

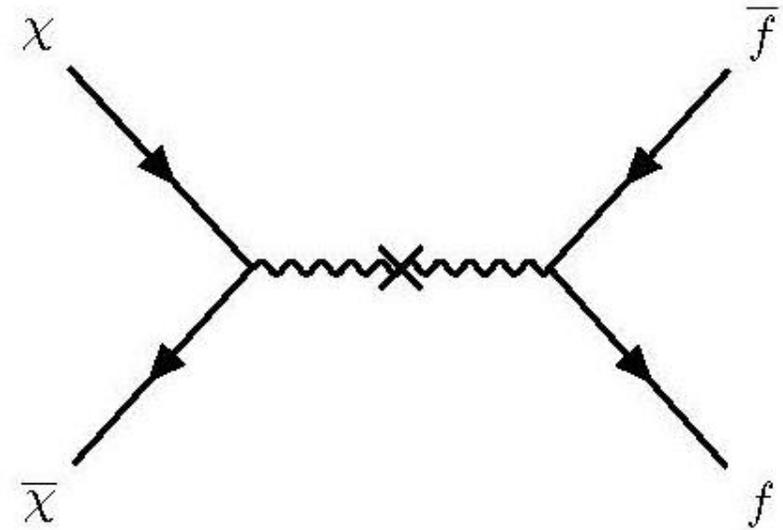
Kaluza Klein Portal Matter

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BASED ON ARXIV:2205.11545 BY G. N. WOJCIK

Vector Portal/Kinetic Mixing DM

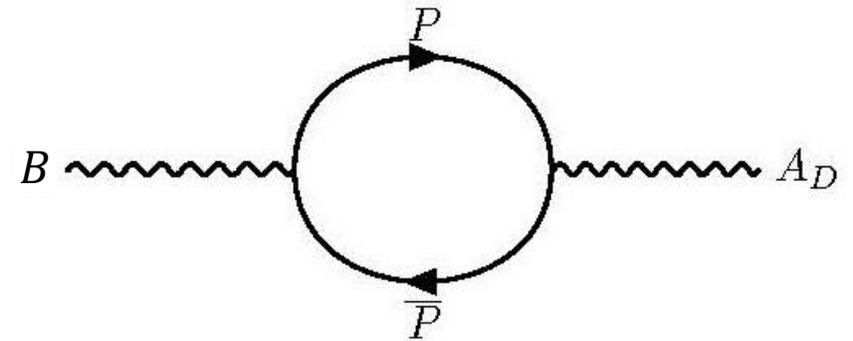
- WIMP miracle: $\langle\sigma v\rangle \sim g_D^4/m_{DM}^2 \sim 10^{-26} \text{ cm}^3 \text{ s}^{-1}$
- If $g_D \sim g_{EW}, m_{DM} \sim \text{TeV}$
- Expand this parameter space: DM is a particle χ that interacts only with a new “dark force” given by the gauge group $U(1)_D$. The portal is the dark $U(1)_D$ gauge boson A_D .
- The Standard Model (SM) is entirely uncharged under $U(1)_D$.
- So, SM now couples with strength $\sim \epsilon e Q$ to A_D .
- $m_{DM}, m_{A_D} \sim 0.1 - 1 \text{ GeV}, \epsilon \sim 10^{-(3-4)}$ reproduces the correct relic abundance without running afoul of other experimental constraints.



