

# Imprints of Early Universe Cosmology on Gravitational waves

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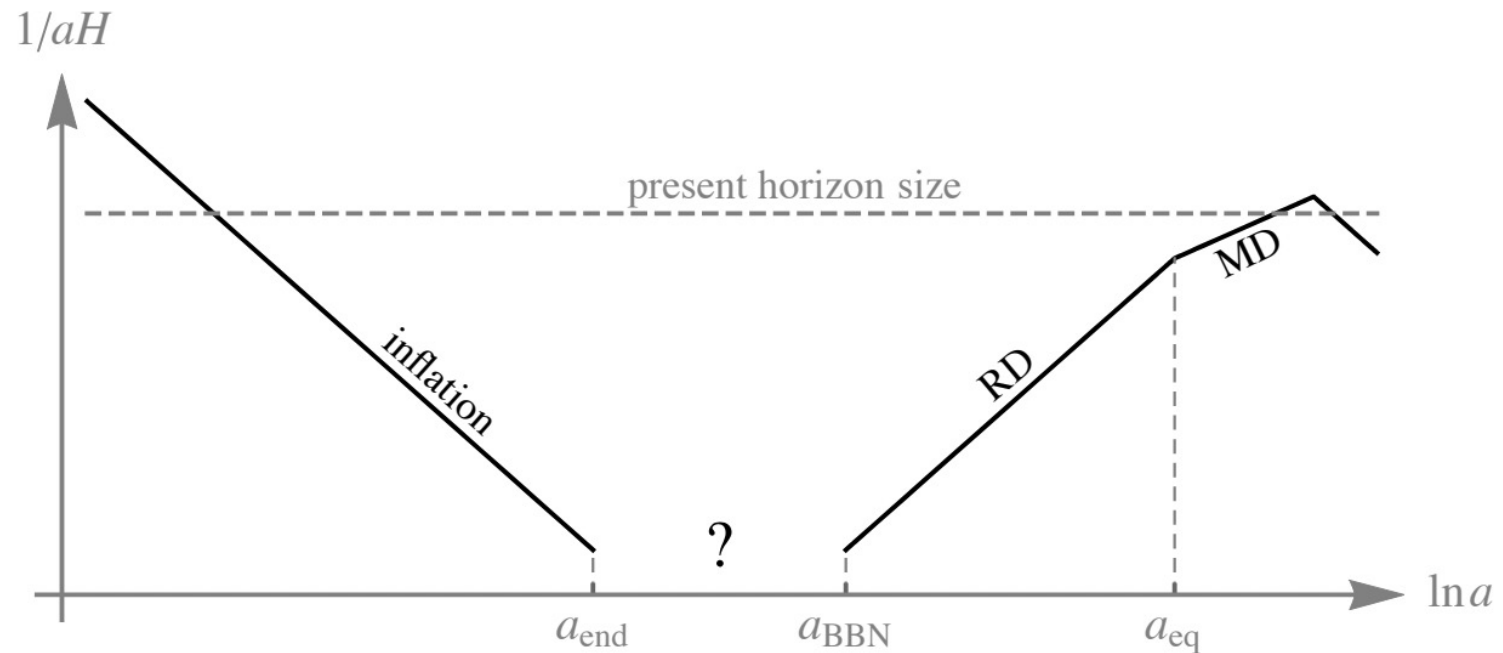
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# Motivation

- Physics before BBN ( $T \sim \text{MeV}$ ) is not well understood due to lack of observational data.
- Gravitational waves can be a natural way to probe this epoch between end of inflation and BBN.



# Introduction

- ▶ Energy/entropy injection before BBN has been discussed extensively:
  - ▶ Fluctuations generated during inflation and later reentry [Carr & Lidsey, ....]
  - ▶ Collapse of domain walls [Cai et al, ...]
  - ▶ PBH reheating [Bernal et al, ...]
  - ▶ Bubble collisions during phase transition [Kodama et al, ...]
  - ▶ Temperature increase during reheating [Co et al, ...]
- ▶ The rate of energy injection can be either be fast where the field remains stuck as the temperature rises or can be slow where the field tracks its T dependent minima.
- ▶ We consider its effects on the GW spectrum for a wide range of FOPT in hidden sectors.

# Cosmological setup

- ▶ We consider the scenario where the hidden sector is thermally decoupled to the SM.
- ▶ We assume that the SM makes up bulk of the energy density of the universe.
- ▶ The ratio of hidden sector temperature and that of SM is given by  $\xi = \frac{T_h}{T_{SM}} < 1$
- ▶ Any small change in the energy density of the universe will thus have more impact on hidden sector as compared to SM and the Hubble will be unaffected.
- ▶ Net energy density of the universe is given as,  $\rho_R(T) = \frac{\pi^2}{30} \left( g_h^*(T) + \frac{g_{SM}^*(T_{SM})}{\xi^4} \right) T^4$

# Model realization

$$V \approx D(T^2 - T_0^2)\phi^2 - ET\phi^3 + \frac{\lambda}{4}\phi^4$$

- Initially, at high  $T$ , the field is in symmetric phase and there's just 1 minima at  $\phi = 0$
- As universe cools,  $T < T_1$ , there exist a second minima

$$T_1^2 = \frac{T_0^2}{1 - \frac{9E^2}{8\lambda D}}, \quad \phi_1 = \frac{3ET_1}{2\lambda}$$

