J/ψ production in pp and Heavy Ion Collisions

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work in progress

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Why do we study J/ψ production in heavy-ion collisions?

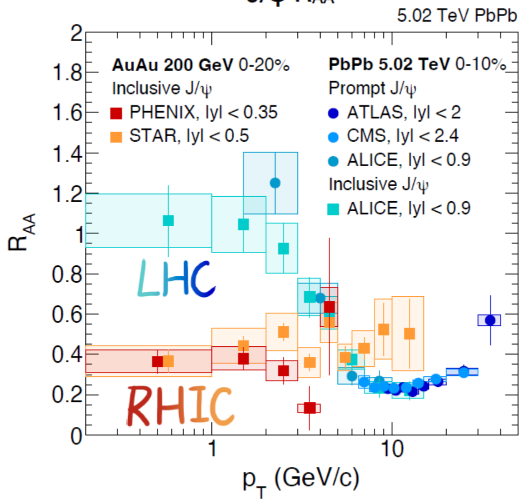
J/ψ mesons

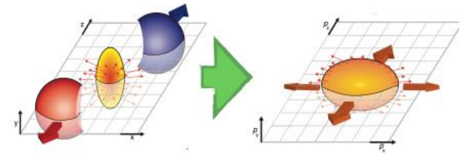
- are a hard probe: tests quark gluon plasma from creation to hadronization
- no consistent microscopical theory available yet

  which are not understood yet
- show quite different results for key observables at RHIC and LHC:

$$R_{AA}(p_T) = \frac{dN_D^{AA}/dp_T}{\langle N_{coll}\rangle \, dN_D^{pp}/dp_T}$$

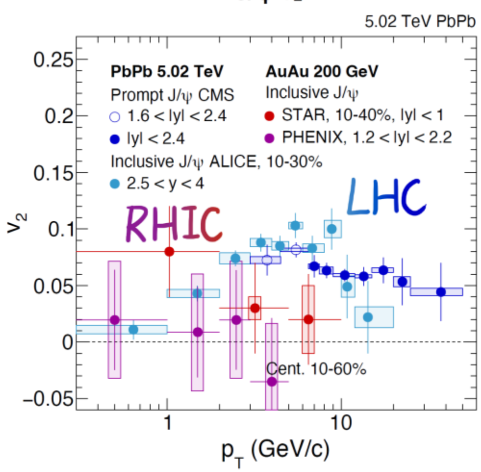
J/ψ R_{AA}





$$\frac{dN}{d\phi} \propto 1 + 2v_1 cos\phi + 2v_2 cos(2\phi).....$$

J/ψ v_2

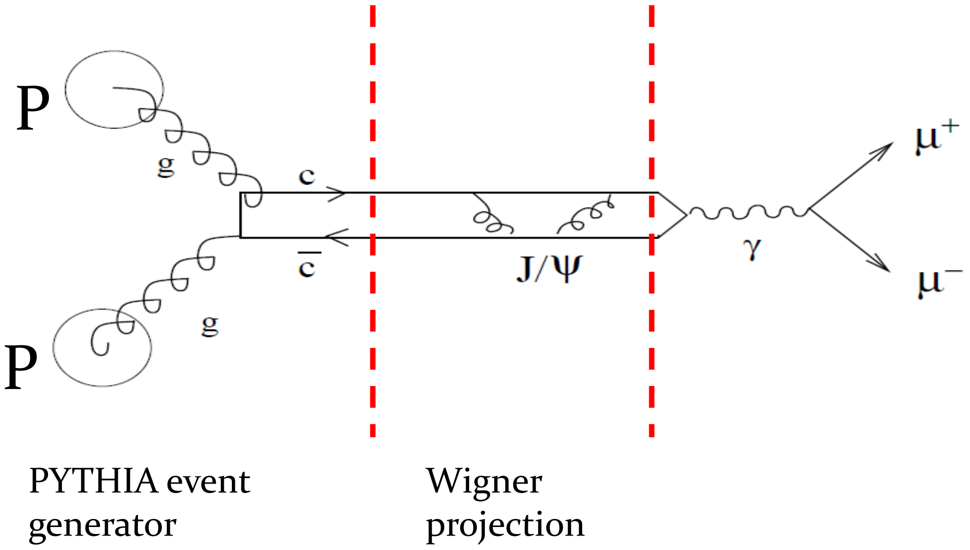


J/ψ production in p+p collisions

How to describe a composite object if perturbative QCD can only deal with quarks and gluons

Need non perturbative input ➔ assumptions.

Our approach: Wigner density formalism (as successful at lower energies)



Wigner Density Formalism

Interaction depends on relative coordinates only, -> plane wave of CM

Starting point: Wave function (w.f.) of the relative motion of state i: $|\Phi_i>$

w.f → density matrix $|\Phi_i><\Phi_i|$

Fourier transform of density matrix in relative coord. → Wigner density of $|\Phi_i>$ (close to classical phase space density)

$$\Phi_i^W(\mathbf{r},\mathbf{p}) = \int d^3y e^{i\mathbf{p}\cdot\mathbf{y}} < \mathbf{r} - \frac{1}{2}\mathbf{y}|\Phi_i><\Phi_i|\mathbf{r} + \frac{1}{2}\mathbf{y}> .$$

$$\mathrm{R} = \frac{\mathrm{r}_1 + \mathrm{r}_2}{2}, \quad \mathrm{r} = \mathrm{r}_1 - \mathrm{r}_2,$$
$$\mathrm{P} = \mathrm{p}_1 + \mathrm{p}_2, \quad \mathrm{p} = \frac{\mathrm{p}_1 - \mathrm{p}_2}{2}.$$

$$n_i(\mathbf{R},\mathbf{P}) = \int d^3r d^3p\ \Phi_i^W(\mathbf{r},\mathbf{p}) n^{(2)}(\mathbf{r_1},\mathbf{p_1},\mathbf{r_2},\mathbf{p_2})$$

$n^{(2)}(\mathbf{r_1},\mathbf{p_1},\mathbf{r_2},\mathbf{p_2})$ two body c cbar density matrix

pp: In momentum space given by PYTHIA (Innsbruck tune)

In coordinate space $\sim r^2 \exp\left(-\frac{r^2}{2\delta^2}\right)$ $\delta^2 = \langle r^2\rangle/3 = 4/(3m_c^2)$

Wigner Density Formalism

If there are N c cbar pairs in the system the phase space density of states $|\Phi_i>$

$$n_i(\mathbf{R},\mathbf{P}) \;=\; \sum \int \frac{d^3r d^3p}{(2\pi)^3}\, \Phi_i^W(\mathbf{r},\mathbf{p}) \prod_j \int \frac{d^3r_j d^3p_j}{(2\pi)^3}$$

$$n^{(N)}(\mathbf{r_1},\mathbf{p_1},\mathbf{r_2},\mathbf{p_2},...,\mathbf{r_N},\mathbf{p_N}) \qquad (5)$$

Sum over all possible ccbar pairs after integration of the relative coordinates
Integration over all N-2 left particles.

Multiplicity of $|\Phi_i>$

$$P_i = \int \frac{d^3R d^3P}{(2\pi)^3} n_i(\mathbf{R},\mathbf{P})$$

Momentum distribution

$$\frac{dP_i}{d^3P} = \int \frac{d^3R}{(2\pi)^3} n_i(\mathbf{R},\mathbf{P})$$

Wigner Density Formalism

The Wigner density of the state $|\Phi_i>$ is different for S and P states

We choose the simplest possible parametrization

$$\Phi_S^W(\mathbf{r},\mathbf{p}) = 8\frac{D}{d_1 d_2}\exp\left[-\frac{r^2}{\sigma^2}-\sigma^2 p^2\right],$$

$$\Phi_P^W(\mathbf{r},\mathbf{p}) = \frac{16}{3}\frac{D}{d_1 d_2}\left(\frac{r^2}{\sigma^2}-\frac{3}{2}+\sigma^2 p^2\right) \times \exp\left[-\frac{r^2}{\sigma^2}-\sigma^2 p^2\right],$$

$$r = r_c - r_{\bar{c}}$$

$$p = \frac{p_c - p_{\bar{c}}}{2}$$

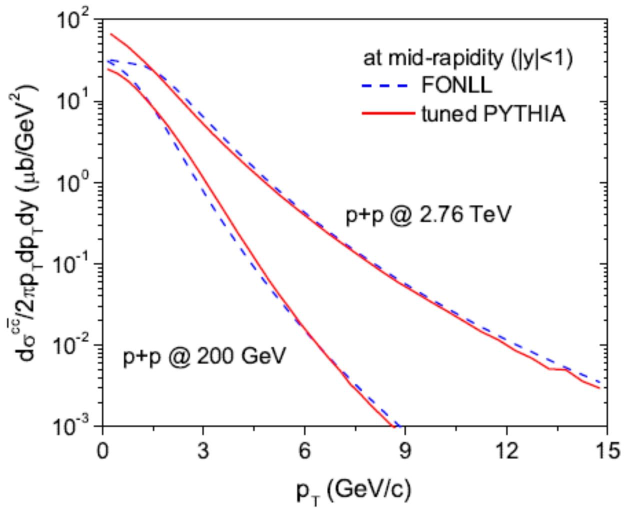
D : degeneracy of Φ
d_1 : degeneracy of c
d_2 : degeneracy of cbar
σ ~ radius of Φ

