

Euler-Lagrange equations for high energy effective actions in QCD and in gravity

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1 High energy hadron scattering

Froissart bound for total cross sections

$$\lim_{s \rightarrow \infty} \sigma_t(s) = \Im \frac{A(s, 0)}{2s} < c \ln^2 s, \quad s = 4E^2$$

Pomeranchuk theorem

$$\lim_{s \rightarrow \infty} \frac{\sigma_t^{pp}(s) - \sigma_t^{p\bar{p}}(s)}{\sigma_t^{pp}(s)} = 0$$

Pomeron contribution to scattering amplitudes

$$A(s, t) \approx i s \gamma_P^2(t) s^{\Delta_P + \alpha'_P t}, \quad \Delta_P \ll 1, \quad \alpha'_P \ll 1/m^2$$

Gribov's Pomeron effective action

$$S = \int dy d^2\rho \left(\phi^* (\partial_y + \alpha'_P \vec{\partial}^2 - \Delta_P) \phi + i\lambda(\phi^2 \phi^* + \phi \phi^{*2}) + \dots \right)$$

2 Gluon reggeization

QCD Born amplitude at high energies $s \gg t$

$$M_{AB}^{A'B'}|_{Born} = 2s g T_{A'A}^c \delta_{\lambda_{A'}\lambda_A} \frac{1}{t} g T_{B'B}^c \delta_{\lambda_{B'}\lambda_B}$$

Leading Logarithmic Approximation

$$M(s, t) = M|_{Born} s^{\omega(t)}, \quad \alpha_s \ln s \sim 1, \quad \alpha_s = \frac{g^2}{4\pi} \ll 1$$

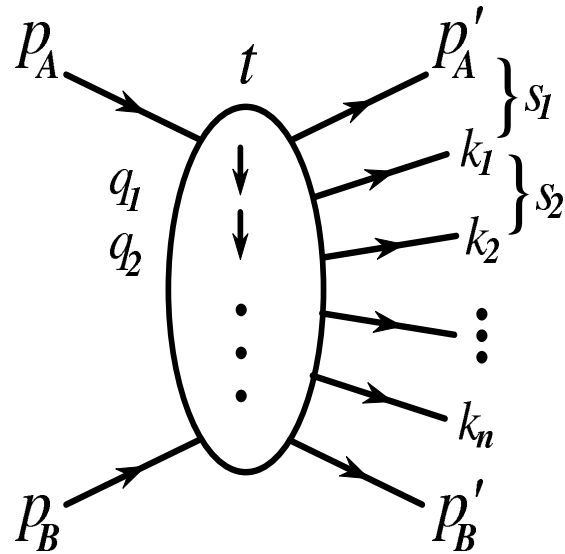
Gluon Regge trajectory in LLA

$$\omega(-|q|^2) = -\frac{\alpha_s N_c}{4\pi^2} \int d^2k \frac{|q|^2}{|k|^2 |q-k|^2} \approx -\frac{\alpha_s N_c}{2\pi} \ln \frac{|q|^2}{\lambda^2}$$

Bootstrap equation (FKL (1975))

$$\omega f = 1 + \left(\omega(-|k|^2) + \omega(-|q-k|^2) + \hat{K}_8 \right) f, \quad f = \frac{1}{\omega - \omega(-|q|^2)}$$

3 Production in multi-Regge kinematics



$$M_{2 \rightarrow 2+n}^{FKL} \sim \frac{s_1^{\omega_1}}{|q_1|^2} gT_{c_2 c_1}^{d_1} C(q_2, q_1) \frac{s_2^{\omega_2}}{|q_2|^2} \dots gT_{c_{n+1} c_n}^{d_n} C(q_{n+1}, q_n) \frac{s_{n+1}^{\omega_{n+1}}}{|q_{n+1}|^2},$$

$$C(q_2, q_1) = \frac{q_2 q_1^*}{q_2^* - q_1^*}, \quad \sigma_t = \sum_n \int d\Gamma_n |M_{2 \rightarrow 2+n}^{BFKL}|^2$$