$$\sim \frac{\epsilon}{2 c_W} B_{\mu\nu} A_D^{\mu\nu}$$

Portal Matter: Origins of ϵ

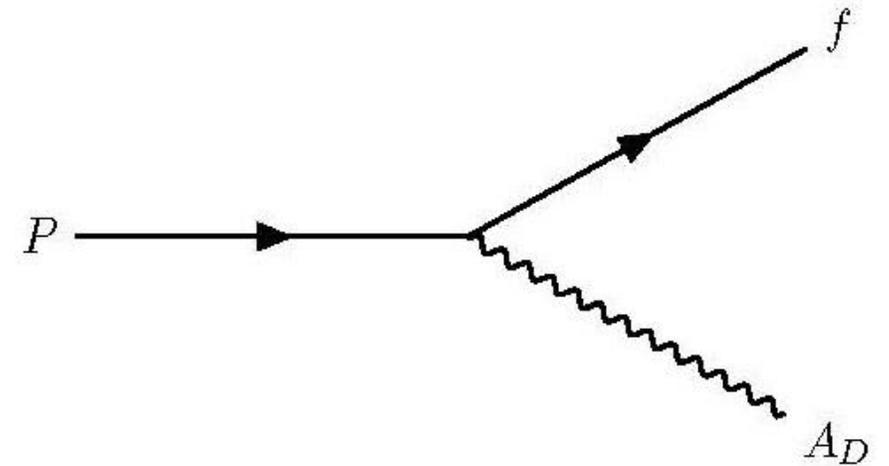
- The minimal setup: DM and a dark photon A_D . The small parameter ϵ is added by hand.
- By asking where ϵ comes from, can we get a window into higher-energy physics?
- A natural source for $\epsilon \sim 10^{-(3-4)}$ would be *portal matter*: Heavy particles charged under both $U(1)_D$ and SM hypercharge. Call the particle P , for portal matter.
- To get *finite* and *calculable* ϵ , we need $\sum Q_{Y_i} Q_{D_i} = 0$ to eliminate dependence on the renormalization scale μ .



$$\epsilon \propto \sum_i Q_{Y_i} Q_{D_i} \log \left(\frac{m_i^2}{\mu^2} \right)$$

Portal Matter Phenomenology 101

- What can we say about portal matter? To keep it simple, we'll assume it's fermionic.
- Precision electroweak constraints, $H \rightarrow gg$ branching ratio $\rightarrow P$ is **vector-like**.
- Portal matter has SM charge, so it has to decay quickly (cosmological measurements).
- If portal matter P is a copy of an SM particle f , it can mix with that particle and decay via $P \rightarrow f A_D$. Depending on parameter space, A_D decays to either dark matter or charged SM fermions.
- Simple portal matter can give atypical signatures (e.g., displaced lepton-jets).^{1,2} **Highly suppressed** decay via regular vector-like fermion decay channels, e.g. $P \rightarrow f Z$



¹T. G. RIZZO, PHYS. REV D **99**, NO.11, 115024 (2019) [ARXIV:1810.07531 [HEP-PH]]

²T. D. REUTER AND T. G. RIZZO, PHYS. REV. D **101**, NO.1, 015014 (2020) [ARXIV:1909.09160 [HEP-PH]]

Non-Minimal Portal Matter

- Minimal portal matter:
 - No structure for requiring ϵ finite
 - Only additional gauge group is $U(1)_D$
- Why does the SM have vector-like copies??
 - Extend the dark group to a larger gauge group $\mathcal{G}_D \supset U(1)_D$ and assume portal matter and an SM field are part of **the same multiplet** of \mathcal{G}_D .
- Group structure of \mathcal{G}_D might guarantee finite mixing
- But then... Why are the $U(1)_D$ -neutral ones chiral & light? Why are the others vector-like?
- How do we ensure light chiral states for one set of charge states, and heavy vector-like states for others?
- How do we break $\mathcal{G}_D \rightarrow U(1)_D$?

$$\begin{aligned}\mathcal{G}_{SM} \times SU(2)_I \times U(1)_I &\rightarrow \mathcal{G}_{SM} \times U(1)_D \\ (\mathbf{R}, \mathbf{2})_{-\frac{1}{2}} &\rightarrow \mathbf{R}_0 + \mathbf{R}_{-1} \\ (\bar{\mathbf{R}}, \mathbf{1})_{+1} &\rightarrow \bar{\mathbf{R}}_{+1}\end{aligned}$$

R: Some representation of \mathcal{G}_{SM}

Black: SM fermion

Blue: Portal matter fermion

Portal Matter in an Extra Dimensions

- In theories of extra dimensions, we already see light fermion modes + heavy vector-like particles.
- Imagine a 5 D theory compactified on an $S_1/(Z_2 \times Z'_2)$ orbifold
- Fields can have Neumann (+) or Dirichlet (-) boundary conditions at boundaries $\phi = 0, \pi$. In 5D to have a light mode, you need Neumann condition at both boundaries.
- On the right: Chiral SM fermion plus vector-like copies!

$$\mathcal{G}_{SM} \times SU(2)_D \rightarrow \mathcal{G}_{SM} \times U(1)_D$$

$$(\mathbf{F}, \mathbf{3}) \rightarrow \mathbf{F}^0 + \mathbf{F}^{+1} + \mathbf{F}^{-1}$$

Fermion	$\phi = 0$	$\phi = \pi$
F_L^0	+	+
F_R^0	-	-
$F_L^{\pm 1}$	+	-
$F_R^{\pm 1}$	-	+

