


# Benchmarks on Double Higgs production for Singlet Extension



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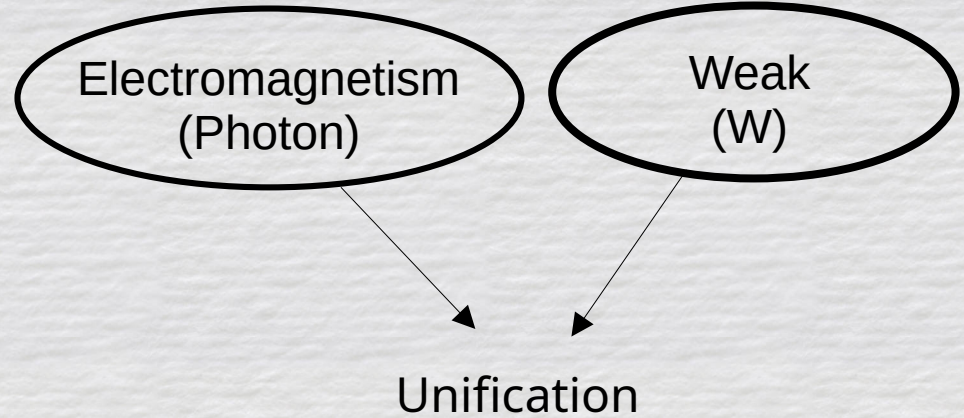
Collaborating with: I. Lewis; M. Sullivan; J. Scott

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# Going Beyond The Standard Model of Elementary Particles

three generations of matter (fermions)			interactions / force carriers (bosons)		
	I	II	III		
mass	$\approx 2.2 \text{ MeV}/c^2$	$\approx 1.28 \text{ GeV}/c^2$	$\approx 173.1 \text{ GeV}/c^2$	0	$\approx 124.97 \text{ GeV}/c^2$
charge	$\frac{2}{3}$	$\frac{2}{3}$	$\frac{2}{3}$	0	0
spin	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	1	0
	<b>u</b> up	<b>c</b> charm	<b>t</b> top	<b>g</b> gluon	<b>H</b> higgs
	<b>d</b> down	<b>s</b> strange	<b>b</b> bottom	<b><math>\gamma</math></b> photon	
	<b>e</b> electron	<b><math>\mu</math></b> muon	<b><math>\tau</math></b> tau	<b>Z</b> Z boson	
	<b><math>\nu_e</math></b> electron neutrino	<b><math>\nu_\mu</math></b> muon neutrino	<b><math>\nu_\tau</math></b> tau neutrino	<b>W</b> W boson	



SCALAR BOSONS

$m=?$

0

0

**S**

**Higgs 2**

- Help in further understanding of EW phase transition
- Accuracy for double Higgs production
- Possible DM candidate?

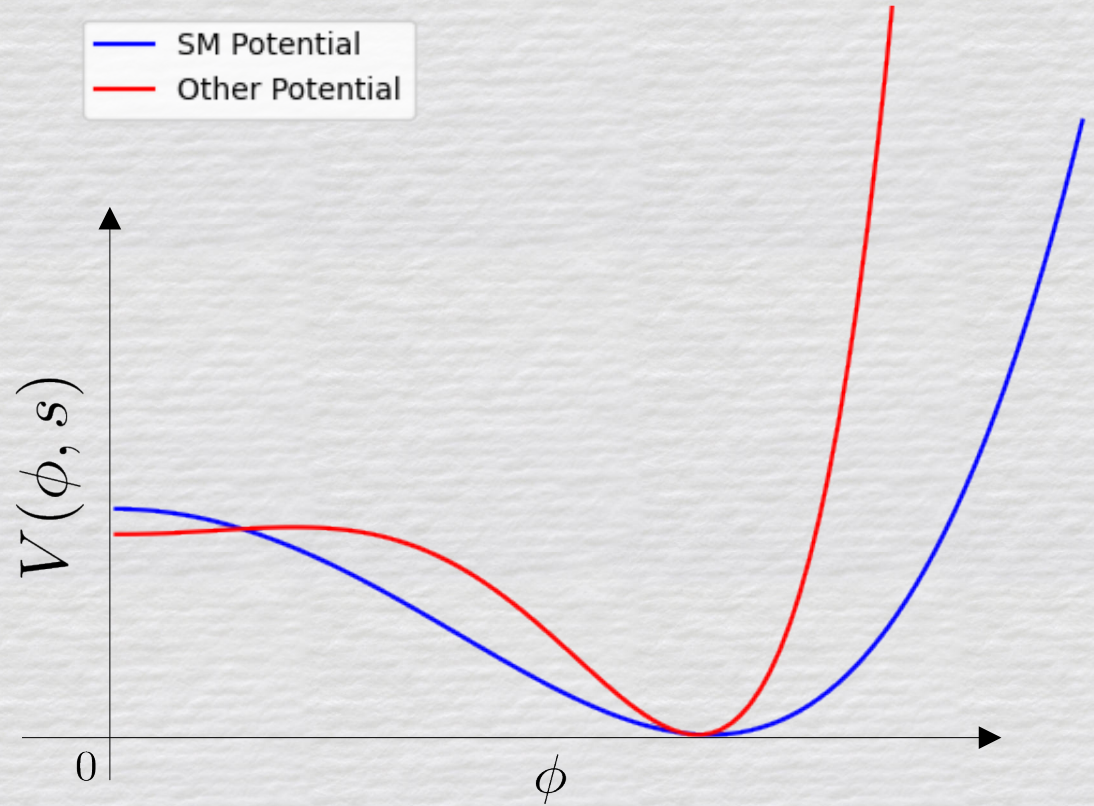
# Model Extension

All interactions embedded  
into the Lagrangian

$$\mathcal{L}_{\text{SM}}$$

The extension must hold  
the already known Global  
minimum

$$\mathcal{L}_{\text{new}} = \mathcal{L}_{\text{SM}} + V(H, S)$$



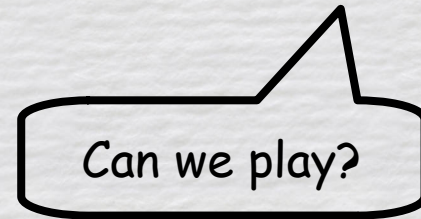
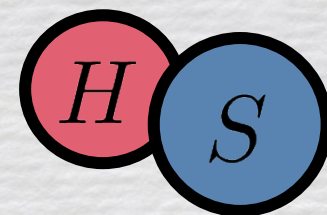
# Model Extension

$$H = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ \phi_0 + v \end{pmatrix}$$

