

Search for dark photons in heavy-ion collisions

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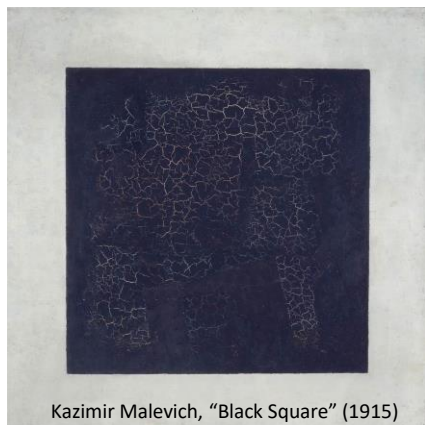
&

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(Havana Uni. & GSI, Darmstadt)



Kazimir Malevich, "Black Square" (1915)

**10th international conference on
High Energy Physics in the LHC Era
(HEP-2023)**

9-13 January 2023, Valparaíso, Chile

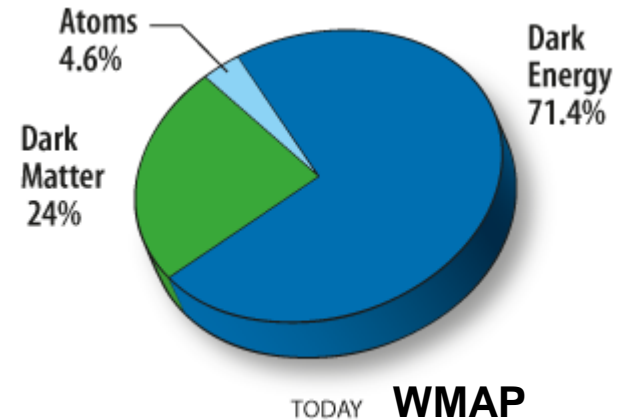
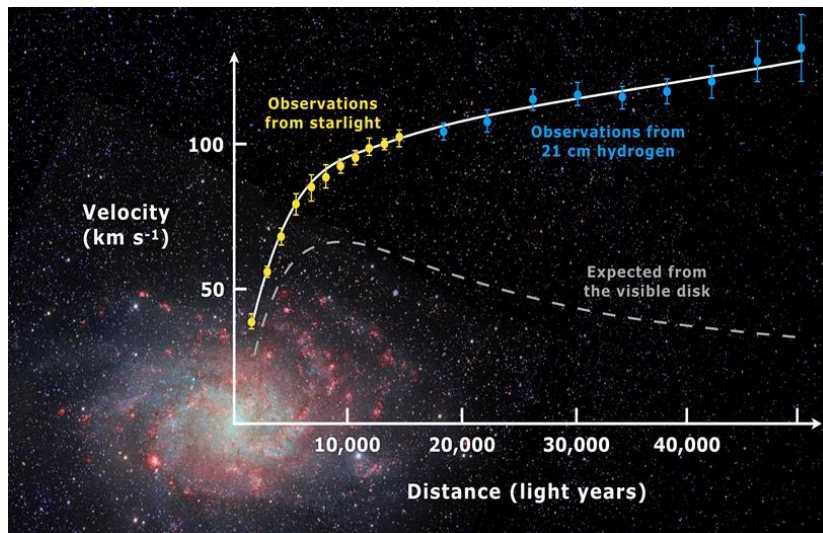


Structure of Universe

- ❑ Dark matter (DM) ~24%
- ❑ DM detected by astrophysical observations based on **gravitational** effects:

1933: F. Zwicky: observation of galaxy clusters

1970: V. Rubin: rotation anomalies in galaxies

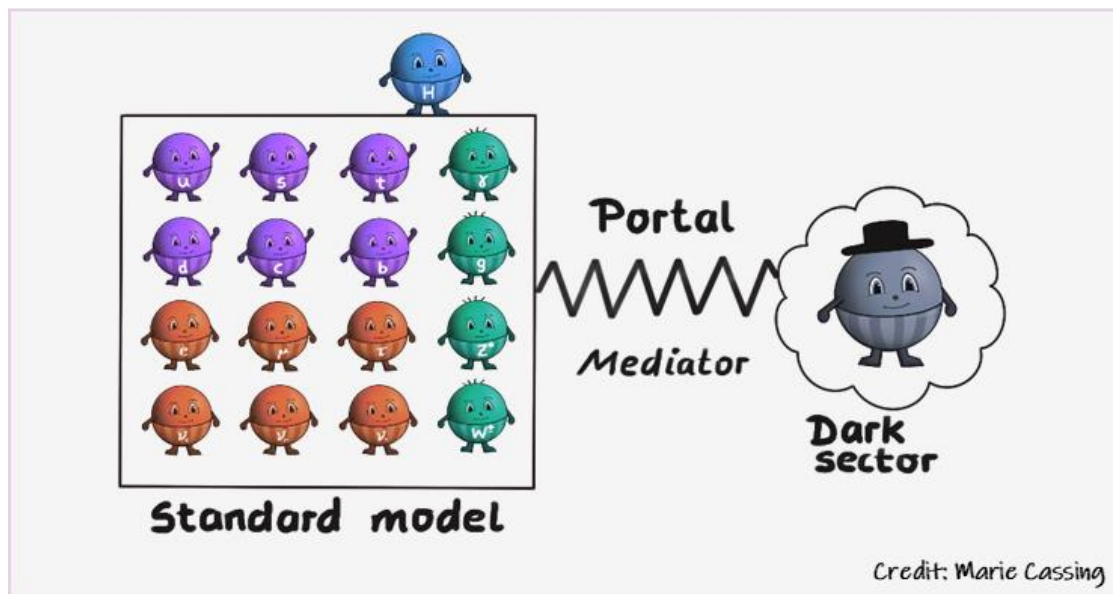


Hubble: gravitational lensing



Dark matter portals

Search for **non-gravitational** dark matter (DM) **interactions with normal matter**, i.e. with standard model (SM) particles



$$\mathcal{L} \supset \begin{cases} -\frac{\epsilon}{2 \cos \theta_W} B_{\mu\nu} F'^{\mu\nu}, & \text{vector portal} \\ (\mu\phi + \lambda\phi^2) H^\dagger H, & \text{Higgs portal} \\ y_n L H N, & \text{neutrino portal} \\ \frac{a}{f_a} F_{\mu\nu} \tilde{F}^{\mu\nu}, & \text{axion portal.} \end{cases} \quad \leftarrow$$

