

# Perturbative and non-perturbative renormalization for quasi-PDFs

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in collaboration with:

*C. Alexandrou, K. Cichy, K. Hadjyiannakou, K. Jansen*

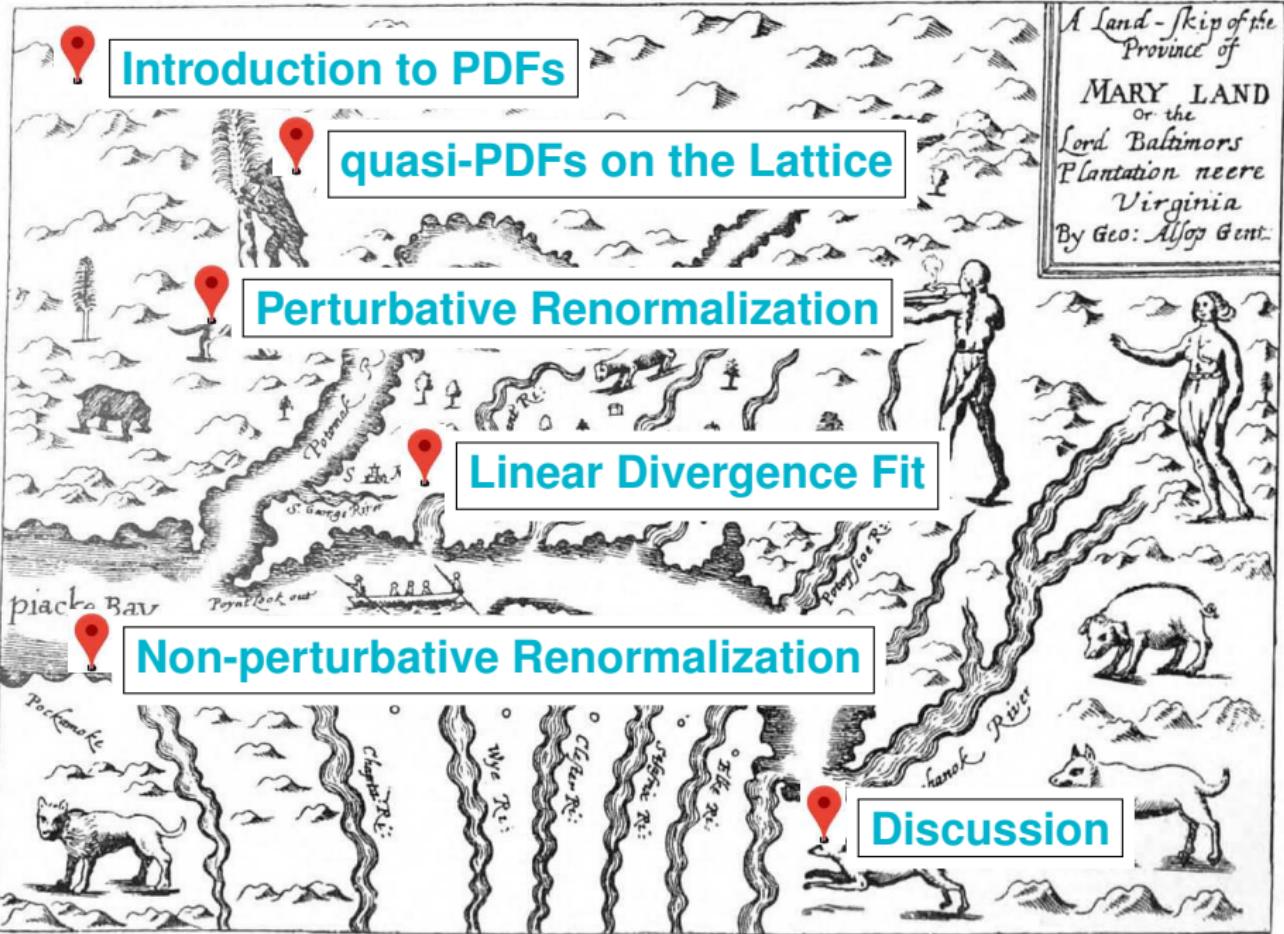
*H. Panagopoulos, A. Scapellato, F. Steffens*

*Lattice PDF Workshop*

*Maryland, USA*

*April 6, 2018*

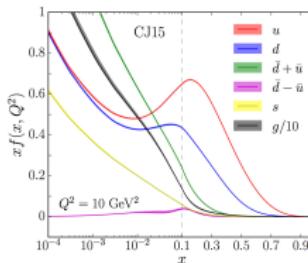
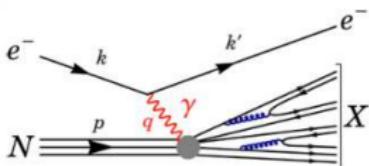
# ROADMAP OF TALK



A

# Introduction to PDFs

# Probing Nucleon Structure



CJ15 PDFs

## Parton Distribution Functions

- ★ powerful tool to describe the structure of a nucleon
- ★ necessary for the analysis of Deep inelastic scattering (DIS) data
- ★ Parametrization of off-forward matrix of a bilocal quark operator (light-like)

$$F_\Gamma(x, \xi, q^2) = \frac{1}{2} \int \frac{d\lambda}{2\pi} e^{ix\lambda} \langle p' | \bar{\psi}(-\lambda n/2) \textcolor{red}{O} \underbrace{\mathcal{P} e^{-\lambda/2}}_{\text{gauge invariance}} \psi(\lambda n/2) | p \rangle$$

$q = p' - p$ ,  $\bar{P} = (p' + p)/2$ ,  $n$ : light-cone vector ( $\bar{P} \cdot n = 1$ ),  $\xi = -n \cdot \Delta/2$

# PDFs on the Lattice

- ★ first principle calculations of PDFs are necessary
- ★ On the lattice: long history of moments of PDFs

$$f^n = \int_{-1}^1 dx x^n f(x)$$

- ★ rely on OPE to reconstruct the PDFs (**difficult task**):
  - signal-to-noise is bad for higher moments
  - $n > 3$ : operator mixing (unavoidable!)
- ★ Alternative investigation of PDFs ?

