



Multiplicity Dependence of Non-extensive Parameters for Strange and Multi-Strange Particles in Proton-Proton Collisions at $\sqrt{s} = 7$ TeV at the LHC

(Eur. Phys. J. A (2017) 53: 103)

9th International Workshop on Multiple Partonic Interactions at the LHC, Shimla

Arvind Khuntia
Indian Institute of Technology Indore, India

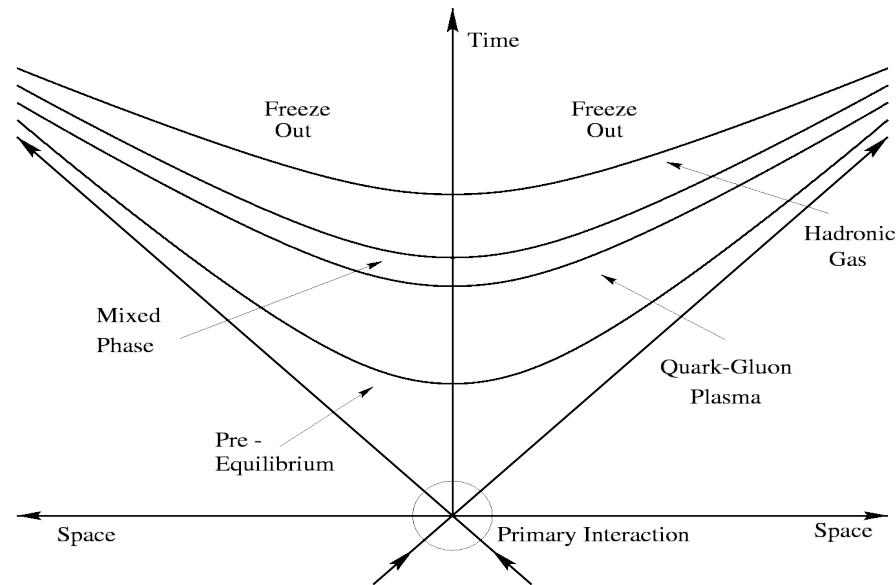
Collaborators: Sushanta Tripathy,
Prof. Raghunath Sahoo and
Prof. Jean Cleymans

OUTLINE

- ❖ Introduction
- ❖ Motivation
- ❖ Tsallis Non-extensive statistics in high energy collisions
- ❖ Particle Spectra in high energy collisions
- ❖ Multiplicity dependence of thermodynamic parameters
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Introduction

High energy collision \longrightarrow QGP \longrightarrow Hadron



- ❖ Transverse momentum spectra of hadrons are important tools to understand the dynamics of high energy collisions.
- ❖ The produced hadrons from the collisions may carry information about the collision dynamics and the subsequent space-time evolution of the system till the occurrence of the final freeze-out.

Motivation

Why is high multiplicity pp interesting?

