



Testing NSI at colliders



NTN Workshop 05.29.2019
Dorival Gonçalves



Motivation

$$\mathcal{L}_{\text{NSI}} = -2\sqrt{2} G_F \varepsilon_{\alpha\beta}^{fP} (\bar{\nu}_\alpha \gamma_\rho \nu_\beta) (\bar{f} \gamma^\rho P f)$$

Collider Physics

Solar neutrino exp.

Atmospheric
neutrino exp.

NSI

Reactor neutrino exp.

Beam neutrino exp.

see Ivan's talk for global constraints on NSI

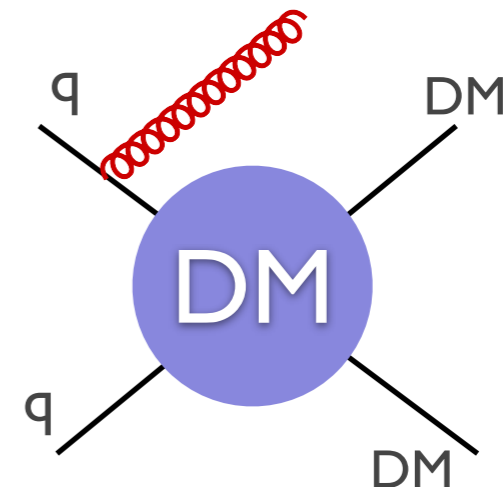
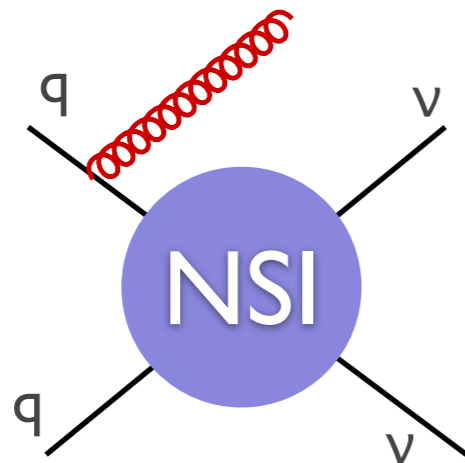
Motivation

$$\mathcal{L}_{\text{NSI}} = -2\sqrt{2} G_F \varepsilon_{\alpha\beta}^{fP} (\bar{\nu}_\alpha \gamma_\rho \nu_\beta) (\bar{f} \gamma^\rho P f)$$

In general, NSI contains both flavor-changing and flavor-diagonal interactions.

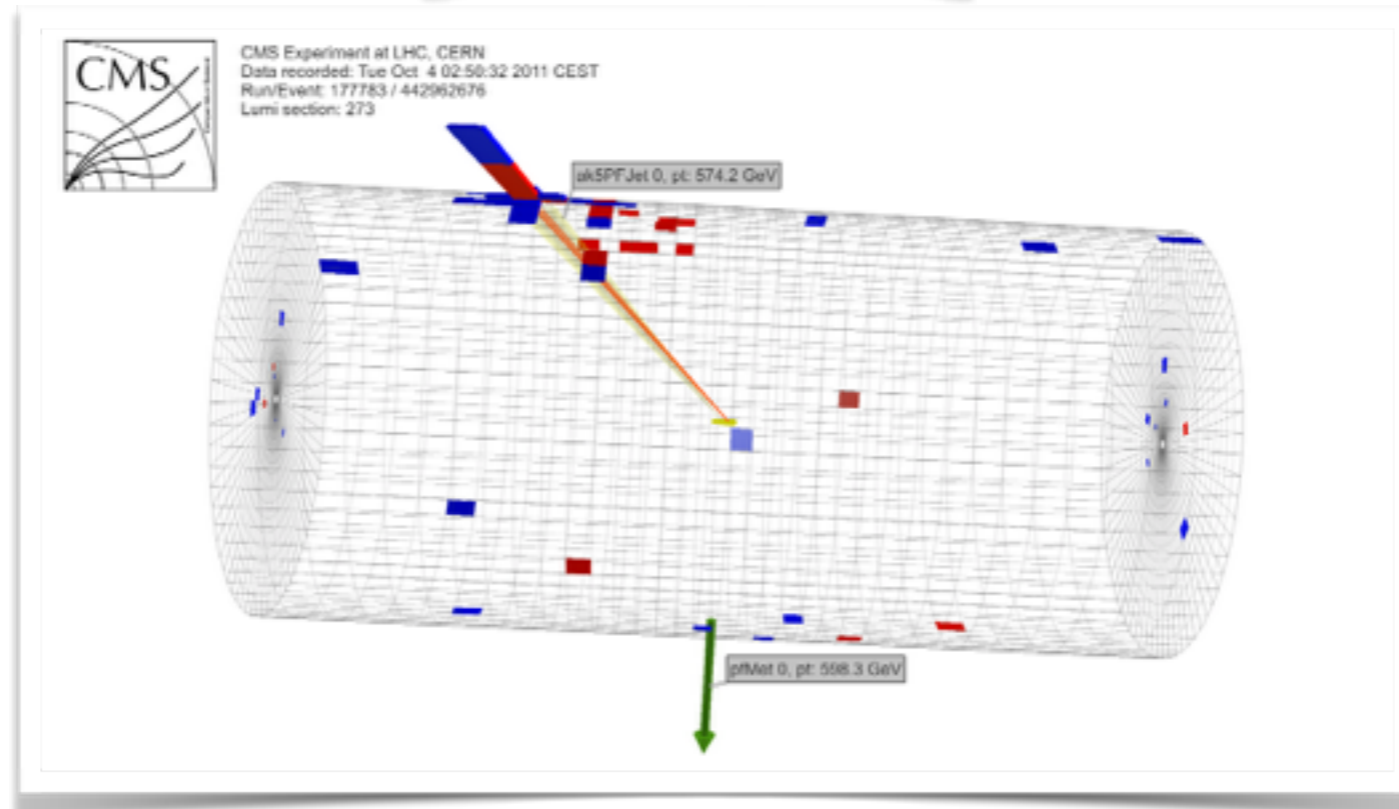
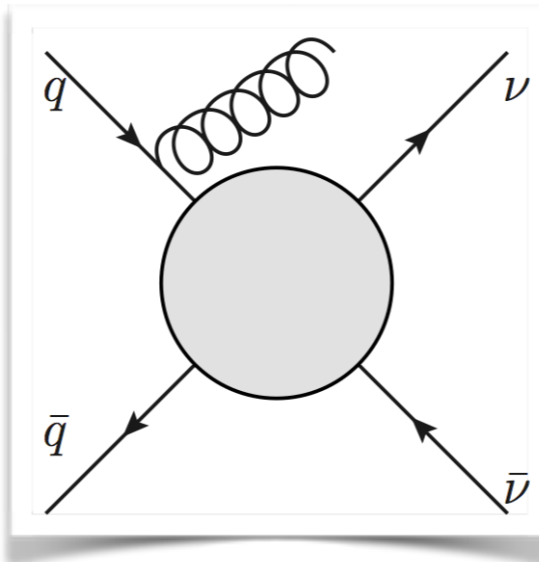
Flavor-changing - No SM analog

Flavor-diagonal - Interfere with SM background (can display non-trivial differences to DM scenario)



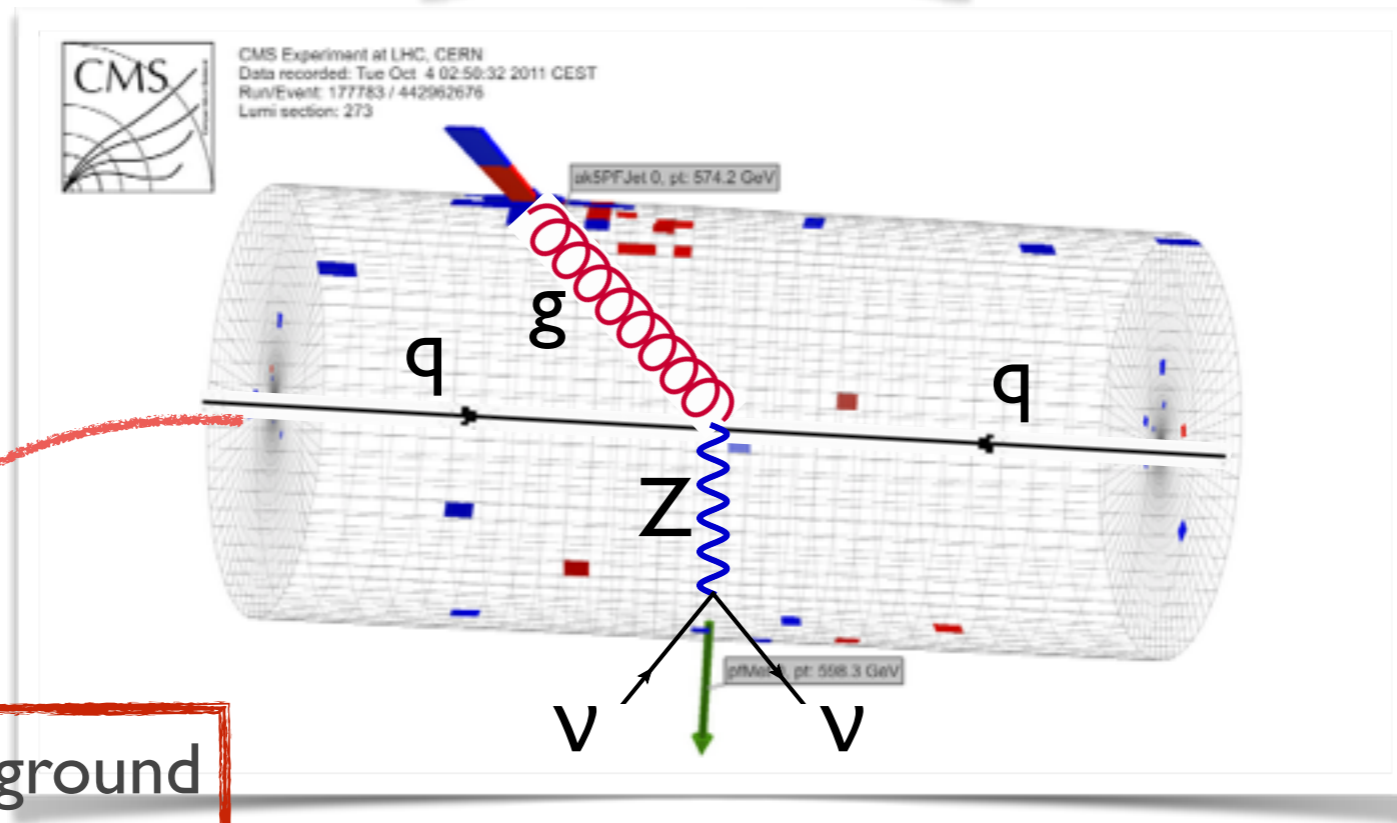
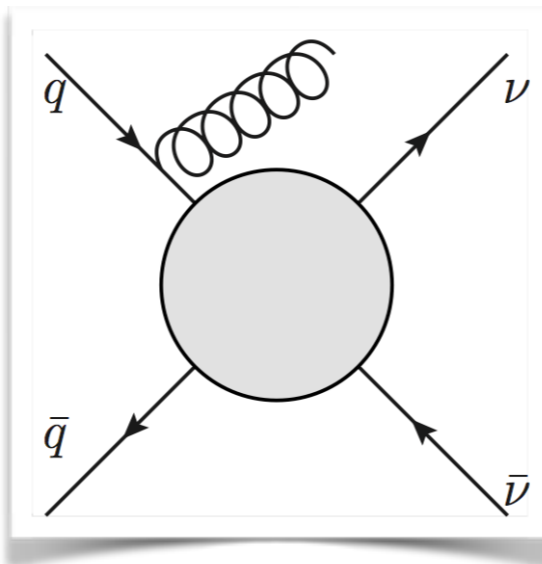
NSI at colliders

$$\mathcal{L}_{\text{NSI}} = -2\sqrt{2} G_F \epsilon_{\alpha\beta}^{fP} (\bar{\nu}_\alpha \gamma_\rho \nu_\beta) (\bar{f} \gamma^\rho P f)$$



