

# The Constructive Method for Massive Particles in QED

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Lai, Liu & Terning [hep-ph/2312.11621](https://arxiv.org/abs/hep-ph/2312.11621)

# Scattering Amplitude Using Field Theory

Scattering of $n$ gluons	$n =$	3	4	5	6	7	...
	#diagrams =	1	3	10	38	154	...

Elvang & Huang [hep-th/1308.1697](https://arxiv.org/abs/hep-th/1308.1697)

Lagrangian involving magnetic charge (Zwanziger 1971):

$$\mathcal{L}_{\text{vis}} = -\frac{n^{\alpha}}{2n^2} \left[ n^{\mu} g^{\beta\nu} \left( F_{\alpha\beta}^A F_{\mu\nu}^A + F_{\alpha\beta}^B F_{\mu\nu}^B \right) - \frac{n_{\mu}}{2} \varepsilon^{\mu\nu\gamma\delta} \left( F_{\alpha\nu}^B F_{\gamma\delta}^A - F_{\alpha\nu}^A F_{\gamma\delta}^B \right) \right] \\ - e J_{\mu} A^{\mu} - \frac{4\pi}{e} K_{\mu} B^{\mu} ,$$

Terning & Verhaaren [hep-th/1809.05102](https://arxiv.org/abs/hep-th/1809.05102)

# Constructive Method

On-shell particles with complex momenta  $P^2 = m^2$ , four-momentum conserved.

back to real momenta

$$A_3 \times B_3 = A_3(p_1, p_2, P) \times B_3(p_3, p_4, P) \xrightarrow{\text{back to real momenta}} R(p_1, p_2, p_3, p_4)$$

$$\mathcal{M}_4 = \frac{R(p_1, p_2, p_3, p_4)}{P^2 - m^2}$$

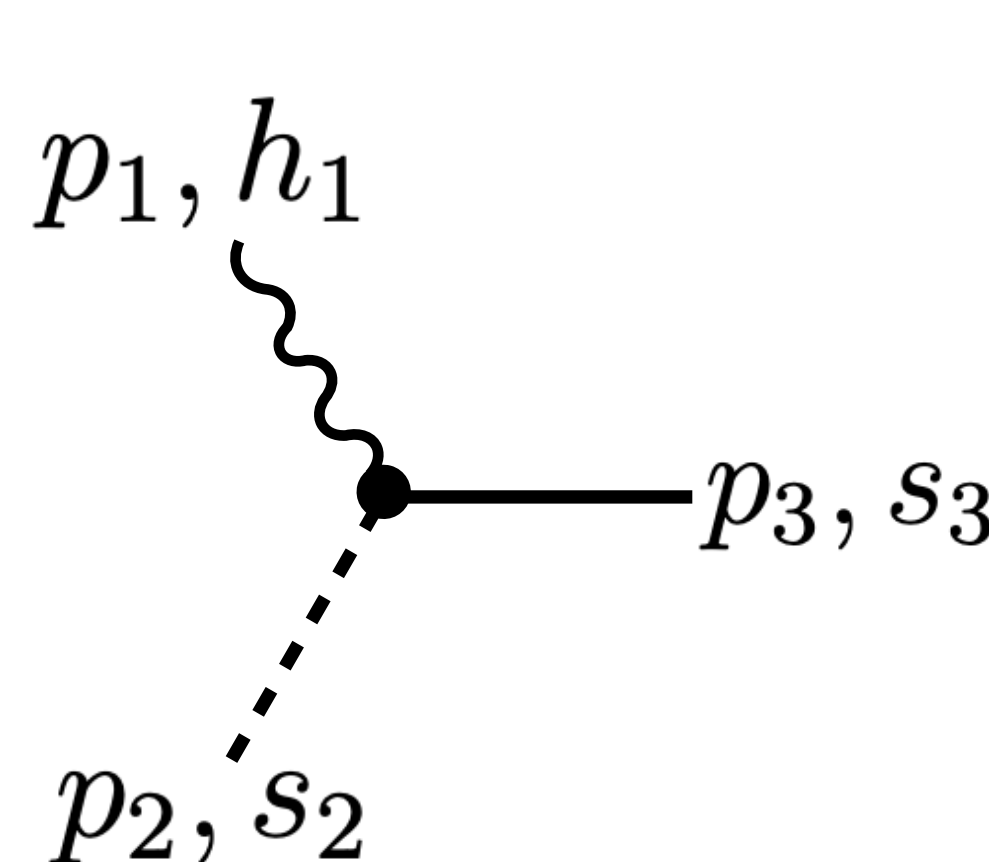
real  $P^2 \neq m^2$

Use 3-point amplitudes as building blocks, bypass field theory.