Types:

- Unpolarized      (vector current)
- Polarized        (axial current)
- Transversity     (tensor current)

# PDFs on the Lattice

Novel direct approach: [X.Ji, arXiv:1305.1539]

- ★ compute **quasi-PDF** on the lattice
- ★ contact with physical PDFs in two steps:
  1. Renormalization of quasi-PDFs in Lattice Regularization
  2. Matching procedure (LaMET)

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Exploratory studies are maturing:

- [X. Xiong et al., arXiv:1310.7471], [H-W. Lin et al., arXiv:1402.1462], [Y. Ma et al., arXiv:1404.6860],
- [Y.-Q. Ma et al., arXiv:1412.2688], [C. Alexandrou et al., arXiv:1504.07455], [H.-N. Li et al., arXiv:1602.07575],
- [J.-W. Chen et al., arXiv:1603.06664], [J.-W. Chen et al., arXiv:1609.08102], [T. Ishikawa et al., arXiv:1609.02018],
- [C. Alexandrou et al., arXiv:1610.03689], [C. Monahan et al., arXiv:1612.01584], [A. Radyushkin et al., arXiv:1702.01726],
- [C. Carlson et al., arXiv:1702.05775], [R. Briceno et al., arXiv:1703.06072], [M. Constantinou et al., arXiv:1705.11193],
- [C. Alexandrou et al., arXiv:1706.00265], [J-W Chen et al., arXiv:1706.01295], [X. Ji et al., arXiv:1706.08962],
- [K. Orginos et al., arXiv:1706.05373], [T. Ishikawa et al., arXiv:1707.03107], [J. Green et al., arXiv:1707.07152],
- [Y-Q Ma et al., arXiv:1709.03018], [J. Karpie et al., arXiv:1710.08288], [J-W Chen et al., arXiv:1711.07858],
- [C.Alexandrou et al., arXiv:1710.06408], [T. Izubuchi et al., arXiv:1801.03917], [C.Alexandrou et al., arXiv:1803.02685],
- [J-W Chen et al., arXiv:1803.04393]. . . .

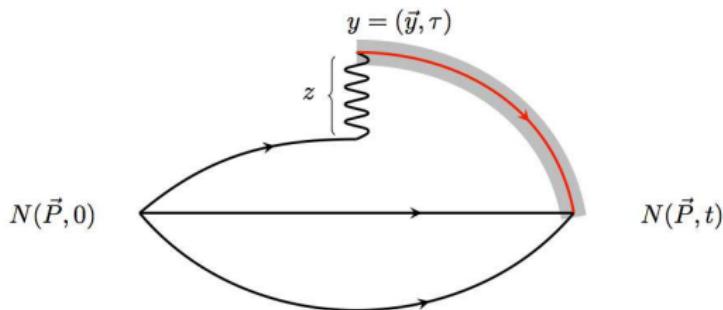
# B quasi-PDFs on the Lattice

# Access of PDFs on a Euclidean Lattice

- ★ quasi-PDF purely spatial for nucleons with finite momentum

$$\tilde{q}(x, \mu^2, P_3) = \int \frac{dz}{4\pi} e^{-i x P_3 z} \langle N(P_3) | \bar{\Psi}(z) \gamma^z \mathcal{A}(z, 0) \Psi(0) | N(P_3) \rangle_{\mu^2}$$

- $\mathcal{A}(z, 0)$ : Wilson line from  $0 \rightarrow z$
- $z$ : distance in any spatial direction (momentum boost in  $z$  direction)



- ★ At finite but feasibly large momenta on the lattice:  
a large momentum EFT can relate Euclidean  $\tilde{q}$  to PDFs through a factorization theorem
- ★ use of Perturbation Theory for the matching

C

History of

Renormalization

# Renormalization History

Prior 2017:

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[M. Constantinou, H. Panagopoulos, Phys. Rev. D 96 \(2017\) 054506, \[arXiv:1705.11193\]](#)
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## 2017:

- ★ Results available for renormalized matrix elements (ME)
- ★ Comparison with phenomenology becomes a real possibility

C

# Perturbative Renormalization

# Perturbative Calculation

## ★ Operators

$$\mathcal{O}_\Gamma^\mu \equiv \bar{\psi}(x) \Gamma \mathcal{P} e^{i g \int_0^z A(\zeta) d\zeta} \psi(x + z \hat{\mu})$$

Scalar, Pseudoscalar, Vector, Axial, Tensor

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### A Dimensional Regularization (DR)

- Compute conversion factor between  $\overline{MS}$  and other schemes

### B Lattice Regularization (LR)

- Extract full pert. Z-factors using Green's functions  
in both DR and LR

## A. Dimensional Regularization



### Features of DR Calculation

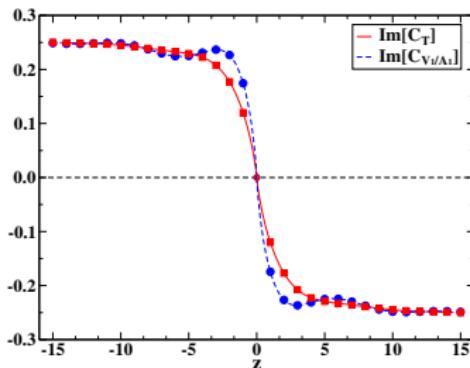
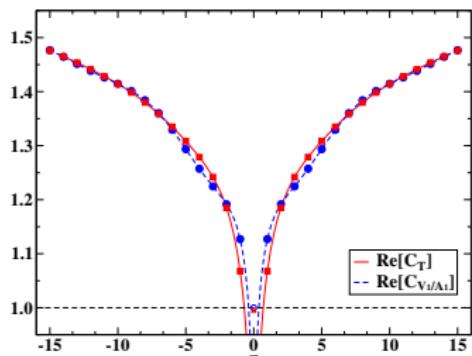
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- ★ Z-factors in  $\overline{\text{MS}}$ : real function
- ★ Conversion factor: a complex function

## A. Dimensional Regularization



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Parameters chosen based on B55 ETMC ensemble, Nucleon momentum:  $P_3=4$

- ★ Necessary ingredient for non-perturbative renormalization

## B. Lattice Regularization



- ★ Linear divergence from tadpole diagram:  $\propto |z|/a$
- ★ To all orders in pert. theory:  $e^{-c \frac{|z|}{a}}$  [Dotsenko et al., NPB169 (1980) 527]
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## B. Lattice Regularization



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- ★ **Extraction of Z-factor**

$$\langle \psi \mathcal{O}_\Gamma \bar{\psi} \rangle_{\text{amp}}^{DR, \overline{\text{MS}}} - \langle \psi \mathcal{O}_\Gamma \bar{\psi} \rangle_{\text{amp}}^{LR} = \frac{g^2 C_f}{16 \pi^2} e^{i q_\mu z} \times \mathcal{F}$$

$$\mathcal{F} = \Gamma \left( c_1 + c_2 \beta + c_3 \frac{|z|}{a} + \log(a^2 \bar{\mu}^2) (4-\beta) \right) + (\Gamma \cdot \gamma_\mu + \gamma_\mu \cdot \Gamma) \left( c_4 + c_5 c_{\text{sw}} \right)$$