Where σ reproduces the rms radius of the vacuum c cbar state $|\Phi_i>$

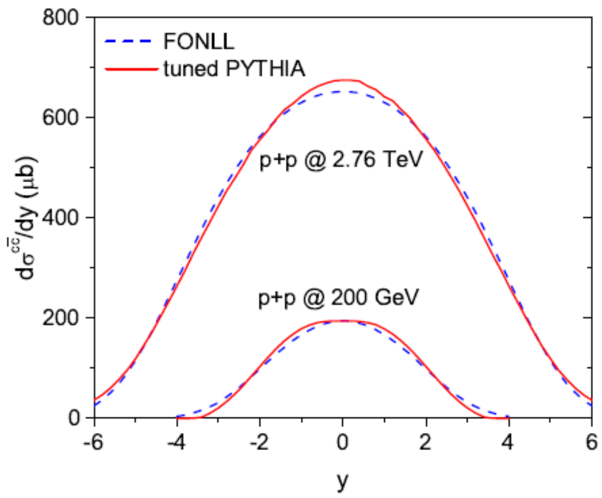
$$\Phi = J/\psi(1S), \qquad \chi_c(1P),\ \psi'(2S)$$

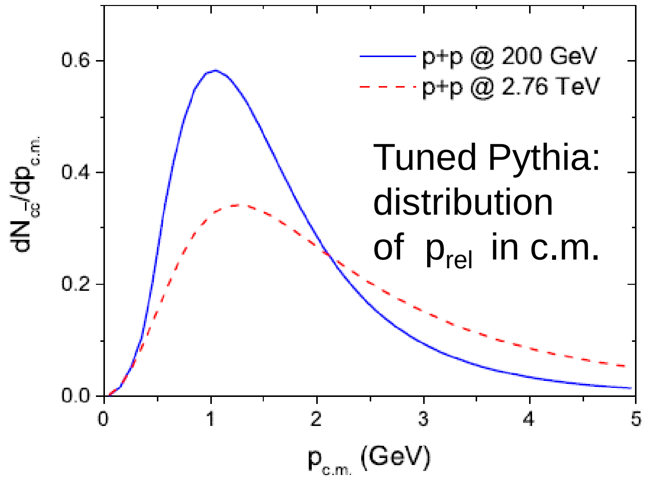
Wigner Density Formalism

The (Innsbruck) tuned PYTHIA reproduces FONLL calculations
but in addition it keeps the ccbar correlation (not known in FONLL)



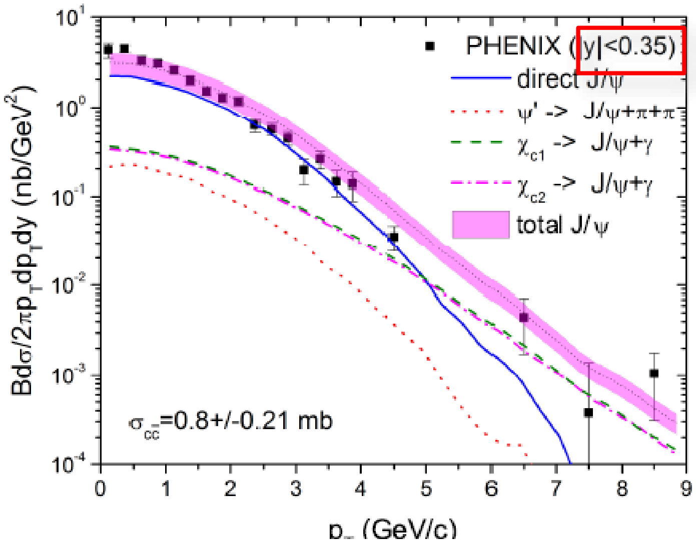
Distribution of charm quarks agrees

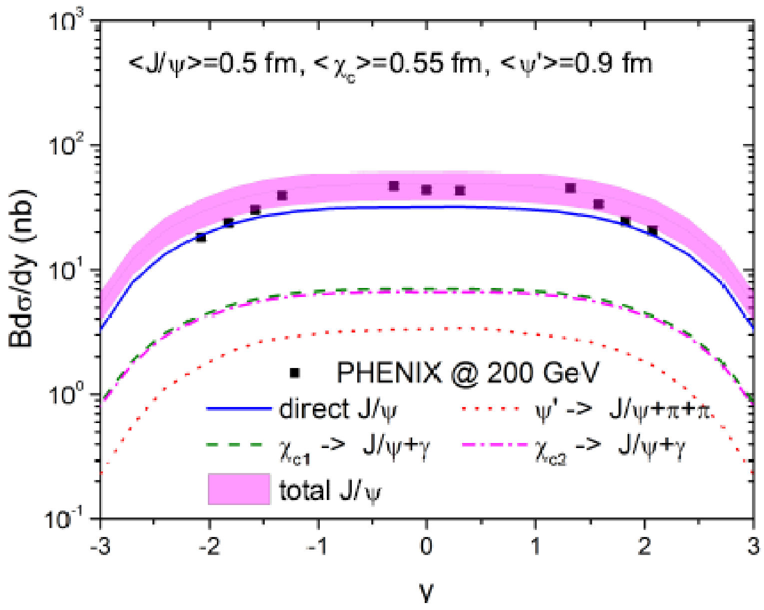


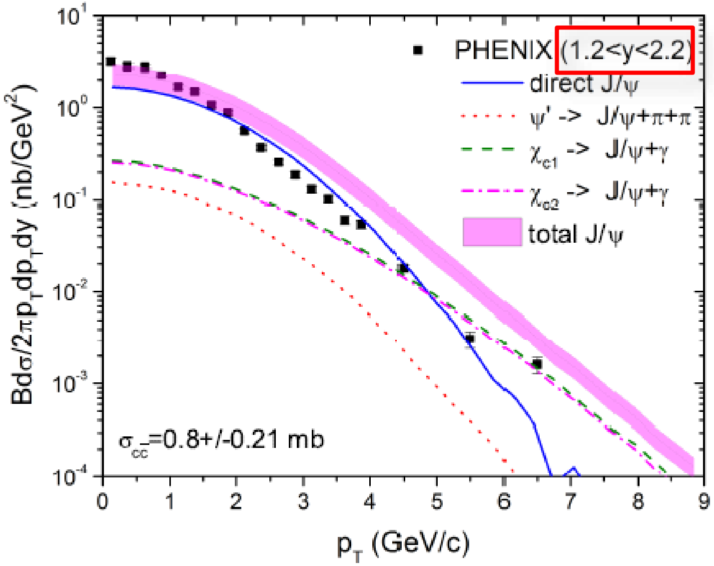


but quite different relative momenta at RHIC and LHC

pp: comparison with Phenix data





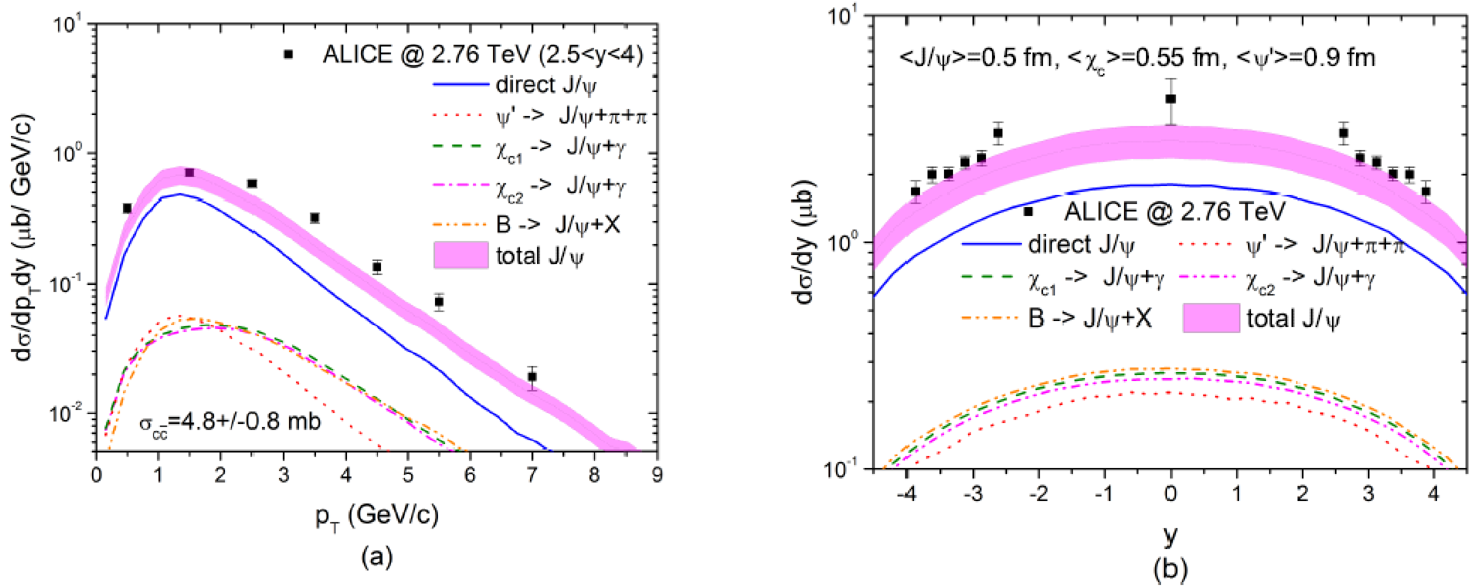


Good agreement for
rapidity spectra
pt spectra at |y| < 0.35
pt spectra at 1.25 < |y| < 2.2

