

Frontiers in Particle Physics - 2024

# Probing New Physics with JWST observations of high redshift massive galaxies

Priyank Parashari

Based on P. Parashari & R. Laha, MNRAS: Letters, 526, L63-L69 (2023) arXiv: 2305.00999

August 09, 2024





# Introduction

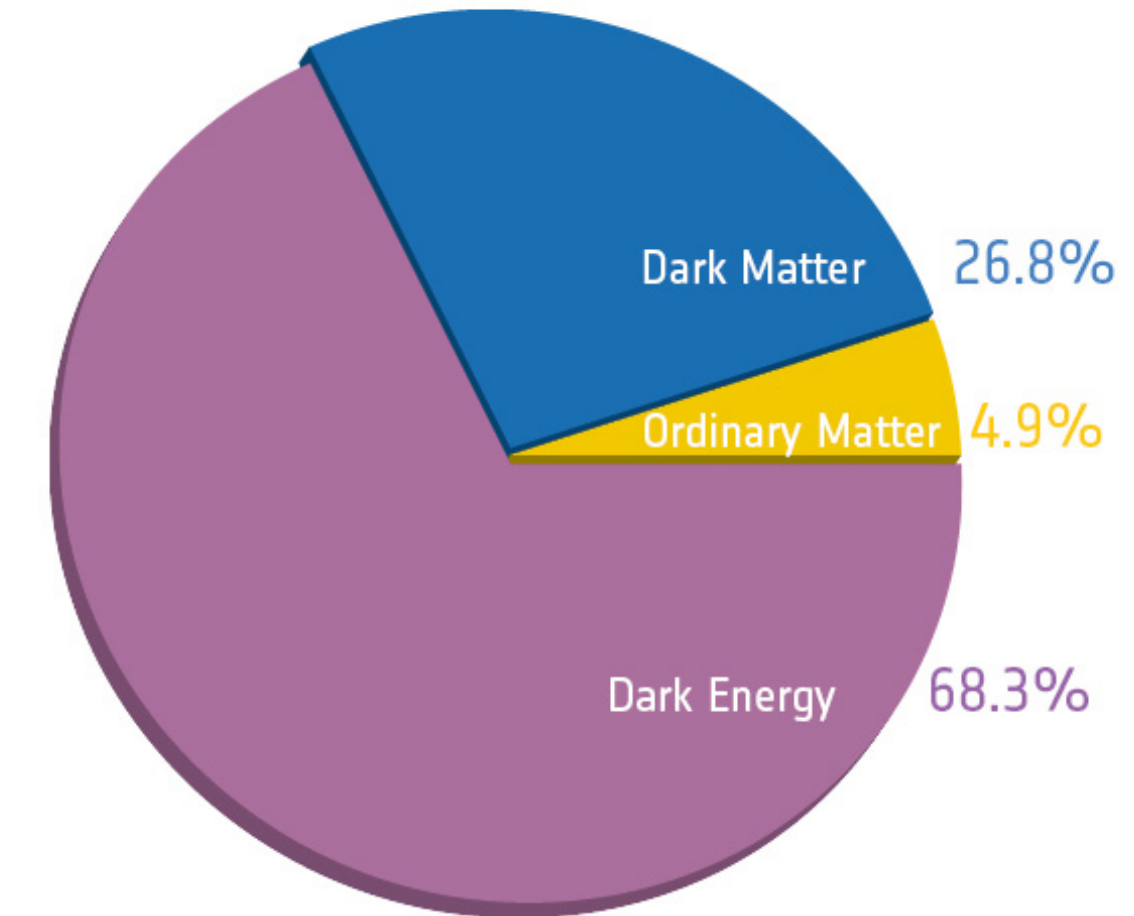
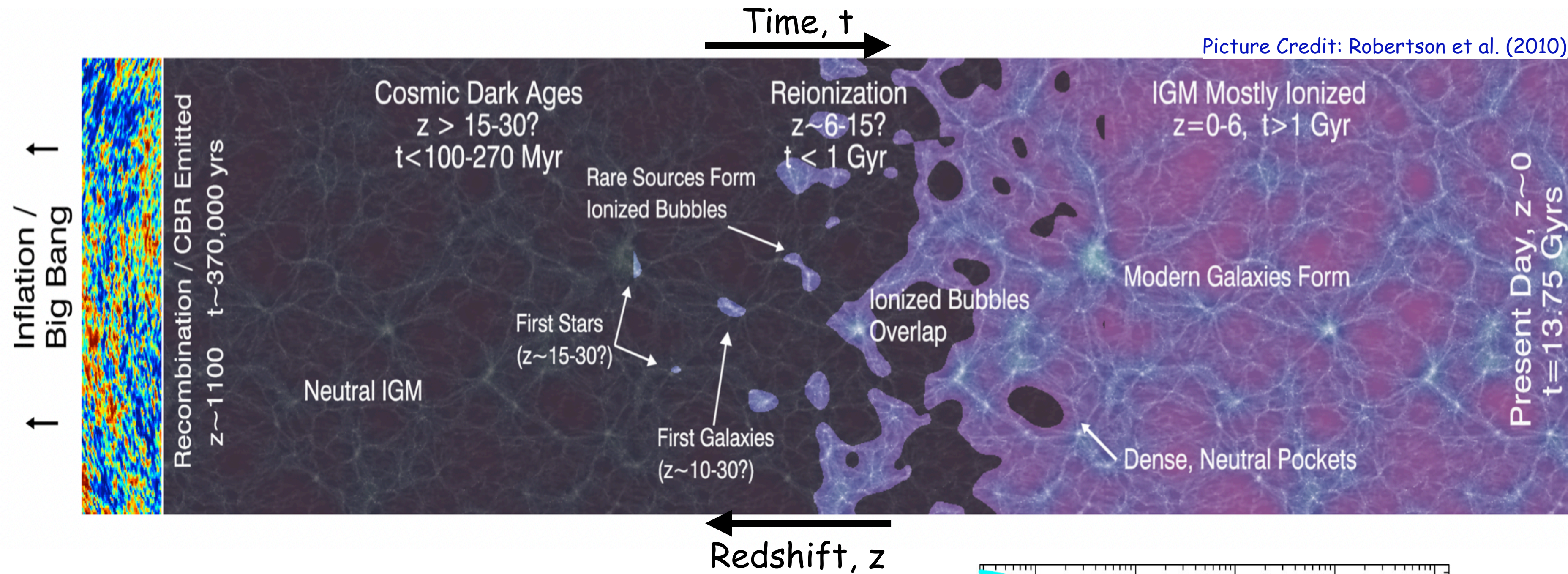
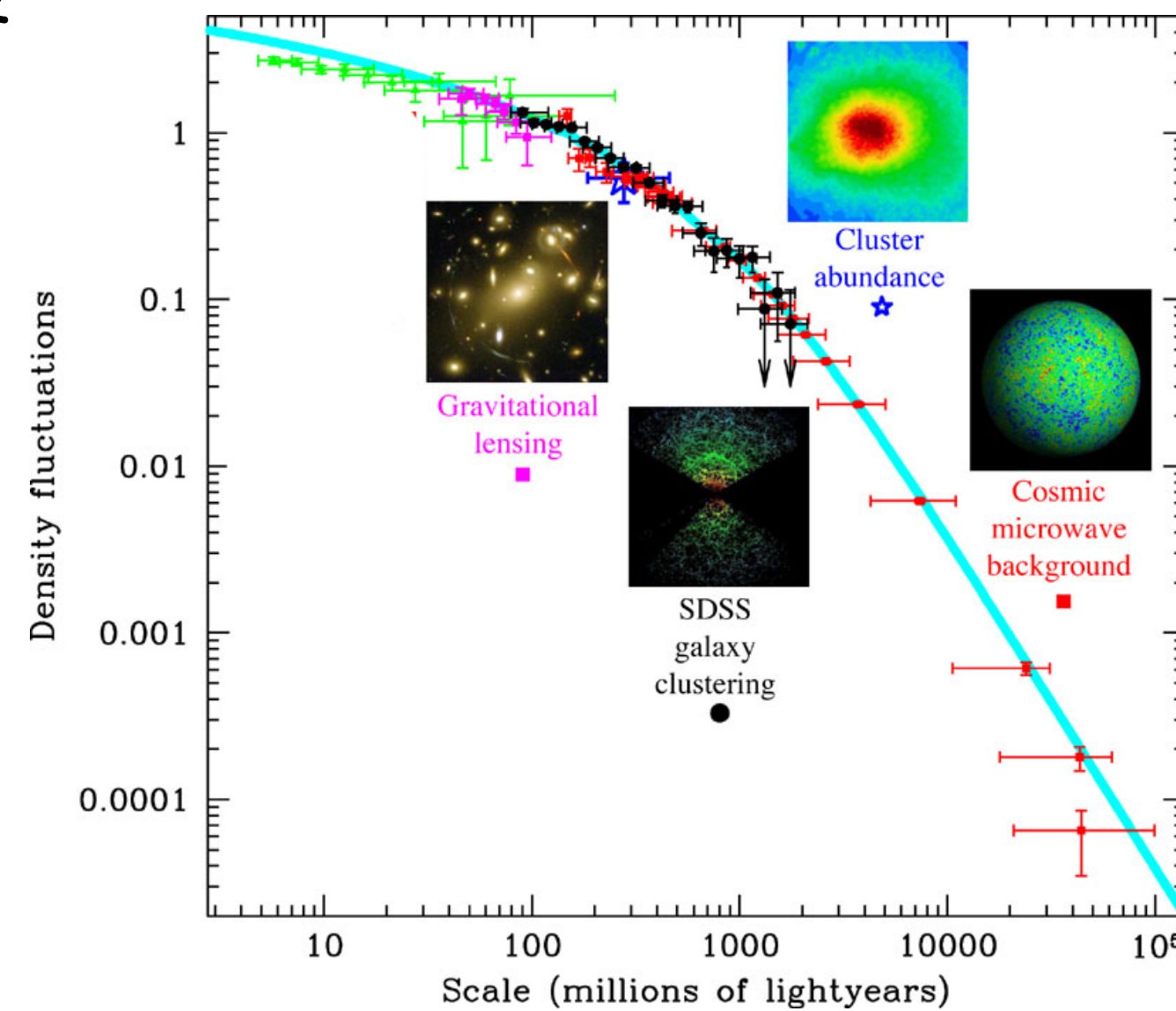
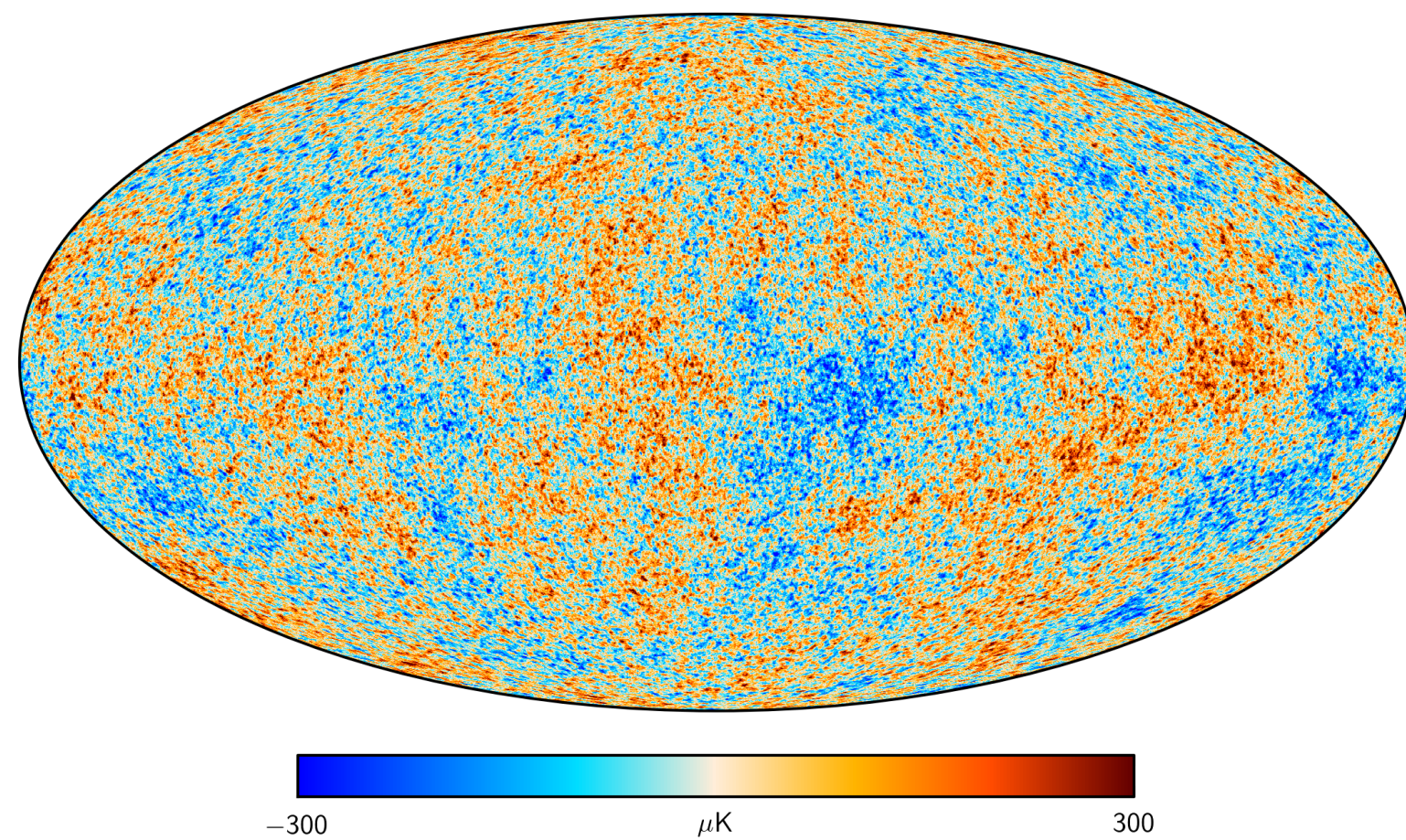


Figure credit: Planck collaboration



## Standard $\Lambda$ CDM Cosmology

- Six parameters model ( $\omega_b, \omega_c, \tau_{\text{reio}}, \theta, A_s, n_s$ )
- Fits the cosmological observations quite well.
- Some discrepancies? ( $H_0, \sigma_8$ )



# Introduction

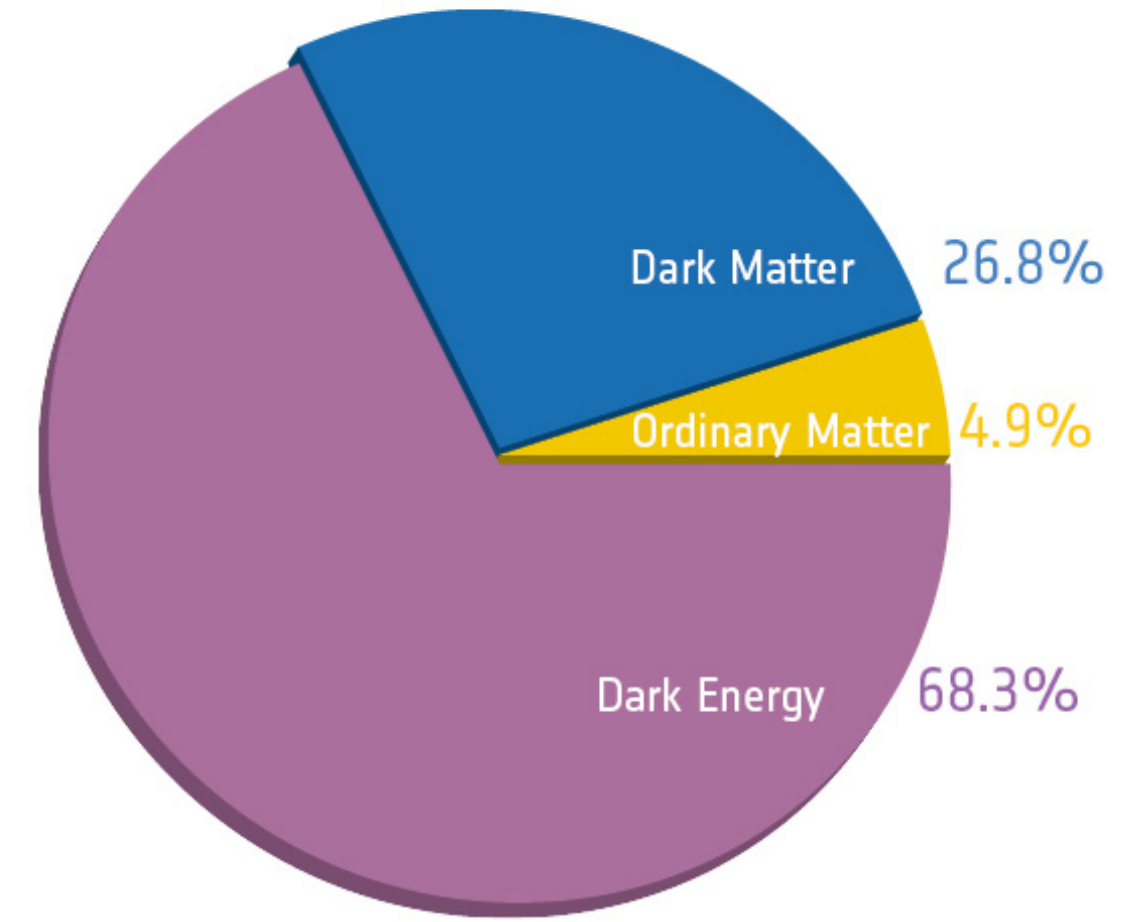
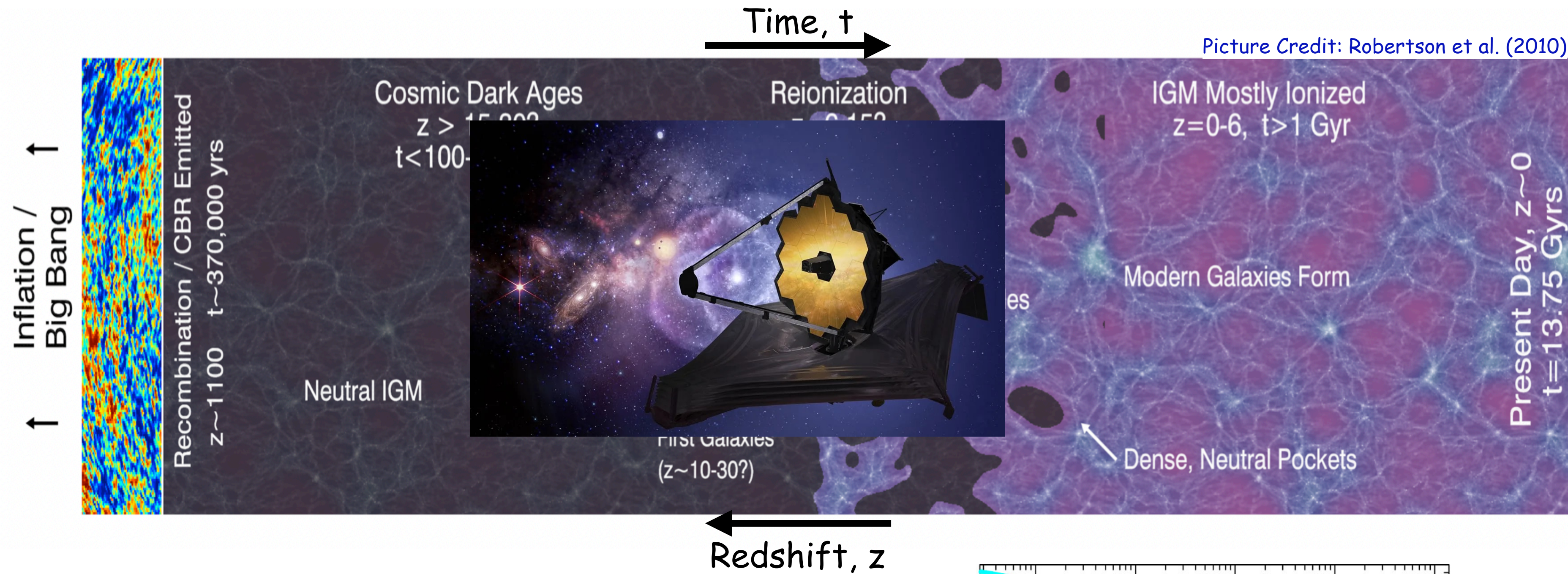
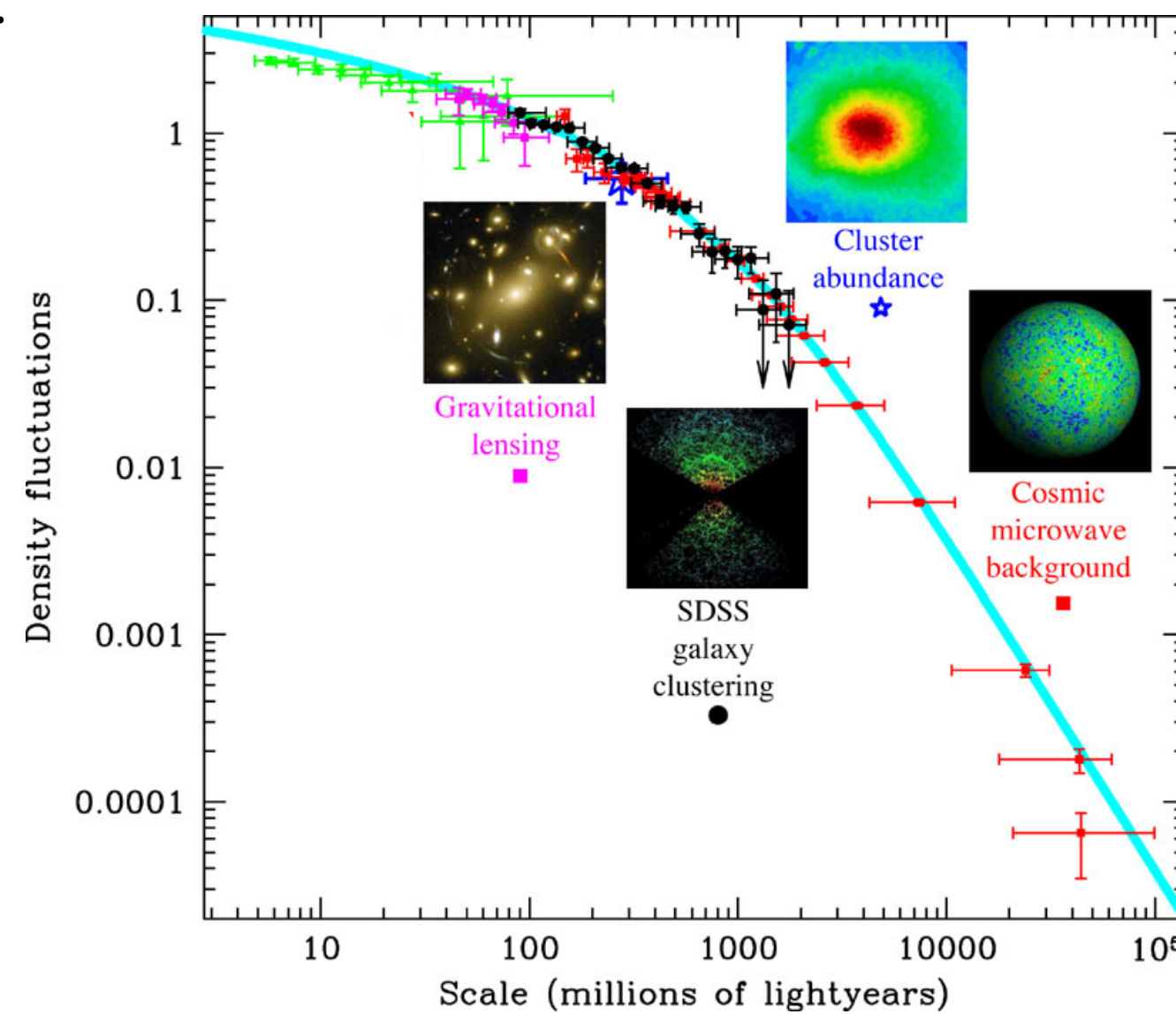
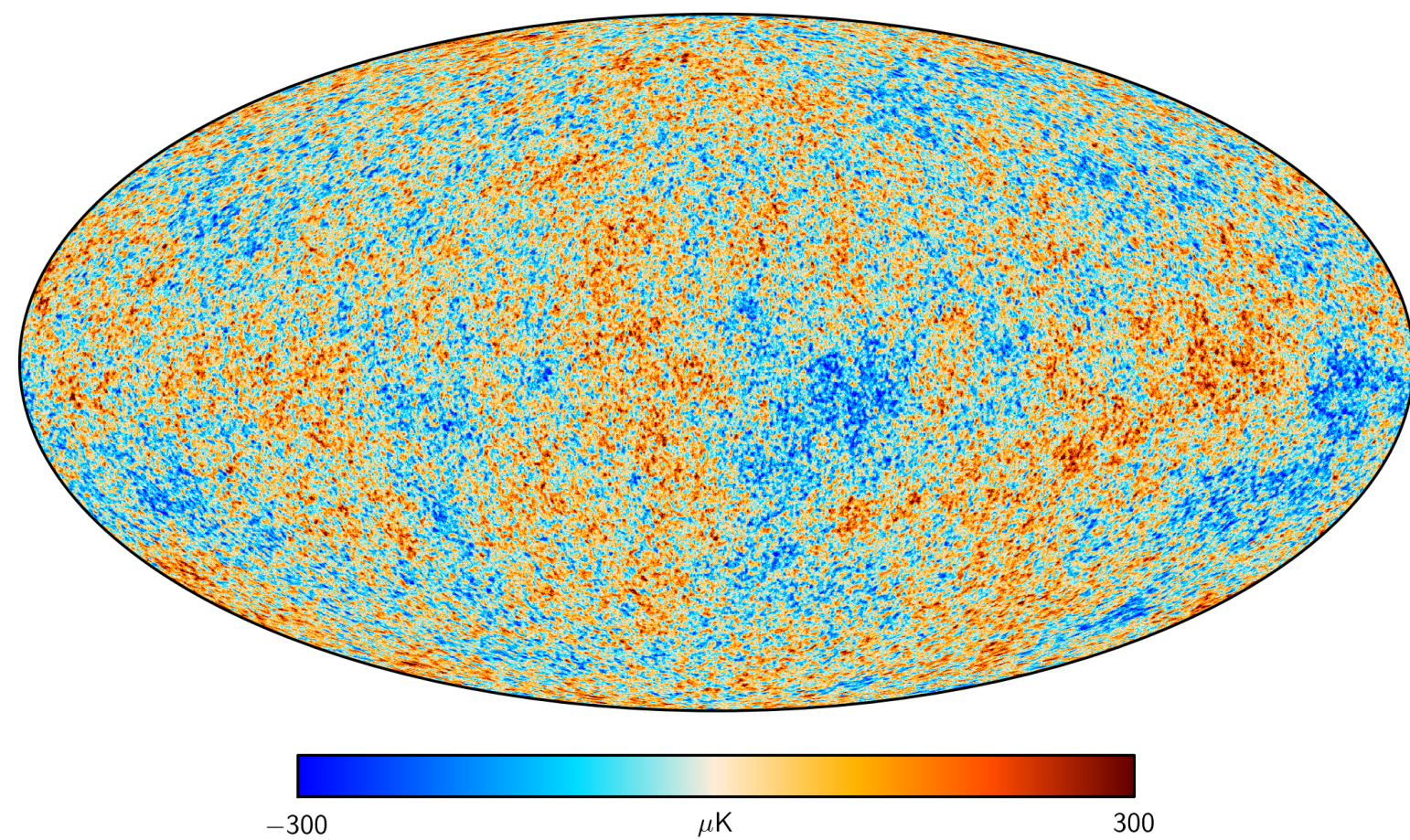