Vector portal

The '**vector portal**' assumes the mixing of SM and DM via a **U(1)-U(1)'** gauge symmetry group mixing

L.B. Okun, Sov. Phys. 56 JETP (1982);
B. Holdom, Phys. Lett. B 166, 196 (1986)

$$\mathcal{L}_{A'} = -\frac{1}{4}F'^{\mu\nu}F'_{\mu\nu} + \boxed{\frac{1}{2} \frac{\epsilon}{\cos \theta_W} B^{\mu\nu} F'_{\mu\nu}} - \frac{1}{2}m_{A'}^2 A'^{\mu} A'_{\mu}$$

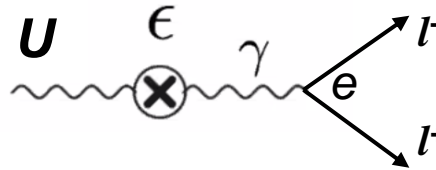
Dark photon field strength:

$$F'_{\mu\nu} \equiv \partial_{\mu}A'_{\nu} - \partial_{\nu}A'_{\mu}$$

SM hypercharge field strength:

$$B_{\mu\nu} \equiv \partial_{\mu}B_{\nu} - \partial_{\nu}B_{\mu}$$

ϵ - kinetic mixing parameter:



$$\epsilon^2 = \alpha'/\alpha$$

Due to the kinetic mixing the dark photon (U-boson) couples to the electromagnetic current with strength **ϵe**

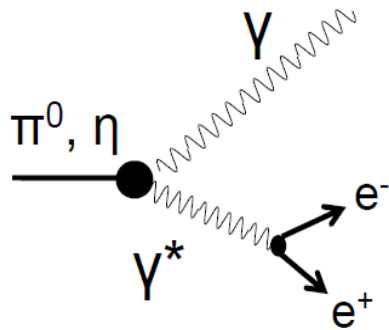
Unknown: kinetic mixing parameter **ϵ** and mass **M_U**

* Notation in literature for the 'dark photon': A' , V , U -boson

Dalitz decay of the dark photon to dileptons

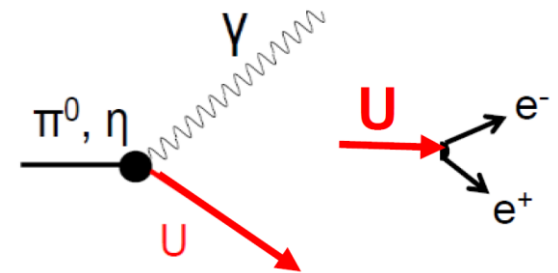
- Dalitz decays of pseudoscalar mesons π^0, η and Δ -resonances to dileptons via the U-boson mediator

Standard model

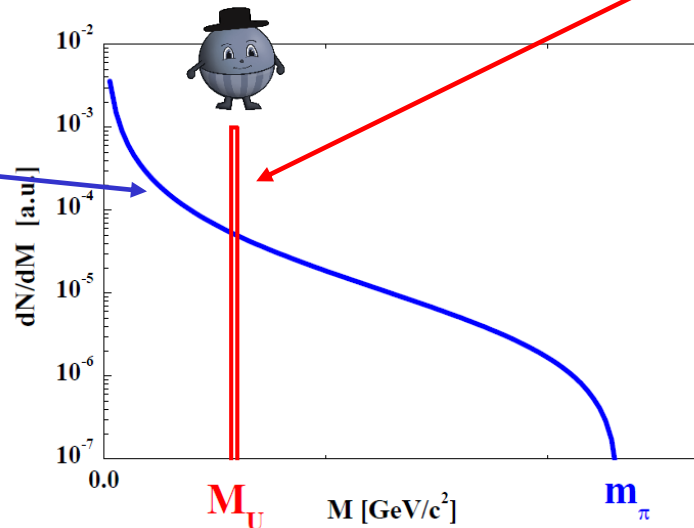


U(1)-U(1)'
kinetic mixing

Beyond SM: DM scenario

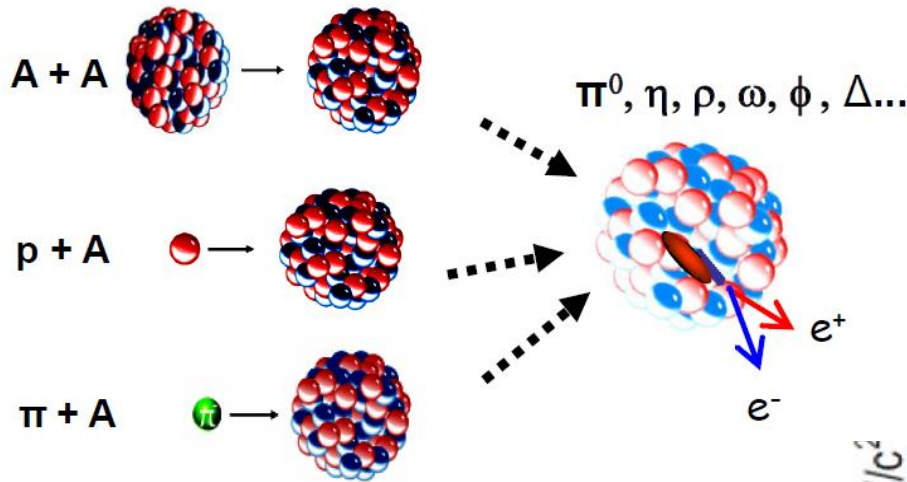


$\pi^0 \rightarrow \gamma + \gamma^*$,
 $\eta \rightarrow \gamma + \gamma^*$, $\gamma^* \rightarrow e^+e^-$
 $\Delta \rightarrow N + \gamma^*$



