

Inflation from Dynamical Projective Connections

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Muhammad Abdullah¹, Calvin Bavor², Biruk Chafamo¹, Xiaole Xiang¹,
Muhammad Hamza Kalim¹, Kory Stiffler^{3,4}, Catherine A. Whiting^{1,2}

¹Bates College, ²Colorado Mesa University, ³Brown University, ⁴University of Iowa

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Cosmological Inflation

Rapid expansion of early universe

Explains problems with Big Bang Model: **Flatness, Horizon, Monopole, etc**

Models of inflation involve scalar field, ϕ :

$$S = \frac{1}{2\kappa_0} \int d^4x \sqrt{|g|} f(\phi) R - \int d^4x \sqrt{|g|} \left[\frac{1}{2} \nabla_a \phi \nabla^a \phi - V(\phi) \right] \quad (1)$$

- $V(\phi)$, kinetic term, coupling to gravity choices guided by theoretical and observational constraints
- often no fundamental principle governing its dynamics

Our approach: Can we explain inflation as particular manifestation of tensor field arising from TW-gravity?

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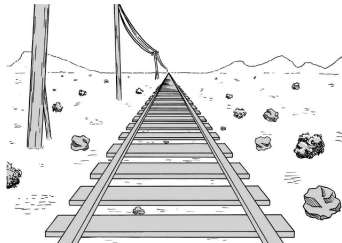
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Projective Geometry



Projective Invariance \Rightarrow General Relativity

$$\mathcal{K}_{bcd}^a = R_{bcd}^a + \delta_{[c}^a \mathcal{P}_{d]b} - \delta_b^a \mathcal{P}_{[cd]}$$

TW - Gravity Action

$$S_{TW} = -\frac{1}{2\tilde{\kappa}_0} \int dl d^d x \sqrt{|G|} \mathcal{K} - \tilde{J}_0 c \int dl d^d x \sqrt{|G|} \left[\mathcal{K}^2 - 4\mathcal{K}_{\alpha\beta} \mathcal{K}^{\alpha\beta} + \mathcal{K}^\alpha_{\beta\mu\nu} \mathcal{K}^{\beta\mu\nu} \right] \quad (2)$$

Integrate $\ell = \frac{\lambda}{\lambda_0}$



$$S_{TW} = -\frac{1}{2\tilde{\kappa}_0} \int d^d x \sqrt{|g|} (R + (d-1)\mathcal{P} + 2\Lambda_0) + J_0 c \int d^d x \sqrt{|g|} \left[\lambda_0^2 K_{abc} K^{abc} - \mathcal{P}_{ab} \tilde{\mathcal{P}}_*^{ab} \right] - J_0 c S_{GB} \quad (3)$$

$$K_{abc} = \nabla_{[b} \mathcal{P}_{c]a}, \quad \tilde{\mathcal{P}}_*^{ab} \equiv (d-1)g^{ab}\tilde{\mathcal{P}} - 2(2d-3)\tilde{\mathcal{P}}^{ab} \quad (4)$$

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Coupling Constants

$d = 4$, $c = \hbar = 1$ define dimensionless couplings:

$$\kappa_0 = n_\kappa M_p^{-2}, \quad \lambda_0 = n_\lambda M_p^{-1}, \quad J_0 = n_J \hbar = n_J$$

- n_κ = gravitational coupling
- n_λ = projective length scale
- n_J = angular momentum parameter

By definition: $n_\kappa, n_\lambda, n_J > 0$

$$\hat{n} = \frac{n_\lambda}{n_J n_\kappa} \tag{5}$$

Realizing Inflation

$$\mathcal{P}_{ab} = \frac{M_p}{n_\lambda} \phi g_{ab} \quad (6)$$

$$S = -\frac{M_p^2}{2n_\kappa} \int d^4x \sqrt{|g|} f(\phi) R + 12n_J \int d^4x \sqrt{|g|} \left[\frac{1}{2} \nabla_a \phi \nabla^a \phi - V(\phi) \right] \quad (7)$$

Non-Minimally coupled Inflation with linear coupling (Brans-Dicke-like):

$$f(\phi) = 1 + \frac{8}{\hat{n} M_p} \phi \quad (8)$$

Predicted Inflaton Potential

$$V(\phi) = \frac{M_p^4}{32n_J^2 n_\kappa^2} (f(\phi)^2 - 1) = \frac{M_p^2}{2n_\lambda^2} (4\phi^2 + M_p \hat{n} \phi) \quad (9)$$