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linear divergence

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linear divergence      mixing term

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linear divergence

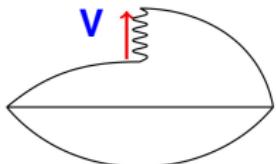
$c_i$  depend on gluonic action

mixing term

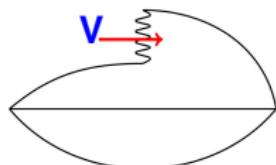
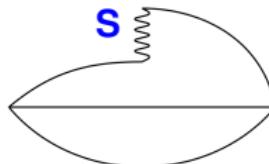
mixing vanishes if  $c_{\text{SW}} = -c_4/c_5$

# Mixing pattern

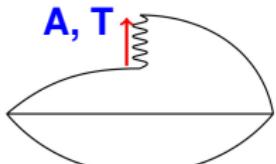
Depends on the relation between the current & Wilson line direction



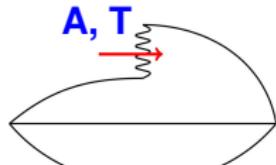
mixing with



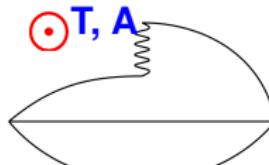
no mixing



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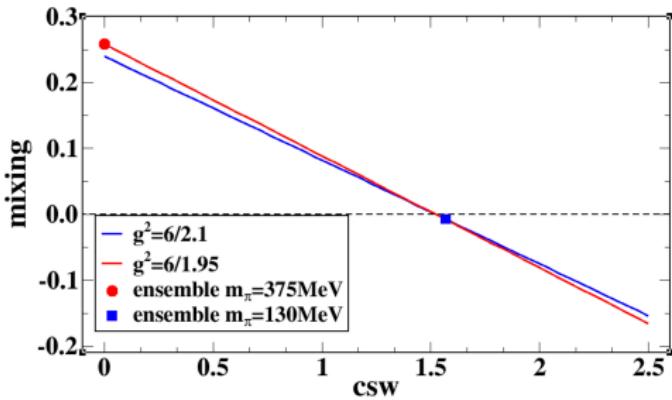
mixing with



: Wilson line direction  
↑ : Current insertion direction

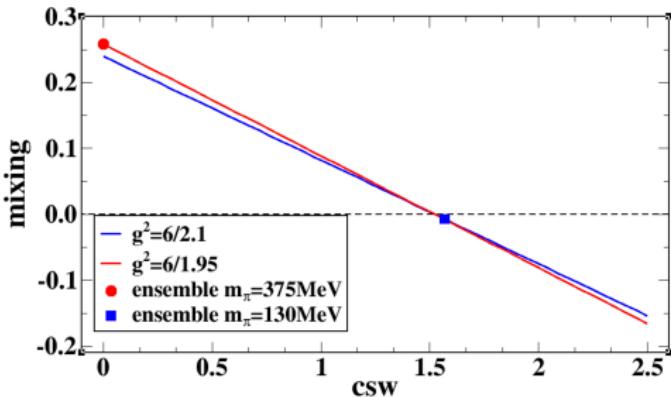
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- ⇒ If mixing present, it must be taken into account (e.g.  $\gamma_3$ )



- ★ Mixing may be suppressed because:
  - mixing vanishes at some value of  $c_{SW}$  ( $=-c_4/c_5$ )
  - numerical application show no mixing for  $c_{SW} \sim 1$  (for Wilson-like actions)

D

Linear

Divergence

# Linear Divergence

Absence of mixing:

$$\mathcal{R} = \frac{q(P_3, z)}{q(P'_3, z')} = e^{\left(-\frac{c}{a} + c_0\right)(|z| - |z'|)} \left(\frac{P_3}{P'_3}\right)^{-6\frac{g^2 C_f}{16\pi^2}}, \quad z P_3 = z' P'_3$$

presence of  $c_0$ :  
[\[R. Sommer, arXiv\[1501.03060\]\]](#)

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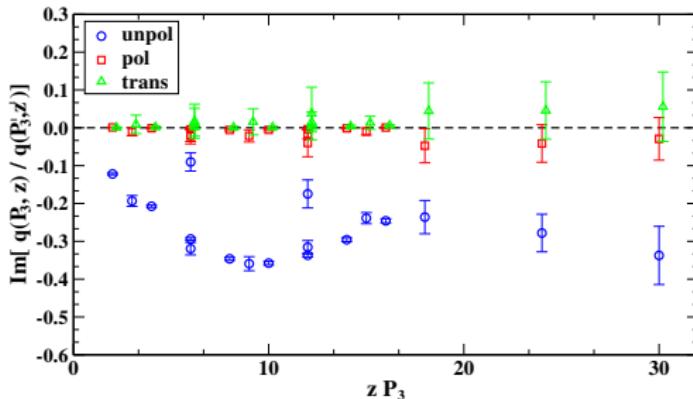
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open symbols:

Twisted Mass

$m_\pi = 375 \text{ MeV}$



- ★ Mixing must be treated for the unpolarized case

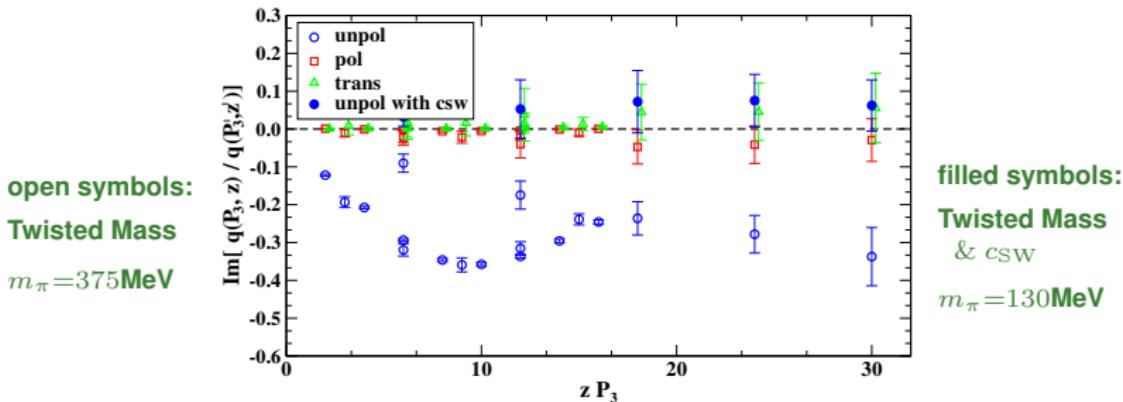
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- ★ Mixing must be treated for the unpolarized case
- ★ Presence of  $c_{\text{sw}}$  suppresses mixing