$\pi^0 \rightarrow \gamma + U$,
 $\eta \rightarrow \gamma + U$, $U \rightarrow e^+e^-$
 $\Delta \rightarrow N + U$

Possible dark photon observation by dilepton experiments



□ Dilepton spectra from SM sources are well studied by dilepton experiments from SIS to LHC energies

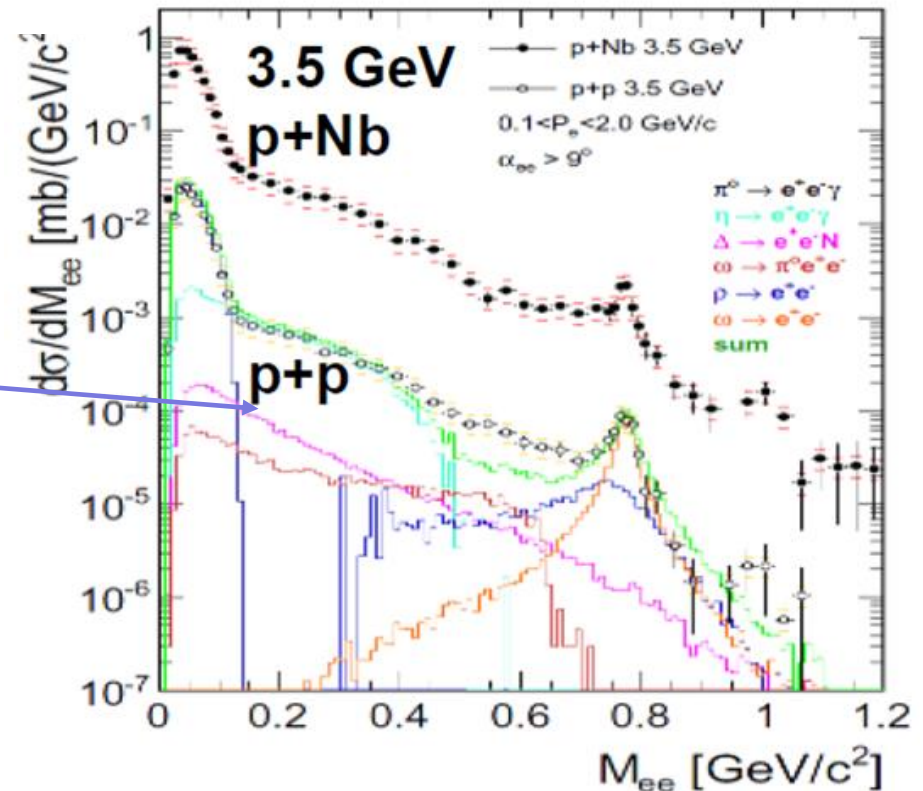
□ Hadron production by p+p, p+A, A+A

□ Hadron decay to dileptons

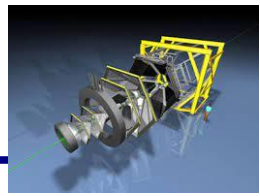
□ Dalitz π^0, η and Δ decays are the dominant dilepton sources at low M

→ Possibility for an experimental observation of dark photons by electromagnetic decays $U \rightarrow e^+e^-$ in heavy-ion experiments

Dilepton spectra at low M ('cocktail')



Search for dark photons with HADES (GSI)

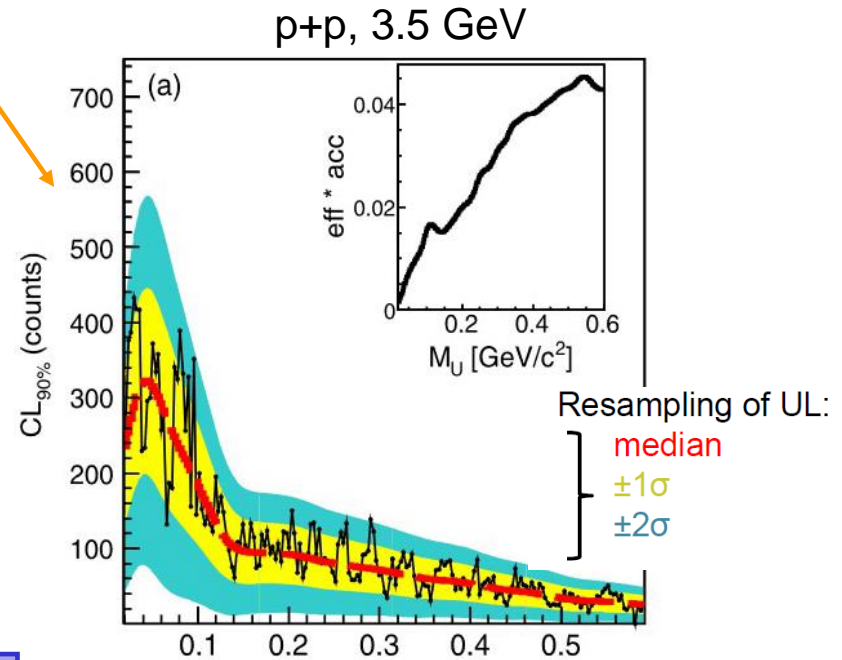
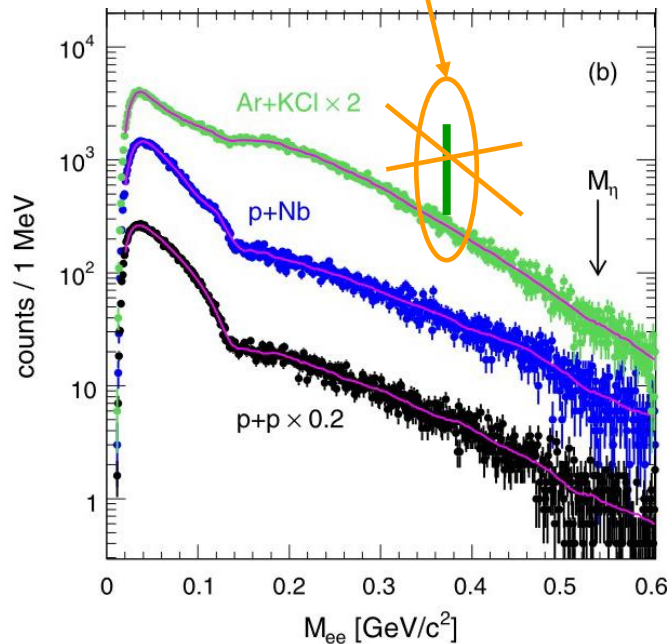