Feeding at RHIC not very important

pp: comparison with ALICE data

same charmonia radii as at RHIC



Important contribution of feeding

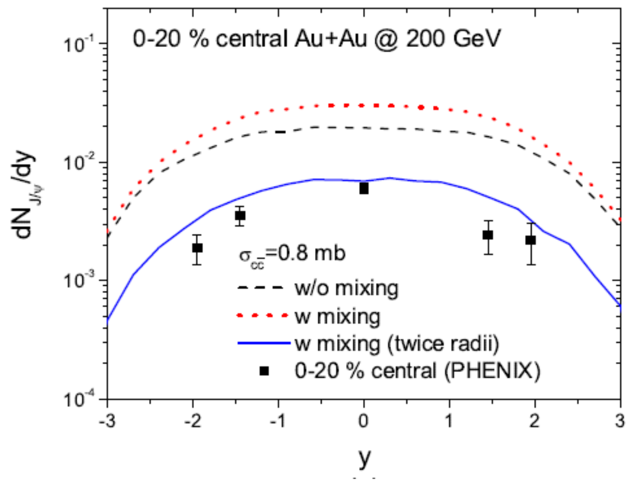
The observed J/ψ data in pp at RHIC and LHC can be well described by Wigner dens.

AA collisions

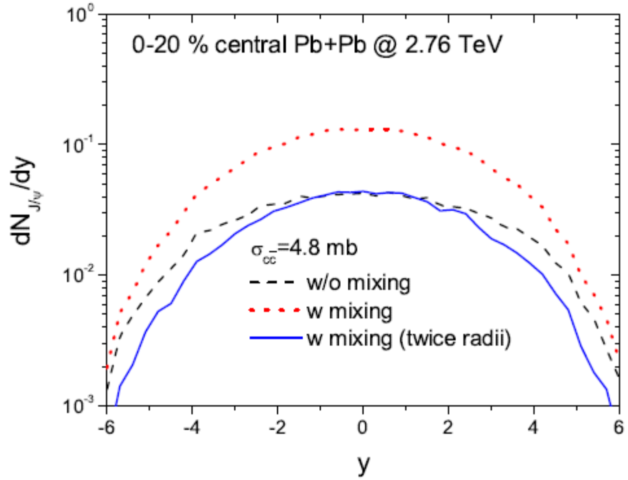
AA: without any QGP

Without the formation of a QGP we expect a (large) enhancement of the J/ψ production because c and cbar from different vertices can form a J/ψ.

enhancement 46%



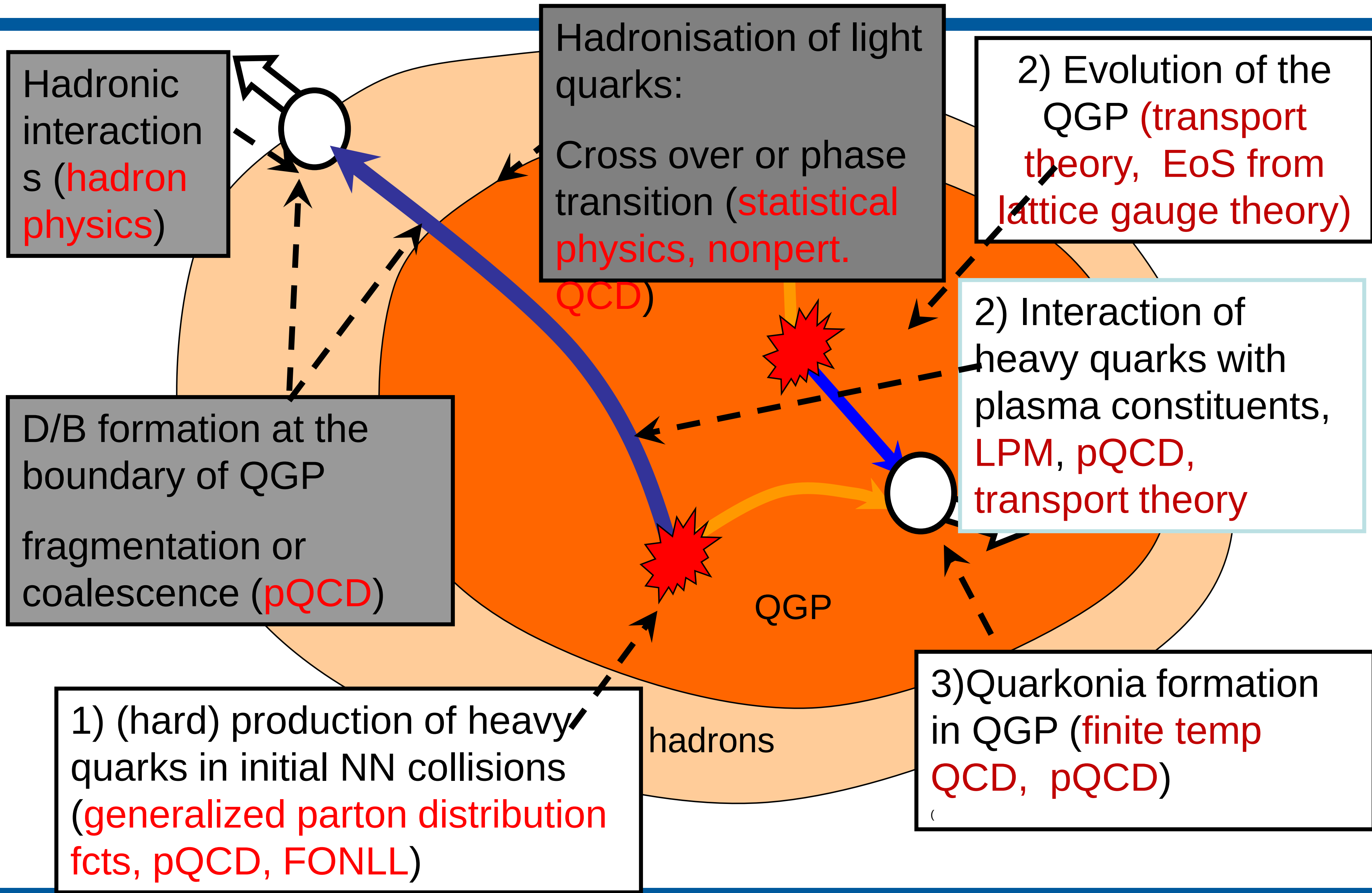
enhancement 270%



but experiments show suppression

Reason: J/ψ production in HI collisions is a very complex process

Complexity of heavy quark physics in HI reactions



The different processes which influences the J/ψ yield

- Creation of heavy quarks (shadowing)
- J/ψ are first unstable in the quark gluon plasma and are created later
- c and cbar interact with the QGP
- c and cbar interact among themselves (<-lattice QCD)
- If QGP arrives at the dissociation temperature T_{diss} , stable J/ψ are possible
- J/ψ creation ends when the QGP hadronizes
- J/ψ can be further suppressed or created by hadronic interaction (task for the future -> Torres-Rincon)
- There are in addition J/ψ from the corona (do not pass the QGP)

