#### Deuteron Electromagnetic Form Factors in Holographic QCD

Valery Lyubovitskij

Institut für Theoretische Physik, Universität Tübingen Kepler Center for Astro and Particle Physics, Germany



*in collaboration with* Thomas Gutsche Ivan Schmidt Alfredo Vega PRD 91 (2015) 114001

6th International Workshop "HEP 2016", 8 January 2016, Valparaiso, Chile

### Introduction

- AdS/QCD = Holographic QCD (HQCD) approximation to QCD: attempt to model Hadronic Physics in terms of fields/strings living in extra dimensions – anti-de Sitter (AdS) space
- Motivation: AdS/CFT correspondence 1998 (Maldacena, Polyakov, Witten et al) Extra-Dimensional (ED) theories including gravity are holographically equivalent to gauge theories living on boundary of ED space
- Symmetry arguments: Conformal group acting in boundary theory isomorphic to SO(4, 2) the isometry group of AdS<sub>5</sub> space
- HQCD models reproduce main features of QCD at low and high energies: chiral symmetry, confinement, power scaling of hadron form factors, Regge trajectories for mass spectrum see papers by Brodsky-Teramond, our group, ...
- PDFs, GPDs, Structure functions
- LFWFs from matching AdS/QCD and LF QCD

- We first propose an effective action describing the dynamics of the deuteron in the presence of an external vector field
- Based on this action the deuteron electromagnetic form factors are calculated, displaying the correct  $1/Q^{10}$  power scaling for large  $Q^2$  values.
- This finding is consistent with quark counting rules and the earlier observation that this result holds in confining gauge/gravity duals.
- The  $Q^2$  dependence of the deuteron form factors is defined by a single and universal scale parameter  $\kappa$ , which is fixed from data.

- Effective action in terms of AdS fields  $d^M(x,z)$  and  $V^M(x,z)$
- $d^M(x,z)$  dual to Fock component contributing to deuteron with twist  $\tau = 6$
- $V^M(x,z)$  dual to the electromagnetic field

$$S = \int d^{4}x dz \, e^{-\varphi(z)} \left[ -\frac{1}{4} F_{MN}(x, z) F^{MN}(x, z) - D^{M} d_{N}^{\dagger}(x, z) D_{M} d^{N}(x, z) \right. \\ \left. - \, i c_{2} F^{MN}(x, z) d_{M}^{\dagger}(x, z) d_{N}(x, z) \right. \\ \left. + \, \frac{c_{3}}{4M_{d}^{2}} \, e^{2A(z)} \, \partial^{M} F^{NK}(x, z) \left( -d_{M}^{\dagger}(x, z) \stackrel{\leftrightarrow}{D}_{K} d_{N}(x, z) + \text{H.c.} \right) \\ \left. + \, d_{M}^{\dagger}(x, z) \left( \mu^{2} + U(z) \right) d^{M}(x, z) \right]$$

•  $A(z) = \log(R/z)$ , R is the AdS radius  $F^{MN} = \partial^M V^N - \partial^N V^M$  - stress tensor of vector field  $D^M = \partial^M - ieV^M(x, z)$  - covariant derivative  $\mu^2 R^2 = (\Delta - 1)(\Delta - 3)$  - five-dimensional mass  $\varphi(z) = \kappa^2 z^2$  is the background dilaton field  $\Delta = 6 + L$  is the dimension of  $d^M(x, z)$ 

 $\boldsymbol{L}$  is the maximal value of orbital angular momentum

 $U(z) = U_0 \varphi(z)/R^2$  is the confinement potential

 $U_0$  is constant fixed the deuteron mass.