4 BFKL Pomeron

Balitsky-Fadin-Kuraev-Lipatov equation (1975)

$$E \Psi(\vec{\rho}_1, \vec{\rho}_2) = H_{12} \Psi(\vec{\rho}_1, \vec{\rho}_2), \quad \sigma_t \sim s^\Delta, \quad \Delta = \frac{4\alpha N_c}{\pi} \ln 2$$

BFKL Hamiltonian in the operator form

$$H_{12} = \frac{1}{p_1 p_2^*} (\ln |\rho_{12}|^2) p_1 p_2^* + \frac{1}{p_1^* p_2} (\ln |\rho_{12}|^2) p_1^* p_2 + \ln |p_1 p_2|^2 - 4\psi(1),$$

Holomorphic separability

$$H_{12} = h_{12} + h_{12}^*, \quad [h_{12}, h_{12}^*] = 0, \quad \rho_{12} = \rho_1 - \rho_2, \quad \rho_r = x_r + iy_r$$

Möbius invariance (L. (1986))

$$\rho_k \rightarrow \frac{a\rho_k + b}{c\rho_k + d}, \quad \Psi = \left(\frac{\rho_{12}}{\rho_{10}\rho_{20}} \right)^m \left(\frac{\rho_{12}^*}{\rho_{10}^*\rho_{20}^*} \right)^{\tilde{m}}, \quad m = \gamma + \frac{n}{2}, \quad \gamma = \frac{1}{2} + i\nu$$

5 BKP equation in LLA

Bartels-Kwiecinski-Praszalowicz equation (1980)

$$E \Psi(\vec{\rho}_1, \dots, \vec{\rho}_n) = H \Psi(\vec{\rho}_1, \dots, \vec{\rho}_n), \quad H = \sum_{k < l} \frac{\vec{T}_k \vec{T}_l}{-N_c} H_{kl}$$

Holomorphic separability at large N_c (L. (1988))

$$H = \frac{1}{2} (h + h^*), \quad [h, h^*] = 0, \quad h = \sum_{k=1}^n h_{k,k+1},$$

Monodromy matrix and integrability (L. (1993))

$$t(u) = \prod_{k=1}^n \begin{pmatrix} u + \rho_k p_k & p_k \\ -\rho_k^2 p_k & u - \rho_k p_k \end{pmatrix} = \begin{pmatrix} A(u) & B(u) \\ C(u) & D(u) \end{pmatrix},$$

$$T(u) = A(u) + D(u), \quad [T(u), T(v)] = [T(u), h] = 0$$

6 Effective action for reggeized gluons

Locality in the rapidity space

$$y = \frac{1}{2} \ln \frac{\epsilon_k + |k|}{\epsilon_k - |k|}, \quad |y - y_0| < \eta, \quad \eta \ll \ln s$$

Anti-hermitian gluon and reggeized gluon fields

$$v_\mu(x) = -iT^a v_\mu^a(x), \quad A_\pm(x) = -iT^a A_\pm^a(x), \quad \delta A_\pm(x) = \partial_{\mp} A_\pm(x) = 0$$

Effective action for the reggeon interactions (L. (1995))

$$S = \int d^4x (L_{QCD} + Tr(V_+ \partial_\mu^2 A_- + V_- \partial_\mu^2 A_+) + 2Tr \partial_\sigma A_+ \partial_\sigma A_-)$$

Effective currents and gauge invariance

$$V_\pm = \frac{1}{g} \partial_\pm O(x^\pm), \quad O(x^\pm) \equiv -(D_\pm)^{-1} \overleftarrow{\partial}_\pm, \quad D_\pm = \partial_\pm + gv_\pm$$

7 Euler-Lagrange equations in QCD

Euler-Lagrange equation for high energy QCD

$$[D_\mu, G^{\mu\nu}] = j^\nu, \quad j^\pm = O(x^\pm)(\partial_\sigma^2 A^\pm)O^\pm(x^\pm), \quad j_\perp^\nu = 0$$

