

A  
Presentation  
on

# Effects of Quantum Gravity in the Kerr Black Hole Paradigm

17<sup>TH</sup> INTERNATIONAL CONFERENCE ON INTERCONNECTIONS BETWEEN PARTICLE PHYSICS AND COSMOLOGY

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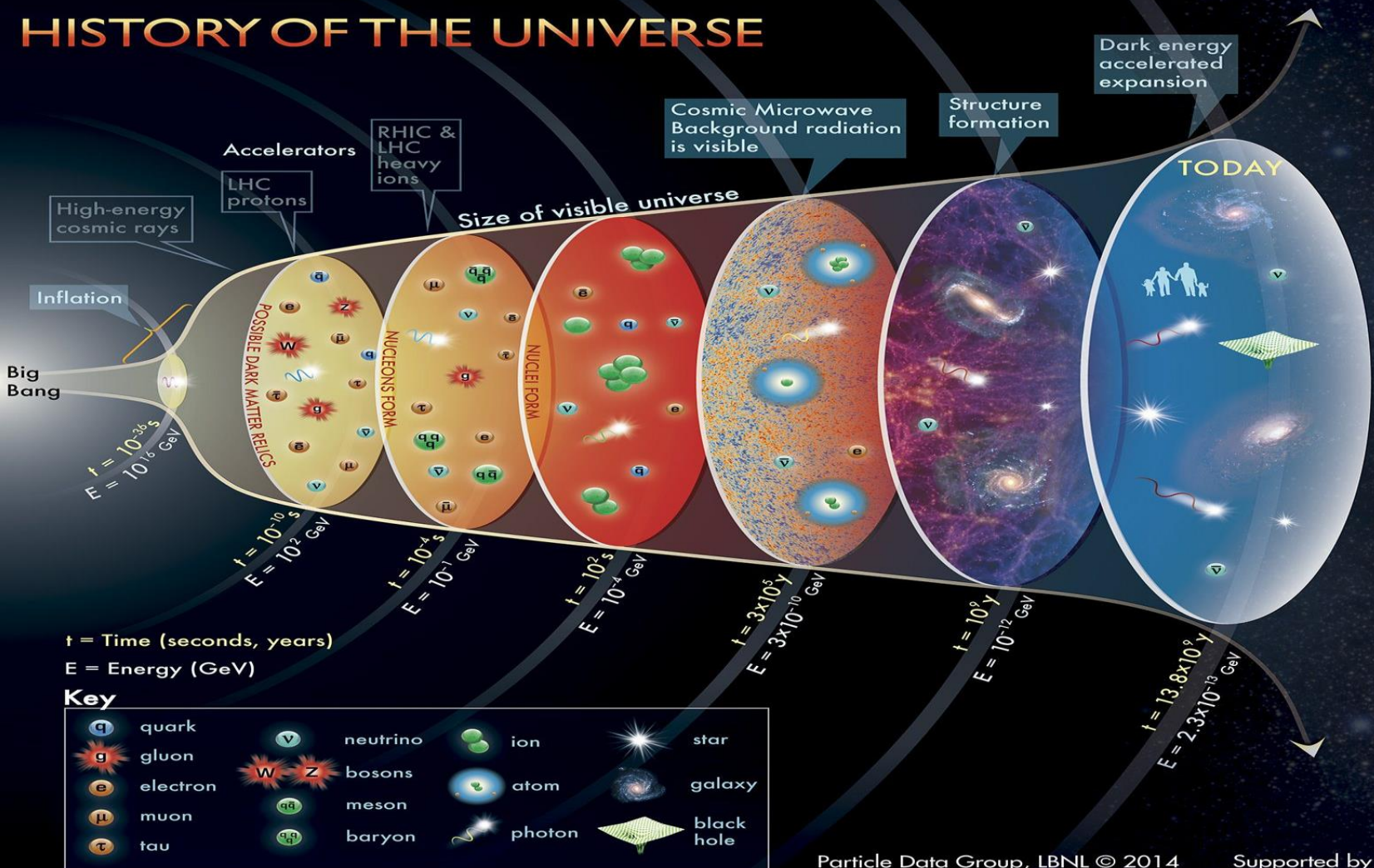
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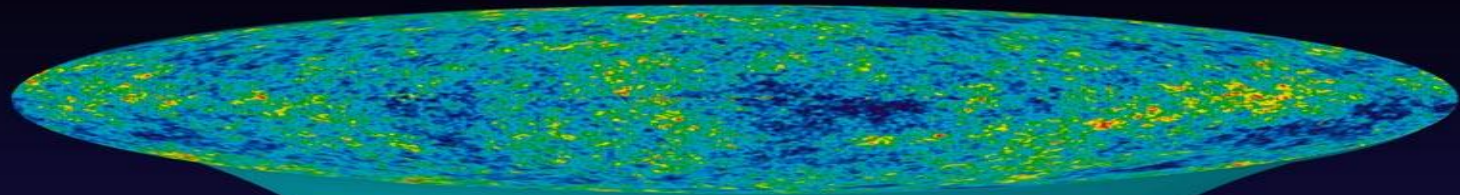
# PLAN OF TALK

