Testable Baryogenesis and Leptonic CP Violation in Seesaw Models

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Based on arXiv:1606.06719 and arXiv:1611.05000

Neutrinos are massive



6.5 0.2 0.25

0.35

0.4 sin² 0,2

[M.C. Gonzalez-Garcia et. al. JHEP 01 (2017) 087 www.nu-fit.org]

0.02 0.025

sin² 0₁₃

Seesaw Type-I

$$\mathcal{L} = \mathcal{L}_{SM} - \sum_{\alpha,i} \bar{L}^{\alpha} Y^{\alpha i} \tilde{\Phi} \nu_R^i - \sum_{i,j=1}^3 \frac{1}{2} \bar{\nu}_R^{ic} M_N^{ij} \nu_R^j + h.c..$$

- Neutrino Masses suggest a NEW PHYSICS scale.
- We focus in the simplest model type-l seesaw
- $Y_N = O(1) \rightarrow \text{Hierarchy}$ problem and not testable.

$$m_{\nu} = Y_N^T \frac{v^2}{M_N} Y_N$$

Minkowski; Yanagida; Glashow; Gell-Mann,

Ramond Slansky; Mohapatra, Senjanovic...



Low-Scale Seesaw (testable)

- Testable in collider (SHIP and FCC)
- Interesting cosmological implications: Leptogenesis (Akhmedov, Rubakov, Smirnov)
- Potential implications in neutrino-less double β decay.
- Small Yukawas, but not so far from the SM.

$$m_{\nu} = Y_N^T \frac{v^2}{M_N} Y_N$$



Parametrization of the Model

- ▶ We set the minimal models with 2 heavy sterile neutrinos.
- Casas-Ibarra parametrization helps us to parametrize the model easily imposing what we know about the light sector.

 $U_{\alpha h} = i U_{\rm PMNS} \sqrt{m_l} \ R^{\dagger}(z) M^{-1/2}$

► We set U_{PMNS} angles and m_l square differences to the current best fit(NOT the PHASE, ORDERING).

 $\{\delta,\phi_1,\gamma,\theta,M_1,M_2\}$

Cosmological Implications

- The heavy steriles are too heavy and unstable to directly affect any CMB or LSS cosmology.
- BUT before EW phase transition we have:
 - Baryon number violation, in our case lepton number violation
 + sphalerons
 - ► Loss of thermal equilibrium, the steriles interact very weekly.
- ► C and CP violation processes, here they come from oscillations. The model acomplish three Sakharov conditions (Andrei Sakharov 1967)

therefore it will produce primordial asymetry

(Akhmedov, Rubakov, Smirnov)

Formalism: Oscillations plus interactions in a primordial plasma Rafelt-Sigle equations for neutrinos plus the lepton chemical potentials.

$$\begin{split} xH_{u} \frac{dr_{N}}{dx} &= -i[\langle H \rangle, r_{N}] - \frac{\langle \gamma_{N}^{(0)} \rangle}{2} \{Y^{\dagger}Y, r_{N} - 1\} + \langle \gamma_{N}^{(1)} \rangle Y^{\dagger}\mu Y - \frac{\langle \gamma_{N}^{(2)} \rangle}{2} \{Y^{\dagger}\mu Y, r_{N}\} \\ xH_{u} \frac{dr_{\bar{N}}}{dx} &= -i[\langle H^{*} \rangle, r_{\bar{N}}] - \frac{\langle \gamma_{N}^{(0)} \rangle}{2} \{Y^{T}Y^{*}, r_{\bar{N}} - 1\} + \langle \gamma_{N}^{(1)} \rangle Y^{T}\mu Y^{*} \\ &+ \frac{\langle \gamma_{N}^{(2)} \rangle}{2} \{Y^{T}\mu Y^{*}, r_{\bar{N}}\}, \\ xH_{u} \frac{d\mu_{B/3-L_{\alpha}}}{dx} &= \frac{\int_{k} \rho_{F}}{\int_{k} \rho_{F}'} \left\{ \frac{\langle \gamma_{N}^{(0)} \rangle}{2} (Yr_{N}Y^{\dagger} - Y^{*}r_{\bar{N}}Y^{T})_{\alpha\alpha}. \\ &+ \mu_{\alpha} \left(\frac{\langle \gamma_{N}^{(2)} \rangle}{2} (Yr_{N}Y^{\dagger} + Y^{*}r_{\bar{N}}Y^{T})_{\alpha\alpha} - \langle \gamma_{N}^{(1)} \rangle \mathrm{Tr}[YY^{\dagger}I_{\alpha}] \right) \right\}, \\ \mu_{\alpha} &= -\sum_{\beta} C_{\alpha\beta}\mu_{B/3-L_{\beta}}. \end{split}$$

- Thermal quantum field theory effects.
- ► Almost all depends on time "x": Numerically hard problem!

Considerations in the derivation

- ► Fermi-Dirac or Bose-Einstein statistics is kept throughout.
- ► Collision terms include 2 ↔ 2 scatterings at tree level with top quarks and gauge bosons, as well as 1 ↔ 2 scatterings including the resummation of scatterings mediated by soft gauge bosons.
- Leptonic chemical potentials are kept in all collision terms to linear order
- Include spectator processes

If we have tools to solve the equations we can compare with observations

Baryon asymetry is very well measured:

 $Y_B^{\text{exp}} \simeq 8.65(8) \times 10^{-11}.$

- Essentially the same technical problem that doing a CMB-LSS precision cosmology fit.
- ► TOOLS:
 - 1. Bayesian Inference(Multinest)
 - 2. Hard Numerical Computations (SQuIDS)



https://github.com/jsalvado/SQuIDS

https://github.com/JohannesBuchner/MultiNest

Flat priors in the Casas-Ibarra parameters $\{\delta, \phi_1, \gamma, \theta, M_1, M_2\}$ except:

1. Flat prior in $\log_{10}\left(\frac{M_1}{\text{GeV}}\right)$ and $\log_{10}\left(\frac{M_2}{GeV}\right)$ 2. Flat prior in $\log_{10}\left(\frac{M_1}{\text{GeV}}\right)$ and $\log_{10}\left(\frac{|M_2-M_1|}{\text{GeV}}\right)$





What if we really measure something* in SHiP

Forecast for SHIP (with and without a measure δ_{CP}



What else can we learn: CP discovery potential

Assumptions:

- 1. SHiP or FCC measures one of the steriles.
- 2. The steriles are the main contribution in generating the light neutrino masses by the low scale seesaw.

Goal:

 Quantify the potential to exclude any of the CP conserving points(NULL HYPOTHESIS):

 $(\delta,\phi_1) = \{(0,0), (0,\pi), (\pi,0), (\pi,\pi)\}$

5σ contours for different gammas, Inverted and Normal



 R_{CP} in terms of $M_1 U^2$ and M_1 for SHiP and FCC-ee.



Conclusions

- We used state of the art computation and numerical tools in low scale leptogeneis.
- ► The precise measurement of the asymmetry is very relevant:
 - Not very degenerate values are preferred.
 - Important overlapping with the sensitivity regions of SHiP and FCC.
- ► A measurement of SHiP, δ_{CP} and $m_{\beta\beta}$ may give a prediction for the primordial asymmetry(Testable)
- A SHiP or FCC measurement has strong power to prove CP violation in the leptonic sector.