- The simplest extension is the addition of a gauge real singlet

$$S = s + x$$

$$V(H, S) = -\mu^2 H^\dagger H + \lambda (H^\dagger H)^2 + \frac{a_1}{2} H^\dagger H S + \frac{a_2}{2} H^\dagger H S^2 + b_1 S + \frac{b_2}{2} S^2 + \frac{b_3}{3} S^3 + \frac{b_4}{4} S^4.$$



- Only coupling to Higgs doublet  $H$ .
- Real minimums only

$$(v_{EW}, 0),$$

$$(v_{\pm}, x_{\pm}), \quad (0, x_{1,2,3})$$

- Scan to see where is global

C.-Y. Chen, S. Dawson, and I. M. Lewis  
arXiv:1410.5488

No  $Z_2$  symmetry  
 More freedom in parameters

- Manipulate to get rid of as many constants, like

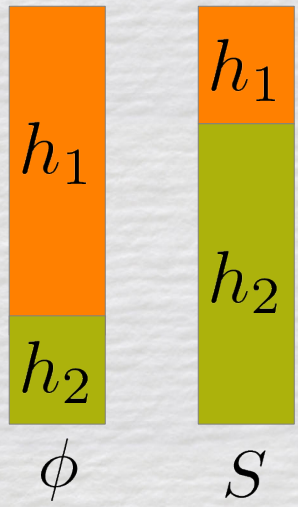
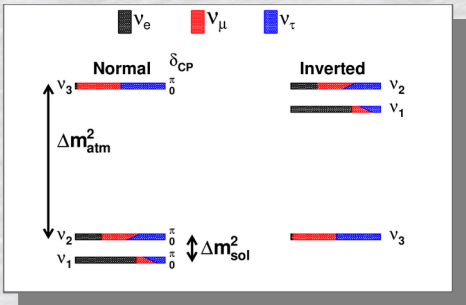
$$\mu^2 = \lambda v_{EW}^2, \quad b_1 = -\frac{v_{EW}^2}{4} a_1 \quad \text{from} \quad V(v_{EW}, 0)$$

- Further manipulation. Transform to the mass eigenstates

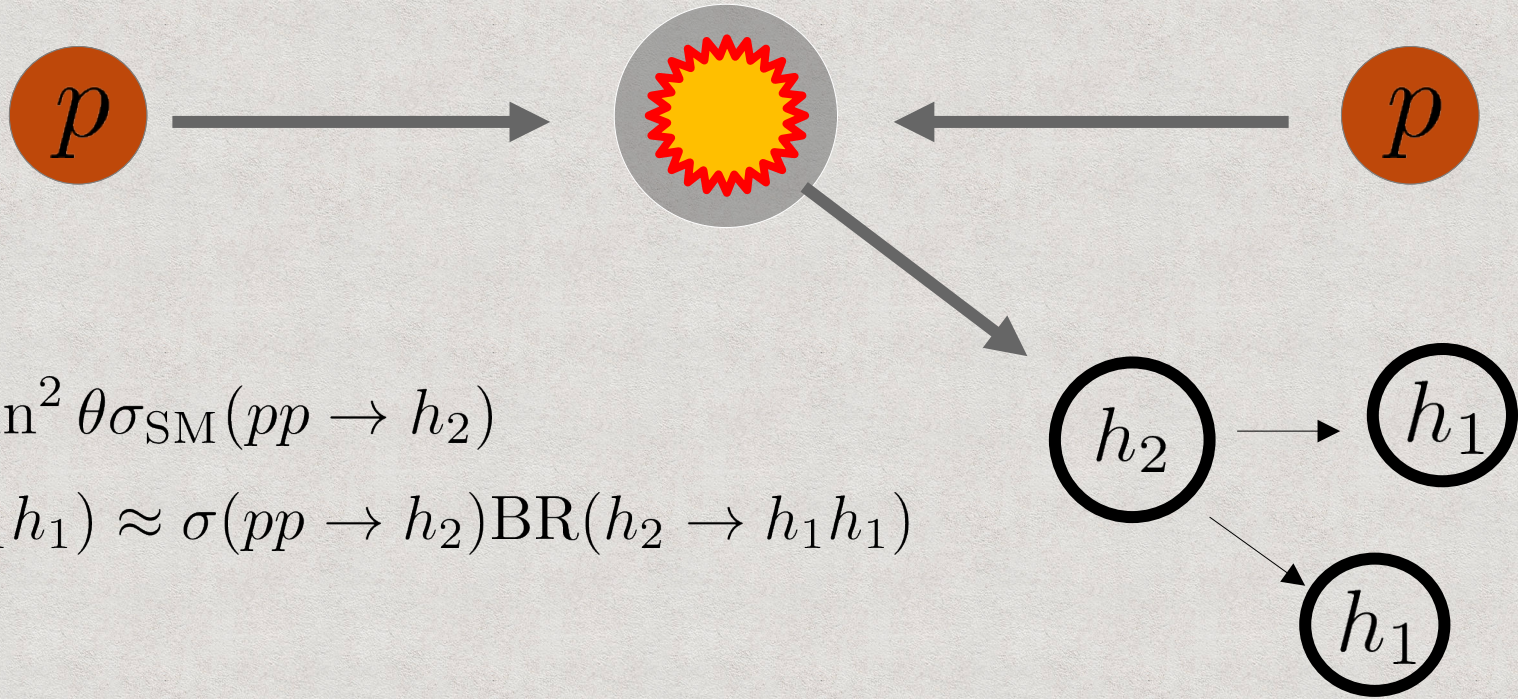
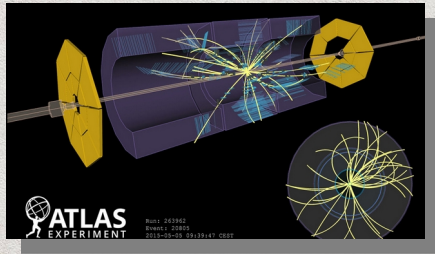
$$V_m(h_1, h_2) \Rightarrow \begin{pmatrix} h_1 \\ h_2 \end{pmatrix} = R(\theta) \begin{pmatrix} \phi \\ s \end{pmatrix}$$

