Studies of nonperturbative structure of hadrons from lepton production experiments

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Based on:

PRD 96 (2017), 096006, PRD 95 (2017), 013004, PRD 91 (2015) 073002, PRD 89 (2014) 053001

PRD 87 (2013) 033008, PRD 86 (2012) 113018



ASPECT OF QFT AND LOEWE'S 65 FEST

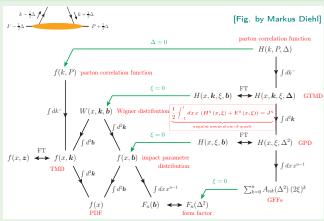
Felix Cumpleaños, Marcelo!



Nucleon (hadron) structure

- ullet Formidable theoretical problem (nonperturbative strongly interacting $ar{q}qg$ ensemble)
- Parton distributions: convenient interface between theory and experiment

Relations between parton distributions



- Helicity of partons/target might be flipped
- Each distribution might depend on flavor

Factorization theorem



Bjorken kinematics

$$Q^2 \to \infty, x_{\rm B} = {\rm const}$$

- $\mathcal{A} \sim \mathcal{C}_{\text{process}} \otimes \mathcal{H}_{\text{target}}$
- Multiparton distributions are suppressed in this

kinematics

Curse of dimensionality

● Only formfactors and PDFs are reasonably measured.

Nucleon (hadron) structure

- Formidable theoretical problem (nonperturbative strongly interacting $\bar{q}qg$ ensemble)
- Parton distributions: convenient interface between theory and experiment

Relations between parton distributions [Fig. by Markus Diehl] parton correlation function $\Delta = 0$ $H(k, P, \Delta)$ f(k, P) parton correlation function $W(x,k,b) \text{ Wigner distribution } \underbrace{\frac{1}{2} \int_{-1}^{1} dx \, x \, (H^q(x,\xi) + E^q(x,\xi)) = J^q}_{\text{angular momentum of quark}} \int d^2k \, \underbrace{\xi = 0}_{\text{FT}} H(x,\xi,b) \xrightarrow{\text{FT}} H(x,\xi,\Delta^2) \text{ GPD}$ Helicity of partons/target might be flipped Each distribution might depend on flavor

Factorization theorem



Bjorken kinematics

$$Q^2 \to \infty, x_{\rm B} = {\rm const}$$

ullet $\mathcal{A} \sim \mathcal{C}_{\mathrm{process}} \otimes \mathcal{H}_{\mathrm{target}}$

Challenge

- No first principle parametrization
- Extracted objects are subject to nontrivial physical constraints

GPD: formal definitions, models, constraints

Constraints on GPD parametrizations:

$$\lim_{\Delta \to 0} H(x, \xi, t) = q(x), \int_{-1}^{1} dx \, H(x, \xi, t) = F(t)$$

$$\int_{-1}^{1} dx \, x^{n-1} H(x, \xi, t) = \sum_{k=1}^{n} a_{kn}(t) \xi^{k}$$

Positivity (in impact space): For $\forall p_{\sigma}(x)$, $\int_{\substack{(1-\xi,\vec{\Delta}_{\perp}/2)\\ |\xi| < x < 1}} \frac{\frac{dx}{(1-x)^5}}{\frac{dx}{(1-x)^5}} p_{\sigma}^*\left(\frac{1-x}{1-\xi}\right) p_{\lambda}\left(\frac{1-x}{1+\xi}\right) F_{\sigma\lambda}\left(x,\xi,\frac{1-x}{1-\xi^2},b_{\perp}\right) \gtrsim 0$

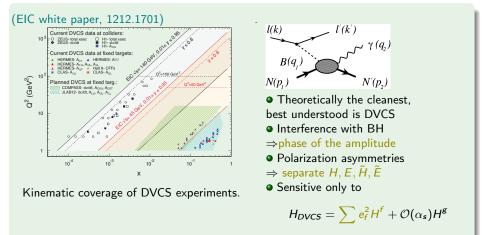
GPDs from nonperturbative models of nucleon structure (χ QSM, NJL, AdS/CFT, ...)

• If "wave functions" of quarks are known, evaluation is quite straightforward:

$$\Rightarrow H_i\left(x,\xi,t\right) \sim \sum_{q} \int d^2 p_\perp \bar{\Phi}_q\left(x-\xi,\vec{p}+\frac{\vec{\Delta}_\perp}{2}\right) \Gamma_i \Phi_q\left(x+\xi,\vec{p}_\perp - \frac{\vec{\Delta}}{2}\right)$$

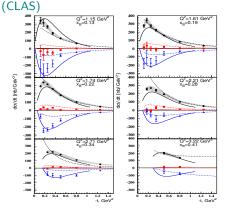
- Advantage: Automatically satisfy polynomiality & positivity constraints
- Yet agreement with experimental data is marginal.
- ⇒Phenomenological approach predominant in the literature (EPJC 59, 809; EPJC 39, 1; PRD 72, 054013; NPB 841, 1, ...)