Equation for a quasi-elastic kinematics

$$[D_\mu, G^{\mu\nu}] = \delta_+^\nu O(x^+)(\partial_{\perp\sigma}^2 A^+(x^-, x_\perp))O^+(x^+)$$

Light-cone gauge $v'_+ = 0$

$$v'_+ = V^{-1}(v_+)(v_\mu + \frac{\partial_\mu}{g})V(v_+), \quad V(v_+) = e^{-\frac{g}{2} \int^{x^+} dx'^+ v_+}$$

Classical solution as a superposition of shock waves

$$\tilde{v}_\nu(x) = \delta_\nu^- \int \frac{d^2 z}{4\pi} \ln(|x - z|^2) \partial_\sigma^{\perp 2} A^+(x^-, z^\perp) = \delta_\nu^- A^+(x^-, x_\perp)$$

8 BFKL equation in $N = 4$ SUSY

BFKL kernel eigenvalue in two loops (F., L. (1998))

$$\omega = 4 \hat{a} \chi(n, \gamma) + 4 \hat{a}^2 \Delta(n, \gamma), \quad \hat{a} = g^2 N_c / (16\pi^2), \quad \gamma = i\nu + 1/2$$

Hermitian separability in $N = 4$ SUSY (K.,L. (2000))

$$\Delta(n, \gamma) = \phi(M) + \phi(M^*) - \frac{\rho(M) + \rho(M^*)}{2\hat{a}/\omega}, \quad M = \gamma + \frac{|n|}{2},$$

$$\rho(M) = \beta'(M) + \frac{1}{2}\zeta(2), \quad \beta'(z) = \frac{1}{4} \left[\Psi' \left(\frac{z+1}{2} \right) - \Psi' \left(\frac{z}{2} \right) \right]$$

Maximal transcendentality (K.L. (2002))

$$\phi(M) = 3\zeta(3) + \Psi''(M) - 2\Phi(M) + 2\beta'(M) \left(\Psi(1) - \Psi(M) \right)$$

9 Pomeron and graviton in N=4 SUSY

Diffusion approximation for the BFKL kernel

$$j = 2 - \Delta - \Delta \nu^2, \quad \gamma = 1 + \frac{j-2}{2} + i\nu$$

AdS/CFT relation with the graviton Regge trajectory

$$j = 2 + \frac{\alpha'}{2} t, \quad t = E^2/R^2, \quad \alpha' = \frac{R^2}{2} \Delta, \quad \lambda = g^2 N_c$$

Large coupling asymptotics for γ and Δ (KLOV, BPST, KL, GKL)

$$\gamma = 1 - \sqrt{1 + (j-2)/\Delta}, \quad \Delta = 2\lambda^{-1/2} + \lambda^{-1} - 1/4 \lambda^{-3/2} - 2(1+3\zeta_3)\lambda^{-2}$$

Exact expression for the slope of γ (KLOV, V., Basso)

$$\gamma'(2) = -\frac{\lambda}{24} + \frac{1}{2} \frac{\lambda^2}{24^2} - \frac{2}{5} \frac{\lambda^3}{24^2} + \frac{7}{20} \frac{\lambda^4}{24^4} - \frac{11}{35} \frac{\lambda^5}{24^5} + \dots = -\frac{\sqrt{\lambda}}{4} \frac{I_3(\sqrt{\lambda})}{I_2(\sqrt{\lambda})}$$

10 High energy action in gravity

Locality in the rapidity space

$$y = \frac{1}{2} \ln \frac{\epsilon_k + |k|}{\epsilon_k - |k|}, \quad |y - y_0| < \eta, \quad \eta \ll \ln s$$

Constraints for reggeized graviton fields

$$\delta A^{\pm\pm}(x) = 0, \quad \partial_{\pm} A^{\pm\pm}(x) = 0$$

Effective action for the high energy gravity (L. 2011)

$$S = -\frac{1}{2\kappa^2} \int d^4x \left(\sqrt{-g} R + \frac{\partial_+ j^-}{2} \partial_{\mu}^2 A^{++} + \frac{\partial_- j^+}{2} \partial_{\mu}^2 A^{--} \right)$$