HADES data:

p+p, p+Nb at 3.5 GeV, Ar+KCl at 1.76 A GeV

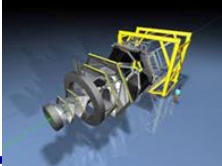
G. Agakishiev et al. (HADES), Phys. Lett. B 731, 265 (2014)

- 1) **Search for a peak structure** in the raw dN/dM spectrum taking into account mass resolution: fit with 5th-order polynomial + Gauss peak for each fixed M bin
- 2) If no peak found, get an **UL (upper limit)** on peak



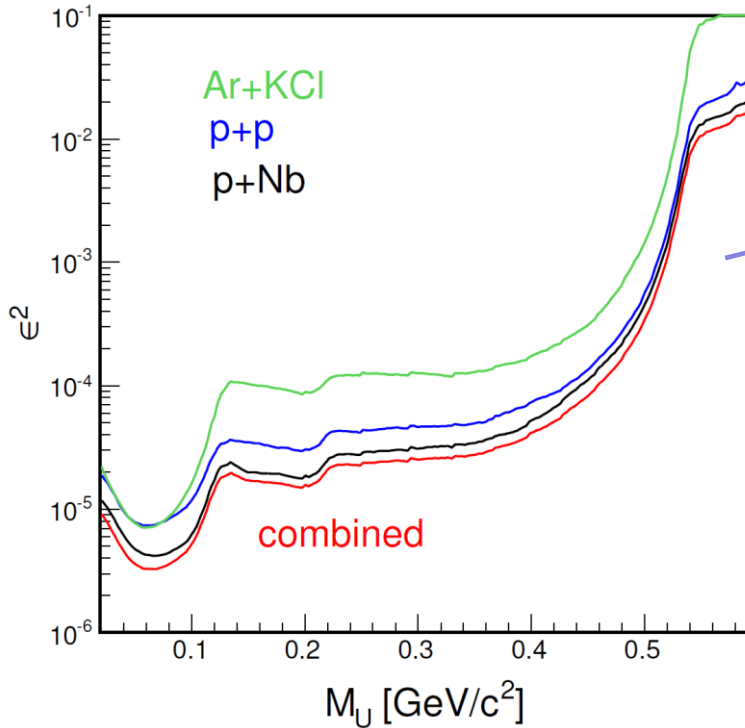
- 3) Transform this UL into an UL on the **mixing parameter ϵ^2** based on the **modelling of the U-boson production rate** (B. Batell, M. Pospelov, and A. Ritz, PRD 80,095024 (2009))

Search for dark photons with HADES



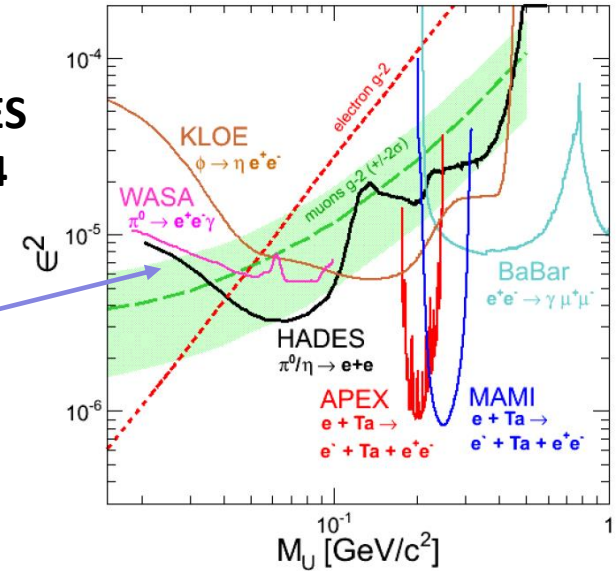
G. Agakishiev et al. (HADES), Phys. Lett. B 731, 265 (2014)

Upper limit on the mixing parameter ϵ^2



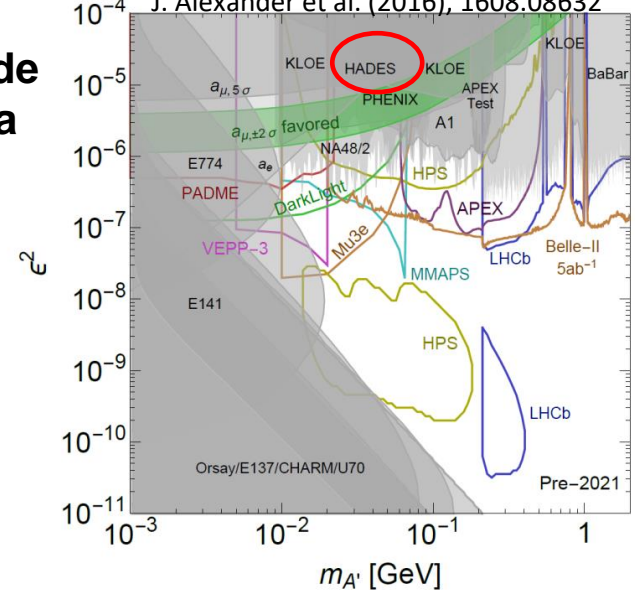
HADES
2014

HADES: Phys. Lett. B 731 (2014) 265



World wide
exp. data
2016

J. Alexander et al. (2016), 1608.08632



HADES coverage : $0.02 < M_U < 0.6$ GeV

Dark photons are not observed so far!