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# James Webb Space Telescope (JWST)

Launch date:  
25 December 2021

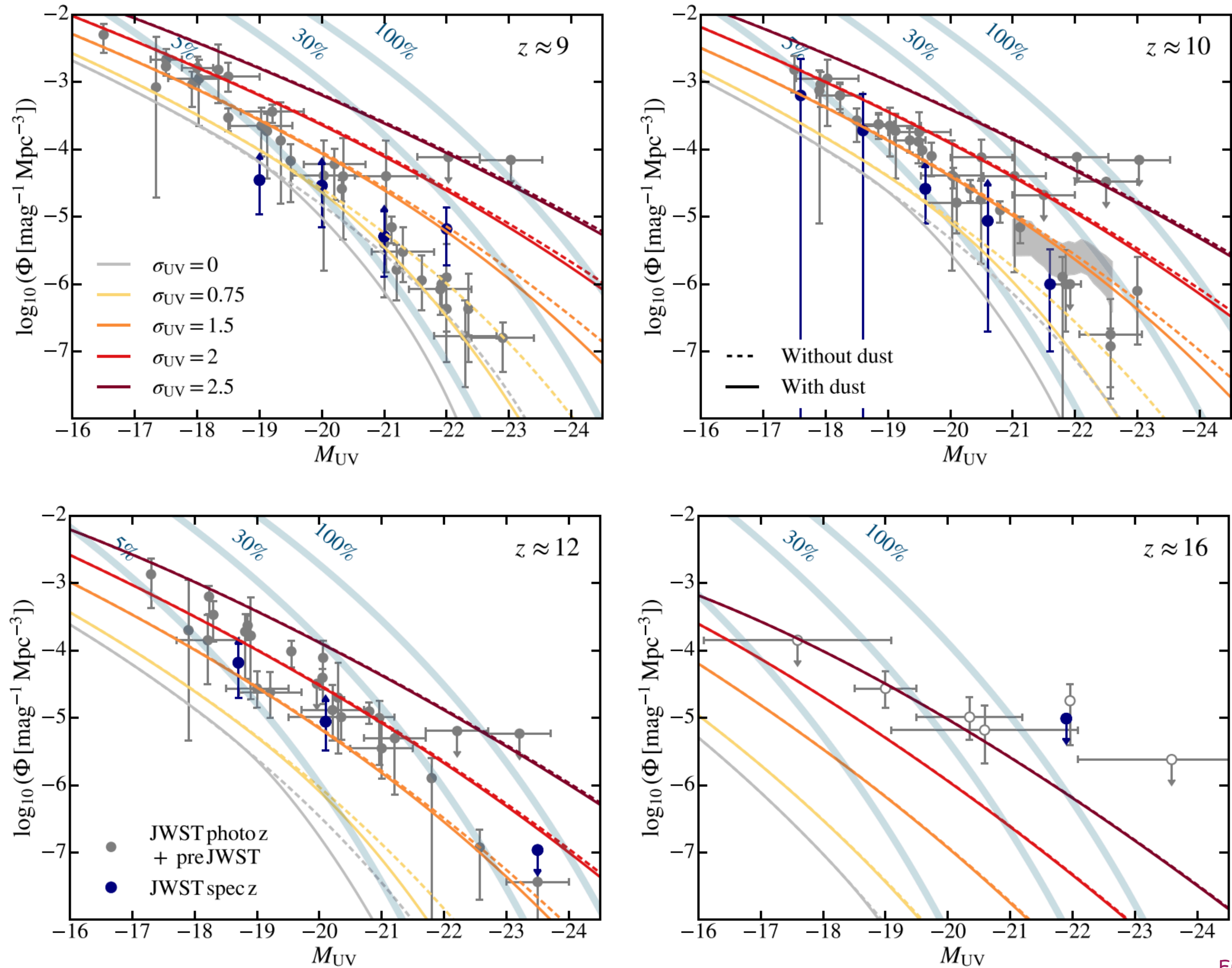
$$\lambda_{\text{obs}} = (1 + z)\lambda_{\text{emit}}$$

- Young stars emits in UV range.
- JWST can observe the high redshift objects (higher than HST).





# JWST Observations: $\Lambda$ CDM tension?



- Early data releases of JWST have revealed several high redshift galaxy candidates. (Castellano et al. 2022; Finkelstein et al. 2022; Naidu et al. 2022b; Adams et al. 2023b; Atek et al. 2023; Bouwens et al. 2023a; Donnan et al. 2023; Harikane et al. 2023b; Robertson et al. 2023; Yan et al. 2023))
- A large population of ultra-violet (UV)-bright galaxies at  $z \approx 10$ . (Finkelstein et al. 2023; Harikane et al. 2023b).
- Possible tension with standard galaxy formation models??

Most of them are observed by photometry and many of them have been confirmed spectroscopically.

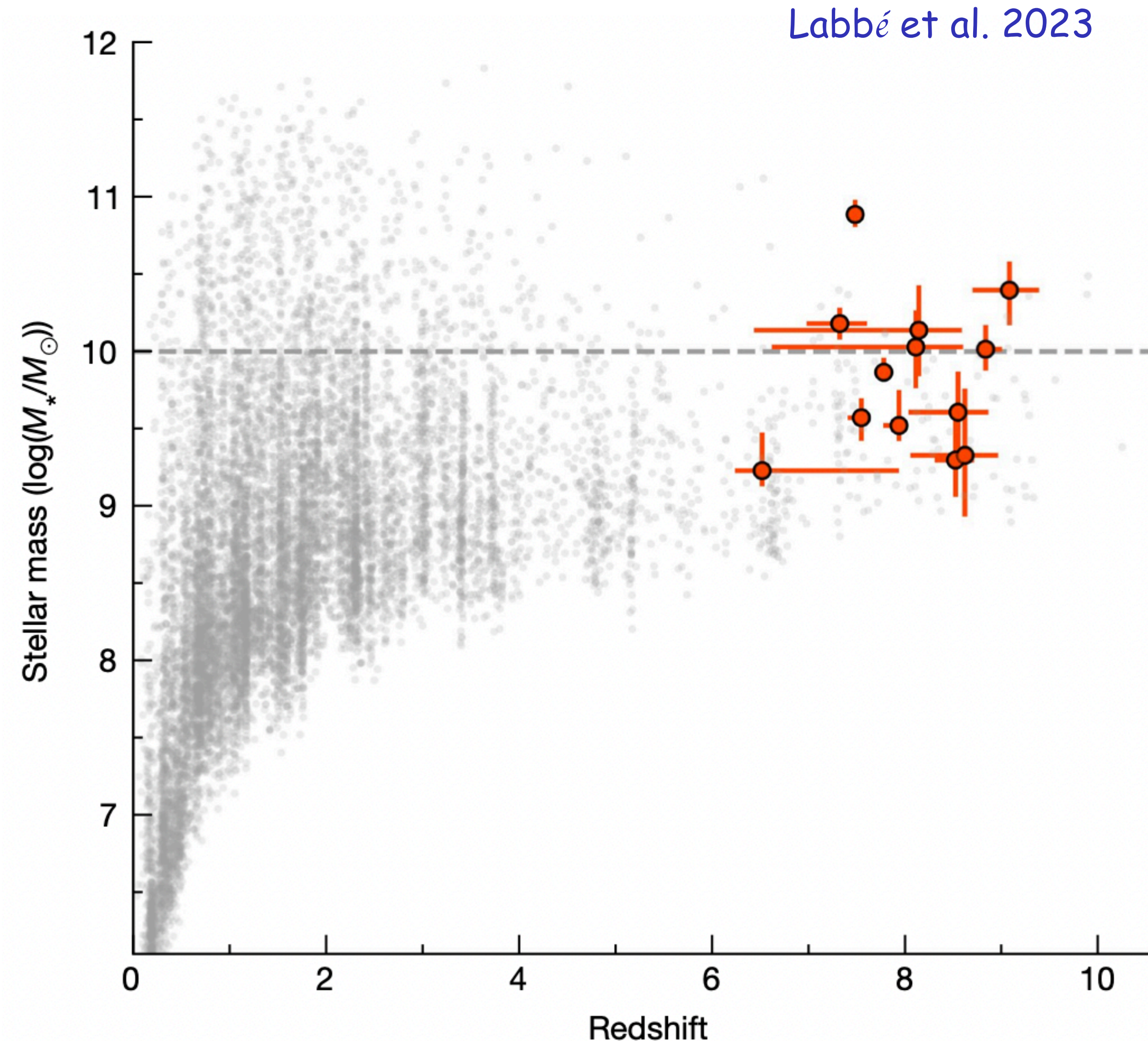


# JWST Observations: $\Lambda$ CDM tension?

Early data releases of JWST have revealed several high redshift surprisingly massive galaxy candidates.

- [Labbé et al. 2023](#) have found 13 galaxy candidates ( $6.5 \leq z \leq 9.1$ ) using JWST data released within the Cosmic Evolution Early Release Science (CEERS) program.
- Six candidates found to have stellar mass  $M_* > 10^{10} M_\odot$ .

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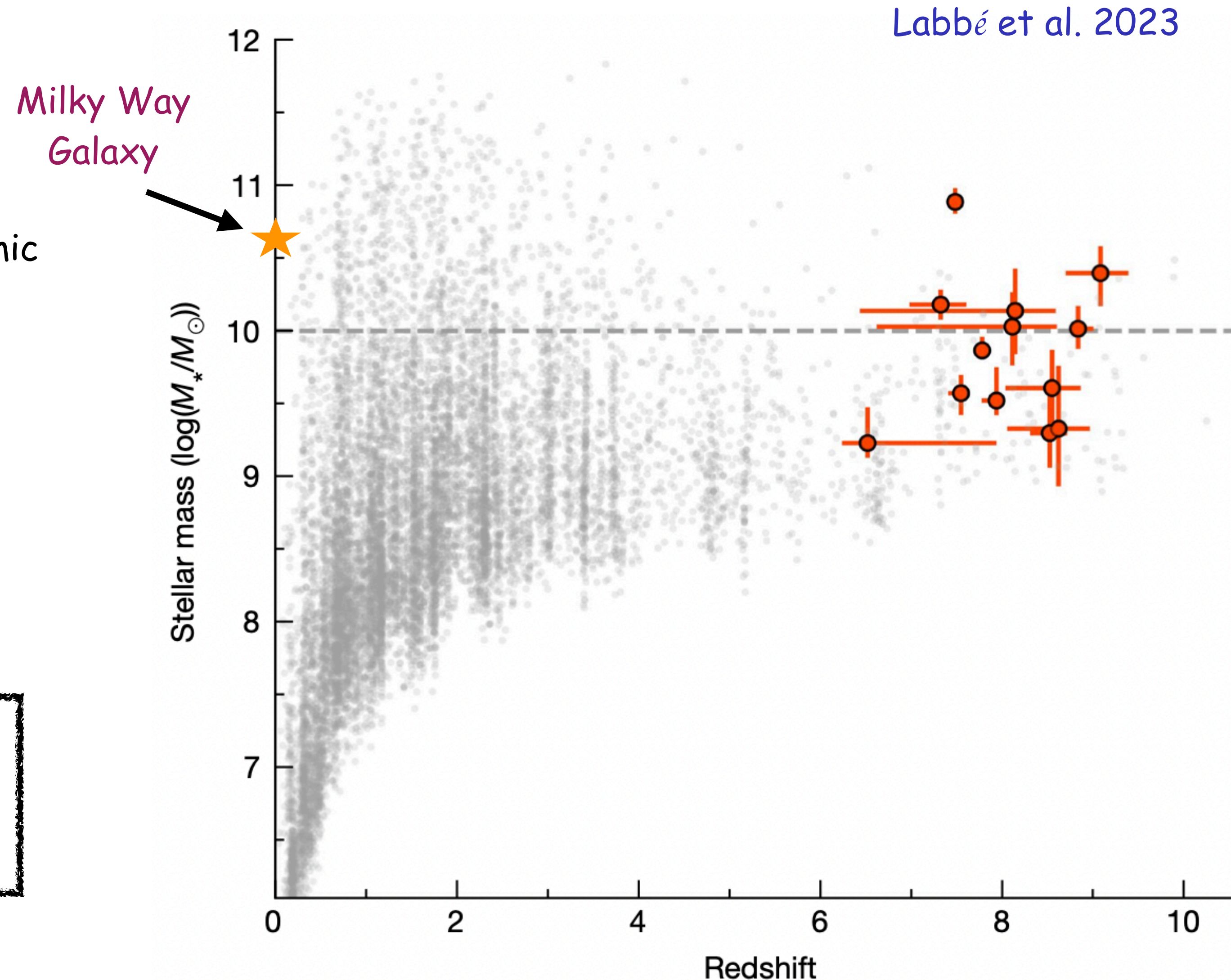


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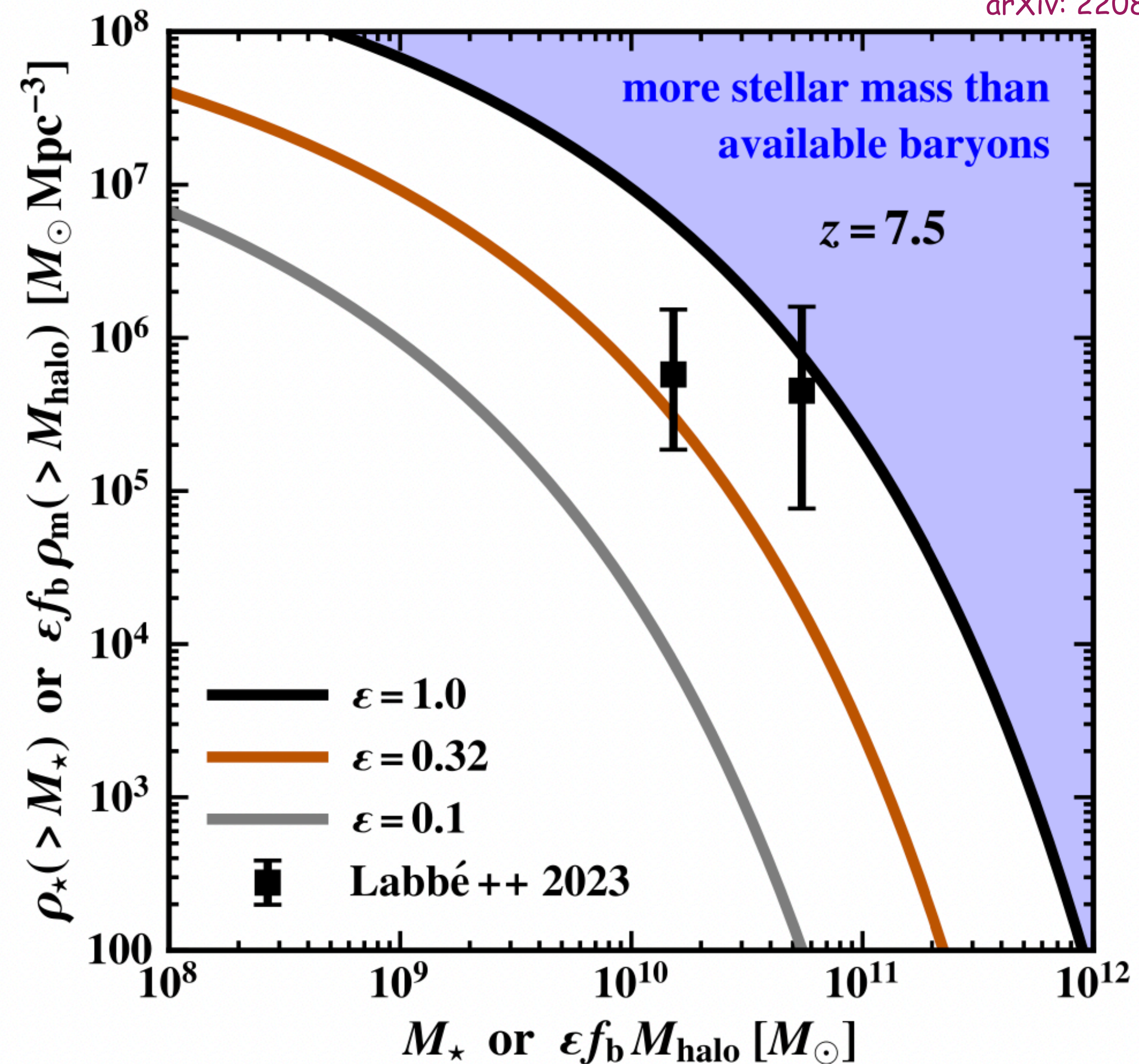
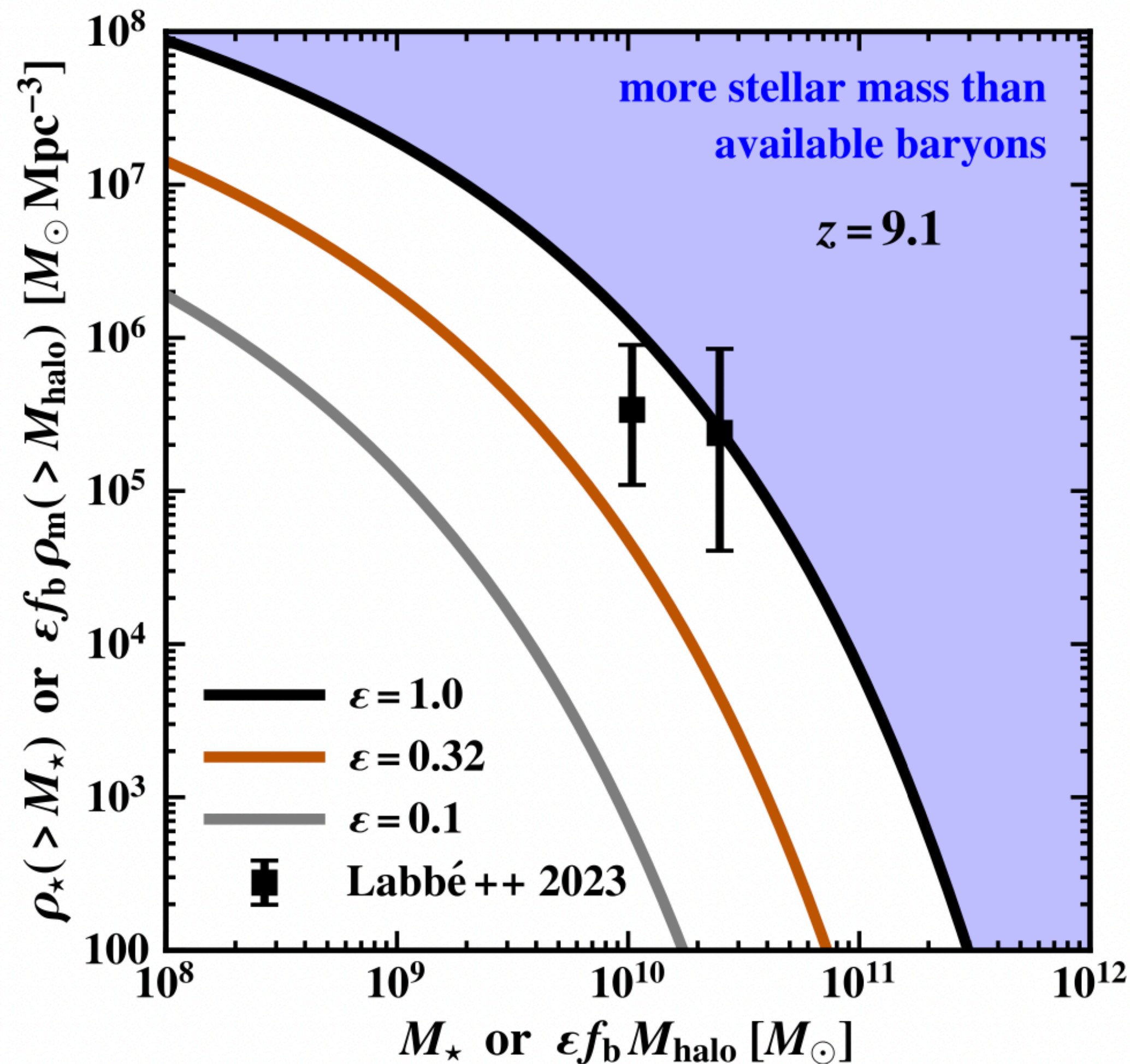