- Similar as neutrinos having flavor and mass eigenstates
- Leaving total of 4 free parameters

$$a_2, b_3, b_4, m_2, \theta$$



# PRODUCTION



$$\sigma(pp \rightarrow h_2) = \sin^2 \theta \sigma_{\text{SM}}(pp \rightarrow h_2)$$

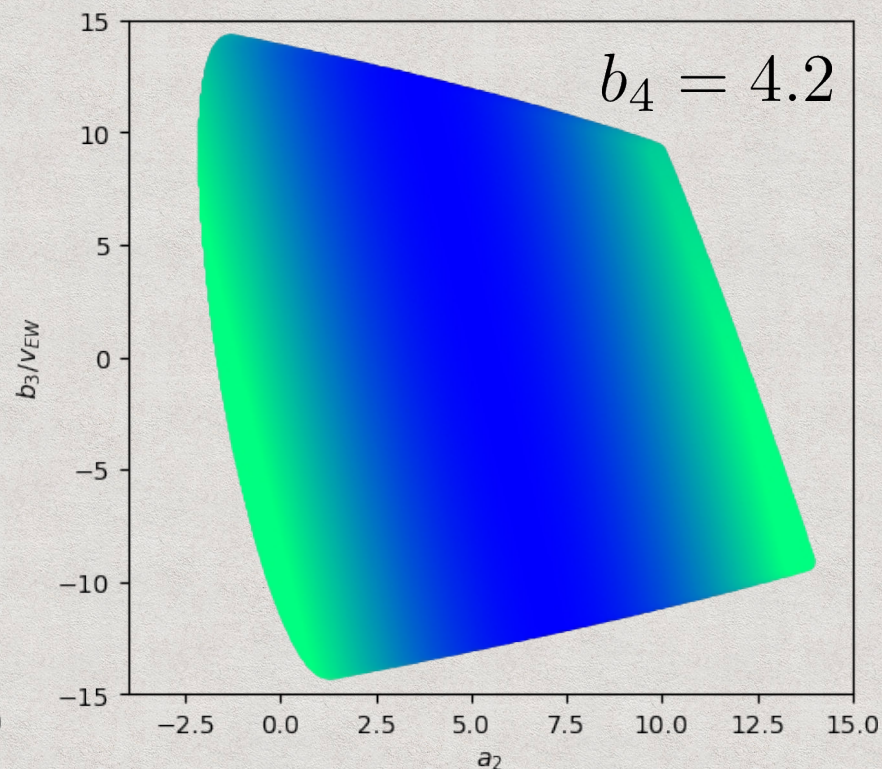
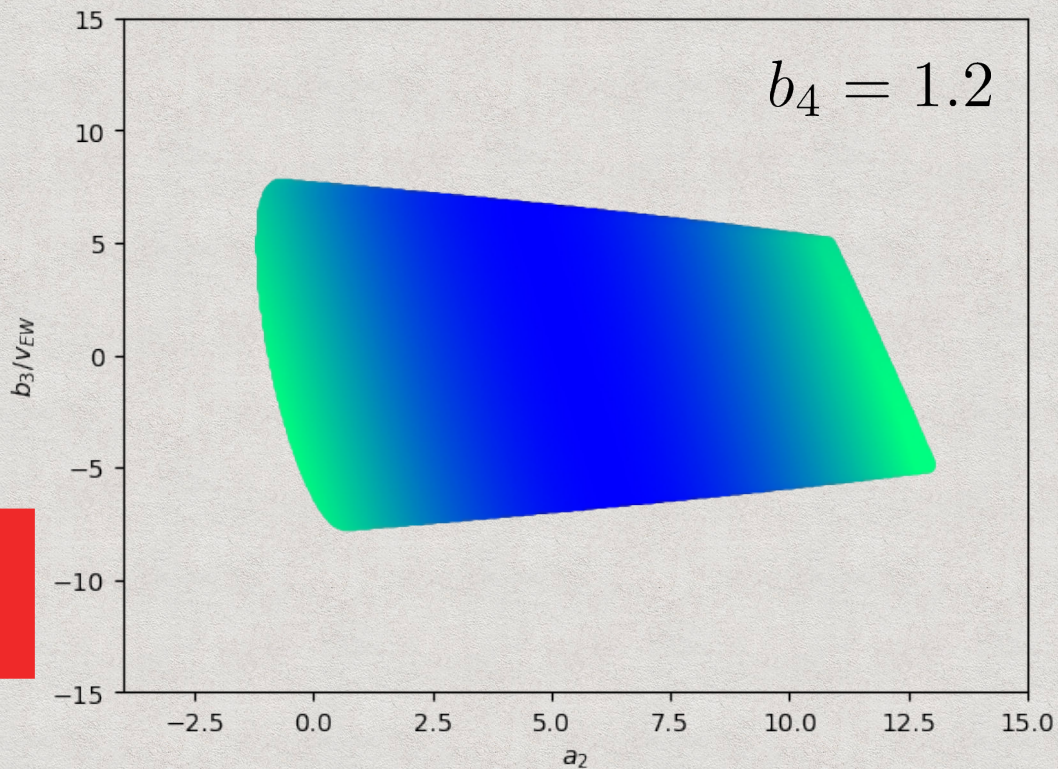
$$\sigma(pp \rightarrow h_2 \rightarrow h_1 h_1) \approx \sigma(pp \rightarrow h_2) \text{BR}(h_2 \rightarrow h_1 h_1)$$

Combine to obtain:

$$\frac{\sigma(pp \rightarrow h_2 \rightarrow h_1 h_1)}{\sigma_{\text{SM}}(pp \rightarrow h_2)} \approx \sin^2 \theta \text{BR}(h_2 \rightarrow h_1 h_1)$$