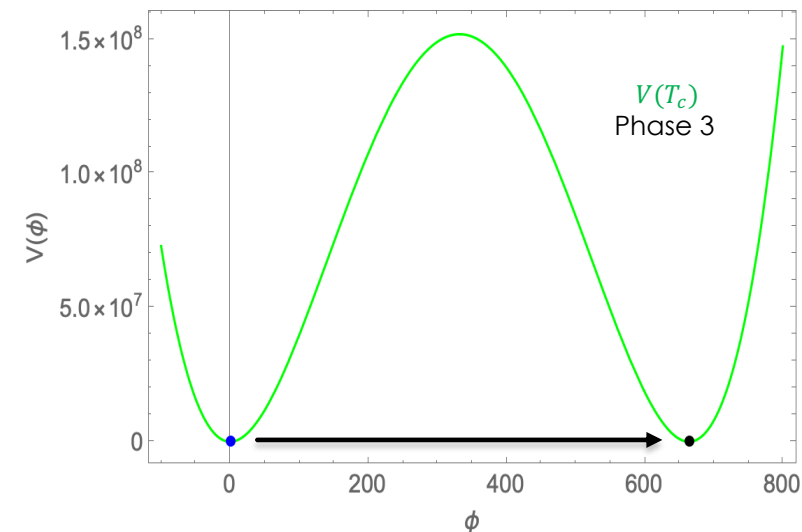
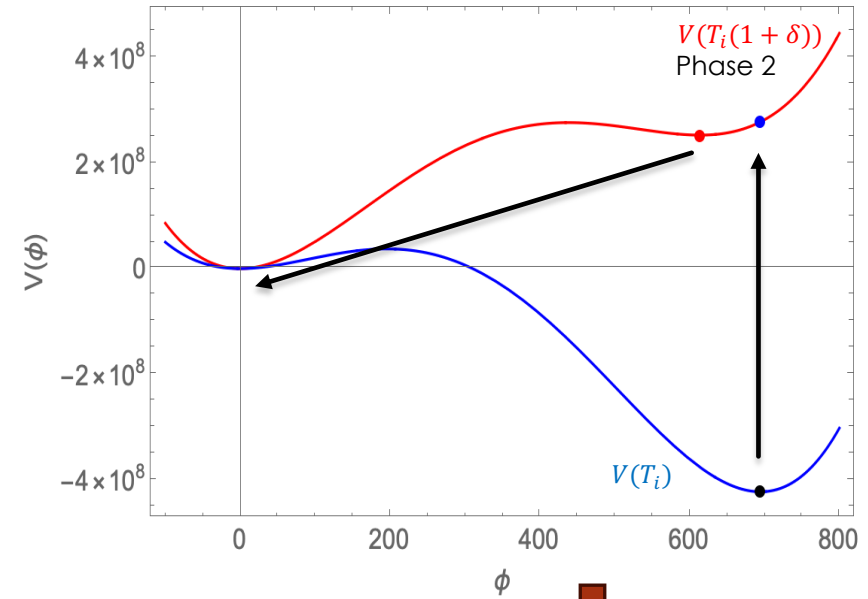
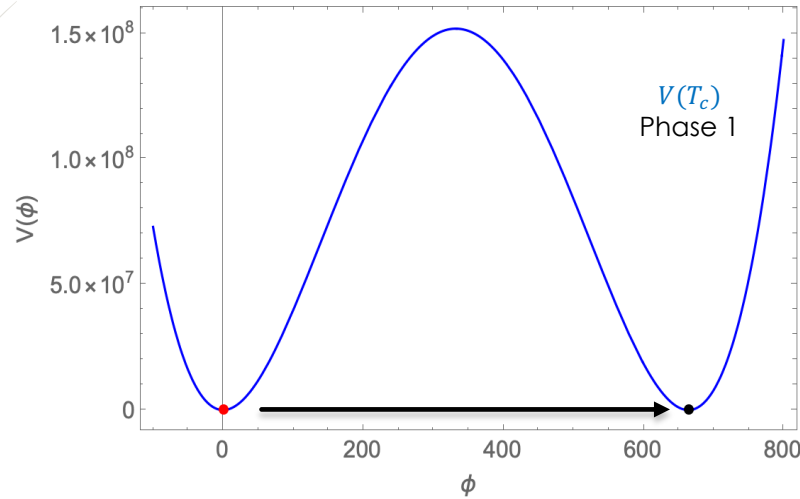
- As it further cools, these two minima become equi-potential and we have an onset of phase transition,

$$V(0, T_c) = V(\phi_c, T_c) \quad T_c^2 = \frac{T_0^2}{1 - \frac{E^2}{\lambda D}}, \quad \phi_c = \frac{2ET_c}{\lambda}$$