- ❖ OBJECTIVE
- ❖ INTRODUCTION
- ❖ SETTING THE STAGE
- ❖ CALCULATION
- ❖ CONCLUSION
- ❖ REFFERENCES

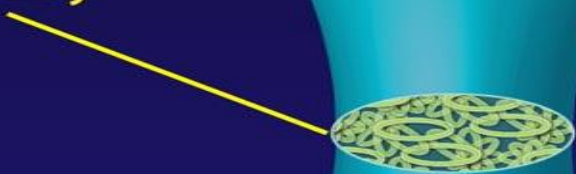
# HISTORY OF THE UNIVERSE



# Cosmic microwave background



Quantum  
Geometry



Bounce

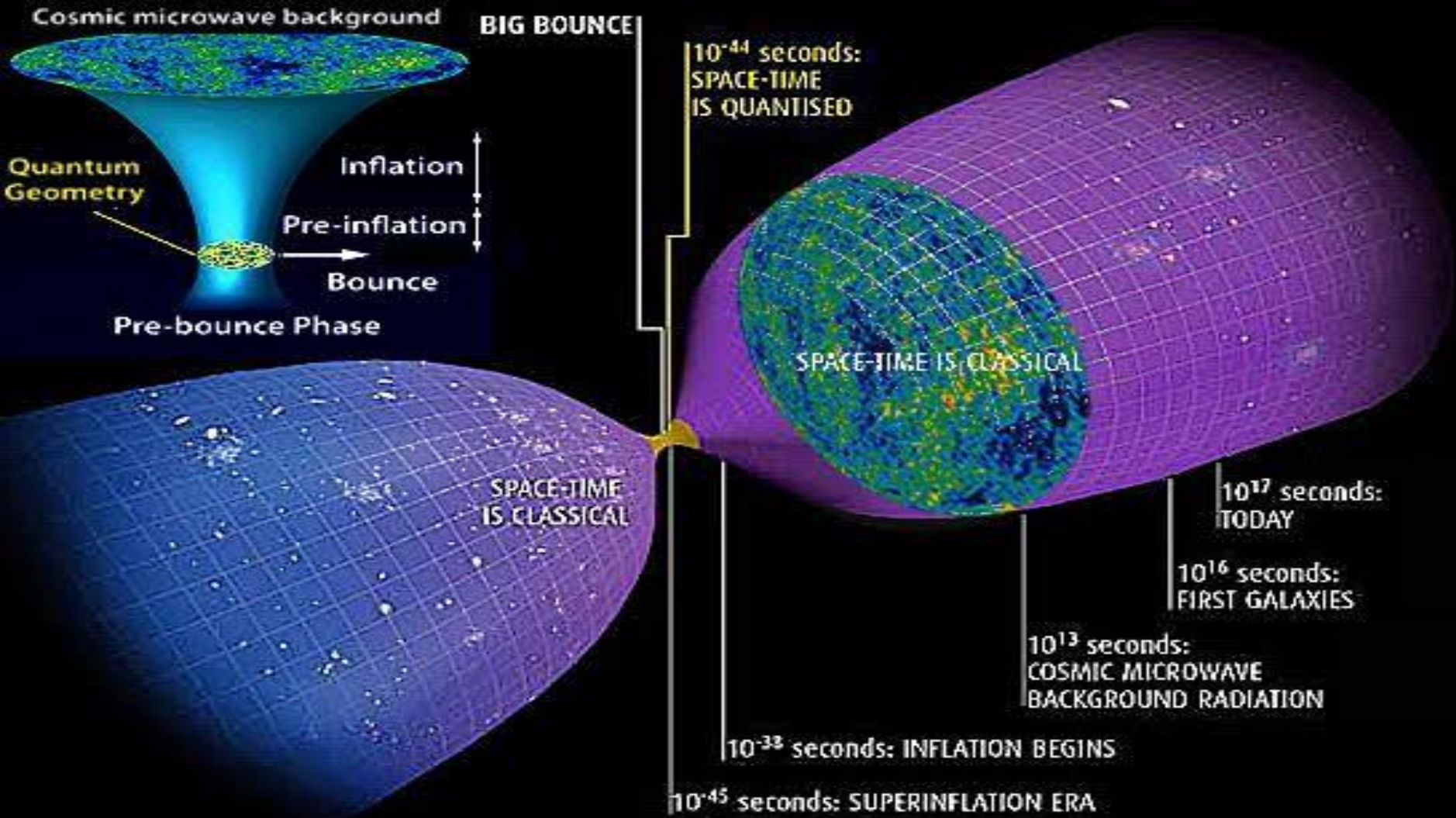
Inflation

Pre-inflation

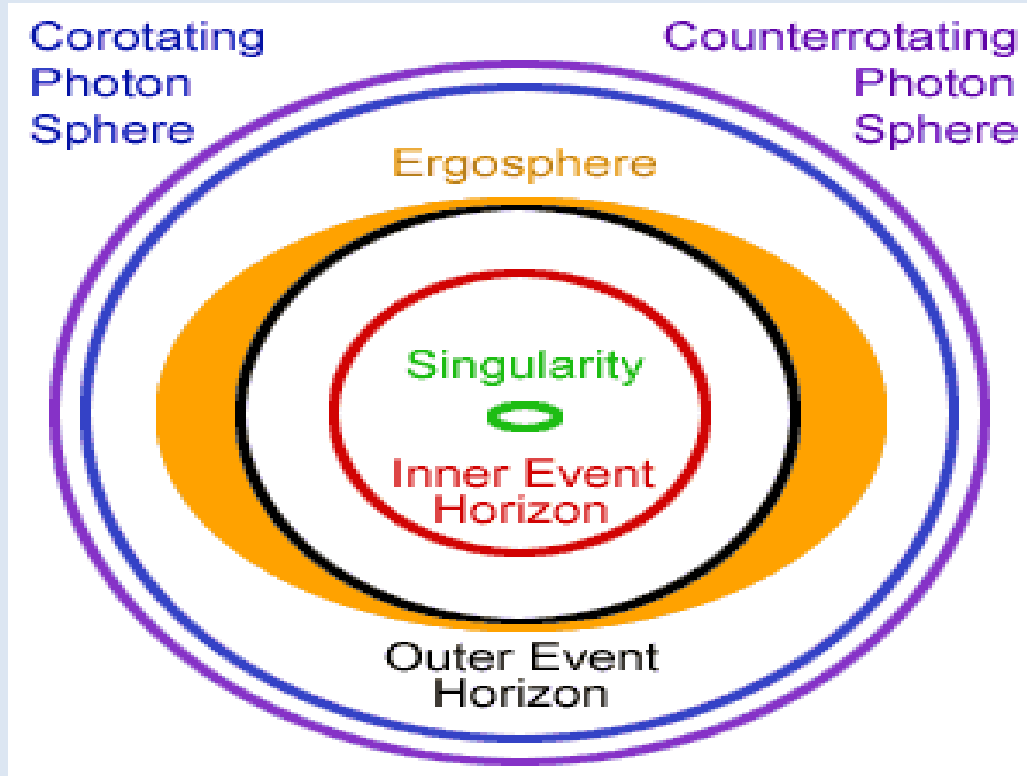


Pre-bounce Phase  
of the Universe





# Schematic Diagram For Rotating Black hole



# OBJECTIVE

- To study the accretion and evaporation of the rotating black holes using Loop Quantum Cosmology (LQC).
- To examine the nature of supermassive rotating black hole in the background of LQC.

# Introduction

- ❖ Loop Quantum Gravity (LQG) is one of the special features of quantum gravity theories which is completely non-perturbative, explicit background independent approach to quantum gravity.
- ❖ Generally the implication of LQG on cosmology for the study of our universe is called as loop quantum cosmology (LQC).
- ❖ Kerr explains successfully regarding the geometry of uncharged black holes. As we know from the no hair theorem, a spinning uncharged black hole is only described by its two quantities. One is its mass ( $M$ ) and another is respective angular momentum ( $J$ ).
- ❖ Depending upon these various quantities such as mass, angular momentum, accretion efficiency ( $f$ ), energy density ( $\rho$ ) and dark energy parameter ( $\omega$ ), we can calculate the evolution of the rotating Kerr black hole.