The model we developed follows the time evolution of all c and cbar quarks

Is based, as our pp calculation, on the Wigner density formalism

assumes that

all c and cbar interact with the medium as those observed finally as D-mesons

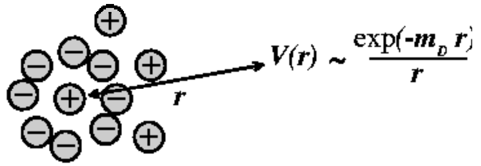
all c and cbar interact among themselves

uses EPOS2 to describe the expanding QGP

HQ interactions with the QGP *Phys.Rev.C* 78 (2008) 014904

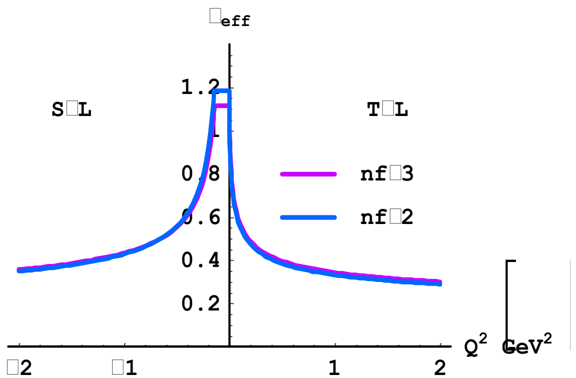
The interaction between HQ and q and g is described by Born type diagrams

$$\frac{d\sigma_F}{dt} = \frac{g^4}{\pi(s-M^2)^2}\left[\frac{(s-M^2)^2}{(t-\kappa m_D^2)^2} + \frac{s}{t-\kappa m_D^2} + \frac{1}{2}\right]$$



q/g is randomly chosen from a Fermi/Bose distribution with the hydro cell temperature

coupling constant and infrared screening are input



Peshier 0801.0595
based on universality
constraint of Dokshitzer

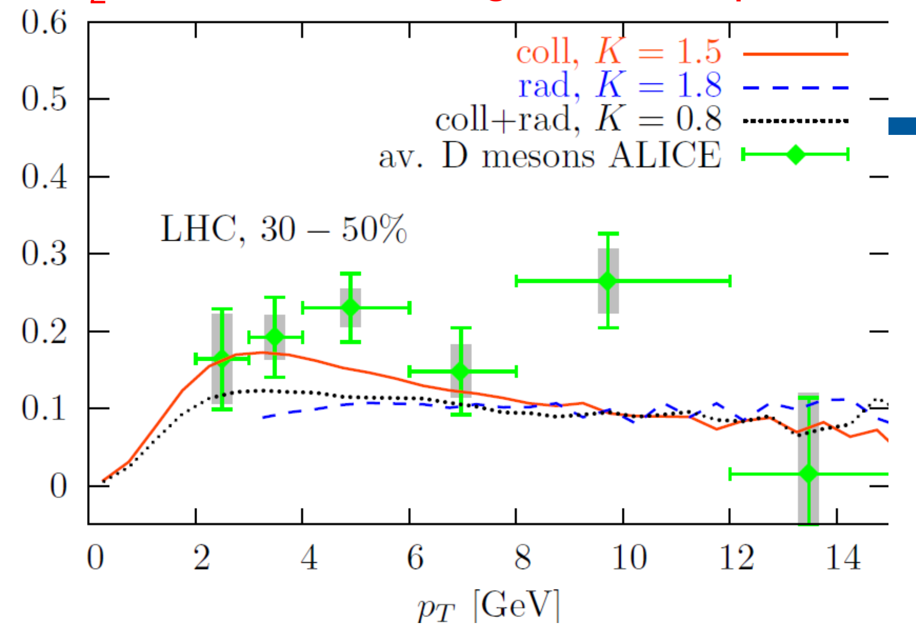
If t is small (<<T) : Born has to be replaced by a hard thermal loop (HTL) approach
For t>T Born approximation is (almost) ok

(Braaten and Thoma PRD441298,2625) for QED:
Energy loss indep. of the artificial scale t* which separates the regimes
Extension to QCD (PRC78:014904)

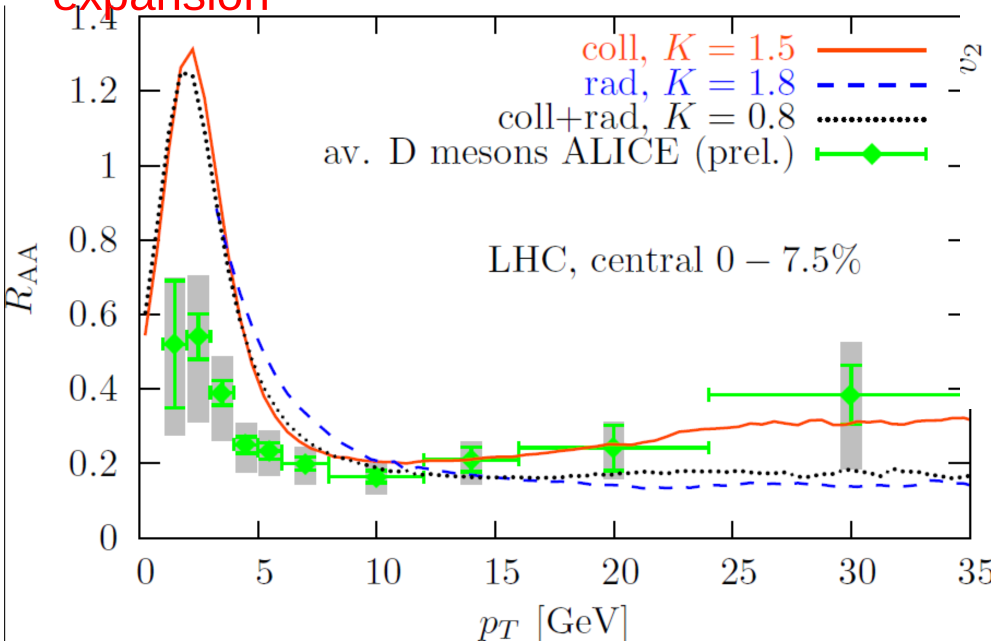
$\kappa \approx 0.2$

D meson results

v_2 tests the late stage of the expansion

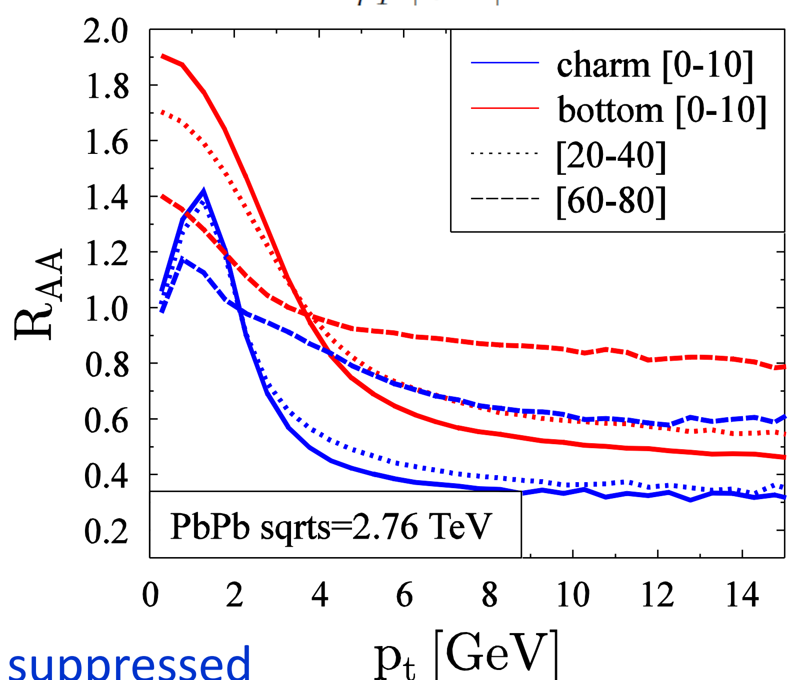


Energy loss tests the initial phase of the expansion



3 options : Collisions only K factor = 1.5
Collision and radiation K = 0.8
Radiation only K= 1.8

R_{AA} and v_2 for coll and coll + radiative give about the same result



b much more suppressed

J/ψ creation in heavy ion collisions

Starting point: von Neumann equation for the density matrix of all particles

$$\partial\rho_N/\partial t = -i[H,\rho_N] \quad \text{with} \quad H = \Sigma_i K_i + \Sigma_{i>j} V_{ij}$$

gives the probability that at time t the state Φ is produced:

$$P^\Phi(t) = \mathrm{Tr}[\rho^\Phi \rho_N(t)] \qquad \rho^\Phi = |\Psi^\Phi >< \Psi_\Phi|$$

This is the solution if we could calculate the quantal $\rho^N(t)$

In our semiclassical approach (correlations are lost) preferable to calculate the rate

$$\Gamma^\Phi(t) = \frac{dP^\Phi}{dt} = \frac{d}{dt}\mathrm{Tr}[\rho^\Phi \rho_N(t)] \qquad P^\Phi(T) = \int_0^T \Gamma^\Phi(t)dt$$