# JWST Observations: $\Lambda$ CDM tension?

These JWST observed massive galaxy candidates requires to have a very high star formation efficiency within standard cosmological model.

Boylan-Kolchin M., 2023, *Nature Astronomy*,

arXiv: 2208.01611



Actively debated in the literature and subject to systematic uncertainties (e.g. Larson et al. 2022; Chen, Mo & Wang 2023; Endsley et al. 2023; Prada et al. 2023; Steinhardt et al. 2023).



# $\Lambda$ CDM tension? $\longrightarrow$ Hint for new physics?

**Implications:** Either the inferred galaxy properties are wrong (**systematic or requires modification on the astrophysical side**) or there is an issue with our successful **standard cosmological model**



- Various solutions involving modification on the astrophysical side or **beyond standard cosmological model** have been explored.

Early dark energy component (Shen et al. 2023; Boylan-Kolchin 2023), presence of primordial black holes or axion miniclusters (Liu & Bromm 2022; Hütsi et al. 2023; Yuan et al. 2023; Dolgov 2023), fuzzy dark matter & warm dark matter (Gong et al. 2023; Bin Liu et al. 2024), primordial non-Gaussianity (Biagetti et al. 2023), cosmic strings (Jiao et al. 2023).



# Halo Mass Function

**Semi-analytical method:** We utilise the extended Press-Schechter formalism to compute the statistics of non-linear density field from the linear power spectrum.

Computed using HMF code

The halo mass function is defined as the number density of DM haloes per unit mass:

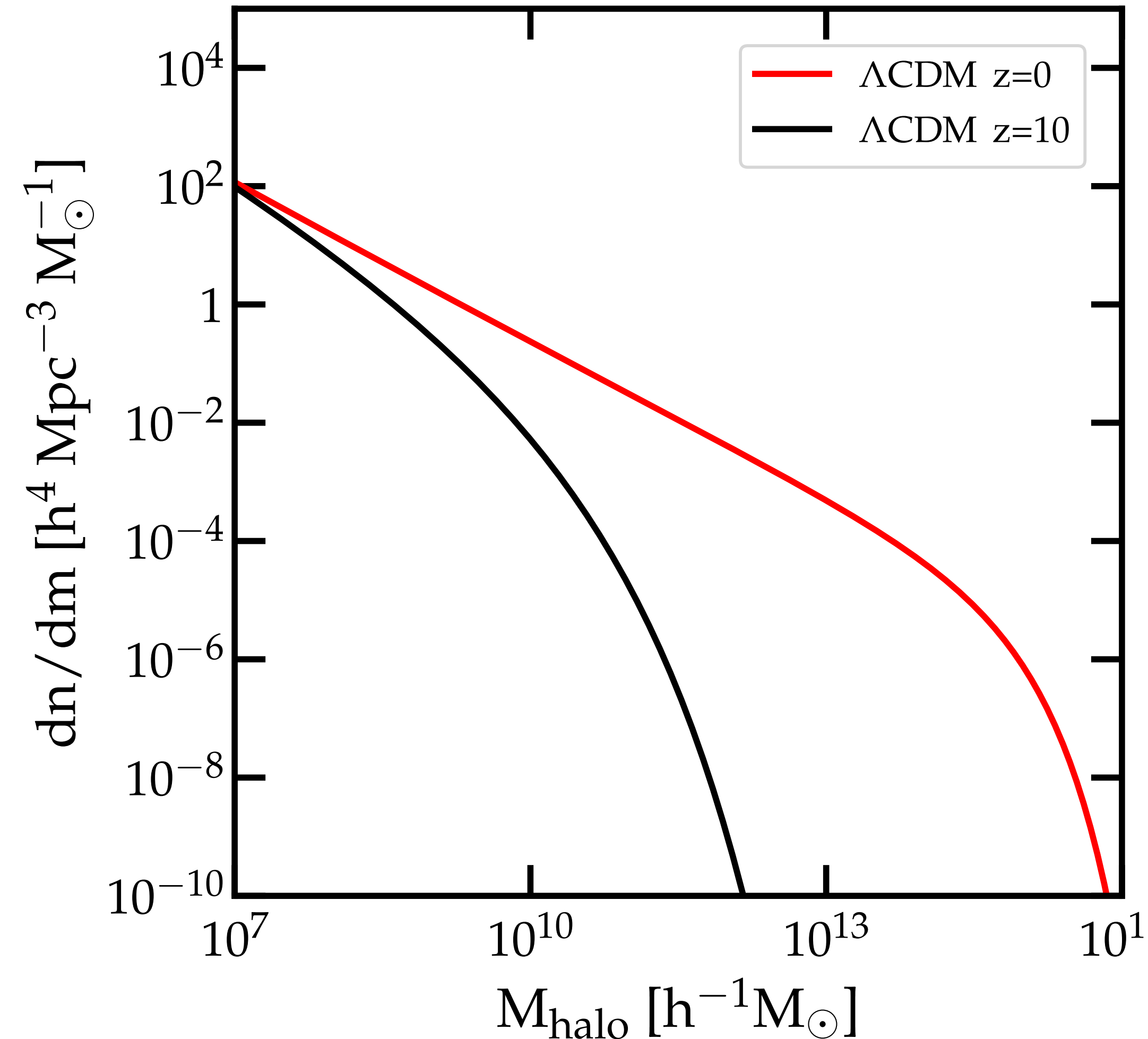
$$\frac{dn}{d \ln M} = M \frac{\rho_0}{M^2} f(\sigma) \left| \frac{d \ln \sigma}{d \ln M} \right|$$

$$\text{Where } \sigma^2(R) = \frac{1}{2\pi^2} \int_0^\infty k^2 P(k) W^2(kR) dk \quad \text{and} \quad M = \frac{4\pi\rho_0}{3} R^3$$

For  $f(\sigma)$ , we use the Sheth-Tormen fitting function:

$$f(\sigma) = A \sqrt{\frac{2a}{\pi}} \left[ 1 + \left( \frac{\sigma^2}{a\delta_c^2} \right)^p \right] \frac{\delta_c}{\sigma} \exp \left[ -\frac{a\delta_c^2}{2\sigma^2} \right]$$

$\delta_c$  = critical overdensity for collapse,  $A = 0.3222$ ,  $a = 0.707$ , and  $p = 0.3$





# Important quantities: Cumulative Comoving Number and Mass Densities & UV luminosity function

The cumulative comoving galaxy number density with stellar masses above some threshold

$M_*$  as

$$n_*( > M_*, z) = \int_{M_{\text{halo}}}^{\infty} dM \frac{dn(M, z)}{dM},$$

where  $M_{\text{halo}} = \frac{M_*}{\epsilon f_b}$  and  $f_b = \Omega_b / \Omega_m$

and the corresponding cumulative comoving stellar mass density

$$\rho_*( > M_*, z) = \epsilon f_b \int_{M_{\text{halo}}}^{\infty} dM M \frac{dn(M, z)}{dM}.$$

- $\epsilon$  is the star formation efficiency, and satisfies  $\epsilon \leq 1$
- Exact value depends on star formation physics

UV Luminosity Function:

$$\Phi_{\text{UV}} = \frac{dn}{dM} \frac{dM}{dM_{\text{UV}}}$$

To compute UV Luminosity function, we need to know the  $M_{\text{UV}} - M_{\text{halo}}$  relation.



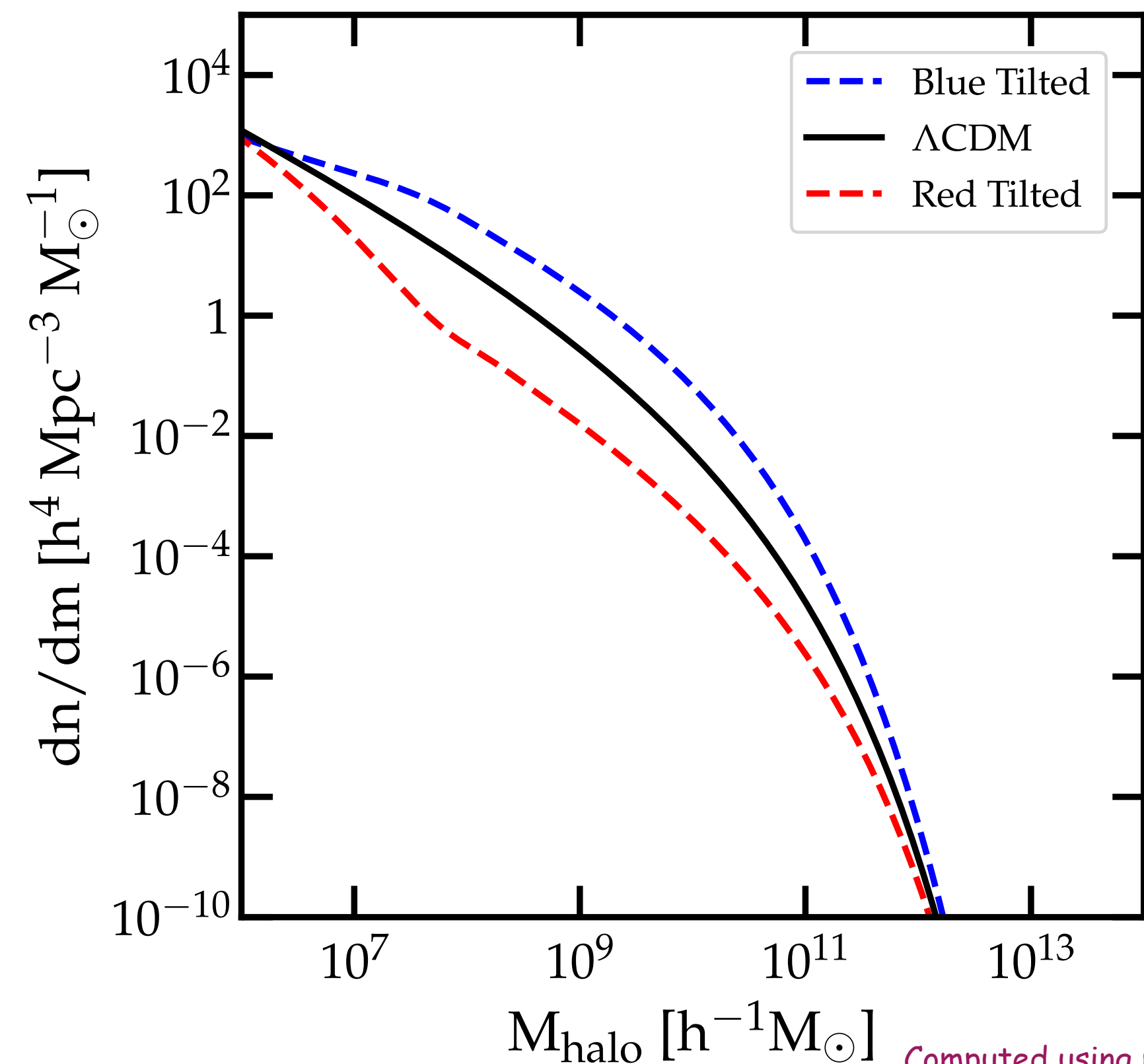
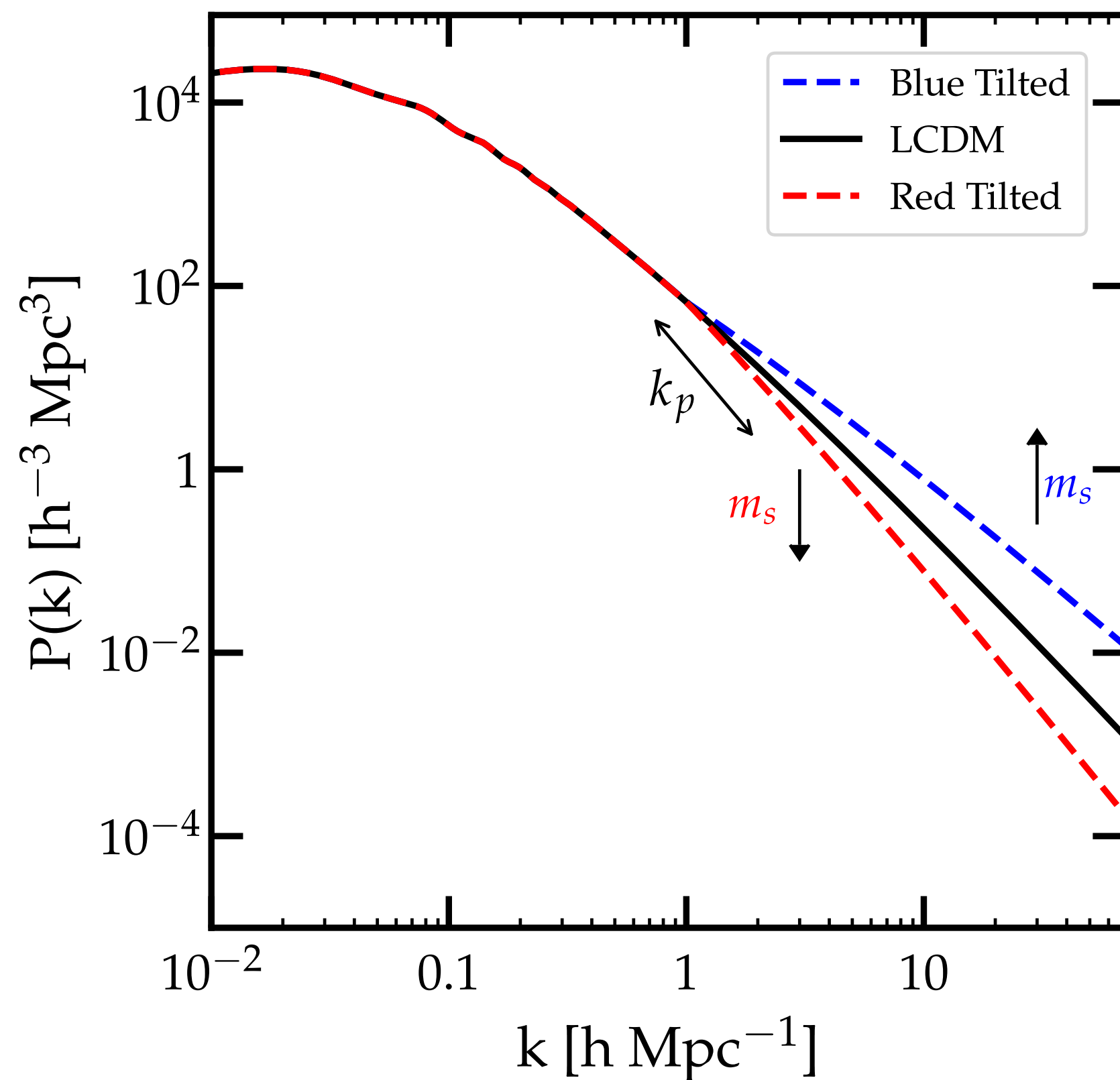
# Modified Primordial Power Spectrum

We study a modified primordial power spectrum where it deviates from the standard primordial power spectrum at small length scales with a model agnostic form:

$$P_{\text{prim}}(k) \propto k^{n_s}, \quad \text{for } k < k_p,$$
$$\propto k_p^{n_s - m_s} k^{m_s}, \quad \text{for } k > k_p.$$

P. Parashari & R. Laha, MNRAS: Letters, 526, L63-L69 (2023) arXiv: 2305.00999

For  $m_s > n_s$ , the power spectrum will be blue tilted on scales  $k > k_p$ , and it is red tilted if  $m_s < n_s$ .