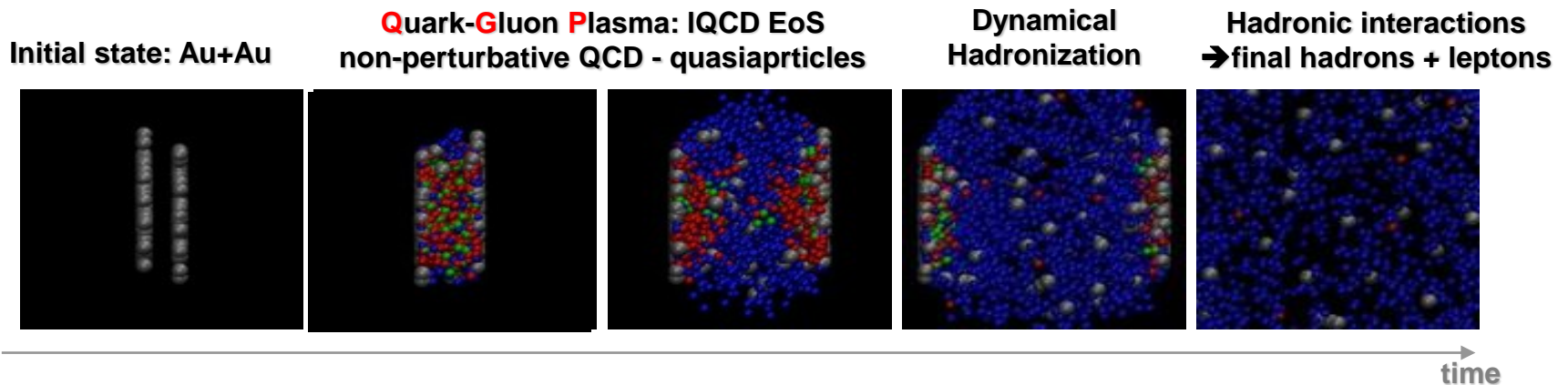


Theoretical modeling of U-boson production

Goal: estimate the upper limit for the kinetic mixing parameter $\varepsilon^2(M_U)$ of the U-boson **from the theoretically calculated dilepton spectra** using the microscopic **PHSD** transport approach

Parton-Hadron-String Dynamics (PHSD) is a **non-equilibrium microscopic transport approach** for the description of strongly-interacting hadronic and partonic matter created in heavy-ion collisions

Dynamics: based on the solution of generalized off-shell transport equations derived from Kadanoff-Baym many-body theory



→ **PHSD** provides a good description of ‘bulk’ hadronic observables as well as **dilepton spectra** from SIS to LHC energies



Light dark photon production in PHSD

Production of hadron \rightarrow decay to U \rightarrow dilepton yield from U-boson decay of mass M_U :

$$\begin{aligned} \pi^0 &\rightarrow \gamma + U, \\ \eta &\rightarrow \gamma + U, \quad U \rightarrow e^+e^- \\ \Delta &\rightarrow N + U, \end{aligned}$$

$$\begin{aligned} N^{U \rightarrow e^+e^-} &= N_{\pi^0 \rightarrow e^+e^-} + N_{\eta \rightarrow e^+e^-} + N_{\Delta \rightarrow e^+e^-} \\ &= Br^{U \rightarrow e^+e^-} (N_{\pi^0 \rightarrow \gamma U} + N_{\eta \rightarrow \gamma U} + N_{\Delta \rightarrow NU}), \end{aligned}$$

I.) $(N_{\pi^0 \rightarrow \gamma U} + N_{\eta \rightarrow \gamma U} + N_{\Delta \rightarrow NU})$, $N_{i \rightarrow \gamma U} = N_i Br_{i \rightarrow \gamma \gamma} \cdot \frac{\Gamma_{i \rightarrow \gamma U}}{\Gamma_{i \rightarrow \gamma \gamma}}$, $i = \pi^0, \eta$ $\leftarrow \underline{\varepsilon^2} = \alpha' / \alpha$

The yield of U-bosons of mass M_U : $N_{\Delta \rightarrow NU} = N_{\Delta} Br_{\Delta \rightarrow N \gamma} \cdot \frac{\Gamma_{\Delta \rightarrow NU}}{\Gamma_{\Delta \rightarrow N \gamma}}$.

Dalitz decay of π^0, η mesons and Δ -resonances to U-bosons and real photons or N

Based on the model: B. Batell, M. Pospelov, and A. Ritz, Phys. Rev. D 79, 115008 (2009);
 (used in the HADES analysis) Phys. Rev. D 80, 095024 (2009)

- Ratio of the partial widths $\pi^0(\eta) \rightarrow \gamma + U$ and $\pi^0(\eta) \rightarrow \gamma + \gamma$:

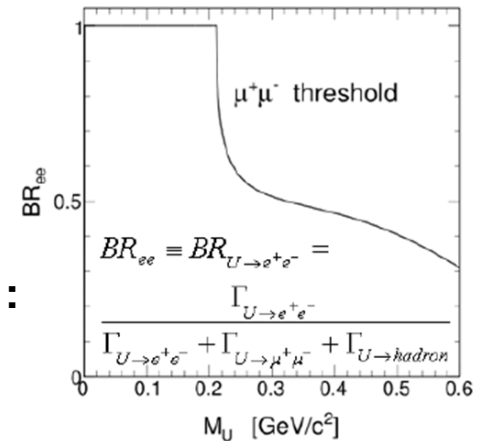
$$\frac{\Gamma_{i \rightarrow \gamma U}}{\Gamma_{i \rightarrow \gamma \gamma}} = 2\varepsilon^2 |F_i(q^2 = M_U^2)| \frac{\lambda^{3/2}(m_i^2, m_\gamma^2, M_U^2)}{\lambda^{3/2}(m_i^2, m_\gamma^2, m_\gamma^2)} \quad i = \pi^0, \eta$$

Formfactor

- II.) $Br^{U \rightarrow e^+e^-}$ Branching ratio for the decay of U-bosons to e^+e^- :

$$Br^{U \rightarrow ee} = \frac{\Gamma_{U \rightarrow e^+e^-}}{\Gamma_{tot}^U} = \frac{1}{1 + \sqrt{1 - \frac{4m_\mu^2}{M_U^2}} \left(1 + \frac{2m_\mu^2}{M_U}\right) (1 + R(M_U))}$$

$$R(\sqrt{s}) = \sigma_{e^+e^- \rightarrow hadrons} / \sigma_{e^+e^- \rightarrow \mu^+\mu^-}$$





Procedure to obtain constraints on $\epsilon^2(M_U)$

- 1) For each bin $[M_U, M_U + dM]$ calculate the **sum of all $U \rightarrow e^+e^-$ contributions** (kinematically possible in this mass bin)

$$\frac{dN^{sumU}}{dM} = \frac{dN_{\pi^0}^{U \rightarrow e^+e^-}}{dM} + \frac{dN_{\eta}^{U \rightarrow e^+e^-}}{dM} + \frac{dN_{\Delta}^{U \rightarrow e^+e^-}}{dM}$$

- 2) Calculate the **sum of all SM contributions and 'dark matter' (DM) contributions** :

$$\frac{dN^{total}}{dM} = \frac{dN^{sumSM}}{dM} + \frac{dN^{sumU}}{dM} = \frac{dN^{sumSM}}{dM} + \epsilon^2 \frac{dN_{\epsilon=1}^{sumU}}{dM}$$

- 3) Obtain **constraints** by requesting that **dN^{total}/dM (SM+DM) cannot exceed the sum of SM channels (i.e. exp. data!) by more than a factor C_U** in each bin dM , i.e.

$$\frac{dN^{total}}{dM} = (1 + C_U) \frac{dN^{sumSM}}{dM}$$

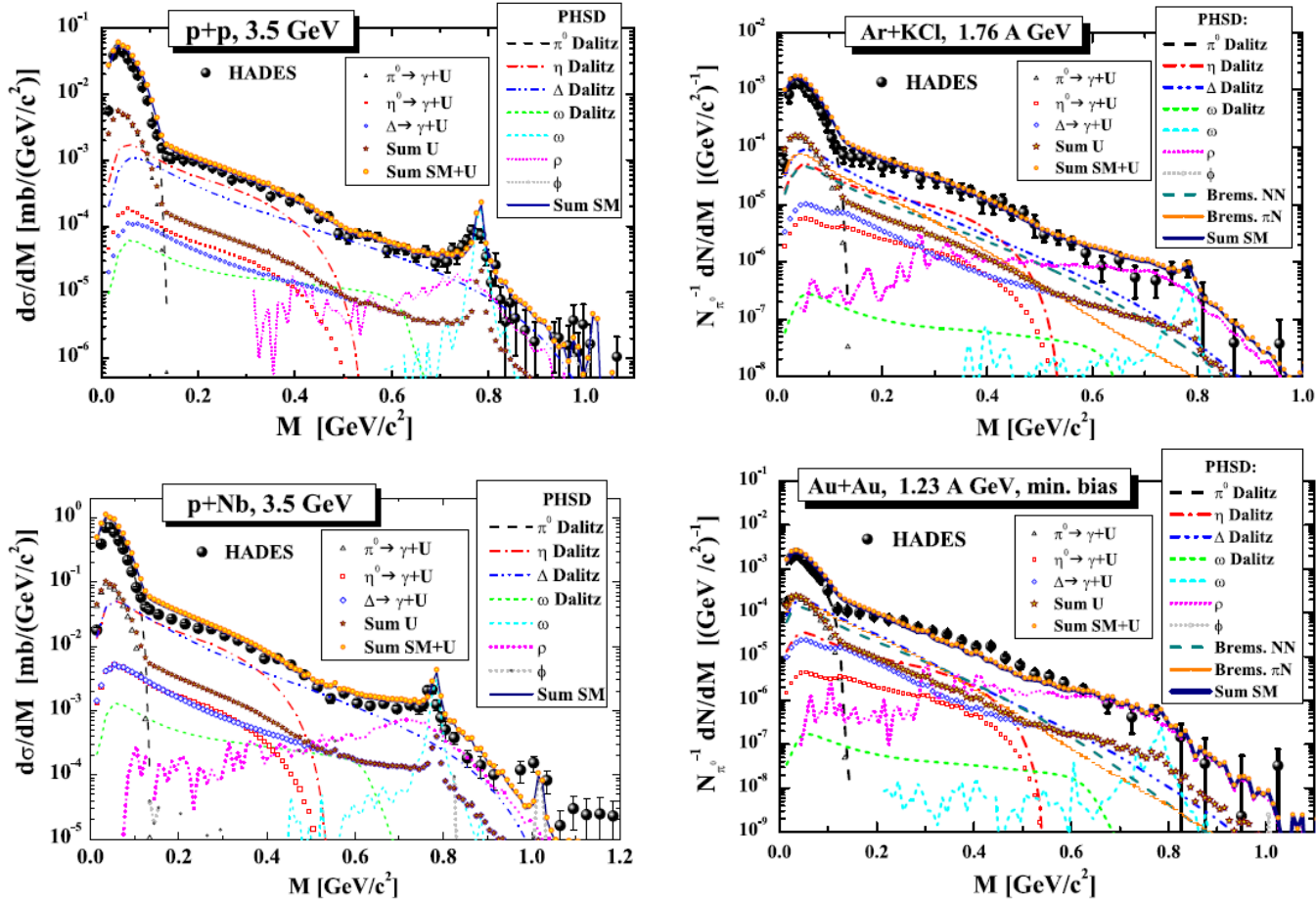


C_U controls the allowed **"surplus"** dilepton yield resulting from dark photons on top of the total SM yield