# SETTING THE STAGE

❖ For a spatially flat FRW universe ( $k=0$ ) filled with dust and dark energy, the Friedmann equation, Raychaudhuri's equation and energy conservation equation in loop quantum cosmology respectively take the form

$$H^2 = \frac{8\pi G}{3} (\rho_\phi + \rho_m) \left\{ 1 - \frac{(\rho_\phi + \rho_m)}{\rho_c} \right\} \quad (1)$$

$$\dot{H} = -4\pi G (\rho_\phi + \rho_m + p_\phi) \left\{ 1 - \frac{2(\rho_\phi + \rho_m)}{\rho_c} \right\} \quad (2)$$

And

$$(\dot{\rho}_\phi + \dot{\rho}_m) + 3H(\rho_\phi + \rho_m + p_\phi) = 0 \quad (3)$$

❖ The expression for density  $\rho(t)$  in radiation and matter dominated era can be calculated as

$$\rho(t)_{t < t_e} = \rho_0 \left[ \frac{\rho_0}{\rho_c} + \left\{ 2\sqrt{\frac{8\pi G}{3}} \rho_0^{\frac{1}{2}}(t - t_e) + \frac{3}{2}\sqrt{\frac{8\pi G}{3}} \rho_0^{\frac{1}{2}}(t_e - t_o) + \left(1 - \frac{\rho_0}{\rho_c}\right)^{\frac{1}{2}} \right\}^2 \right]^{-1} \quad (4)$$

$$\rho(t)_{t > t_e} = \rho_0 \left[ \frac{\rho_0}{\rho_c} + \left\{ \frac{3}{2}\sqrt{\frac{8\pi G}{3}} \rho_0^{\frac{1}{2}}(t - t_o) + \left(1 - \frac{\rho_0}{\rho_c}\right)^{\frac{1}{2}} \right\}^2 \right]^{-1} \quad (5)$$

- ❖ The space time around a rotating black hole with mass M and angular momentum J can be explained by the line element (c=G=1 units)

$$ds^2 = -\left(1 - \frac{2Mr}{\rho^2}\right) dt^2 - \frac{4Ma^*r \sin^2 \theta}{\rho^2} d\phi dt + \frac{\rho^2}{\Delta} dr^2 + \rho^2 d\theta^2 + \left(r^2 + a^{*2} + \frac{2Mr a^{*2} \sin^2 \theta}{\rho^2}\right) \sin^2 \theta d\phi^2$$

- ❖ The equation of state parameter  $\omega$  which we calculated as

$$\omega = 0.01971 \left(\frac{t}{t_0}\right)^2 - 1.0442 \left(\frac{t}{t_0}\right) \quad (6)$$

- ❖ The accretion of black hole mass with the rate is given by

$$\dot{M}_{acc} = 4\pi f R_{BH}^2 \rho_R \quad (7)$$

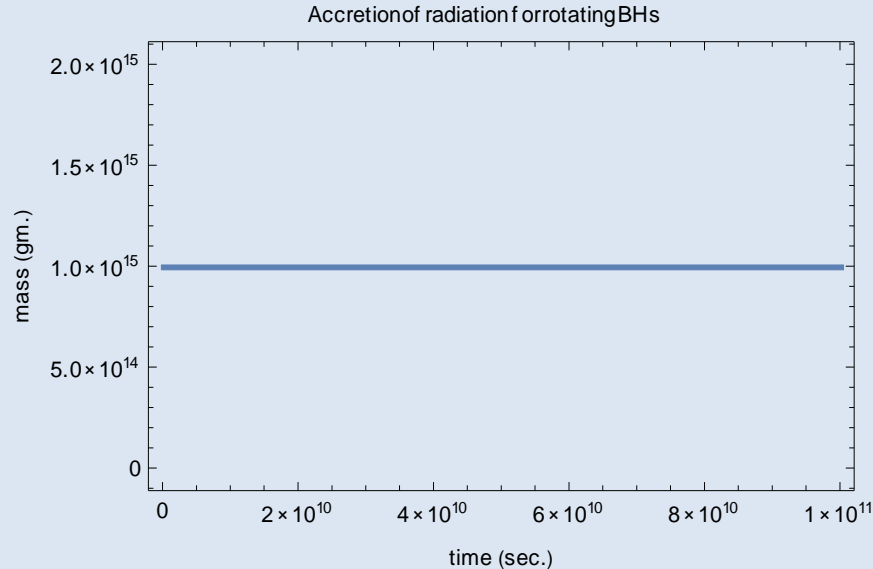
- ❖ The evaporation of black hole mass with the rate is given by