E

# Non-Perturbative Renormalization

# Non-perturbative Renormalization

- ★ RI' scheme
- ★ Use 1-loop conversion factor to convert to the  $\overline{\text{MS}}$  at 2 GeV
- ★ Vertex function has the same divergence as the nucleon matrix element

## No mixing

helicity, transversity,  
unpolarized ( $\gamma_0$ )

$$Z_{\mathcal{O}}(z) = \frac{Z_q}{\mathcal{V}_{\mathcal{O}}(z)}$$

$$\mathcal{V}_{\mathcal{O}} = \frac{\text{Tr}}{12} \left[ \mathcal{V}(p) \left( \mathcal{V}^{\text{Born}}(p) \right)^{-1} \right] \Big|_{p=\bar{\mu}}$$

- ★  $Z_q$ : fermion field renormalization
- ★  $Z_{\mathcal{O}}$  includes the linear divergence

## Mixing

Unpolarized ( $\gamma_3$ )

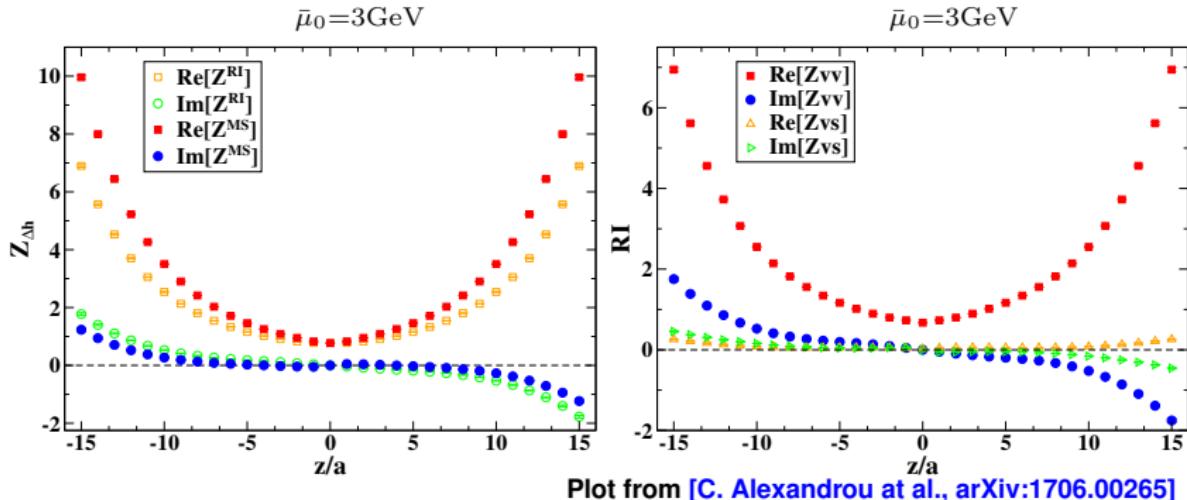
$$\begin{pmatrix} \mathcal{O}_V^R(P_3, z) \\ \mathcal{O}_S^R(P_3, z) \end{pmatrix} = \hat{Z}(z) \cdot \begin{pmatrix} \mathcal{O}_V(P_3, z) \\ \mathcal{O}_S(P_3, z) \end{pmatrix}$$

$$Z_q^{-1} \hat{Z}(z) \hat{\mathcal{V}}(p, z) \Big|_{p=\bar{\mu}} = \hat{1}$$

$$\begin{aligned} h_V^R(P_3, z) &= Z_{VV}(z) h_V(P_3, z) \\ &+ Z_{VS}(z) h_S(P_3, z) \end{aligned}$$

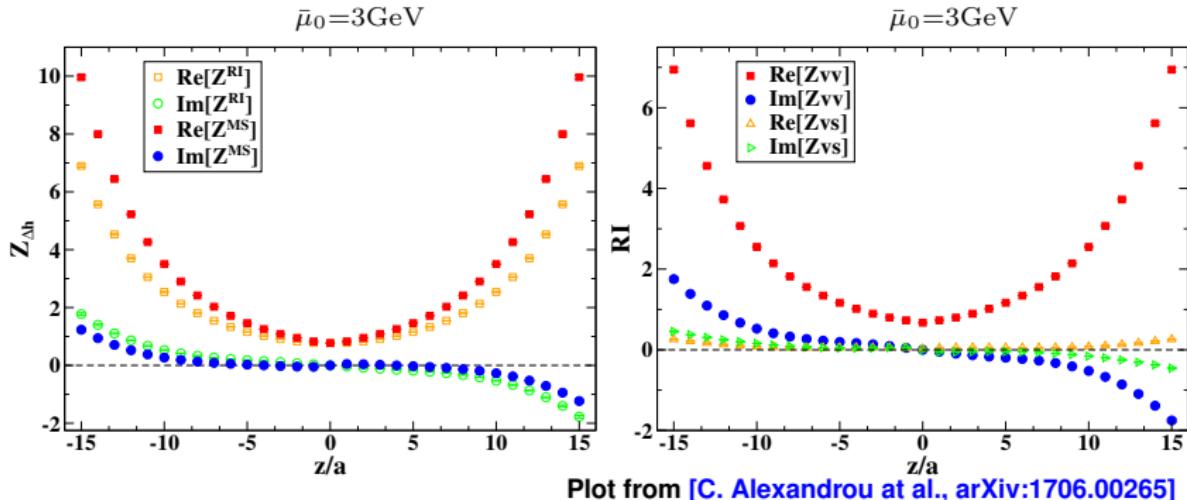
# Numerical Results

- ★ Twisted Mass fermions,  $m_\pi=375\text{MeV}$ ,  $32^3 \times 64$ , HYP smearing
- ★ Conversion & Evolution to  $\overline{\text{MS}}(2\text{GeV})$  (Perturbatively)



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- ★ Z-factors are complex functions
- ★  $\text{Im}[Z_{\mathcal{O}}^{\overline{\text{MS}}}] < \text{Im}[Z_{\mathcal{O}}^{\text{RI}'}]$  (expected from pert. theory)

# Systematic uncertainties

**Ultimate goal: Reliability in final estimates**

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Systematic uncertainties need to be addressed

- ★ Upper bounds estimated in [C. Alexandrou et al., arXiv:1706.00265]
- ★ Both the ME and Z-factors are complex functions, in absence of mixing, e.g. unpolarized with  $\gamma_0$  ( $h \equiv h_{u-d}$ ):

$$\begin{aligned} h^{ren} = Z_h h &= Re[Z_h] Re[h] - Im[Z_h] Im[h] \\ &\quad + I (Re[Z_h] Im[h] + Im[Z_h] Re[h]) \end{aligned}$$