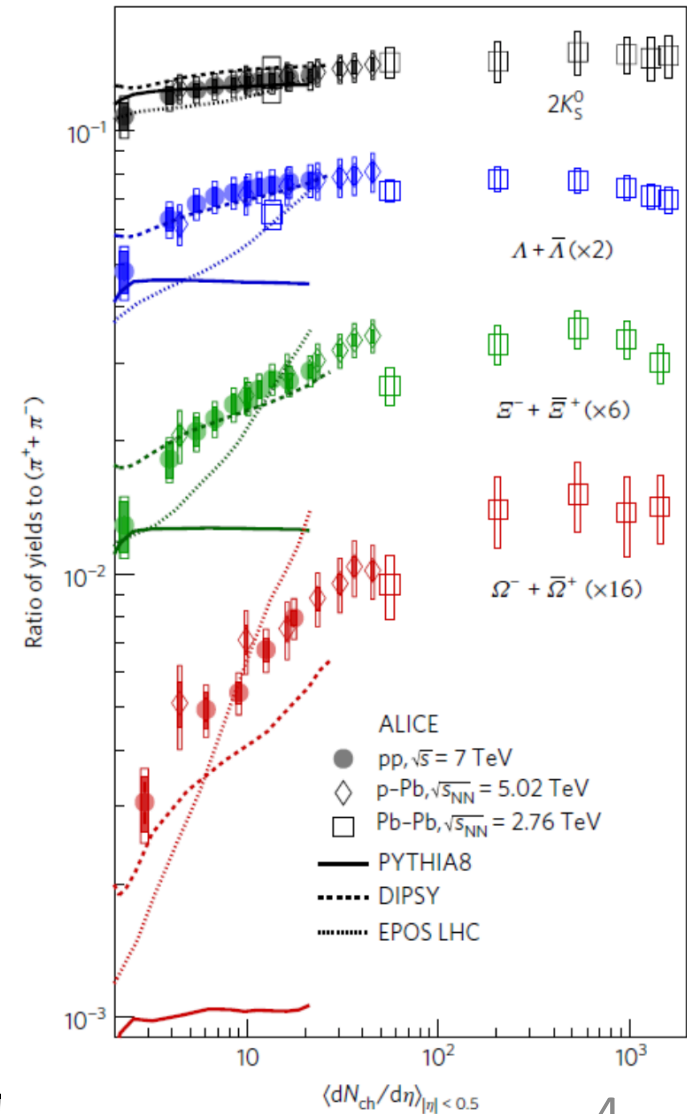
- Probability to have more than one hard scattering is more \rightarrow leads to high multiplicity

Recent results in pp:

In high-multiplicity events strangeness production reaches values similar to those observed in Pb–Pb collisions, where a QGP is formed [1].

- Does the high multiplicity lead to thermalisation in pp collisions?

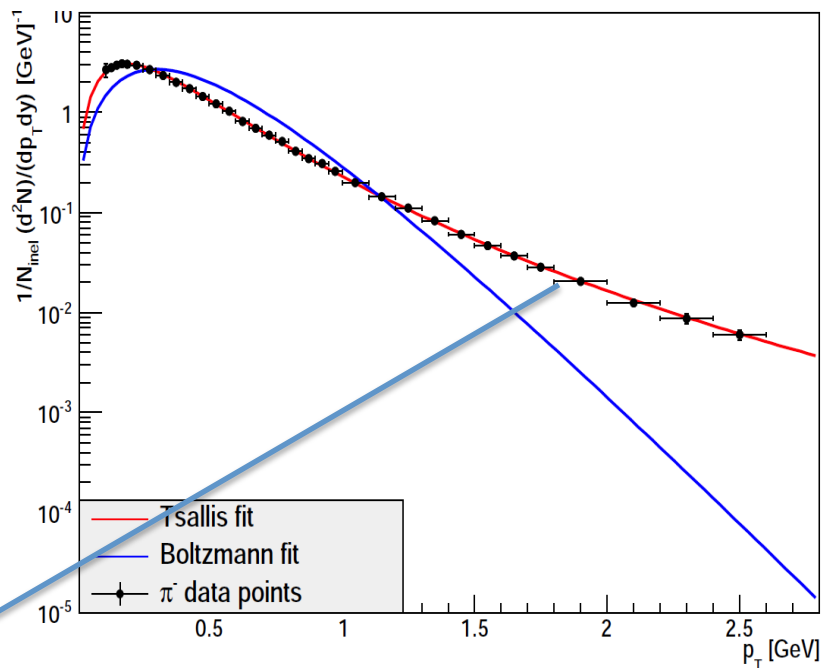
[1] Nature Physics 13, 535–539 (2017)



Transvers Momentum Spectra in High Energy Collisions

- Long back a statistical description of transverse momenta of final state particles produced in high energy collision have been proposed to follow a thermalized Boltzmann type of distribution

$$E \frac{d^3 \sigma}{d^3 p} \simeq C \exp \left(-\frac{p_T}{T} \right).$$



Eur. Phys. J C 71 (2011) 1655

- Experiments at RHIC and LHC observe non-exponential behavior at large transverse momenta