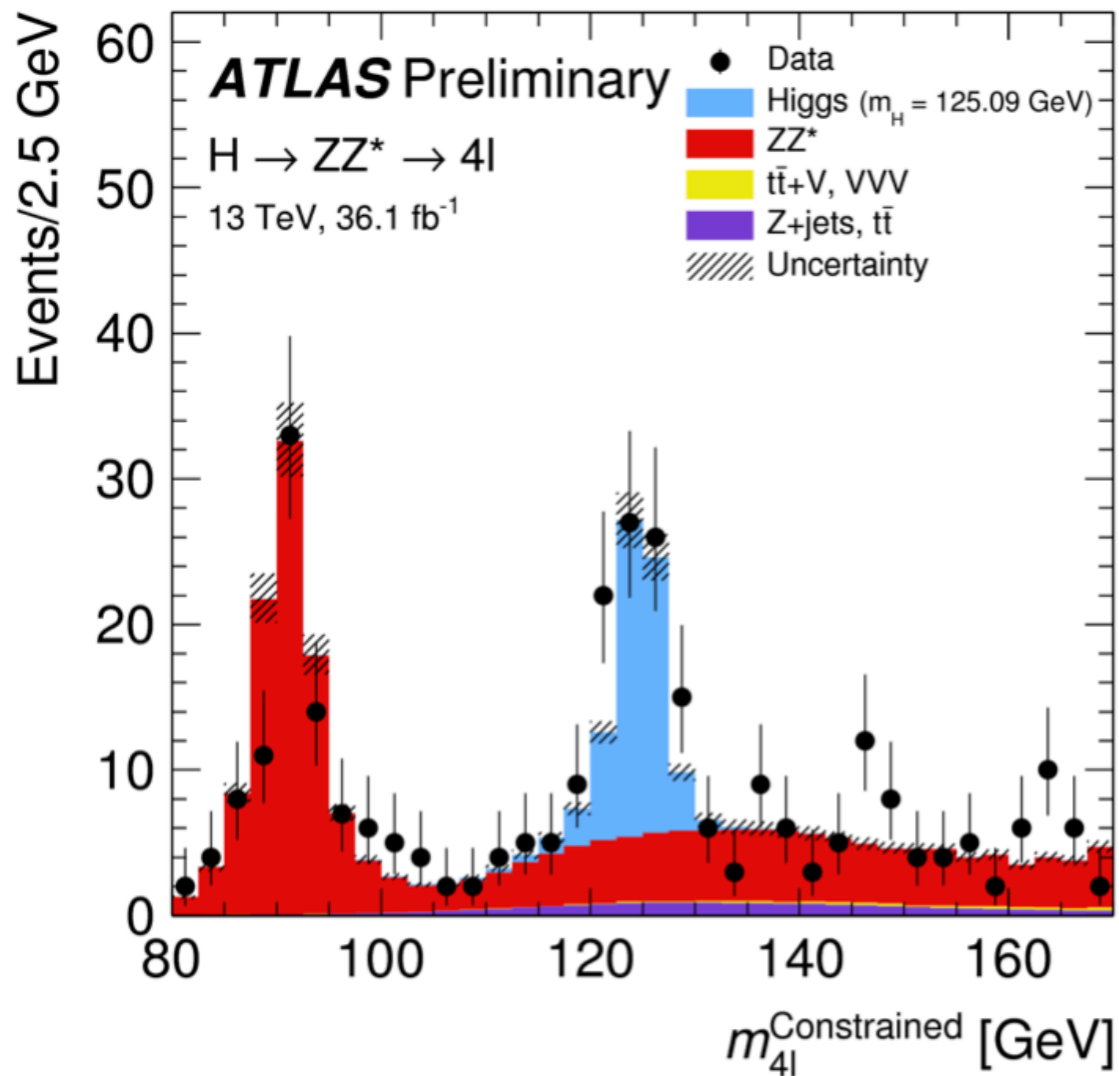
NSI at colliders

$$\mathcal{L}_{\text{NSI}} = -2\sqrt{2} G_F \epsilon_{\alpha\beta}^{fP} (\bar{\nu}_\alpha \gamma_\rho \nu_\beta) (\bar{f} \gamma^\rho P f)$$