- After  $T = T_0$ ,  $\phi = 0$  ceases to be a minima and we are left with,

$$\phi_0 = \frac{3ET_0}{\lambda}$$

# Thermal kick effects



- First transition happens at  $T = T_c$  (Phase 1)
- Due to thermal kick at  $T_i$ , ( $T_c > T_i > T_0$ ),  $T_i \rightarrow T_i(1 + \delta) > T_c$  whereas the field remains stuck at  $\phi_i(T_i)$ , leads to PT from  $\phi_i \rightarrow 0$  (Phase 2)
- As universe cools down, there's another PT from  $0 \rightarrow \phi_c$ , which is like the standard transition but happens at later redshift (Phase 3)

# Strength of GW signal

- Amplitude of GW signal is controlled by strength parameter,  $\alpha$  given as

$$\alpha = \frac{\Delta(V - \frac{1}{4}\partial_T V)}{\rho_R} \Big|_{T=T_N} \quad \rho_R(T_N) = \frac{\pi^2}{30} (g_h^* + \frac{g_{SM}^*}{\xi^4}) T_N^4$$

- For the standard FOPT, we can simplify to get,

$$\alpha|_{0 \rightarrow \phi_c} = \alpha_c \approx \frac{\phi_c^2 (-2 D T_0^2)}{4 \rho_R(T_c)} = \mu^2 \frac{-\phi_c^2}{4 \rho_R(T_c)}$$

- For the PT due to kick, we have

$$\left| \frac{\alpha_i}{\alpha_c} \right| \approx (1 + \delta)^2 \left( \frac{\phi_{min}(\tilde{T}_i) T_c^2}{\phi_c(T_c) \tilde{T}_i^2} \right)^2 \frac{g_{SM}^*(T_c/\xi)}{g_{SM}^*(T_i/\xi)} > 1$$

# Duration of PT

- ▶ This parameter gives a measure of the duration of PT
- ▶ Defined in terms of the Euclidean bounce action as,

$$\frac{\beta}{H_*} = T \frac{d(S_3/T)}{dT} \Big|_{T_*}$$

- ▶ The Euclidean Bounce action is given via,

$$\frac{S_3}{T} = \frac{4.85 M^3}{E^2 T^3} f(\kappa)$$

$$f(\kappa) = 1 + \frac{\kappa}{4} \left( 1 + \frac{2.4}{1-\kappa} + \frac{0.26}{(1-\kappa)^2} \right)$$

$$M^2 = 2D(T^2 - T_0^2), \quad \kappa = \frac{\lambda D}{E^2} \left( 1 - \frac{T_0^2}{T^2} \right)$$

- ▶ For the PT due to kick, the parameters get modified via,

$$\tilde{M}^2(T) = M^2(T) + 3\phi_i(2ET + \lambda\phi_i)$$

$$\tilde{E}T = ET + \lambda\phi_i$$



# $T_N$ and $v_w$

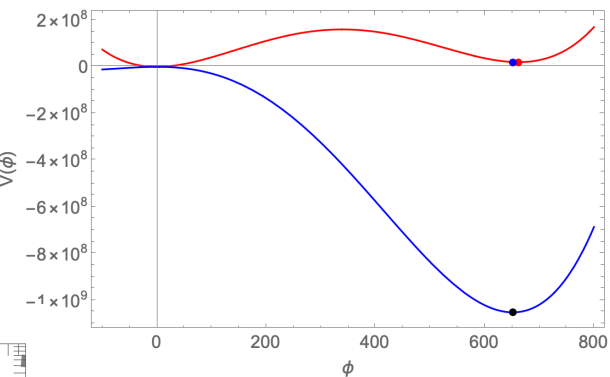
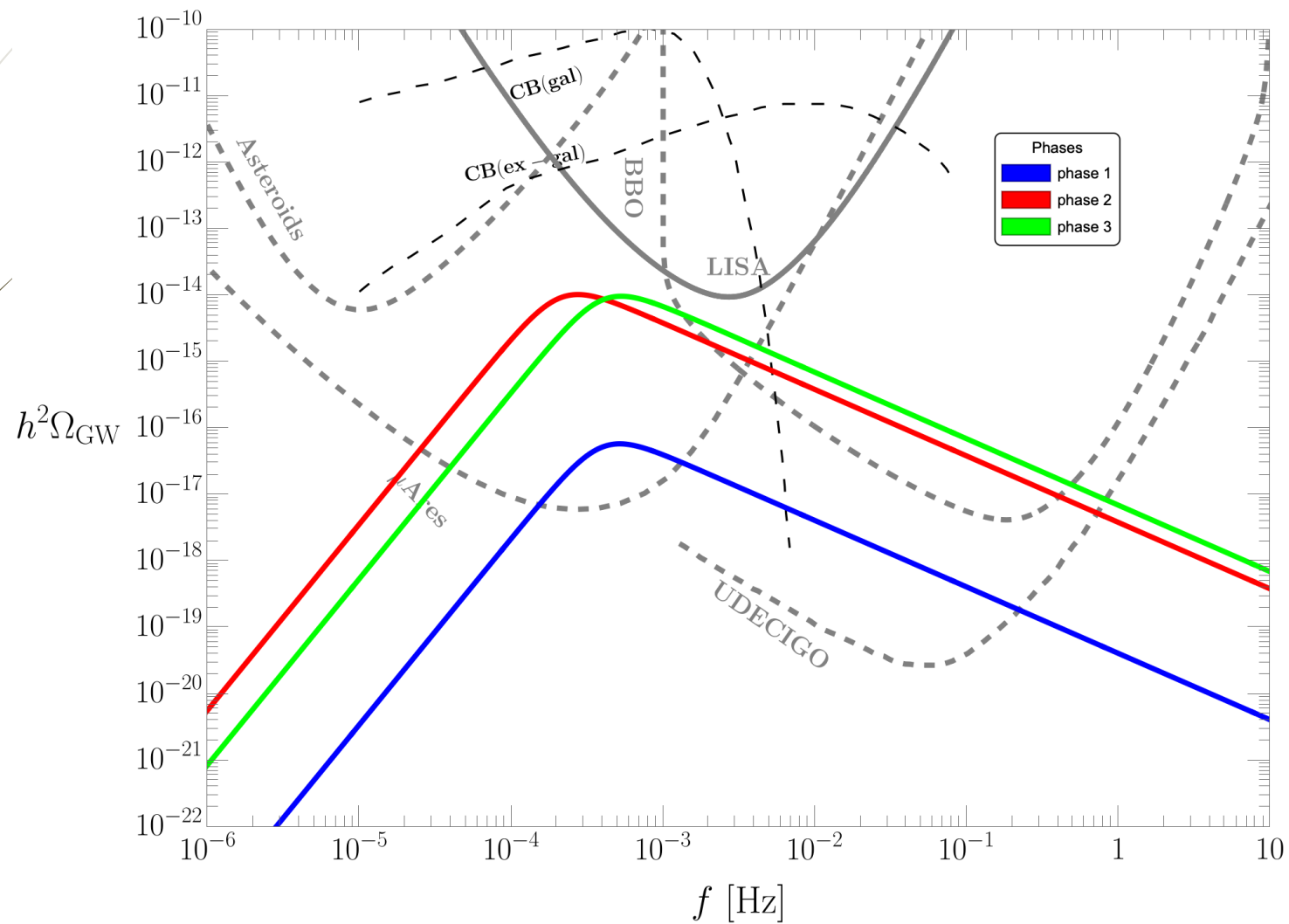
- ▶ Dedicated numerical simulations are needed to calculate the  $T_N$  and  $v_w$
- ▶ We will use analytical approximations for  $T_N$ :