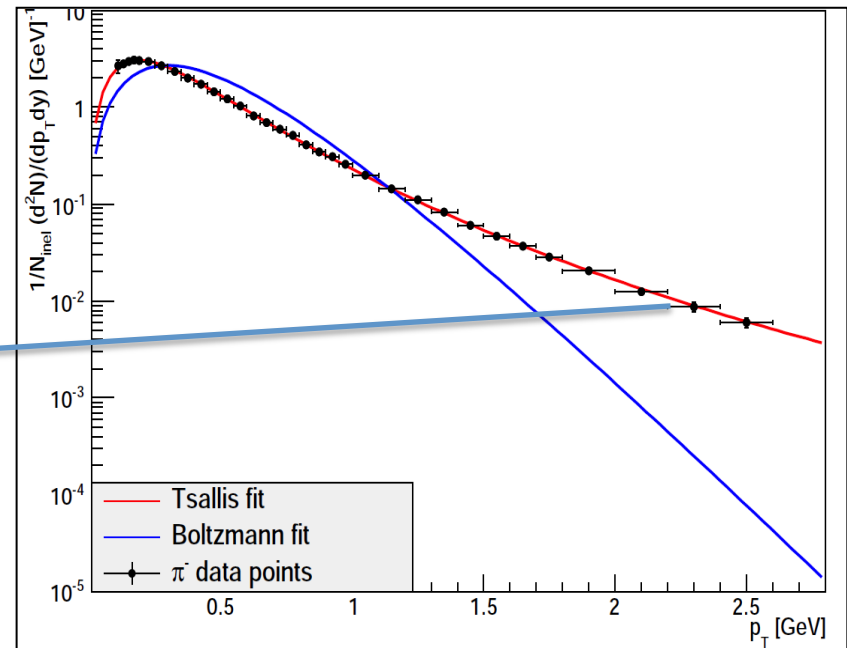
Transverse Momentum Spectra in High Energy Collisions

To account for this non-exponential behaviour, Hagedorn proposed an empirical formula which is given by,

$$E \frac{d^3\sigma}{d^3p} = C \left(1 + \frac{p_T}{p_0}\right)^{-n} \longrightarrow \begin{cases} \exp\left(-\frac{np_T}{p_0}\right) & \text{for } p_T \rightarrow 0, \\ \left(\frac{p_0}{p_T}\right)^n & \text{for } p_T \rightarrow \infty, \end{cases}$$

However, the Tsallis formula based on non-extensive entropy, accounts for the high-energy behaviour of the observed spectra and is given by,

$$f(p_T) = C_q \left[1 + (q-1) \frac{p_T}{T}\right]^{-\frac{1}{q-1}}$$



The Tsallis Non-extensive Statistics

The thermodynamically consistent Tsallis-Boltzmann distribution function is given by,

$$f = \left[1 + (q - 1) \frac{E - \mu}{T} \right]^{-\frac{1}{q-1}} \xrightarrow{(q \rightarrow 1)} e^{-(E - \mu)/T}$$