GPD extraction from DVCS



 DVMP may give access to GPD flavor structure, but theoretically is more complicated

Challenges in GPD extraction from pion production



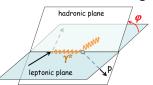
$$\begin{split} \frac{d^4\sigma}{dQ^2dx_Bdtd\phi_\pi} &= \frac{\Gamma(Q^2,x_B,E)}{2\pi} (\sigma_T + \epsilon\sigma_L) \\ &+ \epsilon\cos2\phi_\pi\sigma_{TT} \\ &+ \sqrt{2\epsilon(1+\epsilon)}\cos\phi_\pi\sigma_{LT}) \end{split}$$

Tw-2 contribution is small, probes

$$\sigma_L \sim \left| \{ \tilde{H}, \tilde{E} \} \otimes \phi_{\mathbf{2};\pi} \right|^2$$

 \Leftarrow Dependence on azimuthal angle ϕ_{π} between ep and πp planes

Should not exist in leading twist



 Signals that tw-3 contributions are pronounced

$$\sigma_{TT} \sim \left| \left\{ H_T, E_T \right\} \otimes \phi_{3;\pi} \right|^2$$

$$\sigma_{LT} \sim \left| \left\{ H_T, E_T \right\} \otimes \phi_{3;\pi} \right|^2$$

 \Rightarrow This channel requires significantly larger Q^2 to access GPDs



Challenge in GPD extraction from vector mesons

• Probe unpolarized GPDs H, E, smaller tw-3 contributions

Vector meson wave function unknown

• controlled by confinement (not SCSB), depends heavily on the model

Popular parametrizations:

AdS/CFT wave function

$$\begin{split} \varphi_q^{(i)}(x,\mathbf{k}_\perp) \; &= \; N_q^{(i)} \, \frac{4\pi}{\kappa} \, \sqrt{\frac{\log(1/x)}{1-x}} \, \sqrt{f_q^{(i)}(x) \bar{f}_q(x)} \\ & \times \; \exp\!\left[-\frac{\mathbf{k}_\perp^2}{2\kappa^2} \frac{\log(1/x)}{(1-x)^2} \, \bar{f}_q(x) \right]. \end{split}$$

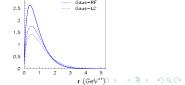
• $f_q(x)$, $f_{\overline{q}}(x)$ -unknown functions, can be fixed from (hypothetical) DIS on ρ -mesons

Uncertainty in WF translates into significant uncertainty in extraction of GPDs from this channel

Boosted Gaussian WF

$$\begin{split} \varphi_q^{\text{BG}}\left(x,\,\mathbf{k}_\perp\right) &= \mathcal{N}_\lambda 2[x(1-x)]^{b_\lambda/2} \sqrt{2\pi R_\lambda^2} \exp\left(\frac{m_f^2 R_\lambda^2}{2}\right) \\ &\times \exp\left(-\frac{\mathbf{k}_\perp^2 + m_f^2}{8[x(1-x)]^{b_\lambda}} R_\lambda^2\right) \end{split}$$

everything except x and k_{\perp} are free parameters



What we suggest?

• Charged current π/K -production

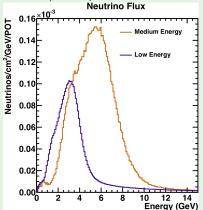
Advantages

- ullet V-A structure of interaction \Rightarrow probes unpolarized ("large") GPDs
- H, E; much smaller contamination by higher twist corrections
- Good knowledge of pion and kaon WF, closeness of wave functions due to SCSB⇒can extract full flavor structure of GPD

Where such processes can be studied?

MINERvA@Fermilab

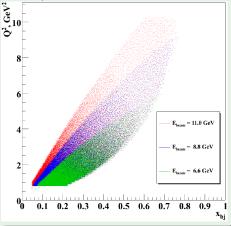
- Extremely large luminosity
- ullet Both u_{μ} and $ar{
 u}_{\mu}$ can be used



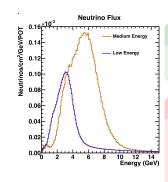
 Analysis of data in Bjorken kinematics has already started UTFSM MINERVA group: Jonathan Miller et al.