$$\frac{S_3}{T_n} \approx 4 \log \frac{T_n}{H}$$

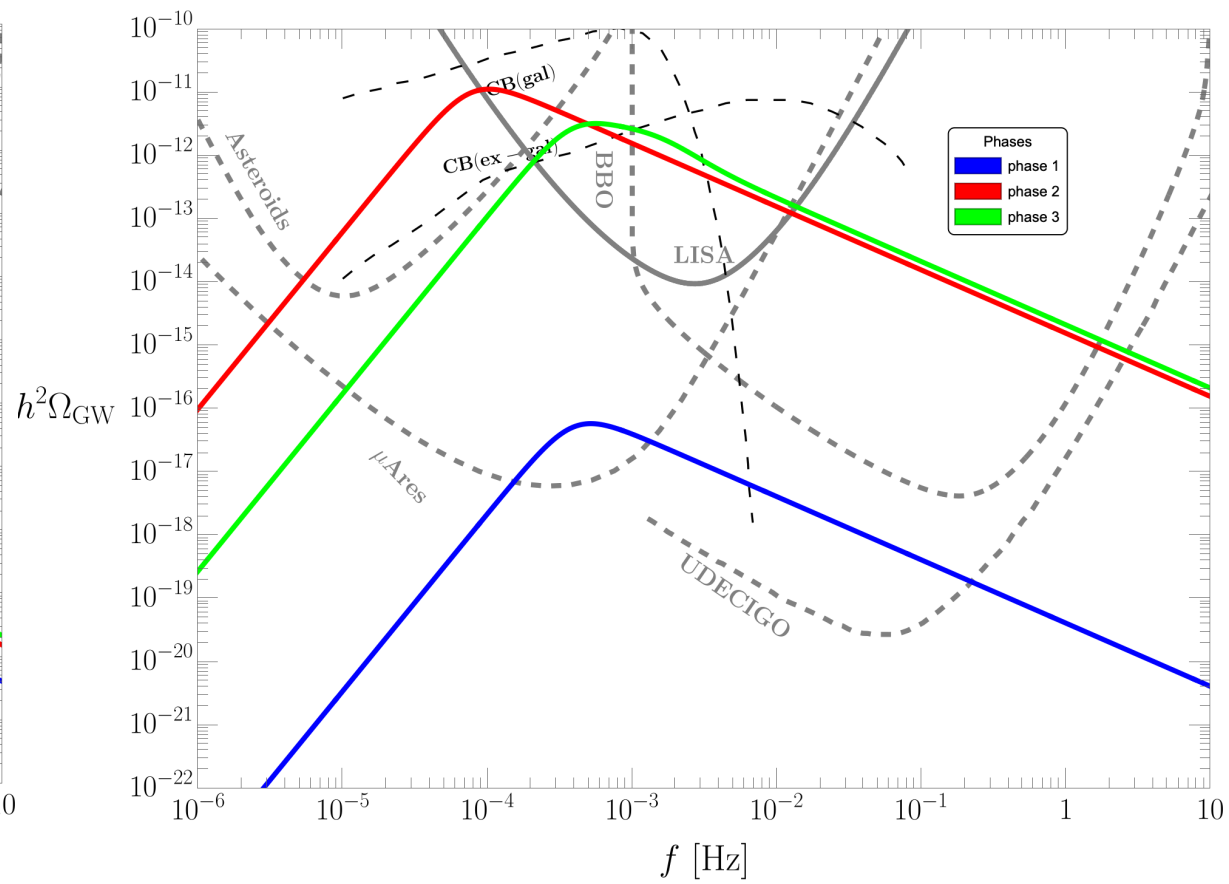
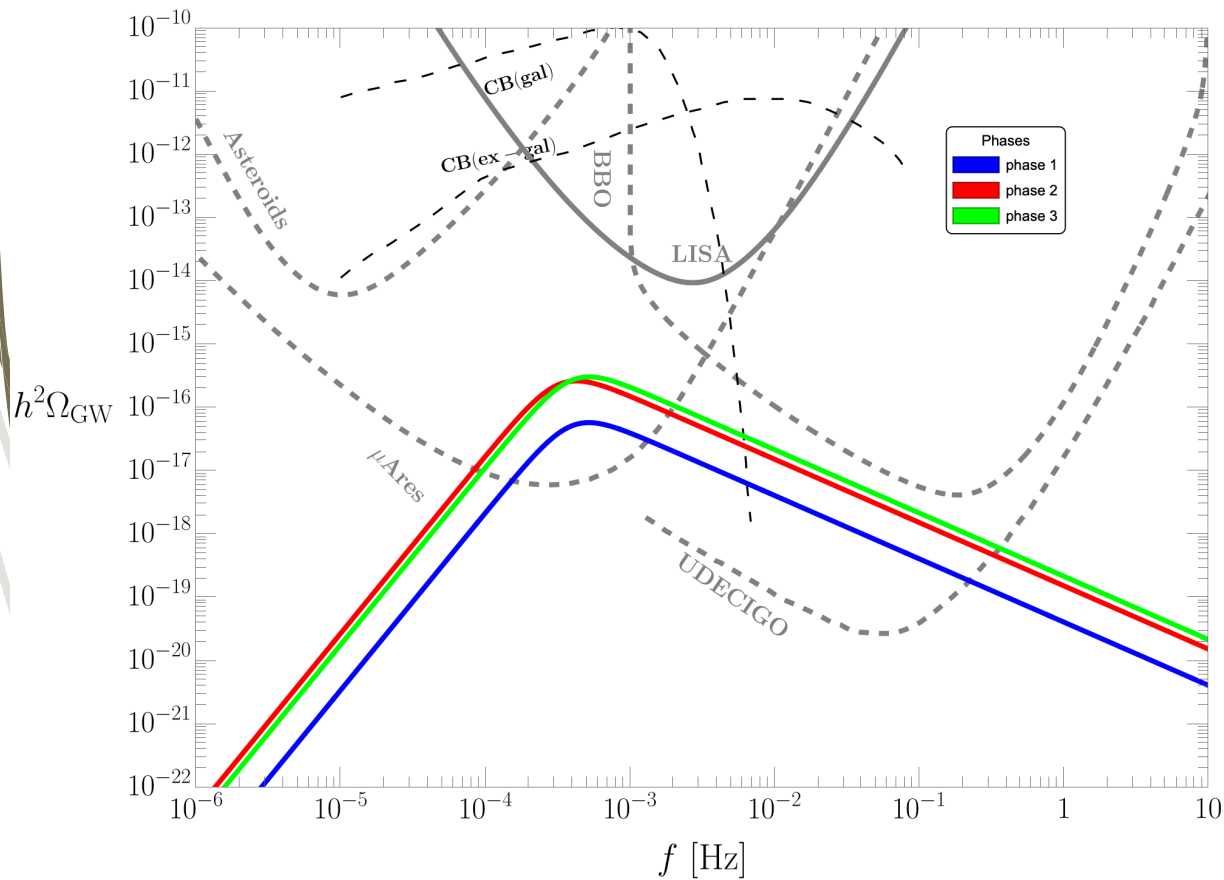
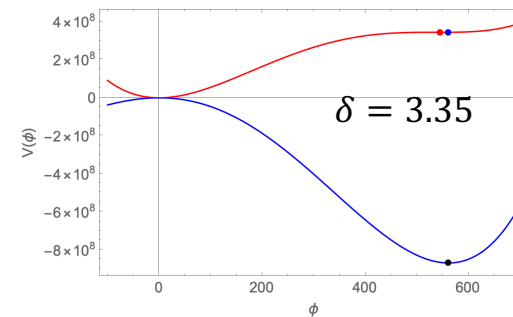
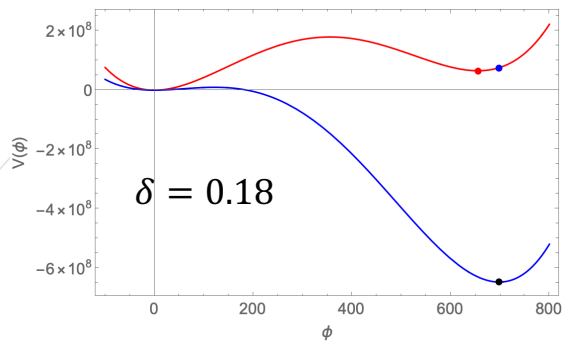
- ▶ For wall velocity, a simple demarcation can be made where weaker FOPT attains terminal velocity and for stronger transition, it can overcome friction and wall becomes ultra-relativistic[1]