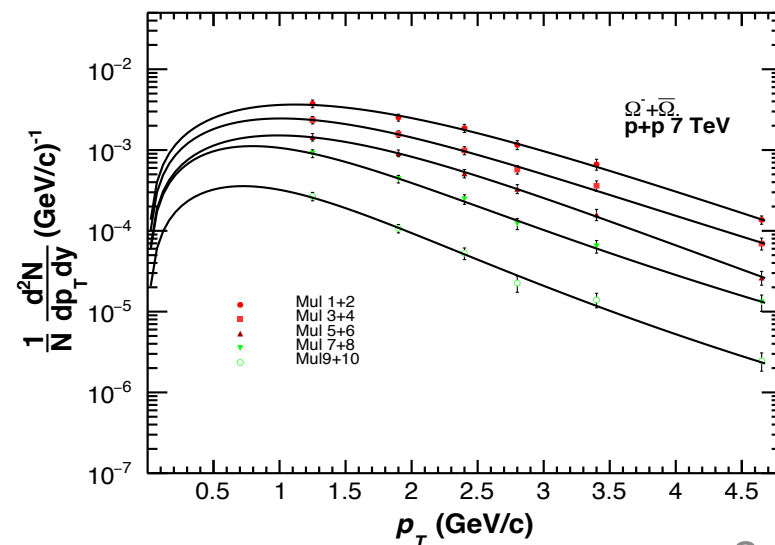
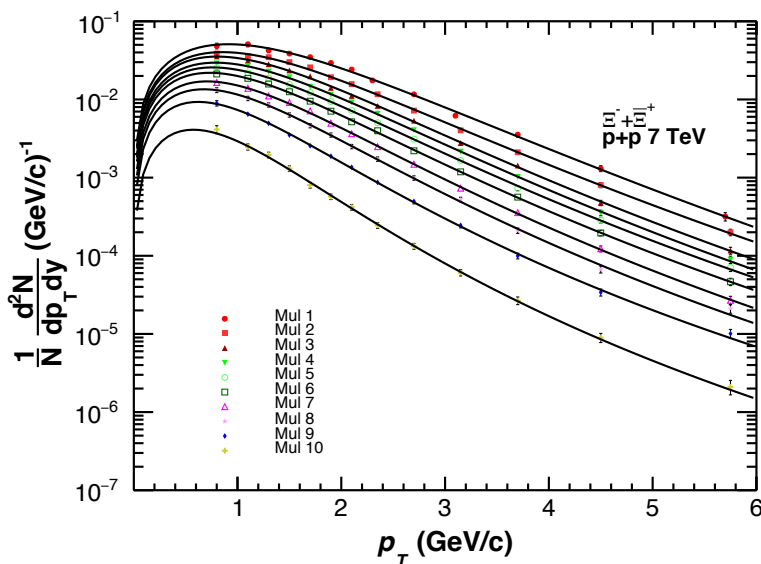
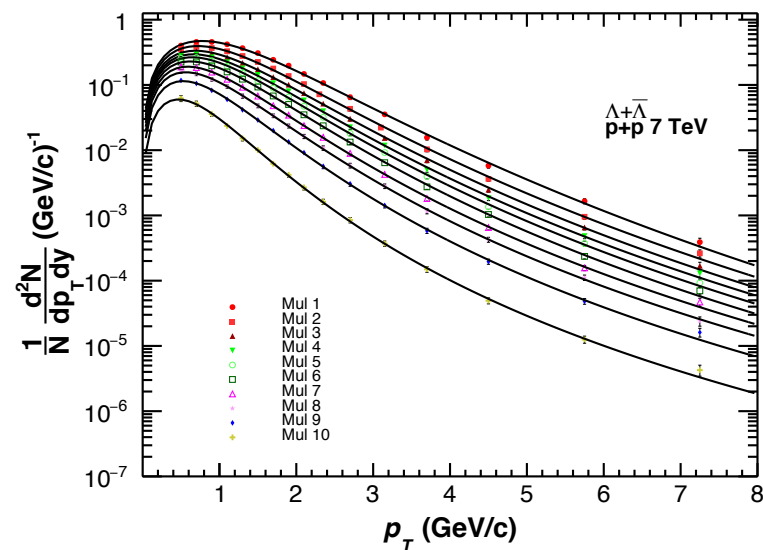
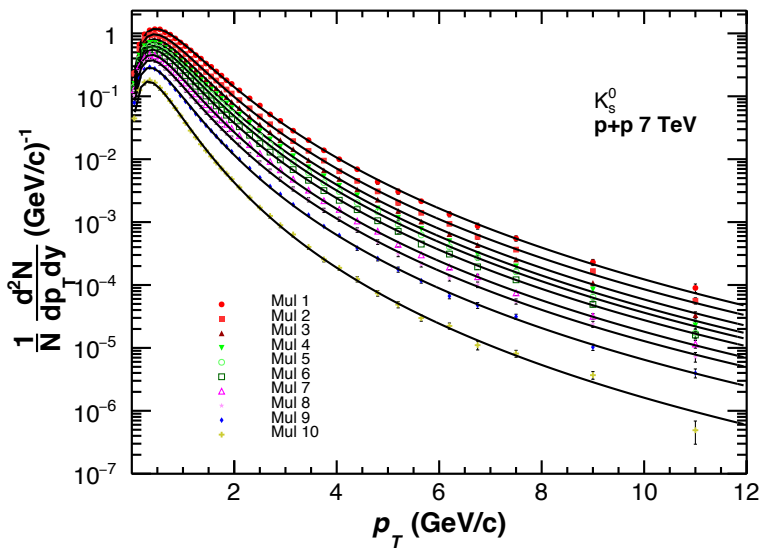
The invariant yield at mid-rapidity is given by,

$$\frac{1}{p_T} \frac{d^2 N}{dp_T dy} \Big|_{y=0} = \frac{gV m_T}{(2\pi)^2} \left[1 + (q - 1) \frac{m_T - \mu}{T} \right]^{-\frac{q}{q-1}}$$