$$\dot{M}_{evp} = -4\pi R_{BH}^2 \sigma_H T_{BH}^4 \quad (8)$$

# Calculation

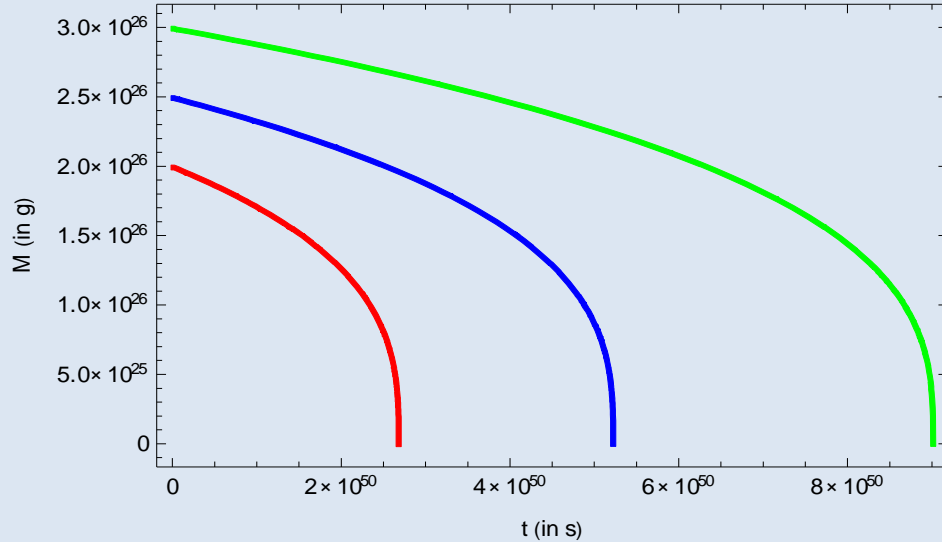
❖ Now the modified accretion of black hole mass in the radiation dominated era is

$$\dot{M}_{acc} = 4\pi f R_{BH}^2 \rho_0 (1 + \omega) \left[ \frac{\rho_0}{\rho_c} + \left\{ 2\sqrt{\frac{8\pi G}{3}} \rho_0^{\frac{1}{2}} (t - t_e) + \frac{3}{2}\sqrt{\frac{8\pi G}{3}} \rho_0^{\frac{1}{2}} (t_e - t_o) + \left(1 - \frac{\rho_0}{\rho_c}\right)^{\frac{1}{2}} \right\}^2 \right]^{-1} \quad (9)$$



❖ Now the evolution of rotating black hole mass in the matter dominated era is

$$\dot{M}_{evo} = \dot{M}_{acc} + \dot{M}_{evp} = 4\pi f R_{BH}^2 \rho_0 (1 + \omega) \left[ \frac{\rho_0}{\rho_c} + \left\{ \frac{3}{2} \sqrt{\frac{8\pi G}{3}} \rho_0^{\frac{1}{2}} (t - t_o) + \left(1 - \frac{\rho_0}{\rho_c}\right)^{\frac{1}{2}} \right\}^2 \right]^{-1} - 4\pi R_{BH}^2 \sigma_H T_{BH}^4 \quad (10)$$