$$v_w \approx \begin{cases} \frac{1}{\sqrt{3}}, & \alpha \lesssim 10^{-2} \\ 1, & \alpha \gtrsim 10^{-2} \end{cases}$$

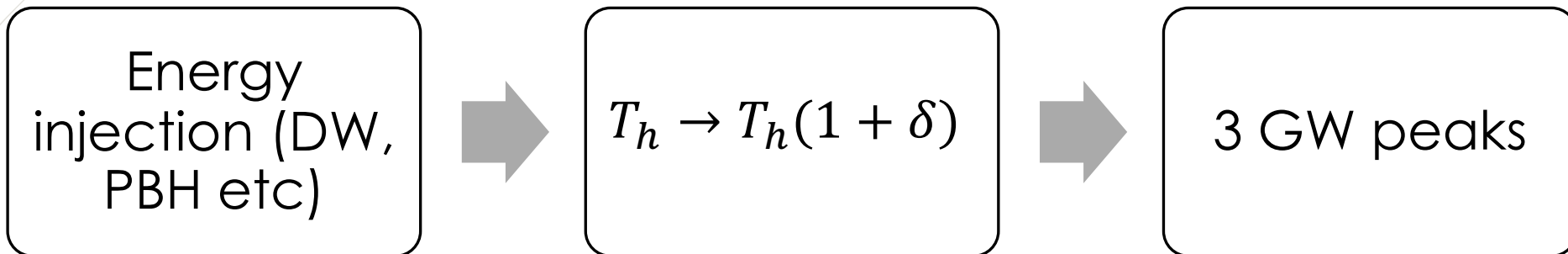
# GW spectrum ( $\mu = 100 \text{ GeV}, \delta = 0.92$ )



# $\delta$ dependence



# Conclusion



- ▶ Energy injection leads to more than one peak frequencies for GW from FOPT in hidden sector.
- ▶ It is fairly independent w.r.to the mass scale of the hidden sector.
- ▶ Even for QCD like transitions, we expect to have multiple peaks due to kick.
- ▶ Hidden sectors with GW can probe a variety of new physics scenario in the pre-BBN era.

THANK YOU!

# BACKUP Slides

# Hidden sector

- For concreteness, we consider a scalar field with U(1) gauge symmetry and a Yukawa like coupling to fermion field,

$$\mathcal{L} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} + \frac{1}{2}D_{\mu}\phi D^{\mu}\phi + i\bar{\psi}\not{D}\psi - \frac{y\phi}{\sqrt{2}}\bar{\psi}\psi - V(\phi)$$

- The tree and thermal potential are given as

$$V_0 = -\frac{1}{2}\mu^2\phi^2 + \frac{1}{4}\lambda\phi^4$$

$$V_{th} = \frac{T^4}{2\pi^2} \left( n_{\phi} J_B \left[ \frac{m_{\phi}^2}{T^2} \right] + n_X J_B \left[ \frac{m_X^2}{T^2} \right] - n_f J_F \left[ \frac{m_f^2}{T^2} \right] \right)$$

# High T Potential

- At high temperatures, the effective potential is given as,

$$V \approx D(T^2 - T_0^2)\phi^2 - ET\phi^3 + \frac{\lambda}{4}\phi^4$$

where,

$$D = \frac{\alpha}{2}, \quad \alpha = \frac{\lambda + g^2}{4} + \frac{y^2}{24}$$

$$E = \frac{1}{12\pi} \left( 3g^3 + \left( 3\lambda + \frac{\alpha T^2 - \mu^2}{\phi^2} \right)^{3/2} \right) \approx \frac{g^3}{4\pi}$$

$$T_0^2 = \frac{\mu^2}{2D}$$



# Gravitational Waves signal

- Differential GW density parameter characterizes them :

$$\Omega_{GW} = \frac{1}{\rho_c} \frac{d\rho_{GW}}{d \log f}, \quad \rho_c = 3M_{pl}^2 H^2$$

- Semi-analytical parametrizations can be used to describe them,

$$\Omega_{GW}^{\text{em}}(f_{\text{em}}) = \sum_{I=\text{BW, SW}} N_I \Delta_I(v_w) \left( \frac{\kappa_I(\alpha_{-1}) \alpha_{\text{tot}}}{1 + \alpha_{\text{tot}}} \right)^{p_I} \left( \frac{H}{\beta} \right)^{q_I} s_I(f_{\text{em}}/f_{p,I}).$$

$$\longrightarrow h^2 \Omega_{GW}^0(f) = h^2 \mathcal{R} \Omega_{GW}^{\text{em}} \left( \frac{a_0}{a_{\text{perc}}} f \right).$$