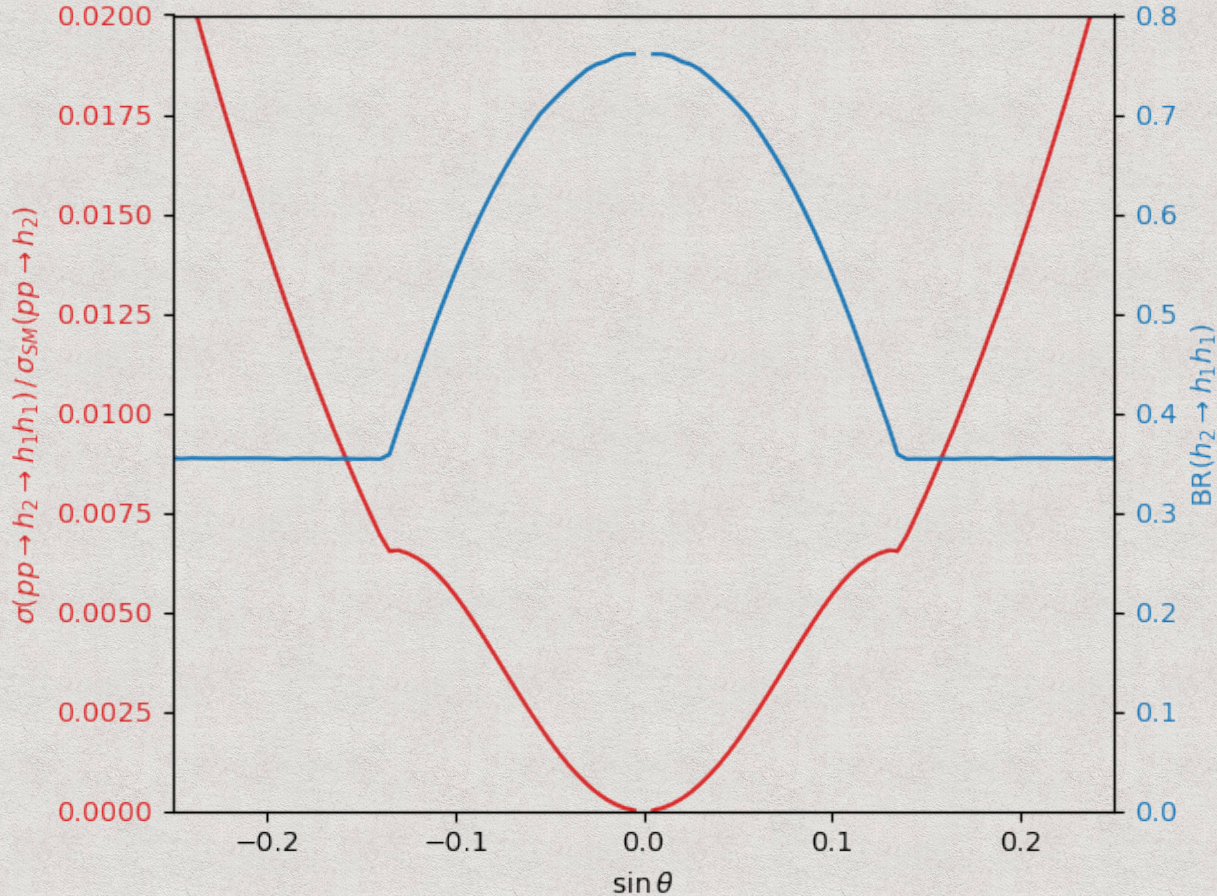
What will be measured

# Valid $a_2$ and $b_3$ regions for $m_2 = 800$ , $\sin \theta = 0.14$



For given angle and mass, max area with  $b_4 = 4.2$   
Now make scan for other parameters

→ Bounded from above by perturbative unitarity.



- Same behaviour for all masses
- In  $\lambda_{211}$ , first there is dominance by  $a_2$  and then by  $\sin \theta$

$$V \supset \frac{\lambda_{211}}{2!} h_2 h_1^2$$

$$\lambda_{211} =$$

$$\sin \theta \left[ -\frac{2m_1^2 + m_2^2}{v_{\text{ew}}} \cos^2 \theta - a_2 v_{\text{ew}} (-3 \cos^2 \theta) + b_3 \sin 2\theta \right]$$