# Little Group Weight & Spinor-Helicity Variables

$$p^\mu (\sigma_\mu)_{a\dot{a}} = |p^I\rangle_a [p_I|\dot{a} = \varepsilon_{IJ} |p^I\rangle_a [p^J|\dot{a} \qquad |p^I\rangle [p_I| = |p^I\rangle (U_I^{\dagger J})(U_J^K)[p_K|$$

$$\text{Massless: } p^\mu (\sigma_\mu)_{a\dot{a}} = |p\rangle_a [p|\dot{a} \qquad |p\rangle [p| = |p\rangle \frac{1}{w} w [p|$$



$$|i\rangle \equiv |p_i\rangle \qquad |\mathbf{i}\rangle |\mathbf{i}\rangle \equiv |i^I\rangle |i^J] + |i^J\rangle |i^I] \text{ symmetrized}$$

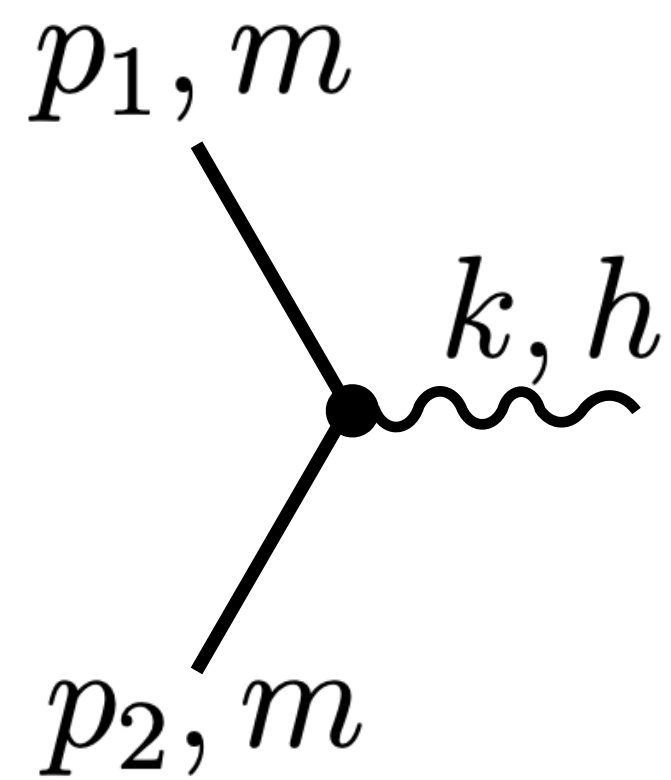
$$\sim |1\rangle^{-h_1 + \frac{n}{2}} |1]^{h_1 + \frac{n}{2}} |\mathbf{2}\rangle^{n_2} |\mathbf{2}]^{2s_2 - n_2} |\mathbf{3}\rangle^{n_3} |\mathbf{3}]^{2s_3 - n_3}$$