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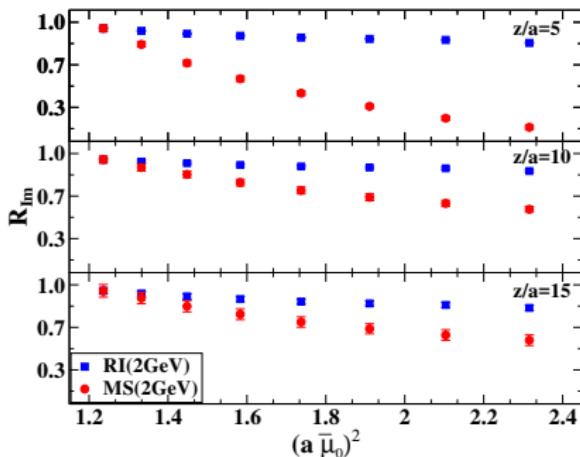
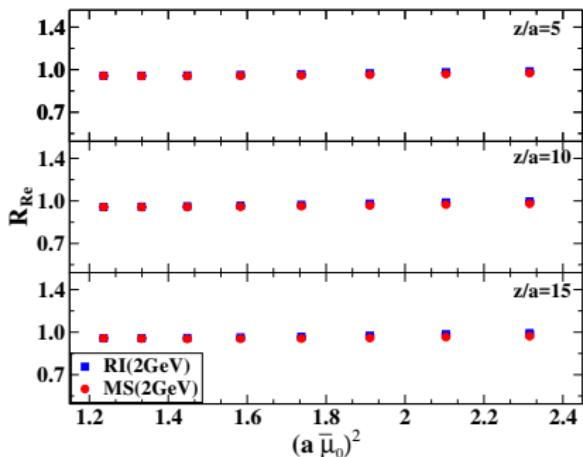
- ★ Uncertainties in Z-factors may have important implications on the final estimates for PDFs

# Systematic uncertainties

## Truncation effects in $C$ :

$$R_{\text{Re}(\text{Im})}^{\text{RI}'(\overline{\text{MS}})}(z, \bar{\mu}_0, \bar{\mu}_0'; \bar{\mu}) \equiv \frac{Z_{\text{Re}(\text{Im})}^{\text{RI}'(\overline{\text{MS}})}(z, \bar{\mu}_0; \bar{\mu})}{Z_{\text{Re}(\text{Im})}^{\text{RI}'(\overline{\text{MS}})}(z, \bar{\mu}_0'; \bar{\mu})}, \quad (\bar{\mu}_0' = 2.67 \text{ GeV})$$

Evolution to 2 GeV in  $\text{RI}'$  and  $\overline{\text{MS}}$  schemes:  
slope in  $R$  reveals truncation effect in conversion factor



★ Effect in Real part:  $\sim 2\%$

★ Effect in Imaginary part:  $\sim 100\%$

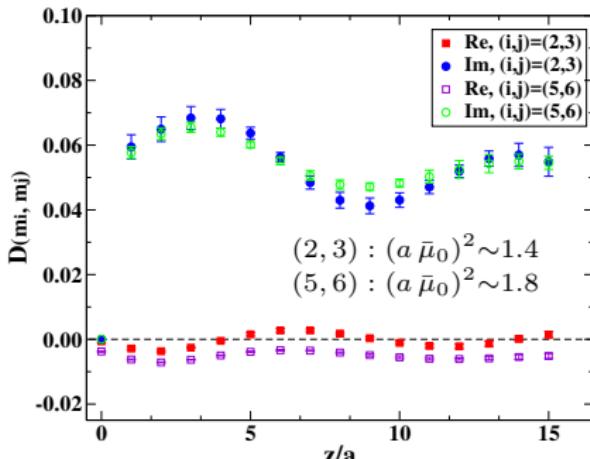
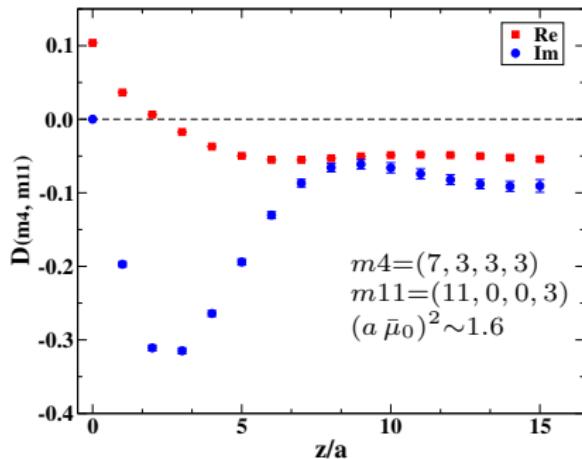
$(\text{Im}[Z^{\overline{\text{MS}}}] = 0 \text{ in Dim. Regul.})$

# Systematic uncertainties

## Lattice artifacts:

$$D_{\text{Re}(\text{Im})}(\bar{\mu}_0, \bar{\mu}'_0) \equiv \frac{Z_{\text{Re}(\text{Im})}^{\text{RI}'}(\bar{\mu}_0) - Z_{\text{Re}(\text{Im})}^{\text{RI}'}(\bar{\mu}'_0)}{Z_{\text{Re}(\text{Im})}^{\text{RI}'}(\bar{\mu}_0)}$$

Comparison of Z-factors on same  $(a \bar{\mu}_0)^2$ , different components:



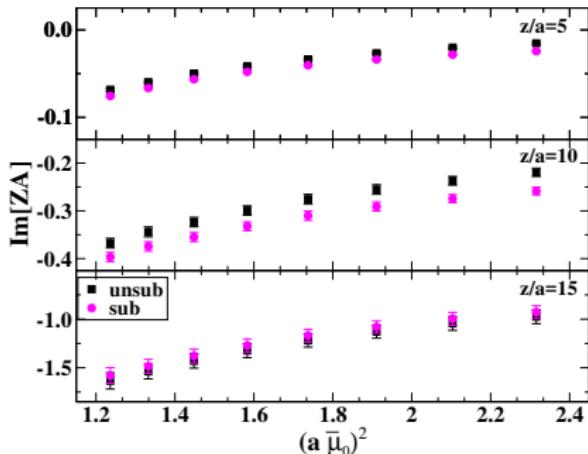
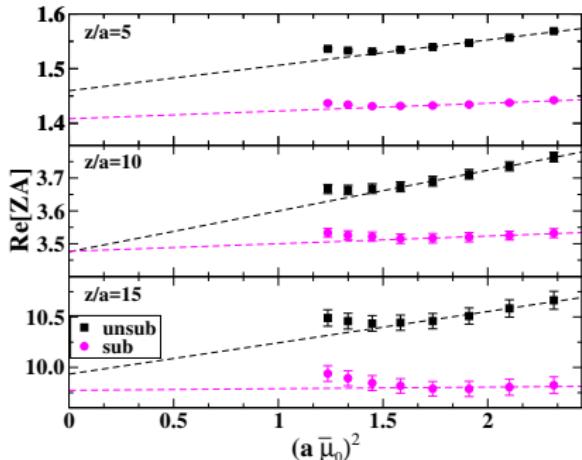
- ★ Effect in Real part:  $\sim 2 - 5\%$
- ★ Effect in Imaginary part:  $\sim 10\%$