Hamilton-Jacobi equation for effective currents $j^{\pm} = 2x^{\pm} - \omega^{\pm}$

$$g^{\mu\nu} \partial_{\mu} \omega^{\pm} \partial_{\nu} \omega^{\pm} = 0, \quad \partial_{\pm} j^{\mp} = h_{\pm\pm} - \left(h_{\rho\pm} - \frac{1}{2} \frac{\partial_{\rho}}{\partial_{\pm}} h_{\pm\pm} \right)^2 + \dots$$

11 Classical equation for effective action

Einstein-Hilbert equation for effective gravity

$$R_{\mu\nu} - \frac{1}{2}g_{\mu\nu} = \theta_{\mu\nu}, \quad \theta_{\mu\nu} = c_{\mu\nu}\partial_\sigma^2 A^{--} + d_{\mu\nu}\partial_\sigma^2 A^{++}$$

Coordinate transformation to the metrics $g'^{\pm\nu} = \eta^{\pm\nu}$

$$g'^{\rho\sigma} = g^{\mu\nu}\partial_\mu x'^\rho\partial_\nu x'^\sigma, \quad T'^{\rho\sigma} = \theta^{\mu\nu}\partial_\mu x'^\rho\partial_\nu x'^\sigma$$

Stress tensor in an arbitrary coordinate system

$$\theta_{\mu\nu} = \partial_\mu x'^+ \partial_\nu x'^+ \partial_\chi^2 A^{--}(x') + \partial_\mu x'^- \partial_\nu x'^- \partial_\chi^2 A^{++}(x')$$

Superposition of shock waves for a quasi-elastic kinematics

$$\tilde{g}_{\mu\nu}(x) = \eta_{\mu\nu} + \partial_\mu x'^+ \partial_\nu x'^+ A^{--}(x'), \quad \tilde{g}_{\mu\nu}^{AS} = \eta_{\mu\nu} + a \delta_\mu^+ \delta_\nu^+ \ln |x_\perp| \delta(x^+)$$

12 Reggeized particle vertices

Graviton-graviton-reggeized graviton vertex (L. (1982))

$$\gamma_{\mu'\nu',\mu\nu}^{+++} = \gamma_{\mu'\mu}^+ \gamma_{\nu'\nu}^+$$

Gluon-gluon-reggeized gluon vertex

$$\gamma_{\mu'\mu}^+ = -k^+ \delta_{\mu\mu'} + k_{\mu'} \delta_{\mu}^+ + k'_{\mu} \delta_{\mu'}^+ + q^2 \frac{\delta_{\mu}^+ \delta_{\mu'}^+}{2k^+}$$

Reggeon-Reggeon-graviton vertex (agrees with Steinmann)

$$\gamma_{\mu\nu} = \gamma_{\mu} \gamma_{\nu} - q_1^2 q_2^2 \left(\frac{n^+}{k^+} - \frac{n^-}{k^-} \right)_{\mu} \left(\frac{n^+}{k^+} - \frac{n^-}{k^-} \right)_{\nu}, \quad k = q_1 - q_2$$

Reggeized gluon - reggeized gluon - gluon vertex

$$\gamma_{\mu} = -q_{1\mu}^{\perp} - q_{2\mu}^{\perp} + n_{\mu}^+ \left(k^- + \frac{q_1^2}{k^+} \right) - n_{\mu}^- \left(k^+ + \frac{q_2^2}{k^-} \right)$$

13 Graviton trajectory at supergravity

Graviton Regge trajectory (L. (1982))

$$j = 2 + \omega, \quad \omega(q^2) = \frac{\alpha}{\pi} \int \frac{q^2 d^2k}{k^2(q-k)^2} f(k, q), \quad \alpha = \frac{\kappa^2}{8\pi^2},$$

$$f(k, q) = (k, q-k)^2 \left(\frac{1}{k^2} + \frac{1}{(q-k)^2} \right) - q^2 + \frac{N}{2} (k, q-k)$$