$$\varepsilon^{ab} |i\rangle_b |j\rangle_a \equiv \langle ij \rangle$$

$$\langle ij \rangle, [ij], [i|p|\mathbf{j}\rangle \quad \checkmark$$

$$[ij\rangle \quad \times$$

# Equal Mass, $x$ -factor

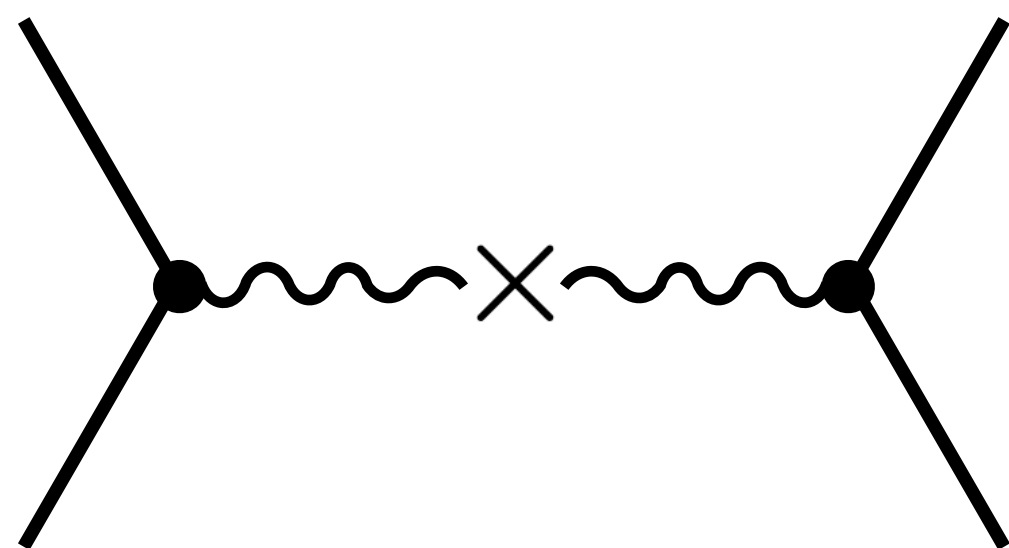


$|k\rangle \propto p_2|k] = -p_1|k]$  are not independent

$$x_{12} \equiv \frac{\langle q|p_2|k]}{m\langle qk]} = -\frac{\langle q|p_1|k]}{m\langle qk]} \quad \mathcal{M}_3 \sim x^h$$

$$\tilde{x}_{12} \equiv \frac{[\tilde{q}|p_2|k\rangle}{m[\tilde{q}k\rangle} = \frac{1}{x_{12}}$$

