

Constraining sleptons at the LHC in a supersymmetric low-scale seesaw scenario

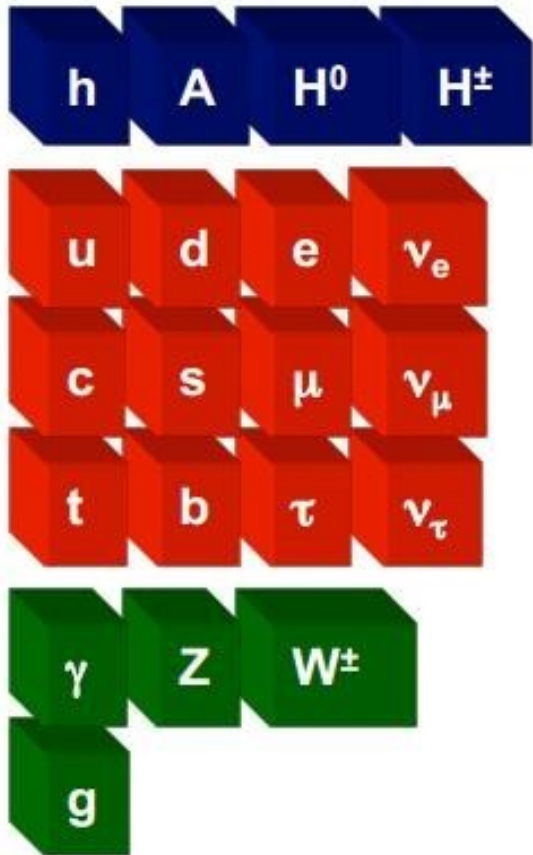
Nhell Cerna V.

HEP - PUCP

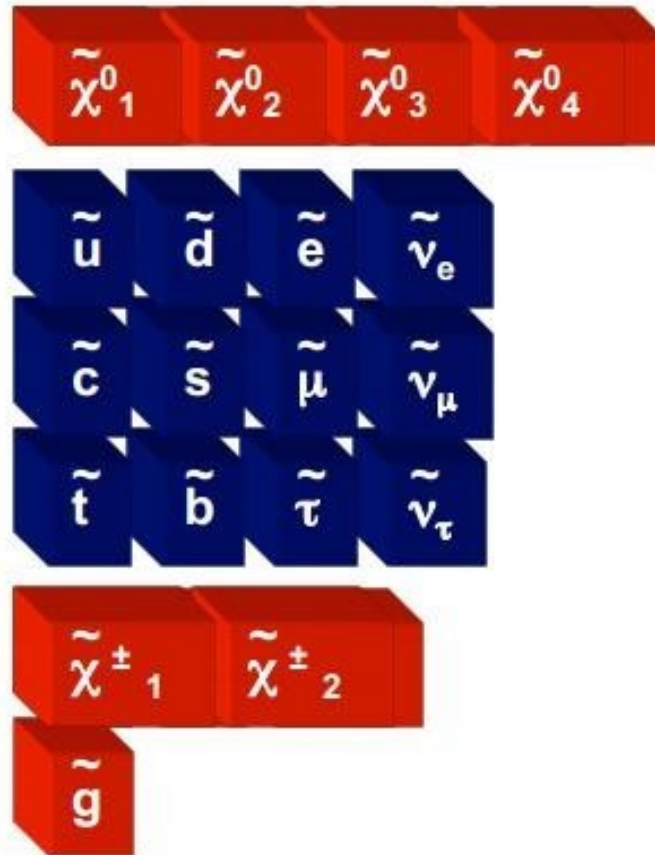
Pontificia Universidad Católica del Perú