- 4) Calculate **$\epsilon^2(M_U)$** by assuming C_U : e.g. $C_U=0.1 \rightarrow 10\%$ DM extra yield to the SM yield

$$\epsilon^2(M_U) = C_U \cdot \left(\frac{dN^{sumSM}}{dM} \right) / \left(\frac{dN_{\epsilon=1}^{sumU}}{dM} \right)$$

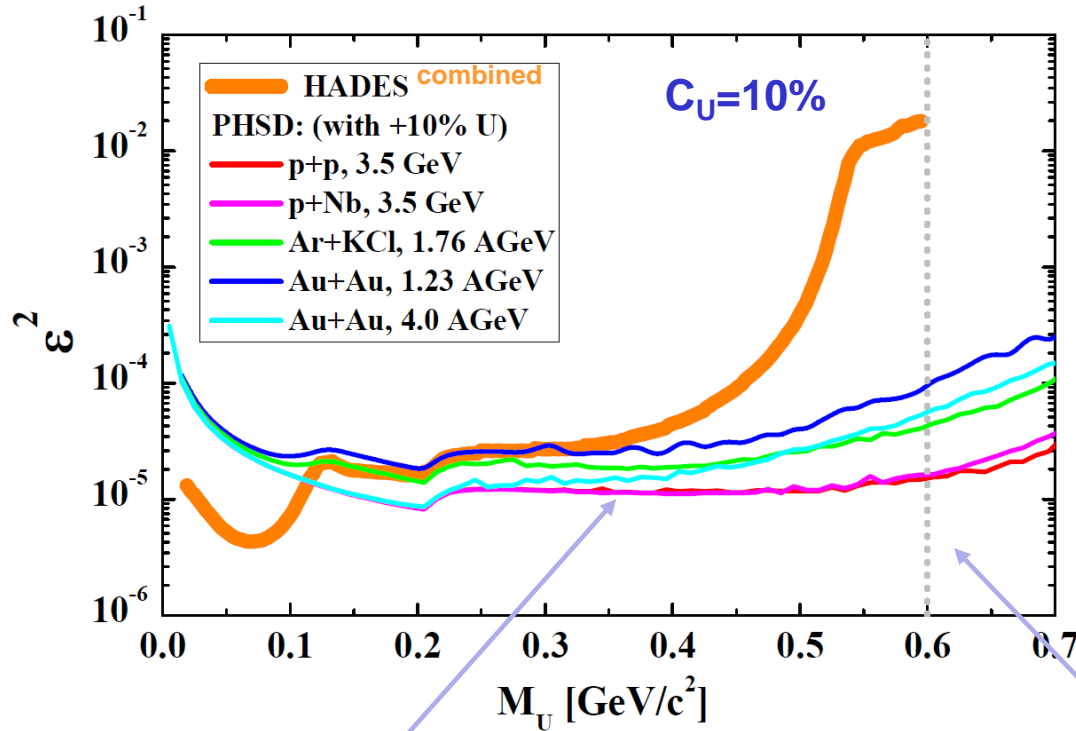
Dileptons yields including dark photons vs. HADES data



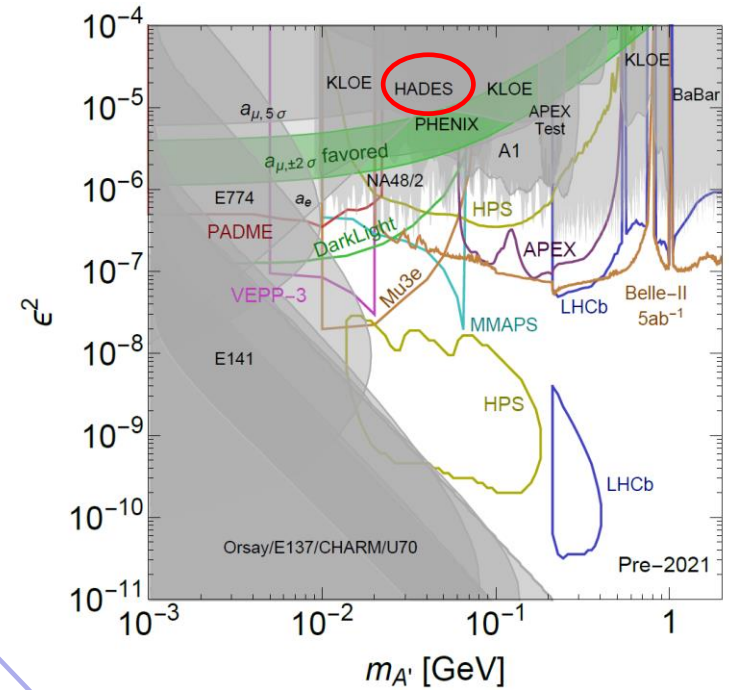
- The HADES data, i.e. **SM contributions** (including exp. acceptance) are well described by the PHSD
- The contributions from **U** → e+e- are added with **C_U=10%** allowed surplus of the total SM yield → the **total sum** is still in a good agreement with exp. data

Mixing parameter $\varepsilon^2(M_U)$

The **upper limit for the kinetic mixing parameter $\varepsilon^2(M_U)$** of light dark photons extracted from the PHSD dilepton spectra - with **10% allowed surplus** of the total SM yield



