
The Inner Dark Matter Distribution in Hydrodynamic Simulations



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Milky Way

Credit: ESO



Synthetic survey of a baryonic simulation

(Sanderson et al. 2020)

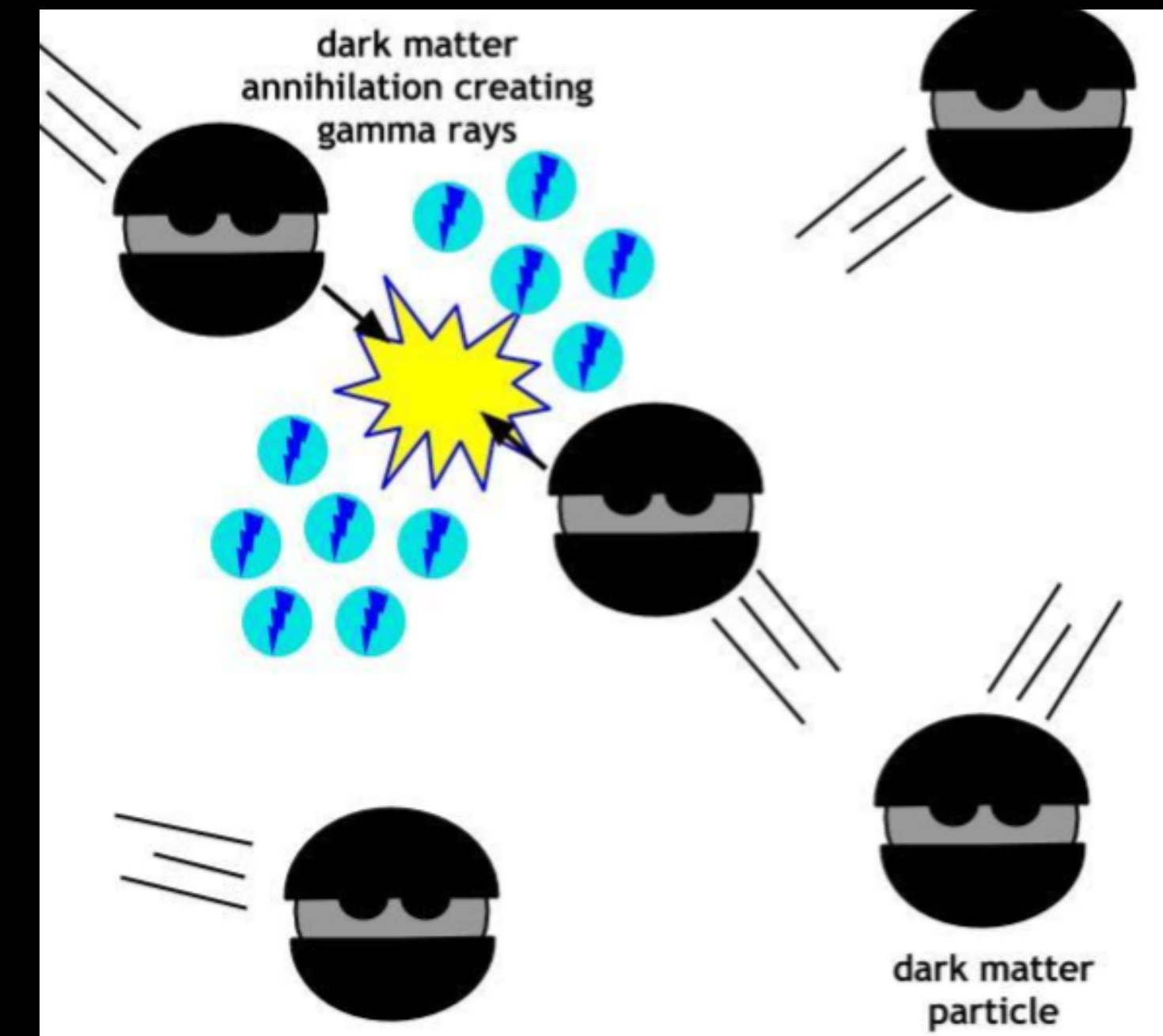


Outline:

- 🌀 **Why care about the density distribution?**
- 🌀 **Cosmological simulations overview**
- 🌀 **DM Distribution in different simulation suites**
- 🌀 **Adiabatic contraction overview**
- 🌀 **Calculation overview**
- 🌀 **Results**

Why care about the density distribution?

- For some DM models (ex: WIMPs) we get γ -ray emission from annihilation (Arcadi et al. 2018).
- The annihilation flux luminosity depends sensitively on ρ_{DM} .
 - $\mathcal{L} \propto \rho_{\text{DM}}^2$
- While traditionally a form for the DM density profile is assumed (NFW, Einasto,...), we can get a more informative result by using the density numerically calculated from the simulation.



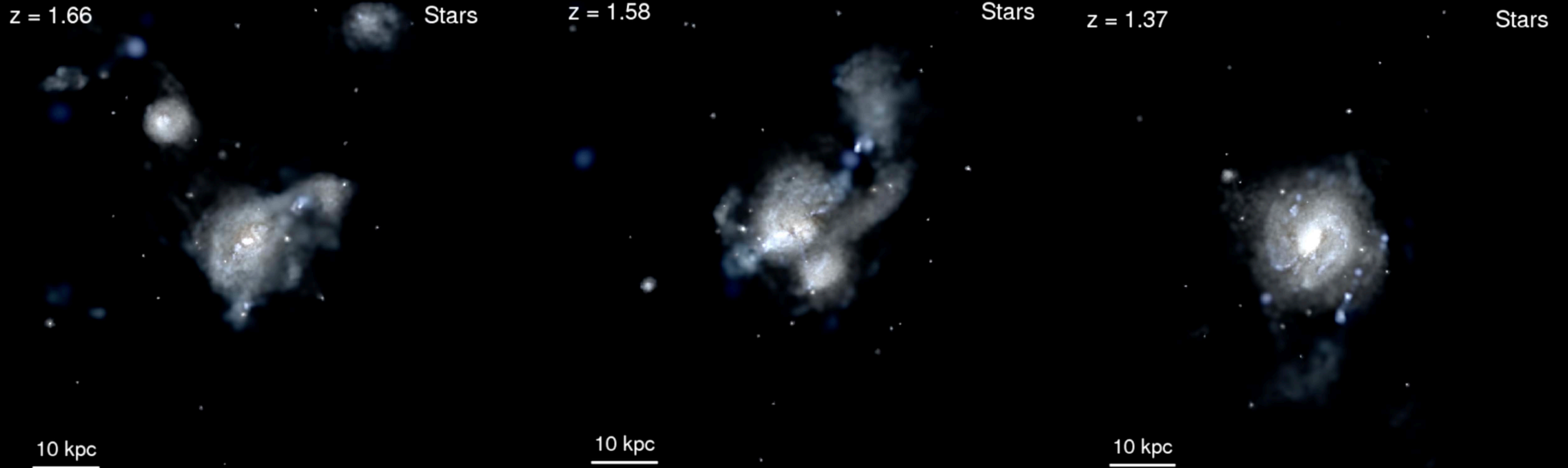
(Credit: Andrea Albert)

Cosmological simulations overview:



(credit: Phil Hopkins)

Cosmological simulations overview:



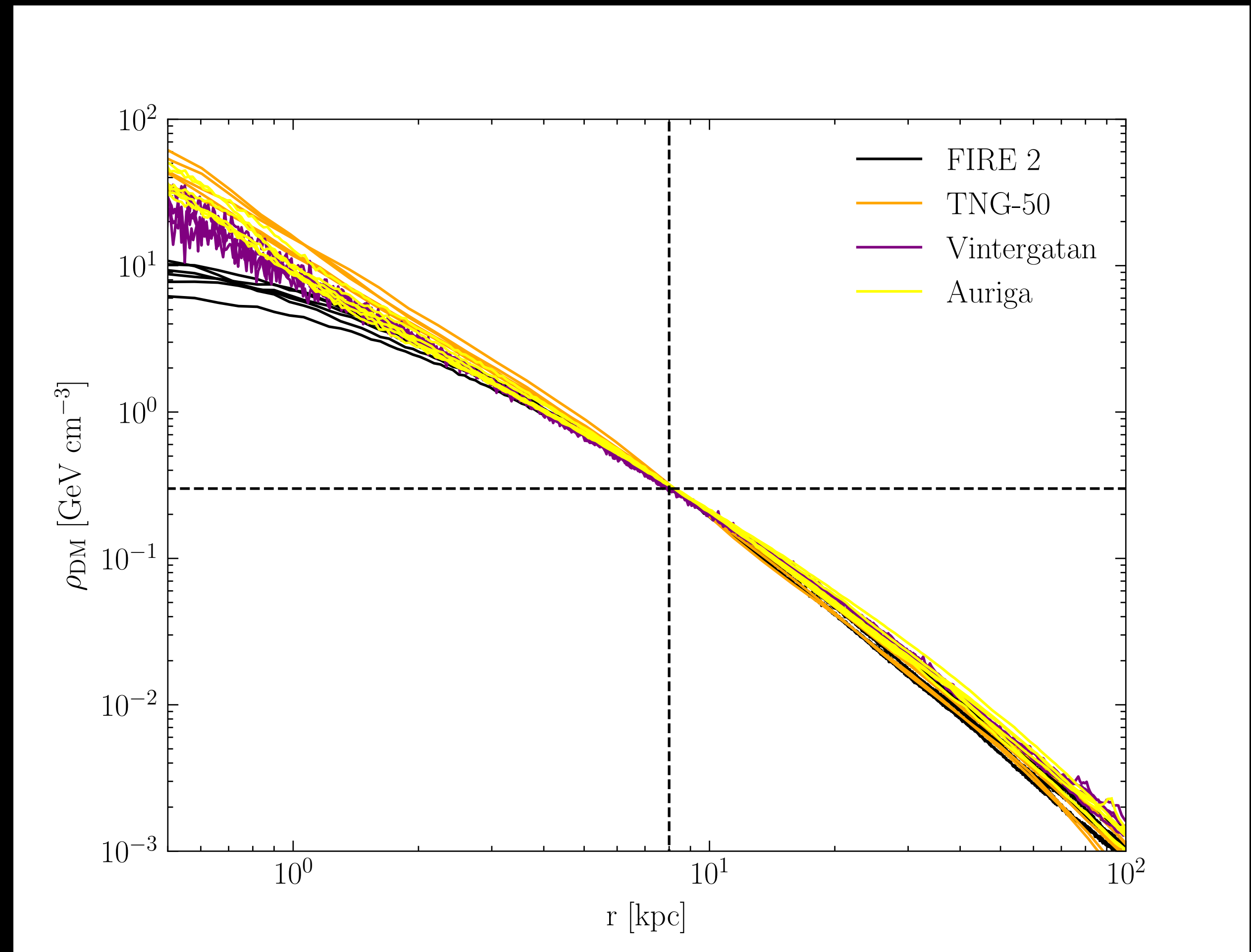
(credit: Phil Hopkins)