Portal Matter in an Extra Dimension II

WHAT HAVE WE GAINED

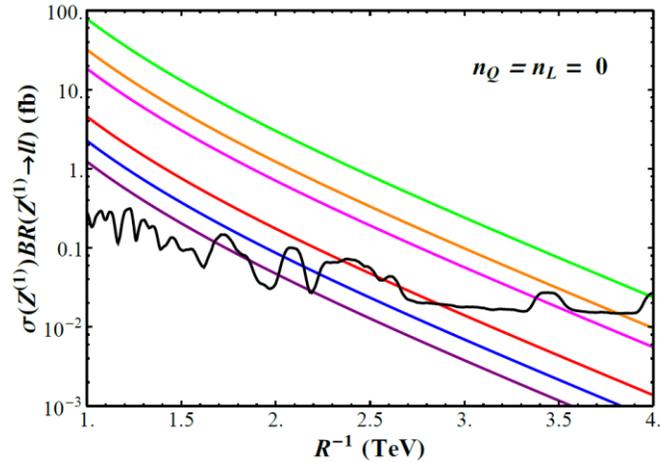
- Consistent boundary conditions also break $SU(2)_D \rightarrow U(1)_D$ at $\phi = \pi$ – there's our high-scale symmetry breaking!
- $U(1)_D$ broken on $\phi = \pi$ brane by scalar: Simpler Higgs sector!
- No-muss, anomaly-free, *simple* non-minimal portal matter– vector-like and chiral fields come out of the same multiplet
- Condition for finite calculable KM is **still** $\sum Q_{Y_i} Q_{D_i} = 0$ (must apply to all sets of fields w/ the same $Z_2 \times Z_2$ parity individually), automatically satisfied.
- Scale motivation?

WHAT HAVE WE ADDED

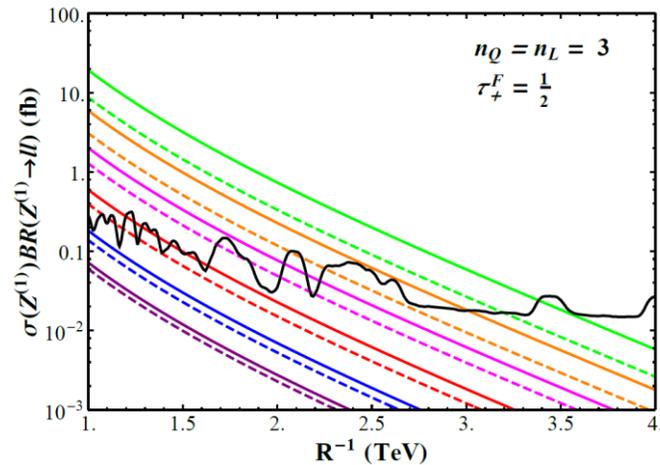
- Kaluza Klein states for SM fields– can be produced directly or affect precision measurements
- Three new gauge bosons: A_D, I^+, I^- . Only A_D has a zero mode.
- **Lighter** Kaluza Klein states for portal matter, I^\pm .
- Need brane-localized kinetic terms for appropriate magnitude KM– we assume BLKT's for all fermions and $SU(2)_D$ gauge bosons. Call these various param's τ 's and ω 's

Phenomenology: SM-Only Constraints

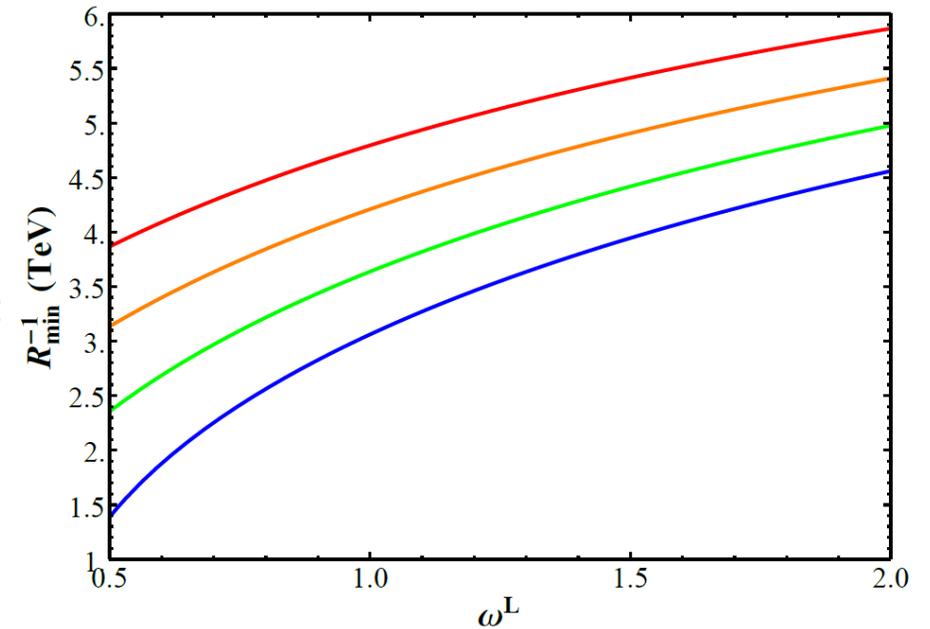
Direct: (No PM)