For time independent ρ^Φ

$$\Gamma^\Phi = Tr(\rho^\Phi d\rho^N(t)/dt) = -iTr(\rho^\Phi[H,\rho^N(t)]) = -iTr(\rho^\Phi[U_{12},\rho^N])$$

$$U_{12} = \sum_{j\leq 3}(V_{1j} + V_{2j})$$

J/ψ creation in heavy ion collisions

Heavy ion studies (BUU,QMD,PHSD) have shown that we obtain very satisfying results if we assume

$$W = \langle W^{classic} \rangle$$

We assume in addition that heavy quarks and QGP partons interact by collisions only

$$\frac{dP^{\Phi}(t)}{dt} = \prod_{j}^{N} \int d^3\mathbf{r}_j d^3\mathbf{p}_j \; W^{\Phi} \frac{d}{dt} W_N^c(t).$$

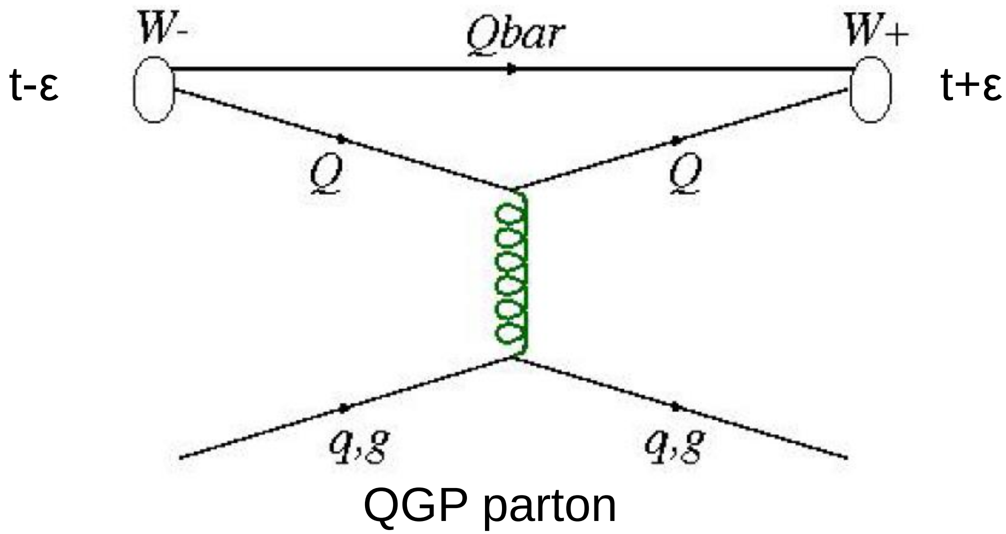
with

$$\begin{aligned} \frac{\partial}{\partial t} W_N^c(t) &= \Sigma_i v_i \cdot \partial_{r_i} W_N^c(\{\mathbf{r}\},\{\mathbf{p}\},t) \qquad (19) \\ &+ \Sigma_{j\geq i} \Sigma_n \delta(t - t_{ij}(n)) \\ &\cdot (W_N^c(\{\mathbf{r}\},\{\mathbf{p}\},t+\epsilon) - W_N^c(\{\mathbf{r}\},\{\mathbf{p}\},t-\epsilon)). \end{aligned}$$

J/ψ creation in heavy ion collisions

If the collisions are point like in time and if $W^{\Phi}(\mathbf{r}_1, \mathbf{r}_2, \mathbf{p}_1, \mathbf{p}_2)$ is time independent

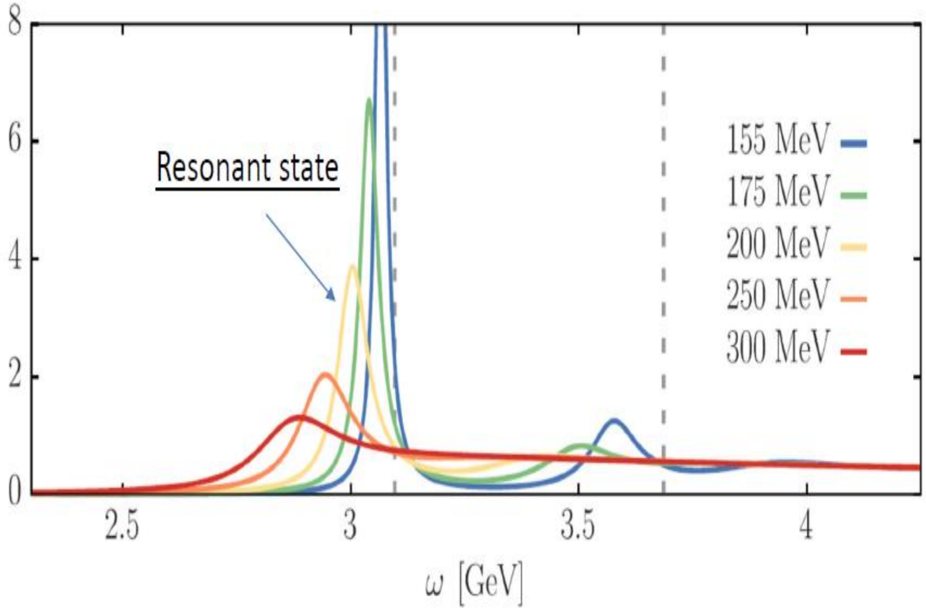
$$\Gamma^{\Phi}(t) = \sum_{i=1,2} \sum_{j\geq 3} \delta(t - t_{ij}(n)) \prod_{k=1}^{N} \int d^3\mathbf{r}_i d^3\mathbf{p}_i$$
$$\cdot \; W^{\Phi}(\mathbf{r}_1, \mathbf{r}_2, \mathbf{p}_1, \mathbf{p}_2)$$
$$\cdot \; [W_N(\{\mathbf{r}, \mathbf{p}\}; t + \epsilon) - W_N(\{\mathbf{r}, \mathbf{p}\}; t - \epsilon)]$$

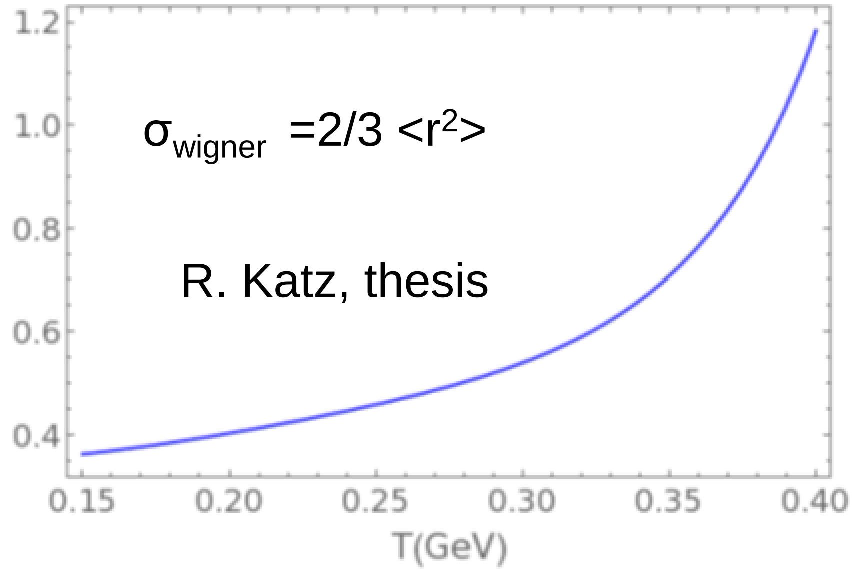


J/ψ creation in heavy ion collisions

Lattice calc: In an expanding QGP $W^{\Phi}(\mathbf{r}_1, \mathbf{r}_2, \mathbf{p}_1, \mathbf{p}_2)$ depends on the temperature and hence on time

Parametrization of the lattice results (Lafferty and Rothkopf PRD 101,056010)





This creates an additional rate, called local rate.