# Refining Renormalization

## ★ Improvement Technique:

- Computation of 1-loop lattice artifacts to  $\mathcal{O}(g^2 a^\infty)$
- Subtraction of lattice artifacts from non-perturbative estimated

## ★ Application to the quasi-PDFs: PRELIMINARY

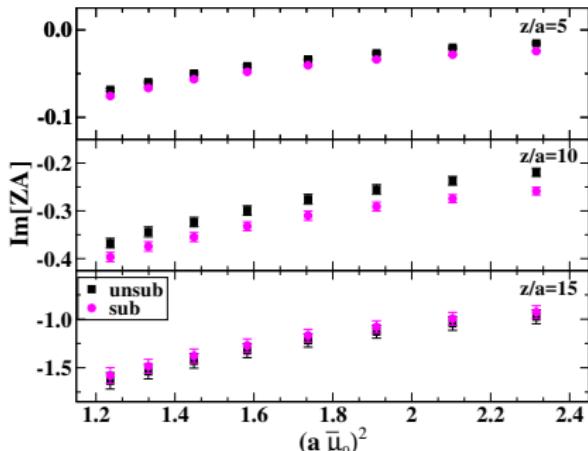
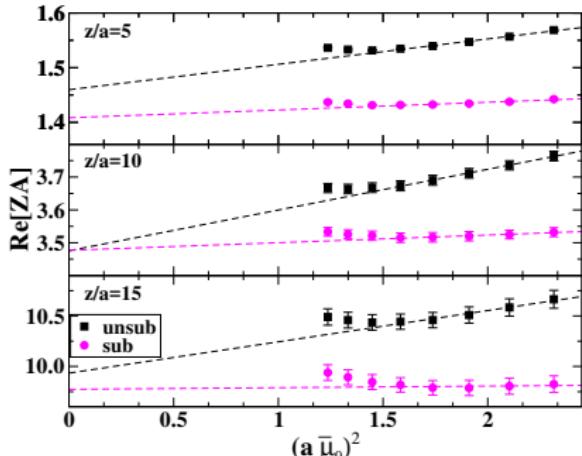


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★ Real part significantly improved

★ Mild change in imaginary part (expected to change with smearing)

- Behavior might be a consequence of absence of smearing in pert. calculation

# Renormalization of ensemble @ physical point

★ Application of stout smearing as in ME (0, 5, 10, 15, 20 steps)

★ Calculation for a range of RI' scales:  $a \bar{\mu}_0 = \frac{2\pi}{48}(\frac{n_t}{2} + \frac{\pi}{2}, n_x, n_y, n_z)$

$(n_t, n_x, n_y, n_z)$	$(a \bar{\mu}_0)^2$	$\bar{\mu}$ (GeV)	$\hat{P}$
(8,3,3,3)	0.772	3.465	0.280
(9,3,3,3)	0.849	3.811	0.306
(10,3,3,3)	0.935	4.195	0.337
(6,4,4,4)	1.003	4.503	0.256
(7,4,4,4)	1.063	4.772	0.251
(8,4,4,4)	1.132	5.079	0.251
(9,4,4,4)	1.209	5.425	0.256
(10,4,4,4)	1.295	5.810	0.268
(11,4,4,4)	1.389	6.233	0.283
(5,5,5,5)	1.415	6.348	0.283
(7,5,5,5)	1.526	6.848	0.261
(9,5,5,5)	1.672	7.501	0.250
(11,5,5,5)	1.852	8.308	0.254
(6,6,6,6)	2.032	9.116	0.285
(8,6,6,6)	2.160	9.692	0.265
(10,6,6,6)	2.323	10.423	0.253
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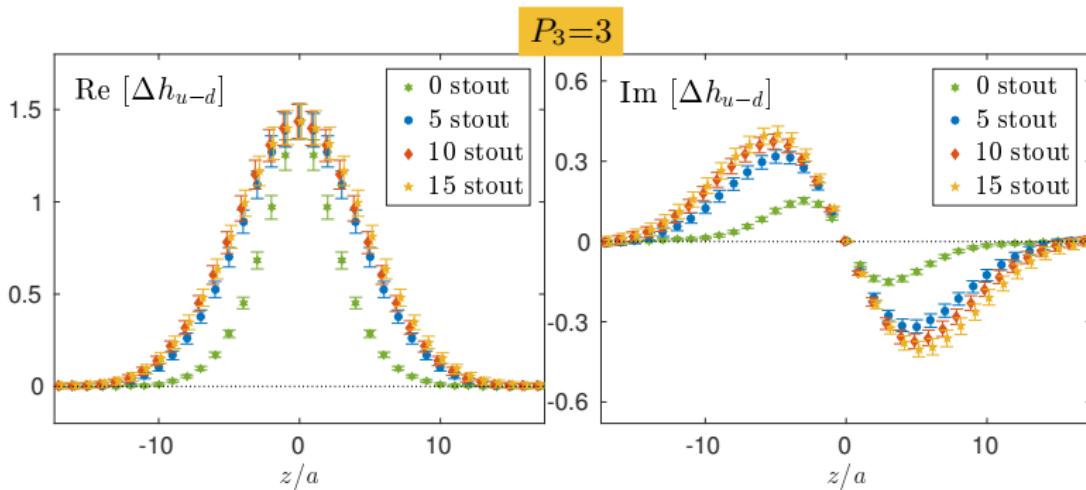
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fit:  
 $Z^{\overline{\text{MS}}} = Z_0^{\overline{\text{MS}}} + (a\bar{\mu}_0)^2 Z_1^{\overline{\text{MS}}}$