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Canonical Scalar Field Action

Perform Conformal transformation from Jordan to **Einstein** frame:

$$g_{ab} = e^{-2\omega} \tilde{g}_{ab} \quad \omega = \ln \sqrt{\frac{f(\phi)}{n_\kappa}} \quad (10)$$

$$h(\phi) = \begin{cases} \sqrt{6}M_p \left[\sqrt{1 + \frac{\hat{n}n_\lambda}{8} f(\phi)} - \tanh^{-1} \sqrt{1 + \frac{\hat{n}n_\lambda}{8} f(\phi)} \right] - \frac{8}{\hat{n}n_\lambda} < f(\phi) < 0 \\ \sqrt{6}M_p \left[\sqrt{1 + \frac{\hat{n}n_\lambda}{8} f(\phi)} - \coth^{-1} \sqrt{1 + \frac{\hat{n}n_\lambda}{8} f(\phi)} \right] f(\phi) > 0 \end{cases} \quad (11)$$

$$S = \int d^4x \sqrt{|\tilde{g}|} \left(-\frac{M_p^2}{2} \tilde{R} + \frac{1}{2} \tilde{\nabla}_a h \tilde{\nabla}^a h - \tilde{V}(\phi(h)) \right) \quad (12)$$

$$\tilde{V}(\phi) = \frac{3M_p^4}{8n_J} \left(1 - \frac{1}{f(\phi)^2} \right) = \frac{6M_p^4 \phi (M_p \hat{n} + 4\phi)}{n_J (M_p \hat{n} + 8\phi)^2} \quad (13)$$

Slow Roll Parameters

Accelerated expansion: $\ddot{a} > 0 \Rightarrow \epsilon < 1, \eta \ll 1$

$$\epsilon = \frac{M_p^2}{2} \left(\frac{\partial \tilde{V} / \partial h}{\tilde{V}} \right)^2, \quad \eta = M_p^2 \frac{\partial^2 \tilde{V} / \partial h^2}{\tilde{V}} \quad (14)$$

$$C(\phi) \equiv \frac{12}{\hat{n}^2 f(\phi)^2} [8 + \hat{n} n_\lambda f(\phi)] = \frac{96}{\hat{n}^2 f(\phi)^2} \left[1 + \frac{\hat{n} n_\lambda}{8} + \frac{n_\lambda}{M_p} \phi \right] \quad (15)$$

$$\epsilon = \frac{M_p^2}{2C} \left(\frac{\partial \tilde{V} / \partial \phi}{\tilde{V}} \right)^2, \quad \eta = \frac{M_p^2}{\sqrt{C} \tilde{V}} \frac{\partial}{\partial \phi} \left[C^{-1/2} \frac{\partial \tilde{V}}{\partial \phi} \right] \quad (16)$$

Model Predictions:

$$\epsilon = \frac{32}{3(f^2 - 1)^2 (8 + \hat{n} n_\lambda f)}, \quad \eta = -\frac{16}{3} \frac{32 + 5\hat{n} n_\lambda f}{(f^2 - 1)(8 + \hat{n} n_\lambda f)^2} \quad (17)$$

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Observational Parameters

Model Predictions

Constraints

Scalar Spectral Index:

$$n_s = 1 - 6\epsilon + 2\eta|_{\phi_*} \quad n_s = 0.9649 \pm 0.0042 \quad (18)$$

Tensor to Scalar Ratio:

$$r = 16\epsilon|_{\phi_*} \quad r < 0.056 \quad (19)$$

Amplitude of Density Perturbations:

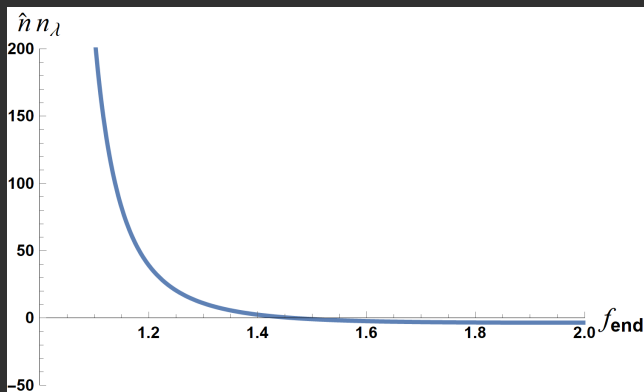
$$A_s = \frac{1}{24\pi^2 M_p^4} \frac{\tilde{V}}{\epsilon} |_{\phi_*} \quad 2 \leq \ln 10^{10} A_s \leq 4 \quad (20)$$