Local Rate

Lattice : J/ψ wavefct is a function of the local QGP temperature

The QGP temperature decreases during the expansion

→ J/ψ wavefct becomes time dependent

creates for $T<T_{diss}$ =400 MeV a local J/ψ prod. rate

$$\Gamma_{loc} = (2\pi\hbar)^3 \int d^3r d^3p \, W_{Q\bar{Q}}(\mathbf{r},\mathbf{p},t)\dot{W}_\Phi(\mathbf{r},\mathbf{p},T(t)).$$

$$= \int d^3r d^3p \quad \frac{16}{(\pi)^3}\dot{\sigma}(T(t))(\frac{\mathbf{r}^2}{\sigma^3(T)} - \frac{\sigma(T)\mathbf{p}^2}{\hbar^2})e^{-(\frac{\mathbf{r}^2}{\sigma^2}+\frac{\sigma^2\mathbf{p}^2}{\hbar})}$$

Total J/ψ multiplicity at time t is then given by

$$P_{Q\bar{Q}}(t) = P^{\text{prim}}(t_{\text{init}}^{Q,\bar{Q}}) + \int_{t_{\text{init}}^{Q,\bar{Q}}}^{t} (\Gamma_{\text{coll},Q\bar{Q}}(t') + \Gamma_{\text{loc},Q\bar{Q}}(t'))dt'$$

For $t \to \infty$ P(t) is the observable J/ψ multiplicity

Interaction of c and cbar in the QGP

V(r) = attractive potential between c and cbar (PRD101,056010)

We work in leading order in γ^{-1}

$$\mathcal{L} = -\gamma^{-1}mc^2 - V(r) \qquad H = \sqrt{m^2 + p_r^2 + \frac{p_\theta^2}{r^2}} + V(r) \qquad p^2 = p_r^2 + p_\theta^2/r^2$$

$$\gamma^{-1} = \sqrt{1 - v^2/c^2} \qquad \frac{\partial \mathcal{L}}{\partial v_i} = p_i = \gamma m v_i$$

Time evolution equation:

$$\dot{r} = \frac{\partial H}{\partial p_r} = \frac{p_r}{\sqrt{m^2 + p_r^2 + \frac{p_\theta^2}{r^2}}}$$

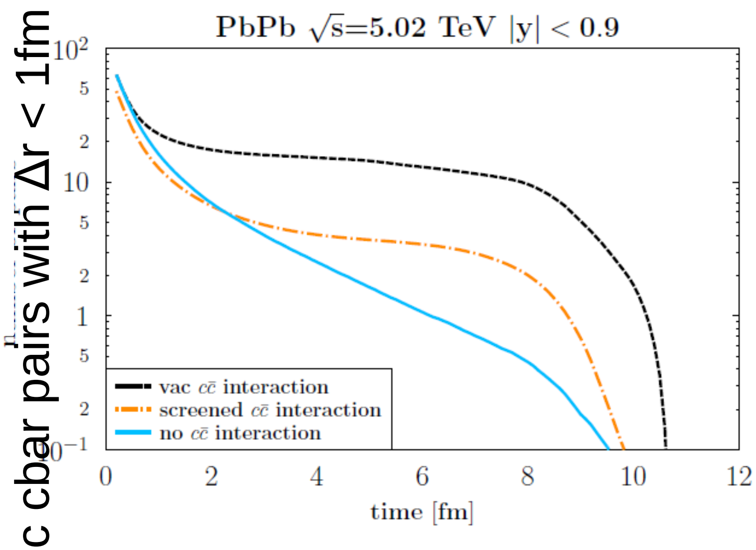
$$\dot{\theta} = \frac{\partial H}{\partial p_\theta} = \frac{p_\theta}{r^2\sqrt{m^2 + p_r^2 + \frac{p_\theta^2}{r^2}}}$$

$$\dot{p}_r = -\frac{\partial H}{\partial r} = \frac{p_\theta^2}{r^3\sqrt{m^2 + p_r^2 + \frac{p_\theta^2}{r^2}}} - \frac{\partial V}{\partial r}$$

$$= \frac{p_\theta \dot{\theta}}{r} - \frac{\partial V}{\partial r}$$

$$\dot{p}_\theta = -\frac{\partial H}{\partial \theta} = 0 \rightarrow p_\theta = \text{const} = L$$

position and momentum of each c cbar pair evolve according to these equations

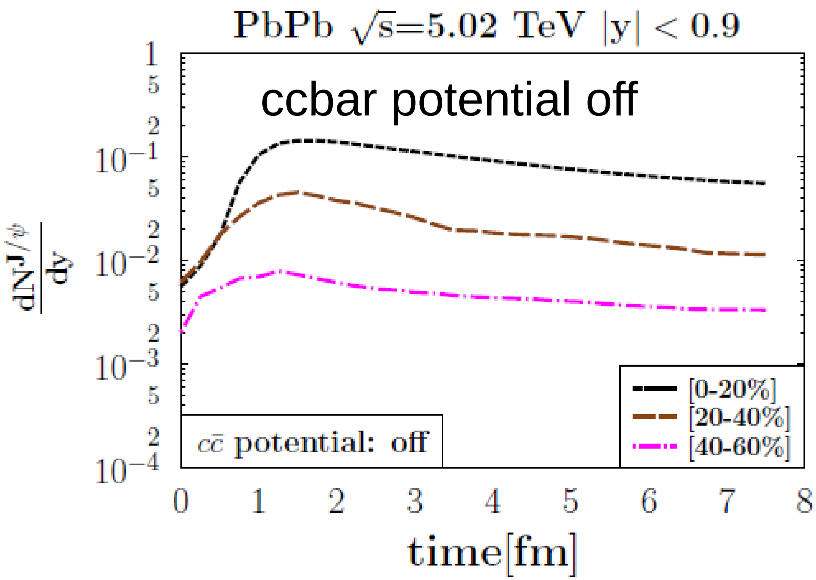


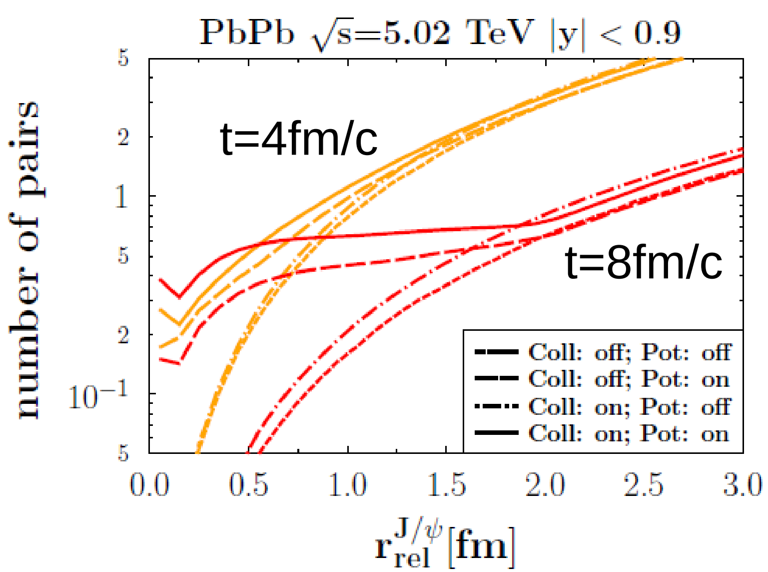
Consequences of the c cbar Interaction

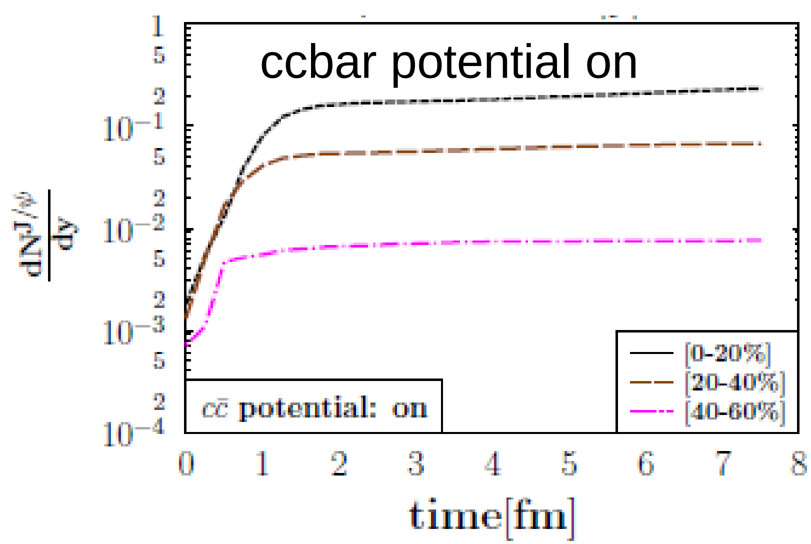
Qq and Qg collisions shift p_T spectra to lower values

(as for D mesons)

QQbar potential interaction increases the production rate



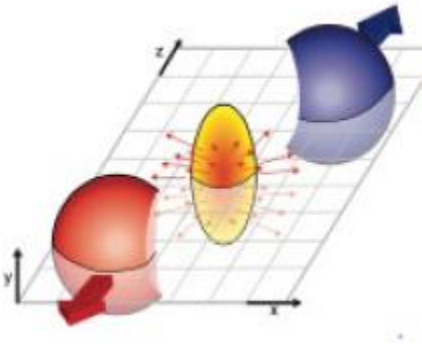




Influence of the Corona

Standard hydrodynamical calculations (EPOS 2) show two classes of particles of initially produced particles:

- Core particles which become part of QGP
- Corona particles from the surface of the interaction zone
  (energy density too low, no collision after production → like pp)
  importent for high pt and for v2



Confirmed by centrality dependence of multiplicity

For elementary particles it is easy to define corona and core particle

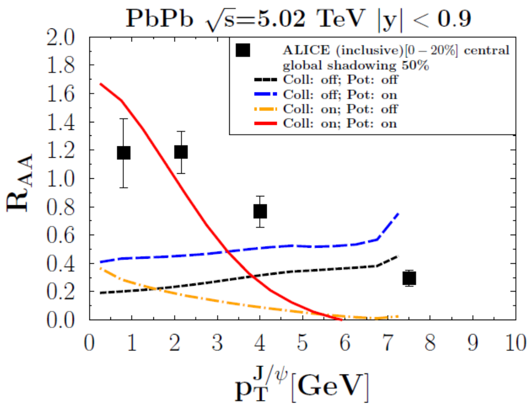
For J/ψ mesons we use working description:

Corona J/ψ are those where none of its constituents suffers from a momentum change of $q > q_{thres}$. Larger q would destroy a J/ψ.