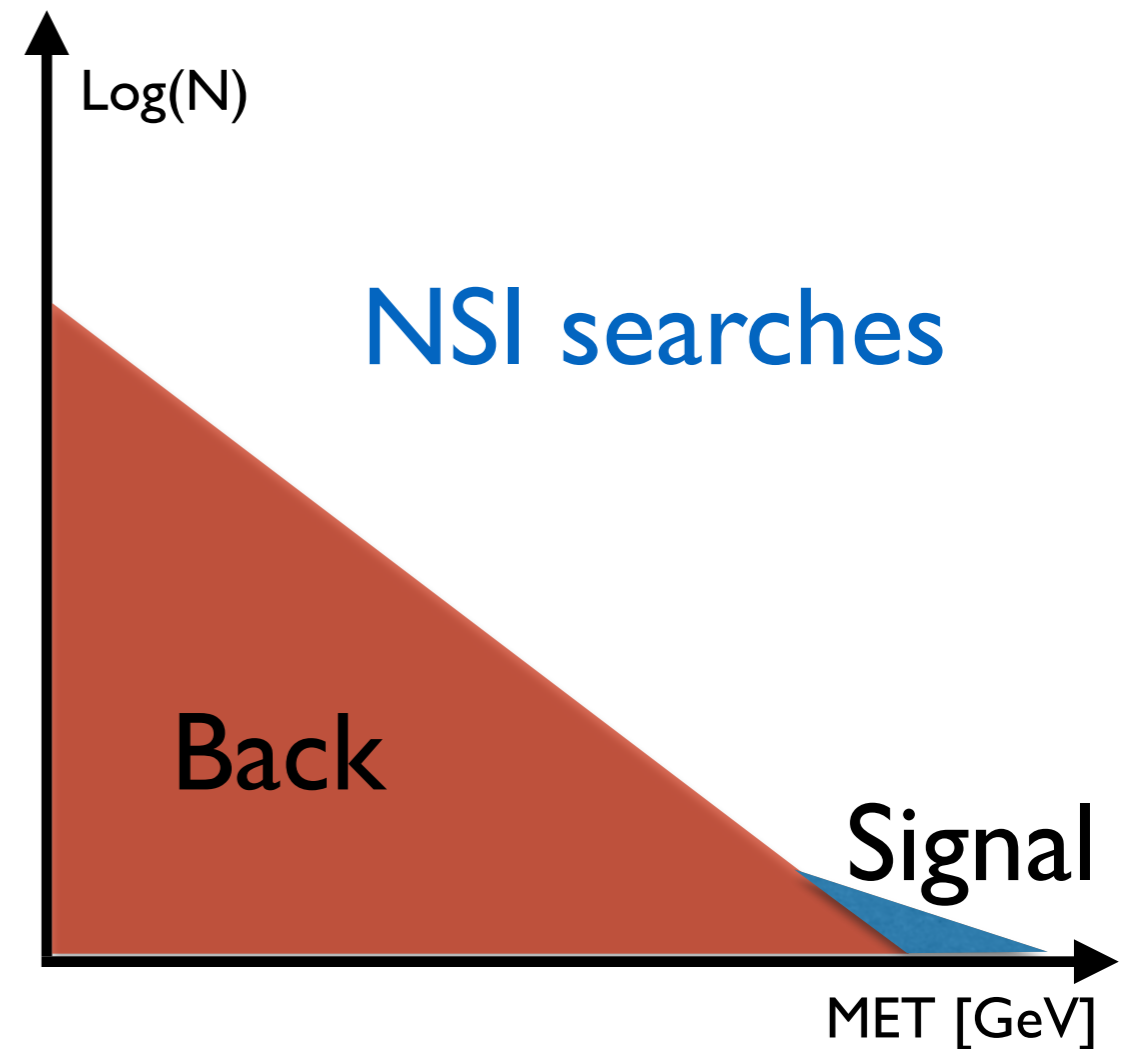


Irreducible SM background

NSI at colliders

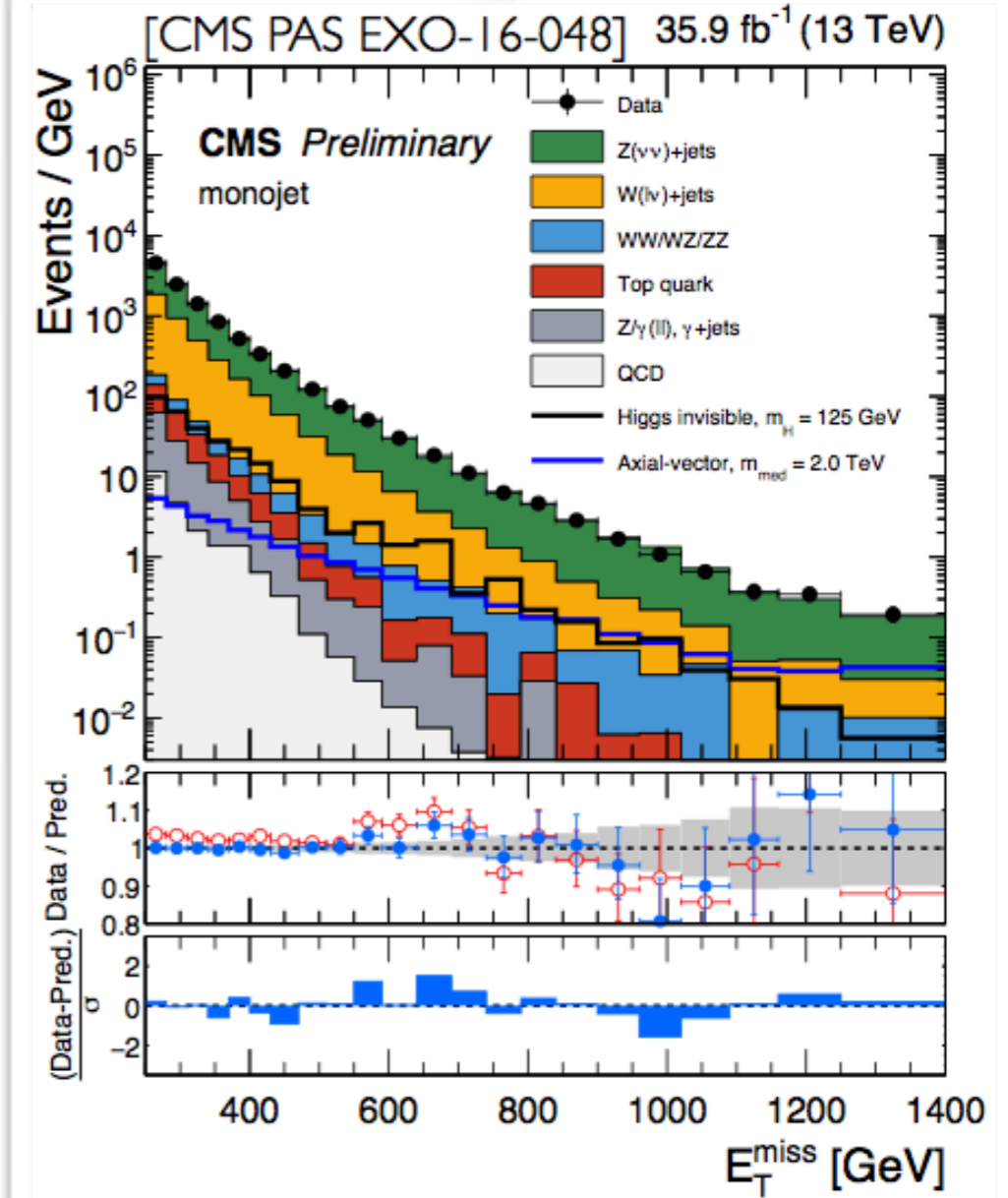
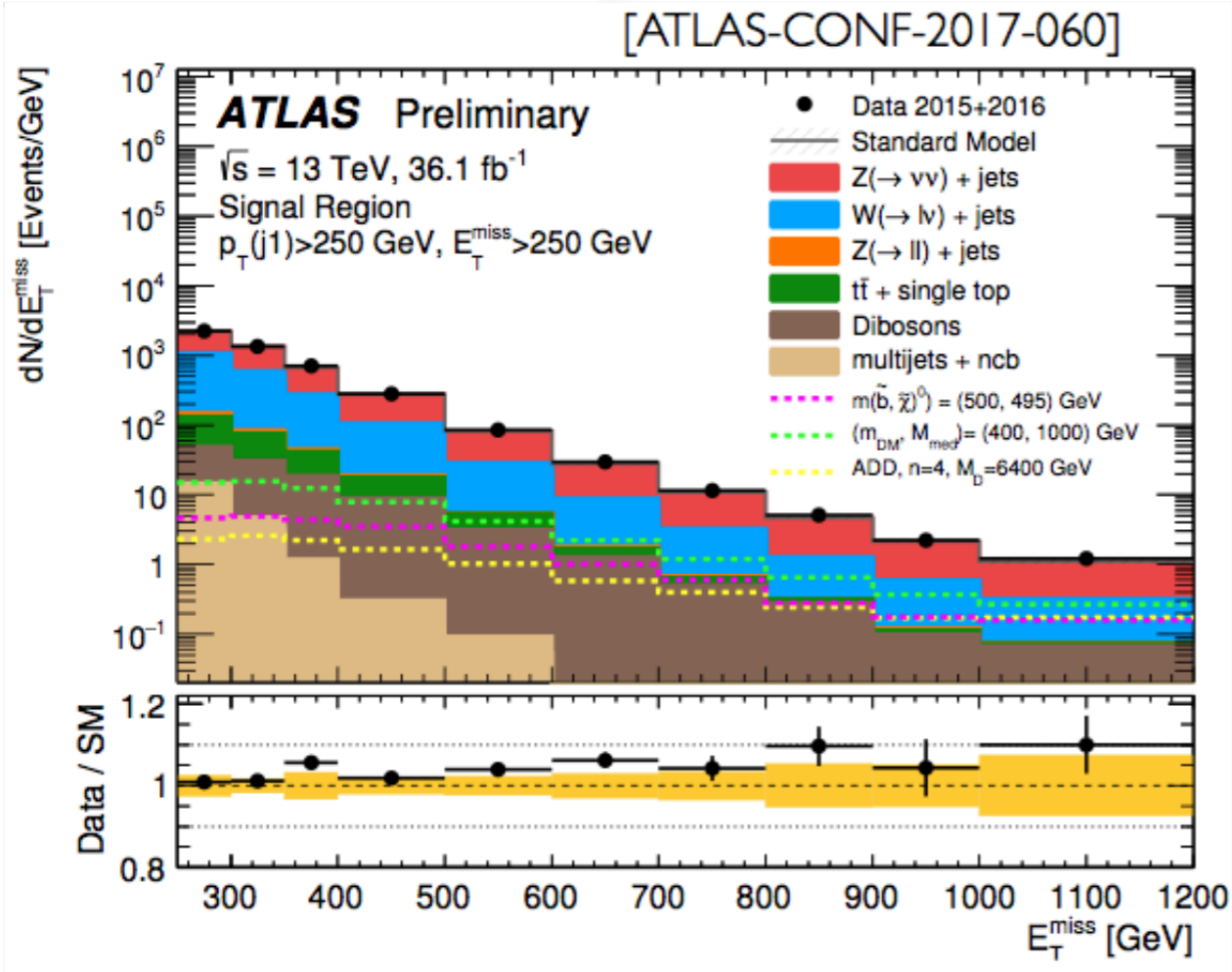


vs.



- Overwhelming SM background
- Signal: small enhancement in the tail of MET distribution
- ➔ Big challenge: requires precise estimation of background

Background composition



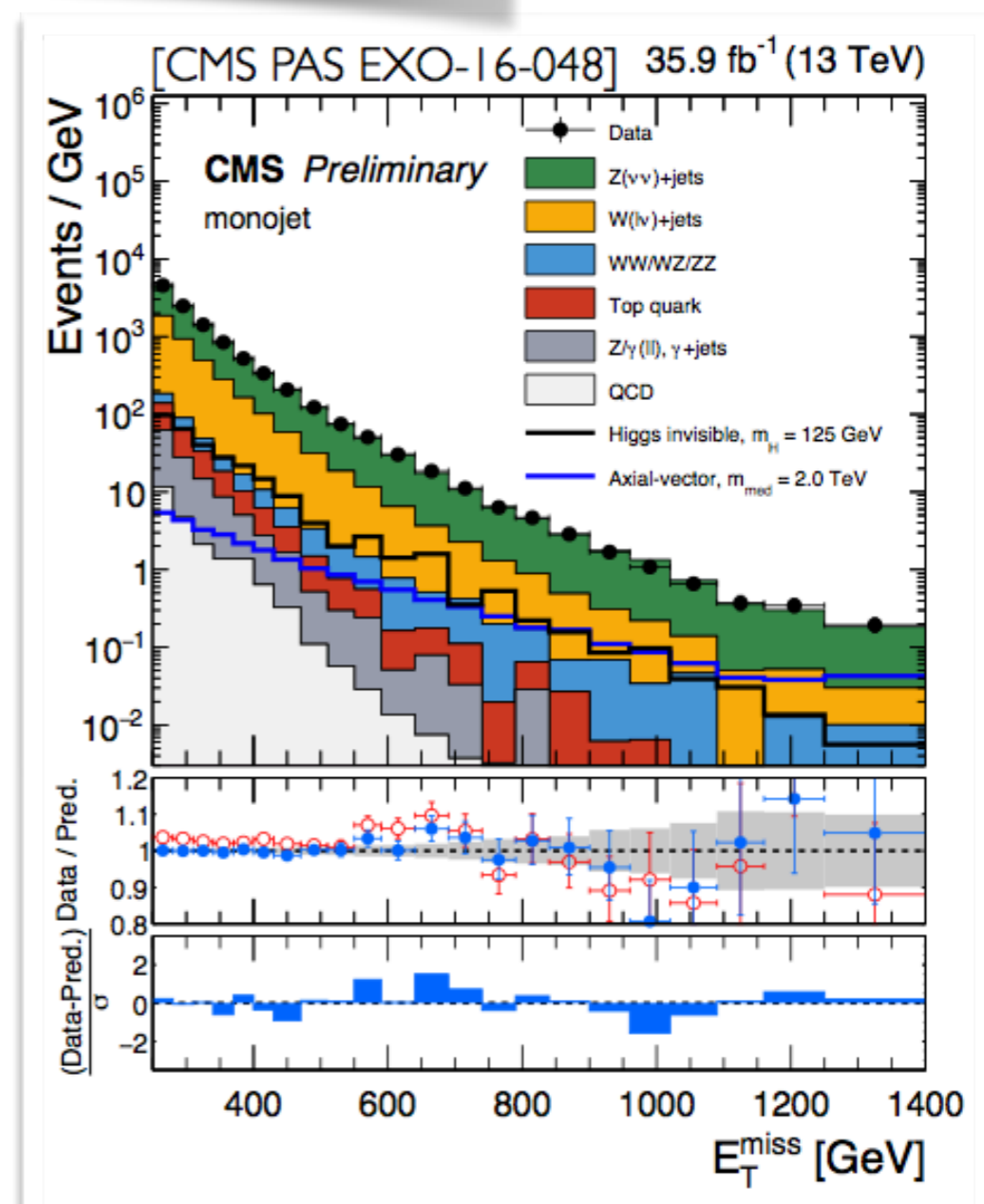
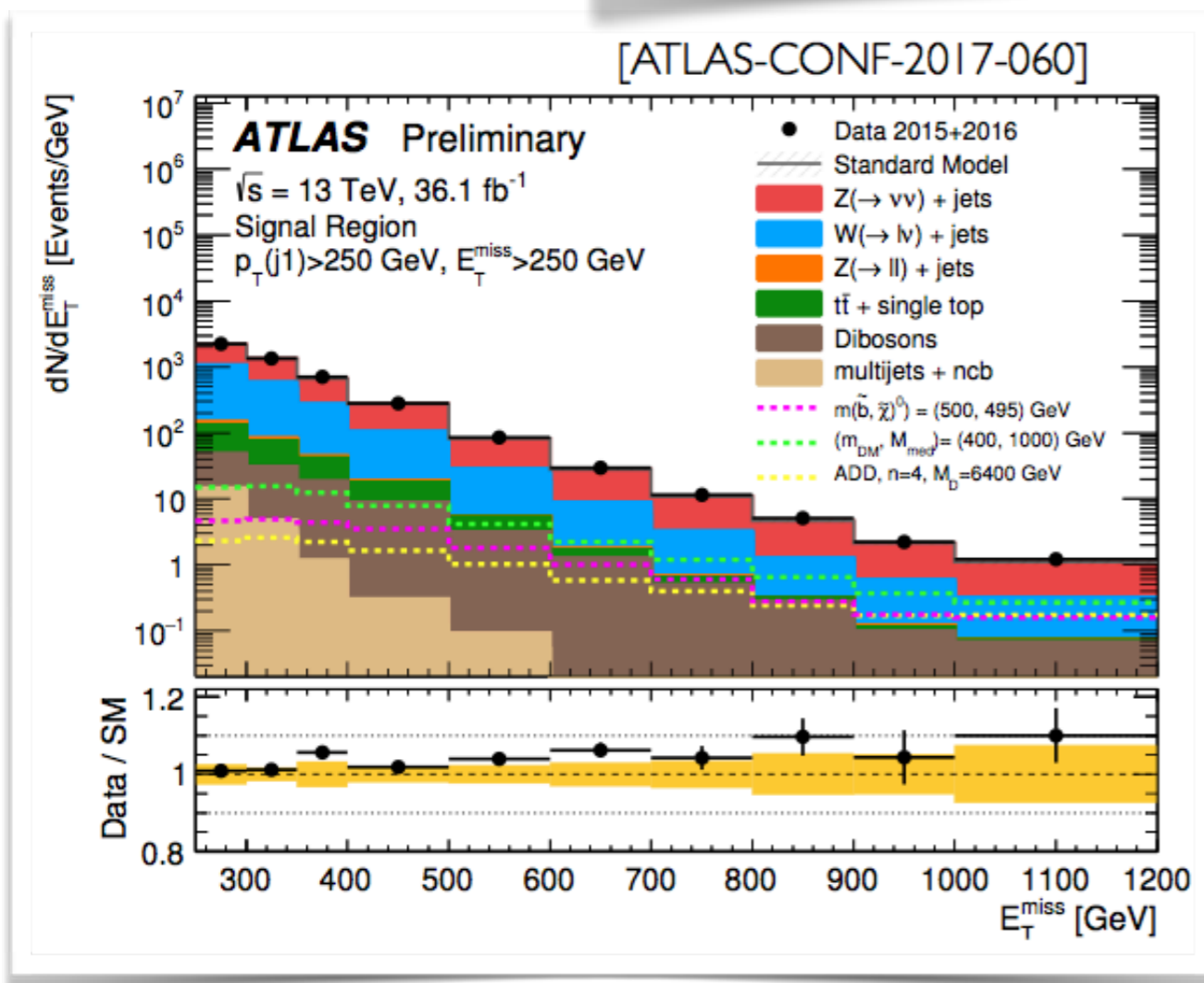
Dominant backgrounds:

$$pp \rightarrow Z(\rightarrow \nu\bar{\nu}) + \text{jets} \Rightarrow \text{MET} + \text{jets}$$

$$pp \rightarrow W(\rightarrow lv) + \text{jets} \Rightarrow \text{MET} + \text{jets} \text{ (lost lepton or hadronic taus)}$$

➔ Major limitation for NSI bounds: large and uncertain backgrounds. Background syst. $\sim 5\%$

Background composition



● Dominant backgrounds:

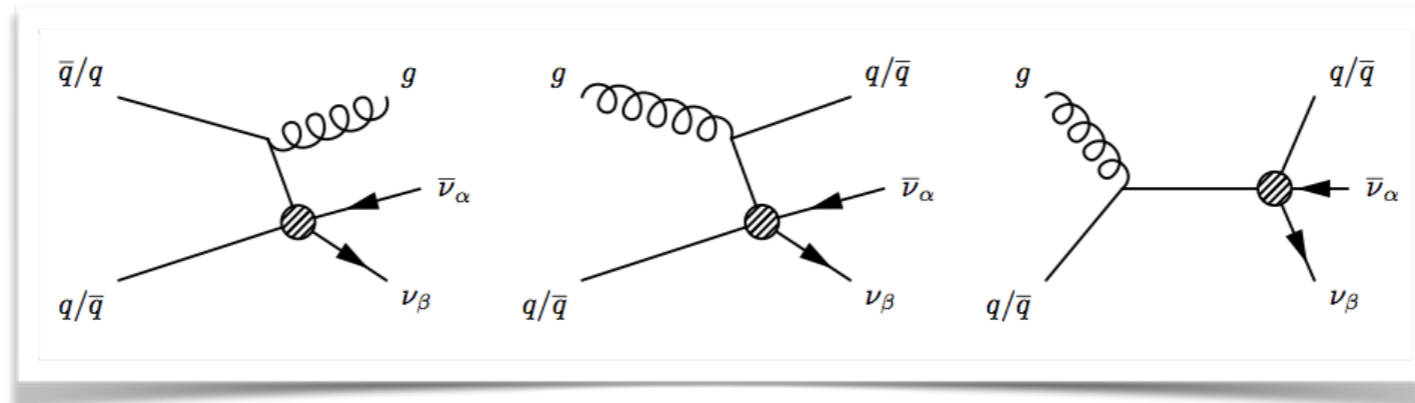
$pp \rightarrow Z(\rightarrow \nu\bar{\nu}) + \text{jets} \Rightarrow \text{MET} + \text{jets}$

$pp \rightarrow W(\rightarrow l\nu) + \text{jets} \Rightarrow \text{MET} + \text{jets}$ (lost lepton or hadronic taus)

➡ New theoretical and experimental efforts are significantly suppressing this syst. constraint

Contact NSI

● Monojet signature: $\mathcal{L}_{NSI} = -2\sqrt{2} G_F \epsilon_{\alpha\beta}^{qX} (\bar{q}\gamma_\mu P_X q) (\bar{\nu}_\alpha \gamma^\mu P_L \nu_\beta) + H.c.$



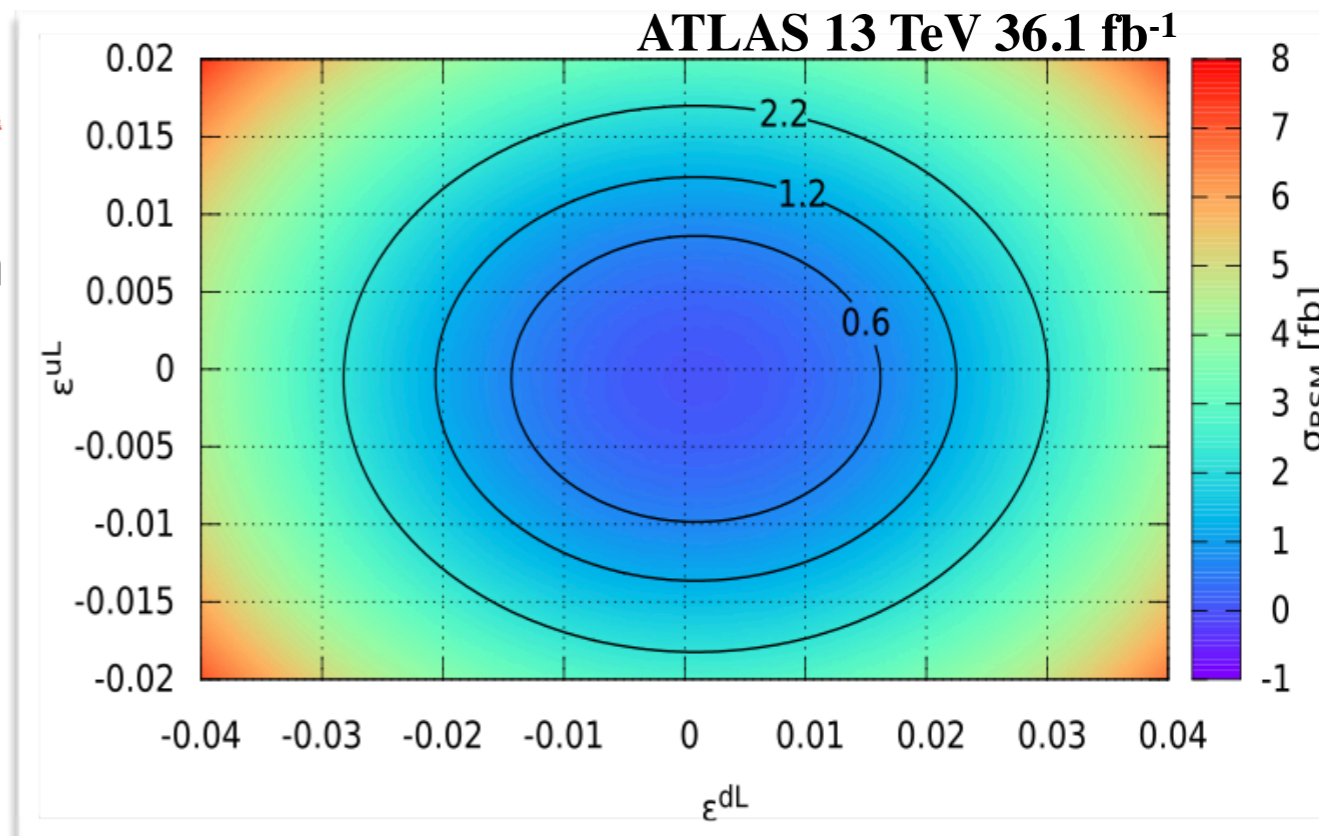
$$\sigma(pp \rightarrow j\bar{\nu}_\alpha\nu_\beta) = \sigma_{SM} + \epsilon\sigma_{int} + \epsilon^2\sigma_{NSI}$$

→ Different choices of (α, β) lead to same observables: indistinguishable from each other