Cosmological simulations overview:

Simulations	Year	Technique	SMBH Feedback	$m_{\text{DM}}(M_{\odot})$	$m_{\text{baryon}}(M_{\odot})$
Auriga L3	2017	Zoom in	Yes	5E+04	6E+03
FIRE-2	2017	Zoom in	No	3.5E+04	7.1E+03
Vintergatan	2020	Zoom in	No	3.5E+04	7.07E+03
TNG-50	2019	Uniform Resolution	Yes	4.5E+05	8E+04

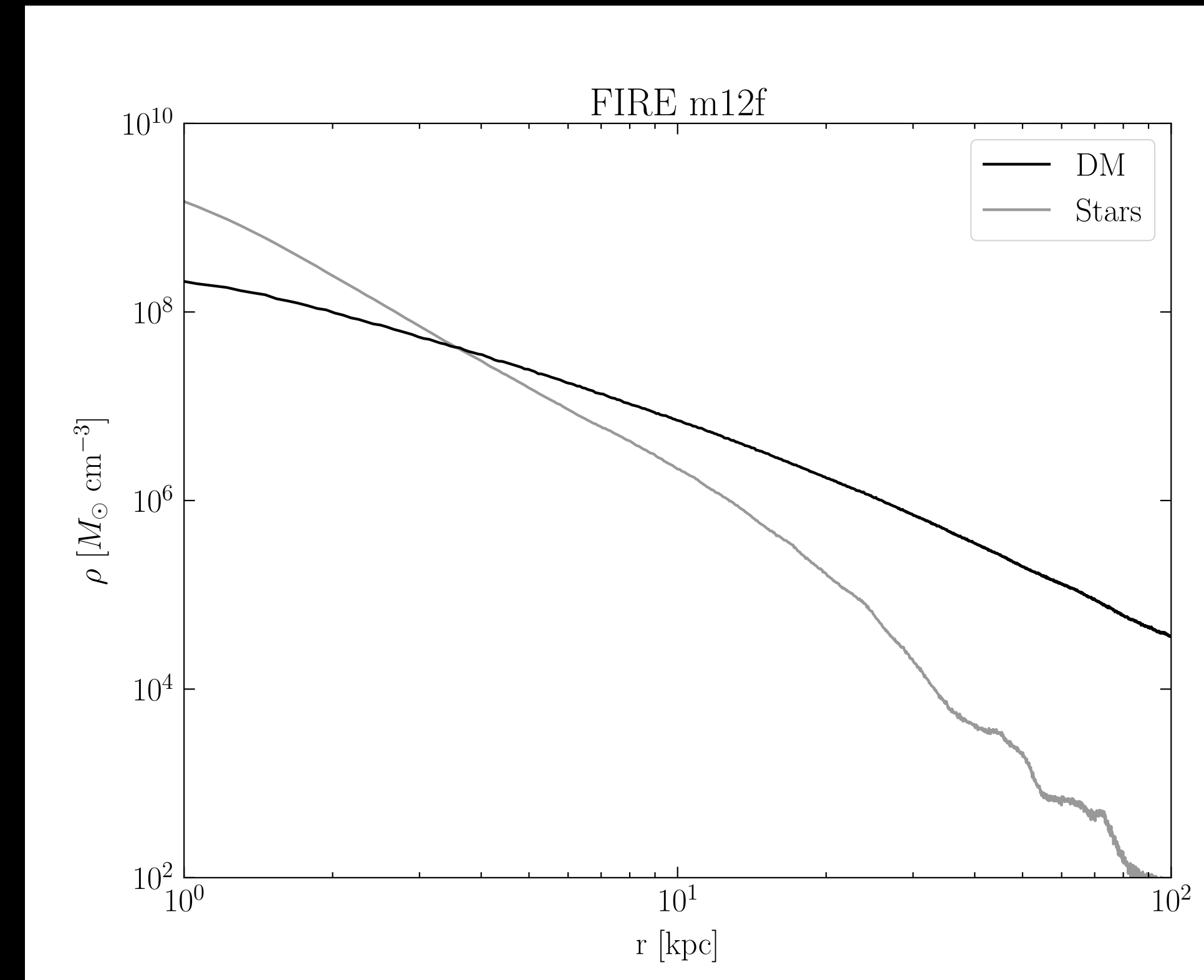
DM Distribution in different simulation suites:

- **Largest difference within 1 kpc**
- **Density similar for $r > 1$ kpc**
- **How can we quantify the difference?**
- **Can any of these be modeled through adiabatic contraction?**



Adiabatic contraction overview:

- **The gravitational field in the central regions of galaxies is dominated by stars.**
- **The conserved quantities for eccentric orbits (Ghigna et al. 1998)**
- **the radial action $I_r \equiv \frac{1}{\pi} \int_{r_p}^{r_a} v_r dr$**
- **(Gnedin et al. 2004) argued that the conserved quantity $r M(\bar{r})$ is a better proxy for the radial action.**



Adiabatic contraction input:

$$M_{\text{DM}}^{\text{initial}}(r_{\text{initial}}) = M_{\text{DM}}^{\text{final}}(r_{\text{final}})$$

$$r_{\text{initial}}(M_{\text{DM}}^{\text{initial}}(\bar{r}_{\text{initial}}) + M_{\text{Stars}}^{\text{initial}}(\bar{r}_{\text{initial}})) = r_{\text{final}}(M_{\text{DM}}^{\text{final}}(\bar{r}_{\text{final}}) + M_{\text{Stars}}^{\text{final}}(\bar{r}_{\text{final}}))$$

➤ Inputs (all $z=0$):

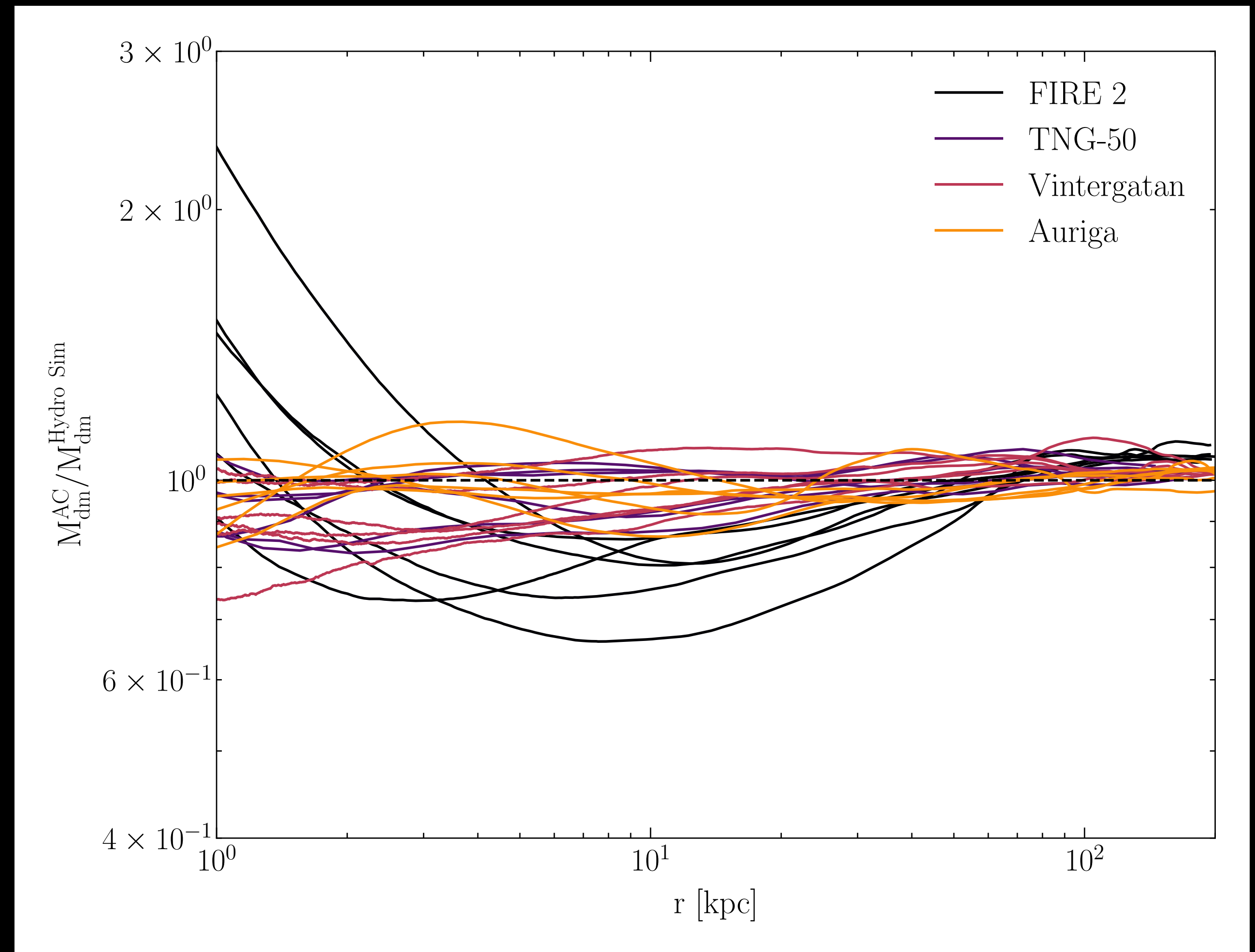
➤ **DM distribution from DMO sim**

➤ **A stellar distribution that is self similar to the DMO distribution**

➤ **Stellar density profile from hydro sim**

Results:

- **The ratio deviates within 10 kpc from 1 for FIRE sims relative to TNG50, Vintergatan and Auriga**
- **Vintergatan, TNG50 and Auriga DM density profiles can be described using adiabatic contraction.**



Conclusion:

- **Two possible solution:**
 - **Adiabatic contraction**
 - **Strong Feedback**
- **We will use AC to model the DM density profile of the MW**
 - **Obtain photon emission from DM annihilation signal.**

Back up

Calculation overview:

$$f_b = \frac{M_{Stars}^{hydro}(r_{200c})}{M_{DM}^{hydro}(r_{200c})} \quad (5)$$

$$f_{norm} = \frac{M_{DM}^{hydro}(r_{200c}) + M_{Stars}^{hydro}(r_{200c})}{M_{DM}^{DMO}(r_{200c})} \quad (6)$$

$$M_{DM}^{initial}(r) = (M_{DM}^{DMO}(r) \cdot f_{norm}) \cdot (1 - f_b) \quad (7)$$

$$M_{Stars}^{initial}(r) = (M_{DM}^{DMO}(r) \cdot f_{norm}) \cdot f_b \quad (8)$$

$$r_{initial}(M_{DM}^{initial}(\bar{r}_{initial}) + M_{Stars}^{initial}(\bar{r}_{initial})) = r_{final}(M_{DM}^{final}(\bar{r}_{final}) + M_{Stars}^{final}(\bar{r}_{final})) \quad (9)$$

$$M_{DM}^{initial}(r_{initial}) = M_{DM}^{final}(r_{final})$$

Find fixed point

$$\bar{r} = r_{vir} A \left(\frac{r}{r_{vir}} \right)^w \quad \text{tested by (Gustafsson et al 2007)}$$

Given such a wide eccentricity distribution, the orbit-averaged radius varies for particles at a given current radius r depending on the orbital phase. Nevertheless, the mean relation can be described by a power law function.

Adiabatic contraction overview:

➤ Assumptions:

➤ Eccentric orbits (Ghigna et al. 1998)

➤ Orbits have a wide distribution of eccentricities which should be taken to account.

➤ Take this distribution into account by averaging over the population of orbits at a given radius:

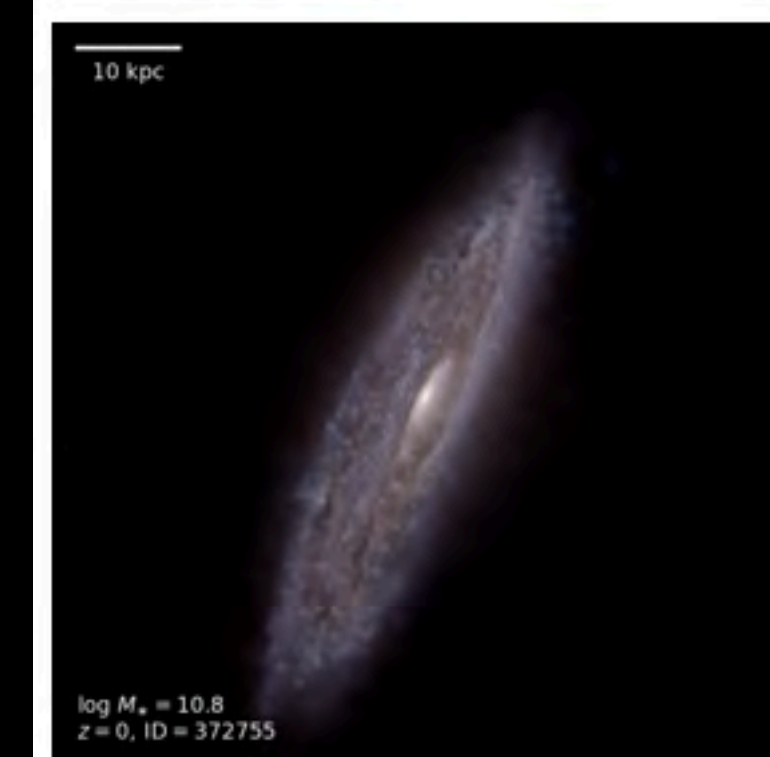
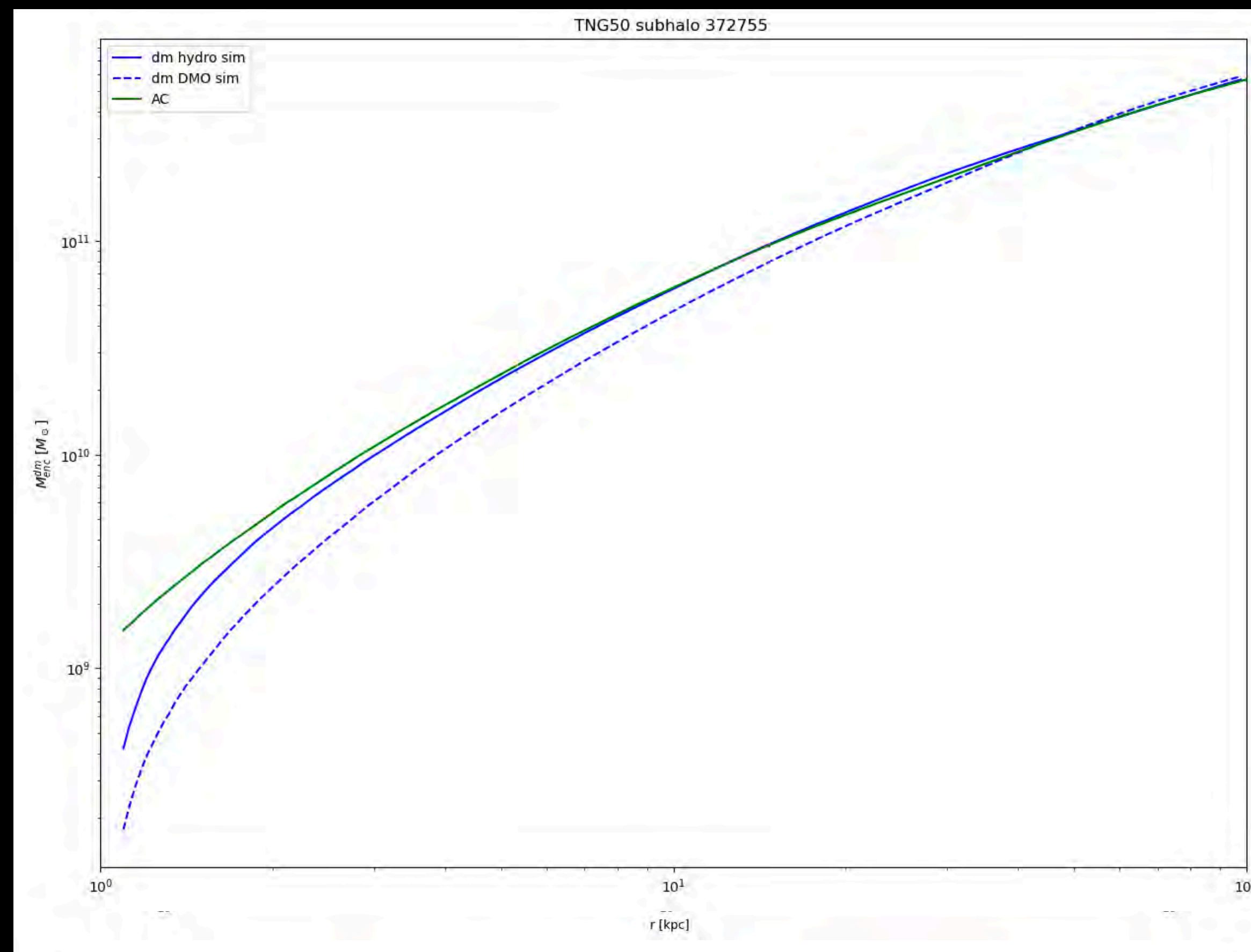
$$\bar{r} = r_{vir} A \left(\frac{r}{r_{vir}} \right)^w$$