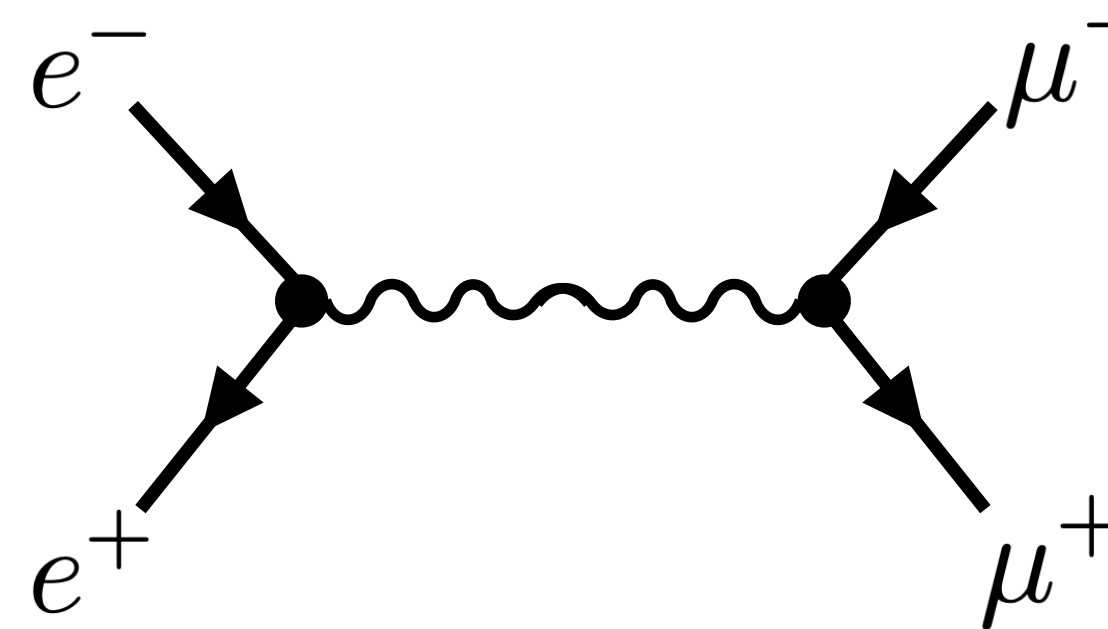
$\tilde{q}$  arbitrary reference spinor



gluing requires removing  $q$  and  $k$  from the expression

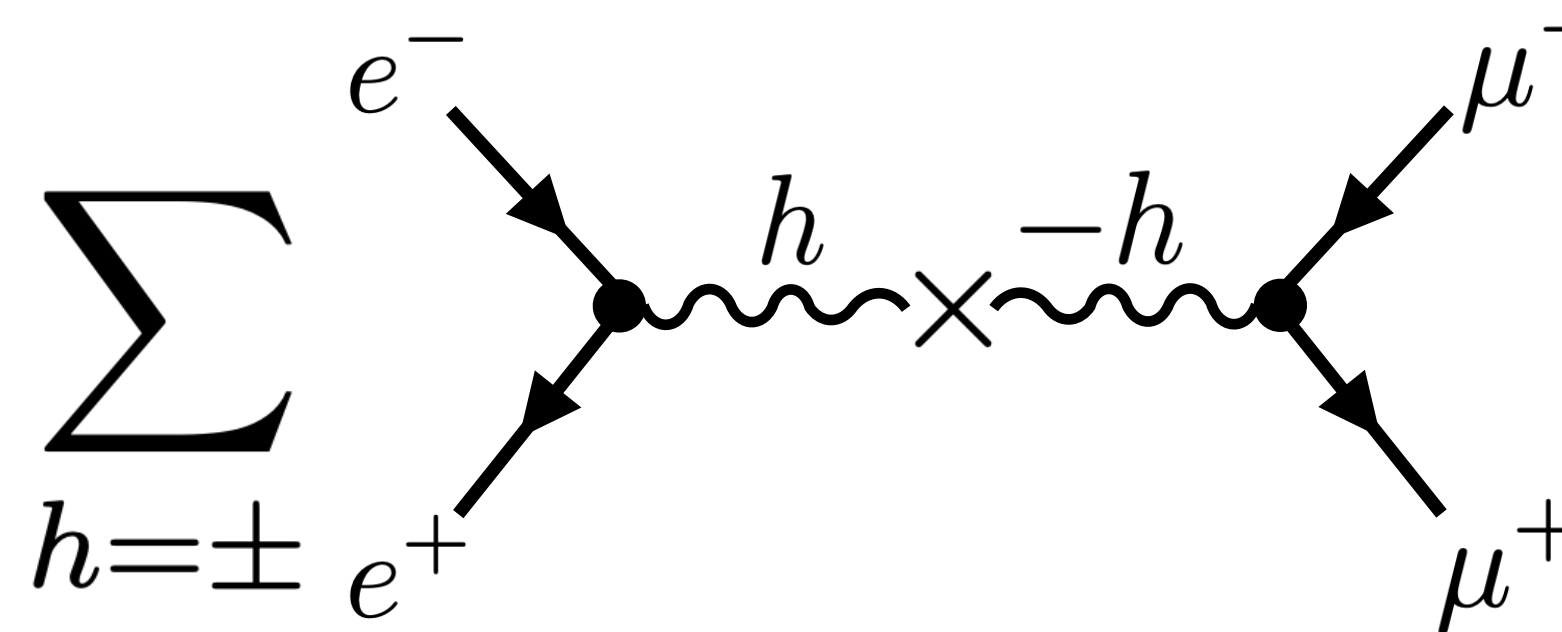
# Challenge with Internal Photon

Using Feynman Rule:



$$= \frac{e^2}{s} (\langle \mathbf{13} \rangle [\mathbf{24}] + [\mathbf{13}] \langle \mathbf{24} \rangle + [\mathbf{14}] \langle \mathbf{23} \rangle + \langle \mathbf{14} \rangle [\mathbf{23}])$$


Constructive Method:



$$\sum_{h=\pm} = \frac{e^2}{s} (x_{34} \tilde{x}_{12} [\mathbf{12}] \langle \mathbf{34} \rangle + x_{12} \tilde{x}_{34} \langle \mathbf{12} \rangle [\mathbf{34}])$$

# Challenge with Internal Photon

Christensen et al. *Nucl.Phys.B* 993 (2023) 116278 [hep-ph/2209.15018](https://arxiv.org/abs/hep-ph/2209.15018)



$$x_{34}\tilde{x}_{12}[\mathbf{12}]\langle\mathbf{34}\rangle + x_{12}\tilde{x}_{34}\langle\mathbf{12}\rangle[\mathbf{34}]$$

$$= \frac{1}{2m_em_\mu} \left[ (u - t + 2m_e^2 + 2m_\mu^2)[\mathbf{12}][\mathbf{34}] + 2([\mathbf{12}][\mathbf{3}|p_2p_1|4] + [\mathbf{1}|p_4p_3|2][\mathbf{34}]) \right]$$

Feynman Rule result

$$\neq [\mathbf{13}]\langle\mathbf{24}\rangle + [\mathbf{14}]\langle\mathbf{23}\rangle + [\mathbf{23}]\langle\mathbf{14}\rangle + [\mathbf{24}]\langle\mathbf{13}\rangle$$

Feynman Rule result  $\rightarrow \mathcal{O}(E^2)$

$$\frac{1}{2m_em_\mu} \left[ (u - t + 2m_e^2 + 2m_\mu^2)[\mathbf{12}][\mathbf{34}] + \dots \right] \rightarrow \mathcal{O}(E^4)$$

# Old Fashioned Perturbation Theory

Particles on-shell, spatial momentum conserved, energy not conserved.