Gravitino action

$$S_{3/2} = -\frac{1}{2} \epsilon^{\mu\nu\rho\sigma} \int d^4x \sum_{r=1}^N \bar{\psi}_\mu^r \gamma_5 \gamma_\nu \partial_\rho \psi_\sigma^r$$

Divergencies of the graviton Regge trajectory

$$\omega(q^2) = -\alpha |q|^2 \left(\ln \frac{|q|^2}{\lambda^2} + \frac{N-4}{2} \ln \frac{|\Lambda|^2}{|q|^2} \right)$$

14 Double-logarithms in gravity

Mellin representation for the scattering amplitude

$$A(s, t) = A_{Born} s^{-\alpha|q|^2 \ln \frac{|q|^2}{\lambda^2}} \Phi(\xi), \quad \Phi(\xi) = \int_{a-i\infty}^{a+i\infty} \frac{d\omega}{2\pi i \omega} \left(\frac{s}{|q|^2} \right)^\omega f_\omega$$

Infrared evolution equation for super-gravity (BLS (2012))

$$f_\omega = 1 + \alpha|q|^2 \left(\frac{d}{d\omega} \frac{f_\omega}{\omega} - \frac{N-6}{2} \frac{f_\omega^2}{\omega^2} \right), \quad \xi = \alpha|q|^2 \ln^2 \frac{s}{|q|^2}$$

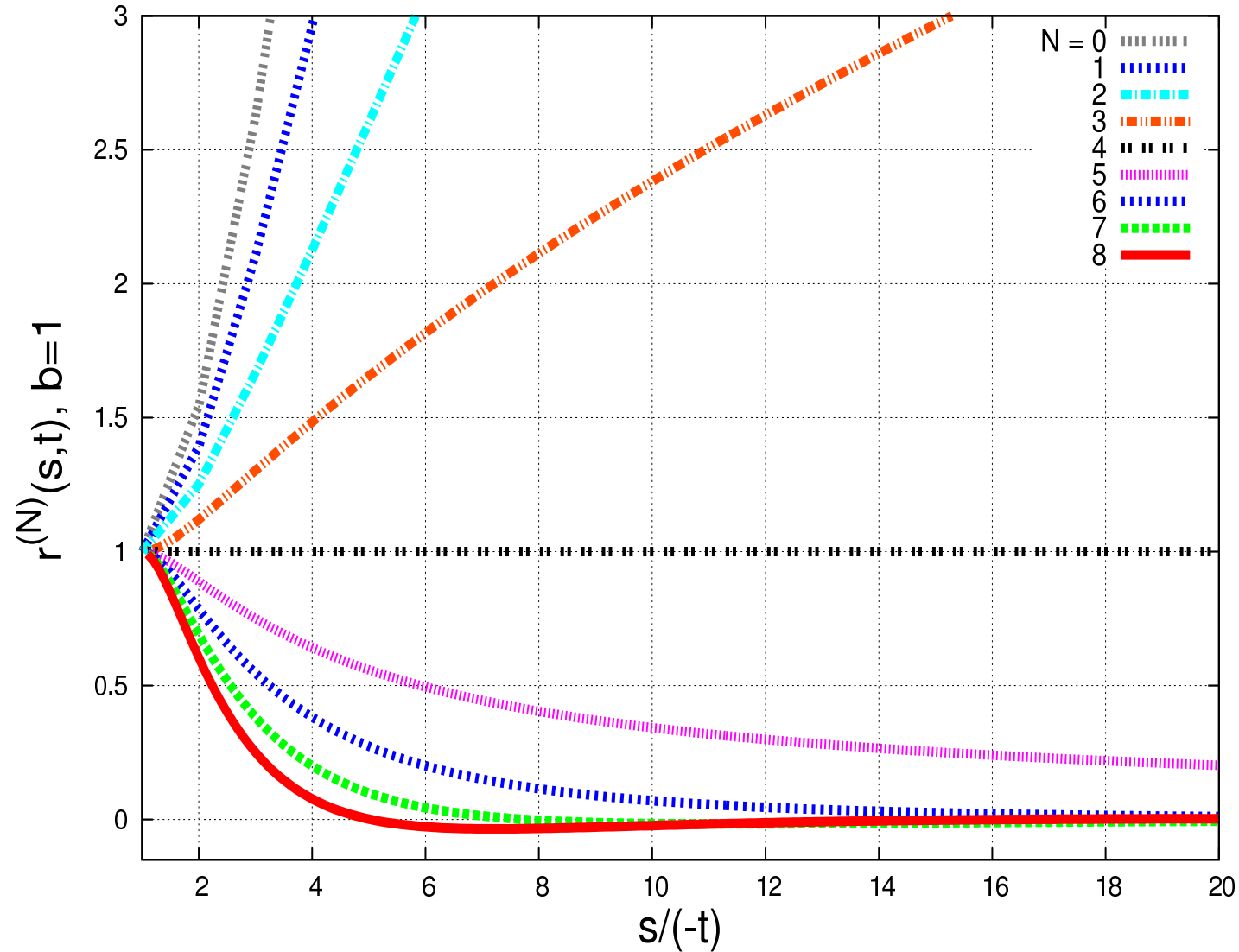
Its solution in terms of the parabolic cylinder function