➤ Spherical symmetry

➤ Conservation of angular momentum

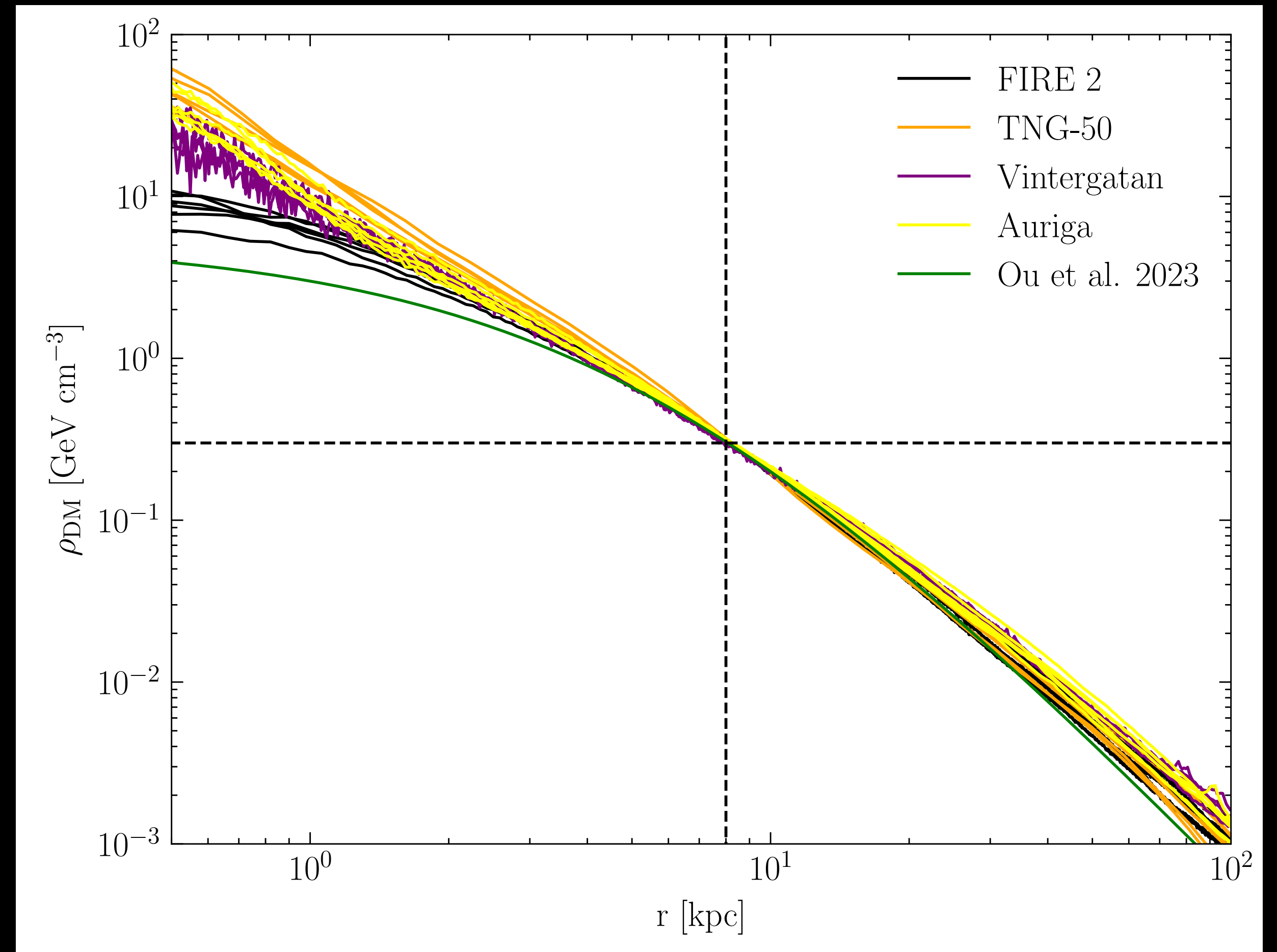
➤ Assume homologous contraction

Adiabatic Contraction in TNG50:

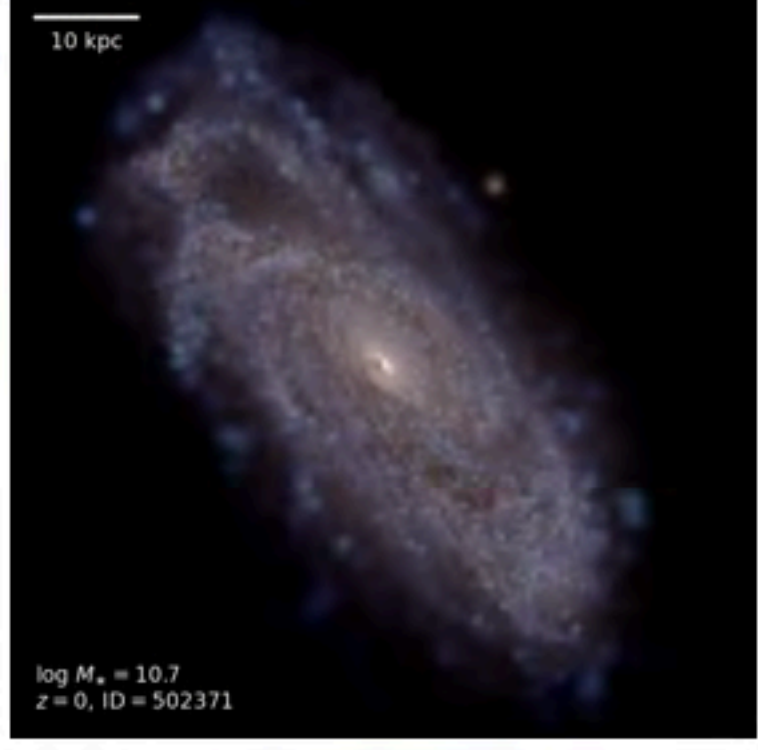
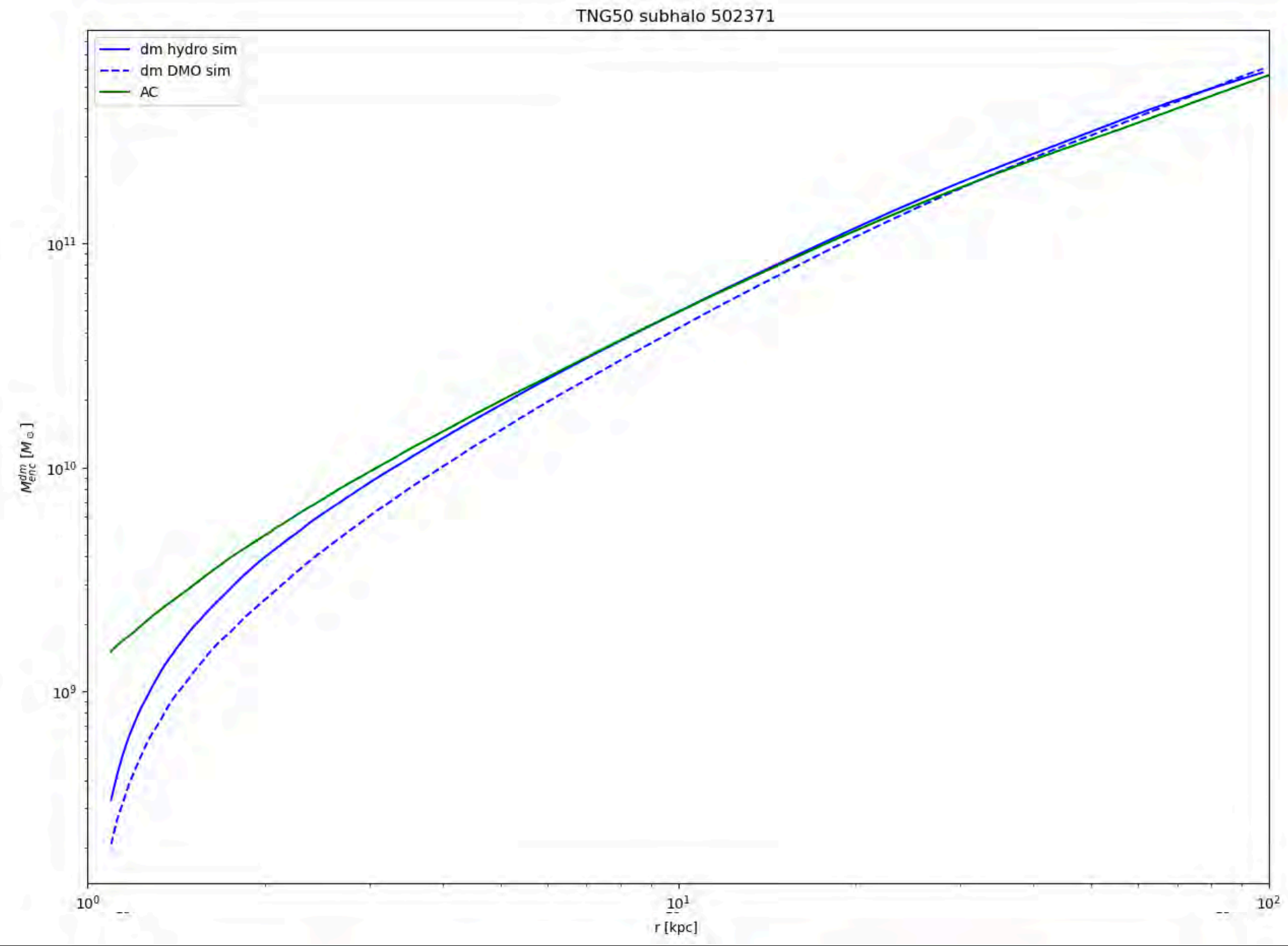


DM Distribution in different simulation suites:

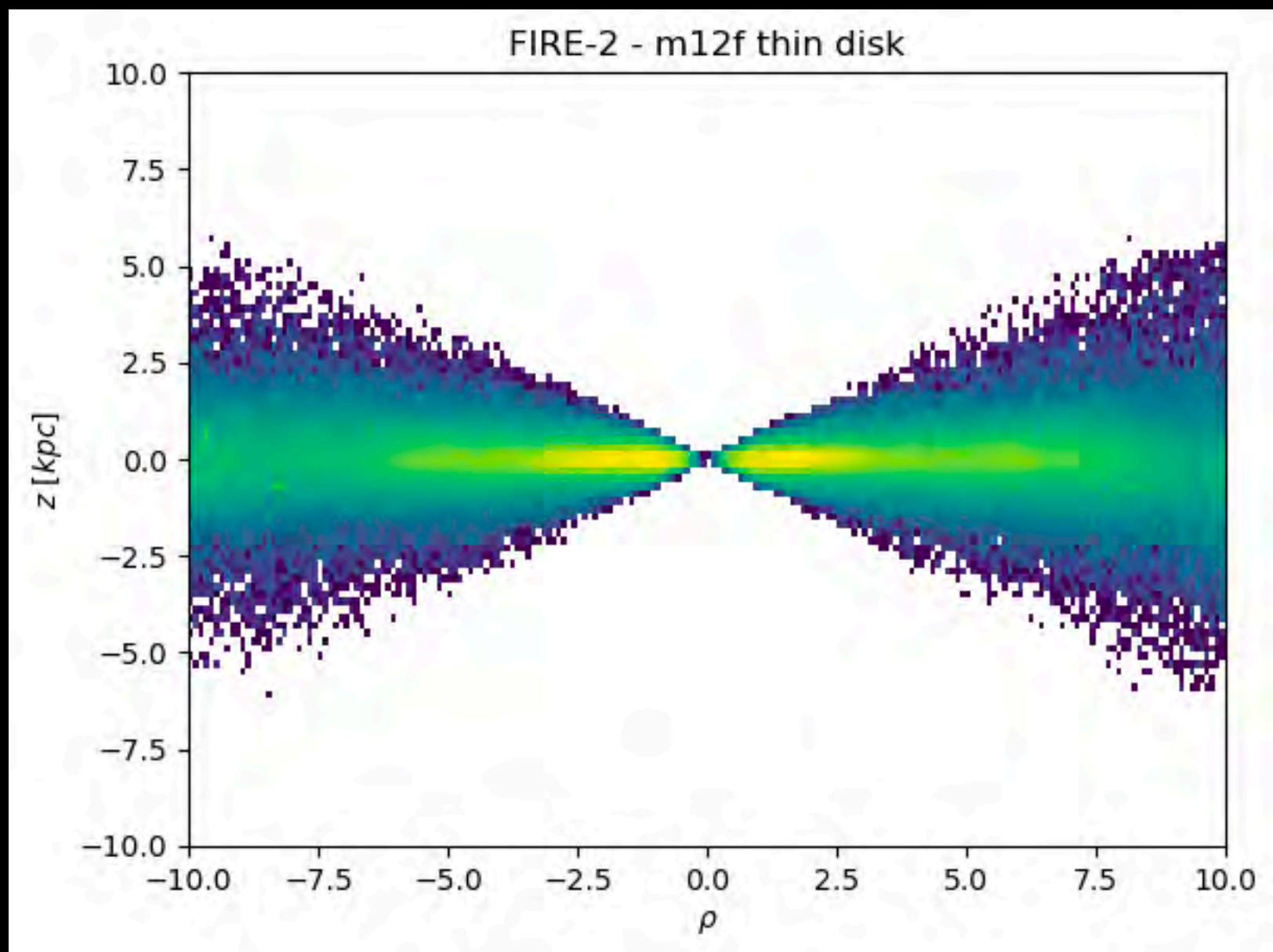
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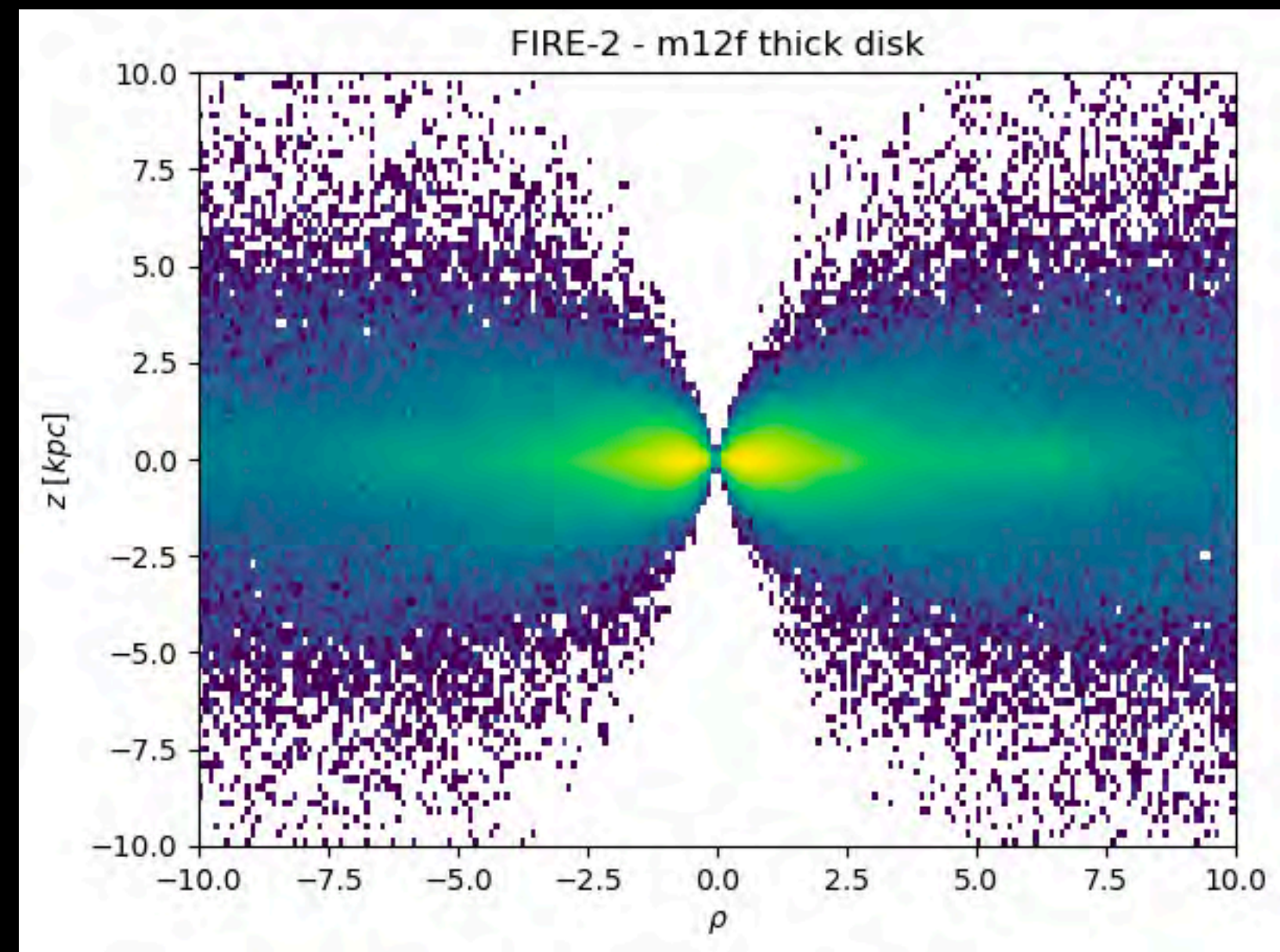
Adiabatic Contraction in TNG50:



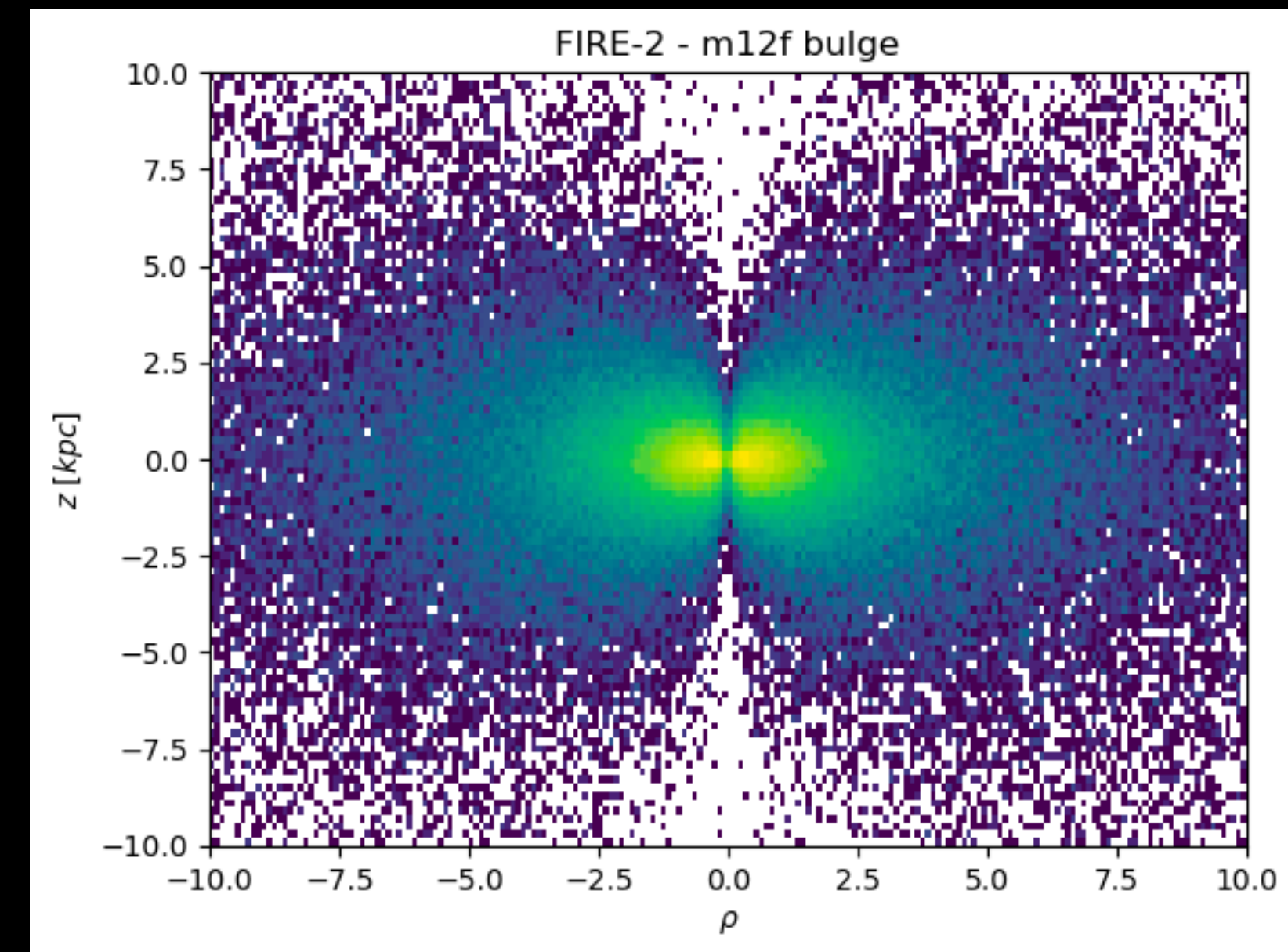
Thin disk



Thick disk

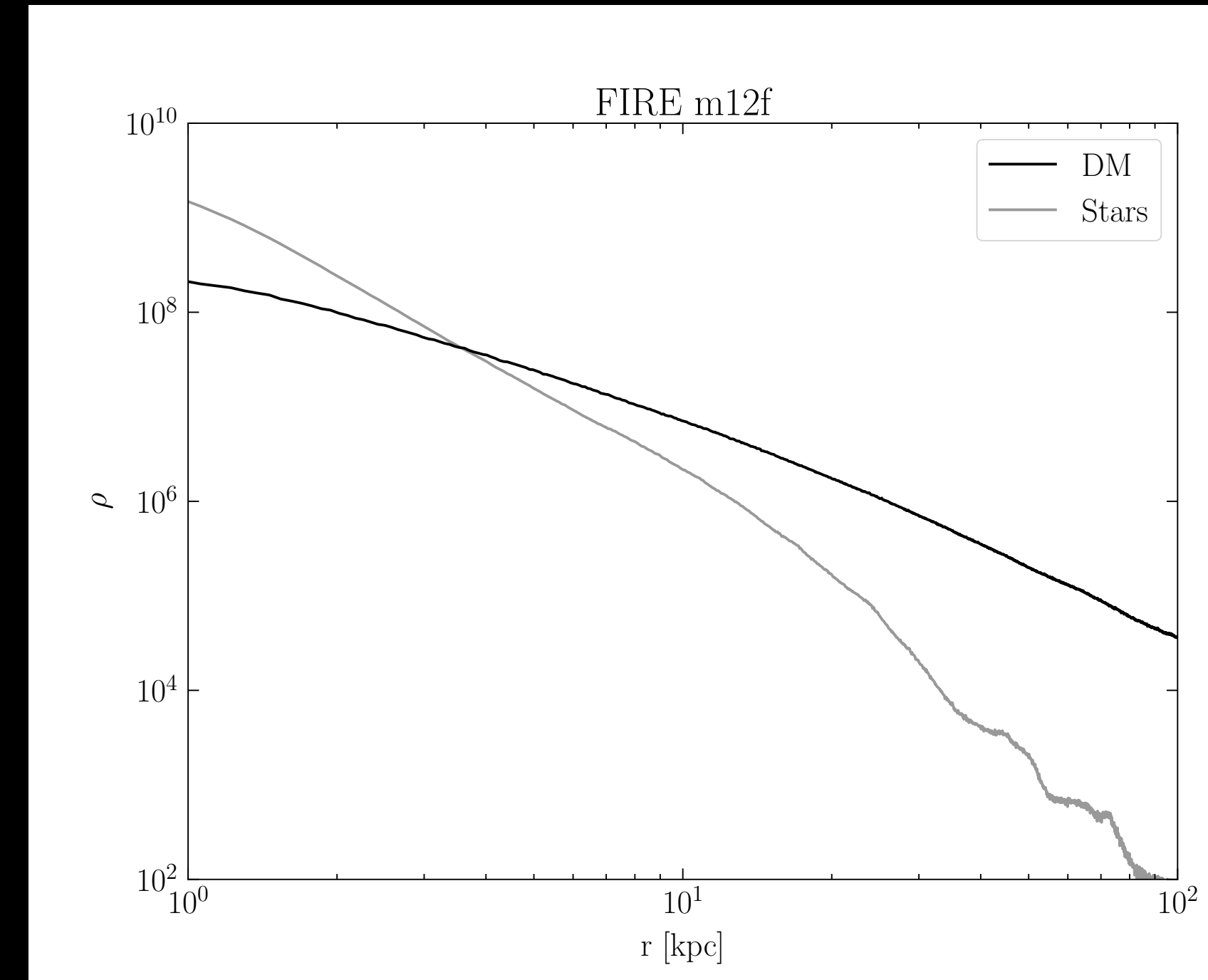


Bulge



Adiabatic contraction overview:

- **Although dark matter exceeds baryonic matter by a factor of $\Omega_b \simeq 5\Omega_{DM}$**
 - **The gravitational field in the central regions of galaxies is dominated by stars.**
 - **As the baryons condense in the center, they pull the dark matter particles inward thereby increasing their density in the central region.**
- **The conserved quantities for eccentric orbits**
 - **the angular momentum J**
 - **the radial action $I_r \equiv \frac{1}{\pi} \int_{r_p}^{r_a} v_r dr$**
 - **(Gnedin et al. 2004) argued that the conserved quantity $r M(\bar{r})$ is a better proxy for the radial action.**

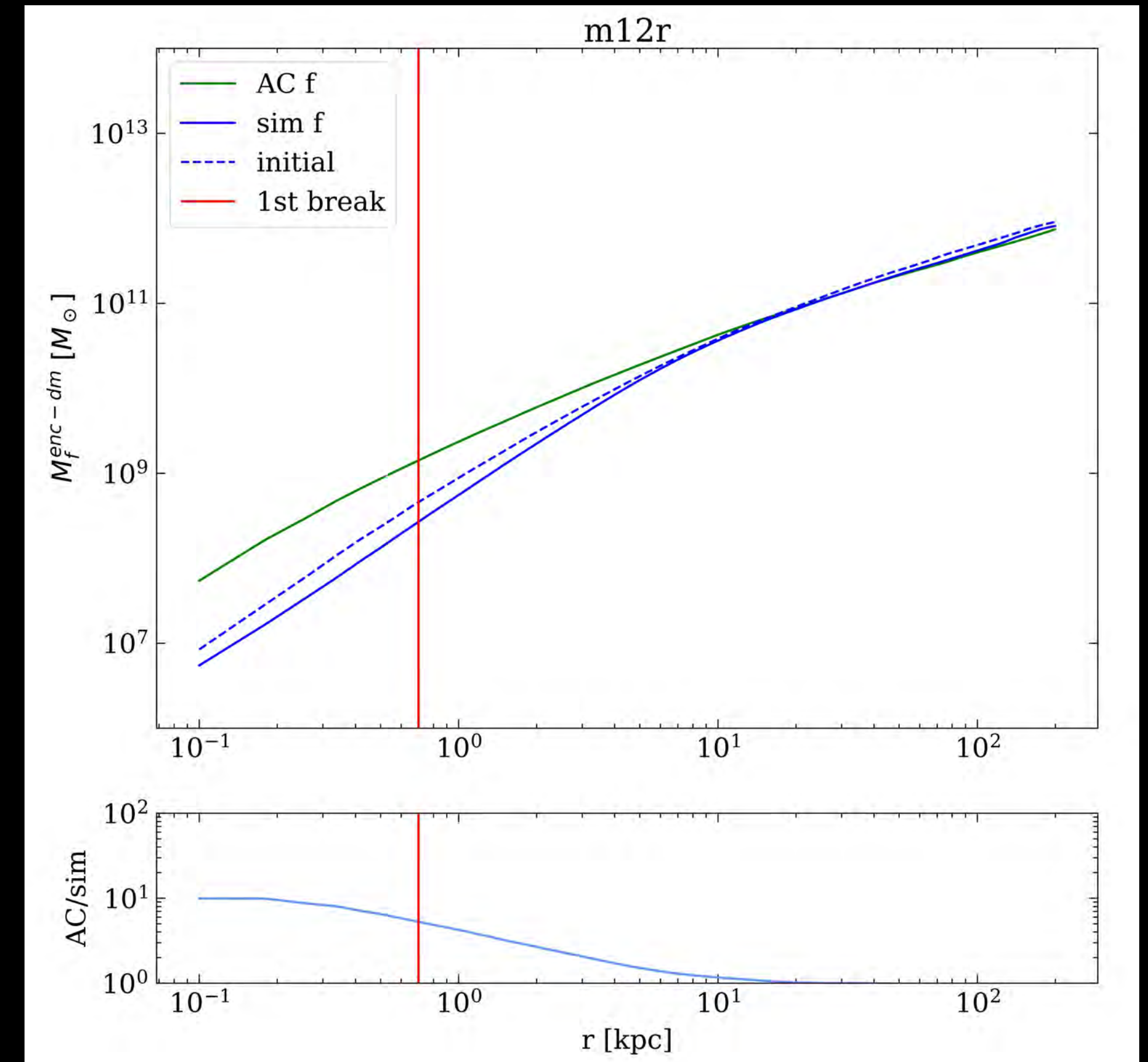
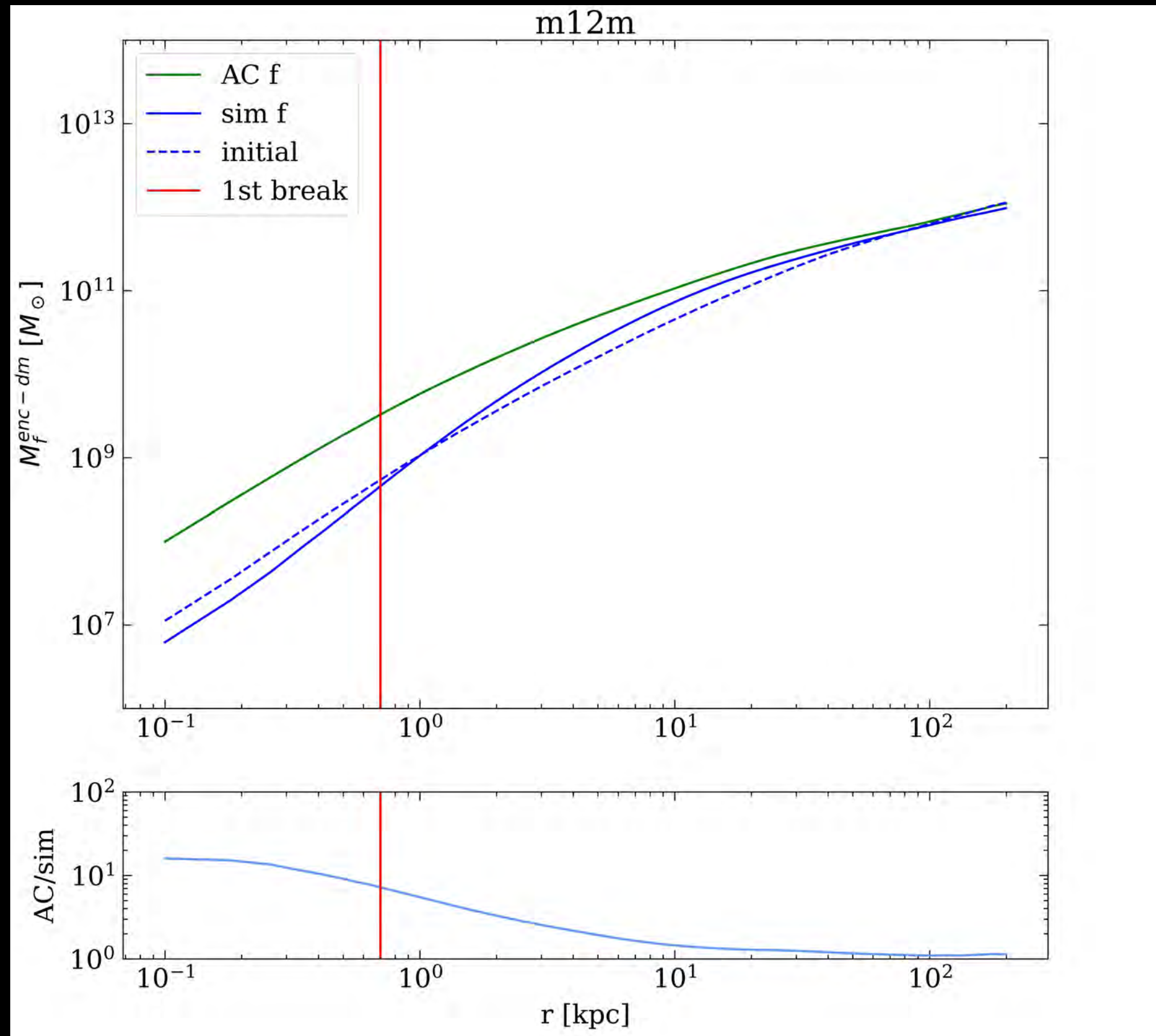