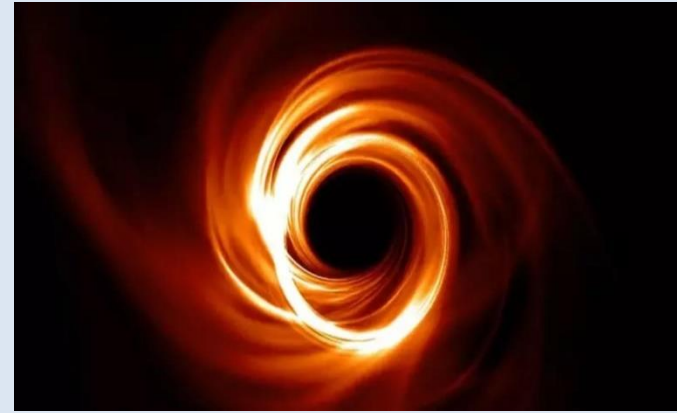
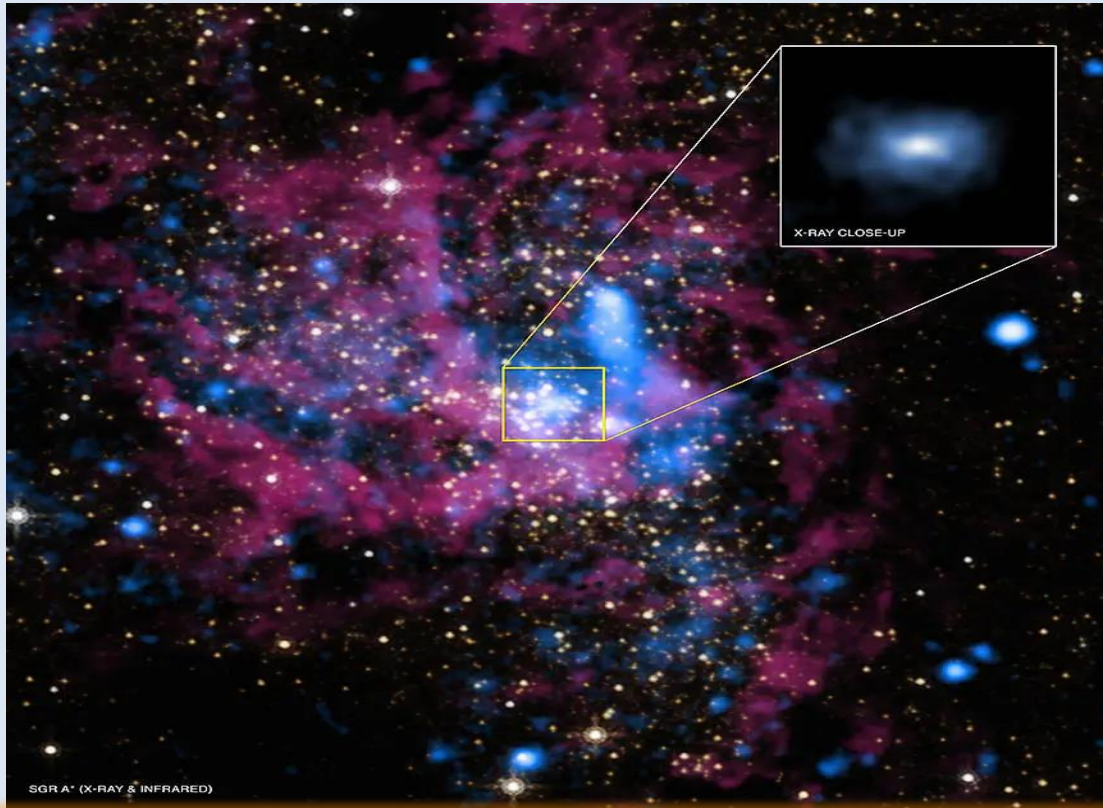


- ❖ In this research work, we have checked the effect of rotating parameter on the evaporation of rotating black holes

Table 2: Variation of the evaporation time with rotating parameter ( $a^*$ ) for fixed value of RBH mass and time, and accretion efficiency ( $f$ ).

$t_i = 10^{-13}$ s, $M_i = 10^{25}$ g, $f = 0.4$	
$(a^*)^2$	$t_{evp}$ (in $10^{18}$ s)
$10^{-12}m_i^2$	33333.32124
$10^{-11}m_i^2$	33333.32124
$10^{-10}m_i^2$	33333.32126
$10^{-9}m_i^2$	33333.32192
$10^{-8}m_i^2$	33333.32274

# Heart of the Milky way



$$\approx 10^6 M_{\odot}$$

Image credit: X-ray: NASA/UMass/D.Wang et al., IR:  
NASA/STScI



# Quasar J0313-1806



$$\approx 10^9 M_{\odot}$$

$$\geq 10^{48} \text{ gm}$$

**Swain, S., Sahoo, G. & Nayak, B.**  
Unveiling the evolution of  
rotating black holes in loop  
quantum cosmology. *Sci Rep* **14**,  
16928 (2024).

<https://doi.org/10.1038/s41598-024-68000-x>

Image credit: NASA/ESA/ESO/Wolfram Freudling et al.  
(STECF)

# Conclusion

- ❖ Accretion of radiation does not affect significantly on the evolution of RBH in loop quantum cosmology.
- ❖ Here we found that RBHs formed in the early radiation dominated era evaporated quickly than the RBHs formed in the later time period.
- ❖ Also, we examined that rotating parameter does not affect vitally on the evaporation of RBHs.
- ❖ we also successfully investigated that those RBHs having mass greater than equal to  $10^{48}$  gm. they all would have been evaporated by the present time due to the accretion of dark energy.

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Thank  
You