Number of e-folds:

$$N \simeq \frac{1}{M_p^2} \int_{h_{\text{end}}}^{h_*} \frac{\tilde{V}}{\partial \tilde{V} / \partial h} dh \quad N \in [50 - 70] \quad (21)$$

Constraining f_{end} with $\epsilon|_{\phi_{end}} = 1$

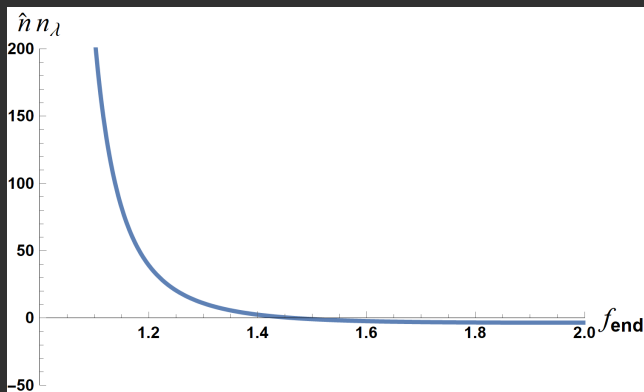
$$\hat{n}n_\lambda = \frac{32}{3f_{end}(f_{end}^2 - 1)^2} - \frac{8}{f_{end}}, \quad \hat{n}n_\lambda > 0 \quad (22)$$



$$1 < f_{end} < \sqrt{1 + 2/\sqrt{3}} \approx 1.47$$

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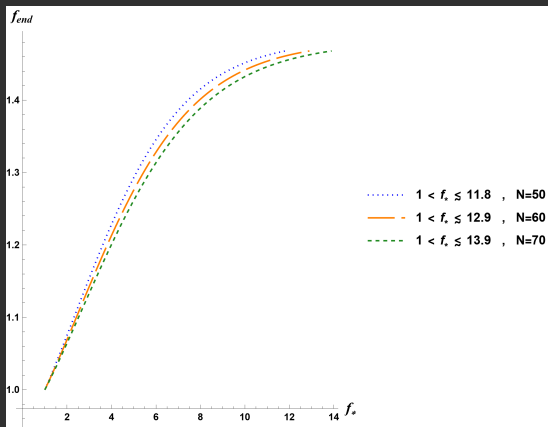
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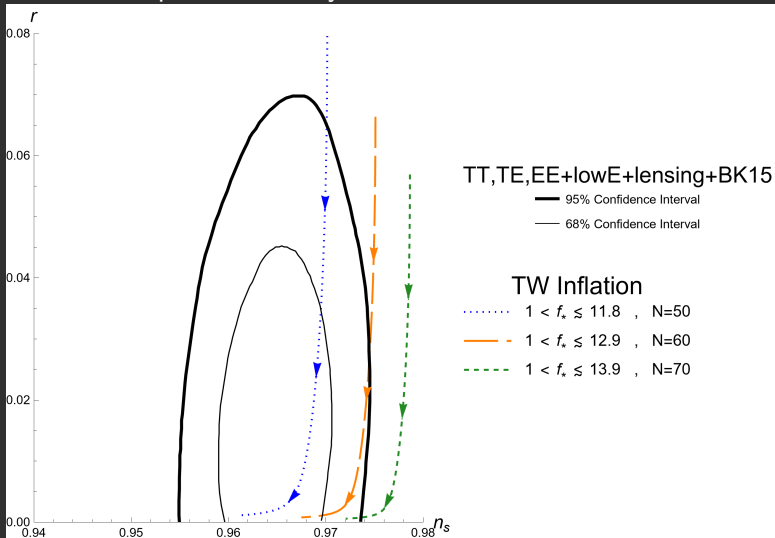
Constraining f_* to f_{end} with N e-folds

$$N \simeq \frac{1}{32} (\hat{n}n_\lambda(\alpha^3 - 1)f_{end}^3 + 12(\alpha^2 - 1)f_{end}^2 - 3\hat{n}n_\lambda(\alpha - 1)f_{end} - 24 \ln \alpha)$$
$$\alpha \equiv \frac{f(\phi_*)}{f(\phi_{end})} \quad , \quad f_{end} = f(\phi_{end}) \quad (23)$$