$$\langle f | S | i \rangle = \langle f | H_{\text{int}} | i \rangle + \sum_n \frac{\langle f | H_{\text{int}} | n \rangle \langle n | H_{\text{int}} | i \rangle}{E_i - E_n} + \dots$$

$P^2 - m^2$  After summing over time-ordering

Equivalent to Feynman Rule Dyson [Phys. Rev. 75 \(1949\) 486](#)

For QED in Coulomb gauge,  $H_{\text{int}} = H_T + H_{\text{Coul}}$

$$H_T = - \int d^3 \mathbf{x} \mathbf{J} \cdot \mathbf{A} \quad H_{\text{Coul}} = \frac{1}{2} \int d^3 \mathbf{x} d^3 \mathbf{y} \frac{J^0(\mathbf{x}) J^0(\mathbf{y})}{4\pi |\mathbf{x} - \mathbf{y}|}$$

Does not contribute to residue when  $s = 0$



# Old Fashioned Perturbation Theory

$$H_T \sim \begin{array}{c} 2 \\ \swarrow \\ \bullet \\ \searrow \\ 1 \\ \uparrow \\ k \uparrow \\ \downarrow \\ h \end{array} = \frac{e}{\sqrt{2\omega_{\mathbf{k}}}} \bar{v}_2 \not{\epsilon}_h u_1 \xrightarrow{s=0} \begin{cases} \frac{e}{\sqrt{\omega_{\mathbf{k}}}} \frac{1}{\langle kq_+ \rangle} (\langle \mathbf{2}q_+ \rangle [k\mathbf{1}] + [\mathbf{2}k] \langle q_+\mathbf{1} \rangle), & h = + \\ \frac{e}{\sqrt{\omega_{\mathbf{k}}}} \frac{1}{[kq_-]} (\langle \mathbf{2}k \rangle [q_-\mathbf{1}] + [\mathbf{2}q_-] \langle k\mathbf{1} \rangle), & h = - \end{cases}$$

Schouten identity


$$|i\rangle \langle jk\rangle + |j\rangle \langle ki\rangle = |k\rangle \langle ji\rangle = (\langle \mathbf{2}q \rangle \langle \mathbf{1}| + \langle q\mathbf{1} \rangle \langle \mathbf{2}|) \frac{p_2 |k\rangle}{m_e \langle qk\rangle}$$

$$\xrightarrow{\text{Schouten}} = \frac{\langle q| p_2 |k\rangle}{m_e \langle qk\rangle} \langle \mathbf{1}\mathbf{2} \rangle = x_{12} \langle \mathbf{1}\mathbf{2} \rangle$$

**Constructive Method Amplitude!**

# Challenge with Internal Photon

Christensen et al. *Nucl.Phys.B* 993 (2023) 116278 [hep-ph/2209.15018](https://arxiv.org/abs/hep-ph/2209.15018)



$$\begin{aligned}
 & x_{34} \tilde{x}_{12} [\mathbf{12}] \langle \mathbf{34} \rangle + x_{12} \tilde{x}_{34} \langle \mathbf{12} \rangle [\mathbf{34}] \\
 &= \frac{1}{2m_e m_\mu} \left[ (u - t + 2m_e^2 + 2m_\mu^2) [\mathbf{12}] [\mathbf{34}] + 2([\mathbf{12}] [\mathbf{3} | p_2 p_1 | \mathbf{4}] + [\mathbf{1} | p_4 p_3 | \mathbf{2}] [\mathbf{34}]) \right] \\
 \text{Our work} & \quad \underbrace{\hspace{10em} \text{Feynman Rule result} \hspace{10em}} \\
 &= [\mathbf{13}] \langle \mathbf{24} \rangle + [\mathbf{14}] \langle \mathbf{23} \rangle + [\mathbf{23}] \langle \mathbf{14} \rangle + [\mathbf{24}] \langle \mathbf{13} \rangle + \frac{s}{2m_e m_\mu} ([\mathbf{12}] [\mathbf{34}] - 2[\mathbf{14}] [\mathbf{23}])
 \end{aligned}$$

On-shell  $s = (p_1 + p_2)^2 = k^2 = 0$ , should drop the  $\mathcal{O}(s)$  term.

⇒ Constructive Method works!

3 months after our paper: Ema et al. [hep-ph/2403.15538](https://arxiv.org/abs/hep-ph/2403.15538) reproduced same result

# Conclusion

Challenge of internal photon in constructive method has been resolved.

OFPT can provide useful insight to on-shell constructive calculation.

