



# Precision Cosmology of Atomic Dark Matter

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Based on [arxiv:2212.02487](https://arxiv.org/abs/2212.02487) w/ Saurabh Bansal, David Curtin, Yuhsin Tsai

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# Some Motivations

## Experiment

- Hubble tension: Additional dark radiation allows for a higher  $H_0$  value from CMB measurements by reducing the sound horizon at recombination.
- $S_8$  tension: Interactions between dark matter and dark radiation can suppress the growth of structure, yielding a lower  $S_8$ .

## Theory

- Approximately  $\mathbb{Z}_2$ -symmetric mirror sectors can address naturalness issues.
- e.g. Mirror Twin Higgs (see Yuhsin's talk).
- Twin/mirror sectors have their own gauge forces, and can have dark matter – dark radiation interactions.
- What happens if we also allow the dark photon coupling to vary from the SM value?

# Atomic Dark Matter

- Dark sector fermions  $p_D$  and  $e_D$  with masses  $m_{p_D}, m_{e_D}$ , oppositely charged under a dark  $U(1)$  gauge force mediated by a massless dark photon with dark fine structure constant  $\alpha_D$ .

$$\mathcal{L}_{aDM} = A_{\mu\nu}A^{\mu\nu} + i\bar{p}_D(\not{D} - m_{p_D})p_D + i\bar{e}_D(\not{D} - m_{e_D})e_D$$

- Asymmetric, overall neutral abundance that accounts for a fraction  $f_D \equiv \frac{\Omega_{aDM}}{\Omega_{DM}} \leq 1$  of the dark matter.
- Temperature ratio of the dark photons to SM photons today  $\xi \equiv T_D/T_{SM}$ . Equivalently,  $\Delta N_D \equiv \left(\frac{8}{7}\right) \left(\frac{11}{4}\right)^{\frac{4}{3}} \xi^4$ , effective number of additional neutrinos at CMB.
- First studied by Kaplan et al. in 0909.0753.

# Cosmology of Atomic Dark Matter: Quick Review

- In the early universe, the dark photons, dark protons, and dark electrons form a tightly coupled plasma.
- The competition between gravity and the pressure support of the dark photons leads to **dark acoustic oscillations** (DAO) in the dark plasma.
- After  $T_D$  falls below  $B_D = \frac{\alpha_D^2 \mu_D}{2}$ , the dark protons and electrons ‘recombine’ into **dark hydrogen**,  $H_D$ .
- When the dark photons no longer efficiently exchange energy with the atomic dark matter, they begin free-streaming.
- The DAOs and additional radiation in the early universe leave distinctive imprints on the CMB and large-scale structure (LSS) of the universe.
- First studied by Cyr-Racine et al. in 1209.5752, 1310.3278.



# Computing Cosmological Evolution of aDM

- The Cosmological Linear Anisotropy Solving System (CLASS) is a numerical code that produces CMB and matter power spectra given cosmological parameter inputs.
- Modified CLASS to include the atomic dark matter sector, and solve for the dark recombination. [https://github.com/jp-barron/class\\_adm-3.1.git](https://github.com/jp-barron/class_adm-3.1.git)
  - Built on top of modifications by Bansal et al. for MTH. [2110.04317]
- After dark recombination, the dark sector perturbation evolution is handled by the Effective Theory of Structure (ETHOS) framework in CLASS. [1512.05344]

ETHOS	DR	$\chi$	$\dot{\kappa}_{DR-DM}$	$\dot{\kappa}_{\chi}$	$c_{\chi}^2$	$\alpha_{\ell=2}$	$\alpha_{\ell \geq 2}$	$\dot{\kappa}_{DR-DR}$	$\beta_{\ell}$
aDM	$\gamma_D$	$H_D$	$-\frac{1}{\tau_D}$	$-\frac{4\rho_{\gamma D}}{3\rho_D}\tau_D^{-1}$	$c_D^2$	9/10	1	0	0

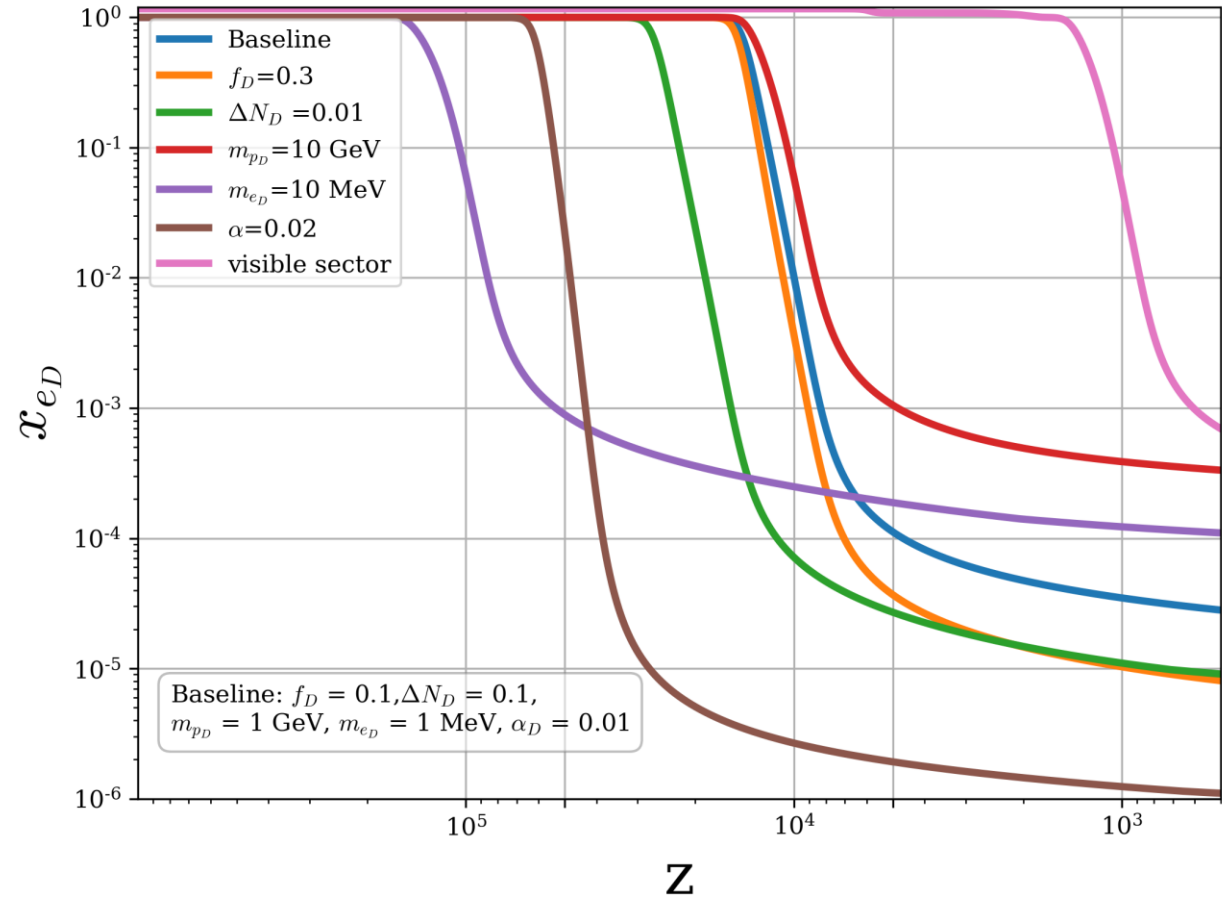
Mapping of aDM parameters to ETHOS parameters

# Dark recombination

- Dark recombination occurs after  $T_D$  falls below  $B_D = \frac{\alpha_D^2 \mu_D}{2}$ .
- Redshift of dark recombination controlled by  $B_D$  and  $\Delta N_D$ .
- In CLASS: Use Saha equation for ionization fraction  $x_D$  above 0.999.

$$\frac{x_D^2}{1 - x_D} = \frac{1}{n_D} \left( \frac{T_{DM} m_{eD}}{2\pi} \right)^{3/2} e^{-B_D/T_{DM}}$$

- Then, solve for  $\frac{dx_D}{dt}$  with effective multi-level atom formalism using HYREC-2. [Lee and Ali-Haïmoud, 2007.14114]
  - Neglects some sub-dominant radiative transfer effects that are only modelled for parameters near SM values.



# Dark decoupling



- When energy transfer rate between dark photons and dark electrons/protons drops below Hubble, they decouple.