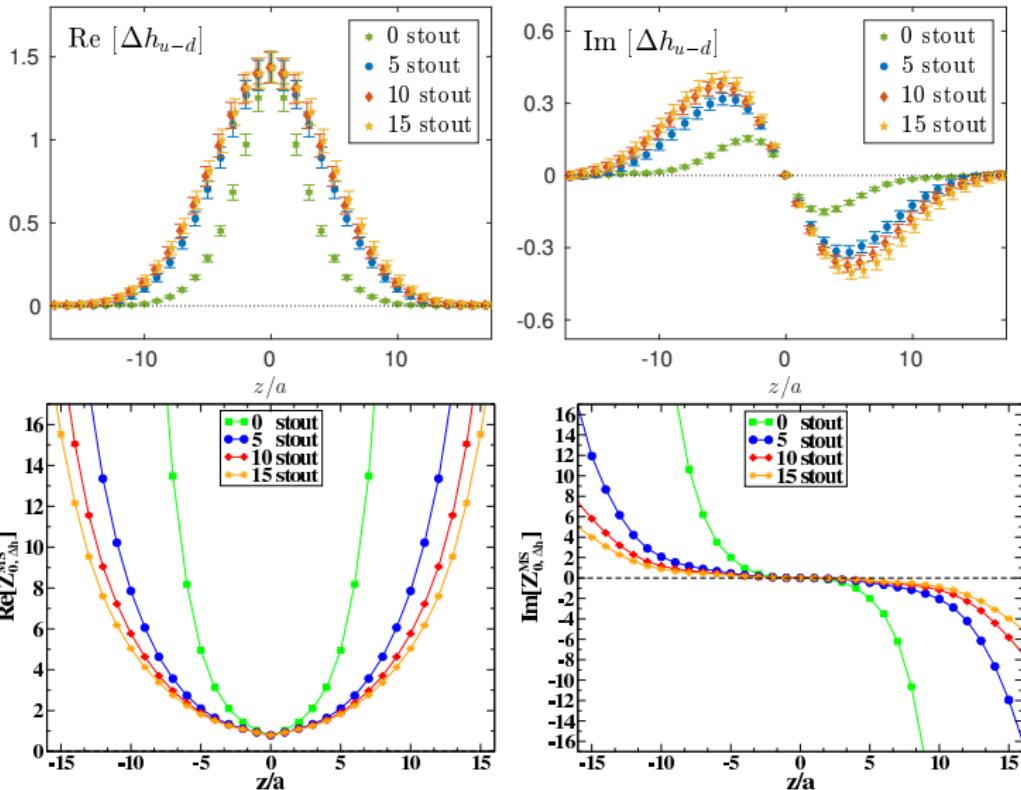
final estimates:  
 $(a\bar{\mu}_0)^2 \epsilon [1 - 2.6]$

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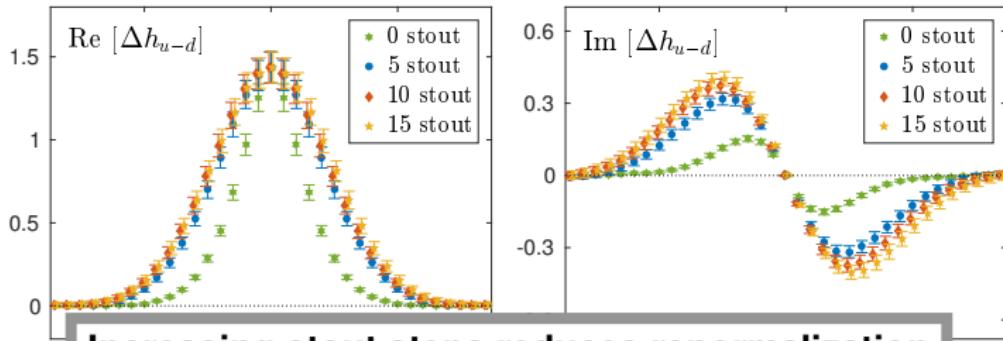
- ★ Smearing improves the signal-to-noise ratio
- ★ Smearing important when renormalization not available (suppresses linear divergence)
- ★ Stout smearing used (up to 20 steps)



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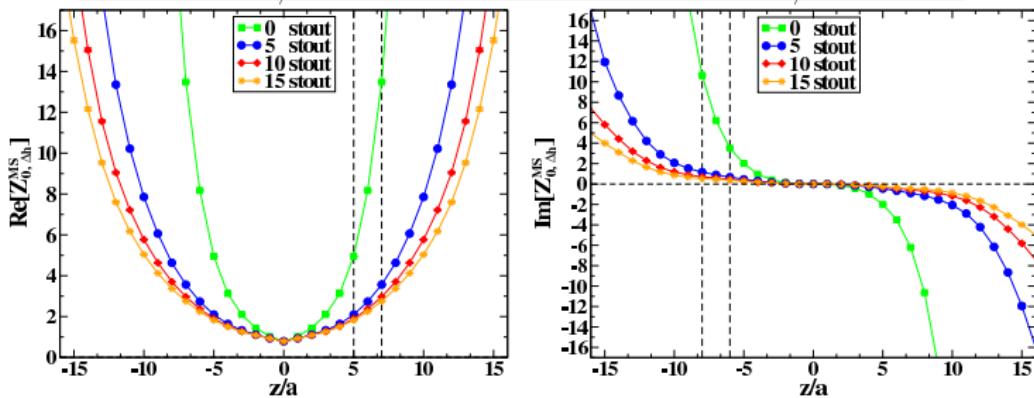