● Limiting single pair of op. at a time

→ $a_u < a_d$ larger densities for u-quark in the proton

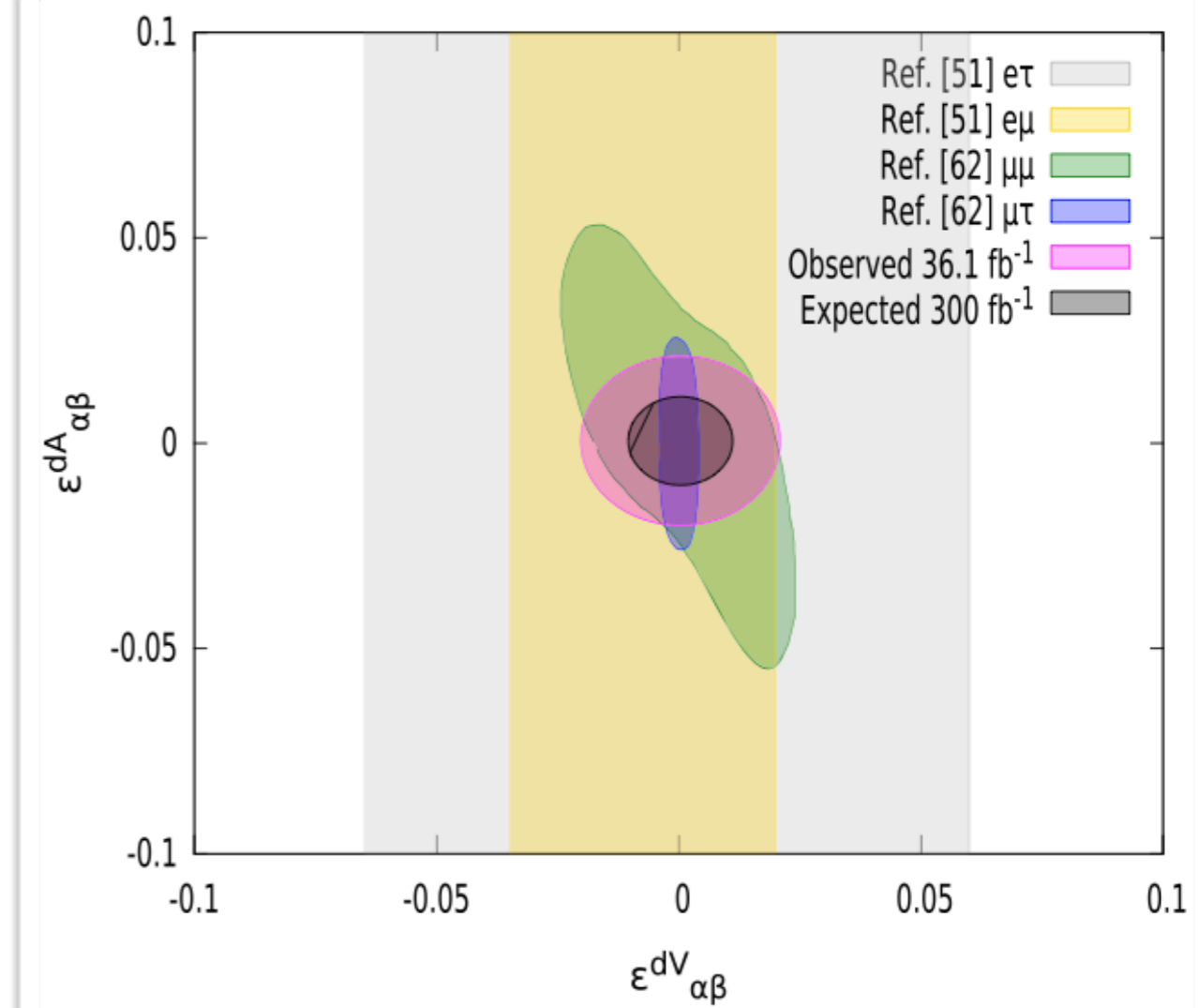
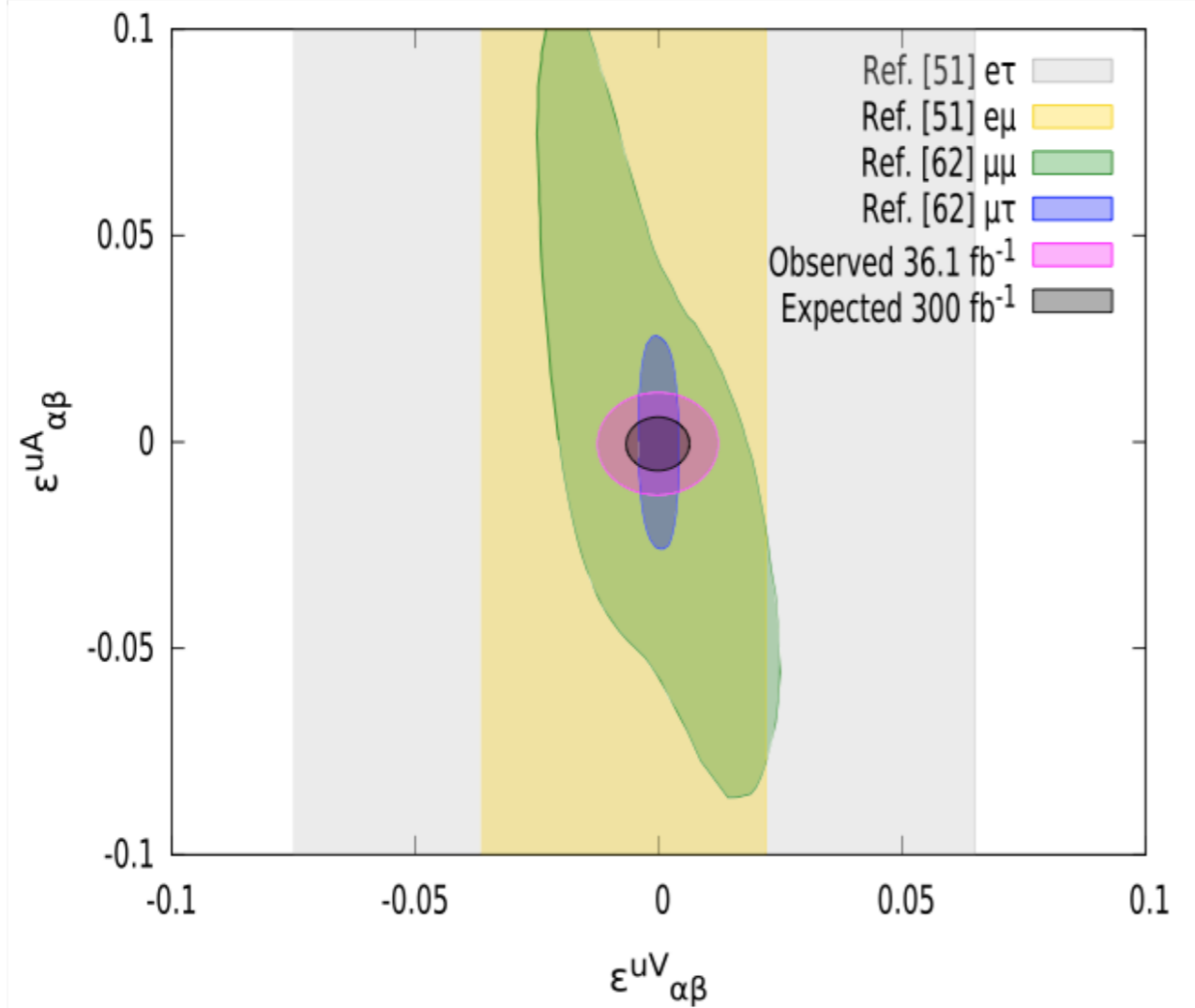
→ $\sigma_{int}/\sigma_{BSM}$ can be as significant as 50%



Choudhury, Ghosh, Niyogi '18; see also Friedland, Graesser, Shoemaker, Vecchi '12; Davidson, Sanz '11

Contact NSI

$$\mathcal{L}_{NSI} = -2\sqrt{2} G_F \epsilon_{\alpha\beta}^{qX} (\bar{q}\gamma_\mu P_X q) (\bar{\nu}_\alpha \gamma^\mu P_L \nu_\beta) + H.c.$$



Choudhury, Ghosh, Niyogi '18

● Ref. [51] neutrino oscillation bounds Coloma, Gonzalez-Garcia, Maltoni, Schwetz '17;

● Ref. [62] global fit atmospheric with accelerator neutrino data Escribuela, Miranda, Tórtola, Valle '11

➔ LHC u-quark currents results are significantly stronger than neutrino exp. d-quark currents are very competitive



Caveat: LHC bounds hold for the contact interaction limit. Light mediator can completely change the picture

Effective Field Theory Description

→ Model independent approach to NSI searches

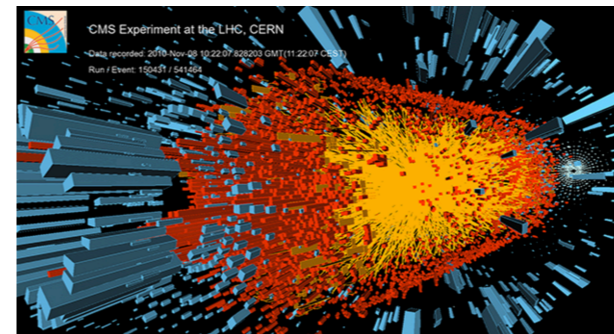
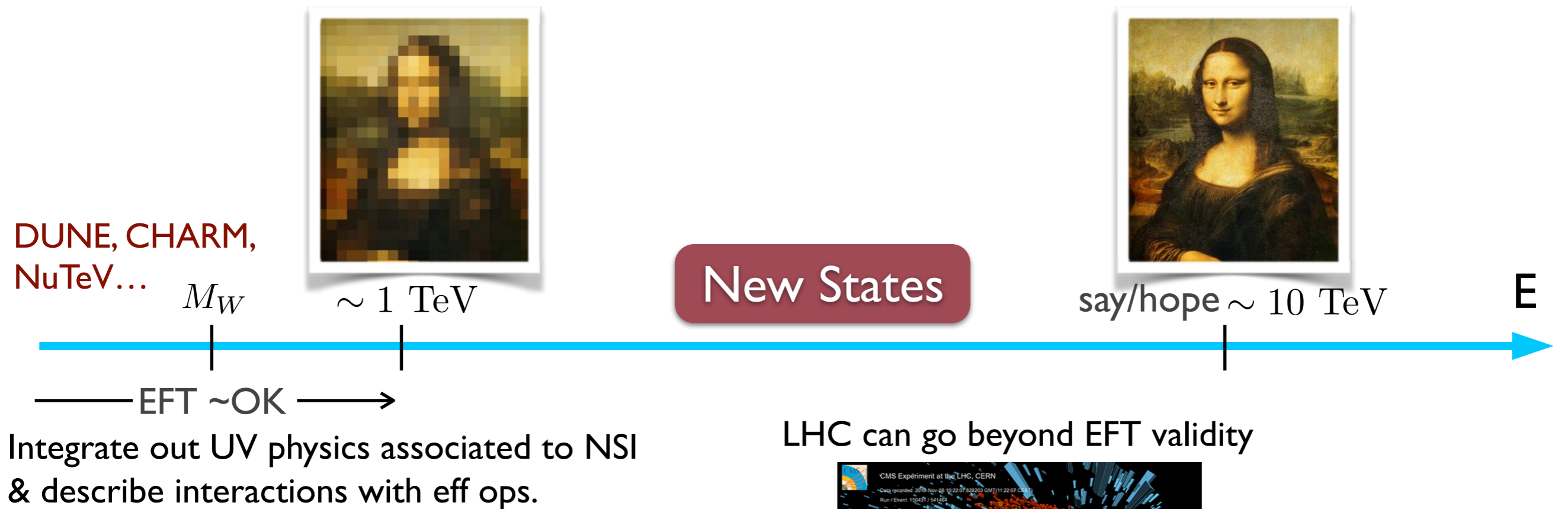
→ NSI bounds with few parameters

→ Valid in the presence of an energy gap: $M_* \gg E$

 
New physics effective scale Relevant process energy scale

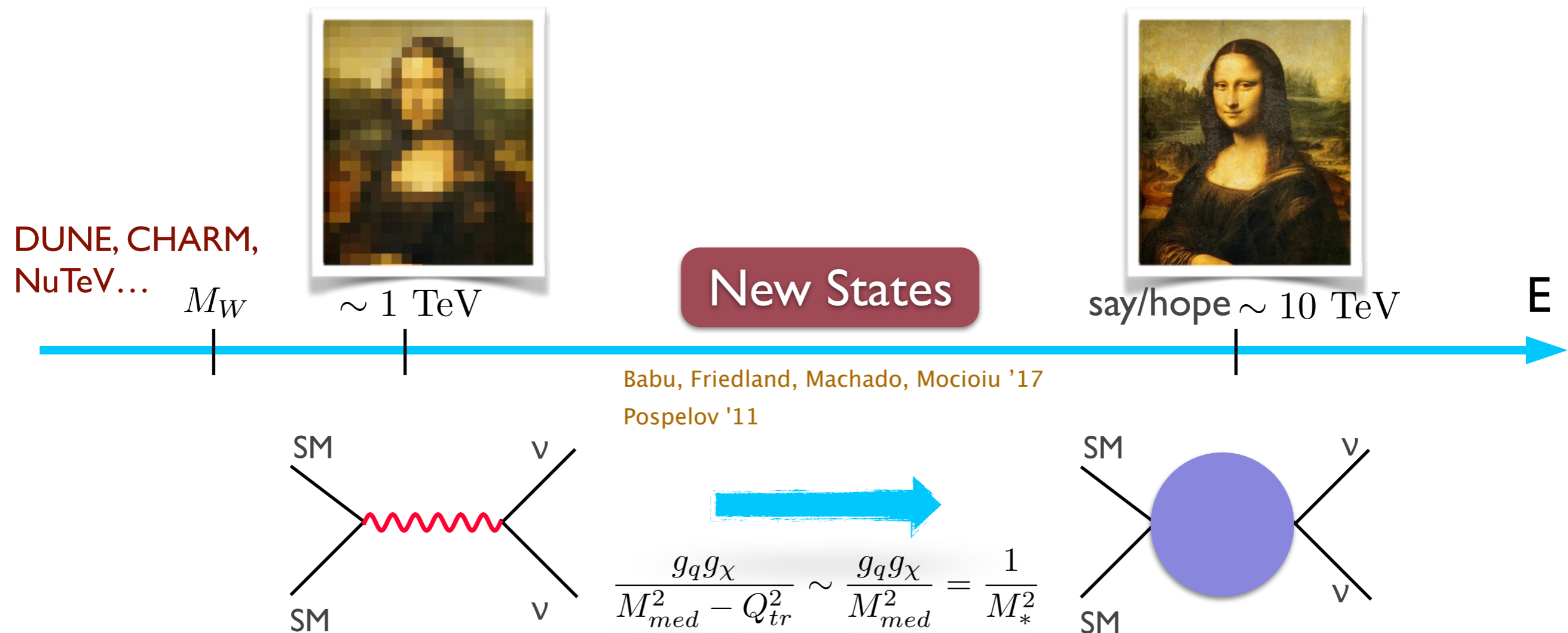
Effective Field Theory Description

- ➔ Model independent approach to NSI searches
- ➔ NSI bounds with few parameters
- ➔ Valid in the presence of an energy gap: $M_* \gg E$