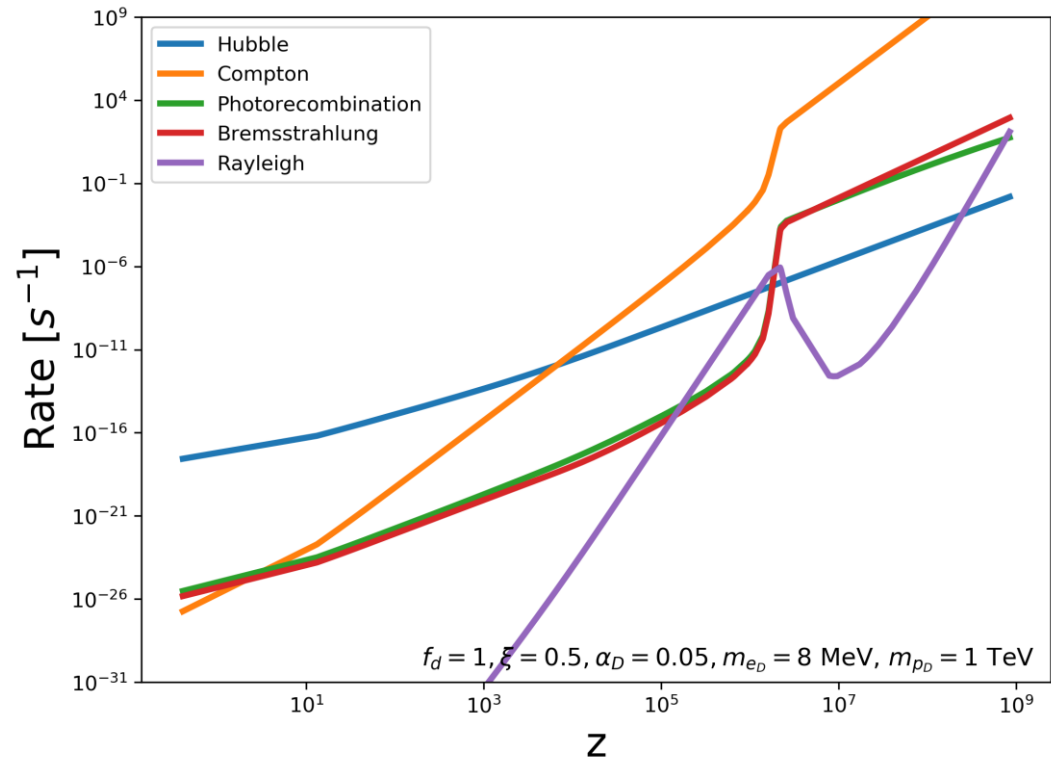
- Dominant energy exchange mechanism is usually Thomson scattering, with cross-section

$$\sigma_{TD} = \frac{8\pi}{3} \left( \frac{\alpha_D}{m_{eD}} \right)^2.$$

- Other processes can dominate over Thomson scattering and delay decoupling

- For high  $\alpha_D$  and  $m_{eD}/m_{HD}$ , Rayleigh scattering.
- For low  $\alpha_D$  and  $\xi$ , bremsstrahlung, photo-recombination/ionization\*.

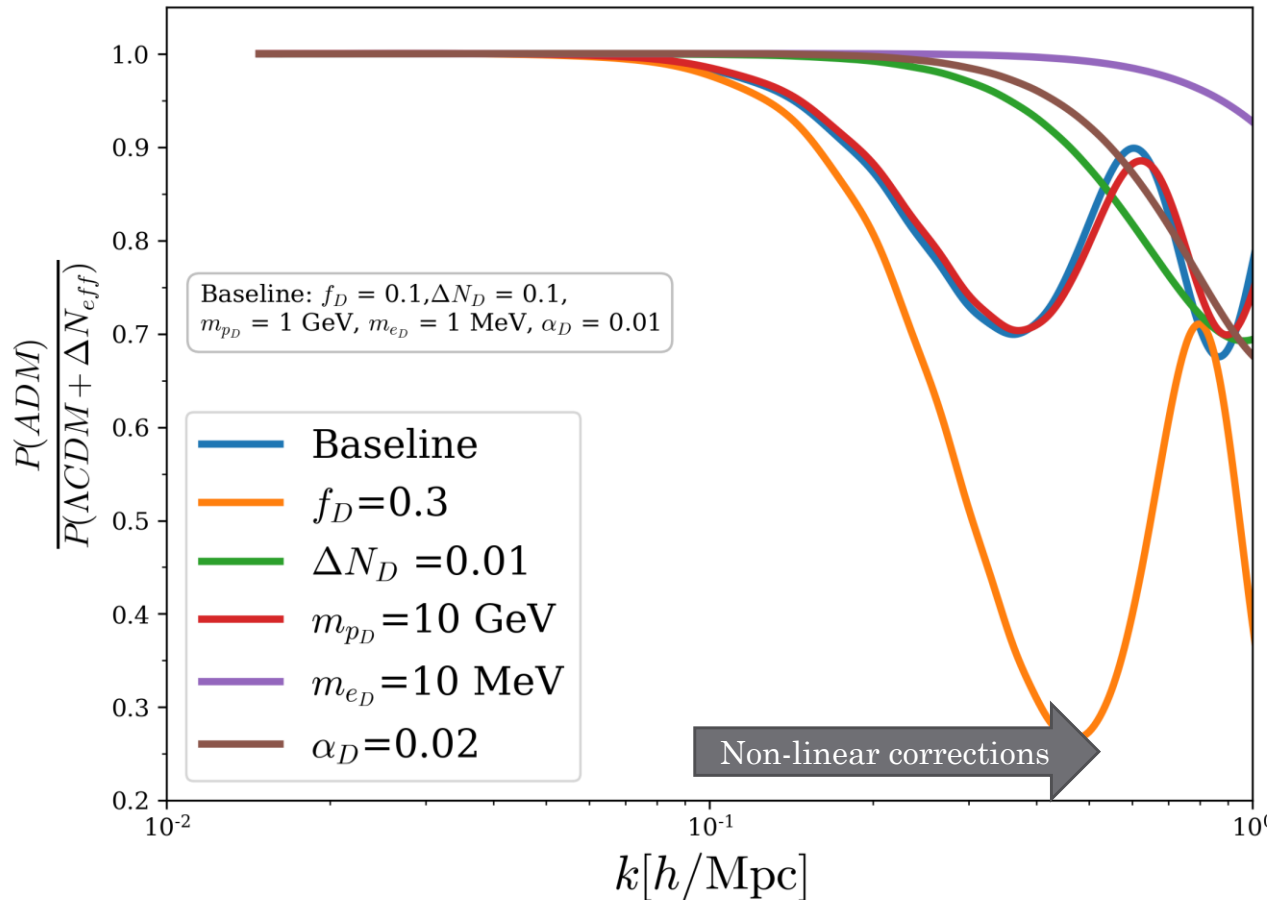
- Opacity of dark plasma set by Thomson, Rayleigh, photo-ionization rates.



\*Photo-heating/cooling processes not yet included in CLASS-aDM due to numerical instabilities – to be added soon. Similar rate to bremsstrahlung.