Direct: (PM)

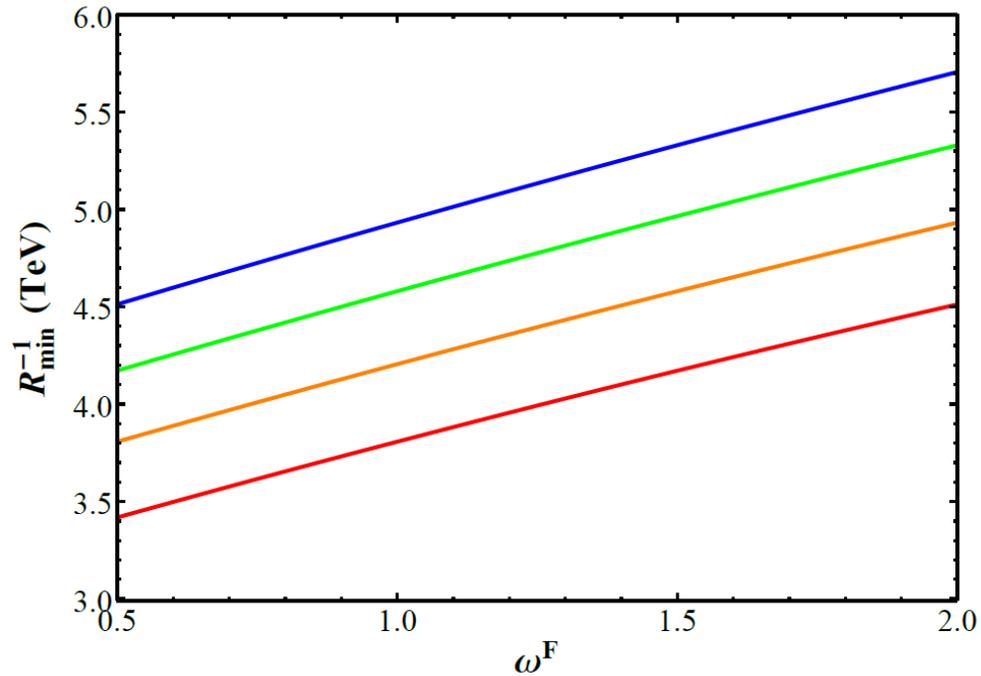


Precision (ρ_0):