Computed using modified HMF code

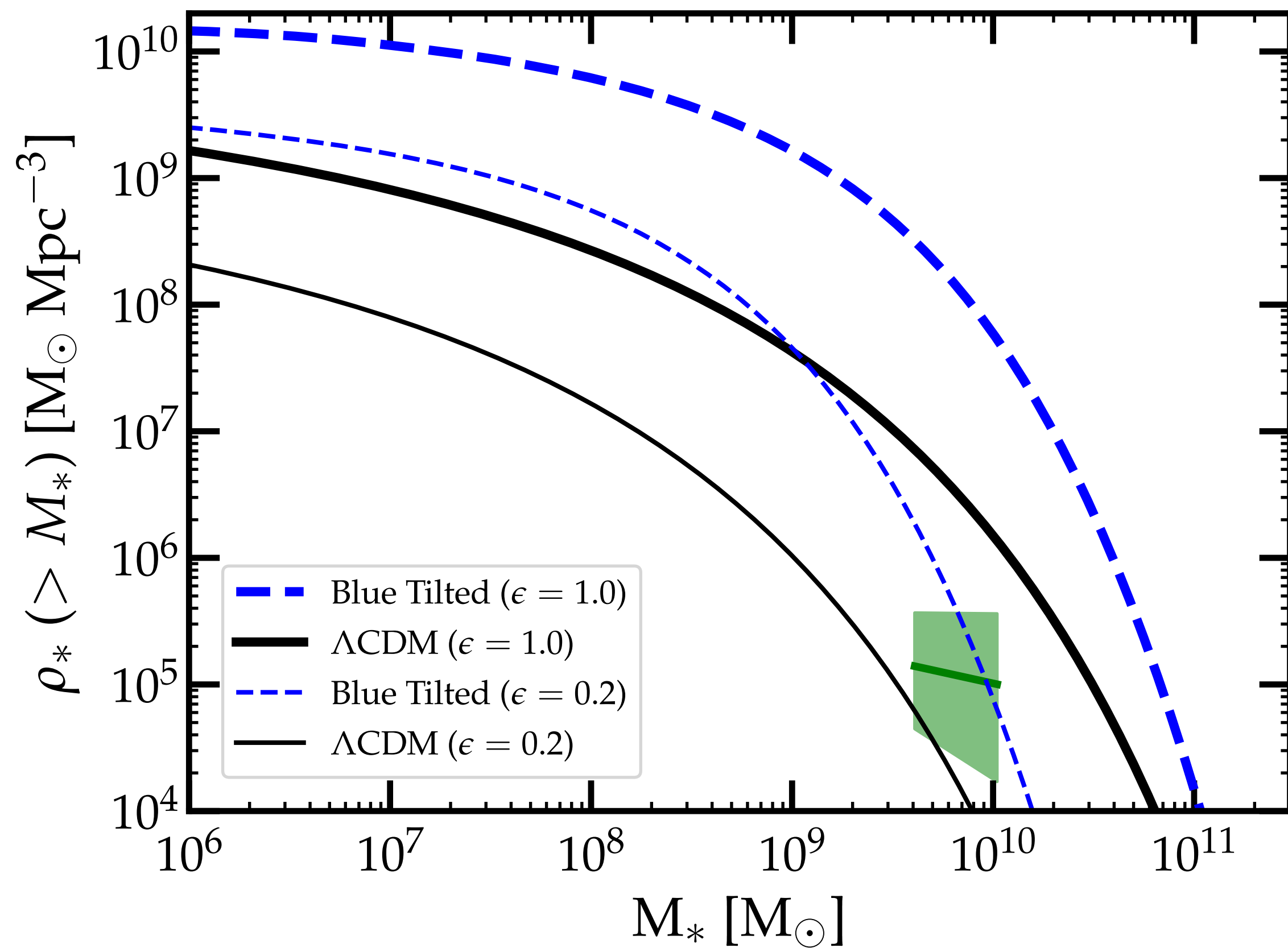


# JWST Observations and Modified Power Spectrum

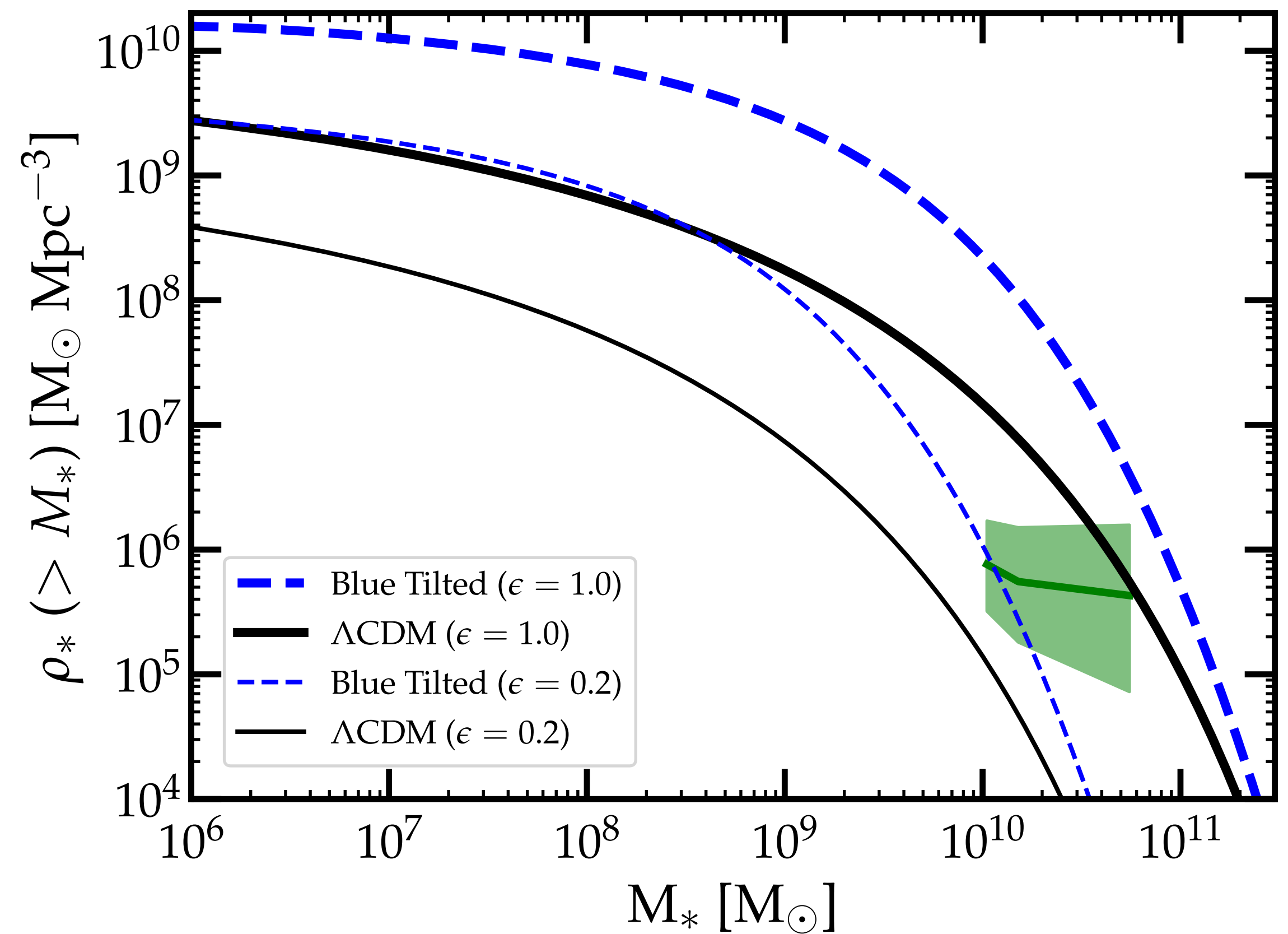
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- Blue-tilted primordial power spectrum can reduce the required star formation efficiency.

$z = 9$



$z = 7.5$

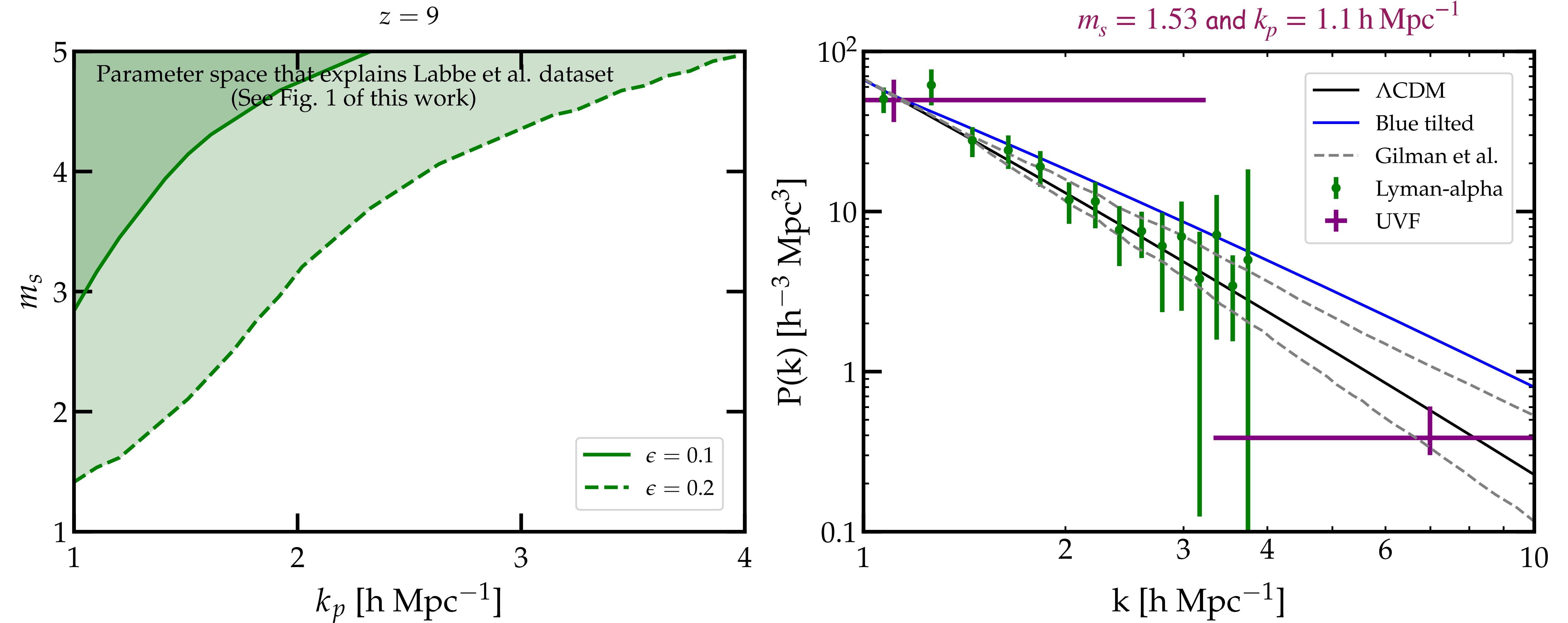


$m_s = 2.0$  and  $k_p = 1 \text{ h Mpc}^{-1}$

Computed using modified HMF code



# JWST Observations and Modified Power Spectrum



- Required parameter space to reduce this tension may be in conflict with earlier observations.



# Other Works with JWST Observations and Modified Power Spectrum

- After our work, [Hirano & Yoshida \(2023\)](#) did simulations with a blue-tilted power spectrum and **found their results consistent with our results.** ([arXiv: 2306.11993](#)). [Padmanabhan H. & Loeb A. \(2023\)](#) also found similar results [arXiv:2306.04684](#).
- In another work, [Sabti et al. \(2023\)](#) performed an analysis by assuming a Gaussian enhancement in the power spectrum and found that the enhancement required to explain [Labbé et al. \(2023\)](#) observations will conflict with previous constraints on these scales by HST, **which is consistent with our result.** ([arXiv: 2305.07049](#))



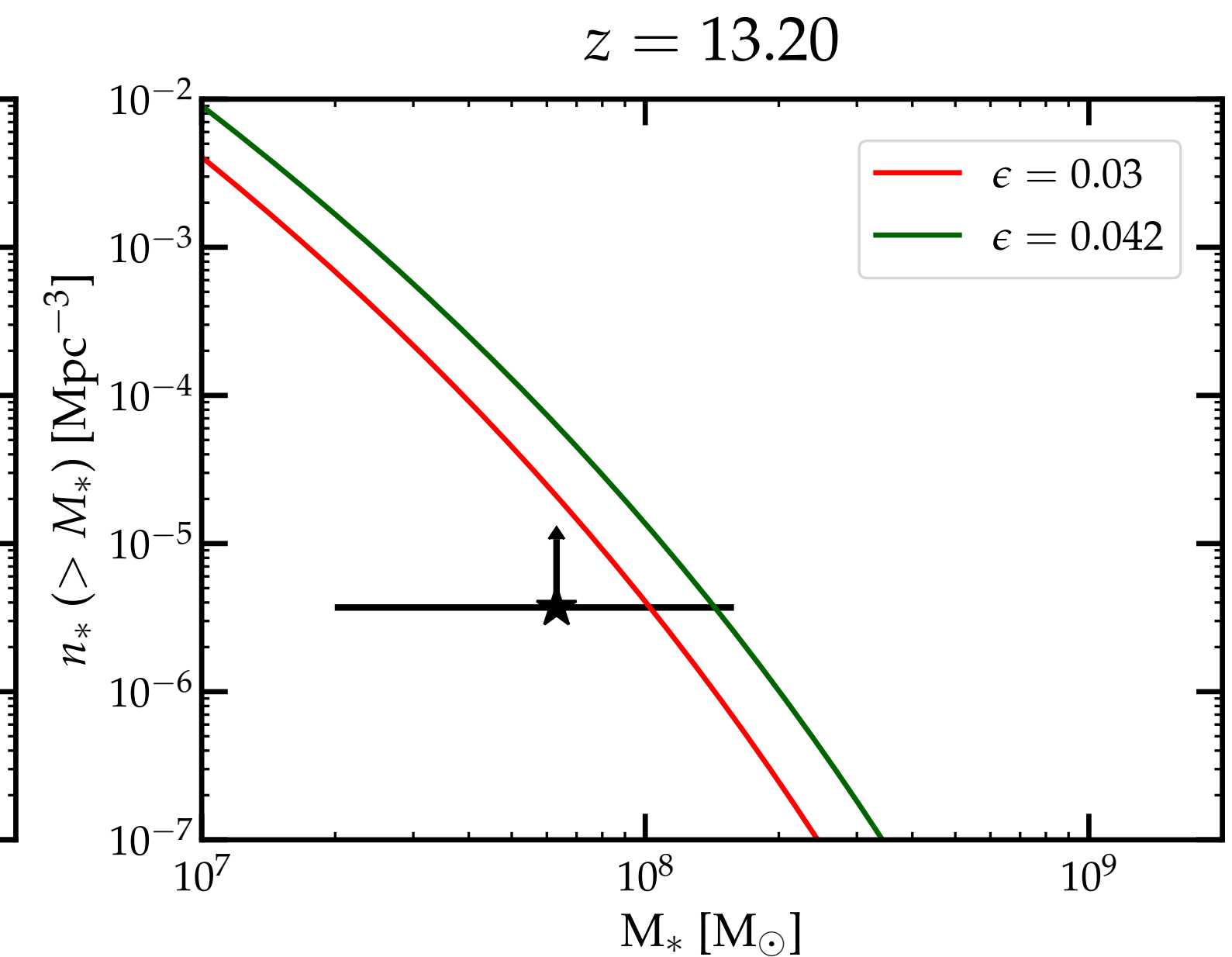
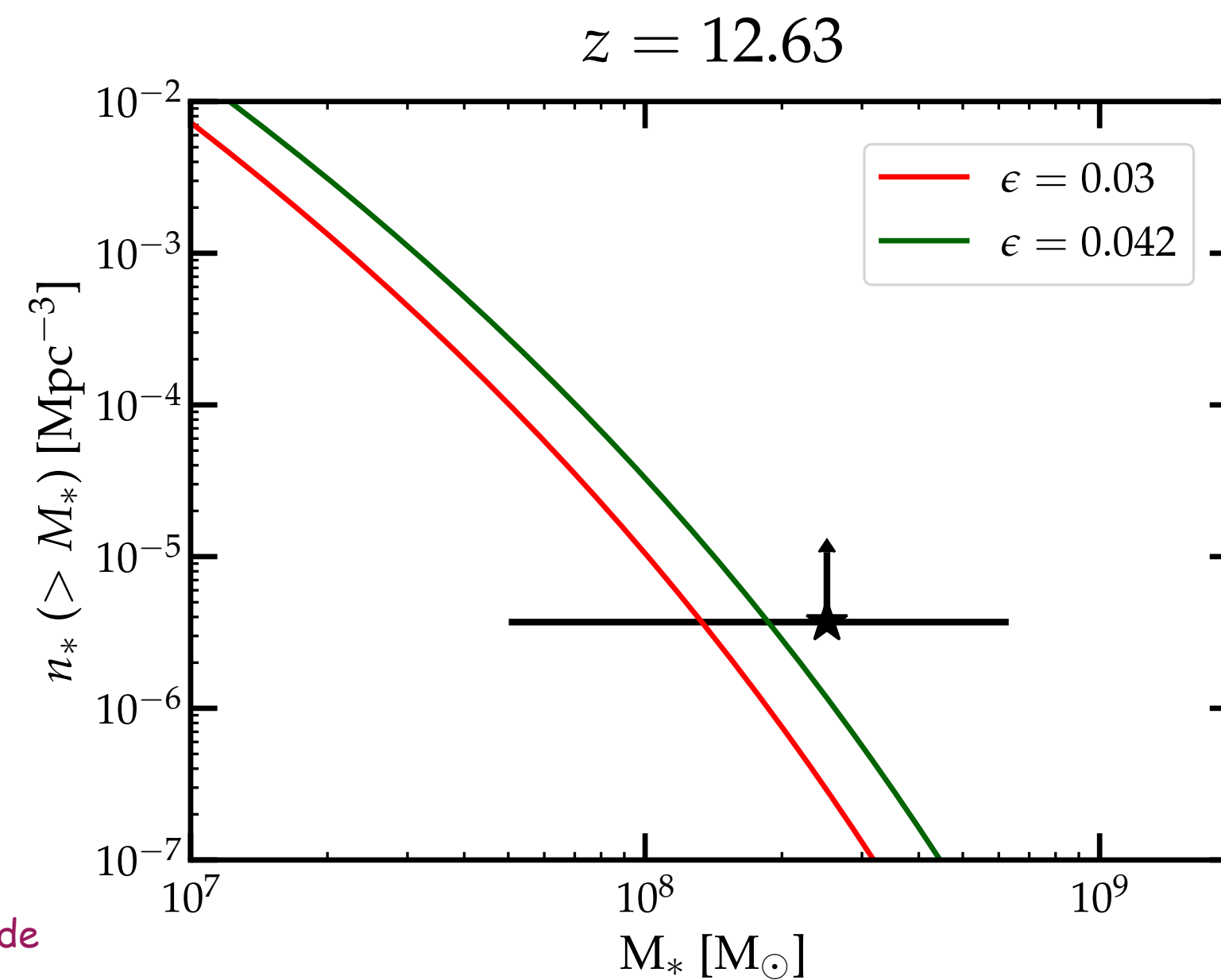
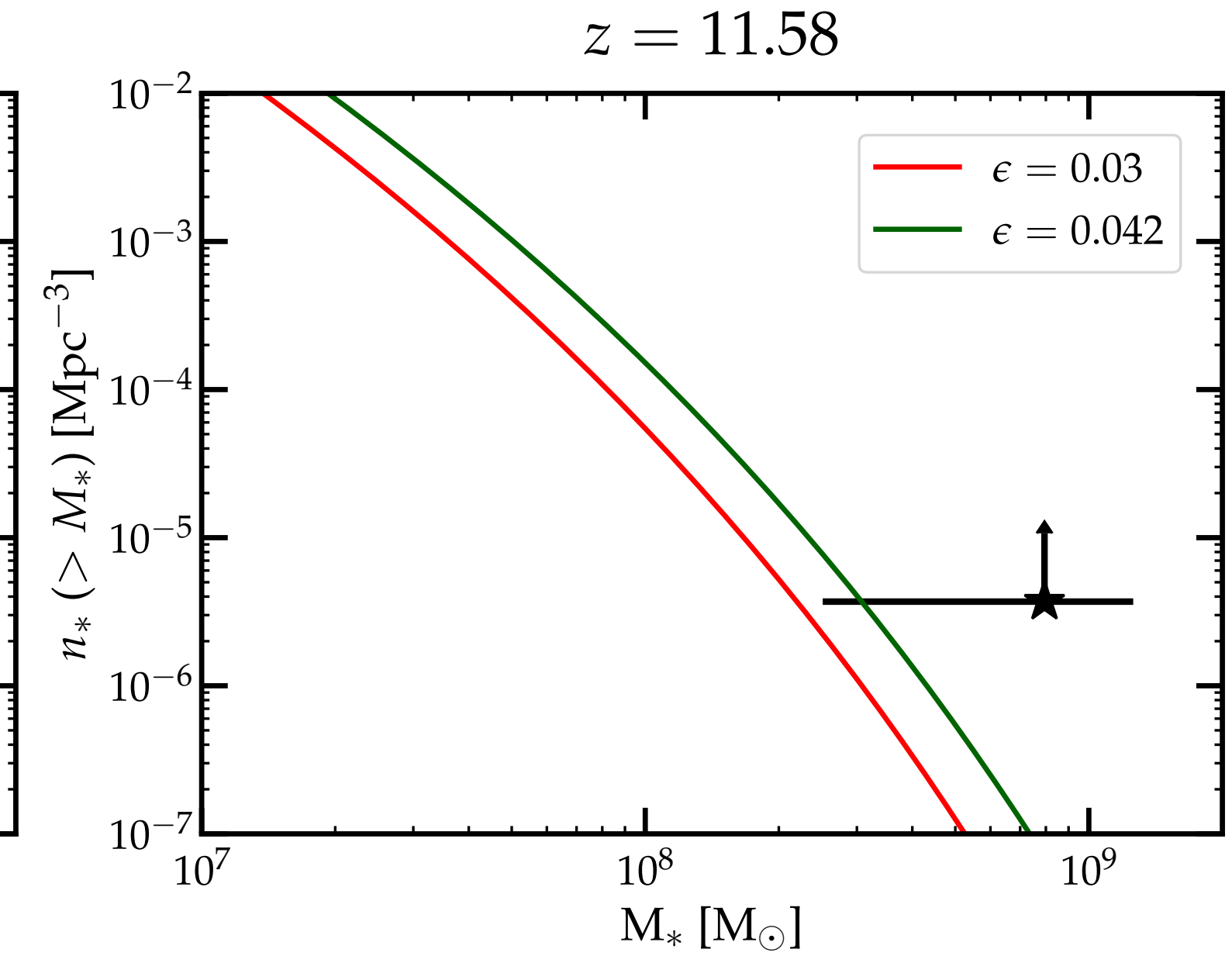
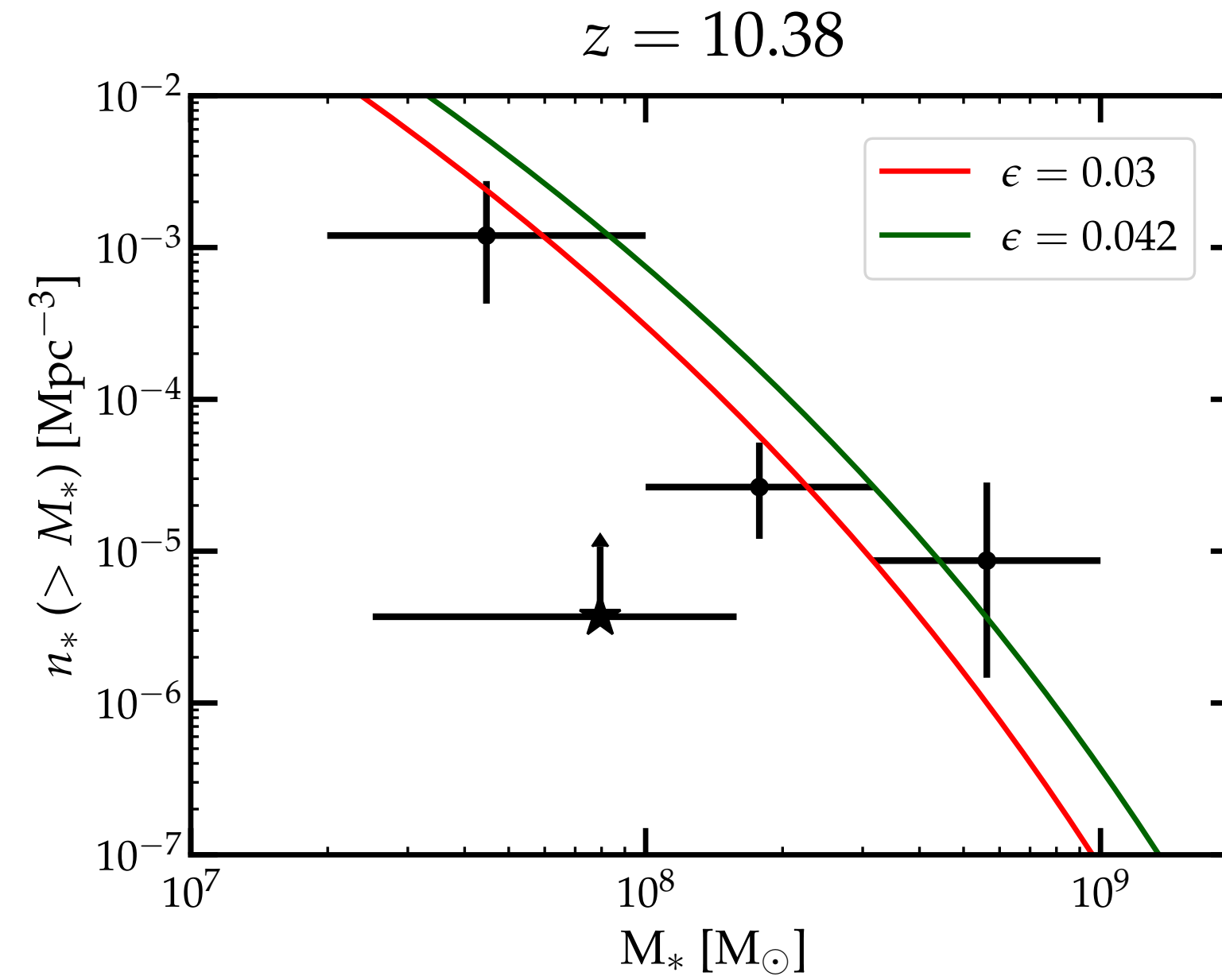
JWST observations from the JADES  
program (spectroscopically confirmed  
galaxies)



# JWST Observations (JADES)

P. Parashari & R. Laha, MNRAS: Letters, 526, L63-L69 (2023) arXiv: 2305.00999

- Recently [Curtis-Lake et al. \(2022\)](#) and [Robertson et al. \(2023\)](#) have reported 4 galaxies from the JADES survey with spectroscopically confirmed redshifts ( $z > 10$ ).
- [Keller et al. \(2023\)](#) reported the lower limit on cumulative comoving galaxy number density inferred from these observations (Black stars).