# Modification of Linear Matter Power Spectrum

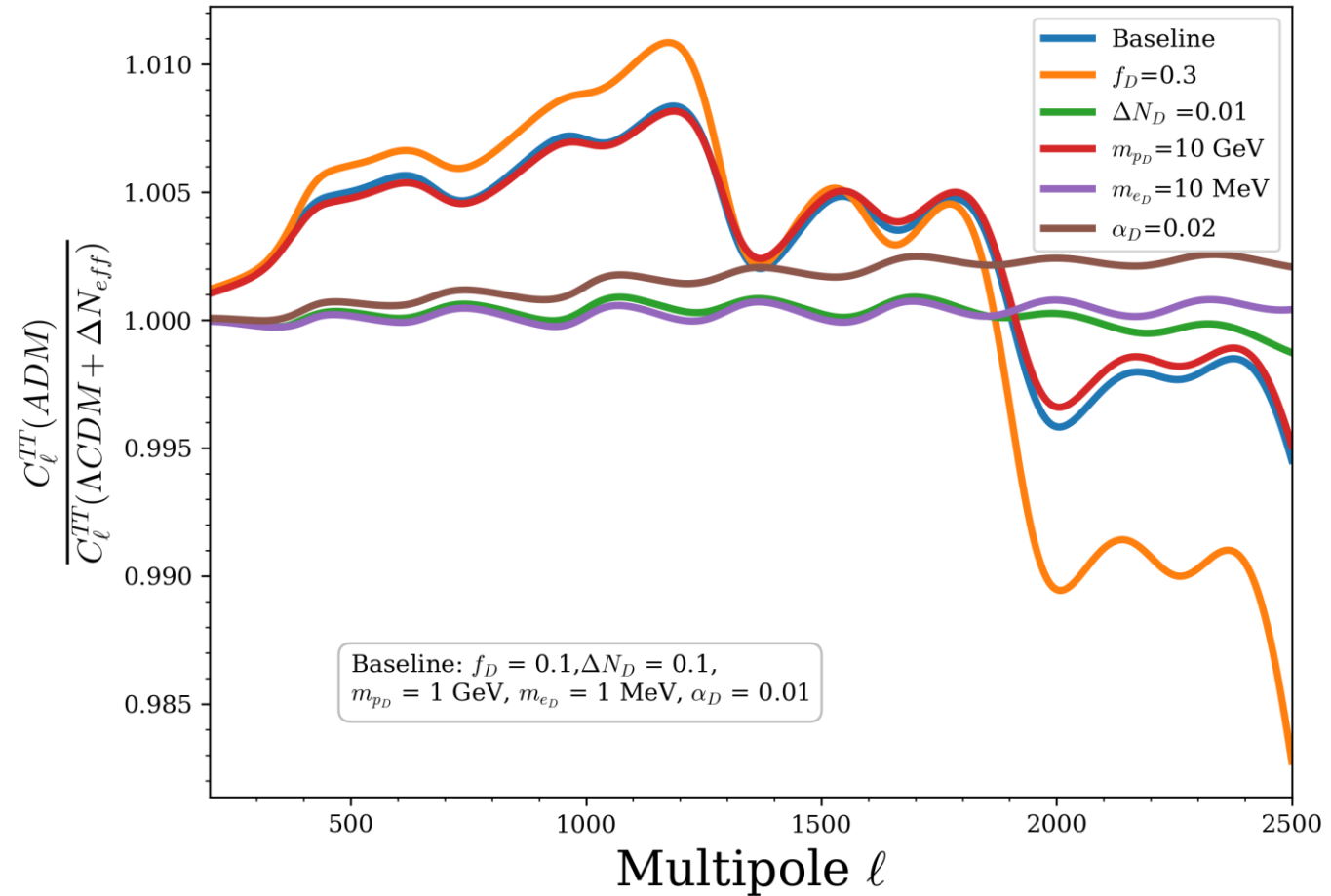
- Suppression and oscillations for  $k$  that enter horizon before dark decoupling.
- Higher  $f_D \rightarrow$  more suppression.
- Lower  $\Delta N_D$ , higher  $B_D \rightarrow$  suppression starts at smaller scales.
- Cannot trust linear power spectrum above  $k \sim 0.1 h \text{ Mpc}^{-1}$ .
- Non-linear evolution smears out oscillations at low redshifts. [2101.12229]



Matter power spectrum relative to  $\Lambda\text{CDM}$  + same amount of  $\Delta N_{eff}$  as  $\Delta N_D$ .

# Modification of CMB Power Spectrum

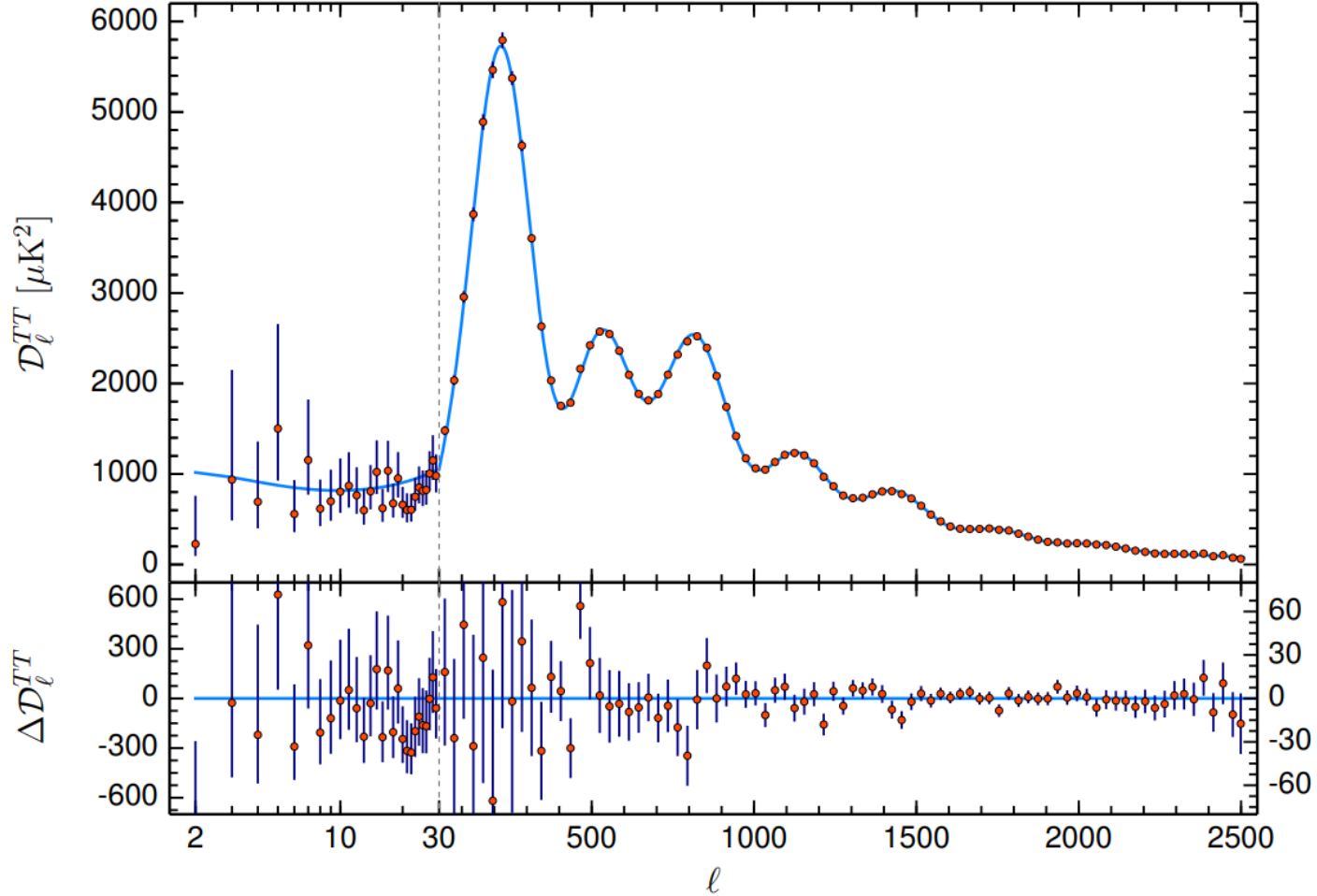
- Additional radiation damps the temperature power spectrum at high  $\ell$ .
- aDM perturbations couple to gravitational potential, which couples to photon perturbations. Complex effect on power spectrum.
- Higher  $f_D \rightarrow$  Larger deviations.
- Lower  $\Delta N_{\text{eff}}$ , higher  $B_D \rightarrow$  smaller deviations.
- Polarization spectrum and lensing also altered.



CMB temperature power spectrum relative to  $\Lambda\text{CDM} +$  same amount of  $\Delta N_{\text{eff}}$  as  $\Delta N_D$ .

# Constraining $\Lambda$ CDM: CMB + BAO

- Planck 2018 TT, TE, EE, high- and low- $l$ , lensing. [1907.12875]
- BAO: Measurements of  $D_V/r_S$  from 6dFGS, SDSS, BOSS. [1106.3366, 1409.3242, 1607.03155]



Planck 2018 temperature power spectrum

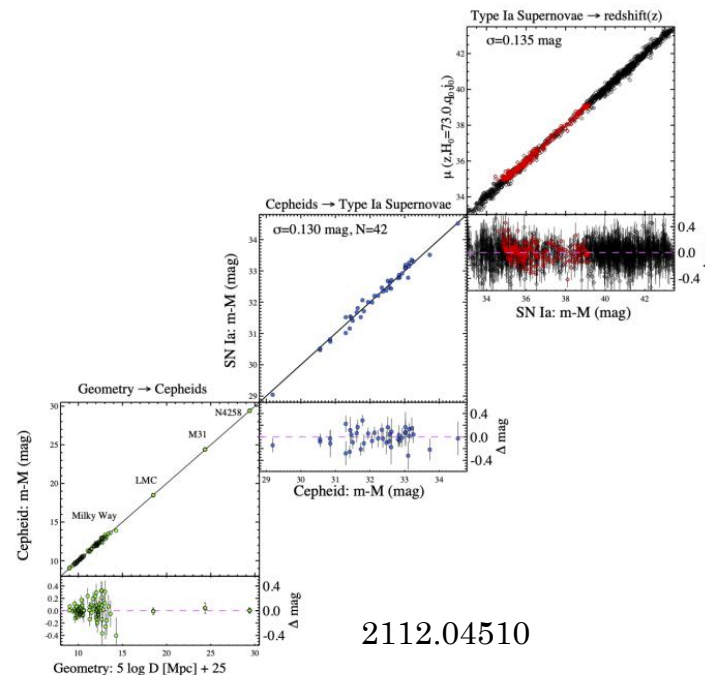


# Constraining aDM: LSS, $H_0$ , $S_8$

- KiDS+VIKING-450 cosmic shear dataset, with cutoff at  $k = 0.3 h \text{ Mpc}^{-1}$ . [1812.06076]
- Pantheon supernova dataset. [1710.00845].
- SH0ES measurement of  $H_0$  using Cepheid variables and Type Ia supernovae. [2112.04510]
- Planck Sunyaev-Zeldovich measurement of  $S_8$ .
  - Gives largest tension, to show how far aDM can pull  $S_8$ . [1303.5080]

Table 4. Fiducial result for the KV450 cosmic shear measurement.