Jefferson Laboratory

- Monochromatic beam, $E_e = 11 \text{ GeV}$
- Luminosity $\mathcal{L} = 10^{36} \mathrm{cm}^{-2} s^{-1}$
- Beam/target can be polarized



MINERvA experiment (neutrinos)



Advantages

- Extremely large luminosity
- ullet Both u_{μ} and $ar{
 u}_{\mu}$ can be used (W^{\pm} -induced production)

Challenges

- Beam not monochromatic, should consider spectrum averaged observables
- Detector=extended nuclear target, nuclear effects are important
- Accessible Q^2 is not very large, loop corrections might be pronounced (will DUNE improve the resolution?)

Flavor combinations of GPDs probed by various processes

(PRD 86 (2012) 113018)

• Experimentally easiest: $\nu p \to \mu^- \pi^+ p$, $\bar{\nu} p \to \mu^+ \pi^- p$ (ongoing analysis by MINERvA group @USM)

$$\mathcal{A}_{\nu p \to \mu^{-}\pi^{+}p}^{(\mathrm{LO})} \sim \int_{-1}^{1} dx \left(\frac{H_{d}\left(x,\,\xi\right)}{x-\xi+i0} + \frac{H_{u}\left(x,\,\xi\right)}{x+\xi-i0} \right)$$

$$\mathcal{A}_{\bar{\nu}\rho\to\mu^{+}\pi^{-}\rho}^{(\mathrm{LO})} \sim \int_{-1}^{1} dx \left(\frac{H_{u}\left(x,\,\xi\right)}{x-\xi+i0} + \frac{H_{d}\left(x,\,\xi\right)}{x+\xi-i0} \right)$$

• Can probe 29 CC processes in total if use SU(3) flavour relations for transition GPDs $H_{p\to Y}$, e.g.

$$H_{p\rightarrow n}=H_{u}\left(x,\xi,t\right) -H_{d}\left(x,\xi,t\right)$$

• In the NLO the coefficient functions $\frac{1}{x+\xi \pm i0}$ get much more complicated

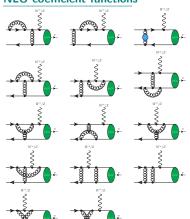


Loop corrections

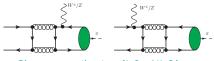


- ep experience: loop corrections are large in this kinematics (JETPL 80, 226; EPJC 52, 933)
- Challenge: separate corrections to coefficient function and GPD/DA $_{\pi}$ evolution kernel (scale μ_F)

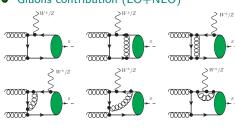
NLO coefficient functions



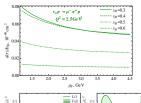
Sea quarks contribution



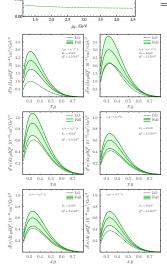
Gluons contribution (LO+NLO)

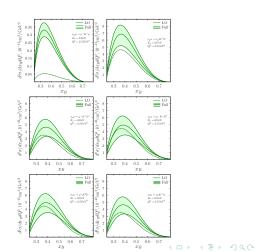


Loop corrections



- ullet Weak dependence on factorization scale for $\mu_{F}\gtrsim 3$ GeV
- Scale choice: $\mu_R = \mu_F = Q$
- ullet Estimates of NNLO corrections: $\mu_R = \mu_F \in (0.5, 2)Q$
- \bullet NLO corrections increase all the cross-sections ${\gtrsim}50\%$
- ⇒NNLO corrections are needed!



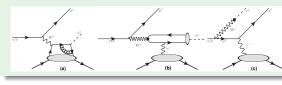


Contaminations by twist-3 & Bethe-Heitler mechanisms

Twist-3 contributions

• Quark spin flip \Rightarrow probe transversity GPDs H_T , E_T , \tilde{H}_T , \tilde{E}_T (large at CLAS6)

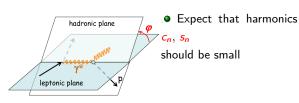
Bethe-Heitler mechanism (diagrams b, c)

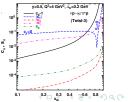


- \bullet formally is suppressed by $\alpha_{\it em}$
- kinematically is enhanced by $Q^2/(t \cdot \alpha_s^2(Q^2))$

Both mechanisms generate azimuthal asymmetry

$$\frac{d^4 \sigma^{(\text{tot})}}{dt \, dQ^2 d \ln \nu \, d\varphi} = \frac{1}{2\pi} \frac{d^3 \sigma^{(DVMP)}}{dt \, dQ^2 d \ln \nu} \times \sum_{n} \left(c_n \cos n\varphi + s_n \sin n\varphi \right)$$