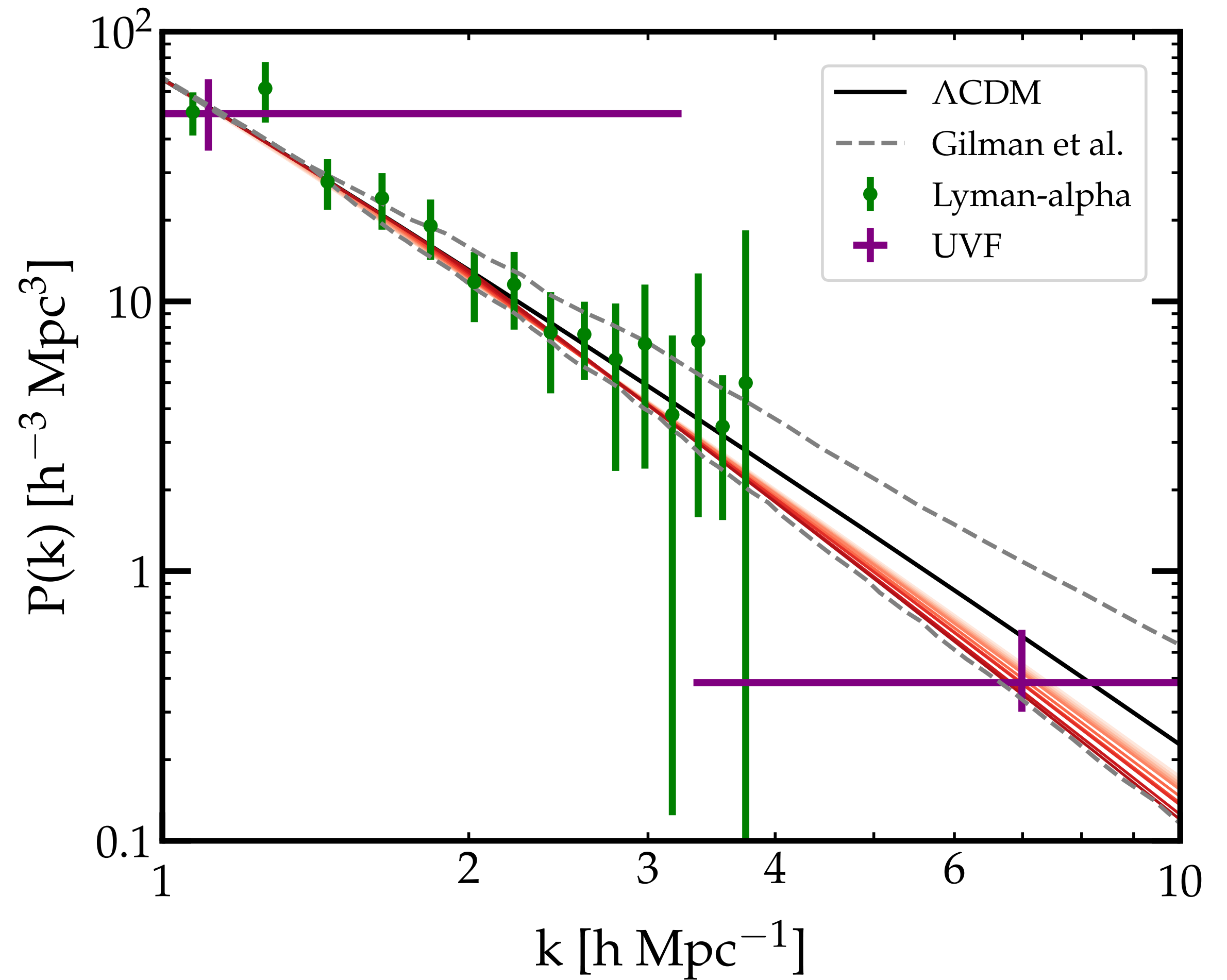
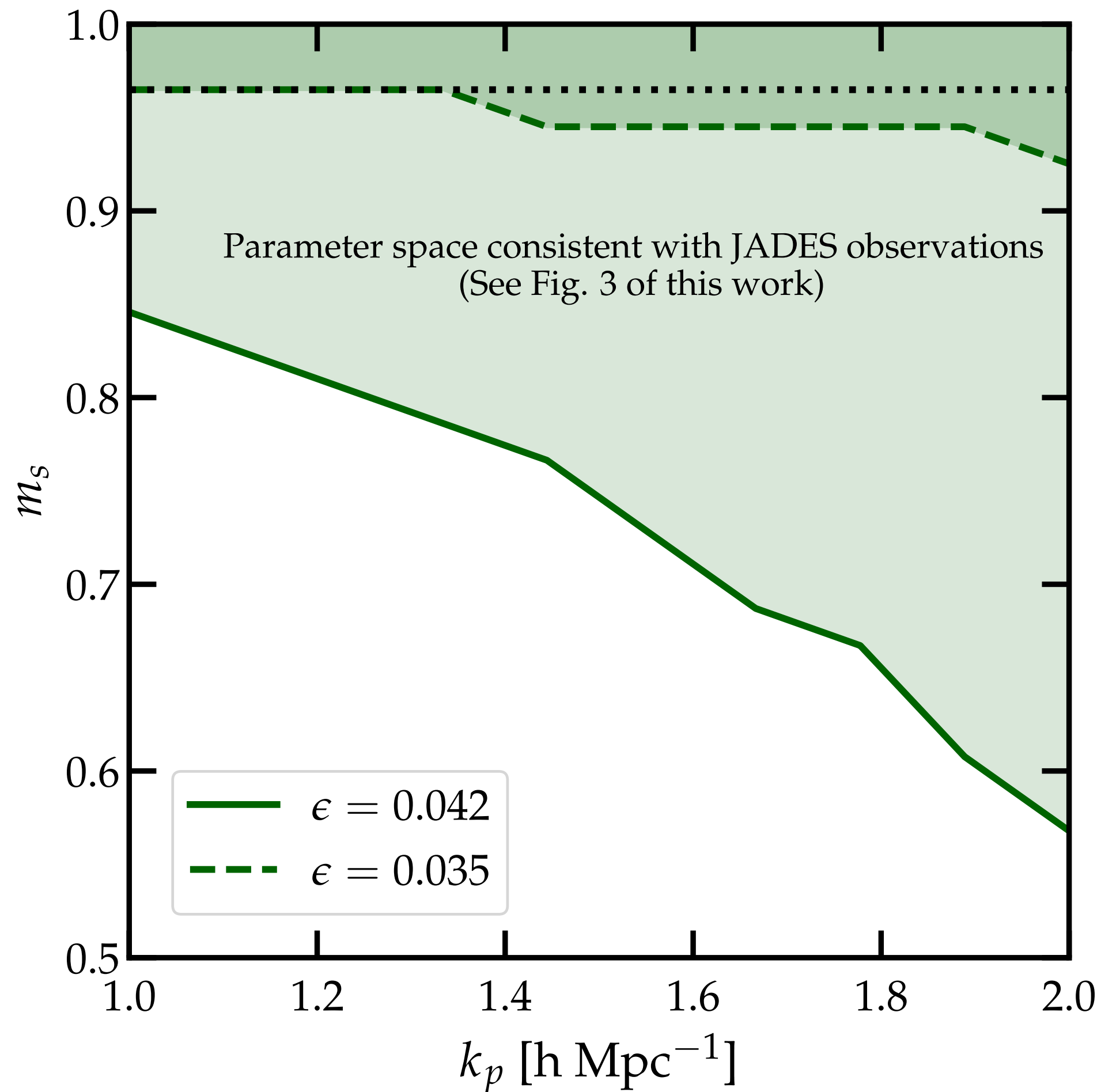


Since a red-tilted primordial power spectrum predicts a smaller cumulative comoving galaxy number density, we can use this observational data to constrain the red-tilt for a given  $\epsilon$ .



# Constraints on Power Spectrum

P. Parashari & R. Laha, MNRAS: Letters, 526, L63-L69 (2023) arXiv: 2305.00999



Computed using modified HMF code

we find the most stringent constraint on the matter power spectrum at scales  $k \sim 2 - 7 h \text{ Mpc}^{-1}$ .



# Summary

- JWST has opened up a new window to probe our Universe.
- JWST has already provided exciting and surprising results by observing several surprisingly UV-bright and massive galaxy candidates at high red shifts.
- High star formation efficiency is required to explain these galaxies within standard cosmology.
- A blue-tiled power spectrum can reduce this tension. However, the required parameter space will conflict with other observations.
- JWST has also reported a few spectroscopically confirmed galaxies. These galaxies can put the stringent constraints on matter power spectrum over scales  $k \sim 2 - 7 h \text{ Mpc}^{-1}$ .
- We show the significance of JWST observations as a potential power spectrum probe.



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Thank you!



# Outline

- Introduction to standard cosmological model
- Matter power spectrum and halo mass function
- JWST observations of high redshift galaxies
- Modified primordial power spectrum and JWST observations
- Summary



# Standard Cosmological Model

## Standard Cosmology: Inflation + $\Lambda$ CDM

- Six parameters model ( $\omega_b, \omega_c, \tau_{\text{reio}}, \theta, A_s, n_s$ )
- Fits the cosmological observations
- Inflation sets the initial condition for the structure formation.
- The minimal single field inflation models predict a scale invariant primordial power spectrum.

$$P_{\text{prim}} = A_s \left( \frac{k}{k_*} \right)^{n_s - 1}$$

Parameters related to inflation

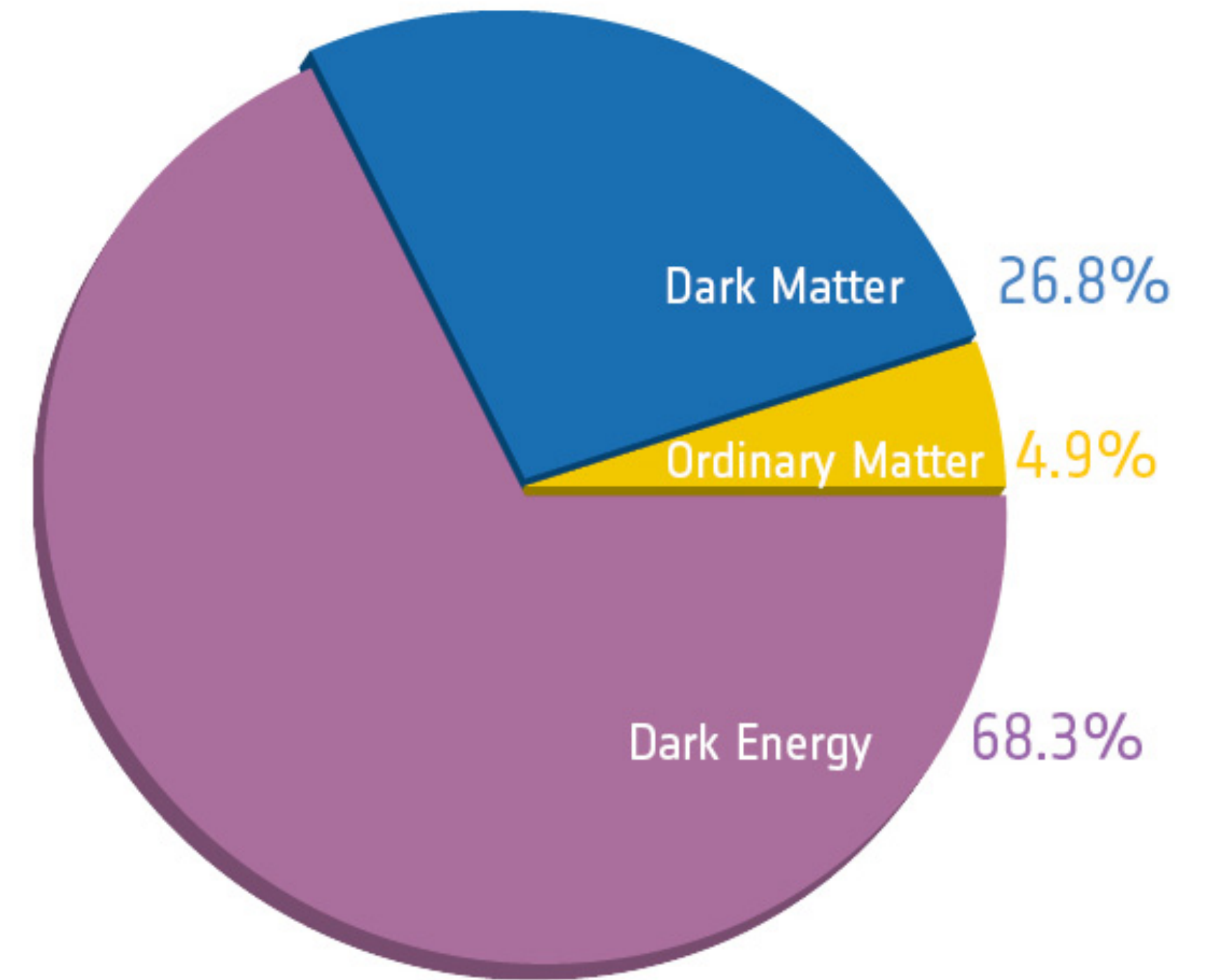


Figure credit: Planck collaboration



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Planck CMB 2018 constraints:

$$n_s = 0.9649 \pm 0.0042$$

$$A_s = (2.099 \pm 0.102) \times 10^{-9}$$

Parameters related to inflation

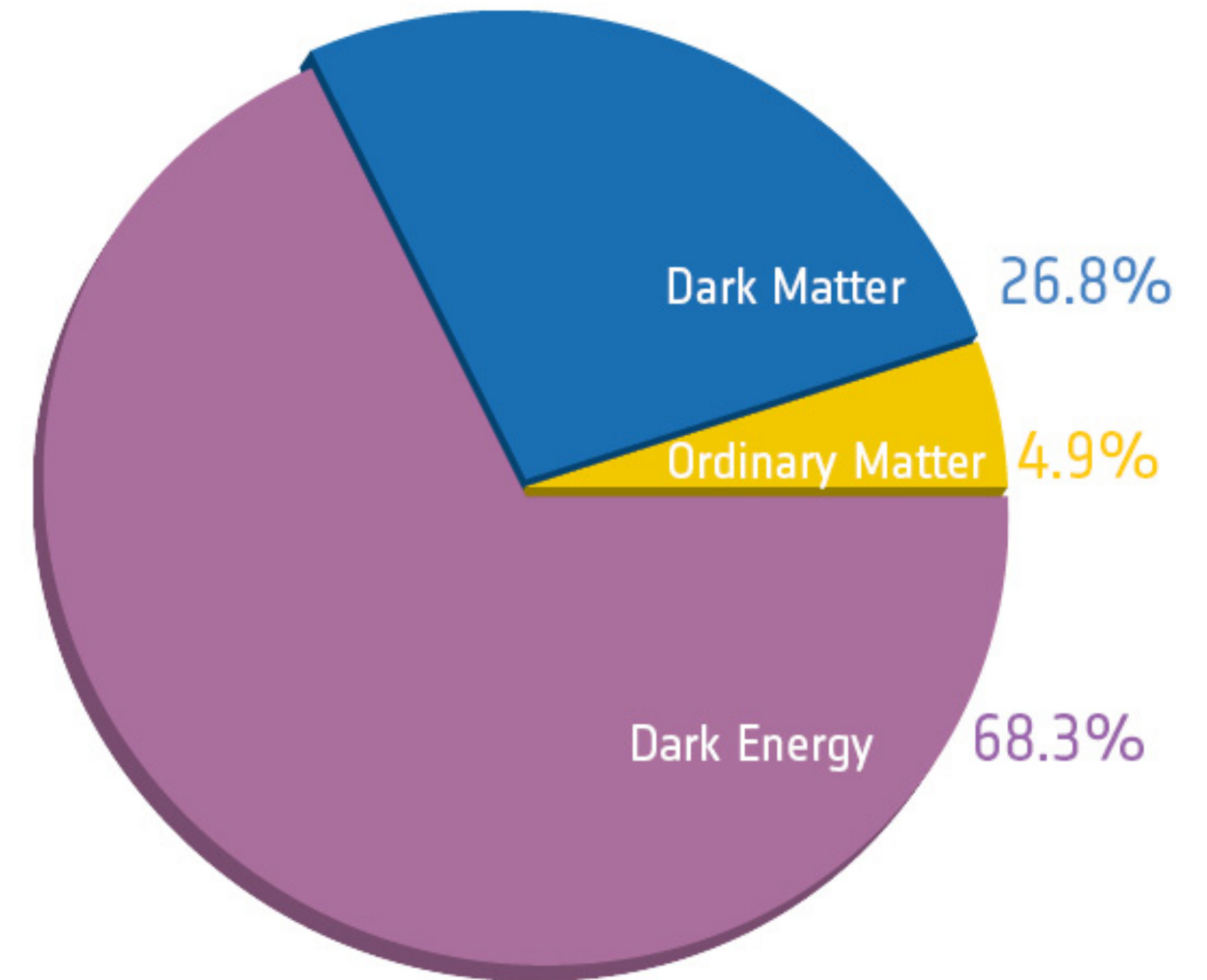


Figure credit: Planck collaboration

- CMB does not probe all the scales.
- Other observations at different scales is needed to probe power spectrum at all scales.



# Matter Power Spectrum

- Tiny fluctuations generated during inflation evolves and form the structures that we observe at present.
- The matter power spectrum: defined as the two point correlation function of the density perturbations.

$$P(k) = P_{\text{prim}}(k)T^2(k)$$

where  $P_{\text{prim}}(k)$  is the primordial power spectrum (depends on inflation model) and  $T(k)$  is the transfer function.