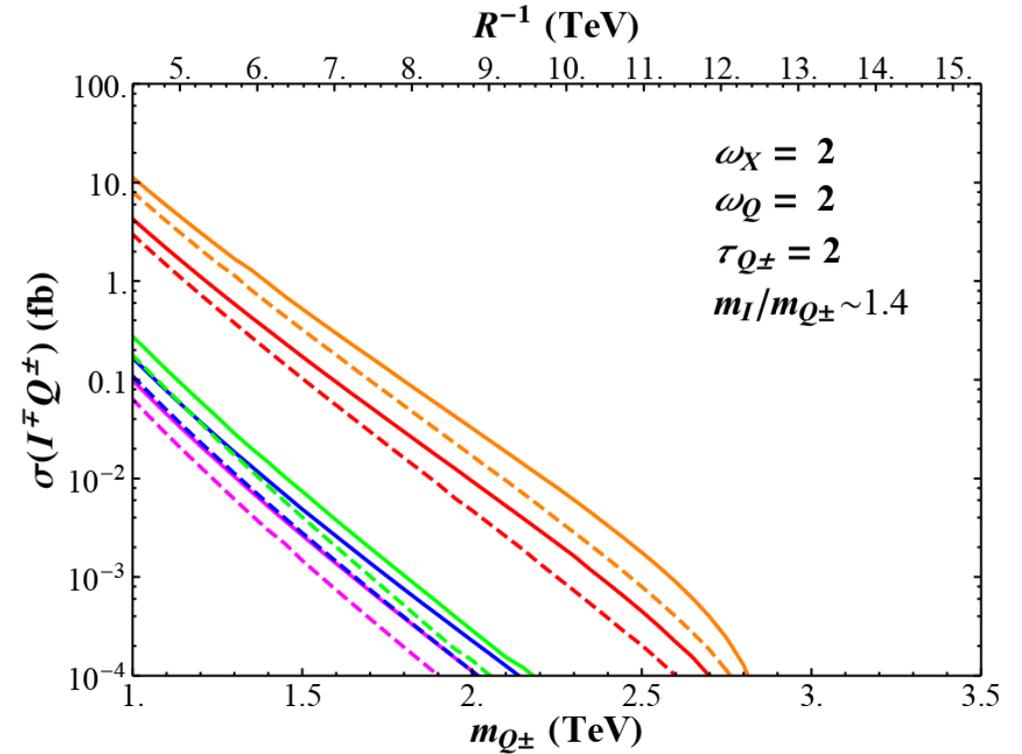


Phenomenology: Portal Matter

HADRONIC PM PAIR PRODUCTION



CROSS SECTION FOR PORTAL MATTER+I



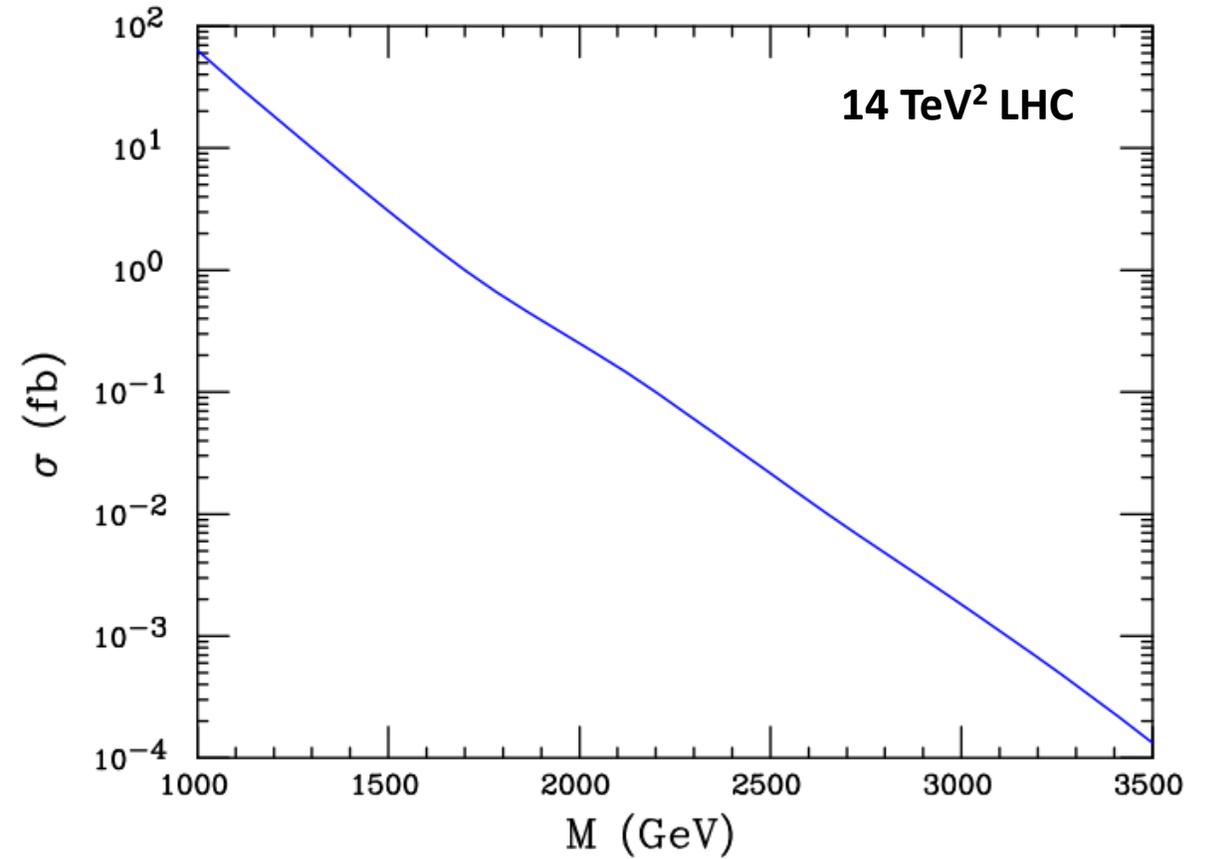
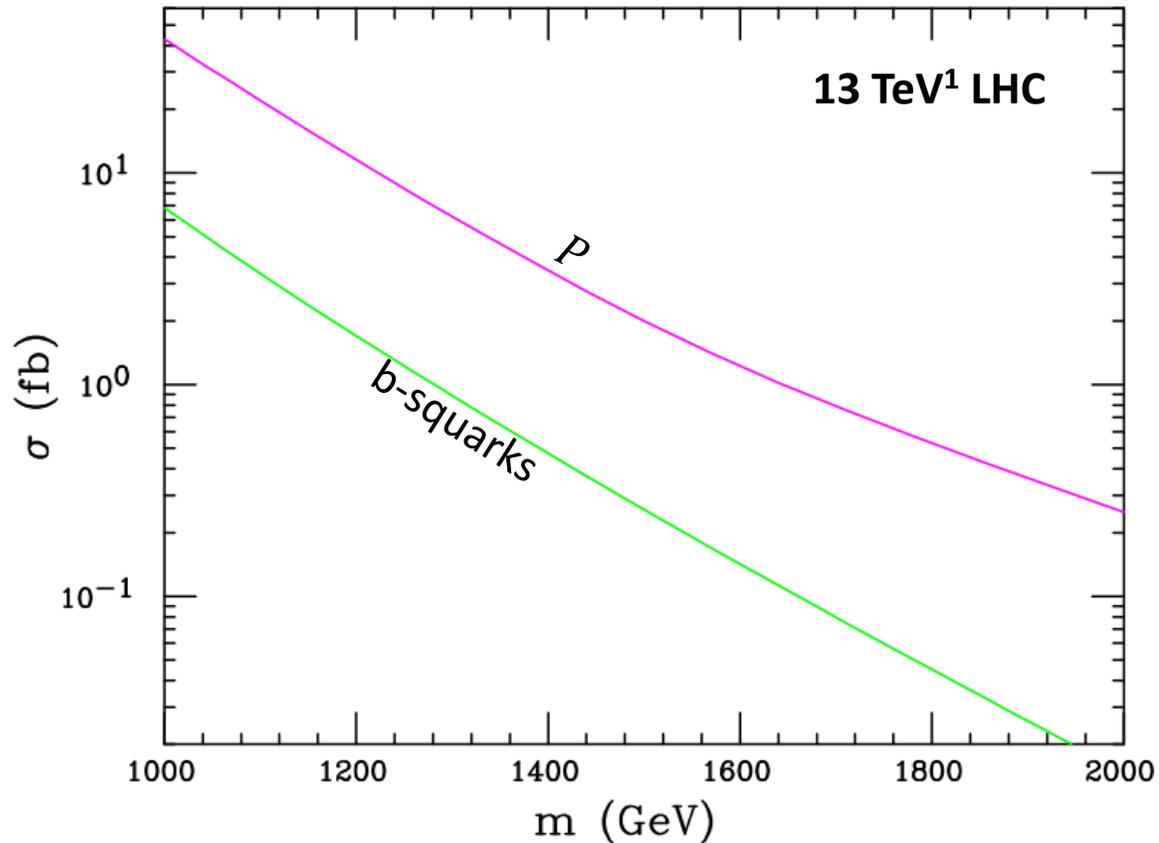
Conclusions

- The scenario where vector-like copies of SM fermions act as portal matter is straightforwardly realized in the context of an extra dimension
- Leads to simpler Higgs sectors and fewer ad hoc particle introductions than equivalent in 4D
- Even in 5D, the naïve 4D condition for finite and calculable KM still works!
- Kaluza Klein portal matter will likely be the **lightest** new KK modes appearing in the theory and their existence weakens other direct KK mode searches— could be the first direct evidence we'd find of a TeV-scale extra dimension!

Thank You!

Backup Slides

QCD Pair Production



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