Effective Field Theory Description

- ➔ Model independent approach to NSI searches
- ➔ NSI bounds with few parameters
- ➔ Valid in the presence of an energy gap: $M_* \gg E$



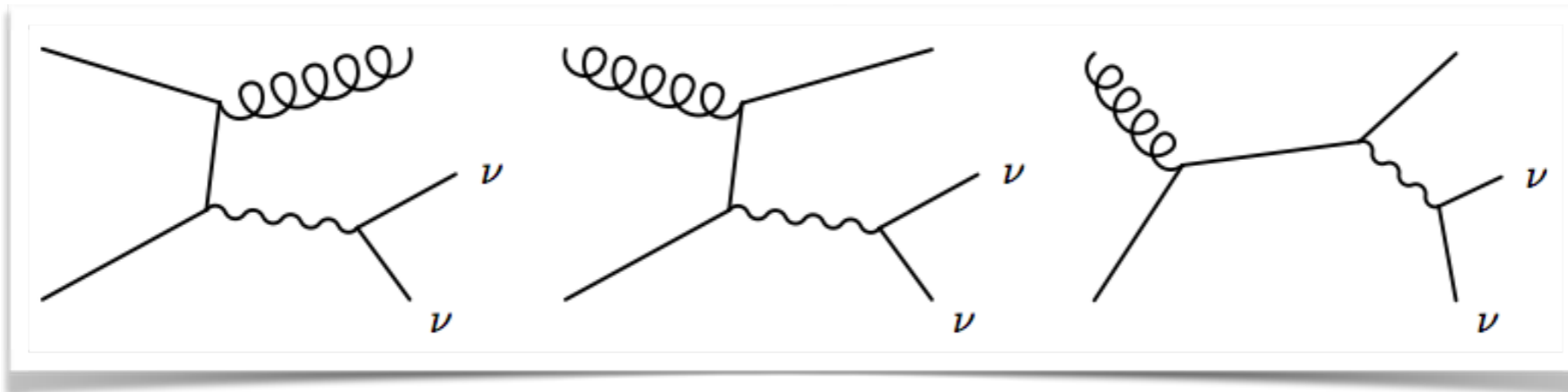
➔ For a given choice of $\sqrt{g_q g_\chi}$, use only events that satisfy $Q_{tr} < M_{med} = \sqrt{g_q g_\chi} M_*$

Simplified Model for NSI

$$\mathcal{L}_{NSI} = -2\sqrt{2}G_F \epsilon_{\alpha\beta}^{fP} (\bar{\nu}_\alpha \gamma_\rho \nu_\beta) (\bar{f} \gamma^\rho P f)$$

$$\mathcal{L}_{NSI} = g_\nu (\bar{\nu} P_L \gamma_\mu \nu) R^\mu + (\bar{q} \gamma_\mu (g_q^V + g_q^A \gamma^5) q) R^\mu$$

Leptophobic Z' model

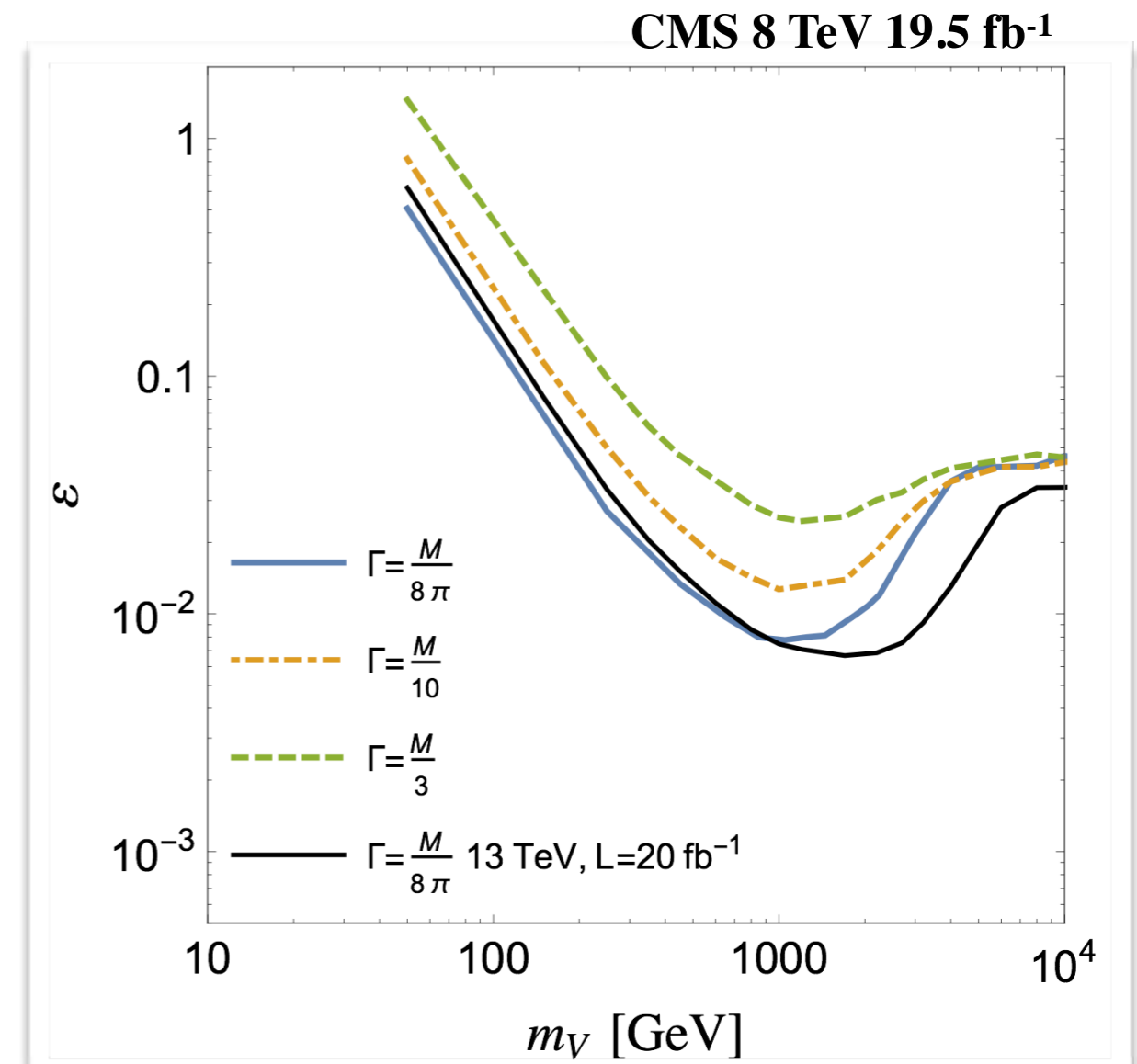


Simplified Model for NSI

$$\mathcal{L}_{NSI} = -2\sqrt{2}G_F\varepsilon_{\alpha\beta}^{fP} (\bar{\nu}_\alpha\gamma_\rho\nu_\beta) (\bar{f}\gamma^\rho P f)$$

$$\mathcal{L}_{NSI} = g_\nu (\bar{\nu}P_L\gamma_\mu\nu) R^\mu + (\bar{q}\gamma_\mu(g_q^V + g_q^A\gamma^5)q) R^\mu$$

- Notice that the bounds saturate for $m_\nu < 1$ TeV. Bounds assume syst. unc. of 5% on background
- Systematic uncertainties saturate the bounds already at the very low luminosity.



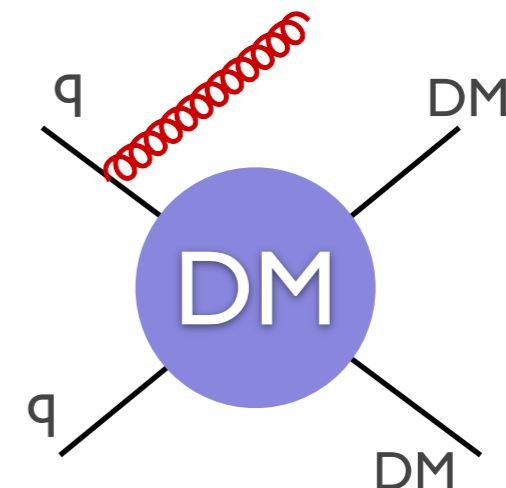
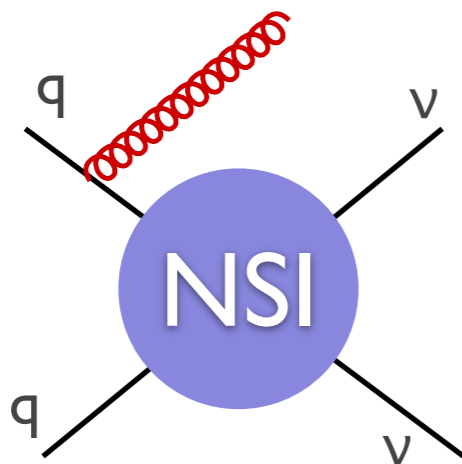
Diogo, Frandsen, Shoemaker '15

Discriminating NSI signature from DM

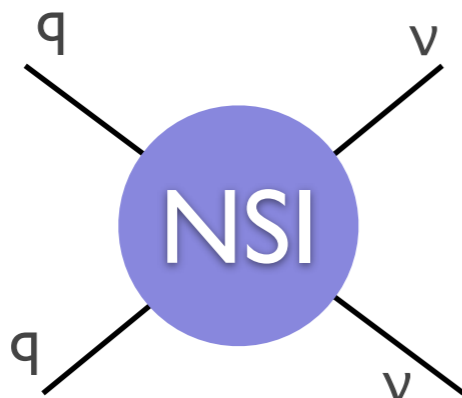
$$\mathcal{L}_{\text{NSI}} = g_\nu (\bar{\nu} P_L \gamma_\mu \nu) R^\mu + (\bar{q} \gamma_\mu (g_q^V + g_q^A \gamma^5) q) R^\mu$$

vs.