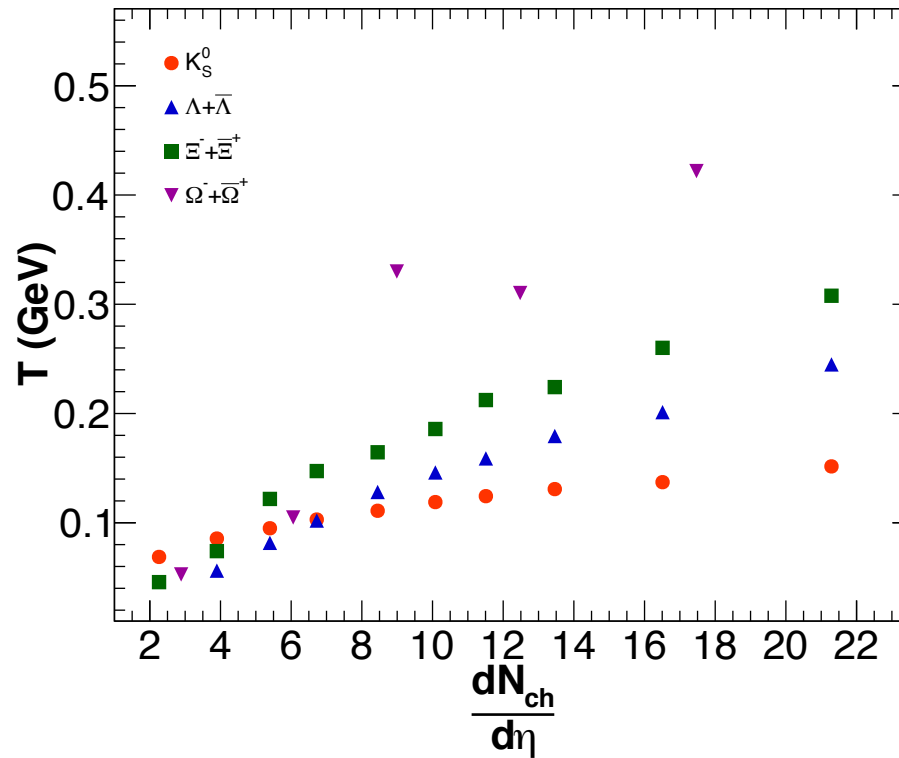
at LHC energies $\mu \sim 0$, so the invariant yield becomes

$$\frac{1}{p_T} \frac{d^2 N}{dp_T dy} \Big|_{y=0} = \frac{gV m_T}{(2\pi)^2} \left[1 + (q - 1) \frac{m_T}{T} \right]^{-\frac{q}{q-1}}$$