Parameter	Symbol	Value
CDM density	$\Omega_{\text{CDM}} h^2$	$0.118^{+0.038}_{-0.066}$
Scalar spectrum ampl.	$\ln 10^{10} A_s$	$3.158^{+1.154}_{-1.426}$
Baryon density	$\Omega_b h^2$	$0.022^{+0.003}_{-0.004}$
Scalar spectral index	$n_s$	$1.021^{+0.149}_{-0.141}$
Hubble parameter	$h$	$0.745^{+0.073}_{-0.043}$
IA amplitude	$A_{\text{IA}}$	$0.981^{+0.694}_{-0.678}$
Baryon feedback ampl.	$B$	$2.484^{+0.189}_{-0.475}$
Constant $c$ -term offset	$\delta c$	$0.000^{+0.0002}_{-0.0002}$
2D $c$ -term amplitude	$A_c$	$1.022^{+0.129}_{-0.125}$
Redshift offset bin 1	$\delta z_1$	$-0.007^{+0.034}_{-0.034}$
Redshift offset bin 2	$\delta z_2$	$-0.010^{+0.019}_{-0.021}$
Redshift offset bin 3	$\delta z_3$	$0.013^{+0.020}_{-0.021}$
Redshift offset bin 4	$\delta z_4$	$0.001^{+0.012}_{-0.011}$
Redshift offset bin 5	$\delta z_5$	$-0.001^{+0.011}_{-0.011}$
Matter density	$\Omega_m$	$0.256^{+0.064}_{-0.123}$
Power spectrum amplitude	$\sigma_8$	$0.836^{+0.132}_{-0.218}$
$\sigma_8 \sqrt{\Omega_m/0.3}$	$S_8$	$0.737^{+0.040}_{-0.036}$



1812.06076

2112.04510

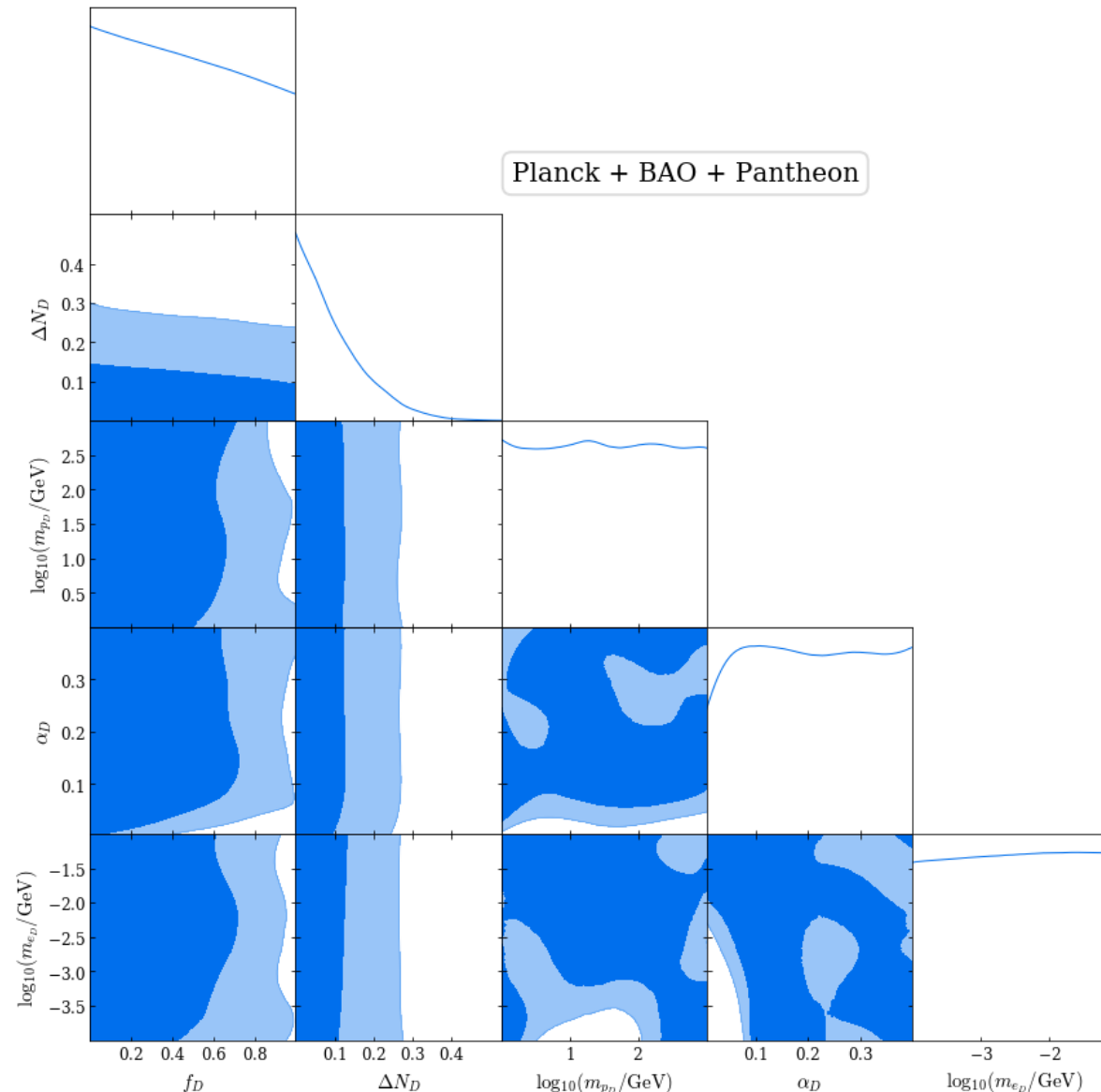


# Constraining aDM parameter space

- Used Markov chain Monte Carlo package MontePython to sample posterior distribution of the model parameters, using publicly available likelihoods for the chosen datasets.
- Generated 95% confidence level contours from posterior.
- Ran  $\geq 6$  chains for each scan.
- Required Gelman-Rubin convergence  $R - 1 < 0.02$  for each parameter.
- CDM parameters:  $\{\omega_b, \omega_{cdm}, h, \ln 10^{10} A_S, n_s, \tau_{reio}\}$ 
  - $h$  instead of  $\theta_s$  because of a bug in CLASS 3.1 shooting procedure.
- Baseline dataset: Planck + BAO + Pantheon. Then add LSS,  $H_0$ ,  $S_8$  measurements.