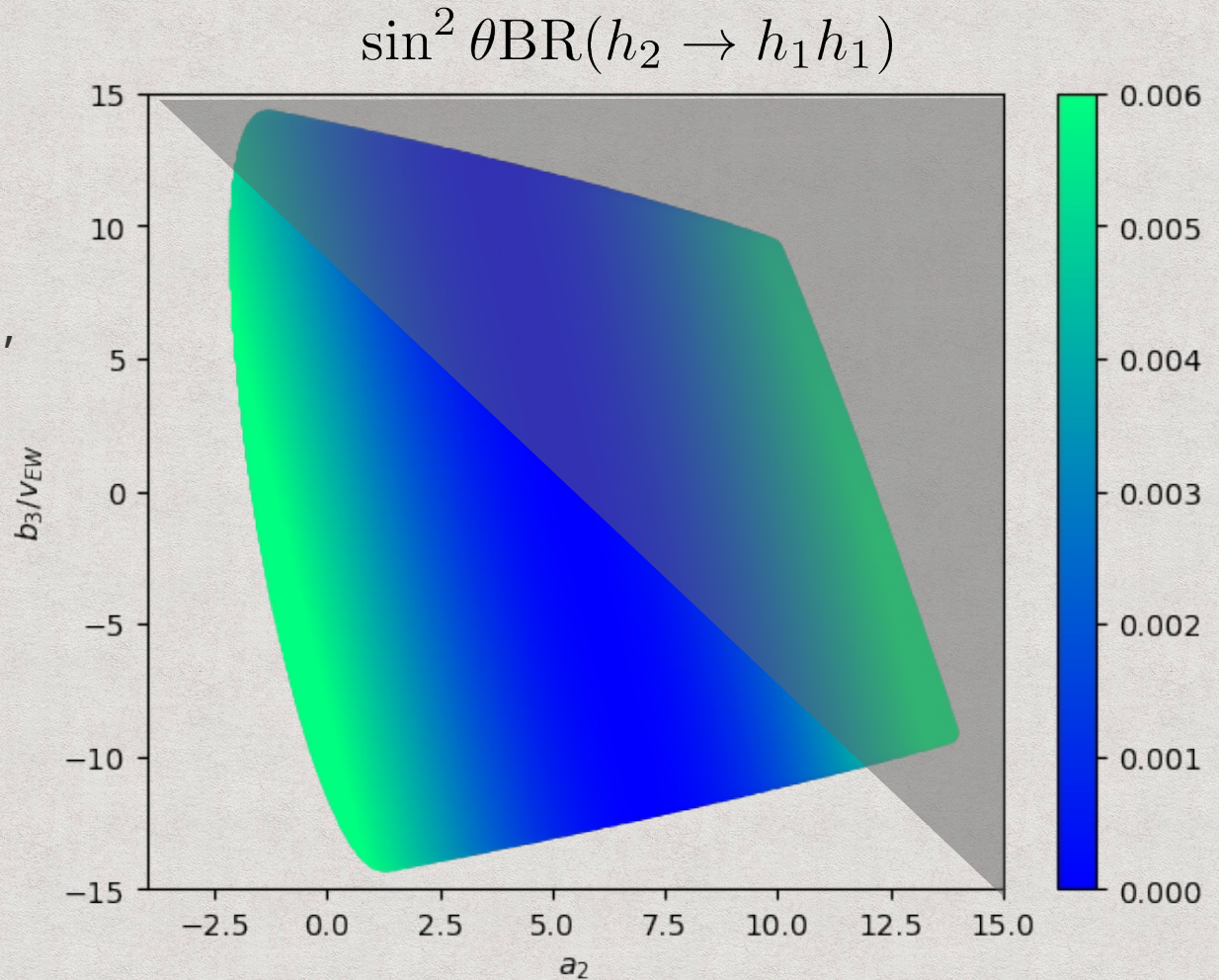
- Narrow width approximation

$$\Gamma(h_2) \leq 0.1 m_2$$

- Entering the Multi-Tev region, the narrow width approximation starts taking effect.

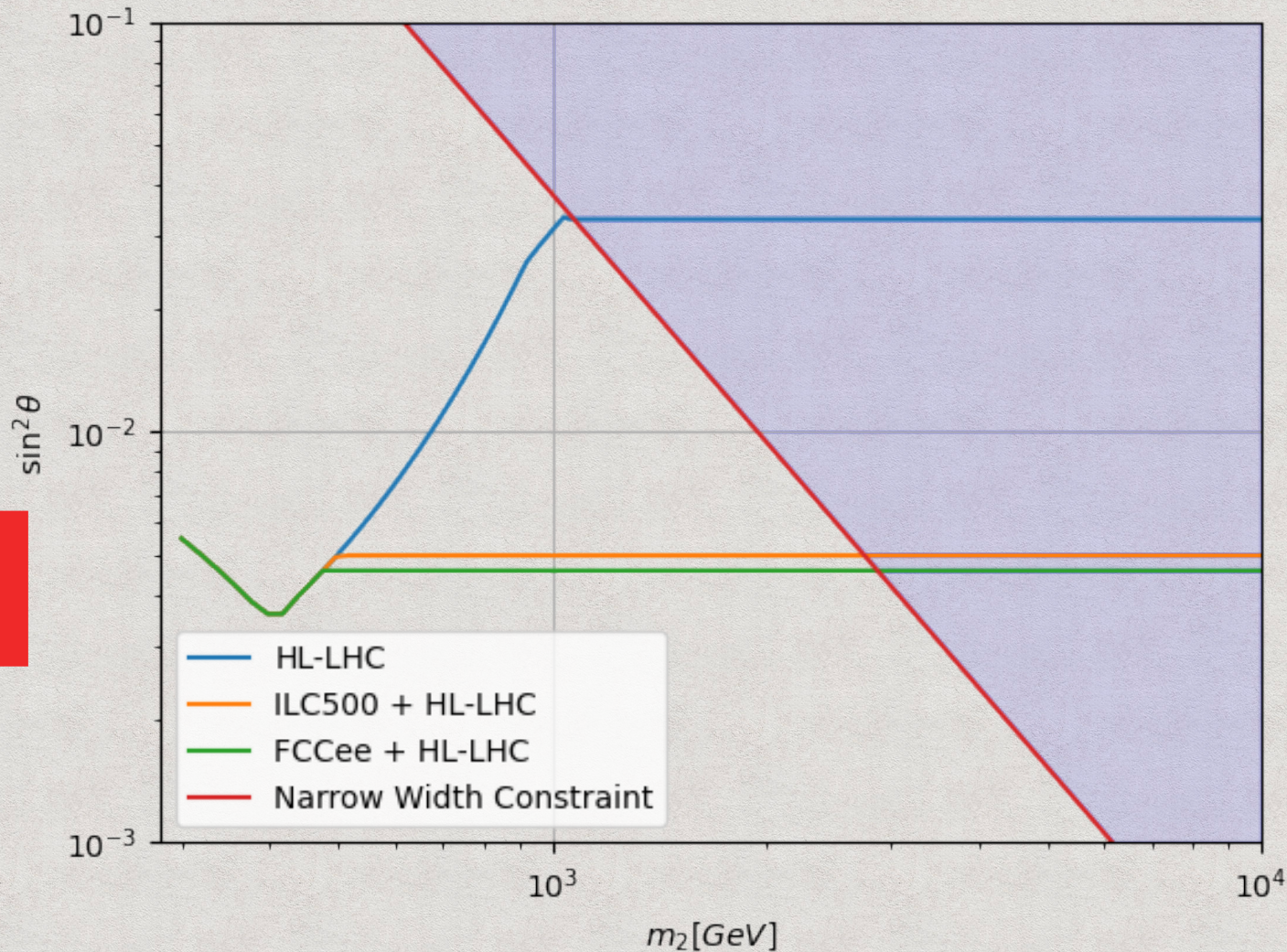
- Constraint in mixing angle:

$$\sin^2 \theta \lesssim \frac{0.2 \pi v_{\text{ew}}^2}{m_2^2}$$



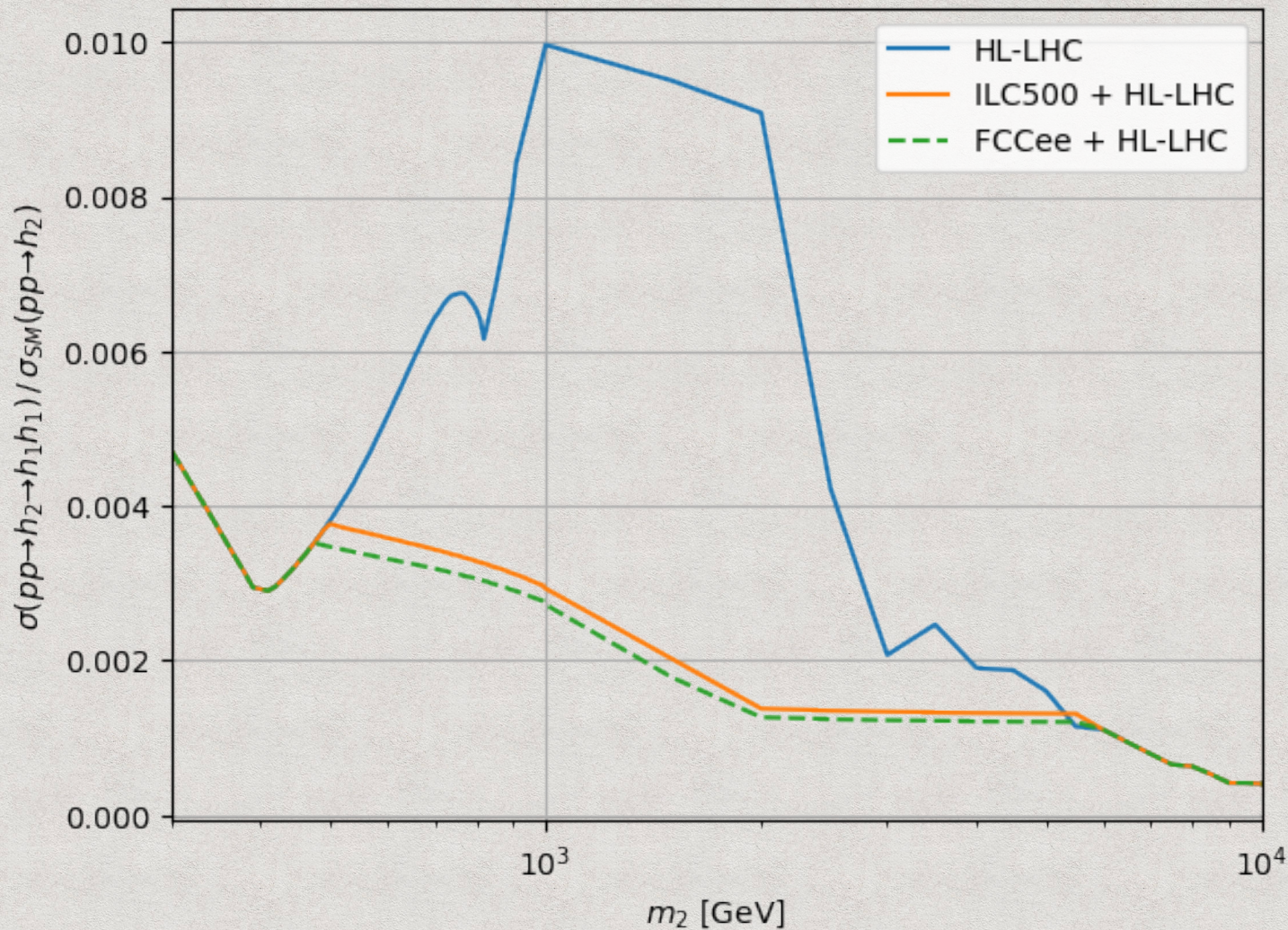
# Colliders benchmarks

arXiv:1910.11775



- Entering high energy region.
- After some point, the approximations takes over

$$\sin^2 \theta \lesssim \frac{0.2\pi v_{ew}^2}{m_2^2}$$



- Maximum production rates constraints

$$\sin^2 \theta_{BR}(h_2 \rightarrow h_1 h_1)$$

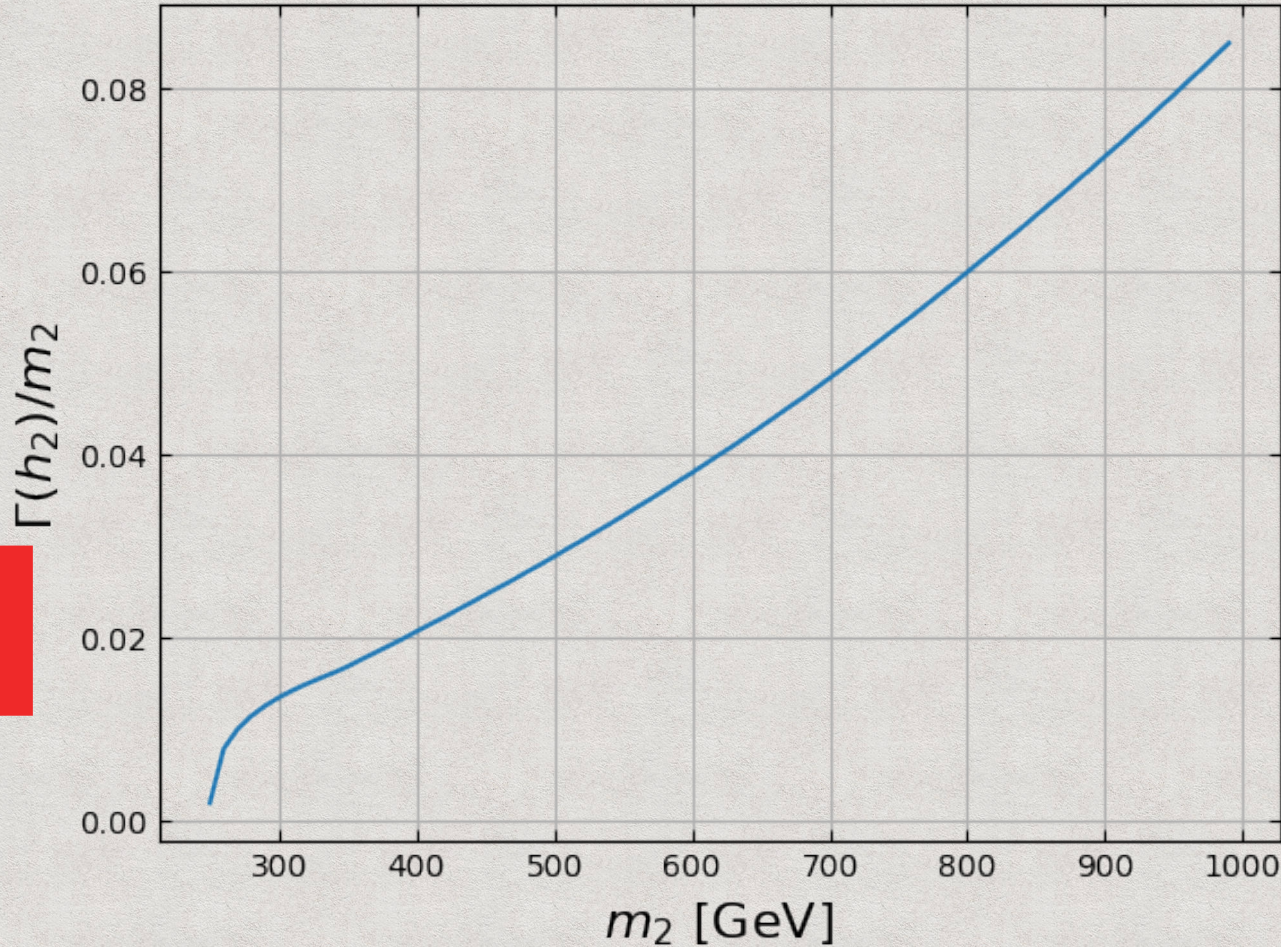
- **Finale:**
- Add real gauge singlet to model
- Identify free parameters and make scan
- Maximize production rate