Fitting with Tsallis Distribution Function

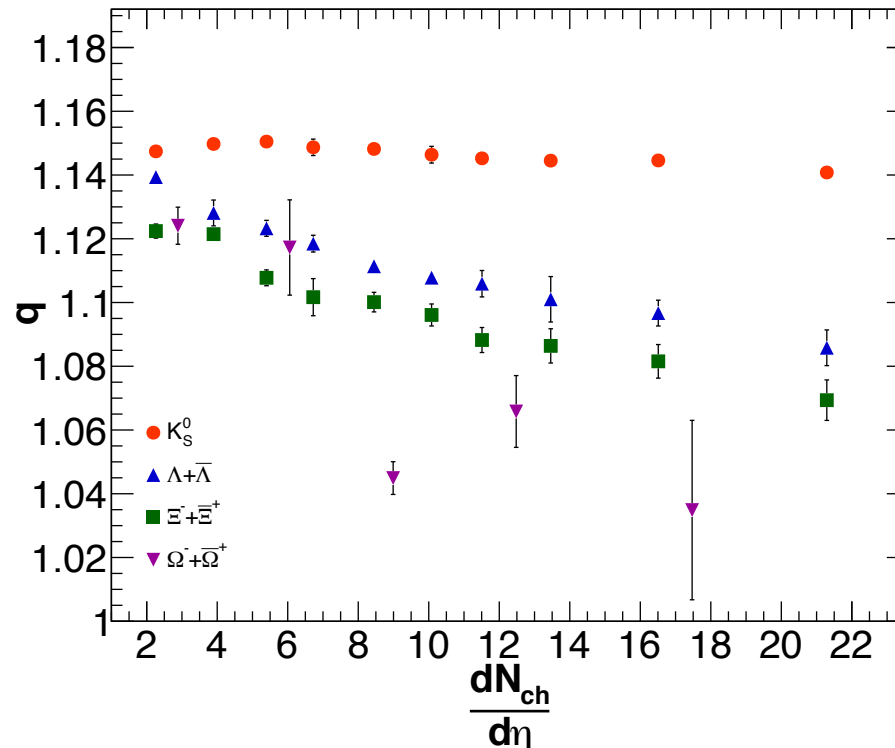


Tsallis Temperature (T)



- The variable T shows a systematic increase with multiplicity, the heaviest baryons showing the steepest increase.
- This is an indication of a mass hierarchy in particle freeze-out: leading to a differential freeze-out scenario (D.Thakur, et al. Adv.High Energy Phys. (2016) 4149352)

Tsallis Non-extensive Parameter (q)



- q decreases towards the value 1 as the multiplicity increases, except for the K_S^0 which shows no clear dependence.
- This shows the tendency of the produced system to equilibrate with higher multiplicities.
- This goes inline with the expected multipartonic interactions, which increase for higher multiplicities in $p + p$ collisions and are thus responsible for bringing the system towards thermodynamic equilibrium.

Summary

- Tsallis distribution provides a very good description of the transverse momentum distributions of strange and multi-strange particles produced in pp collisions at $\sqrt{s} = 7$ TeV
- The variable T shows a systematic increase with multiplicity, the heaviest baryons showing the steepest increase
- q decreases towards the value 1 as the multiplicity increases, except for the K_s^0 which shows no clear dependence
- This shows the tendency of the produced system to equilibrate with higher multiplicities
- This goes inline with the expected multipartonic interactions, which increase for higher multiplicities in p + p collisions and are thus responsible for bringing the system towards thermodynamic equilibrium

Thanks for your attention!

The Tsallis Non-extensive Statistics

Thermodynamic consistency

$$T = \left. \frac{\partial \epsilon}{\partial s} \right|_n \quad \mu = \left. \frac{\partial \epsilon}{\partial n} \right|_s$$

$$n = \left. \frac{\partial P}{\partial \mu} \right|_T \quad s = \left. \frac{\partial P}{\partial T} \right|_\mu$$

Here

$s = S/V$ (entropy density)

$n = N/V$ (particle number density)

$$S_q(A, B) = S_q(A) + S_q(B) + (1 - q)S_q(A)S_q(B)$$

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$$E \frac{d^3 \sigma}{d^3 p} \simeq C \exp \left(-\frac{p_T}{T} \right).$$

Experiments observe non-exponential behavior at large transverse momenta.

$$E \frac{d^3 \sigma}{d^3 p} = C \left(1 + \frac{p_T}{p_0} \right)^{-n} \begin{matrix} \xrightarrow{\hspace{2cm}} \exp \left(-\frac{np_T}{p_0} \right) & \text{for } p_T \rightarrow 0, \\ \left(\frac{p_0}{p_T} \right)^n & \text{for } p_T \rightarrow \infty, \end{matrix}$$

$$f(p_T) = C_q \left[1 + (q-1) \frac{p_T}{T} \right]^{-\frac{1}{q-1}}$$

$$S_q(A, B) = S_q(A) + S_q(B) + (1-q) S_q(A) S_q(B)$$

$$S_q = \frac{1 - \sum p_i^q}{q-1}$$