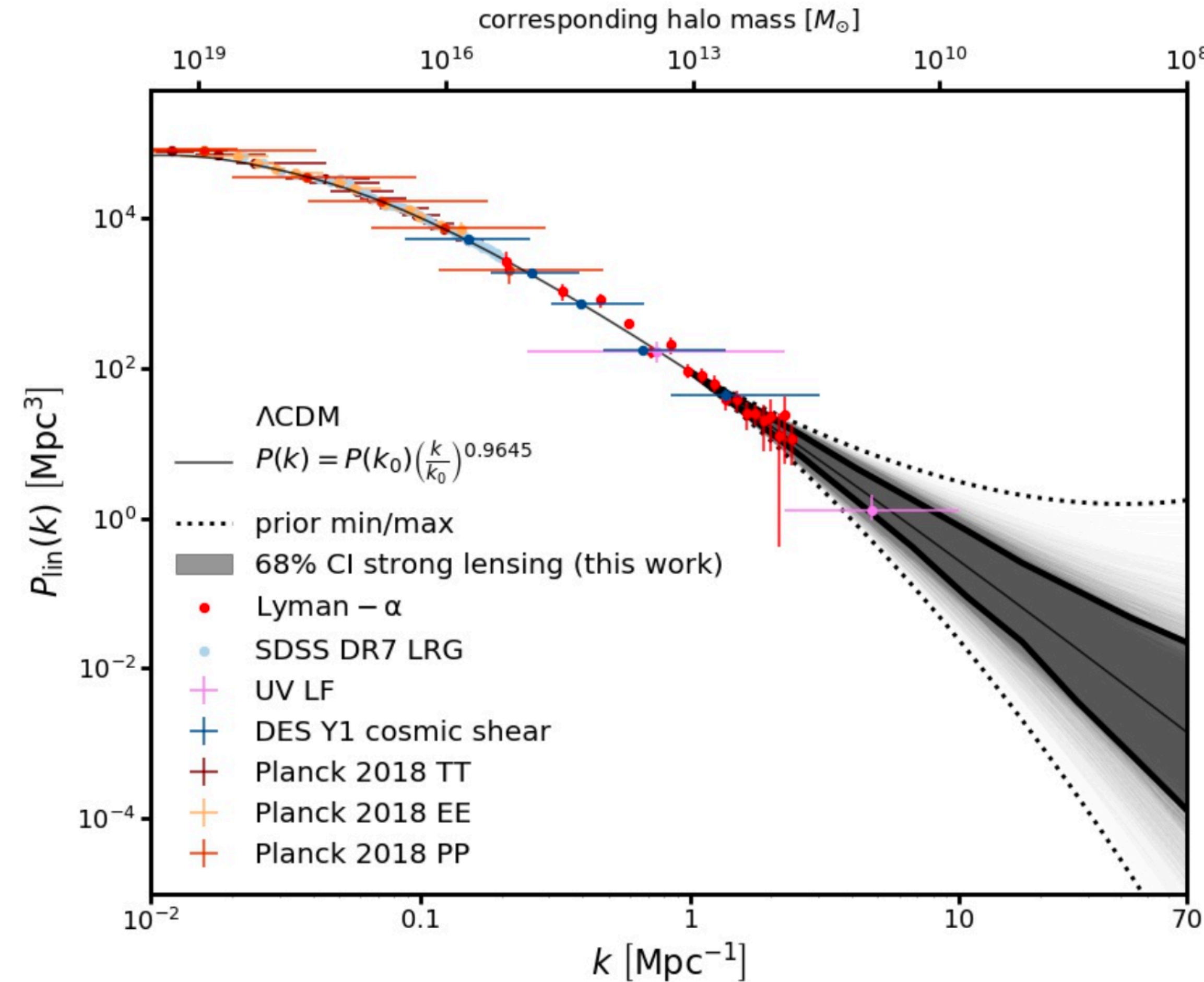
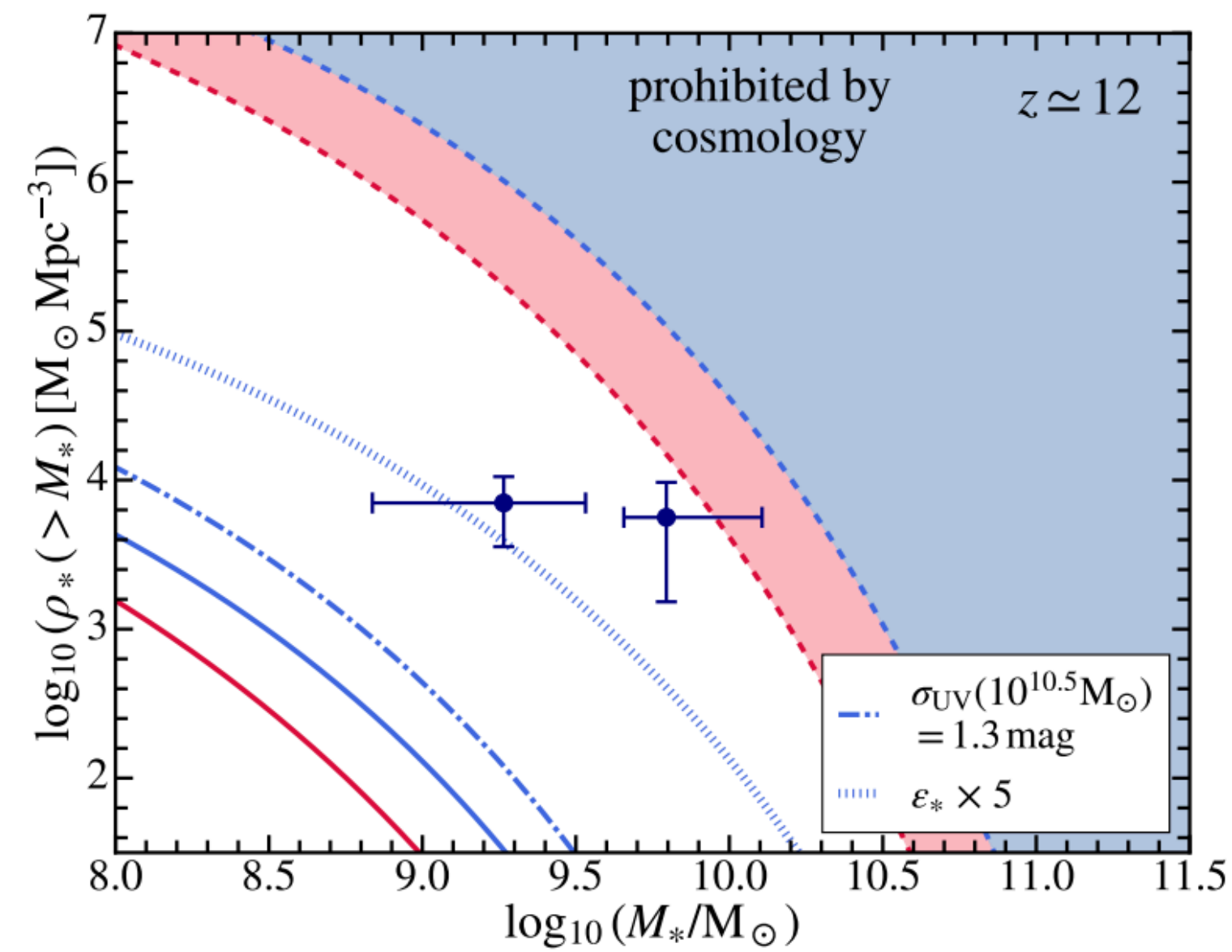
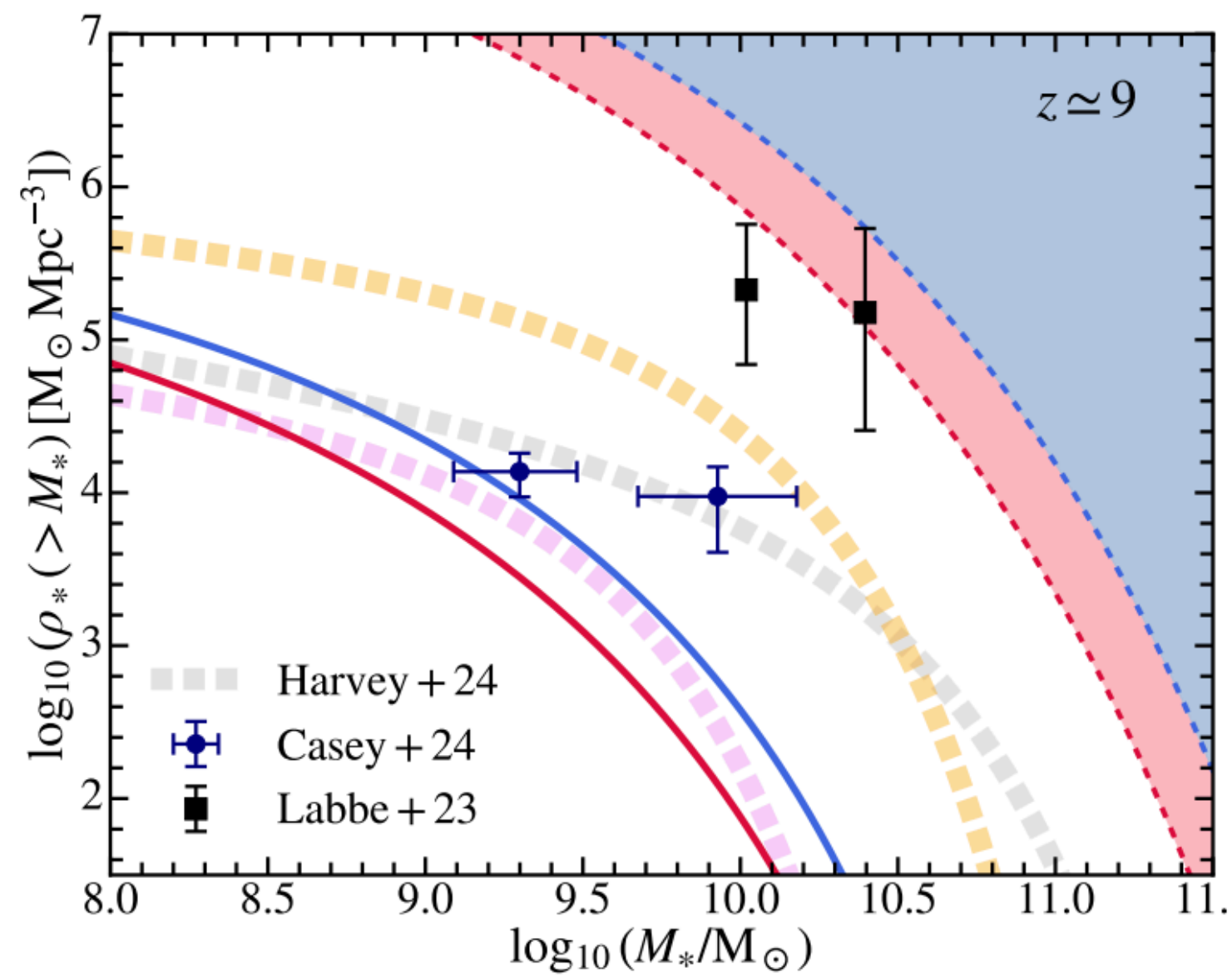
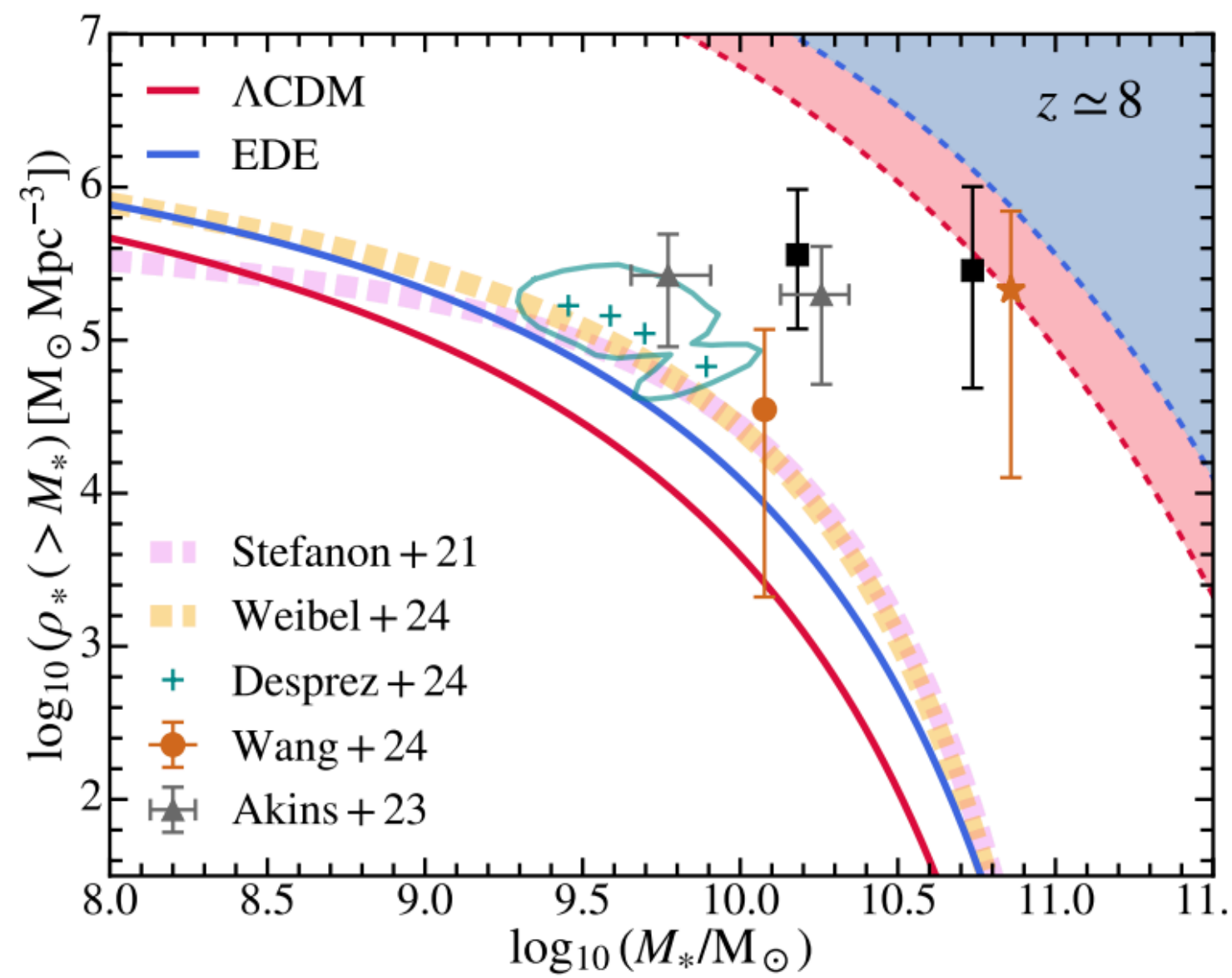


Figure credit: D. Gilman et al. (2022) arXiv:2112.03293





# Universe Timeline

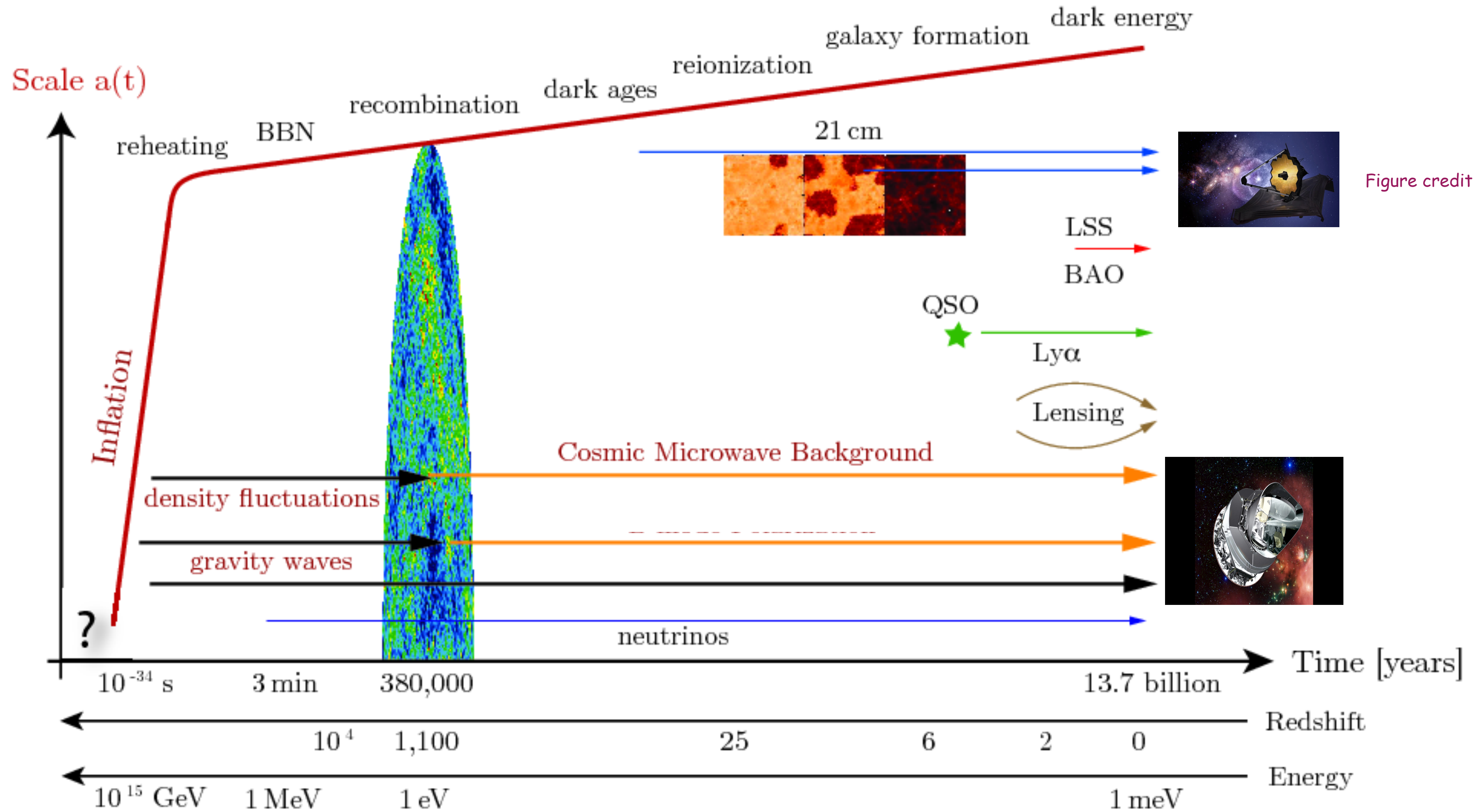


Figure credit: NASA webb

# Cumulative Comoving Number and Mass Densities

Let me define the cumulative comoving number density of haloes with masses above some threshold  $M_{\text{halo}}$  as

$$n(> M_{\text{halo}}, z) = \int_{M_{\text{halo}}}^{\infty} dM \frac{dn(M, z)}{dM},$$

and the corresponding cumulative comoving mass density of haloes

$$\rho(> M_{\text{halo}}, z) = \int_{M_{\text{halo}}}^{\infty} dM M \frac{dn(M, z)}{dM}.$$

Cumulative comoving galaxy number density:

$$n_*( > M_*, z) = n(> M_{\text{halo}}, z)$$

and cumulative comoving stellar mass density:

$$\rho_*( > M_*, z) = \epsilon f_b \rho(> M_{\text{halo}}, z)$$

where stellar mass  $M_* = \epsilon f_b M_{\text{halo}}$

and  $f_b = \Omega_b / \Omega_m$

- The star formation efficiency,  $\epsilon \leq 1$
- Exact value depends on star formation physics