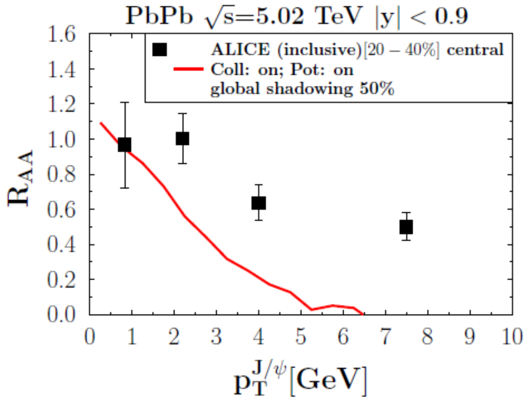
Comparison with ALICE data

Caution: excited state decay, b decay and hadronic rescattering not in yet

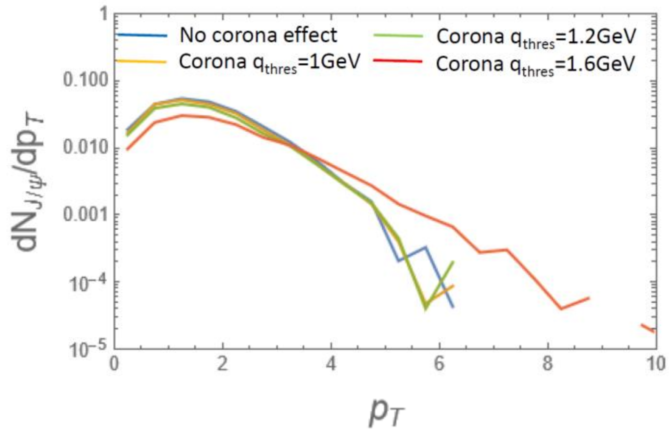
[0-20%]   no corona

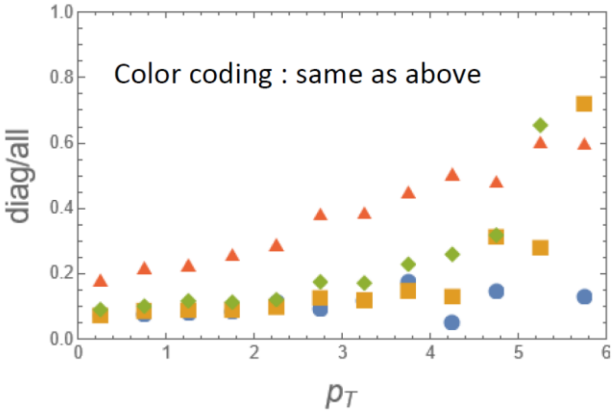


[20-40%] no corona



influence of the corona





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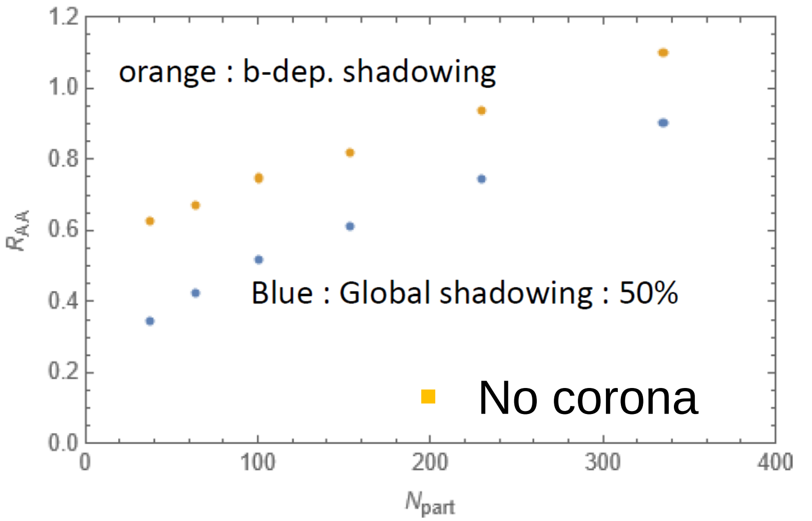
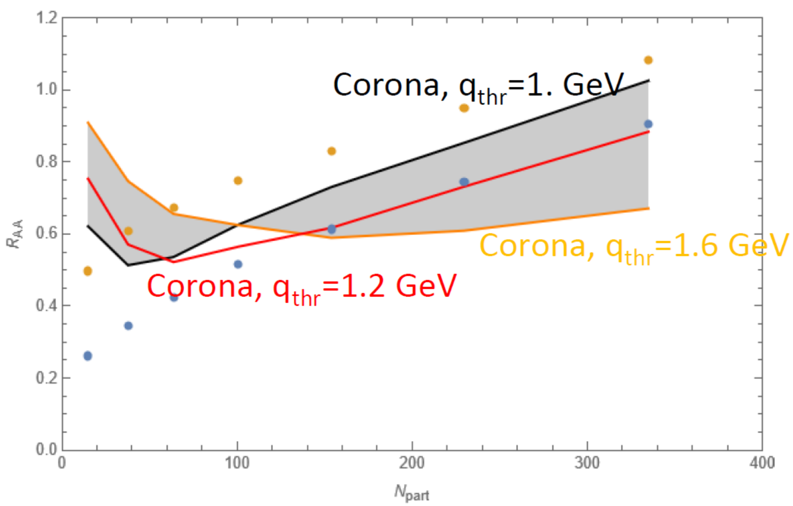
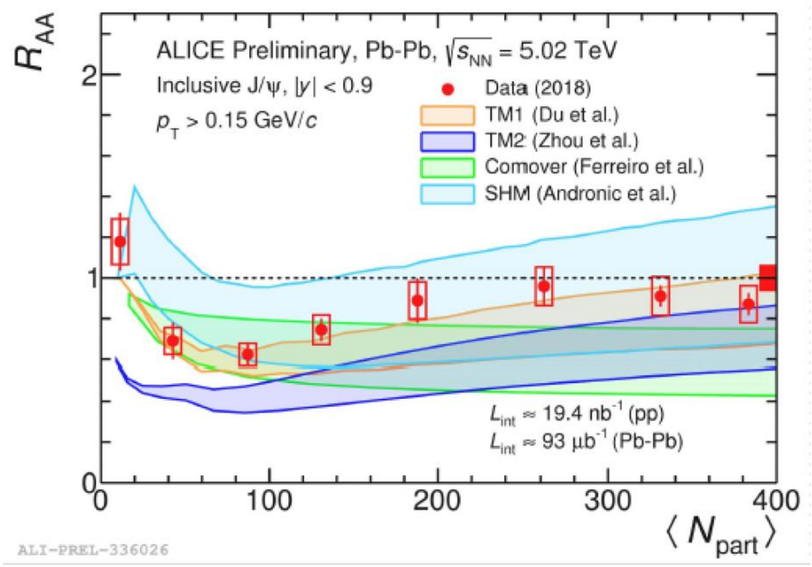
Comparison with ALICE data





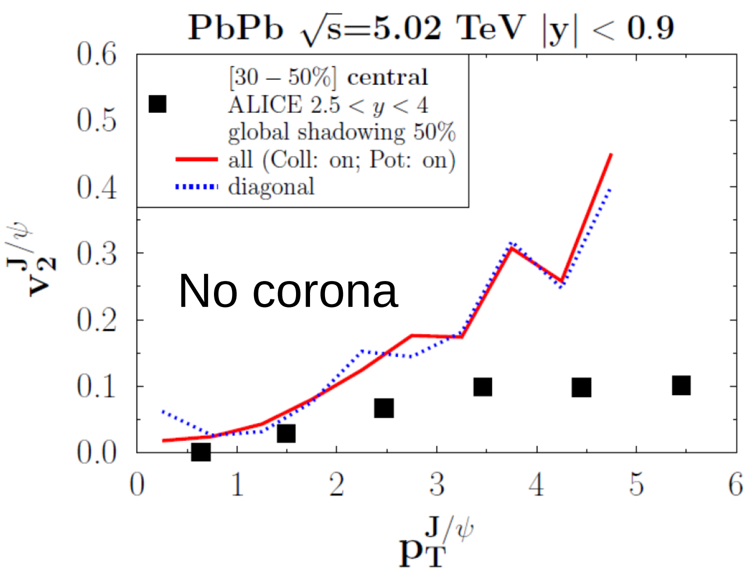
Corona J/ψ bring

- R_{AA} close to one for peripheral reactions
- the participant dependence close to data