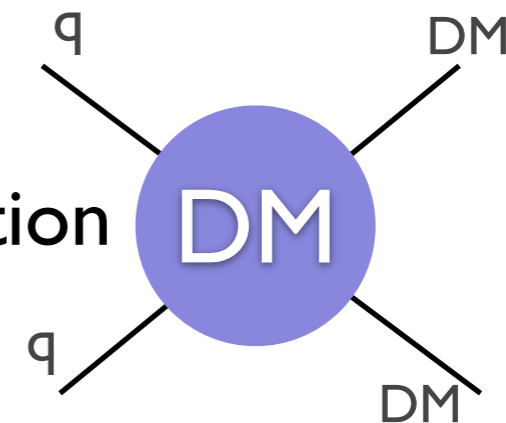
$$\mathcal{L}_{\text{DM}} = g_X (\bar{X} \gamma_\mu X) R^\mu + (\bar{q} \gamma_\mu (g_q^V + g_q^A \gamma^5) q) R^\mu + m_X \bar{X} X$$



Osc. exp.

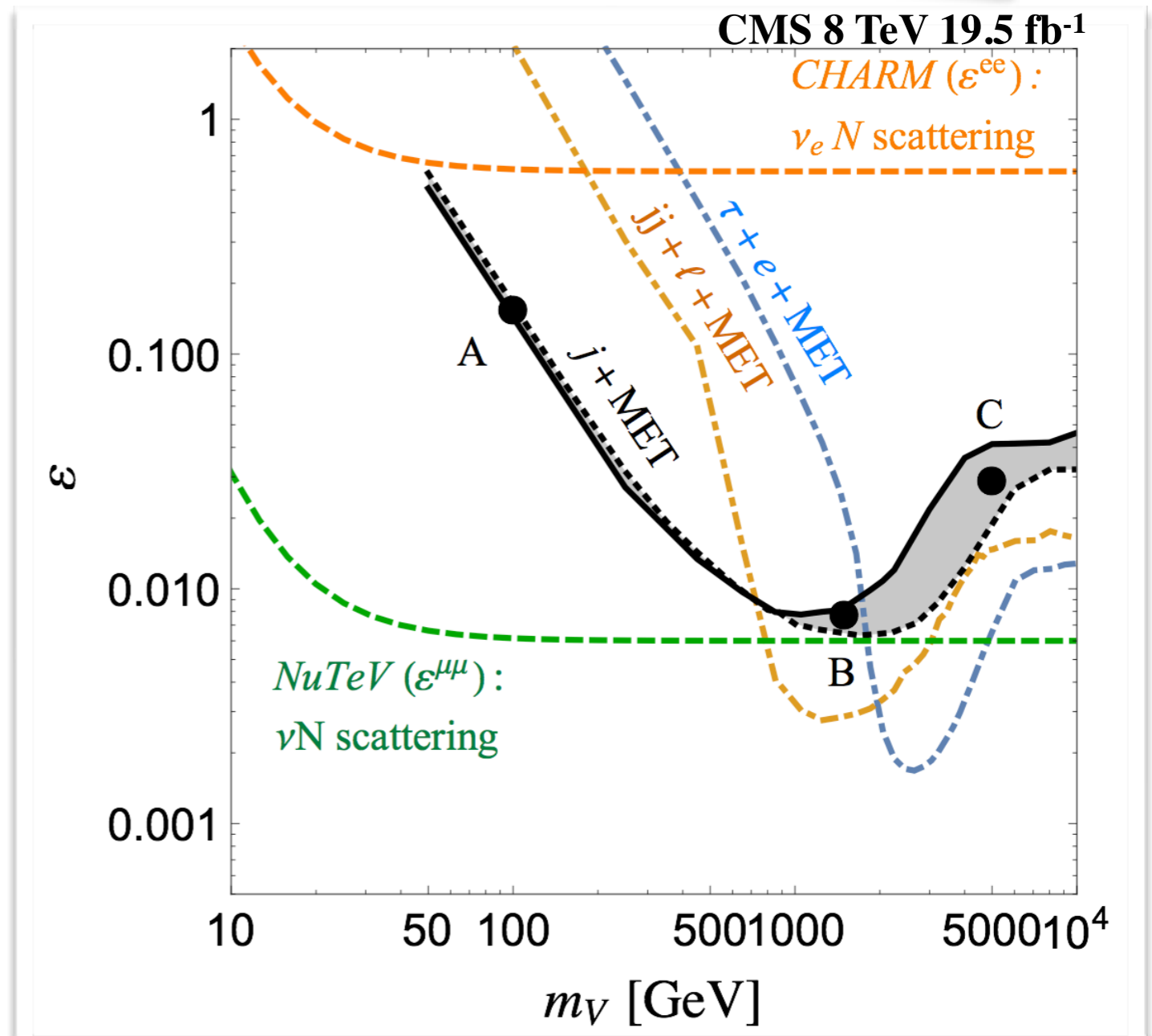
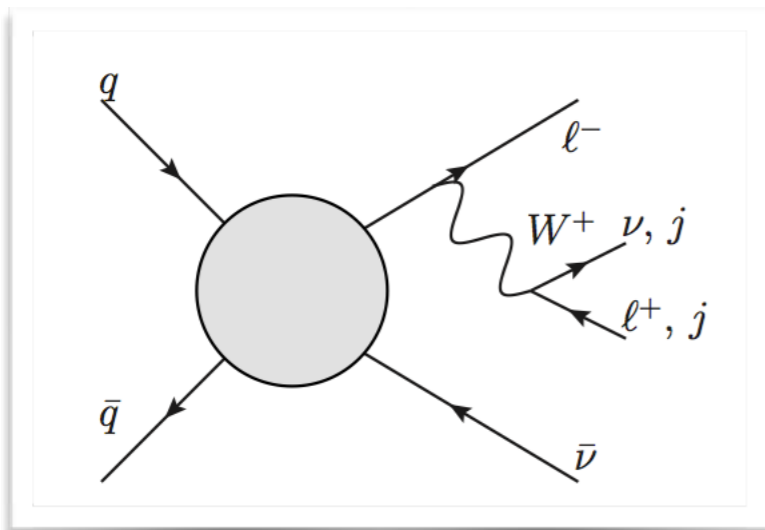


Direct detection



Discriminating NSI signature from DM

Multilepton searches



Diogo, Frandsen, Shoemaker '15

Complementarity:

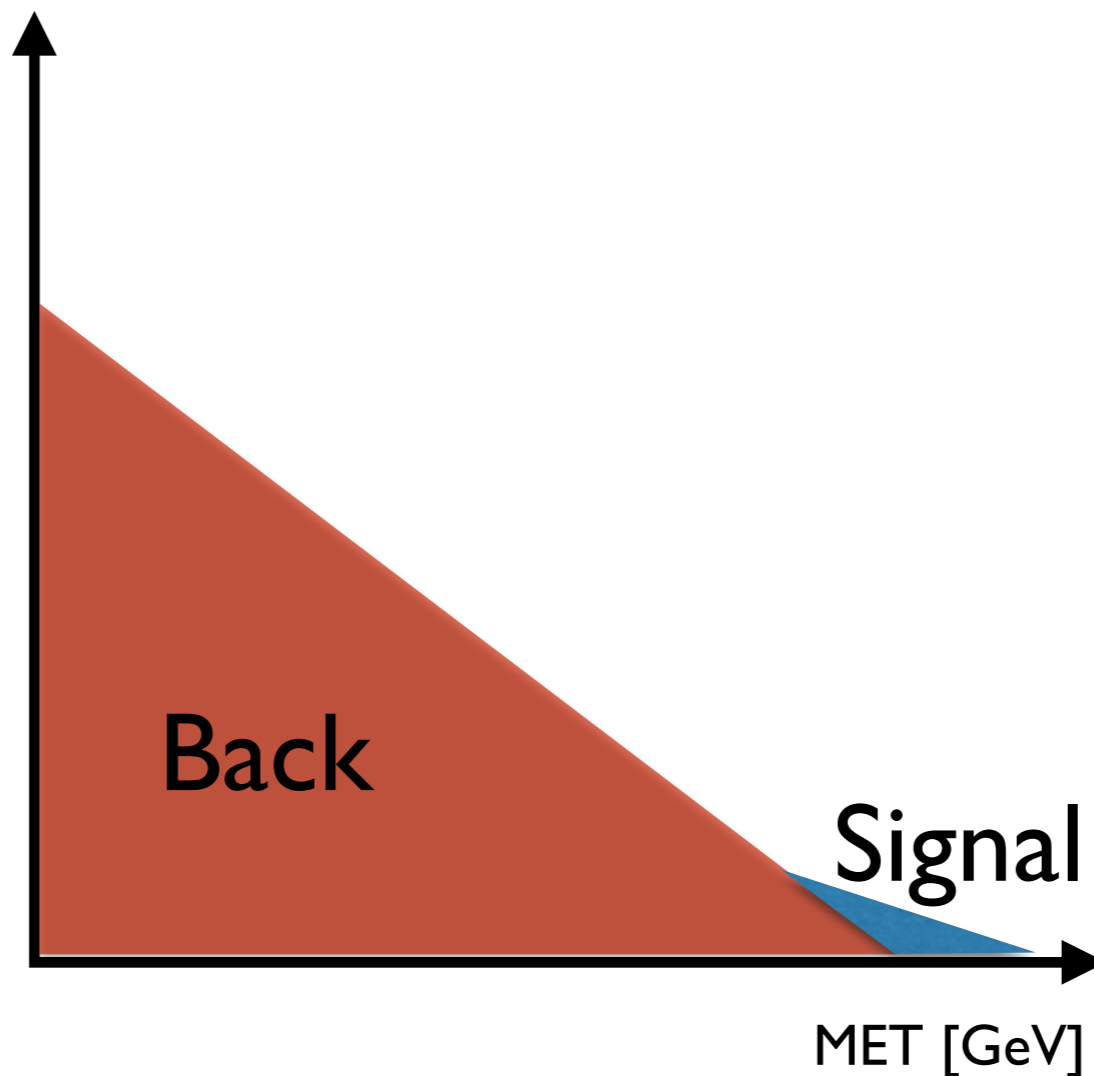
Heavy mediators are strongly constrained by LHC data

Low mediators constrained by low-energy experiments. Bounds from nu-N scattering measurements