**Coming soon...?**



**¡Muchas Gracias!**



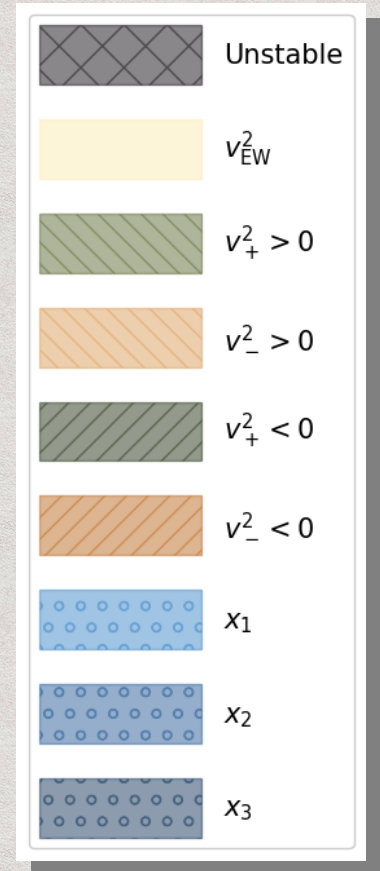
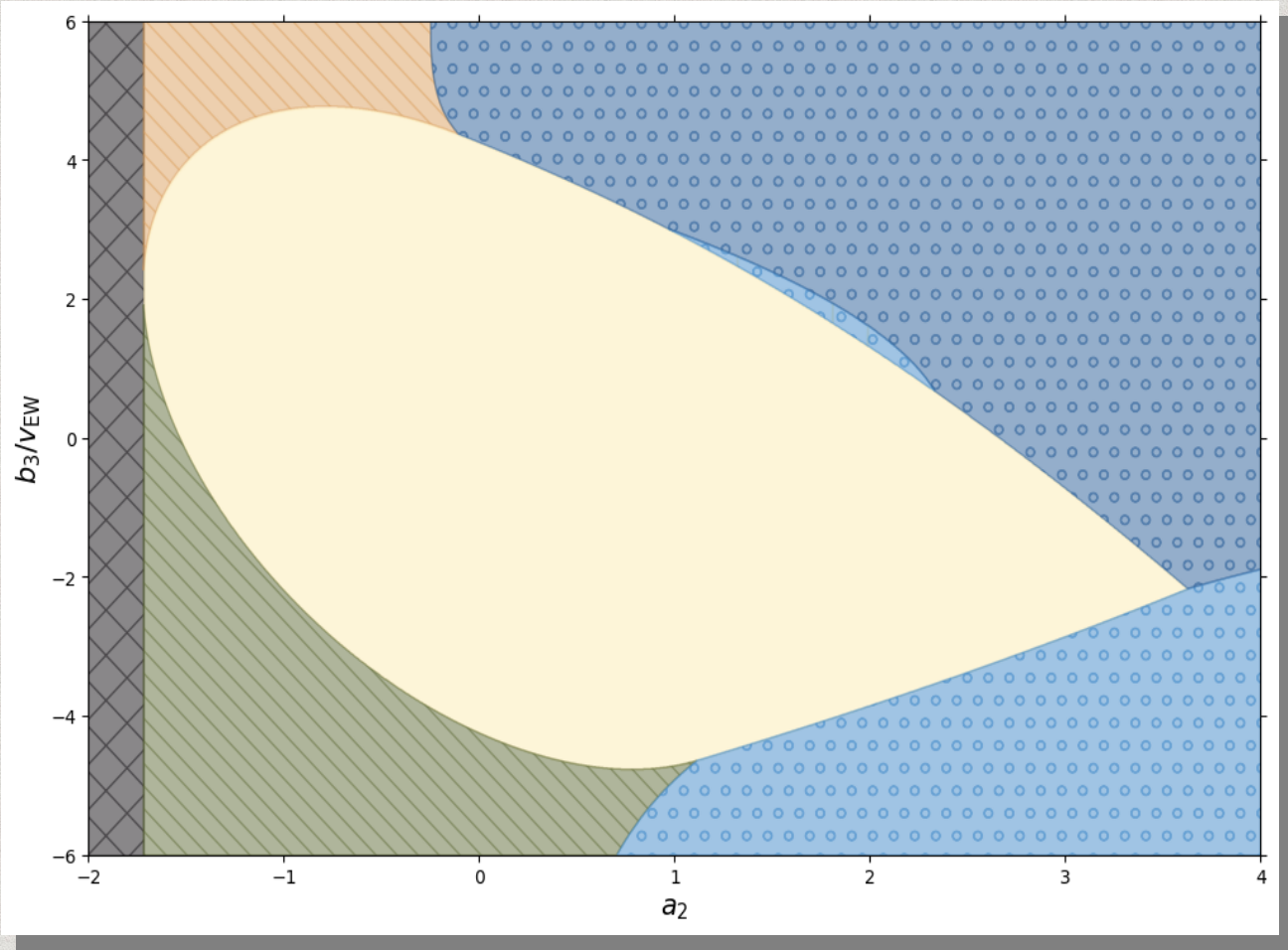
Ratio between maximum width and mass, sticking to the previous constraints.