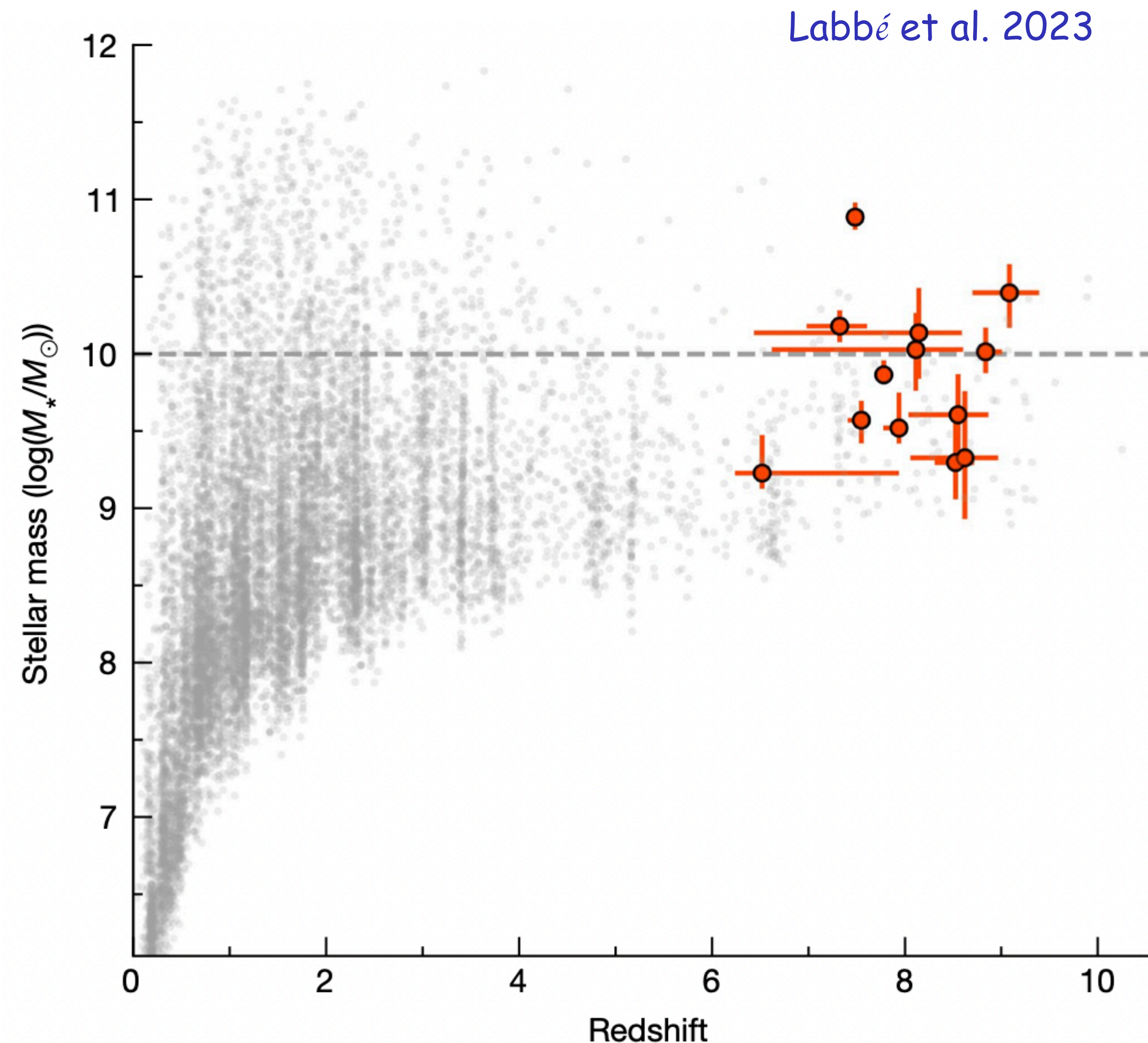


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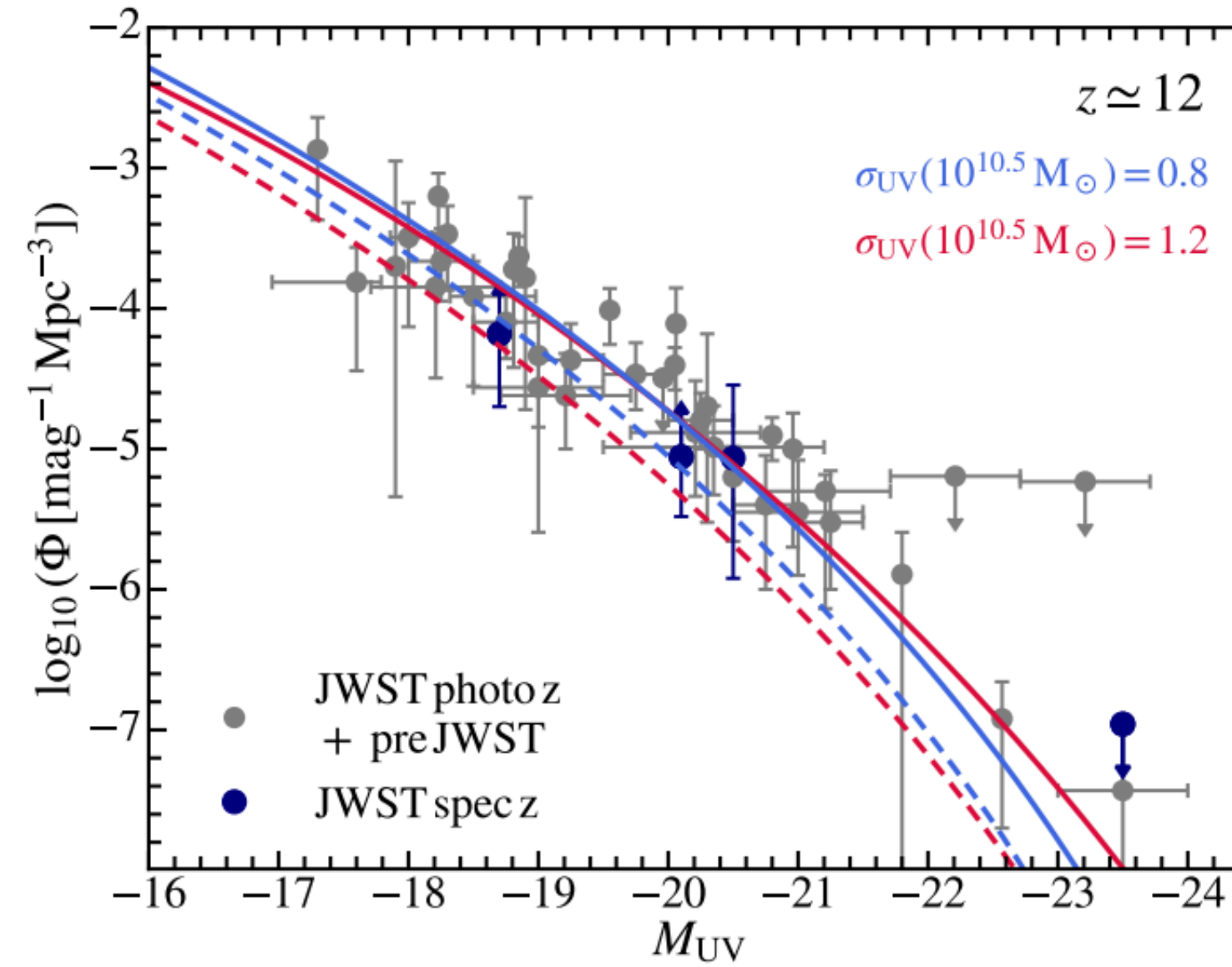
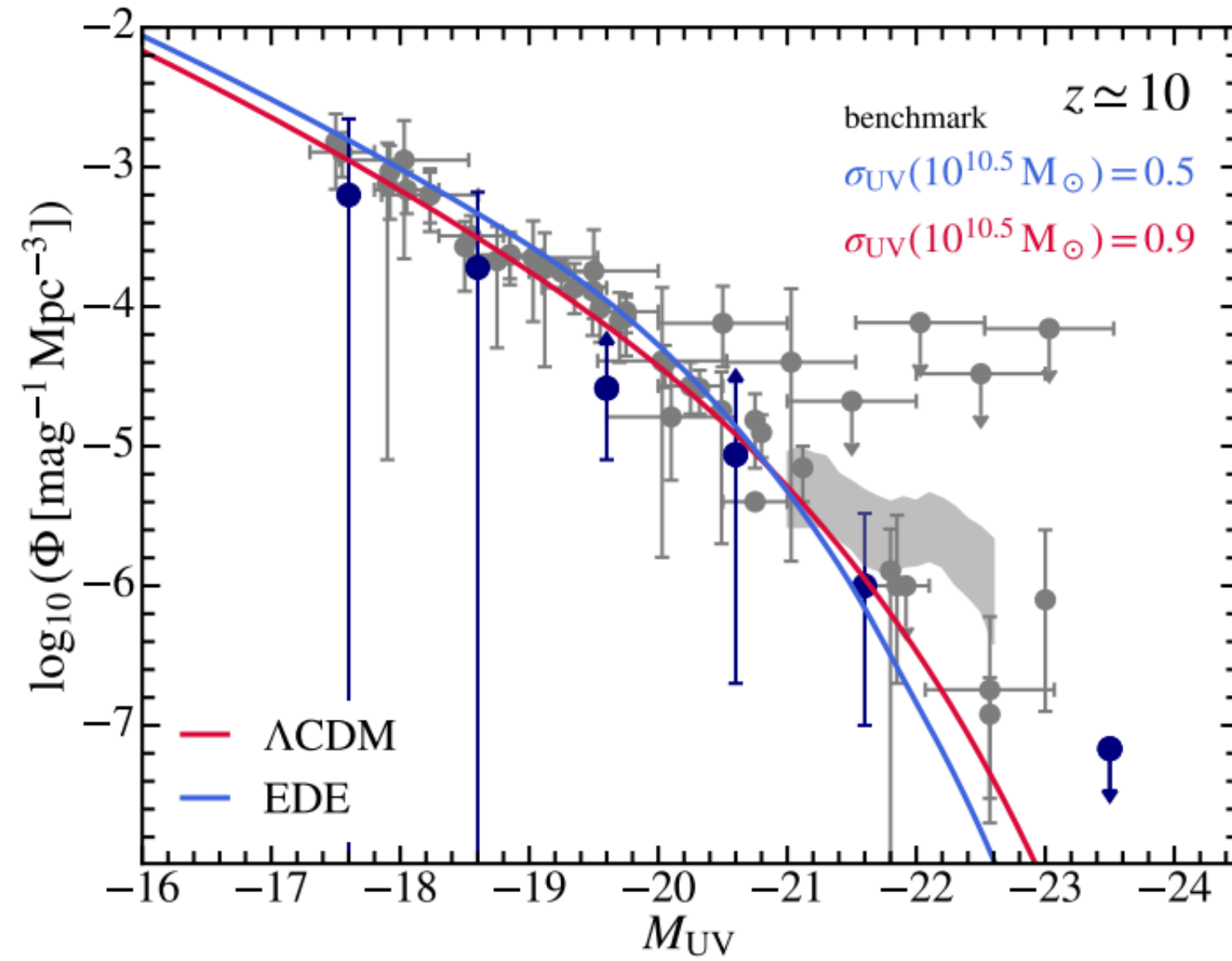
- [Labbé et al. 2023](#) have found 13 galaxy candidates ( $6.5 \leq z \leq 9.1$ ) using JWST data released within the Cosmic Evolution Early Release Science (CEERS) program.
- They identified these candidates by observing two redshifted breaks in their spectral energy distributions (SEDs): 1. Lyman break (1216 Å) and 2. Balmer break (3600 Å).
- Six candidates found to have stellar mass  $M_* > 10^{10} M_\odot$ .

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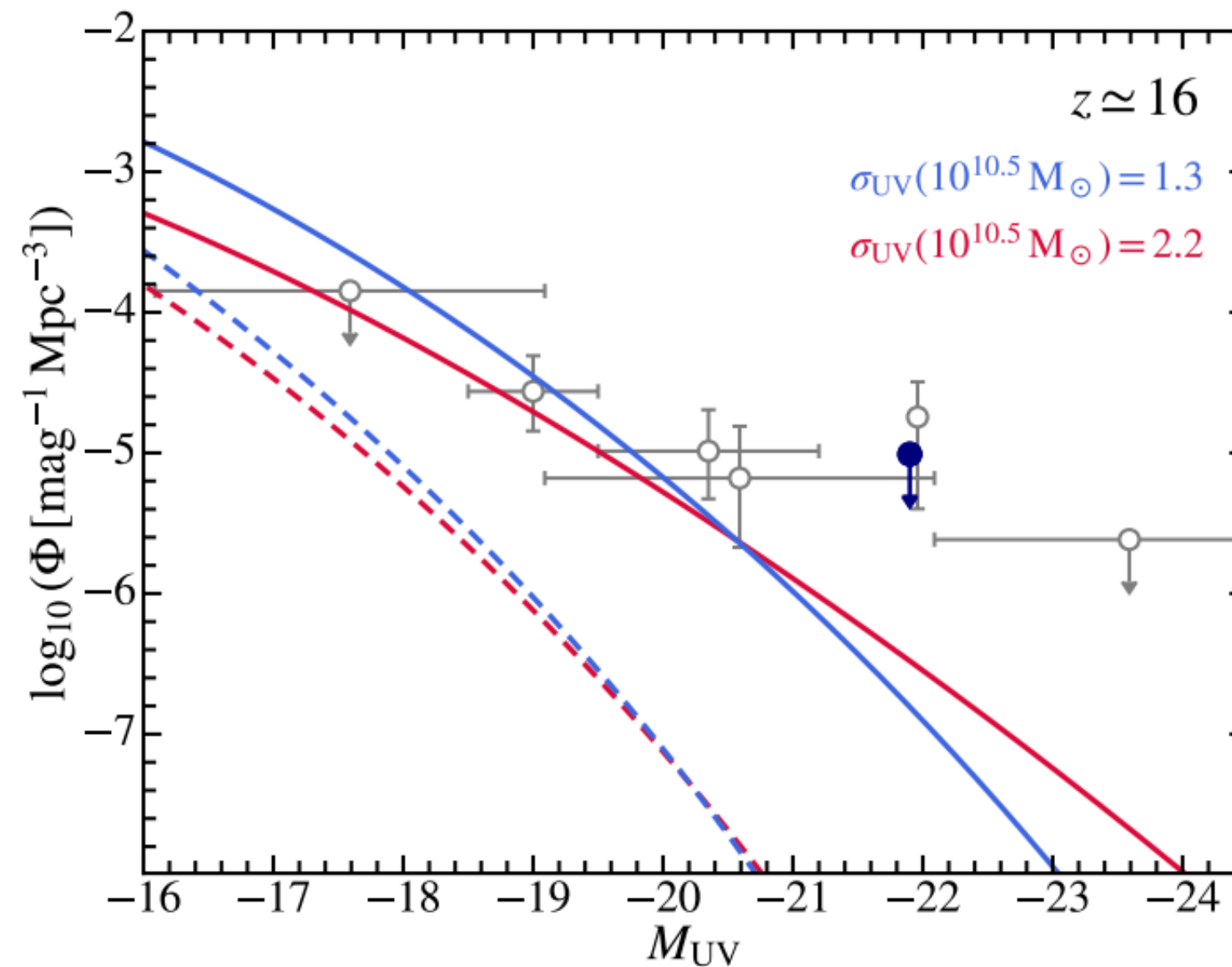
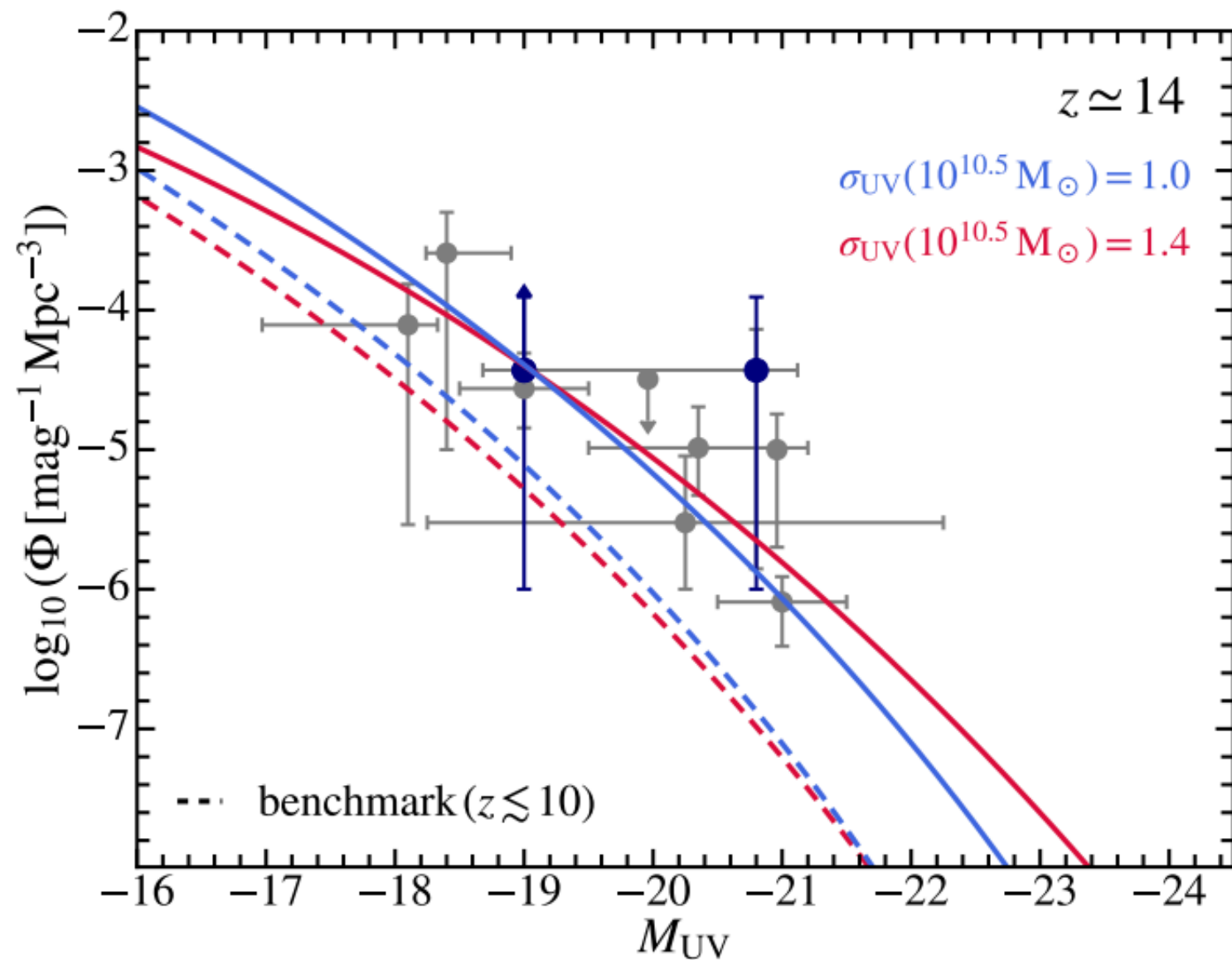