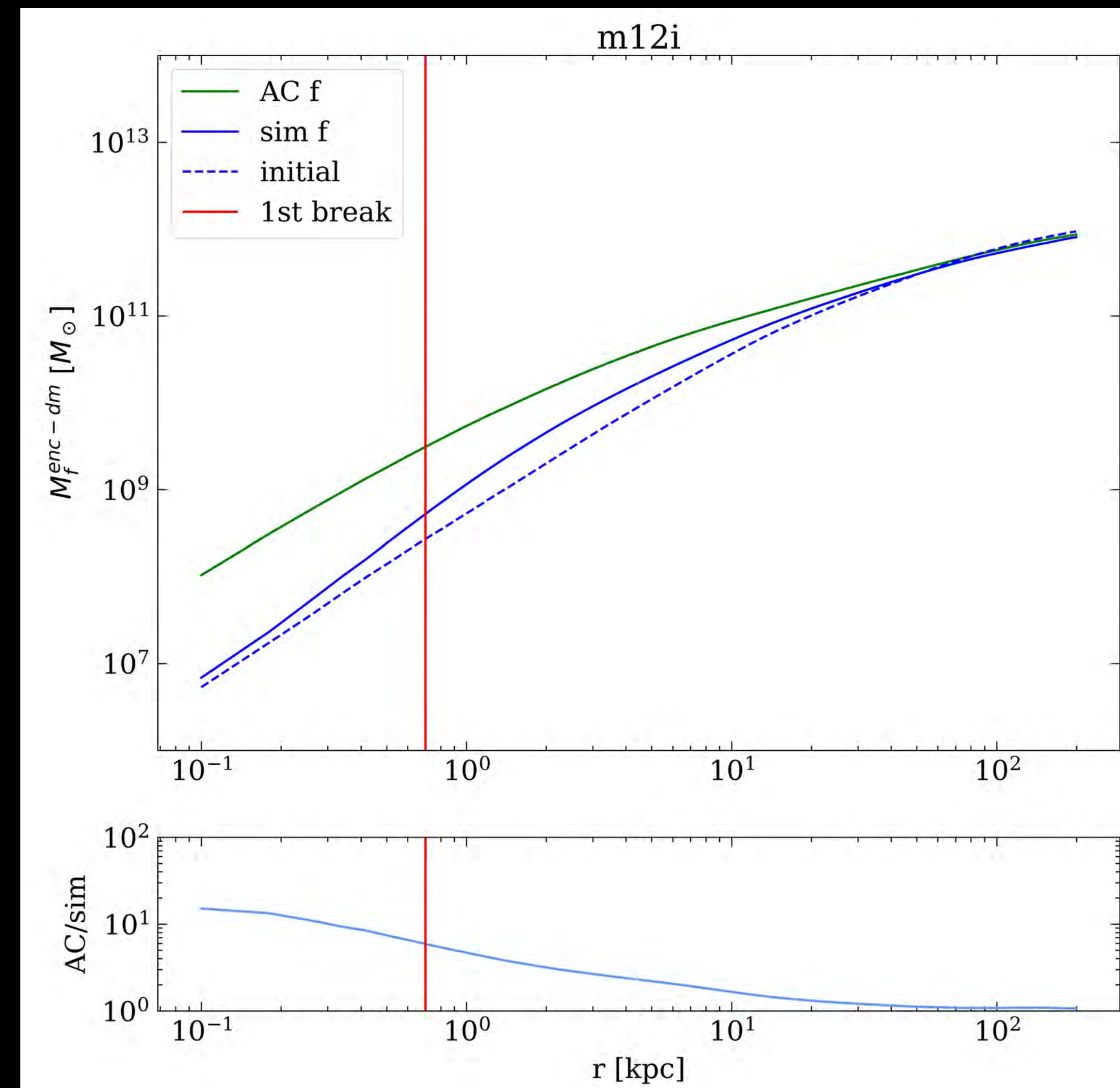
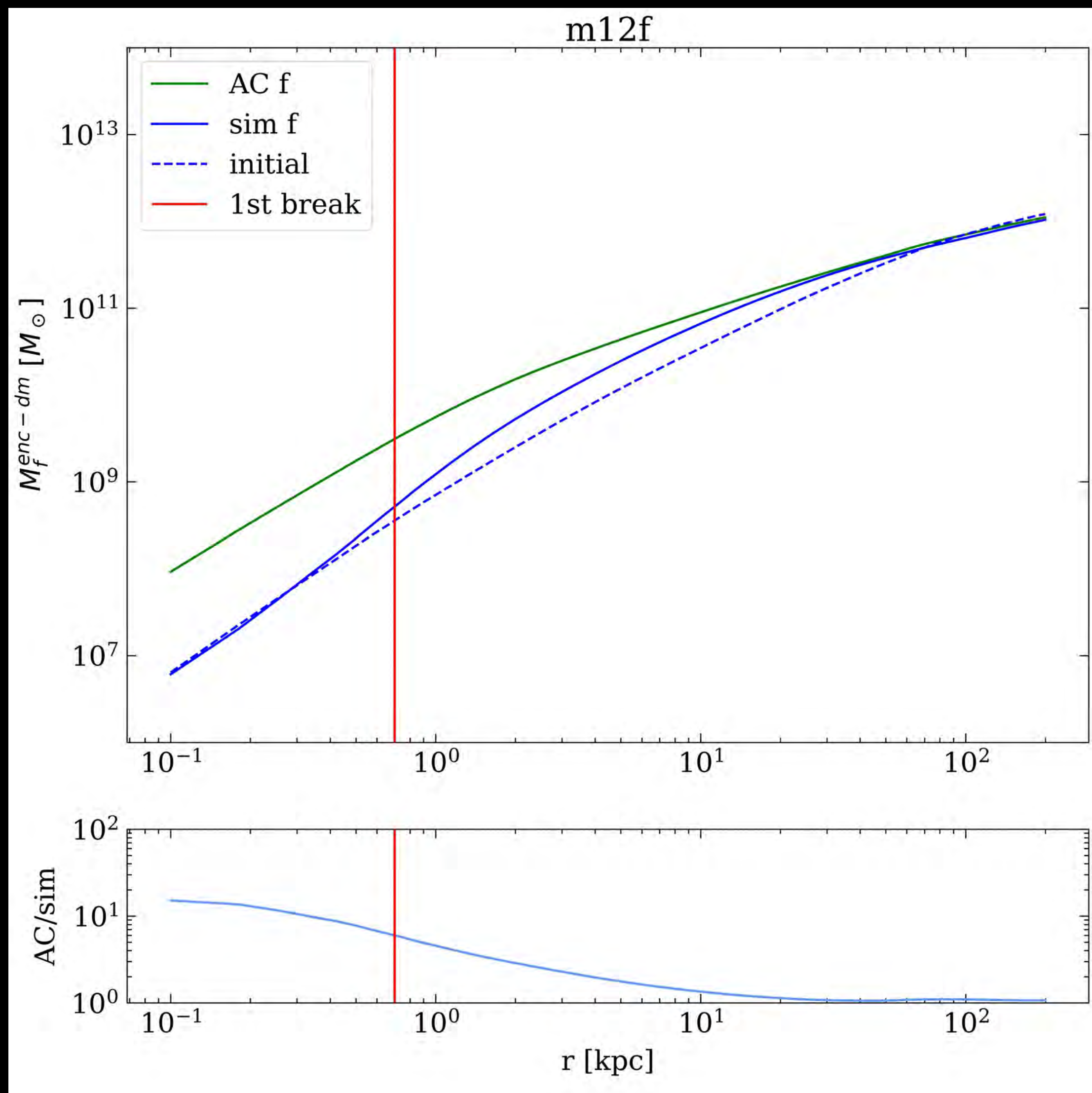
Adiabatic Contraction in FIRE m12s