$$\frac{f_\omega}{\omega} = \frac{2}{6-N} \frac{1}{\sqrt{b}} \frac{d}{dx} \ln \left(e^{\frac{x^2}{4}} D_{\frac{6-N}{2}}(x) \right), \quad x = \frac{\omega}{\sqrt{b}}$$

Perturbative expansion

$$\Phi(\xi) = 1 - \frac{N-4}{2} \frac{\xi}{2} + \frac{(N-4)(N-3)}{2} \frac{\xi^2}{4!} - \frac{N-4}{8} (5N^2 - 26N + 36) \frac{\xi^3}{6!} + \dots$$

15 Amplitudes in DL approximation



16 Double-logarithmic eikonal picture

Scattering amplitude in the eikonal approximation

$$A_{DL}(s, t) = -2is s^{-\alpha|q|^2 \ln \frac{|q|^2}{\mu^2}} \int d^2\rho e^{i\vec{q}\vec{\rho}} \left(e^{i\delta_{DL}(\vec{\rho}, \ln s)} - 1 \right),$$

Double logarithmic approximation for the phase (BLS)

$$\delta_{DL}(\vec{\rho}, \ln s) = \frac{s}{2} \frac{\kappa^2}{(2\pi)^2} \int \frac{d^2q}{|q|^2} e^{-i\vec{q}\vec{\rho}} \Phi(\xi)$$

Eikonal phase in $N = 8$ SUSY at small impact parameters

$$\delta_{DL}^{N=8}(\vec{\rho}, \ln s) = \frac{s}{2} \frac{\kappa^2}{(2\pi)^2} \pi \ln \frac{1}{\lambda^2 \alpha \ln^2(\rho^2 s)}, \quad \rho^2 \ll \alpha \ln^2(\rho^2 s)$$

Agreement with exact calculations (BLS)

$$A_4^{N=8} = \frac{\kappa^2 s^2}{|q|^2} (-i\pi s) \alpha^2 |q|^2 \frac{\ln^3 \frac{s}{|q|^2}}{3}$$

17 Discussion

1. Gluon and graviton reggeization
2. BFKL equation and its Möbius invariance
3. BKP equations and integrability
4. Effective action for high energy QCD
5. Euler-Lagrange equation for effective QCD
6. Pomeron and reggeized graviton in $N = 4$ SUSY
7. Effective action for reggeized gravitons
8. Euler-Lagrange equation for high energy gravity
9. Graviton scattering in the DL approximation.