# Results

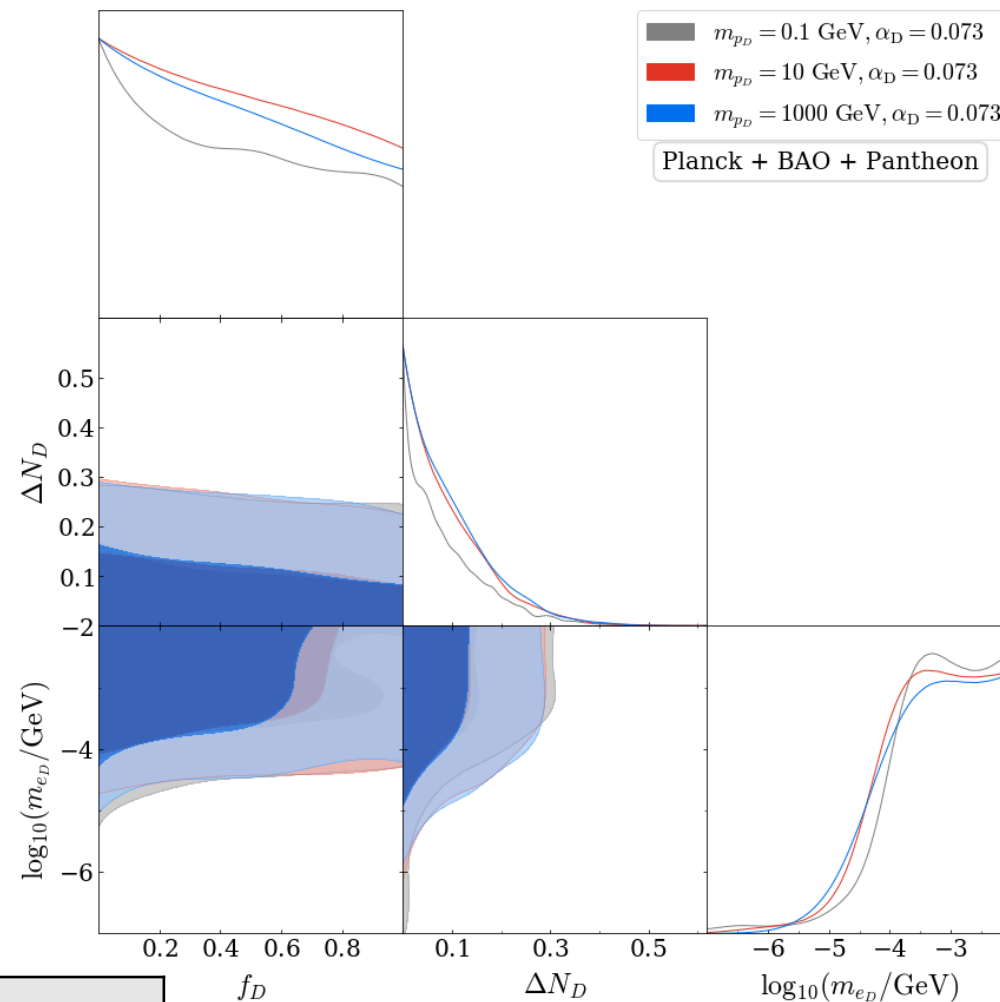
- Allow all five parameters ( $f_D, \Delta N_D, m_{pD}, m_{eD}, \alpha_D$ ) to vary.
- Only  $\Delta N_D$  is robustly constrained.
- Recover Planck limit of  $\Delta N_{\text{eff}} \leq 0.3$ .
- No preference for non-zero  $f_D$ .
- Ranges of priors limited by restriction that dark recombination not be at extreme redshifts ( $> 10^8, \ll 100$ )



Parameter	$f_D$	$\Delta N_D$	$m_{pD}$	$m_{eD}$	$\alpha_D$
Prior	$0 < f_D < 1$	$0 < \Delta N_D < 1$	$0 < \log_{10}\left(\frac{m_{pD}}{\text{GeV}}\right) < 3$	$-4 < \log_{10}\left(\frac{m_{eD}}{\text{GeV}}\right) < -2$	$0.005 < \alpha_D < 0.4$

# 3D scans: constraining $\Delta N_D$ , $f_D$ , and $m_{e_D}$

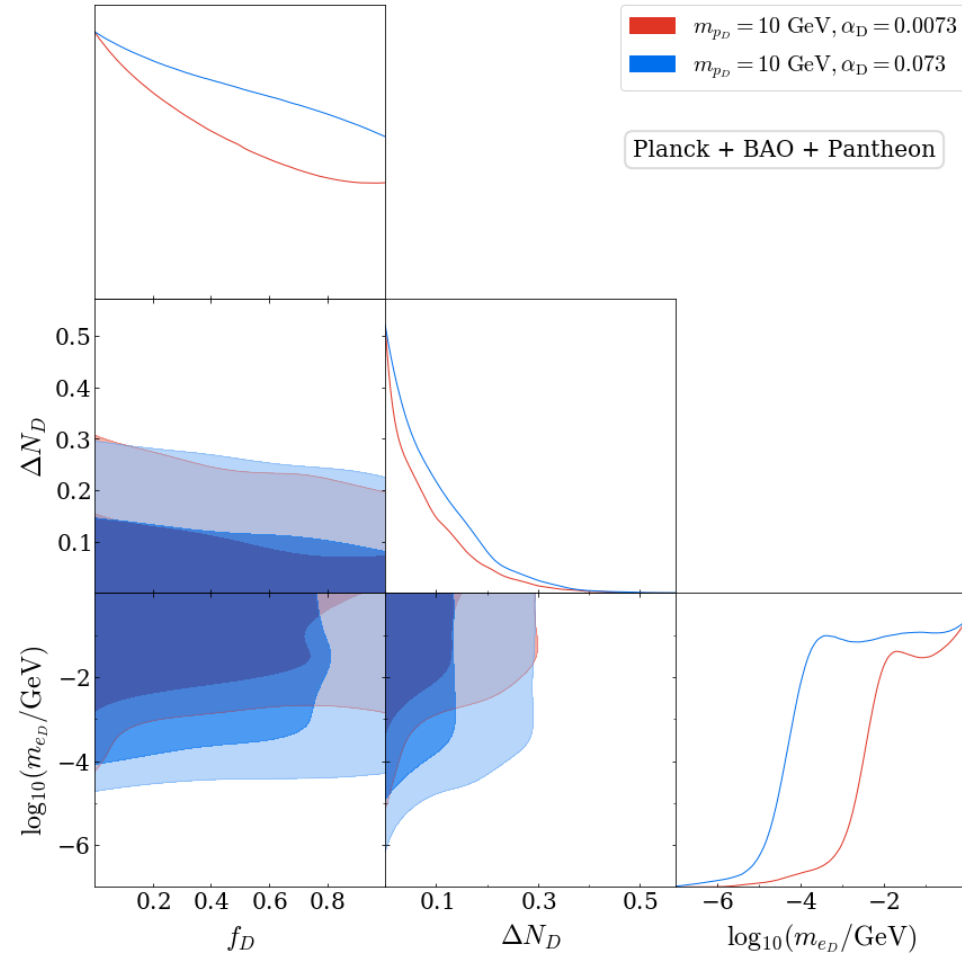
- Varying  $m_{p_D}$  has little effect on constraints on other parameters.
- Constraint on  $\Delta N_D$  sharply transitions from  $< 0.3$  to much tighter as  $m_{e_D}$  ( $\sim B_D, z_{rec}^{dark}$ ) is lowered.



Parameter	$f_D$	$\Delta N_D$	$m_{e_D}$
Prior	$0 < f_D < 1$	$0 < \Delta N_D < 1$	$-7 < \log_{10}(m_{e_D}/\text{GeV}) < 0$

# 3D scans: constraining $\Delta N_D$ , $f_D$ , and $m_{e_D}$

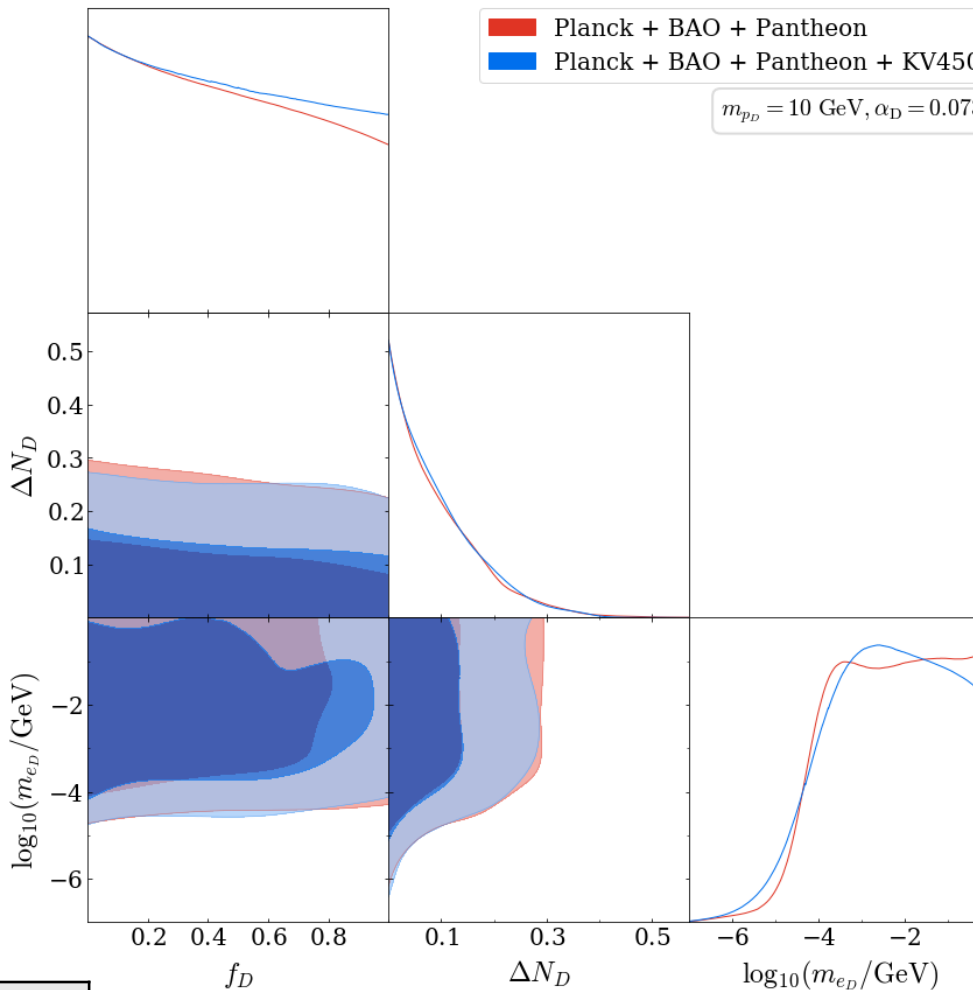
- For higher  $\alpha_D$ , lower bound on  $m_{e_D}$  is lower.
- Large  $f_D$  not ruled out for a range of parameters.