November 26, 2018

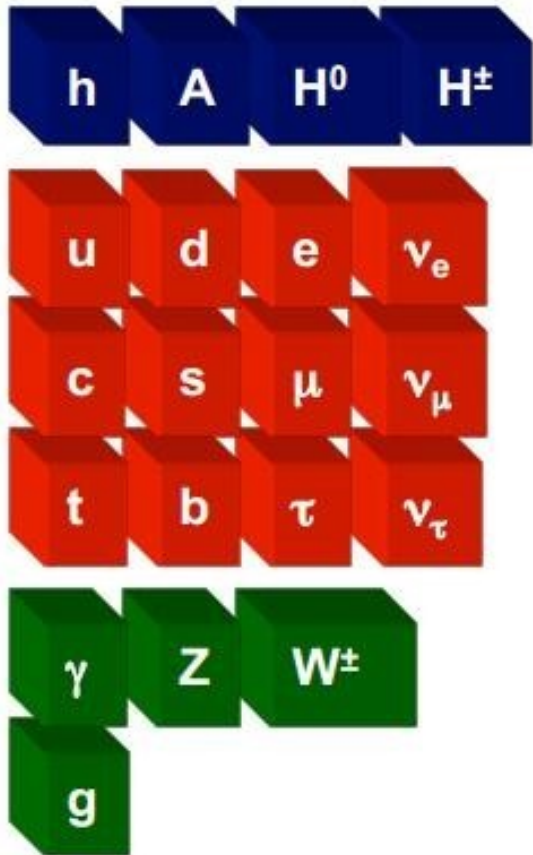
Standard particles



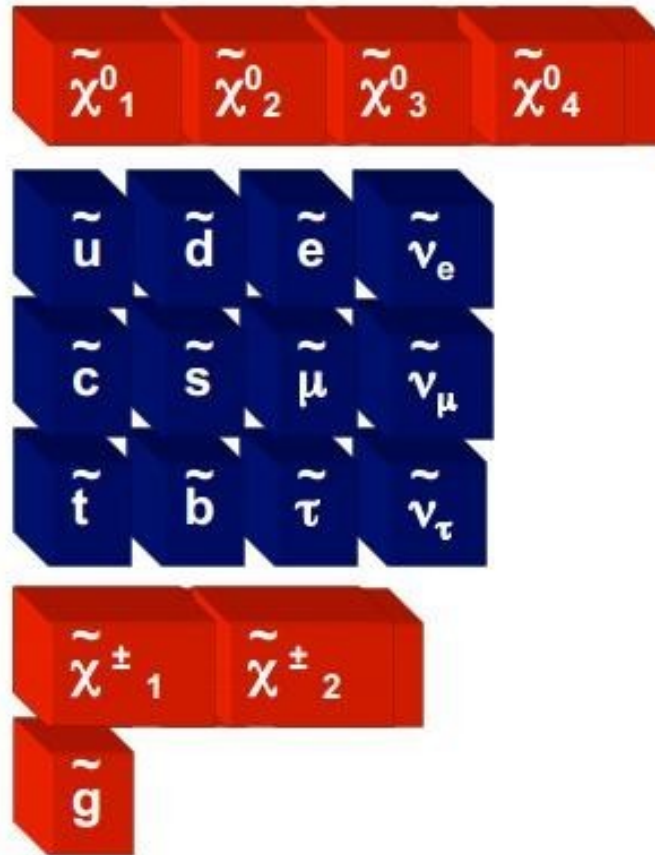
Supersymmetry particles



Standard particles

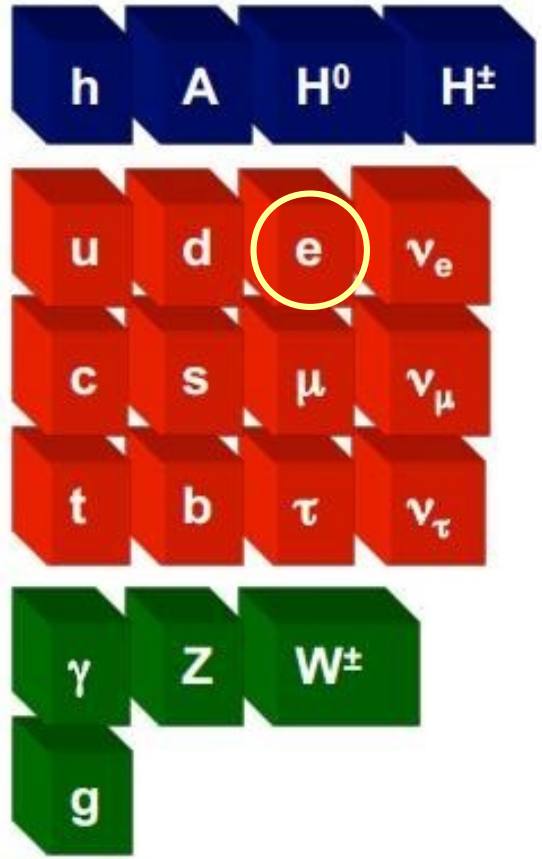


Supersymmetry particles

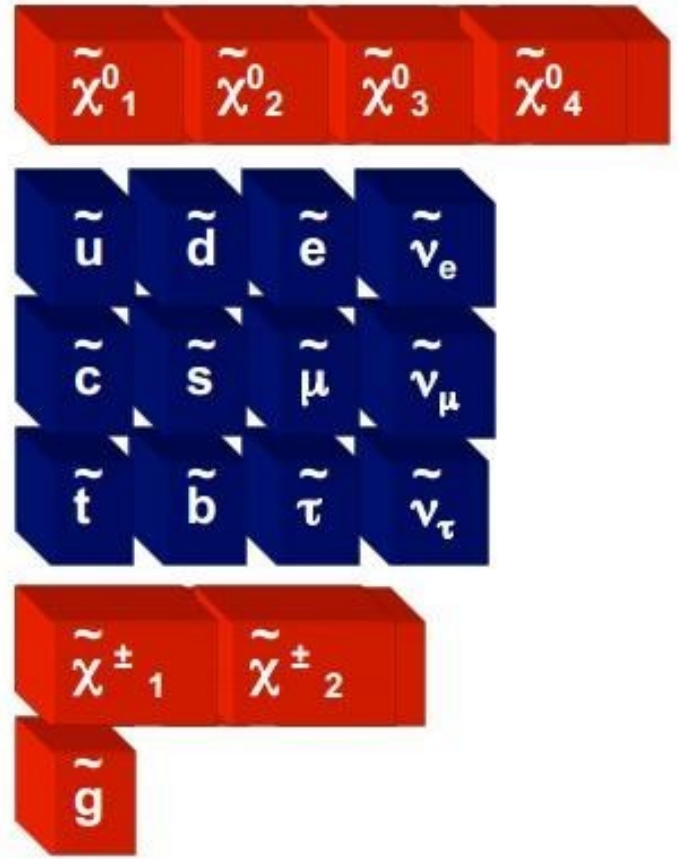


- Solves quadratic divergences
- Predicts new particles (DM candidate).

Standard particles

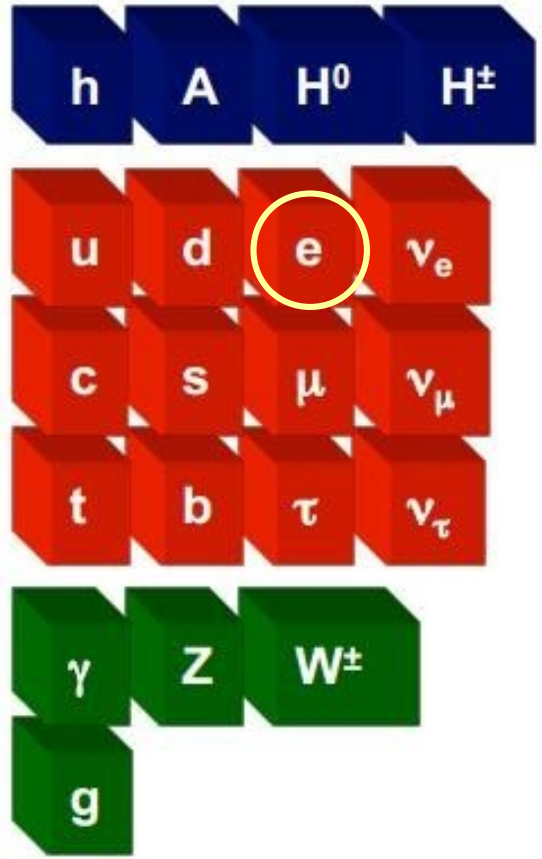


Supersymmetry particles

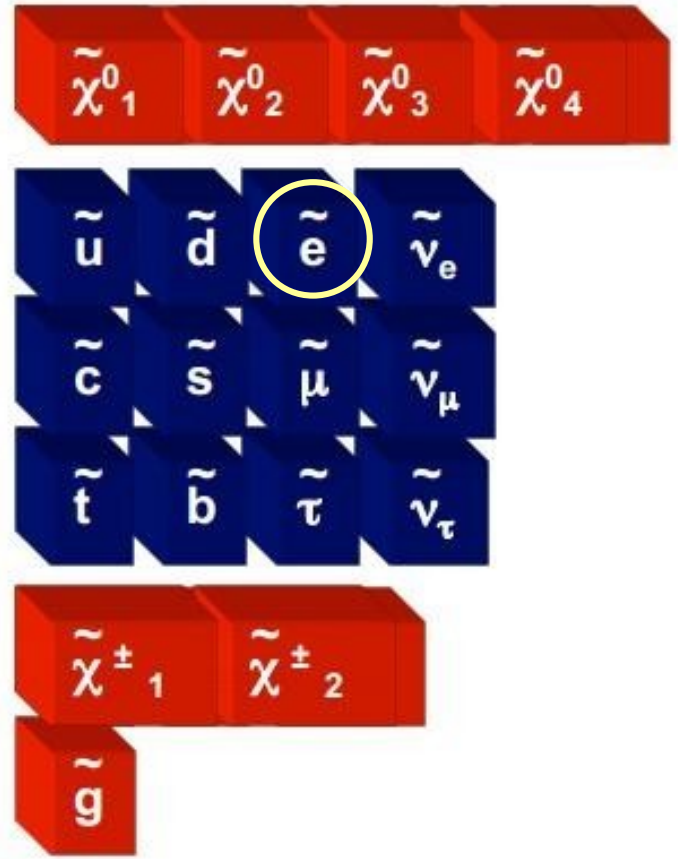


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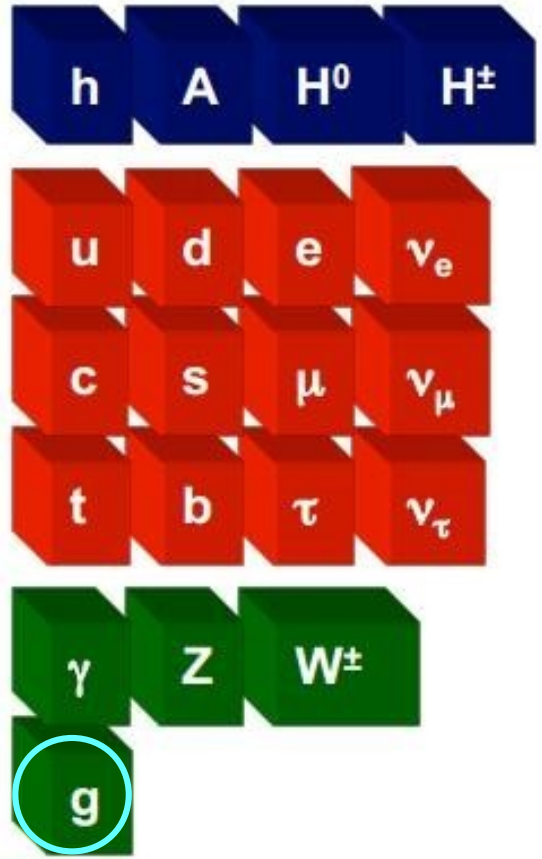


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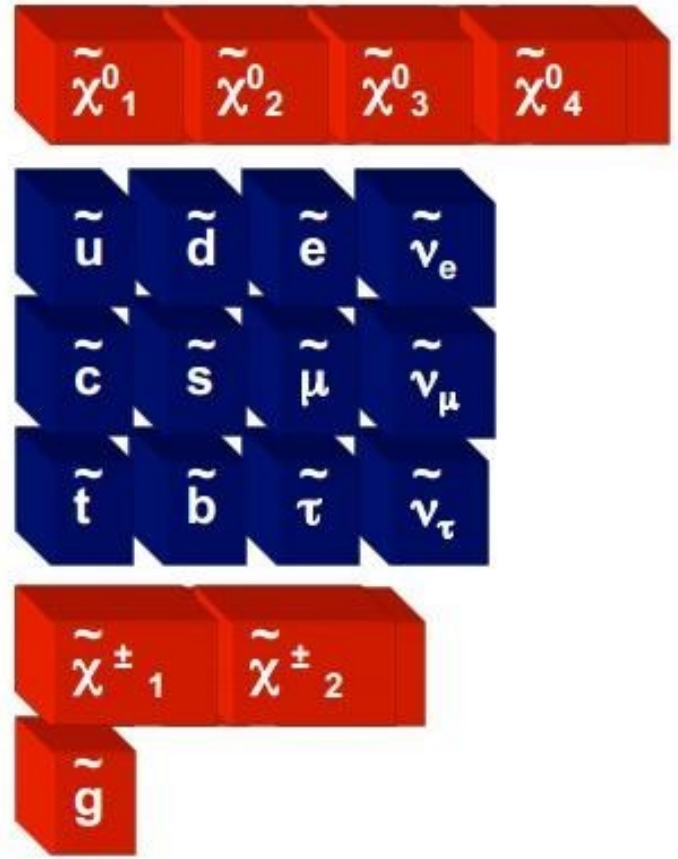