Use axial gauge for both vector fields  $d^{z}(x, z) = 0$  and  $V^{z}(x, z) = 0$ 

 First perform Kaluza-Klein (KK) decomposition for vector AdS field dual to deuteron

$$d^{\mu}(x,z) = \exp\left[\frac{\varphi(z) - A(z)}{2}\right] \sum_{n} d^{\mu}_{n}(x)\Phi_{n}(z),$$

 $d_n^{\mu}(x)$  is the tower of the KK fields dual to the deuteron fields with radial quantum number *n* and twist-dimension  $\tau = 6$ , and  $\Phi_n(z)$  are their bulk profiles. Then we derive the Schrödinger-type equation of motion for the bulk profile

$$\left[-\frac{d^2}{dz^2} + \frac{4(L+4)^2 - 1}{4z^2} + \kappa^4 z^2 + \kappa^2 U_0\right]\Phi_n(z) = M_{d,n}^2 \Phi_n(z)$$

The analytical solutions of this EOM read

$$\Phi_n(z) = \sqrt{\frac{2n!}{(n+L+4)!}} \kappa^{L+5} z^{L+9/2} e^{-\kappa^2 z^2/2} L_n^{L+4}(\kappa^2 z^2),$$
  
$$M_{d,n}^2 = 4\kappa^2 \left[ n + \frac{L+5}{2} + \frac{U_0}{4} \right],$$

where  $L_n^m(x)$  are the generalized Laguerre polynomials

- Restricting to the ground state (n = 0, L = 0) we get  $M_d = 2\kappa \sqrt{\frac{5}{2} + \frac{U_0}{4}}$
- $\kappa = 190$  MeV (fitted from data on electromagnetic deuteron form factors)
- Using central value of deuteron mass  $M_d = 1.875613$  GeV we fix  $U_0 = 87.4494$

- We can compare this value for the deuteron scale parameter to the analogous one of  $\kappa_N$  defining the nucleon properties mass and electromagnetic form factors. In description of nucleon case we fixed the value to  $\kappa_N \simeq 380$  MeV, which is 2 times bigger than the deuteron scale parameter  $\kappa$ .
- Difference between the nucleon and deuteron scale parameters can be related to the change of size of the hadronic systems - the deuteron as a two-nucleon bound state is 2 times larger than the nucleon.

 In case of the vector field dual to the electromagnetic field we perform a Fourier transform with respect to the Minkowski coordinate

$$V_{\mu}(x,z) = \int \frac{d^4q}{(2\pi)^4} e^{-iqx} V_{\mu}(q) V(q,z)$$

where V(q, z) is its bulk profile obeying the following EOM:

$$\partial_z \left( \frac{e^{-\varphi(z)}}{z} \, \partial_z V(q,z) \right) + q^2 \frac{e^{-\varphi(z)}}{z} \, V(q,z) = 0 \; .$$

Its analytical solution can be written in the form of

$$V(Q,z) = \kappa^2 z^2 \int_0^1 \frac{dx}{(1-x)^2} e^{-\kappa^2 z^2 x/(1-x)} x^a, \quad a = \frac{Q^2}{4\kappa^2}, \quad Q^2 = -q^2.$$

The gauge-invariant matrix element describing the interaction of the deuteron with the external vector field (dual to the electromagnetic field) reads

$$M_{\rm inv}^{\mu}(p,p') = -\left(G_1(Q^2)\epsilon^*(p')\cdot\epsilon(p) - \frac{G_3(Q^2)}{2M_d^2}\epsilon^*(p')\cdot q\,\epsilon(p)\cdot q\right)(p+p')^{\mu}$$
$$- G_2(Q^2)\left(\epsilon^{\mu}(p)\,\epsilon^*(p')\cdot q - \epsilon^{*\mu}(p')\,\epsilon(p)\cdot q\right)$$

where  $\epsilon(\epsilon^*)$  and p(p') are the polarization and four-momentum of the initial (final) deuteron, and q = p' - p is the momentum transfer.

- Three EM form factors  $G_{1,2,3}$  of the deuteron are related to the charge  $G_C$ , quadrupole  $G_Q$  and magnetic  $G_M$  form factors by
- Expressions for the form factors

$$G_C = G_1 + \frac{2}{3}\tau_d G_Q$$
,  $G_M = G_2$ ,  $G_Q = G_1 - G_2 + (1 + \tau_d)G_3$ ,  $\tau_d = \frac{Q^2}{4M_d^2}$ .

These form factors are normalized at zero recoil as

$$G_C(0) = 1$$
,  $G_Q(0) = M_d^2 \mathcal{Q}_d = 25.83$ ,  $G_M(0) = \frac{M_d}{M_N} \mu_d = 1.714$ ,

•  $Q_d = 7.3424 \text{ GeV}^{-2}$  and  $\mu_d = 0.8574 - \text{quadrupole}$  and magnetic moments of the deuteron.