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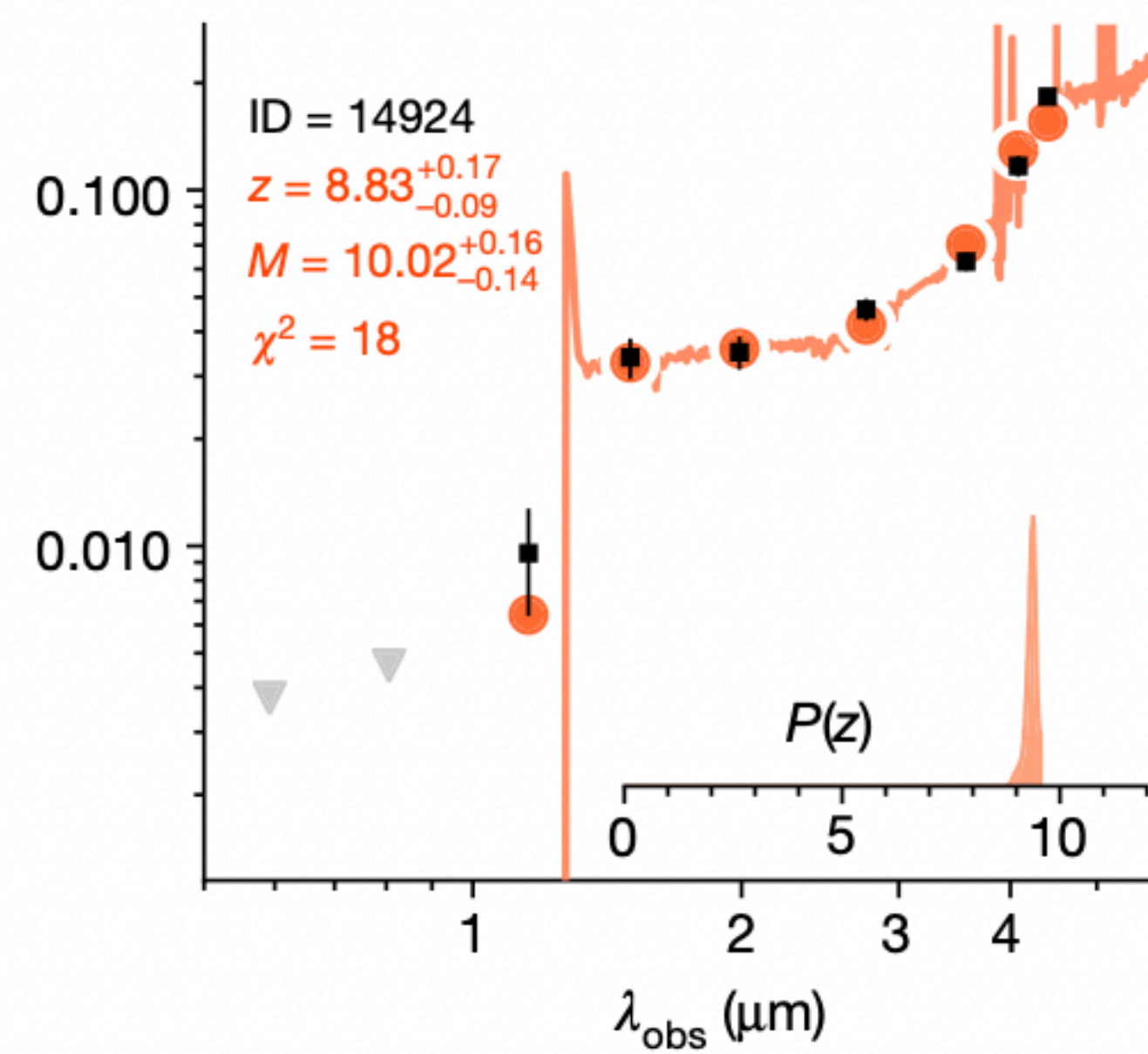
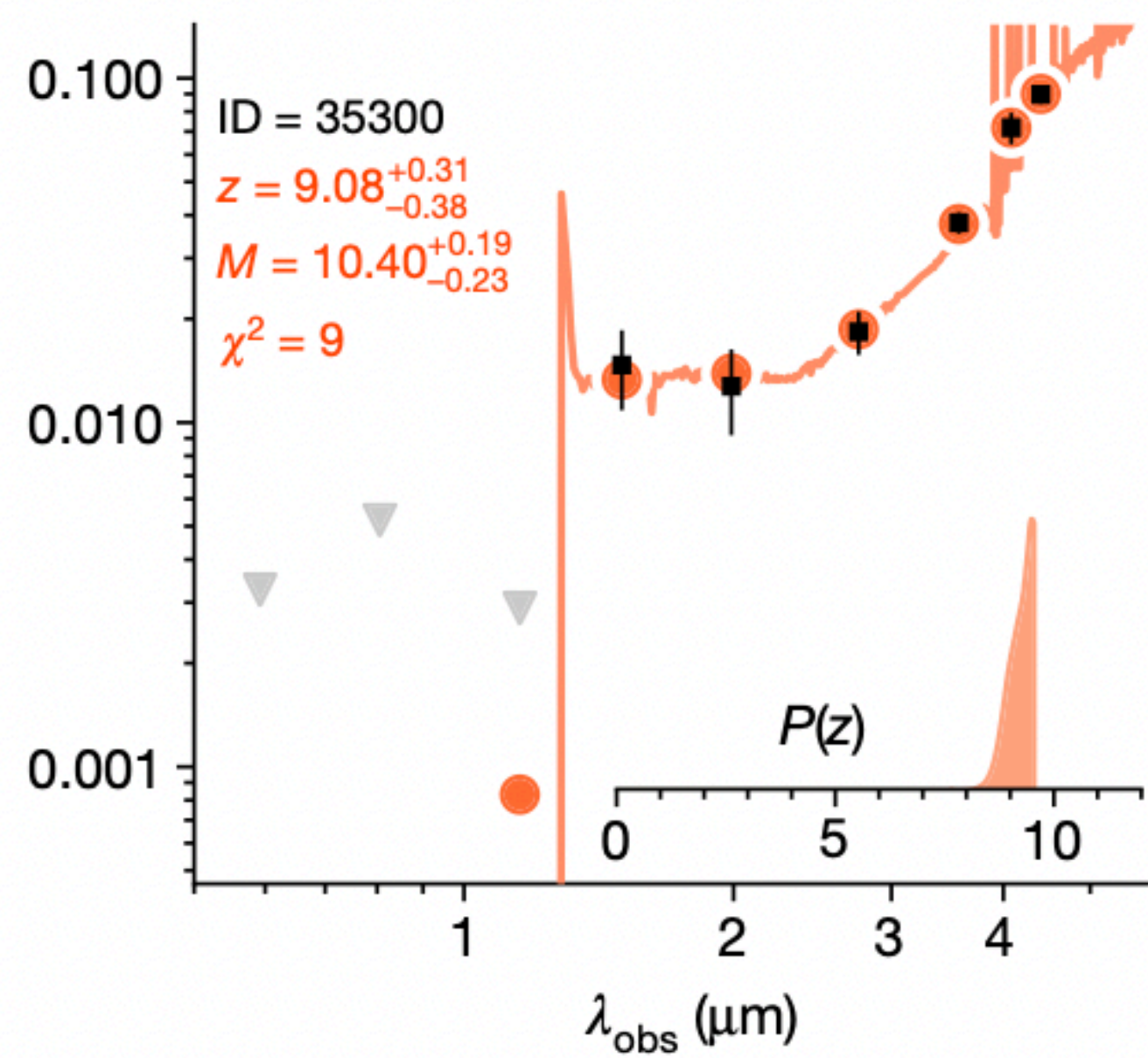
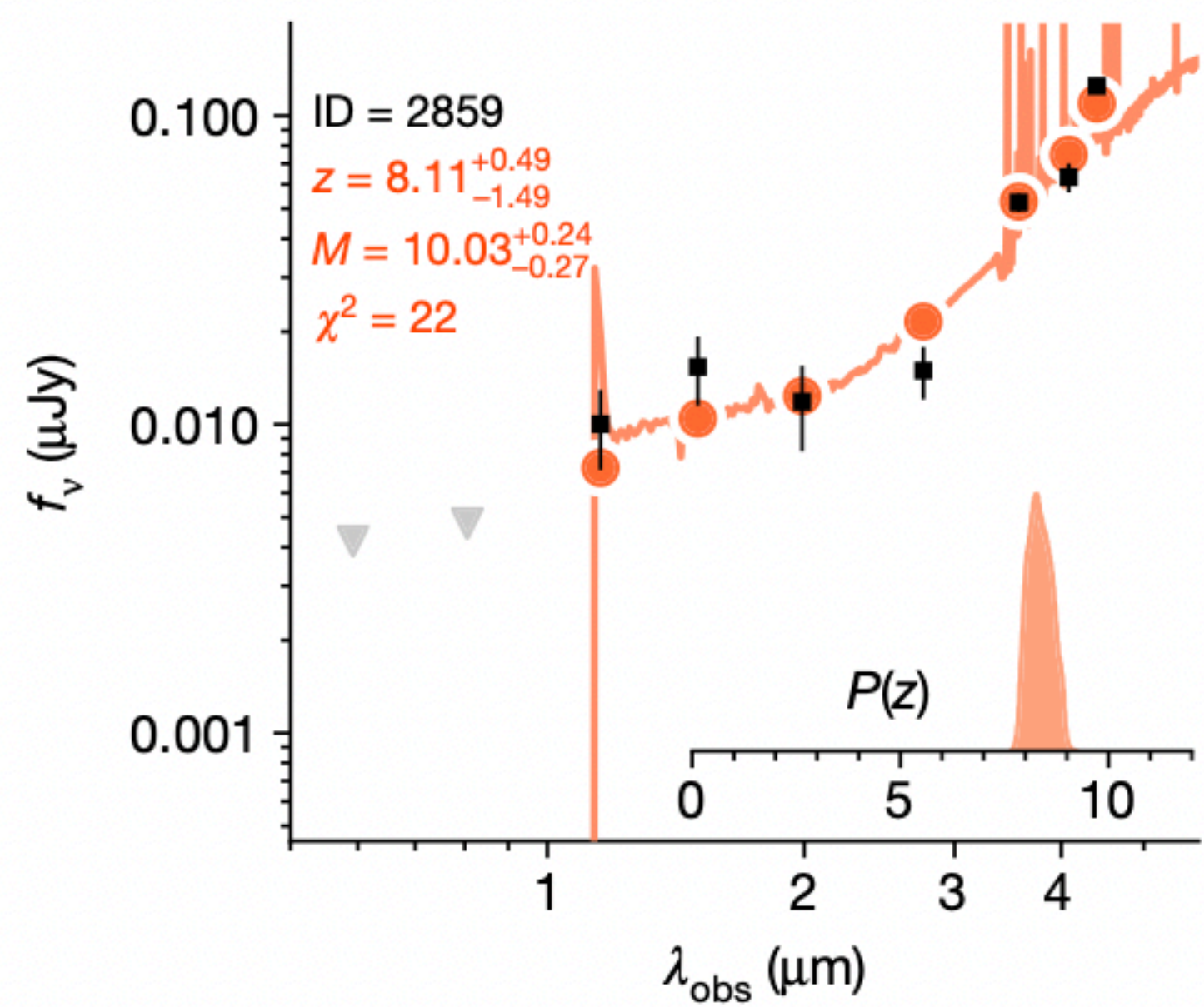
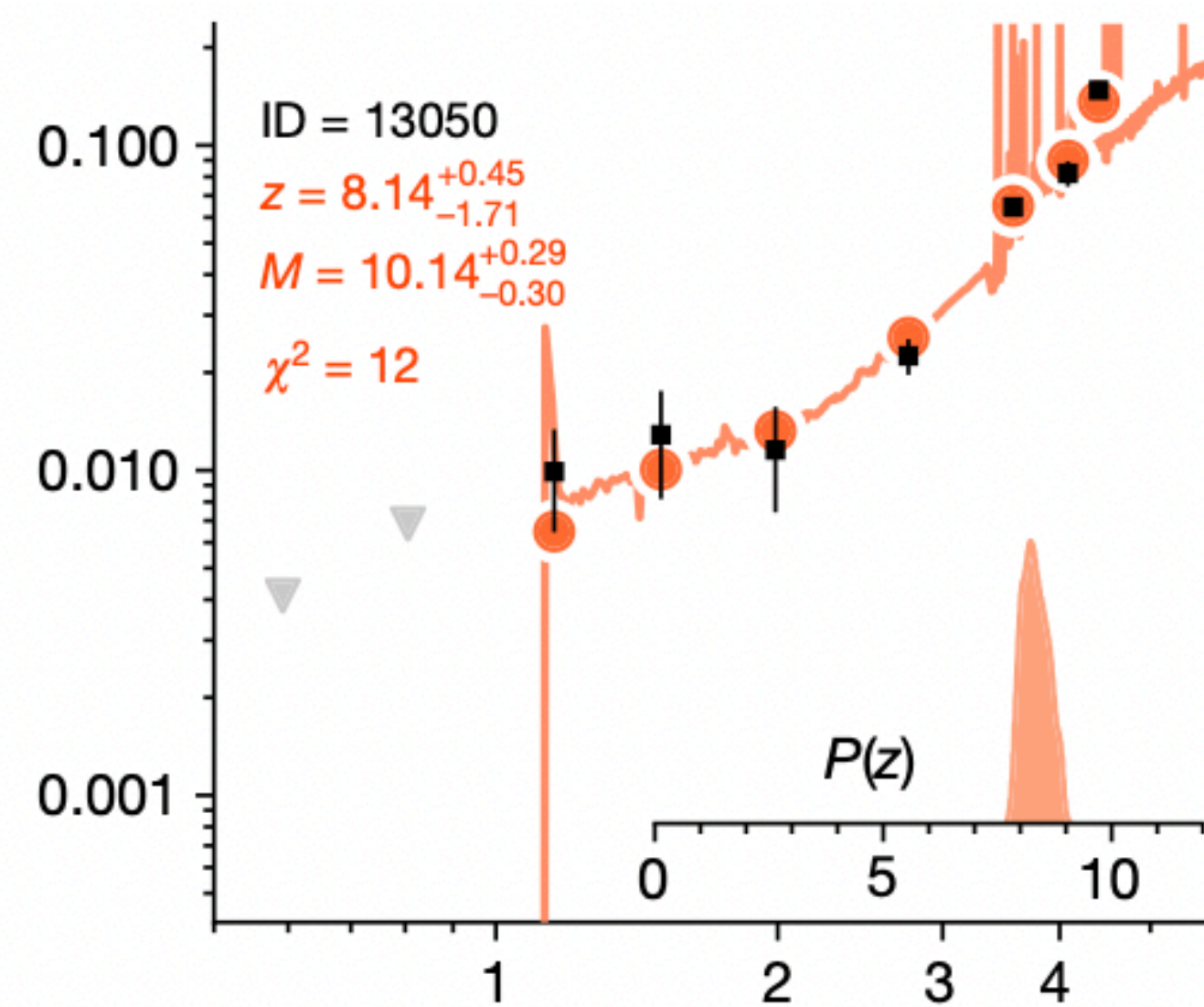
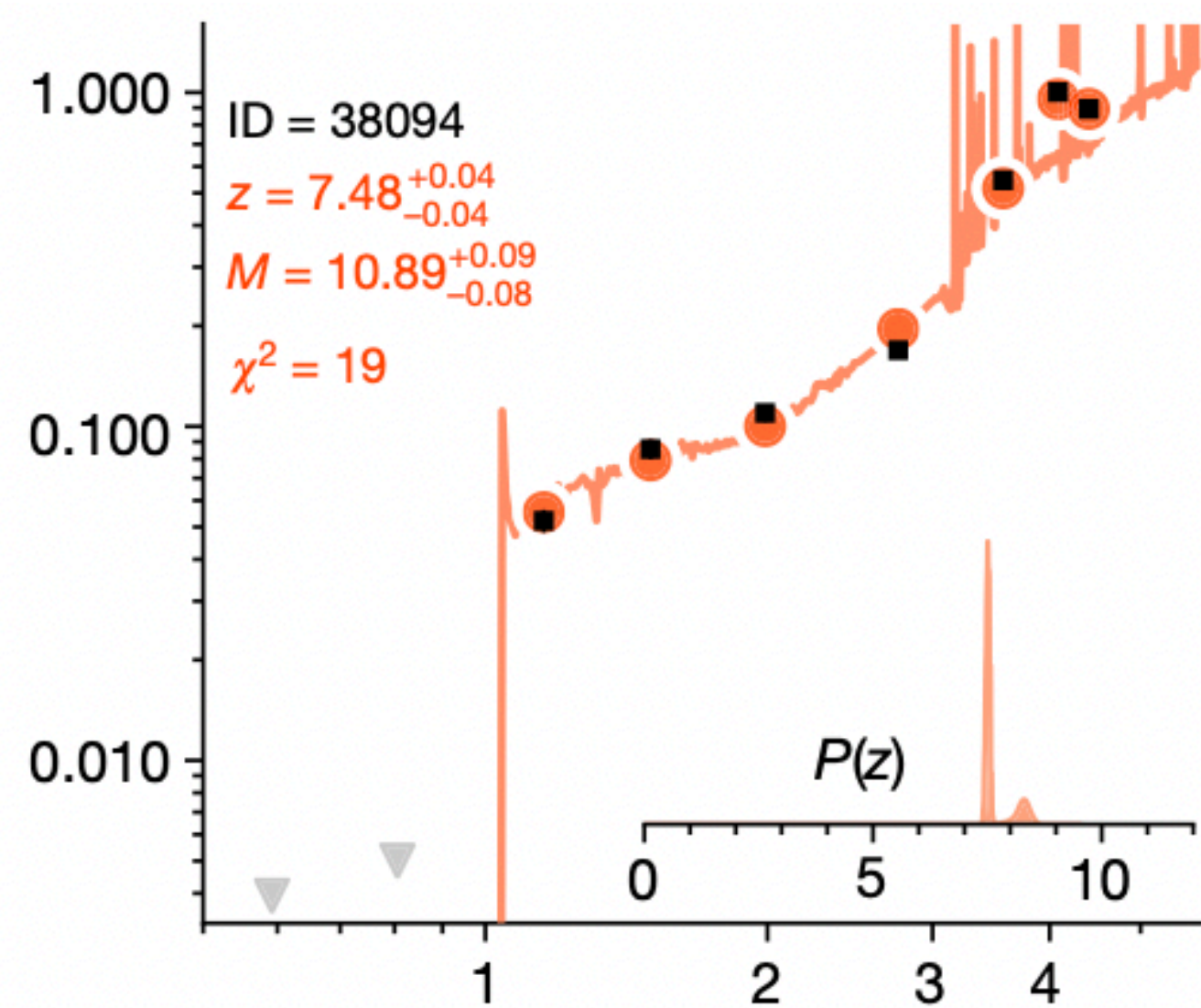
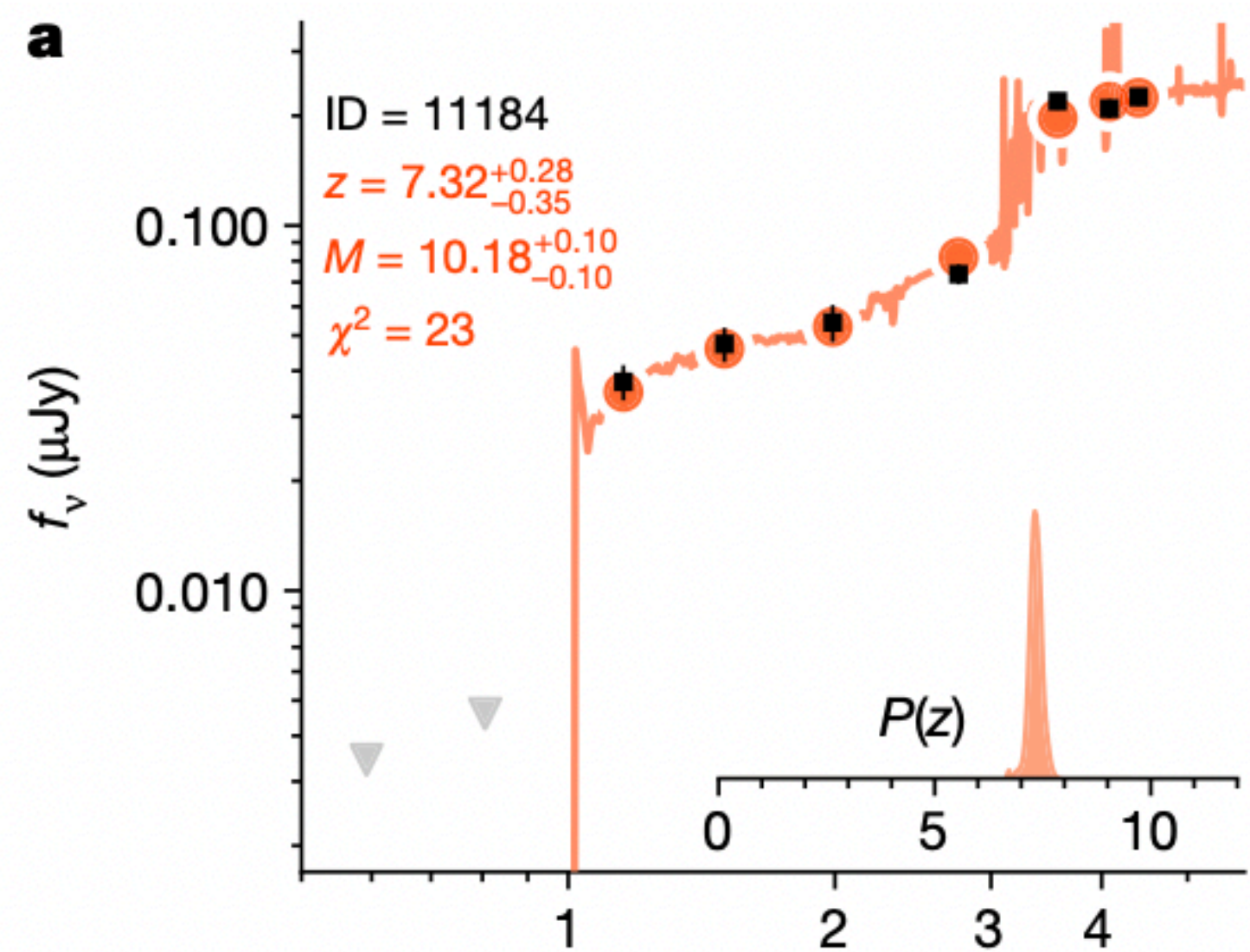


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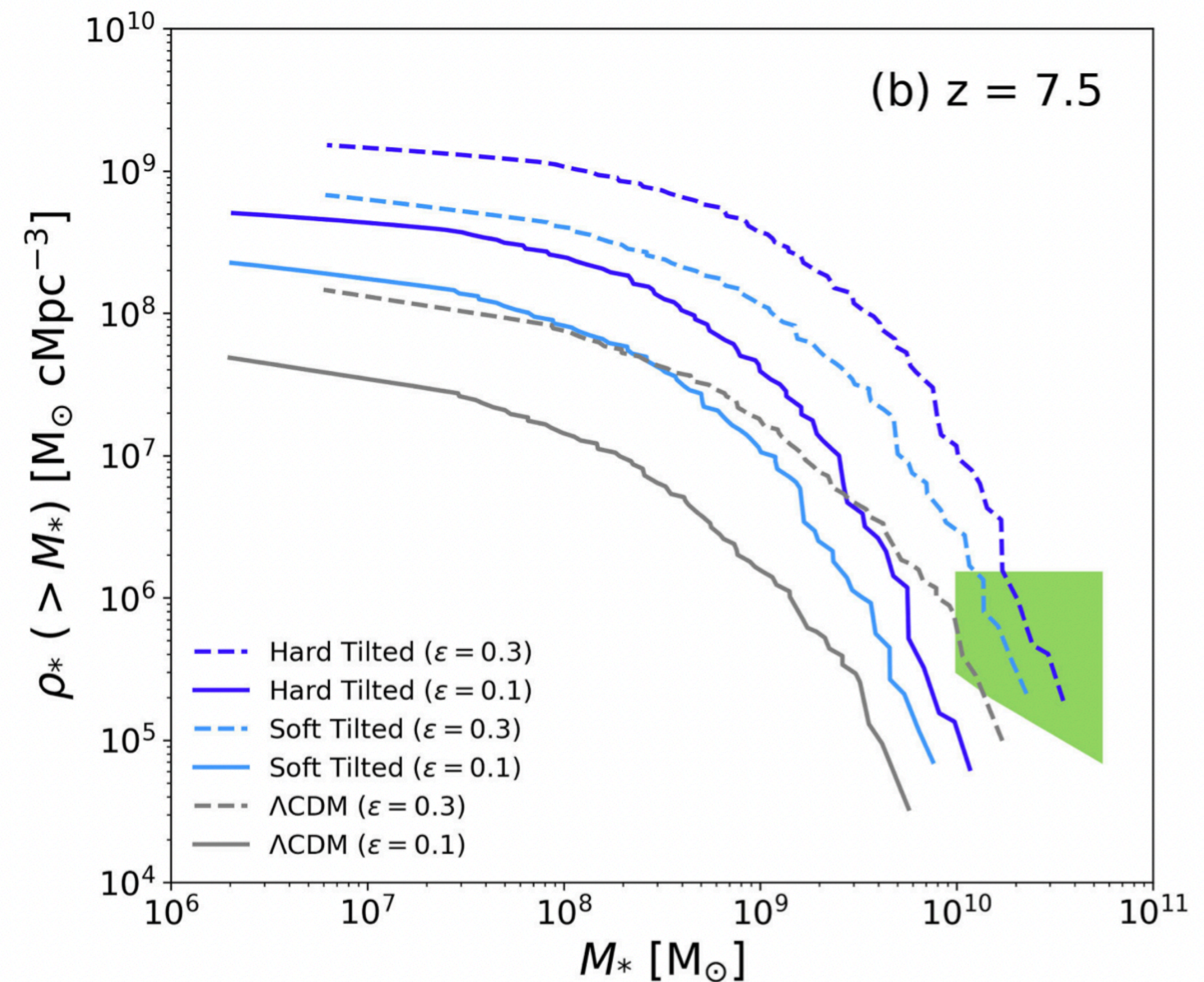
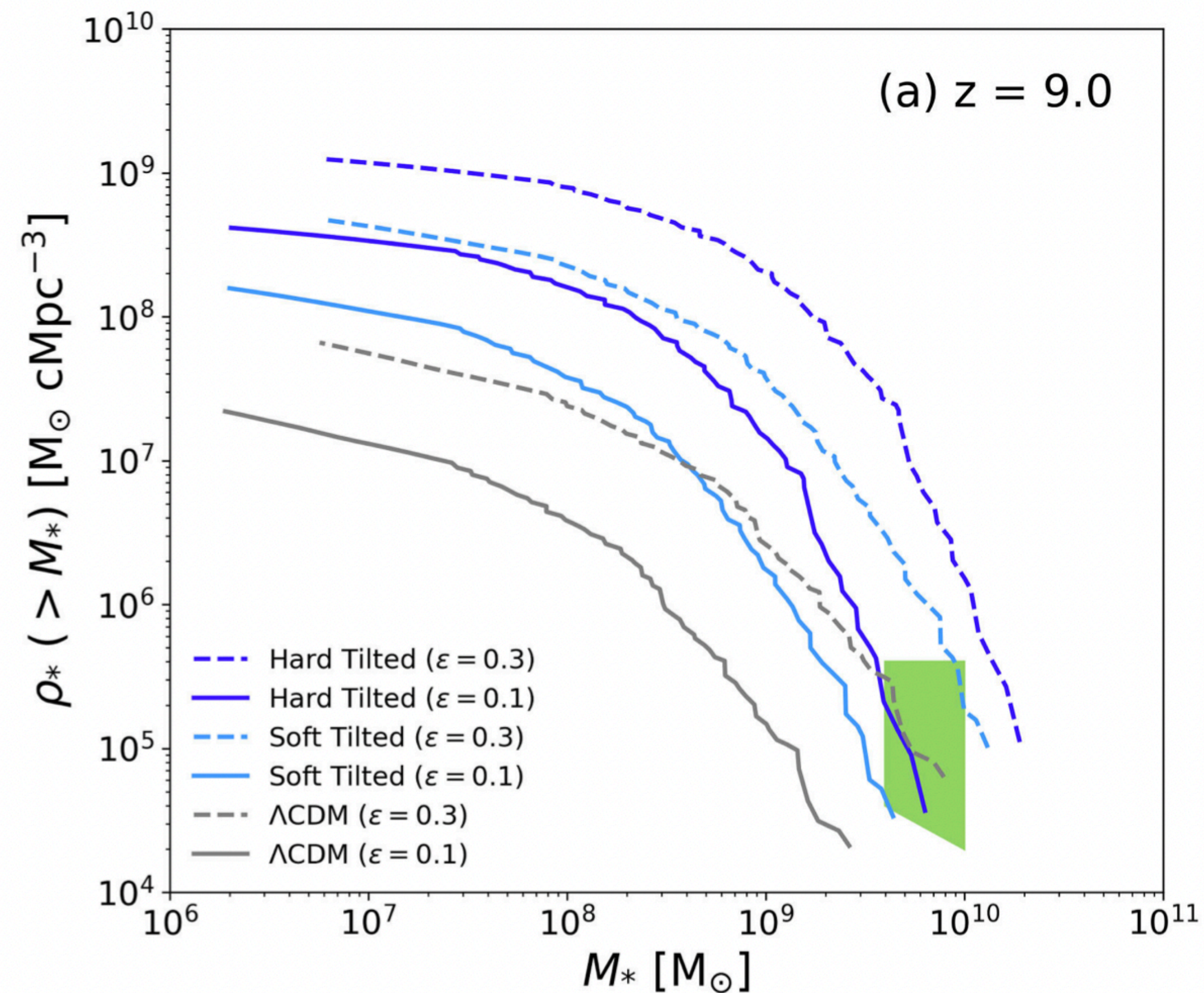




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Priyank Parashari & Ranjan Laha, MNRAS: Letters, 526, L63-L69 (2023)

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