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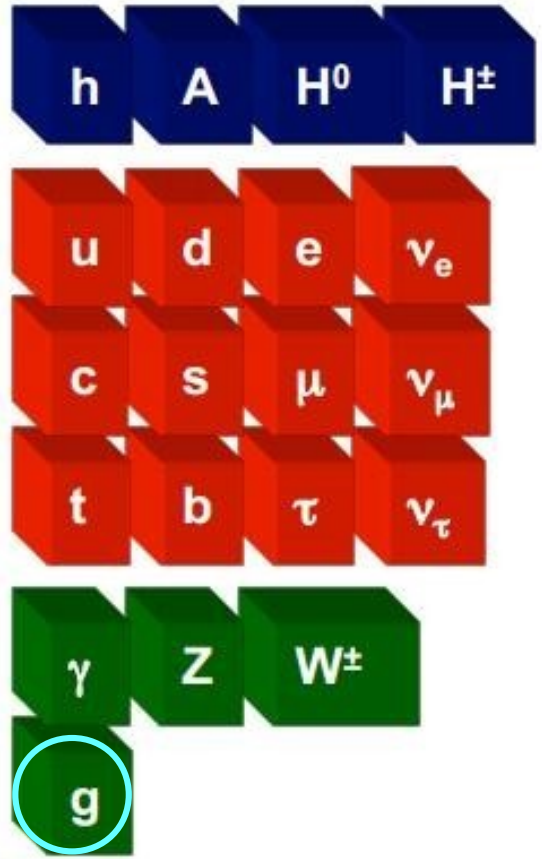


Supersymmetry particles

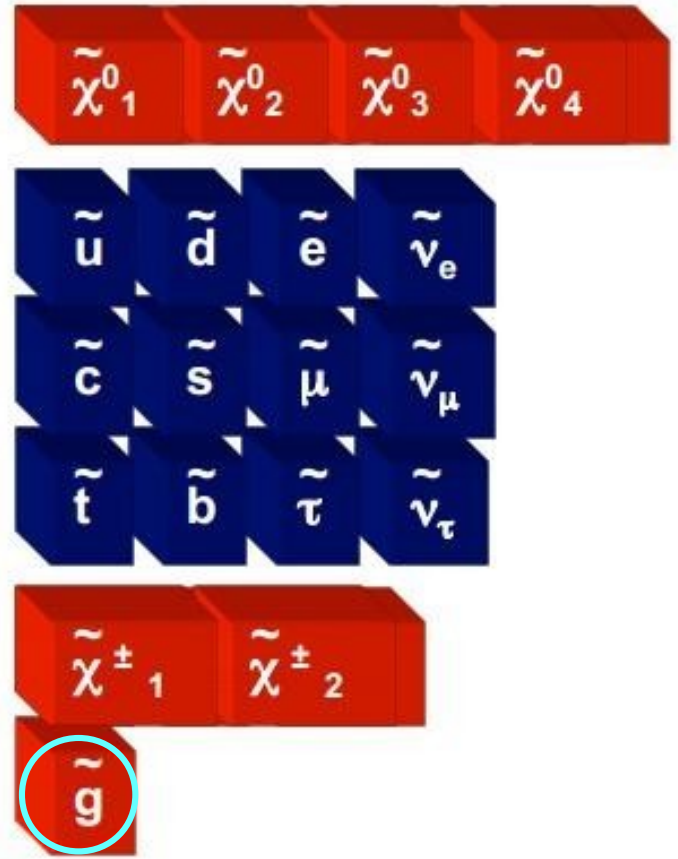


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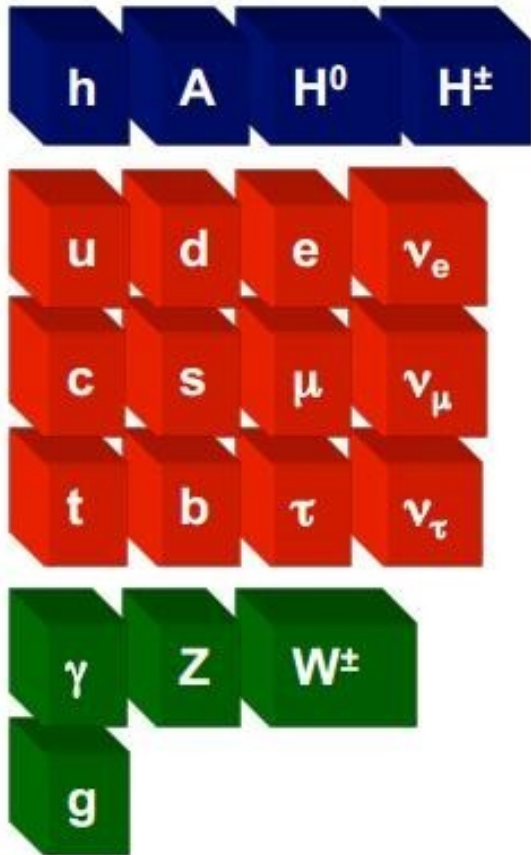


Supersymmetry particles

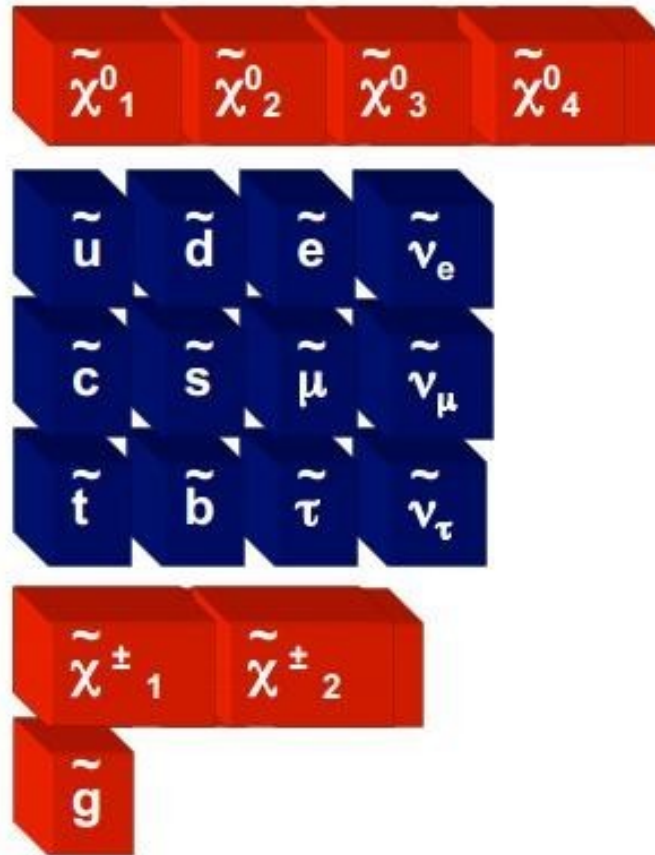


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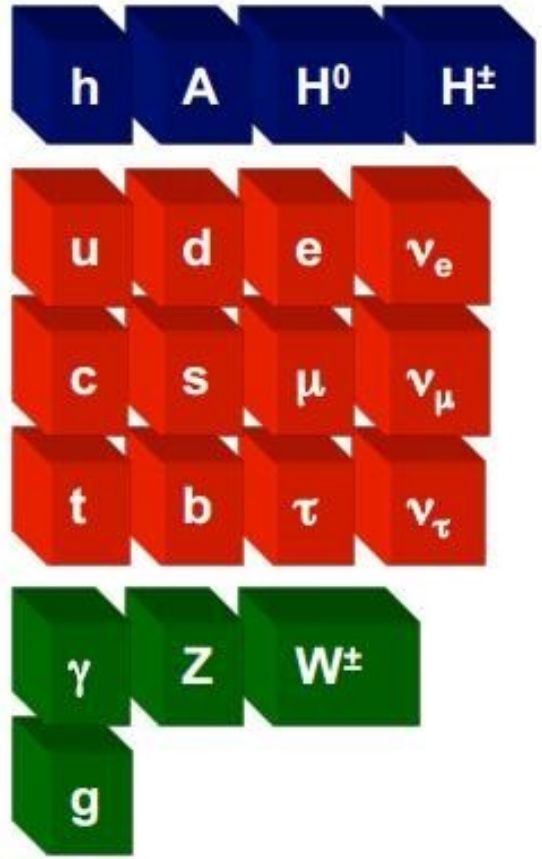
- Solves quadratic divergences
- Predicts new particles (DM candidate).