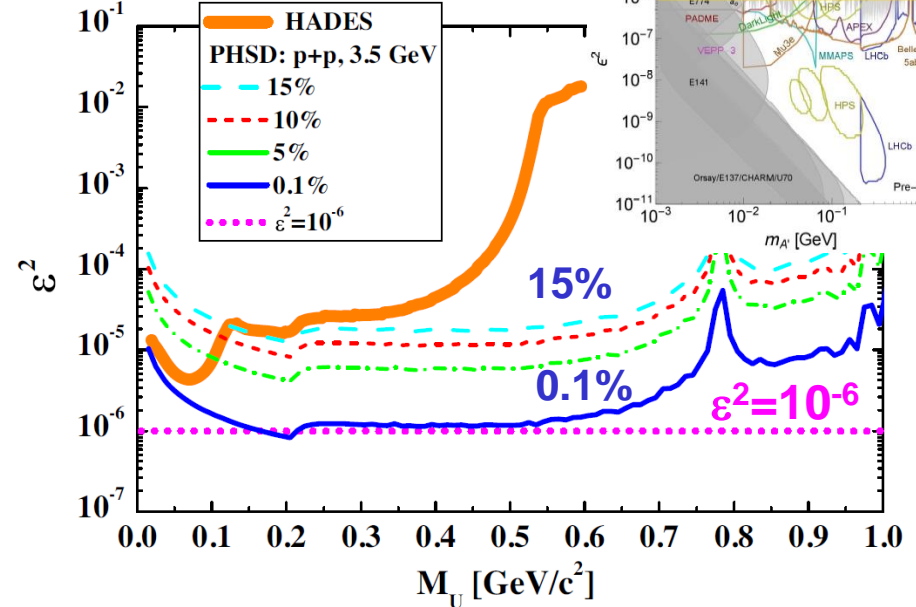
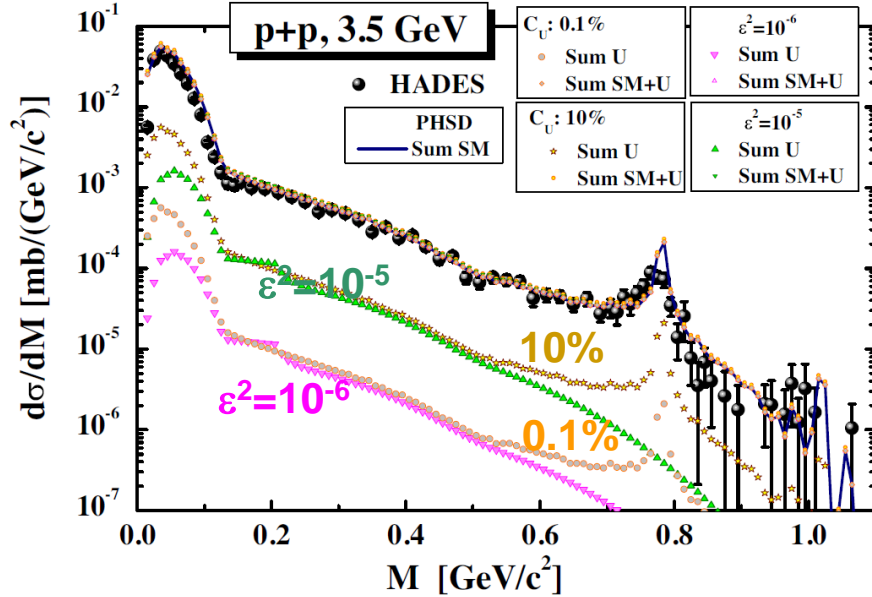
Validity range of extracted kinetic mixing parameter: **$0 < M_U < 0.6 \text{ GeV}$** based on low energy dilepton spectra



Possible contribution from other dark photon channels

Limits for the mixing parameter $\varepsilon^2(M_U)$

- The PHSD predictions for $\varepsilon^2(M_U)$ with 0.1%, 5%, 10%, and 15% allowed surplus of the U-boson contributions over the total SM yield



The **theoretically** extracted upper limit of the kinetic mixing parameter $\varepsilon^2(M_U)$ of light dark photons from Dalitz decays of π^0, η mesons and Δ -resonances:

- strongly reduces by lowering the allowed 'surplus'
- exp. data of high precision is needed to reduce the upper limit for $\varepsilon^2(M_U)$

Summary

- ❑ We presented **first microscopic transport calculations**, based on the PHSD approach, for the **dilepton yield from the decay of hypothetical dark photons** (or U-bosons), $U \rightarrow e^+e^-$ from $p + p$, $p + A$ and heavy-ion collisions at SIS energies
- ❑ For that we incorporated in the PHSD the **production of U-bosons** by the Dalitz decay $\pi^0 \rightarrow \gamma + U$, $\eta \rightarrow \gamma + U$, $\Delta \rightarrow N + U$ with further dilepton decays $U \rightarrow e^+e^-$ based on the theoretical model by Batell, Pospelov and Ritz, which describes the interaction of DM and SM particles by the **$U(1)-U(1)'$ mixing**
- ❑ We **introduced a procedure to define theoretical constraints on the upper limit of the kinetic mixing parameter $\varepsilon^2(M_U)$** :
 Since dark photons are not observed in dilepton experiments so far, we can require that their contribution **can not exceed some limit** which would make them visible in experimental data
- ❑ We found that the **extracted upper limit of $\varepsilon^2(M_U)$ is consistent with** the experimental results of the **HADES experiment** for $0.15 < M_U < 0.4$ GeV, as well as with the world-wide experimental compilation
- ➔ **Proposed theoretical procedure allows:**
 - to check any theoretical ideas on the $\varepsilon^2(M_U)$ independent on exp. data
 - to study the influence of exp. acceptance, system and centrality selection
 - to perform the simulation for testing experimental set-ups for the search of U-bosons