Constraining f_* with r vs. n_s Parameter space

Planck, Bicep2, Keck Array data

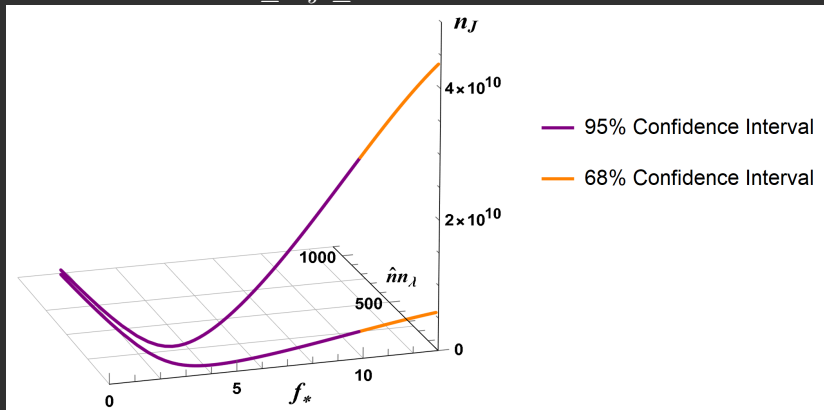


Values of f_* , f_{end} , $\hat{n}n_\lambda$ constrained by observational values of r vs. n_s

		95% CL Boundary			68% CL Boundary		
		f_*	f_{end}	$\hat{n}n_\lambda$	f_*	f_{end}	$\hat{n}n_\lambda$
N	50	1.13	1.01	30300	1.62	1.05	1150
	60	1.70	1.05	1060	9.84	1.44	0.896
	70	11.1	1.45	0.625			

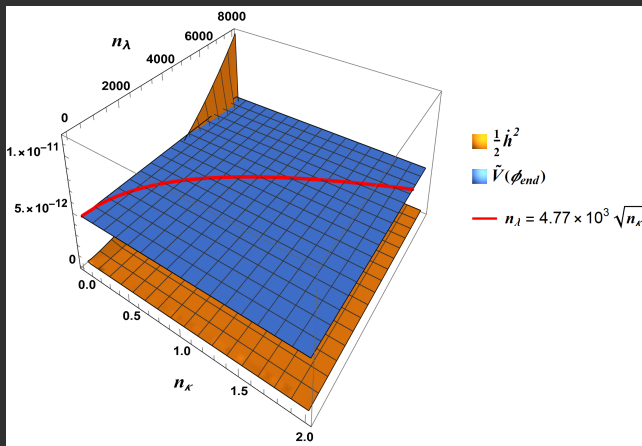
Constraining n_J with A_s

$$N = 60: 1.15 \times 10^8 \leq n_J \leq 4.37 \times 10^{10}$$



Validating Slow Roll

Slow Roll: $\frac{1}{2}\dot{h}^2 \ll \tilde{V}$ for $N = 60$, $n_J = 4.37 \times 10^{10}$, $f_* = 12.9$,
 $f_{end} = 1.47$



Parameter Space

Minimum-Maximum values of three free parameters

$$\begin{aligned} \text{Choosing: } & 0 < n_\kappa < 2 \\ & 1.15 \times 10^8 < n_J < 4.37 \times 10^{10} \\ & 0 < n_\lambda < 1.34 \times 10^6 \end{aligned} \tag{24}$$

Conclusion

- Thomas-Whitehead Gravity predicts a model of inflation when $P_{ab} \sim \phi g_{ab}$ consistent with data for range of parameters $(n_\kappa, n_\lambda, n_J)$

Future Work

$$P_{ab} \sim \phi g_{ab} + w_0 W_{ab}, \quad \text{with} \quad W_{ab} g^{ab} = 0 \quad (25)$$

Scalar field Ansatz for traceless W_{ab} :

$$W_a{}^b = u(t) \begin{bmatrix} 3 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -1 \end{bmatrix} \quad (26)$$

- Can $u(t)$ be curvaton, responsible for reheating mechanism?
- Traceless tensor field W_{ab} : responsible for density perturbations of CMB without quantum fluctuations?
- We assumed connection is Levi-Civita: Investigate Inflation application in Palatini formalism of TW-gravity

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