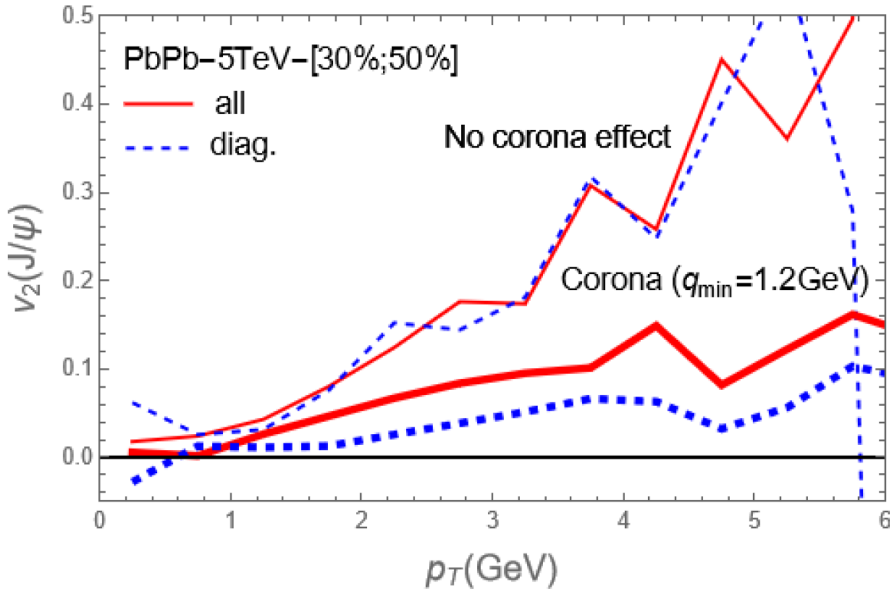


Comparison with ALICE data

[30-50%]



[30-50%]



Corona J/ψ

- bring v_2 closer to the experimental values
- create difference between diagonal and off-diagonal

caution:

comparison of mid and forward rapidities

Summary

New microscopic approach which follows each c and cbar from creation until detection as J/ψ

(no rate equation, no Fokker Planck eq., no thermal assumptions)

- c and cbar are created in initial hard collisions (controlled by pp data)
- when entering the QGP J/ψ become unstable
- c and cbar interact by potential interaction (lattice potential)
  c and cbar interact by collisions with q,g from QGP
- when $T < T_{diss} = 400$ MeV J/ψ can be formed (and later destroyed)
- described by Wigner density formalism (as in pp)

⇩

- Including corona J/ψ, preliminary results agree reasonably with ALICE data for R_{AA} as well as for v_2 .
- The later production (over) compensates the expected multiplicity increase (with respect to pp) due to c and cbar from different vertices
- Has many common features with open quantum system approach (however bottom up)

Outlook; a lot remains to be done.

- Follow the color structure, excited states
- Relativistic kinematics, J/ψ interaction in the hadronic expansion
  Collisions of preformed J/ψ (r < interaction range) with QGP partons (dipole cross section)

Our approach and open quantum systems

Von Neuman eq.

$$\partial\rho_N/\partial t = -\frac{i}{\hbar}[H,\rho_N] \qquad H = H_{1,2} + H_{N-2} + U_{1,2} \qquad U_{1,2} = \Sigma_j V_{1,j} + \Sigma_j V_{2,j}$$

Prob. to find quarkonium $P^\Phi(t) = \mathrm{Tr}[\rho^\Phi \rho_N(t)]$ with $[\rho^\Phi, H_{1,2}] = 0 \quad [\rho^\Phi, H_{N-2}] = 0$

Quarkonium rate: $$\frac{dP^\Phi(t)}{dt} = \Gamma^\Phi(t) = \frac{-i}{\hbar} Tr[\rho^\Phi[U_{1,2},\rho_N(t)]]$$

$$\partial\rho_N(t)/\partial t = -\frac{i}{\hbar}\Sigma_j[K_j,\rho_N(t)] - \frac{i}{\hbar}\Sigma_{k>j}[V_{jk},\rho_N(t)].$$

Interaction: coll. heavy quarks – partons: $$-\frac{i}{\hbar}\Sigma_{k>j}[V_{jk},\rho_N(t)] \equiv \langle \Sigma_{k>j}\Sigma_n \delta(t - t_{jk}(n)) \cdot (W_N^c(\{\mathbf{r}\},\{\mathbf{p}\},t+\epsilon) - W_N^c(\{\mathbf{r}\},\{\mathbf{p}\},t-\epsilon))\rangle.$$

yields

$$\frac{dP^\Phi(t)}{dt} = \Gamma^\Phi(t) = h^3 \frac{d}{dt}\int \prod_j^N d^3r_j d^3p_j W_{12}^\Phi W_N^c(t) = h^3 \int \prod_i^N d^3\mathbf{r}_j d^3\mathbf{p}_j \; W_{12}^\Phi \frac{\partial}{\partial t} W_N^c(t)$$

Lindblad eq. (open quantum systems) in the quantal Brownian motion regime

$$\frac{d}{dt}\rho(t) = -i\left[\frac{p^2}{M} + \Delta H, \rho\right] + \sum_n \int \frac{d^3k}{(2\pi)^3}\left[C_n(\vec{k})\rho C_n^\dagger(\vec{k}) - \frac{1}{2}\left\{C_n^\dagger(\vec{k})C_n(\vec{k}),\rho\right\}\right]$$

Miura, Akamatsu , 2205.15551

new description of c and cbar potential interaction

Not used in the present calculation

Extension to a real relativistic two body kinematics:

Energy and time constraints reduce 8 dim → 6+1 dim phase space

energy constraints

$$\phi_a = \frac{1}{2}(p_{a\mu}p_a^\mu - m_a^2 + \Phi) \approx 0$$

generalized Poisson brackets

$$\{A,B\} = \sum_k \frac{\partial A}{\partial x_k^\mu}\frac{\partial B}{\partial p_{k\mu}} - \frac{\partial B}{\partial x_k^\mu}\frac{\partial A}{\partial p_{k\mu}}$$

which gives the time evolution equations

$$\dot{x}_a^\mu = \{x_a^\mu, \phi_a\} \quad ; \quad \dot{p}_a^\mu = \{p^\mu, \phi_a\}$$

to know what the dot means we need time fixations to the system time

$$\chi_1 = \frac{1}{2}(x_1 - x_2)^\mu U_\mu \quad ; \quad \chi_2 = \frac{1}{2}(x_1 + x_2)^\mu U_\mu - \tau = 0$$

where U is the center of mass velocity

for details: Marty et al. PRC87,034912

new description of c and cbar potential interaction

Fiziev and Todorov (PRD63,104007)

approximation which allows for a separation of CM and relative motion

$$\phi = H = \frac{1}{2\Lambda}(p_{rel}^2 - \mu^2 + \Phi)$$

$$p_{rel}^{cm} = \begin{pmatrix} \frac{s-m_1^2-m_2^2}{2\sqrt{s}} \\ \mathbf{p}_{rel}^{cm} = \nu_2 \mathbf{p}_1^{cm} - \nu_2 \mathbf{p}_1^{cm} \end{pmatrix} \; with \; p_{rel}^{cm} p_{rel}^{cm} = \frac{m_1^2 m_2^2}{s} = \mu_{rel}^2$$

$$\nu_1 - \nu_2 = \frac{m_1^2 - m_2^2}{s}$$

$$\nu_1 + \nu_2 = 1$$

H can be rewritten (for Coulomb)

$$H = \frac{1}{2\lambda}(u_r^2 + \frac{J^2}{r^2} + 1 - (\epsilon^2 + \frac{e^2}{r})^2)$$

J: angular momentum

⍰: const

with the time evolution eqs.

$$\begin{aligned} \dot{r} &= \frac{\partial H}{\partial u_r} = \frac{u_r}{\lambda} \\ \dot{u}_r &= -\frac{\partial H}{\partial r} = \frac{J^2}{\lambda r^3} - \frac{e^2(\epsilon^2 + \frac{e^2}{r})}{\lambda r^2} \\ \dot{\phi} &= \frac{\partial H}{\partial J} = \frac{J}{\lambda r^2} \\ \dot{J} &= -\frac{\partial H}{\partial \phi} = 0 \end{aligned}$$