If DM itself transforms non-trivially under SU(2) the situation is more complex

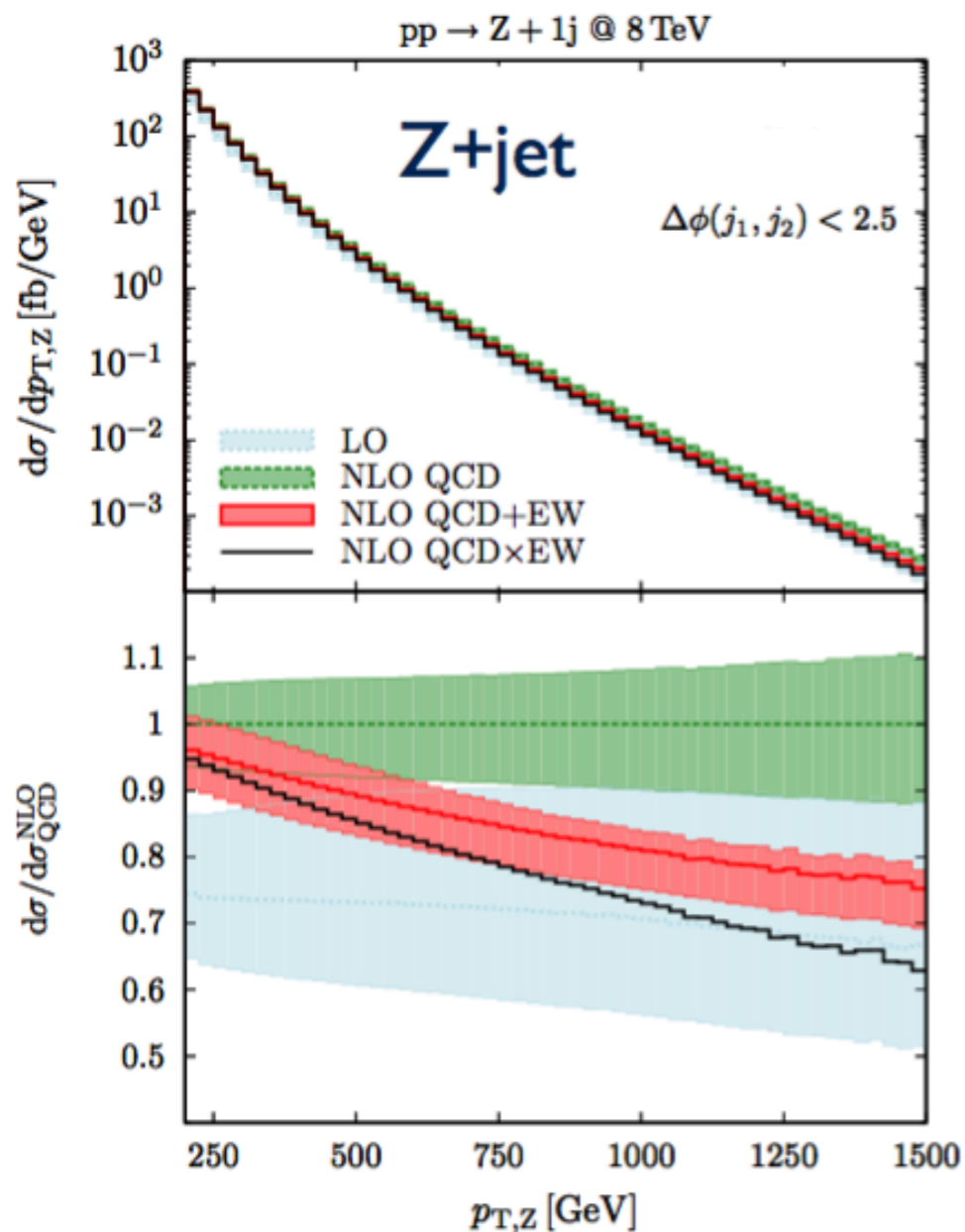
Monojets: systematic uncertainties

- Major limitation for NSI (DM) bounds:
 - Overwhelming SM backgrounds and syst. $\sim 5\%$
 - Signal is a small enhancement in the tail of MET distribution



→ Requires precise estimation of background

Monojets: systematic uncertainties



Lindert, Pozzorini, et al '17

QCD corrections:

Moderate and stable

NLO uncert. 5-10%

EW corrections:

EW corrections $>$ QCD corrections for $p_{T,Z} > 350$ GeV

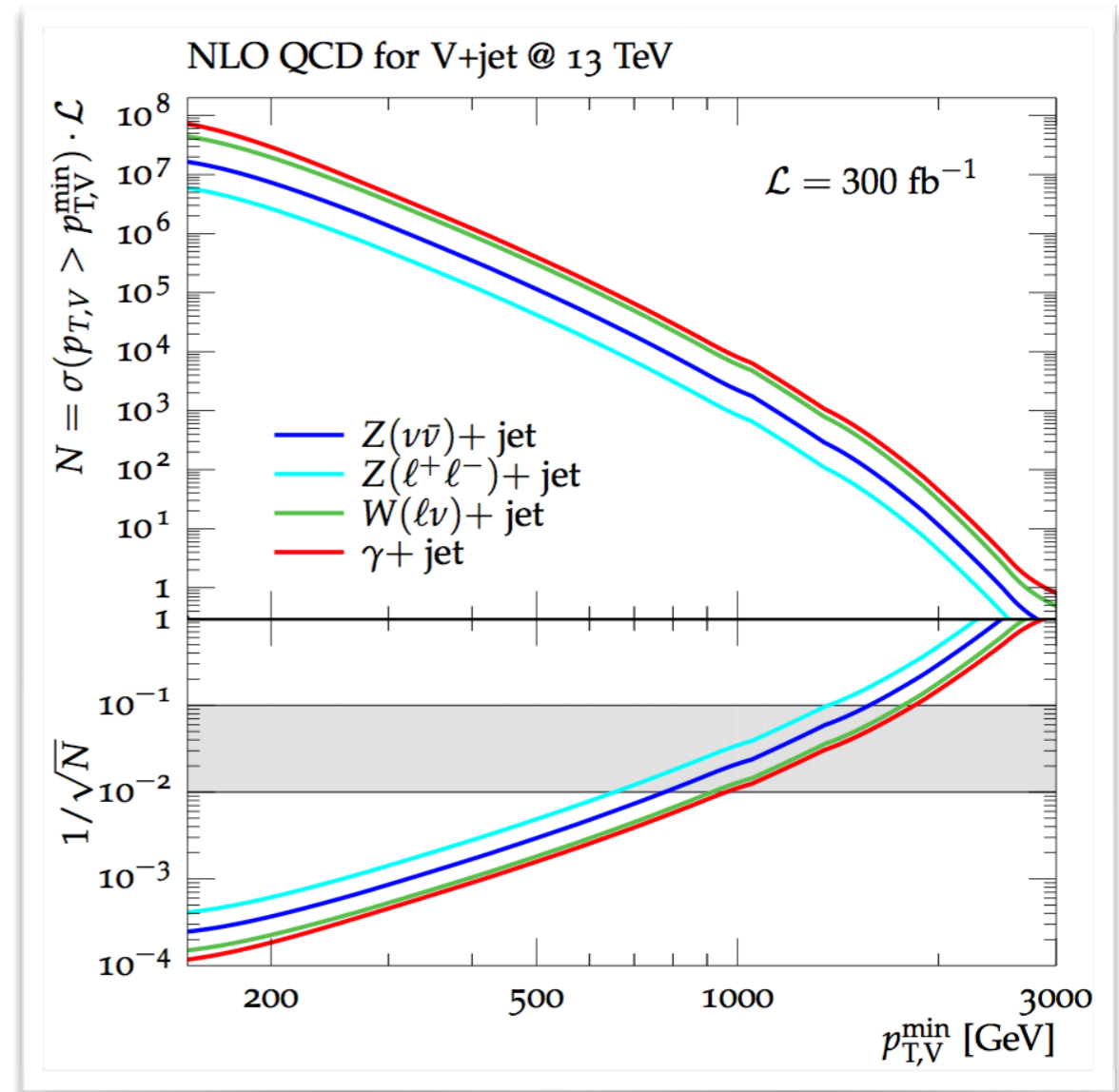
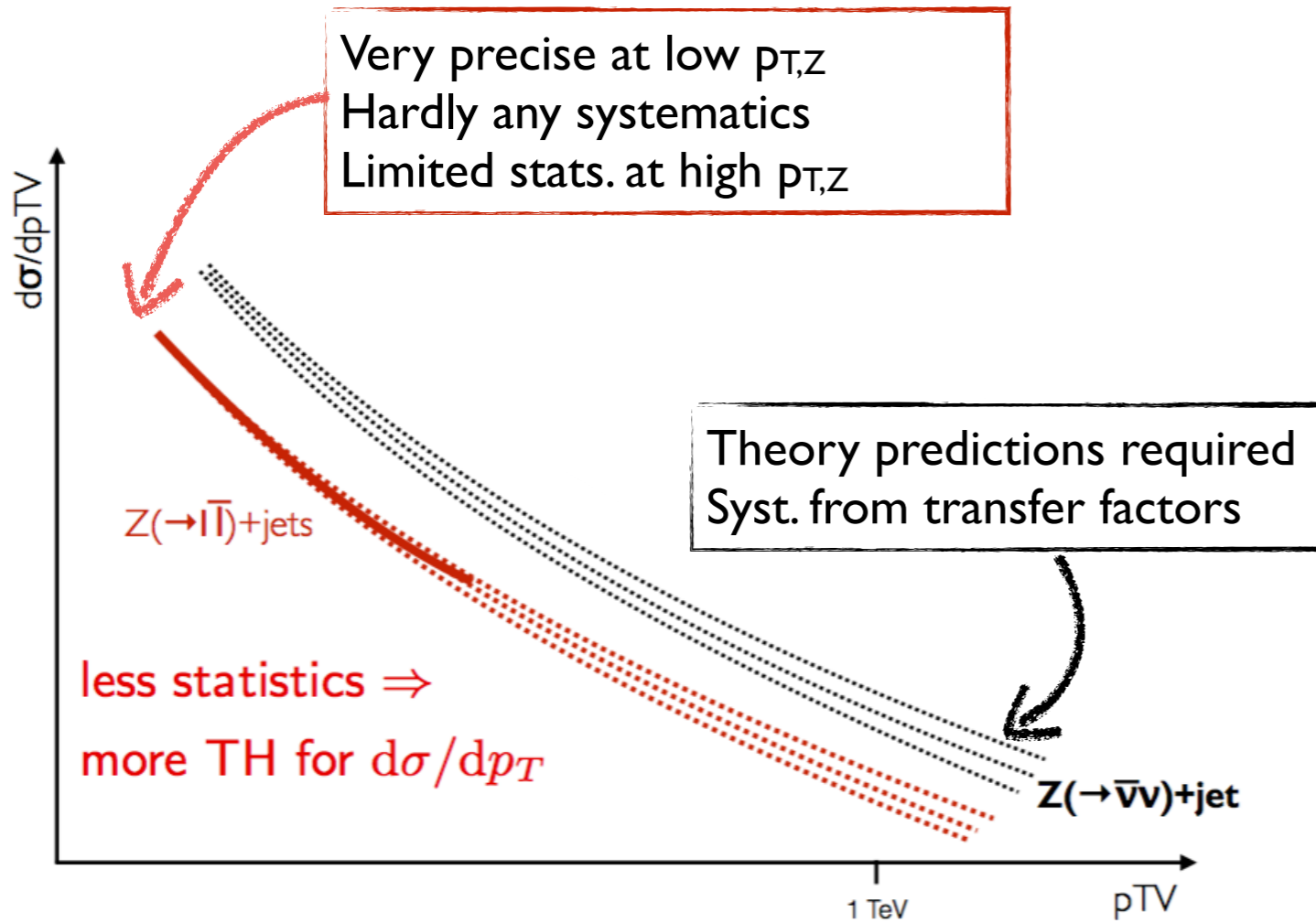
➔ Given that QCD and EW corrections are large, mixed QCD-EW corrections have to be considered

State of the art of MC simulation: NNLO QCD + NLO EW

“Simply” accounting for higher orders is not enough. Uncertainties $O(5\%)$

We need new ingredients to control the errors!

Monojets: systematic uncertainties



Lindert, Pozzorini, et al '17

Summary

- Collider studies result in strong constraints to NSI and are relevant to the global picture
- Monojet and multilepton searches display stronger limits to heavy mediator masses.
They also offer a probe to discriminate NSI from DM (in the case of discovery).
- New theoretical and experimental efforts are significantly improving the background estimation
- ➔ Theory and experimental efforts are transforming the LHC into a precision machine

Thank you for your attention!