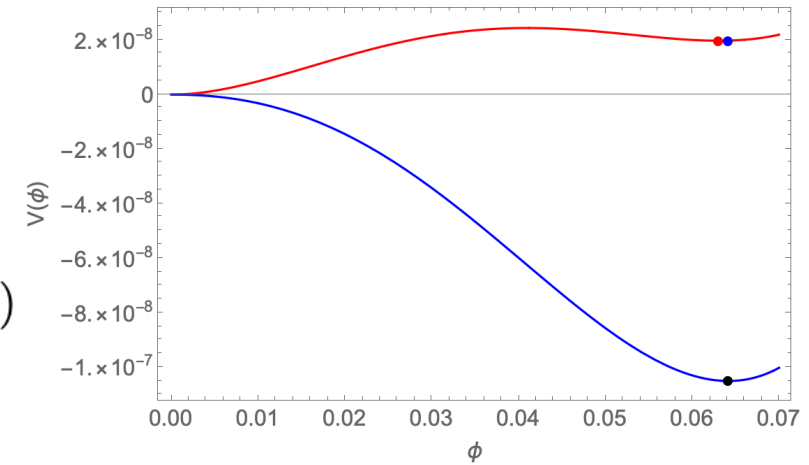
# GW parameters for PT Type 2

- ▶ The difference is due to phase transition between minimas which are not equipotential.
- ▶ Effective potential near transition is given by,

$$V \approx \frac{\tilde{M}^2(T)}{2} \phi^2 - \tilde{E} T \phi^3 + \frac{\lambda}{4} \phi^4$$

$$\tilde{M}^2(T) = M^2(T) + 3\phi_i(2ET + \lambda\phi_i)$$

$$\tilde{E}T = ET + \lambda\phi_i$$



- ▶ The bounce action and the transition strength become,

$$\frac{S_3}{T} = \frac{4.85 \tilde{M}^3(T)}{\tilde{E}^2 T^3} f(\tilde{\kappa}) \quad \left| \frac{\alpha_i}{\alpha_c} \right| \approx (1 + \delta)^2 \left( \frac{\phi_{min}(\tilde{T}_i) T_c^2}{\phi_c(T_c) \tilde{T}_i^2} \right)^2 \frac{g_{SM}^*(T_c/\xi)}{g_{SM}^*(T_i/\xi)} > 1$$

# GW signal parametrization

- Normalization factors and exponents :

$$(N_{\text{BW}}, N_{\text{SW}}) = (1, 0.159) \quad (p_{\text{BW}}, p_{\text{SW}}) = (2, 2) \quad (q_{\text{BW}}, q_{\text{SW}}) = (2, 1)$$

- Potential suppression due to wall velocity :

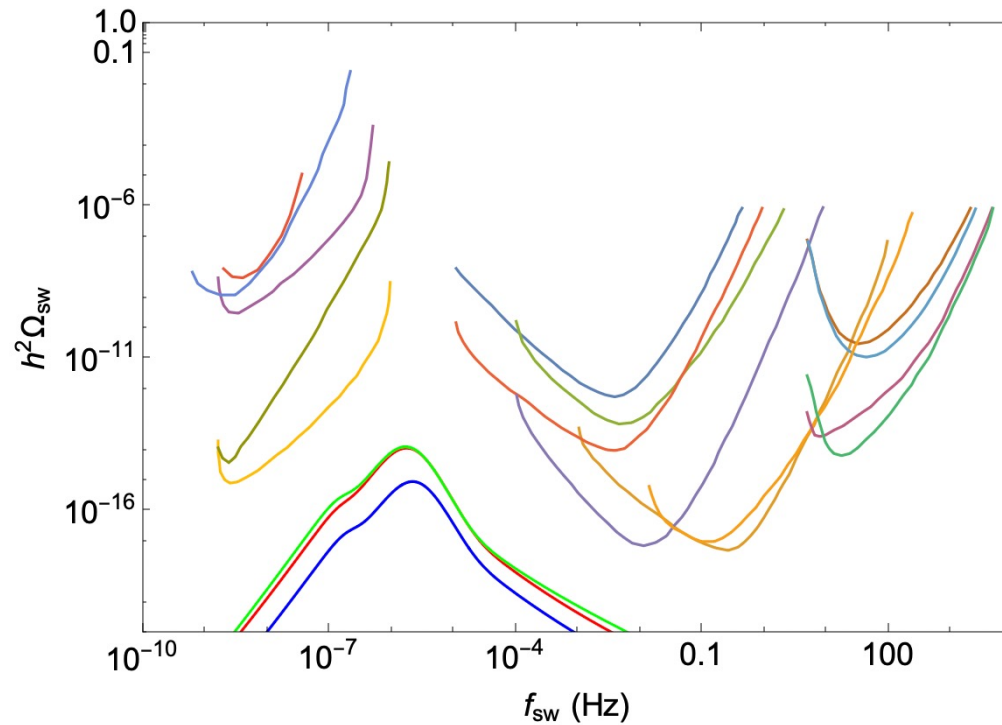
$$(\Delta_{\text{BW}}, \Delta_{\text{SW}}) = \left( \frac{0.11 v_w^3}{0.42 + v_w^3}, 1 \right)$$