• In our approach the deuteron form factor  $G_1(Q^2) = F(Q^2)$ , where  $F(Q^2)$  is the twist-6 hadronic form factor, which is given by the overlap of the square of the bulk profile dual to the deuteron wave function (twist-6 hadronic wave function) and the confined electromagnetic current

$$F(Q^2) = \int_{0}^{\infty} dz \, \Phi_0^2(z) \, V(Q, z) = \frac{\Gamma(6) \, \Gamma(a+1)}{\Gamma(a+6)}$$

where  $a=Q^2/(4\kappa^2).$ 

This formula follows from the result of Brodsky and Teramond – general and universal formula for hadronic FF with twist  $\tau$  in terms of bulk profile  $\phi_{\tau}(z) = \sqrt{\frac{2}{(\tau-2)!}} \kappa^{\tau-1} z^{\tau-3/2} e^{-\kappa^2 z^2/2}$  dual to the hadronic wave function with twist  $\tau$ :

$$F_{\tau}(Q^2) = \int_{0}^{\infty} dz \, \phi_{\tau}^2(z) \, V(Q,z) = \frac{\Gamma(\tau) \, \Gamma(a+1)}{\Gamma(a+\tau)}$$

• By analogy we calculate the other two deuteron FF  $G_2$  and  $G_3$ , expressed in terms of the same universal factor  $F(Q^2)$ 

 $G_i(Q^2) = c_i F(Q^2), \quad i = 2, 3.$ 

The parameters  $c_2$  and  $c_3$  are defined by normalization of the deuteron form factors as

$$c_2 = G_M(0) = 1.714$$
,  $c_3 = G_M(0) + G_Q(0) - 1 = 26.544$ .

- $F(Q^2)$  has the correct power-scaling  $F(Q^2) \sim 1/(Q^2)^5$  at large  $Q^2 \to \infty$
- It can also be written in the Brodsky-Ji-Lepage form derived within pQCD
- Deuteron FF is factorized in terms of the nucleon FF  $F_N(Q^2/4)$  and the so-called "reduced" nuclear form factor  $f_d(Q^2)$  (Brodsky-Ji-Lepage):

$$F_d(Q^2) = f_d(Q^2) F_N^2(Q^2/4)$$

Our result reads

$$F_d(Q^2) \equiv F(Q^2) = \frac{\Gamma(6) \,\Gamma(a+1)}{\Gamma(a+6)} = \frac{5!}{(a+1)\dots(a+5)} = f_d(Q^2) \,F_N^2(Q^2/4)$$

where predictions for  $f_d(Q^2)$  and  $F_N(Q^2/4)$  are

$$f_d(Q^2) = \frac{30(a+1)(a+2)}{(a+3)(a+4)(a+5)}, \quad F_N(Q^2/4) = \frac{2}{(a+1)(a+2)}$$

where  $a=Q^2/(4\kappa^2).$ 

- Numerical results for deuteron FF
- Shaded band corresponds to values of  $\kappa$  in range of 150 MeV <  $\kappa$  < 250 MeV.
- Increase of the parameter  $\kappa$  leads to an enhancement of the form factors.
- The best description of the data on the deuteron form factors is obtained for  $\kappa = 190$  MeV and is shown by the solid line.
- With  $\kappa = 190$  MeV our result for the deuteron charge radius

$$r_C = \sqrt{-6G'_C(0)} = \sqrt{\frac{137}{40\kappa^2} - Q_d} = 1.846$$
 fm close to data

 $r_C = 2.130 \pm 0.010$  fm.







- AdS/QCD = Holographic QCD (HQCD) approximation to QCD: attempt to model Hadronic Physics in terms of fields/strings living in extra dimensions – anti-de Sitter (AdS) space
- HQCD models reproduce main features of QCD at low and high energies
- Soft–wall holographic approach covariant and analytic model for hadron structure with confinement at large distances and conformal behavior at short distances
- Mesons, baryons, exotics from unified point view and including high Fock states
- Thanks to Stan Brodsky, Guy de Téramond, Werner Vogelsang for useful discussions and comments