The model predicts:

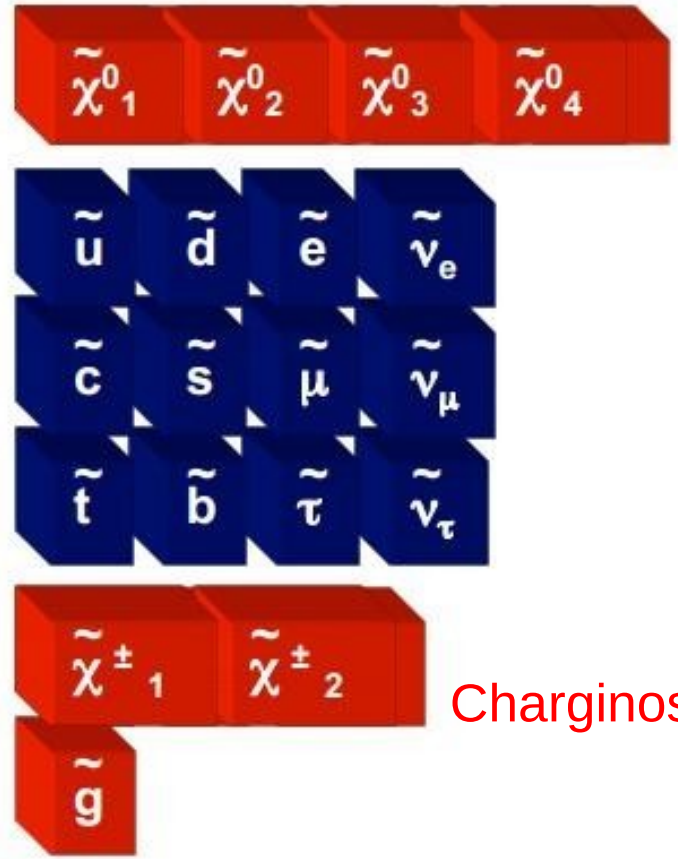
Spín $\frac{1}{2}$ \leftrightarrow Spín 0

Spín 1 or 0 \leftrightarrow Spín $\frac{1}{2}$

Standard particles



Supersymmetry particles



Neutralinos

- Solves quadratic divergences
- Predicts new particles (DM candidate).

The model predicts:

Spín $\frac{1}{2}$ \leftrightarrow Spín 0

Spín 1 or 0 \leftrightarrow Spín $\frac{1}{2}$

Charginos

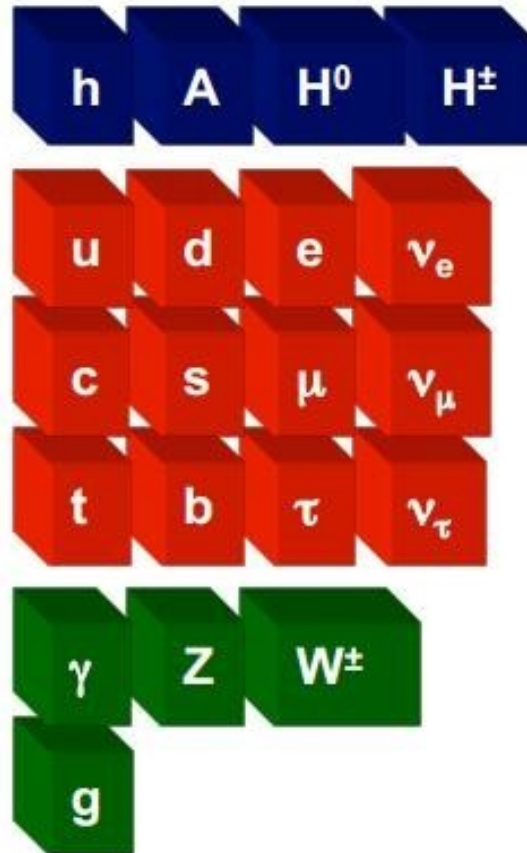
Seesaw mechanism.

This mechanism allows us to give mass to the neutrinos of the standard model, from the addition of new neutrinos (right handed neutrinos).

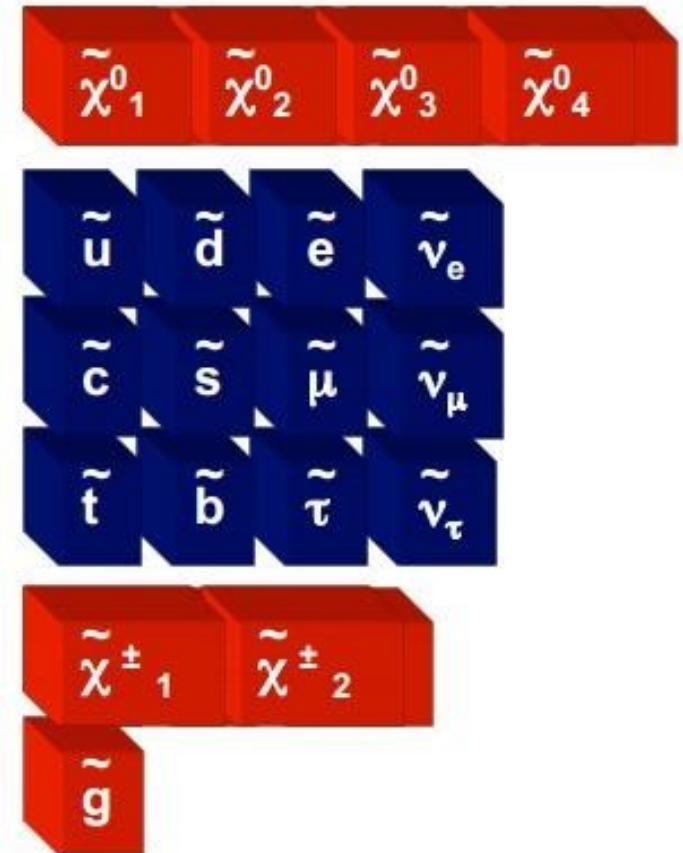
The addition of this mechanism to the MSSM implies:

$$\begin{array}{ccc} \nu_R & \Leftrightarrow & \tilde{\nu}_R \\ \text{R-Neutrino} & & \text{R-sneutrino} \end{array}$$