Parameter	$f_D$	$\Delta N_D$	$m_{e_D}$
Prior	$0 < f_D < 1$	$0 < \Delta N_D < 1$	$-7 < \log_{10}\left(\frac{m_{e_D}}{\text{GeV}}\right) < 0$

# Including KV450 dataset

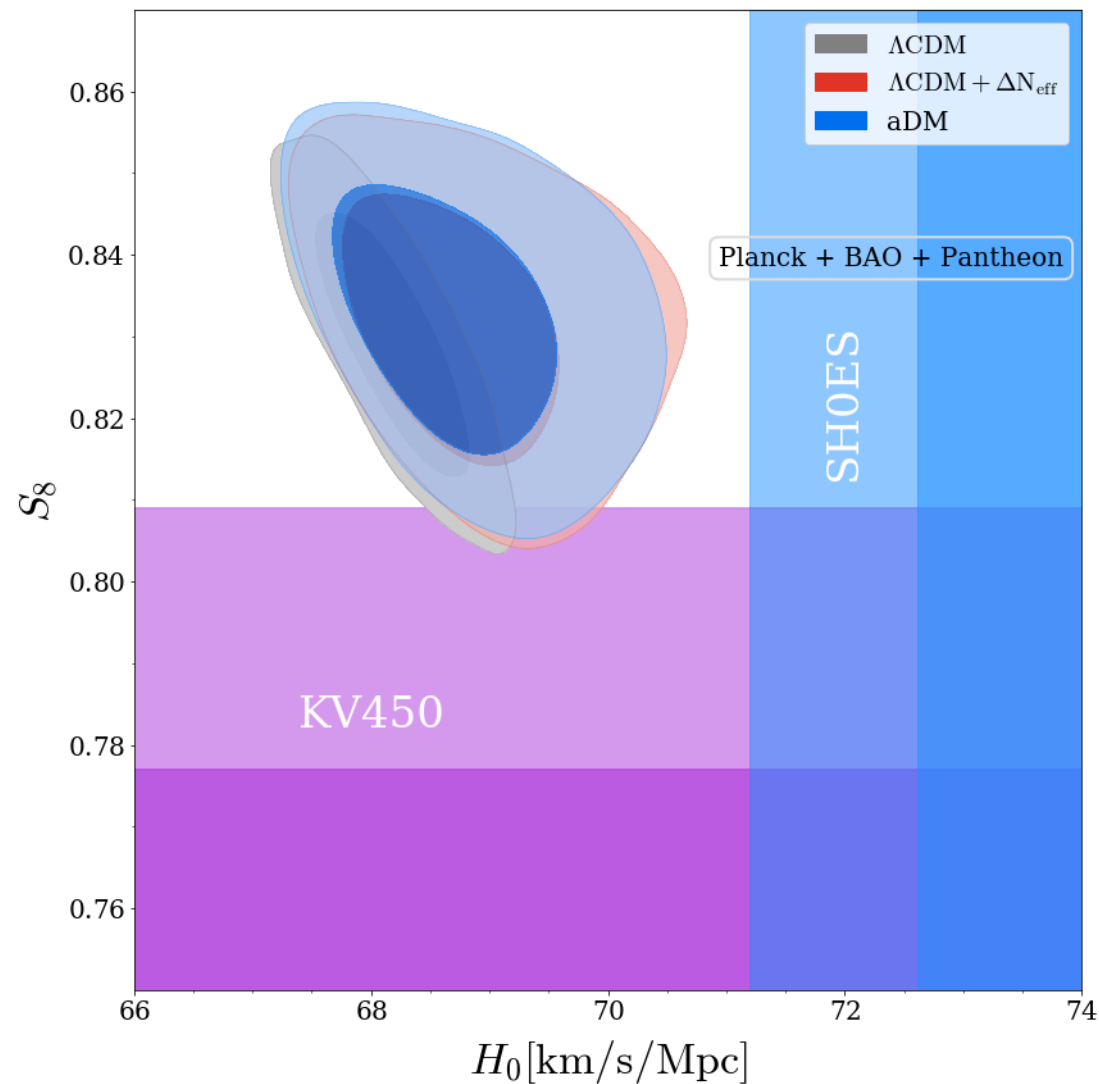
- Constraints almost identical to CMB+BAO.
- Very little constraining power with cut-off in  $k$ .
- Will not include KV450 cosmic shear dataset when comparing with/without local  $H_0$  and  $S_8$  measurements.



Parameter	$f_D$	$\Delta N_D$	$m_{eD}$
Prior	$0 < f_D < 1$	$0 < \Delta N_D < 1$	$-7 < \log_{10}\left(\frac{m_{eD}}{\text{GeV}}\right) < 0$

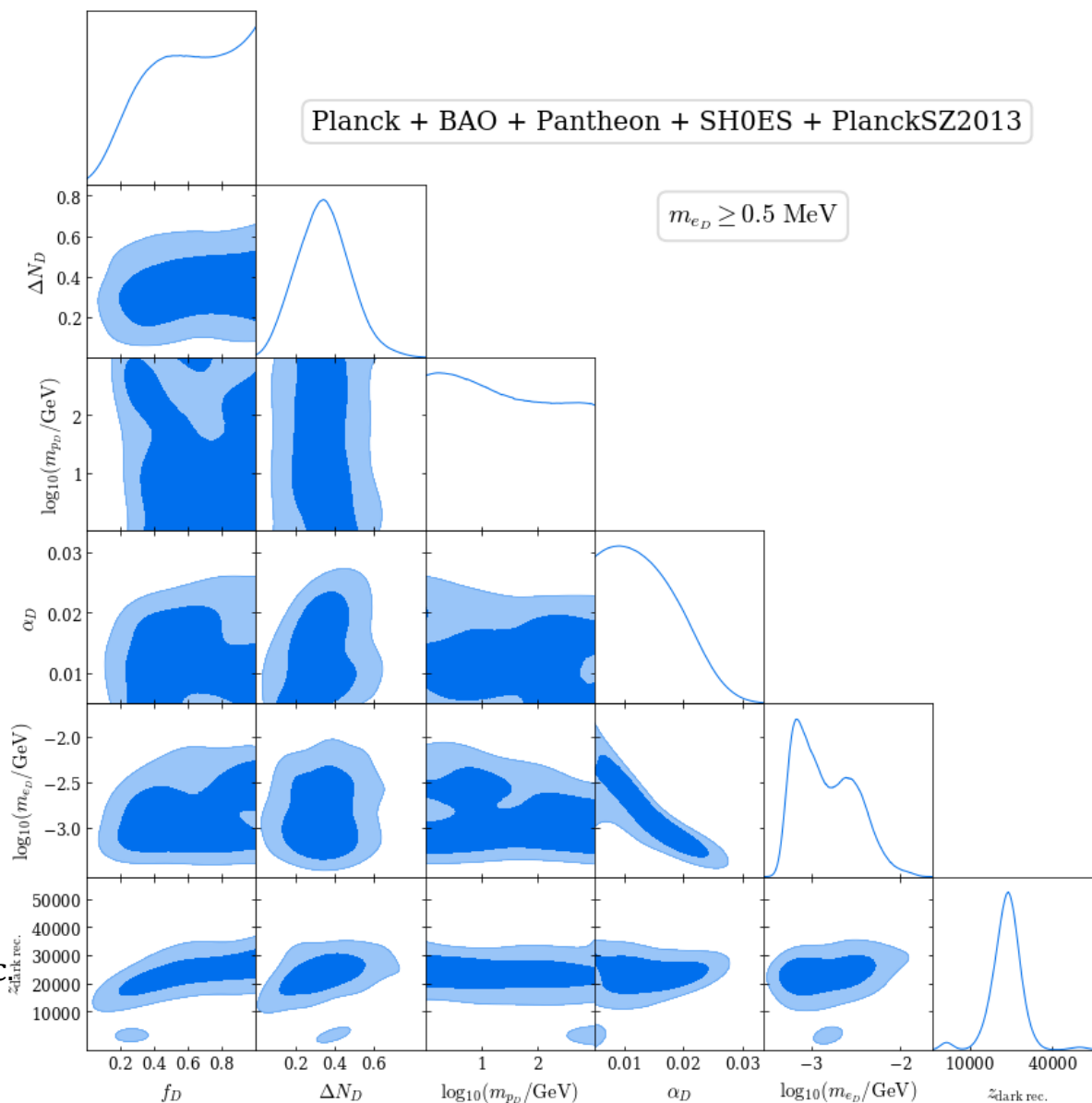
# $H_0$ and $S_8$ : Without direct measurements

- Best-fit region for  $H_0$  extends to higher values than  $\Lambda$ CDM.
- No preference for lower  $S_8$ .
- Very similar contour to  $\Lambda$ CDM +  $\Delta N_{\text{eff}}$ .