(Gnedin et al. 2004)

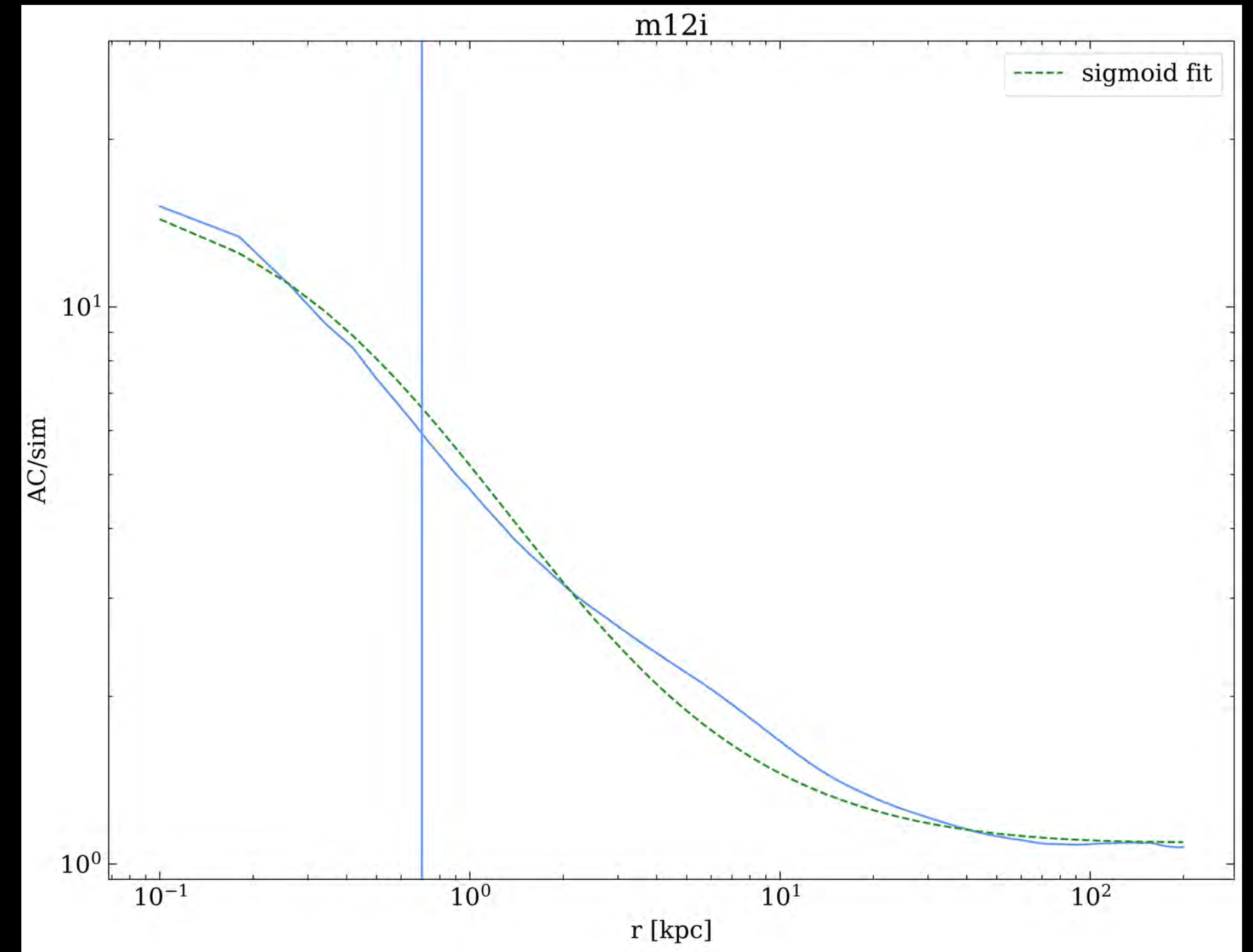
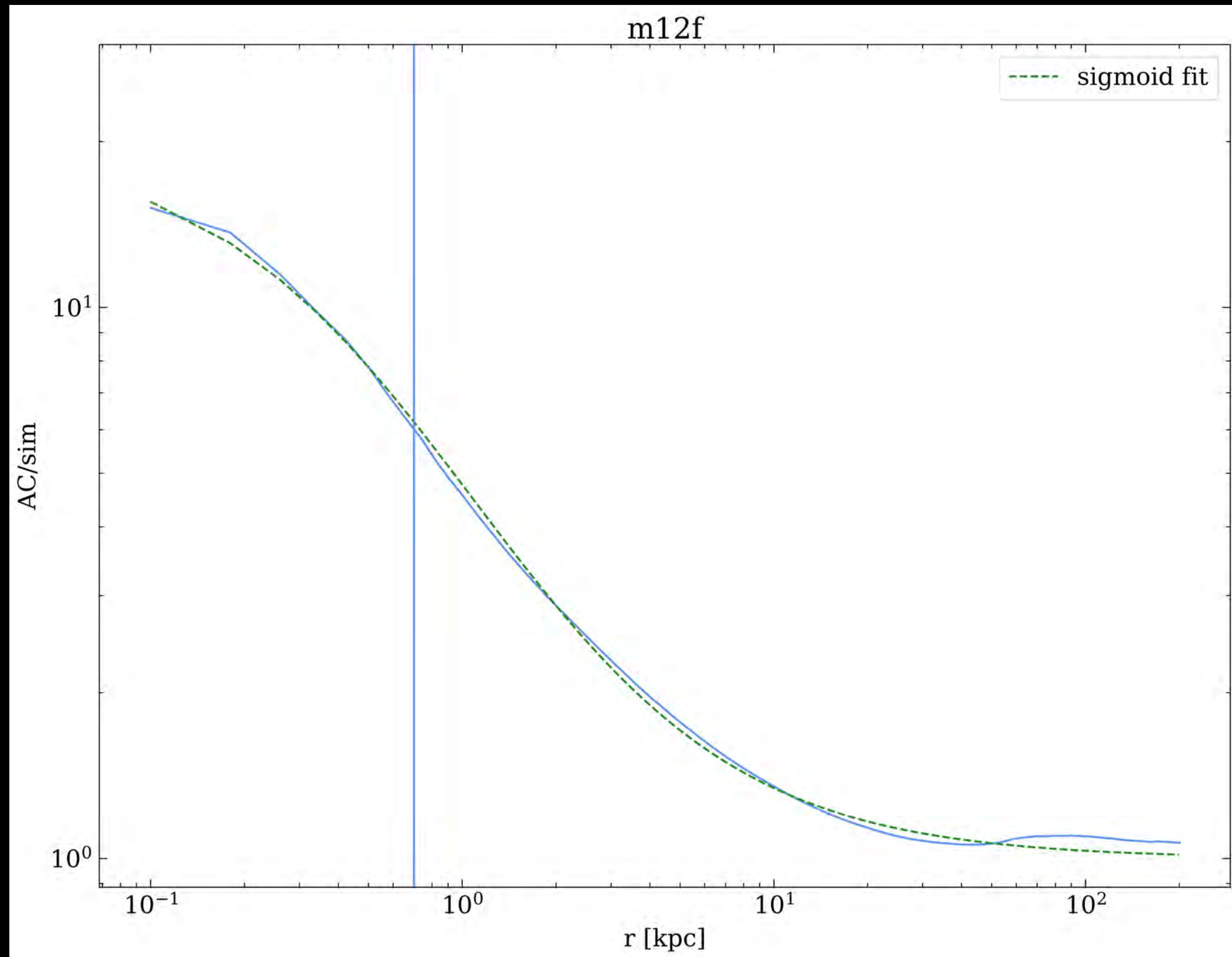
$$r M(\bar{r}) = \text{const}$$

$$\bar{r} = r_{\text{vir}} A \left(\frac{r}{r_{\text{vir}}} \right)^w \quad A = 0.85, w = 0.8$$





Looking at transformation



Adiabatic Contraction in Vintergatan Halo 685