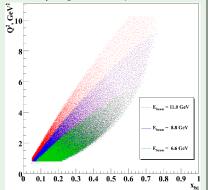


Charged current studies in ep experiments

• (HERA: luminosity insufficient for charged current exclusive processes)

Kinematic coverage of JLAB

- ullet Monochromatic beam, $E_e=11$ GeV
- Luminosity $\mathcal{L} = 10^{36} \mathrm{cm}^{-2} s^{-1}$
- Beam/target can be polarized



Suggested process: $ep o u_e \pi^- p$

• Neutrino ν_e momentum reconstructed via momentum conservation

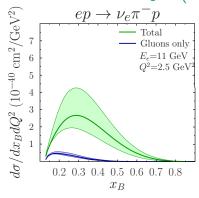
$$p_{
u}=p'+p_{\pi}-p-p_{e}$$

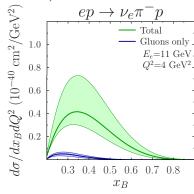
-final hadrons are charged, kinematics resolution should be good.

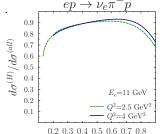
 Potentially can extend also to other members of mesonic and baryonic flavour multiplets,

$$SU(3)_f$$
-relations \Rightarrow GPD flavour combinations

Results for the $e \rightarrow \nu_e M$ (NLO in α_s)

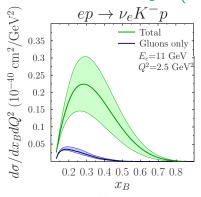


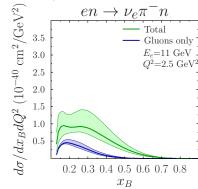


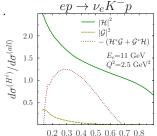


- For pions with beam luminosity $L\sim 10^{35}{\rm cm}^{-2}s^{-1}$ expect \sim 40 events/day/bin (1 GeV bins in Q^2 assumed)
- Gluons give minor contribution and slightly *decrease* the cross-section (interference term q-g is negative)
- Mostly sensitive to GPD H_u , H_d (\gtrsim 80% of result).

Results for the $e \rightarrow \nu_e M$ (NLO in α_s)







- For K-mesons, suppression by an order of magnitude (Cabibbo forbidden), smaller statistics
- \bullet Sizeable $\underline{negative}$ contribution from interference $\mathcal{H}^*\mathcal{G}+\mathcal{G}^*\overline{\mathcal{H}}$
- For neutrons the cross-section is of the same order ($\sim 40\%$ less than in $ep \to \nu_e \pi^- p$), but kinematics reconstruction might be poorer

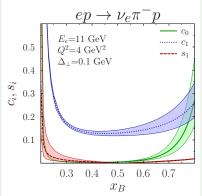
Contaminations by twist-3 & Bethe-Heitler mechanisms

• Generate azimuthal asymmetry, quantify effect in terms of angular harmonics

$$\frac{d^4\sigma^{(\rm tot)}}{dt\,dQ^2d\ln\nu\,d\phi} = \frac{1}{2\pi}\,\frac{d^3\sigma^{(DVMP)}}{dt\,dQ^2d\ln\nu} \times \sum \left({\it C}_{\it n}\cos n\phi + {\it s}_{\it n}\sin n\phi \right) \label{eq:cosnop}$$

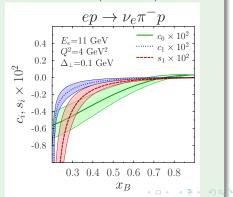
Twist-3 effects

.• Quark spin flip \Rightarrow probe (poorly known) transversity GPDs H_T , E_T , \tilde{H}_T , \tilde{E}_T (large at CLAS6)



Bethe-Heitler mechanism

- interaction with hadron via elastic *t*-channel photon exchange only
- suppressed by α_{em} , kinematically is enhanced by $Q^2/(t \cdot \alpha_s^2(Q^2))$



Summary

- Charged current Deeply Virtual Pion Production can be used as an additional source of information on proton structure (its GPDs)
 - \star Can be studied at νp (ongoing analysis) and ep experiments thanks to large luminosity of modern experiments.
 - \star Has sensitivity to unpolarized GPDs H, E (large components); expect small contamination by higher twist and Bethe-Heitler corrections.
 - \star NNLO analysis of coefficient functions is needed if Q^2 is not very large

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THANK YOU FOR YOUR ATTENTION!