- Spectral shape function and peak frequencies :

$$s_{\text{BW}}(x) = \frac{3.8 x^{2.8}}{1 + 2.8 x^{3.8}}, \quad s_{\text{SW}}(x) = x^3 \left( \frac{7}{4 + 3 x^2} \right)^{7/2},$$

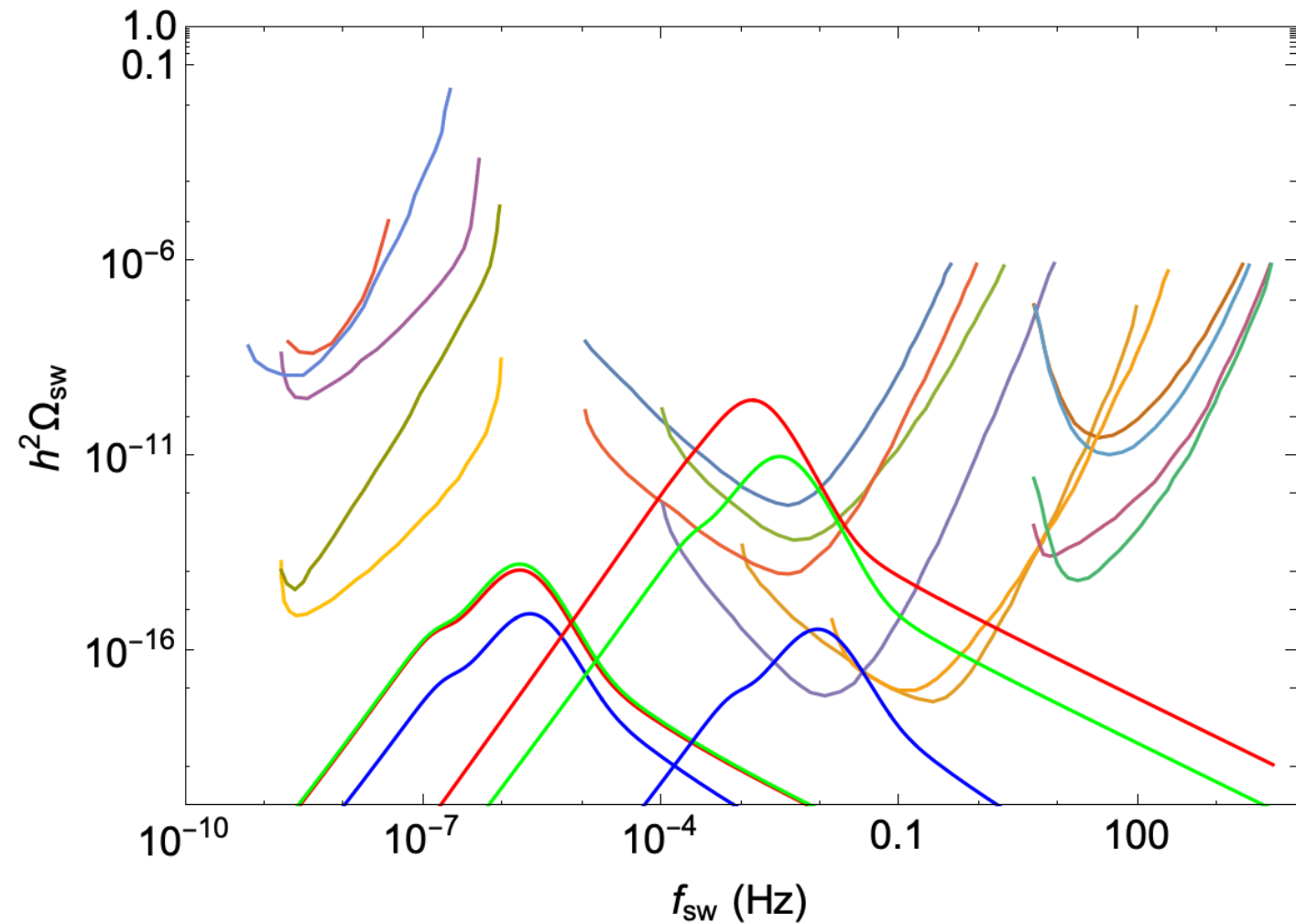
$$f_{\text{p,BW}} = 0.23 \beta, \quad f_{\text{p,SW}} = 0.53 \beta / v_w.$$

# GW spectrum $\mu = 10 \text{ MeV}$



$\{\alpha_{\text{kick}}, \beta_{\text{kick}}/H, T_{\text{SM,kick}}, \theta\} = 0.000521944, 91.9744, 0.108399, 0.397682$   
 $\{\alpha_{\text{late}}, \beta_{\text{late}}/H, T_{\text{SM,late}}\} = 0.000772954, 92.2871, 0.102398$   
 $\{\alpha, \beta/H, T_{\text{SM,first}}\} = 0.000203883, 92.2871, 0.14312$

# GW spectrum : comparing scales



# Slow reheating