Standard particles



Supersymmetry particles

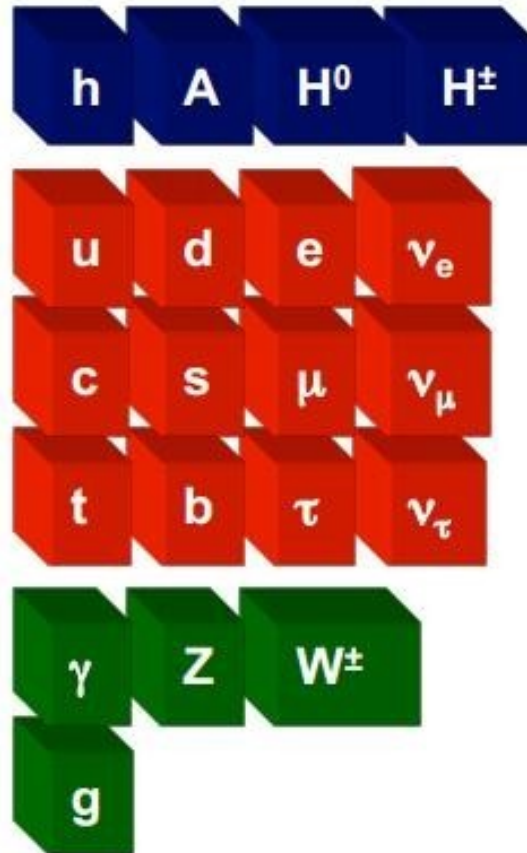


ν_R
Spín $\frac{1}{2}$

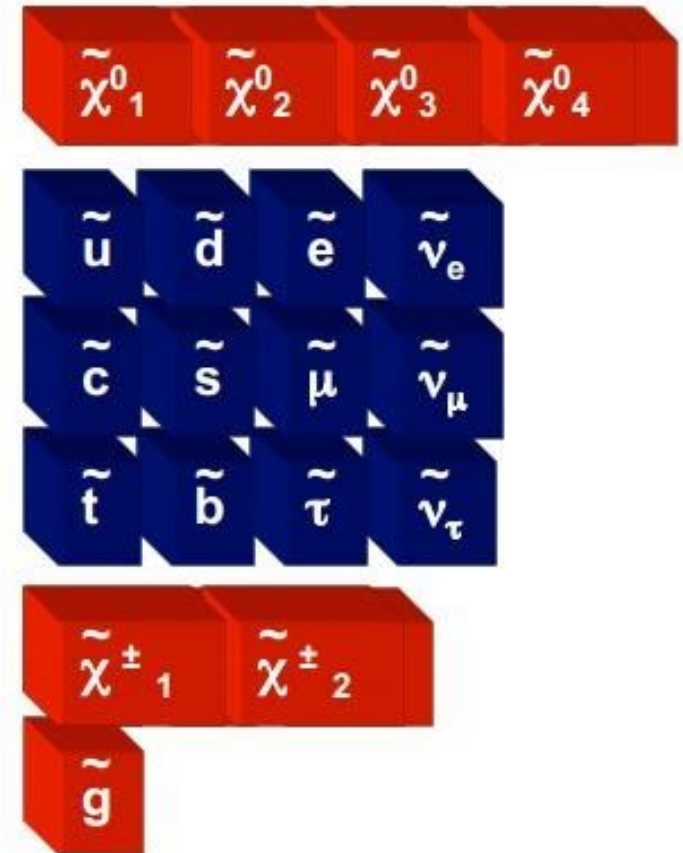


$\tilde{\nu}_R$
Spín 0

Standard particles



Supersymmetry particles



*The work is concentrated in the sleptons sector

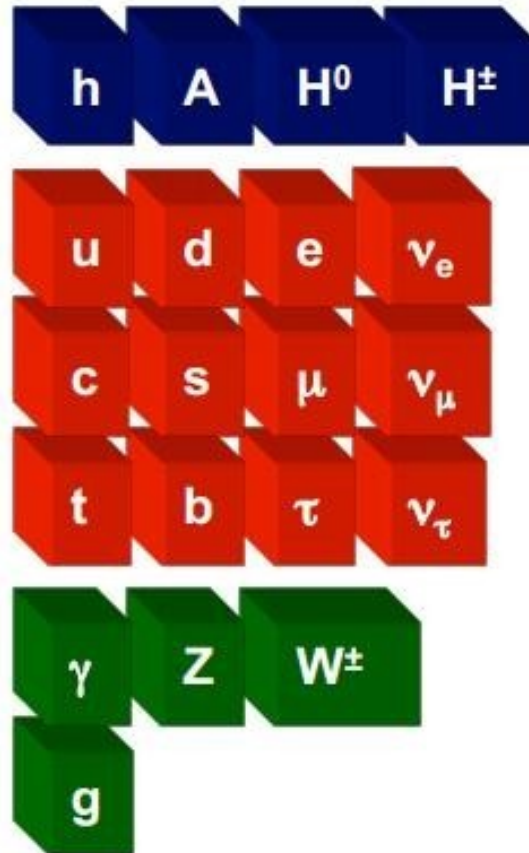
$$v_R \Leftrightarrow \tilde{\nu}_R$$

Spín $\frac{1}{2}$

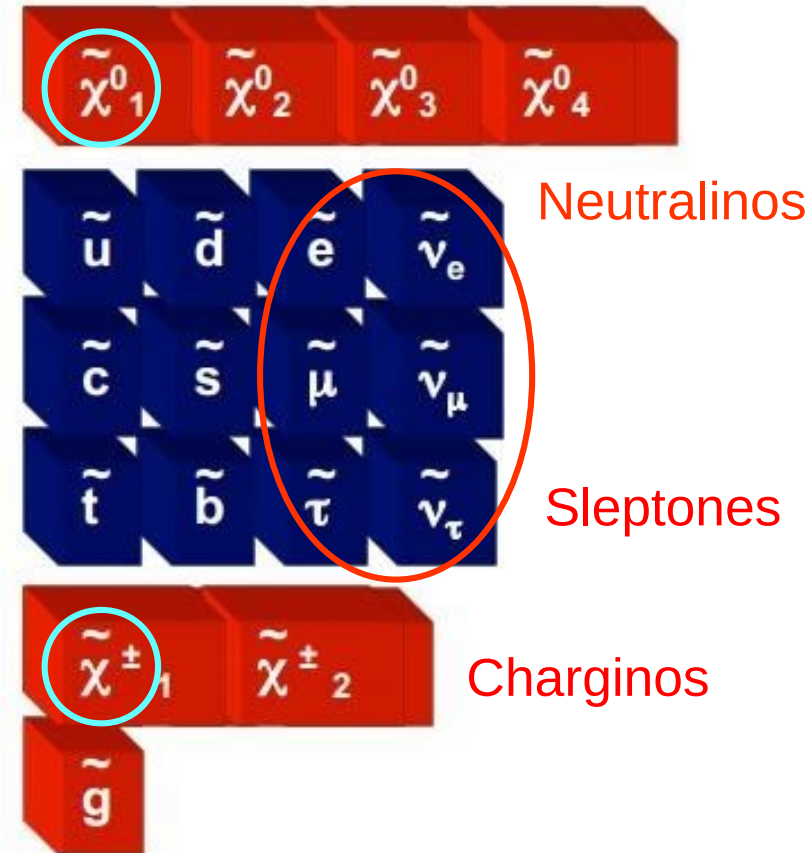
$$\tilde{\nu}_R$$

Spín 0

Standard particles



Supersymmetry particles



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ν_R
Spín $\frac{1}{2}$



$\tilde{\nu}_R$
Spín 0

R-Sneutrinos

By naturalness arguments, the μ term was kept relatively small, such that the lightest neutralinos were higgsino-like ($\tilde{\chi}$).

This means:

$$\mu \ll M_{1,2} \Rightarrow m_{\tilde{\chi}} \approx m_{\tilde{\chi}_{1,2}^0} \approx m_{\tilde{\chi}_{1\pm}}$$

In our work, we choose:

$$\mu \leq 400 \text{ GeV}$$

Usually in the mssm, the lightest particle is the neutralino,

$$\tilde{\chi}^0$$

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$$\tilde{\nu}_R$$

R-Sneutrino like LSP (Lightest and more stable particle)

In our model we consider the following cases:

$$m_{\tilde{\nu}_R} < m_{\tilde{\chi}} < m_{\tilde{l}}$$

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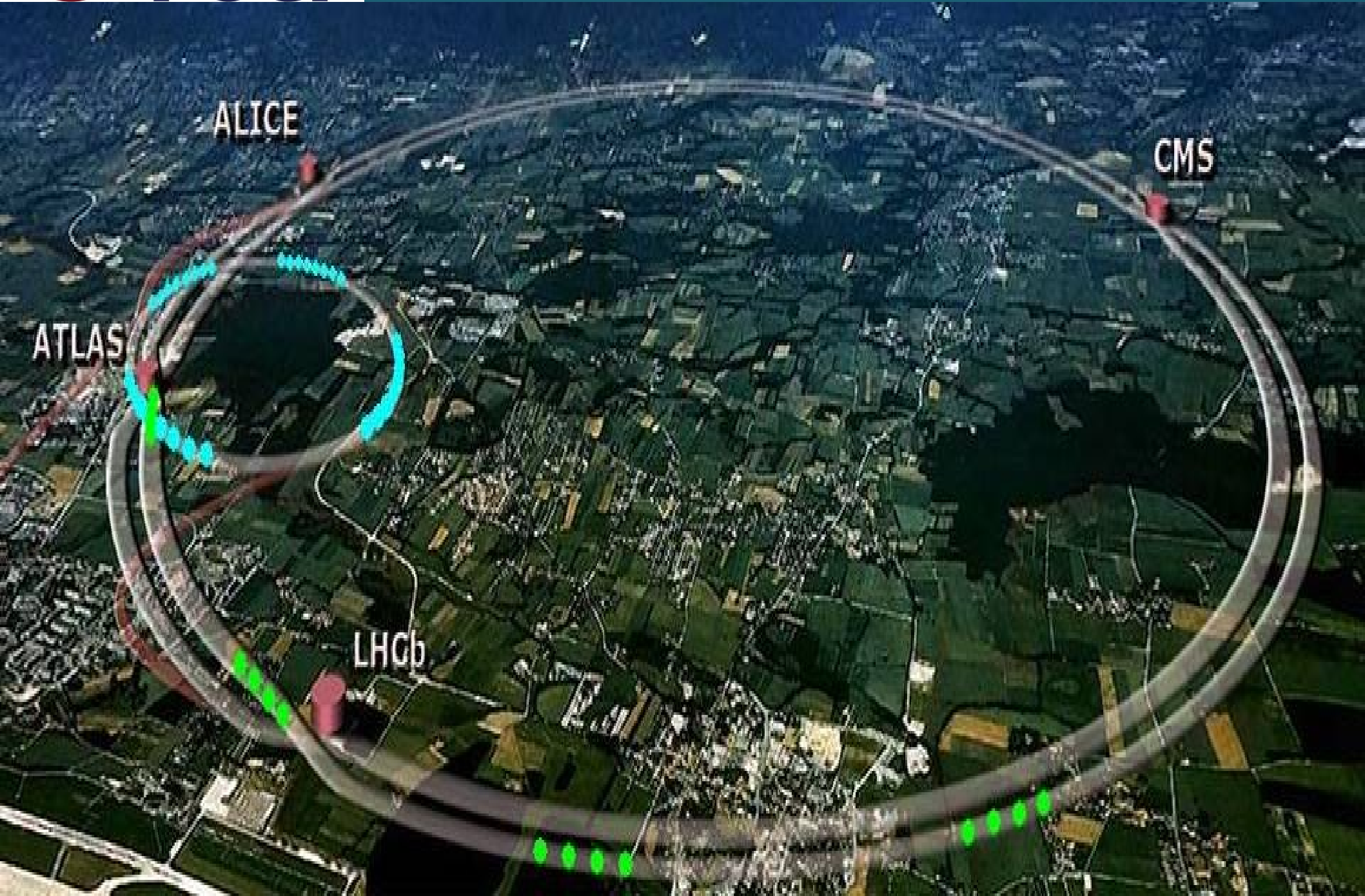
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$$m_{\tilde{\nu}_R} < m_{\tilde{\tau}} < m_{\tilde{\chi}} \quad \text{phenomenology with slepton NLSP}$$



The best way to study the properties of **R-sneutrinos** $\tilde{\nu}_R$, it is through decay in cascade of heavy particles, like as L-sleptons and higgsinos (neutralinos y charginos).