Fair to use then

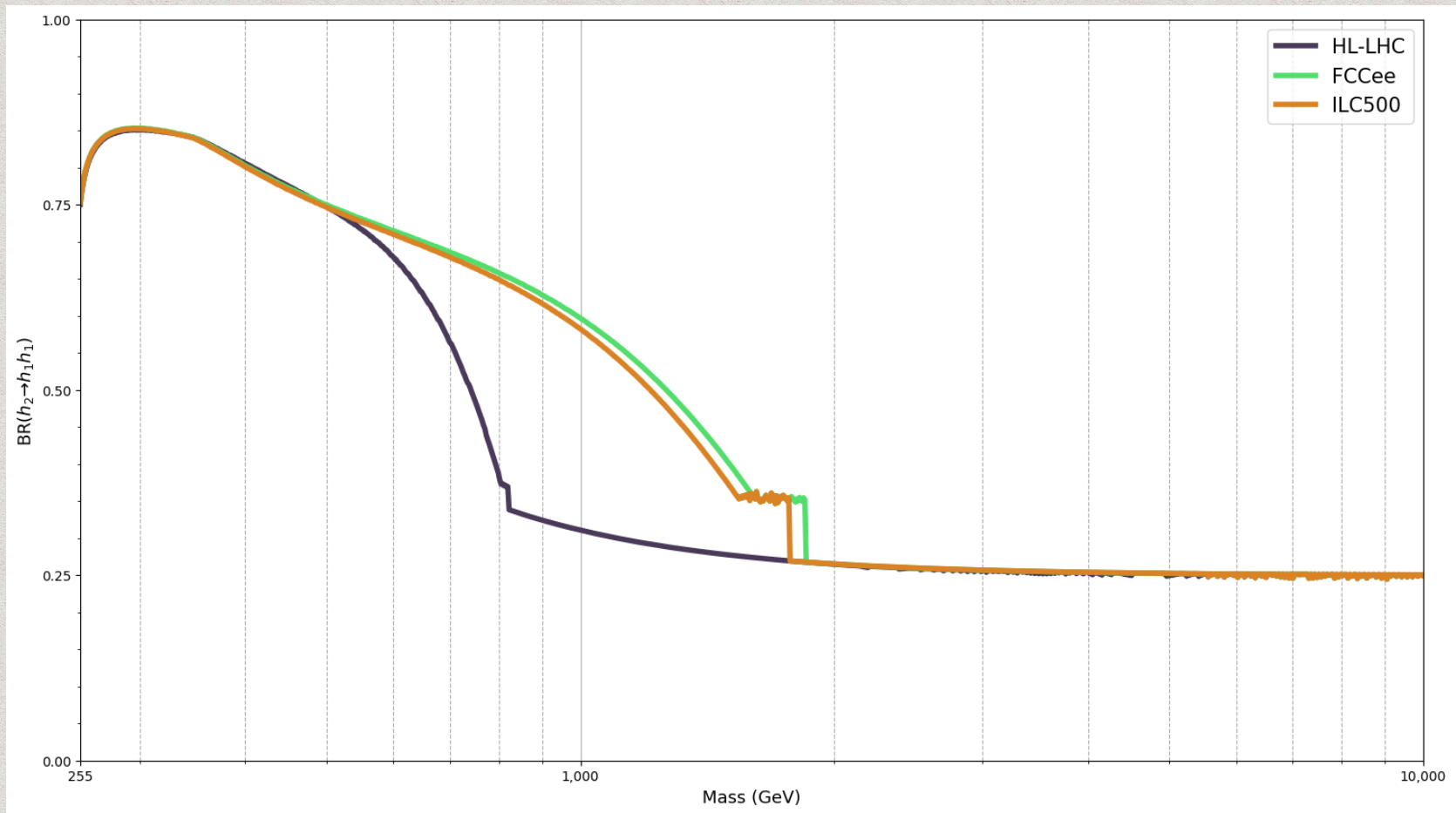
$$\Gamma(h_2) \leq 0.1m_2$$

Allowing to do Narrow width approximation

$$m_2 = 280\text{GeV}, \quad \sin \theta = 0.30$$



- Branching Ratios



By Jacob Scott



## FREE PARAMETERS

- From  $(v, x) = (v_{EW}, 0)$ , it is found

$$\mu^2 = \lambda v_{EW}^2, \quad b_1 = -\frac{v_{EW}^2}{4} a_1$$

- Rewrite in terms of the mass eigenstates. If  $U = (h \ S)$

$$V_m = \frac{1}{2} U M^2 U^T \quad \Rightarrow \quad \begin{pmatrix} h_1 \\ h_2 \end{pmatrix} = R(\theta)^T U^T,$$

- The next constraints can be found

$$a_1 = \frac{m_1^2 - m_2^2}{v_{EW}} \sin 2\theta, \quad \lambda = \frac{m_1^2 \cos^2 \theta + m_2^2 \sin^2 \theta}{2v_{EW}}.$$

$$b_2 = m_1^2 \sin^2 \theta + m_2^2 \cos^2 \theta - \frac{a_2 v_{EW}^2}{2},$$

- Partial width at tree level decay is given by

$$\Gamma(h_2 \rightarrow h_1 h_1) = \frac{\lambda_{211}^2}{32\pi m_2} \sqrt{1 - \frac{4m_1^2}{m_2}}, \quad \Rightarrow \quad m_2 \geq 2m_1$$

- From the scattering  $h_2 h_2 \rightarrow h_2 h_2$ , perturbative unitarity is used

$$\mathcal{M} = 16\pi \sum_i (2i + 1) a_i P_i(\cos \theta), \quad \lambda_{2222} = 6b_4 + \mathcal{O}(\theta^2)$$

- With restriction of  $|a_0| \leq 1/2$

$$a_0 = \frac{3b_4}{8\pi}, \quad \Rightarrow \quad b_4 \leq 4.2$$

## More Constraints

- Vacuum Stability yields

$$V^{(4)} = 4\lambda\phi_0^4 + 2a_2\phi_0^2s^2 + b_4s^4 > 0 \Rightarrow a_2 \geq -2\sqrt{\lambda b_4}.$$

- The following couplings terms will be used

$$V \supset \frac{\lambda_{211}}{2} h_2 h_1^2 + \frac{\lambda_{2222}}{4!} h_2^4$$

- First for  $h_2$  decay, and the second for limit in  $b_4$