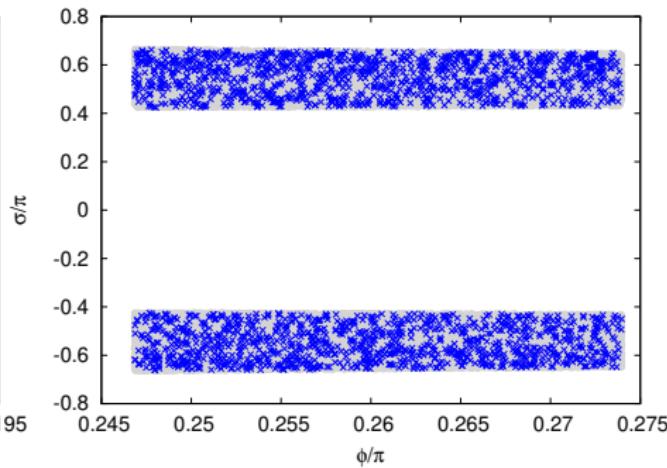
# Thank you!!

# Correction Parameters

Case I:



Case II:



- Better correlation is obtained for  $\alpha_2 = 0$ . Values of  $\alpha_1$  are the same of  $\beta_1$ .
- These regions are consistent with current  $3\sigma$  ranges of the mixing angles

## $\mu - \tau$ symmetric mass matrix

Mass matrix obtained from  $U_{\mu-\tau}$

$$M_v^{\mu-\tau} = \begin{pmatrix} m_1 c_{12}^2 + m_2 s_{12}^2 & -\frac{s_{12}}{\sqrt{8}}(m_1 - m_2) & \sigma \frac{s_{12}}{\sqrt{8}}(m_1 - m_2) \\ -\frac{s_{12}}{\sqrt{8}}(m_2 - m_1) & \frac{1}{2}(m_1 s_{12}^2 + m_2 c_{12}^2 + m_3) & \frac{-\sigma}{2}(m_1 s_{12}^2 + m_2 c_{12}^2 - m_3) \\ \sigma \frac{s_{12}}{\sqrt{8}}(m_1 - m_2) & \frac{-\sigma}{2}(m_1 s_{12}^2 + m_2 c_{12}^2 - m_3) & \frac{1}{2}(m_1 s_{12}^2 + m_2 c_{12}^2 + m_3) \end{pmatrix}$$