# Including local $H_0$ and $S_8$ measurements

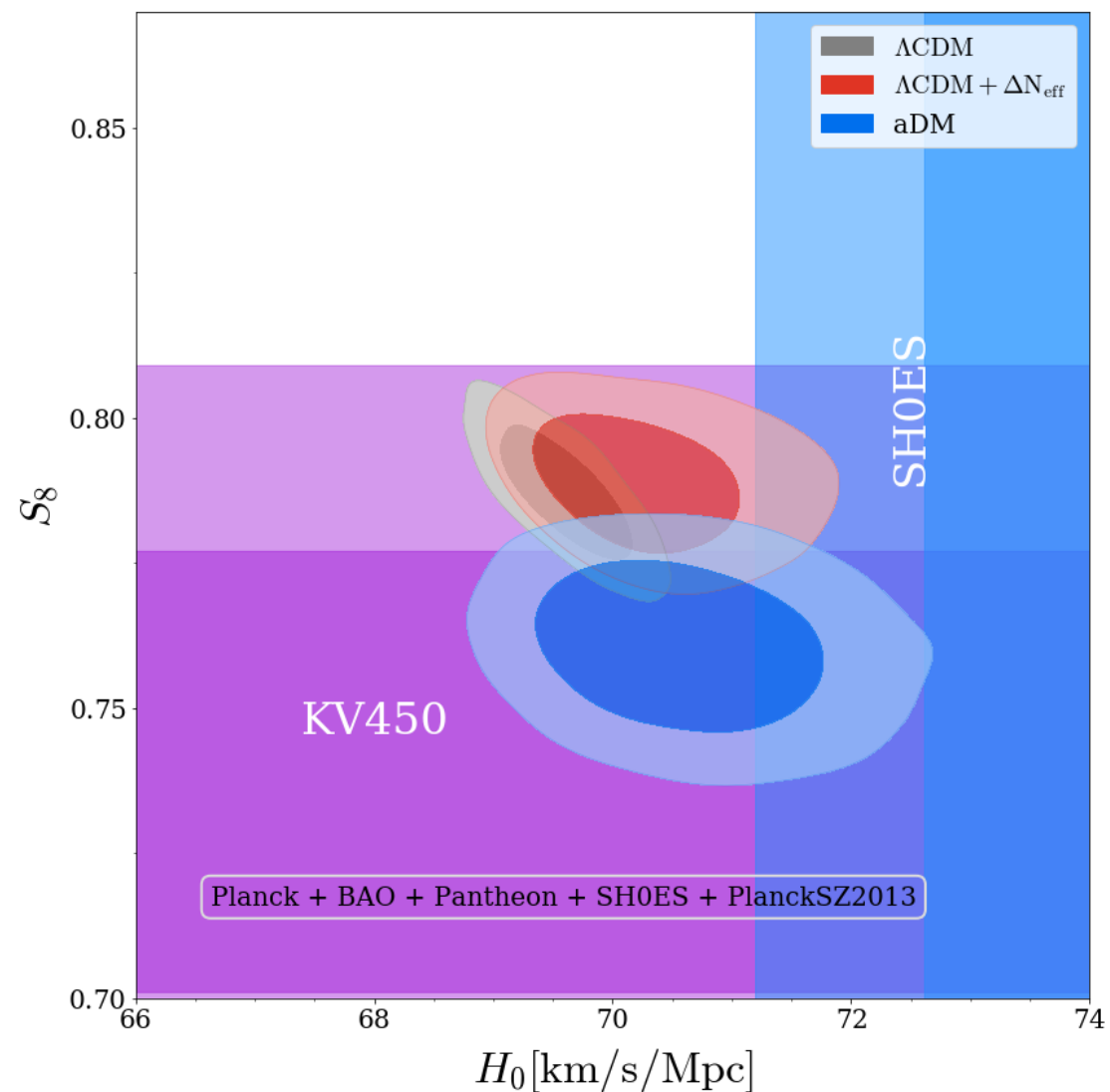
- Preference for non-zero  $\Delta N_D \sim 0.3$  and  $f_D > 0.1$ .
- Strong preference for redshift of dark recombination around  $2 \times 10^3$ .
- No preference for particular value of  $m_{pD}$ .



Lower bound of  $m_{eD} > 0.5$  MeV to avoid violating  $\Delta N_{\text{eff}}$  constraints from BBN. Consistent calculation of  $Y_{\text{He}}$  for arbitrary  $m_{eD}$  in the works now.

# $H_0$ and $S_8$ : With direct measurements

- Larger 95% confidence region.
- Best-fit aDM values pulled further than  $\Lambda$ CDM, with lower minimum  $\chi^2$ .
- $\Lambda$ CDM +  $\Delta N_{\text{eff}}$  unable to accommodate lower  $S_8$ .
- DM-DR interactions necessary to allow lower  $S_8$ .





Param	$\Lambda$ CDM		$\Lambda$ CDM + $\Delta N_{\text{eff}}$		aDM	
	best-fit	mean $\pm\sigma$	best-fit	mean $\pm\sigma$	best-fit	mean $\pm\sigma$
$100 \Omega_b h^2$	2.264	$2.268^{+0.013}_{-0.012}$	2.261	$2.275^{+0.014}_{-0.014}$	2.278	$2.275^{+0.013}_{-0.015}$
$\Omega_{dm} h^2$	0.1165	$0.1163^{+0.00078}_{-0.00079}$	0.1171	$0.1182^{+0.0012}_{-0.0019}$	0.1252	$0.124^{+0.0025}_{-0.0026}$
$h$	0.6956	$0.696^{+0.0036}_{-0.0039}$	0.6958	$0.7027^{+0.0048}_{-0.0069}$	0.7085	$0.7062^{+0.0072}_{-0.009}$
$\ln(10^{10} A_s)$	3.033	$3.024^{+0.014}_{-0.013}$	3.031	$3.026^{+0.015}_{-0.013}$	3.054	$3.052^{+0.015}_{-0.014}$
$n_s$	0.9722	$0.9728^{+0.0037}_{-0.0037}$	0.9724	$0.9764^{+0.0042}_{-0.0047}$	0.973	$0.9721^{+0.004}_{-0.0047}$
$\tau_{\text{reio}}$	0.05129	$0.04859^{+0.0071}_{-0.0069}$	0.04987	$0.0475^{+0.008}_{-0.0068}$	0.05609	$0.05634^{+0.0074}_{-0.0073}$
$H_0$	69.56	$69.6^{+0.36}_{-0.39}$	69.58	$70.27^{+0.48}_{-0.69}$	70.85	$70.62^{+0.72}_{-0.9}$
$\sigma_8$	0.8091	$0.8047^{+0.0052}_{-0.0047}$	0.8098	$0.809^{+0.0058}_{-0.0063}$	0.7629	$0.7679^{+0.0099}_{-0.011}$
$S_8$	0.7921	$0.7871^{+0.0073}_{-0.0078}$	0.7942	$0.7891^{+0.0079}_{-0.008}$	0.7563	$0.7606^{+0.01}_{-0.0098}$
$f_D$	-	-	-	-	0.3265	{0.17, 1}
$\Delta N_D$	-	-	0.026	$0.115^{+0.027}_{-0.11}$	0.3934	$0.339^{+0.13}_{-0.14}$
$\log_{10}(m_{pD}/\text{GeV})$	-	-	-	-	2.987	{0.125, 3}
$\log_{10}(m_{eD}/\text{GeV})$	-	-	-	-	-3.819	{-4, -2.1}
$\alpha_D$	-	-	-	-	0.04041	{.005, .05}
$z_{\text{dark rec.}}$	-	-	-	-	$1.924 \times 10^4$	$22510^{+5.3 \times 10^3}_{-4.2 \times 10^3}$
$\chi^2_{\text{total}}$	3853.12		3852.53		3829.62	
Planck	2777.98		2776.15		2783.61	
Pantheon	1026.24		1026.09		1025.94	
BAO	9.08		8.60		6.16	
SH0ES	10.60		11.038		4.43	
Lensing	10.55		11.19		9.46	
Planck SZ	18.07		19.45		0.017	