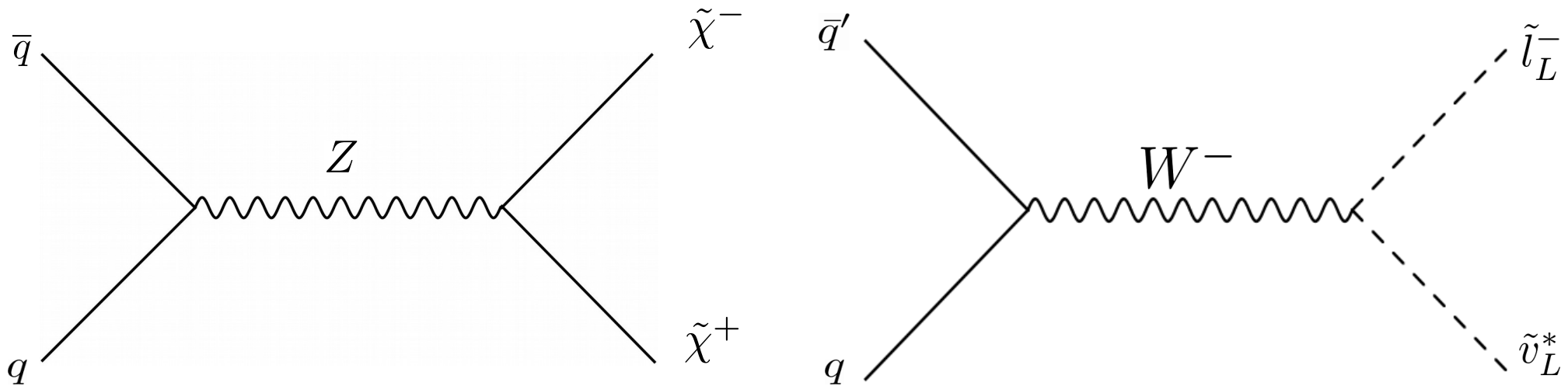
For that, we have to understand how these particles are produced.

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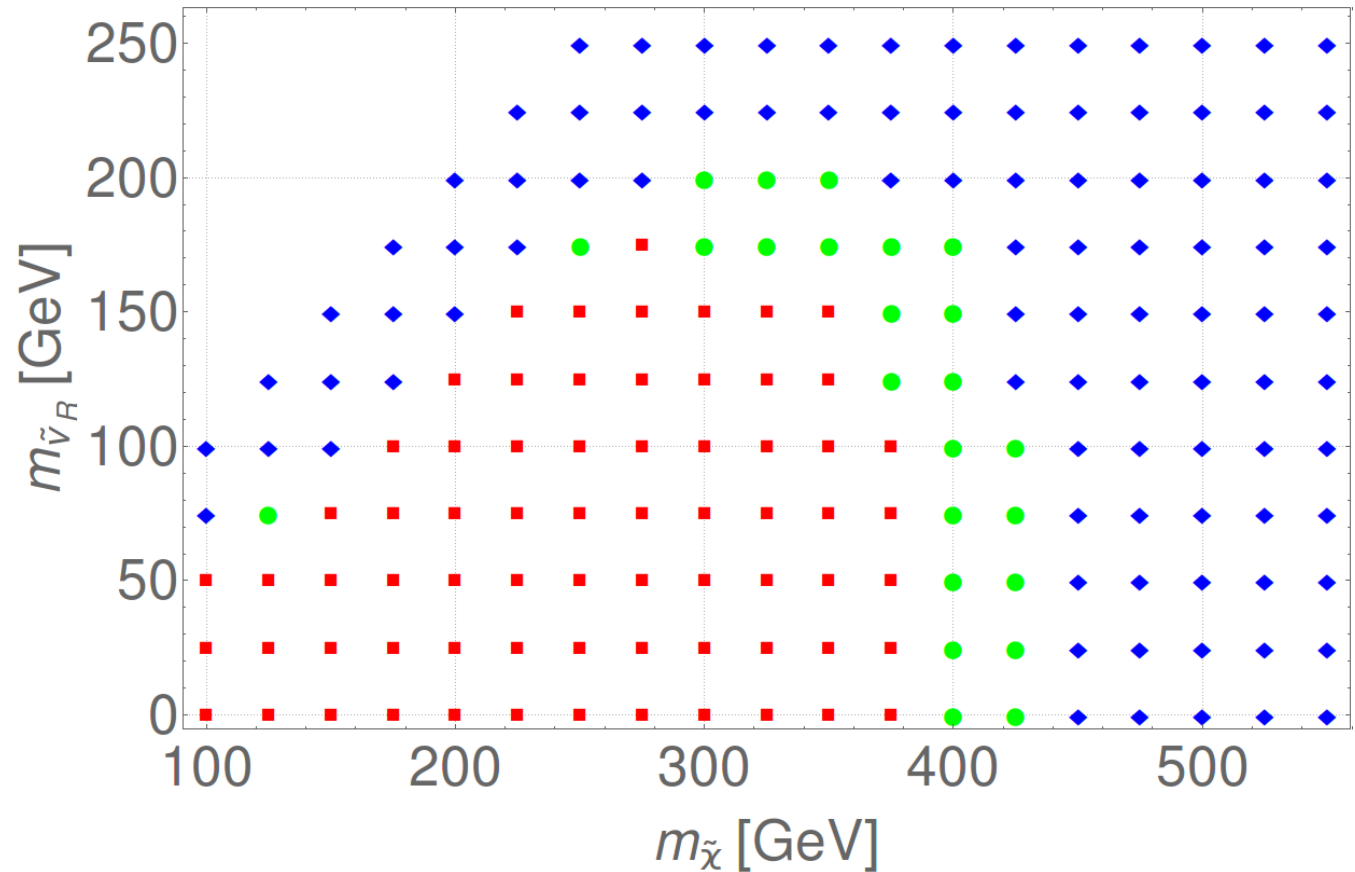
for example:



Electro-Weak Drell-Yan processes.

Signal: $pp \rightarrow \tilde{\chi}^+ \tilde{\chi}^-, \tilde{\chi}^\pm \rightarrow l^\pm \tilde{\nu}_R^{(*)}$

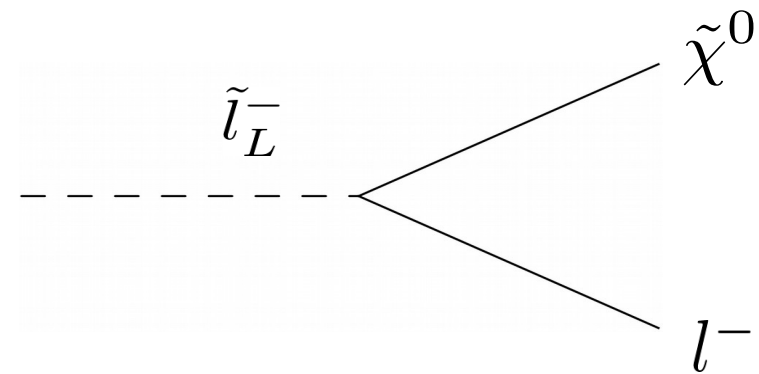
- Excluded
- Allowed
- Ambiguous



$$m_{\tilde{\nu}_R} < m_{\tilde{\chi}} < m_{\tilde{l}}$$

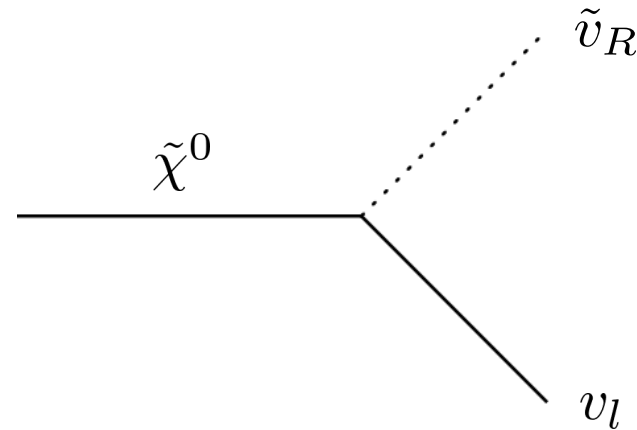
- Sleptons:

$$\tilde{l}_L^- \rightarrow \tilde{\chi}^0 l^-$$



- Neutralinos:

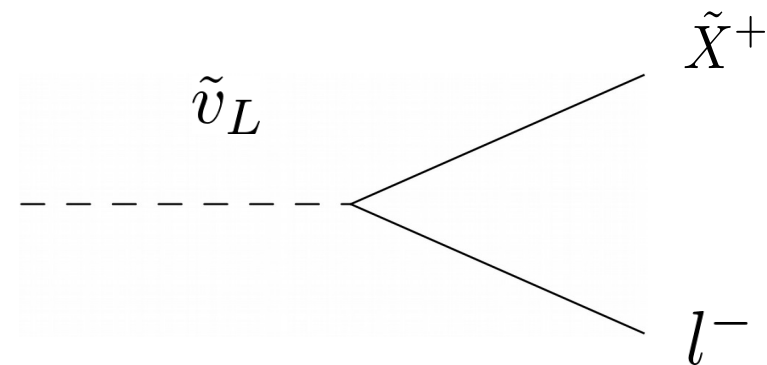
$$\tilde{\chi}^0 \rightarrow \tilde{\nu}_R \nu_l$$



$$m_{\tilde{\nu}_R} < m_{\tilde{\chi}} < m_{\tilde{l}}$$

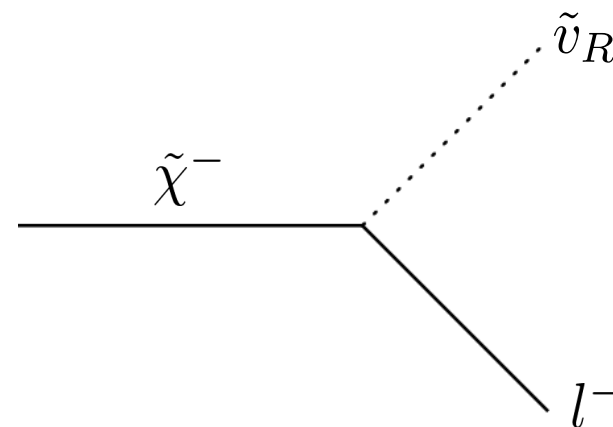
- L-Sneutrinos:

$$\tilde{\nu}_L \rightarrow \tilde{\chi}^+ l^-$$



- Charginos:

$$\tilde{\chi}^- \rightarrow \tilde{\nu}_R l^-$$



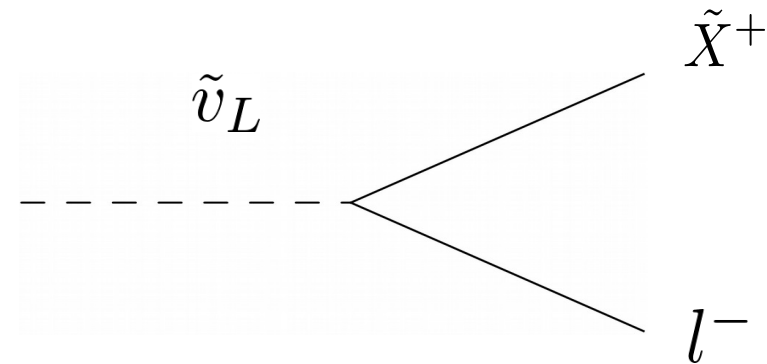
$$m_{\tilde{\nu}_R} < m_{\tilde{\chi}} < m_{\tilde{l}}$$

Signal

$$pp \rightarrow \tilde{l}_L^- \tilde{\nu}_L^* \rightarrow 3l + E_{Tmiss}$$

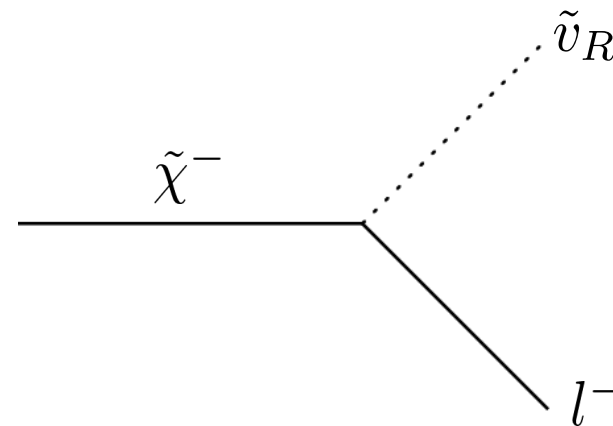
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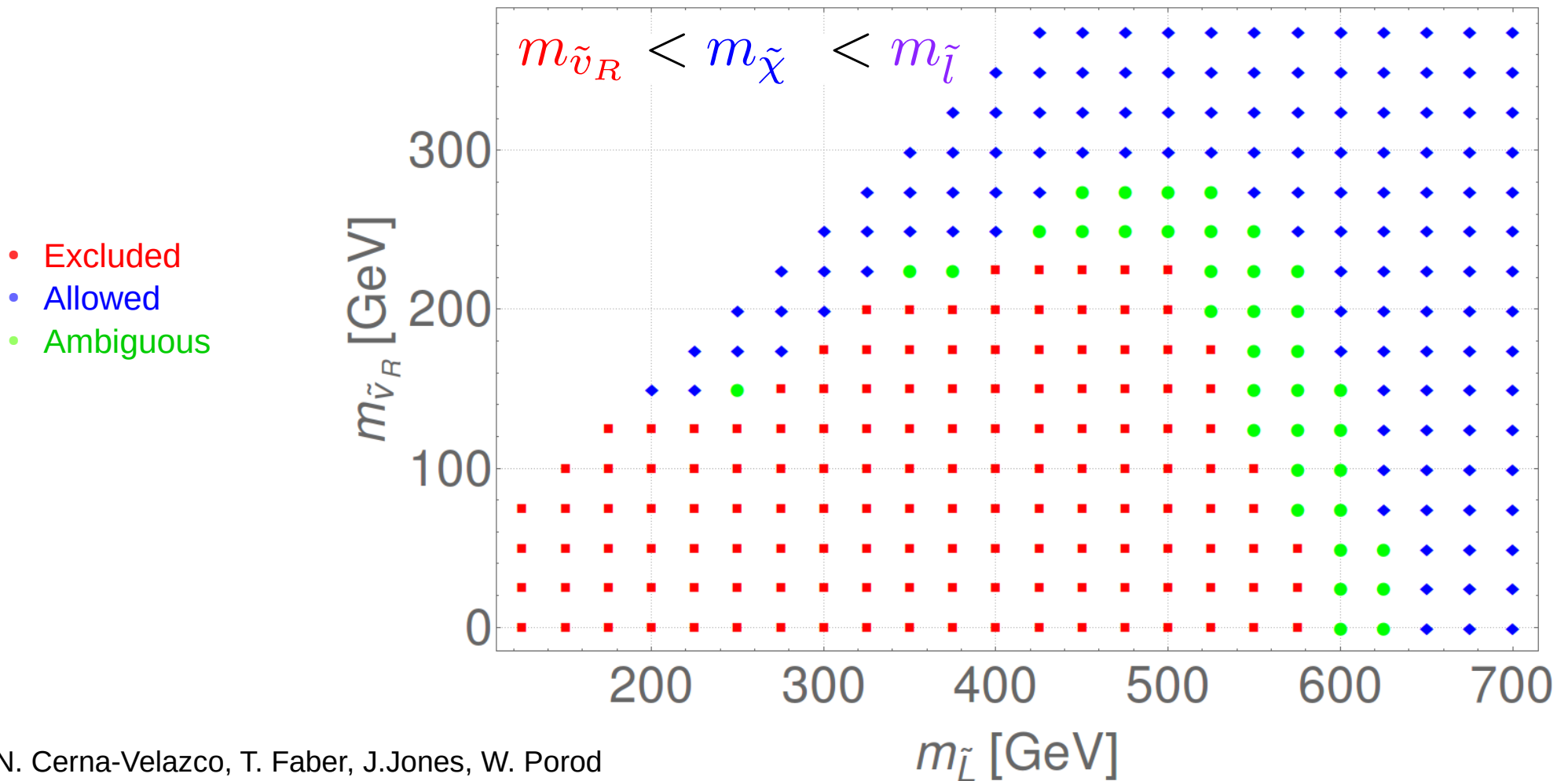
- Charginos:

$$\tilde{\chi}^- \rightarrow \tilde{\nu}_R l^-$$



$$pp \rightarrow \tilde{l}_L^- \tilde{\nu}_L^* \rightarrow 3l + E_{T\text{miss}}$$