# Renormalization of ensemble @ physical point



Increasing stout steps reduces renormalization

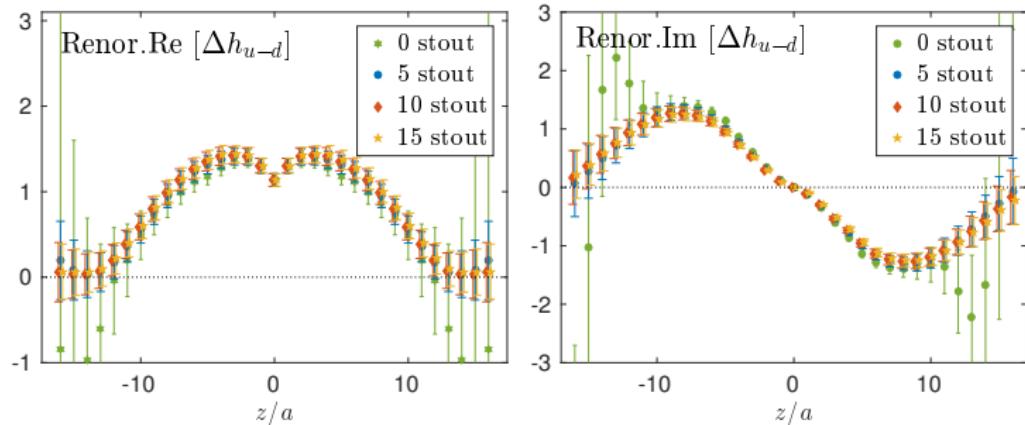
$P_3=3$



# Renormalization of ensemble @ physical point

★ Renormalized ME must be independent of stout steps

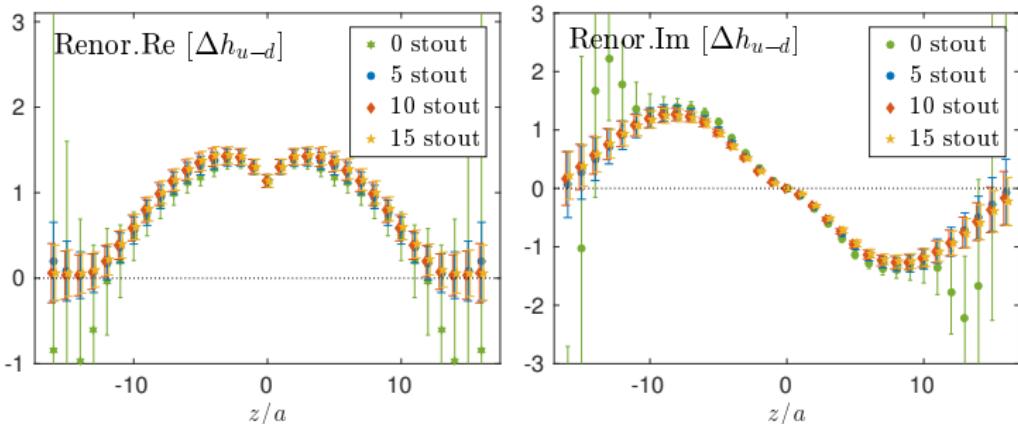
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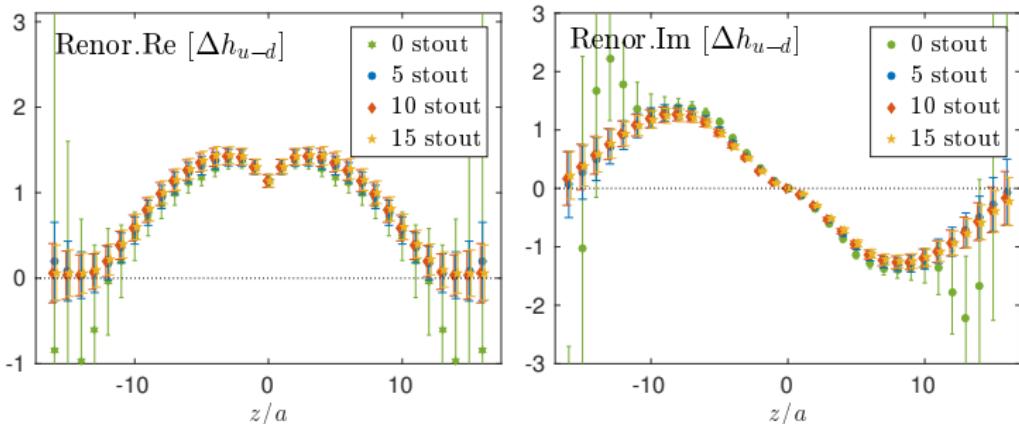


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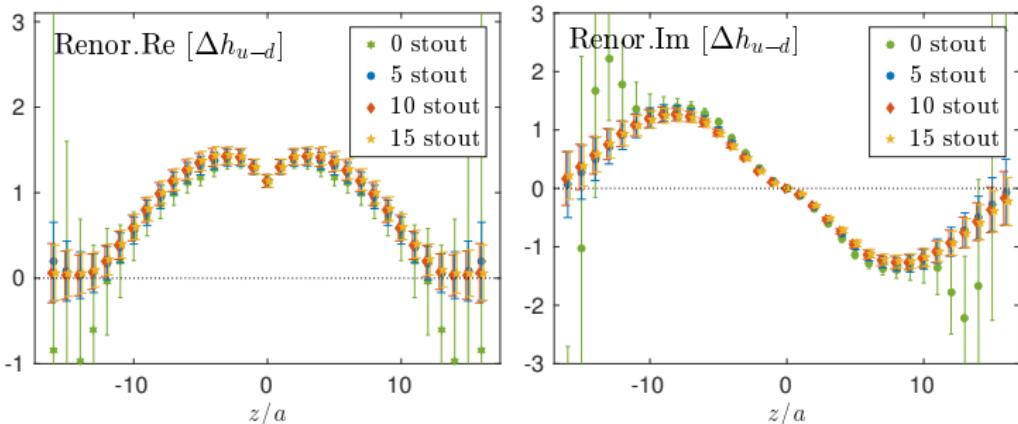


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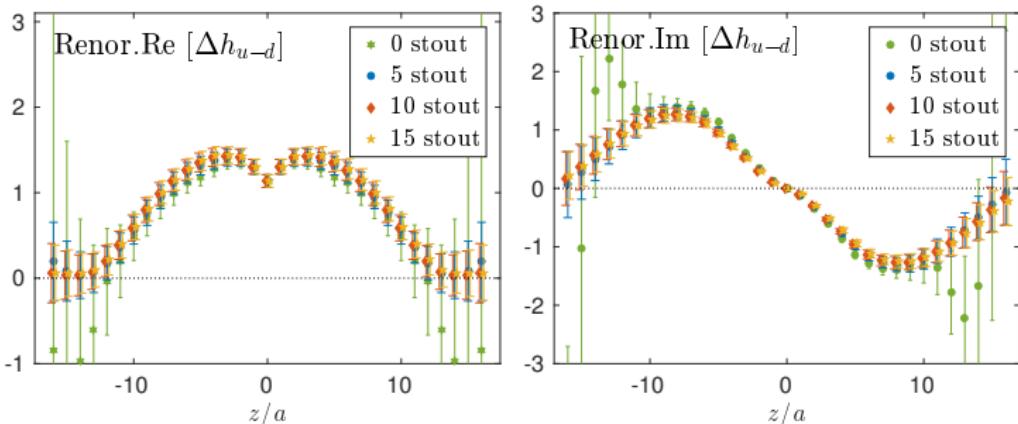


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Beyond the scope of this talk:

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See talks by:

- ★ Fernanda Steffens, Friday @ 16:30pm
- ★ Krzysztof Cichy, Saturday @ 09:40am

E

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# THANK YOU



TMD Topical Collaboration



Grant No. PHY-1714407