Best-fit  $H_0 = 70.62^{+0.72}_{-0.9}$  km/s/Mpc is in  $\sim 2\sigma$  tension with SH0ES measurement of  $H_0 = 73.04 \pm 1.04$  km/s/Mpc.

$$(\chi_{\text{min}}^{\Lambda\text{CDM}})^2 - (\chi_{\text{min}}^{\text{aDM}})^2 = 23.5$$

# Conclusion

- Atomic dark matter is a well-motivated, plausible interacting dark sector scenario.
- It can leave distinctive imprints on cosmological observables.
- By modifying CLASS to compute ADM cosmological evolution in generality, we place new constraints on the atomic dark matter parameter space from precision observations.
- There is no preference for atomic dark matter from CMB observations alone, but large aDM fractions can be allowed from large-scale measurements.
- When including low-redshift measurements of the Hubble constant and structure, a significant fraction of aDM that recombines near  $z = 2 \times 10^3$  is preferred, along with a non-zero amount of dark radiation. The Hubble and  $S_8$  tensions can be eased but not eliminated.





# Future directions

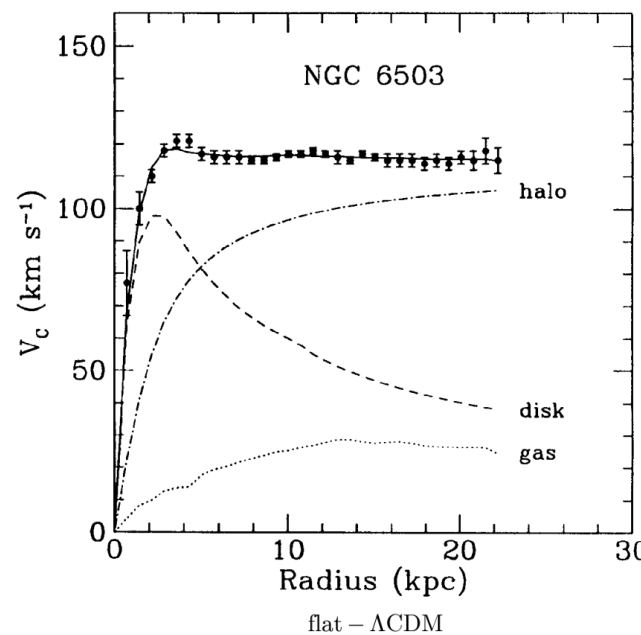
- Adding dark helium to CLASS-aDM, with arbitrary  $Y_{He_D}$ .
- Reliable computation of non-linear matter power spectrum evolution for arbitrary aDM parameters would unlock many constraints at smaller scales. (Halofit for aDM?)
  - Cosmic shear for  $k > 0.1 h \text{ Mpc}^{-1}$ .
  - Lyman- $\alpha$  forest.
  - 21 cm?
- Suggestions?

# Backup Slides

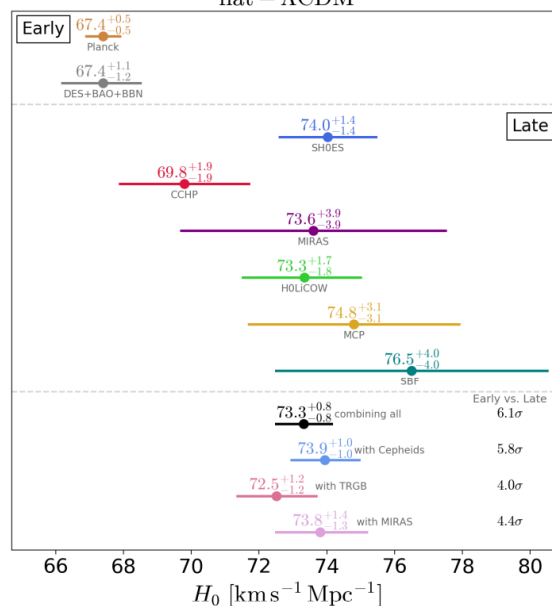


# Dark Matter: CDM or something more?

- The cold, collisionless dark matter paradigm explains many observations extremely well.
  - Galaxy rotation curves.
  - The Bullet cluster.
  - The cosmic microwave background spectrum.
  - Large scale structure in the universe.
- But it's not perfect.
  - Core vs cusp problem.
  - Hubble tension.
  - $S_8$  tension.
- Could a component of dark matter have non-trivial interactions in the dark sector?



Freese, 2008

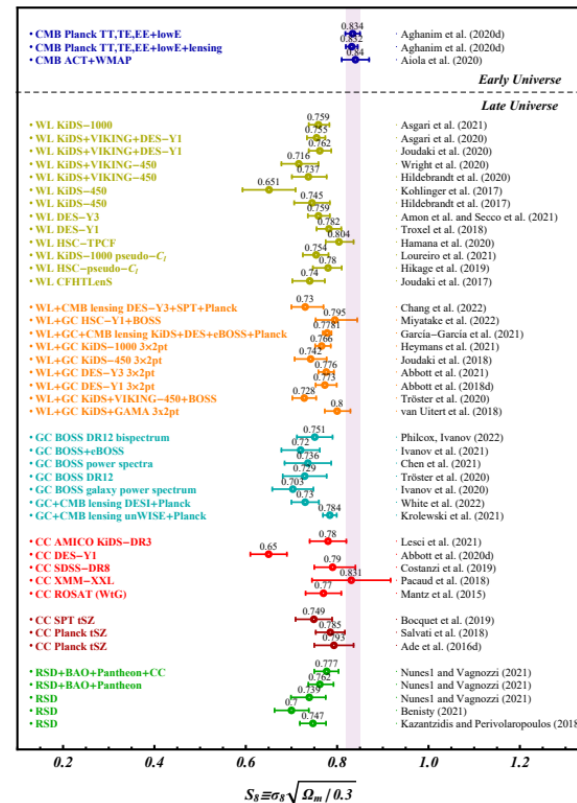
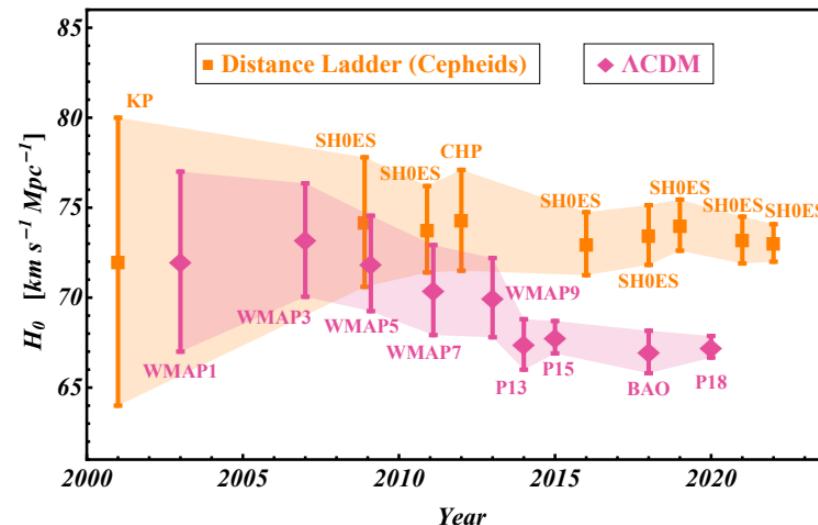


Verde, Treu, Riess, 2019



# Hubble and $S_8$ tensions

- $H_0$  measures the rate of expansion of the universe today.
  - Planck 2018 (CMB):  $H_0 = 67.4 \pm 0.5$  km/s/Mpc
  - SH0ES (distance ladder):  $H_0 = 73.04 \pm 1.04$  km/s/Mpc
- $S_8$  measures the size of fluctuations in the matter density at the scale  $8 h^{-1}$  Mpc.
  - Planck 2018 (CMB):  $S_8 = 0.834 \pm 0.016$
  - KiDS-1000 (weak lensing):  $S_8 = 0.754^{+0.027}_{-0.029}$



Images from 2105.05208

# Dark Acoustic Oscillations