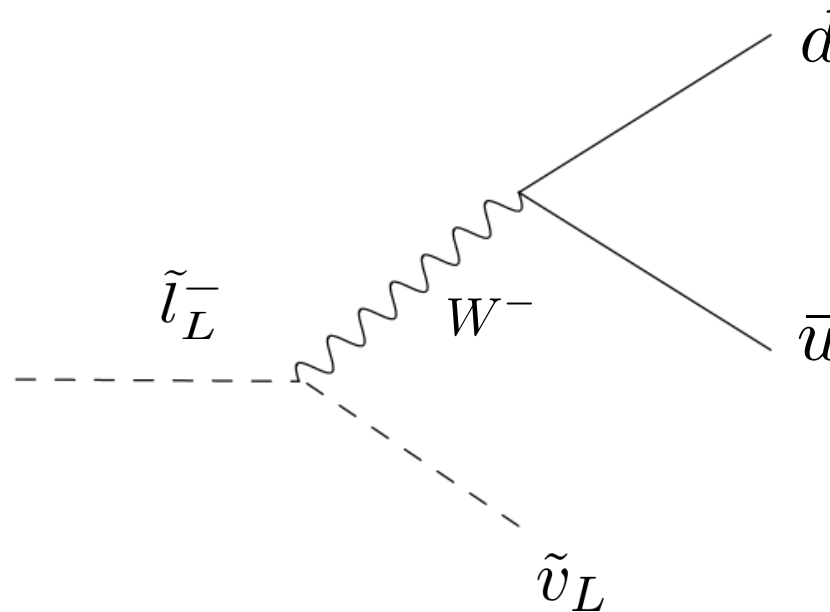
$$m_{\tilde{\chi}^0} = m_{\tilde{\nu}_R} + 25\text{GeV}$$



$$m_{\tilde{\nu}_R} < m_{\tilde{l}} < m_{\tilde{\chi}}$$

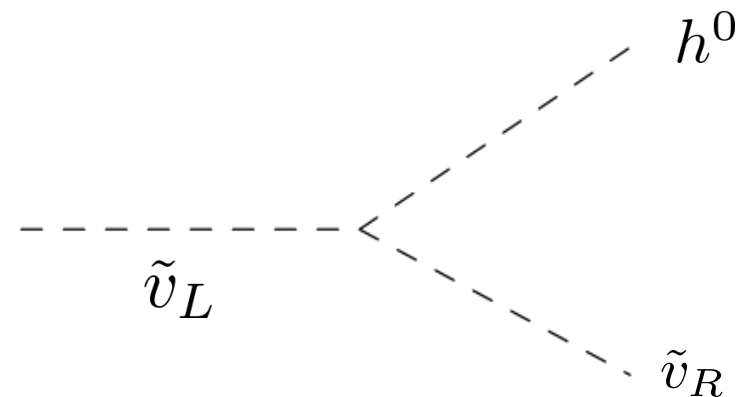
- Sleptons:

$$\tilde{l}_L^- \rightarrow 2j\tilde{\nu}_L$$



- L-Sneutrinos:

$$\tilde{\nu}_L \rightarrow \tilde{\nu}_R h^0$$



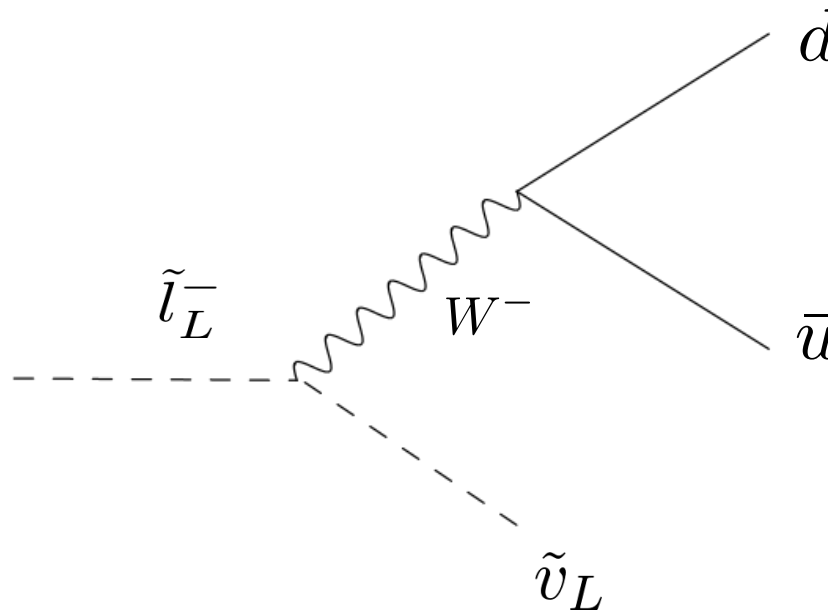
$$m_{\tilde{\nu}_R} < m_{\tilde{l}} < m_{\tilde{\chi}}$$

Signal

$$pp \rightarrow \tilde{l}_L^- \tilde{\nu}_L^* \rightarrow 2h^0 + j's + E_{Tmiss}$$

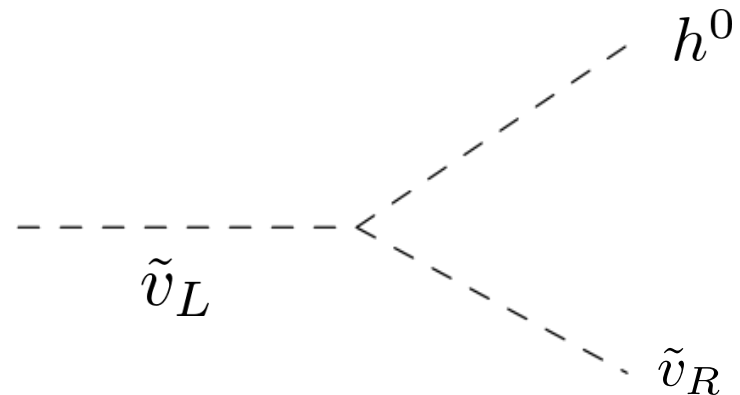
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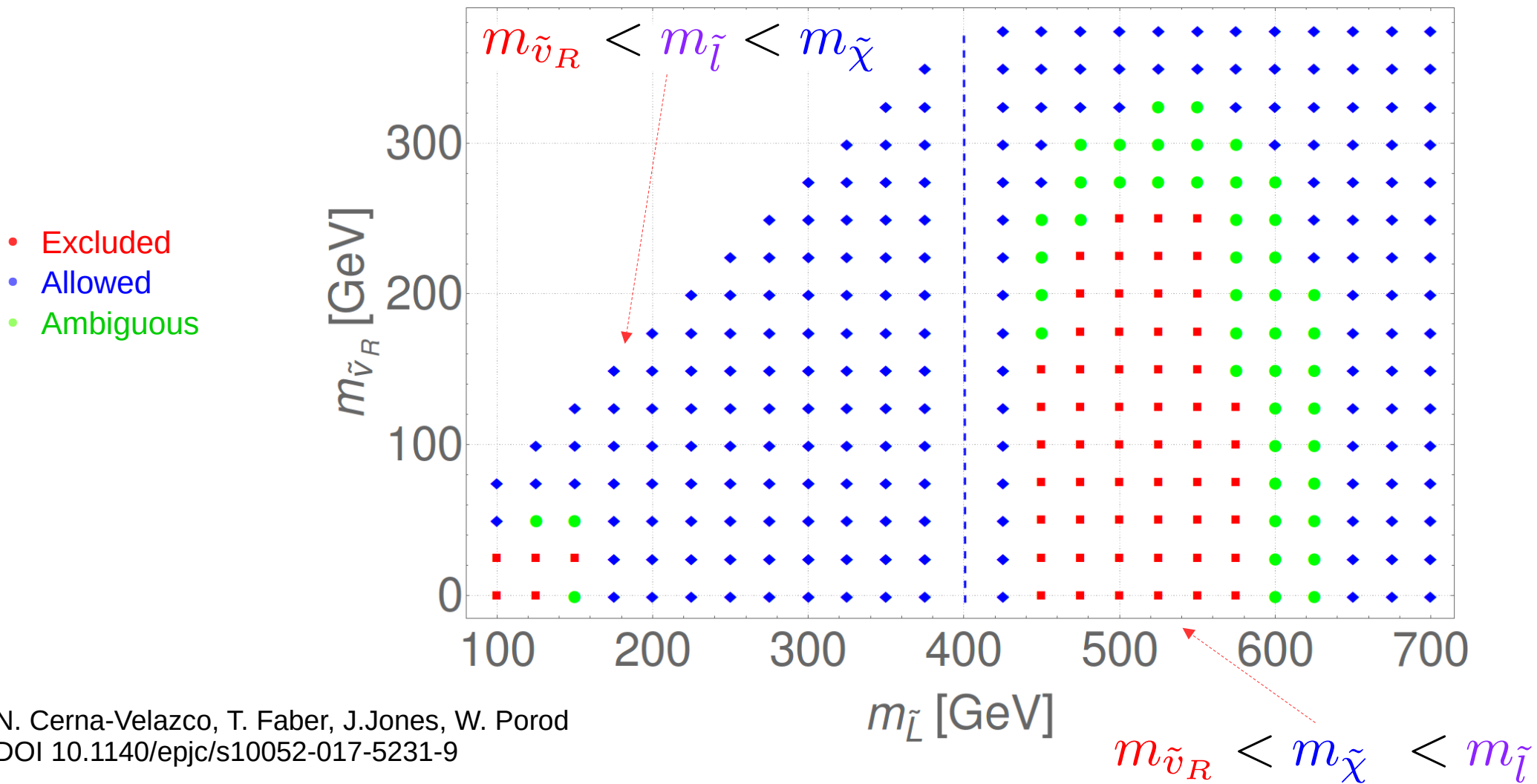
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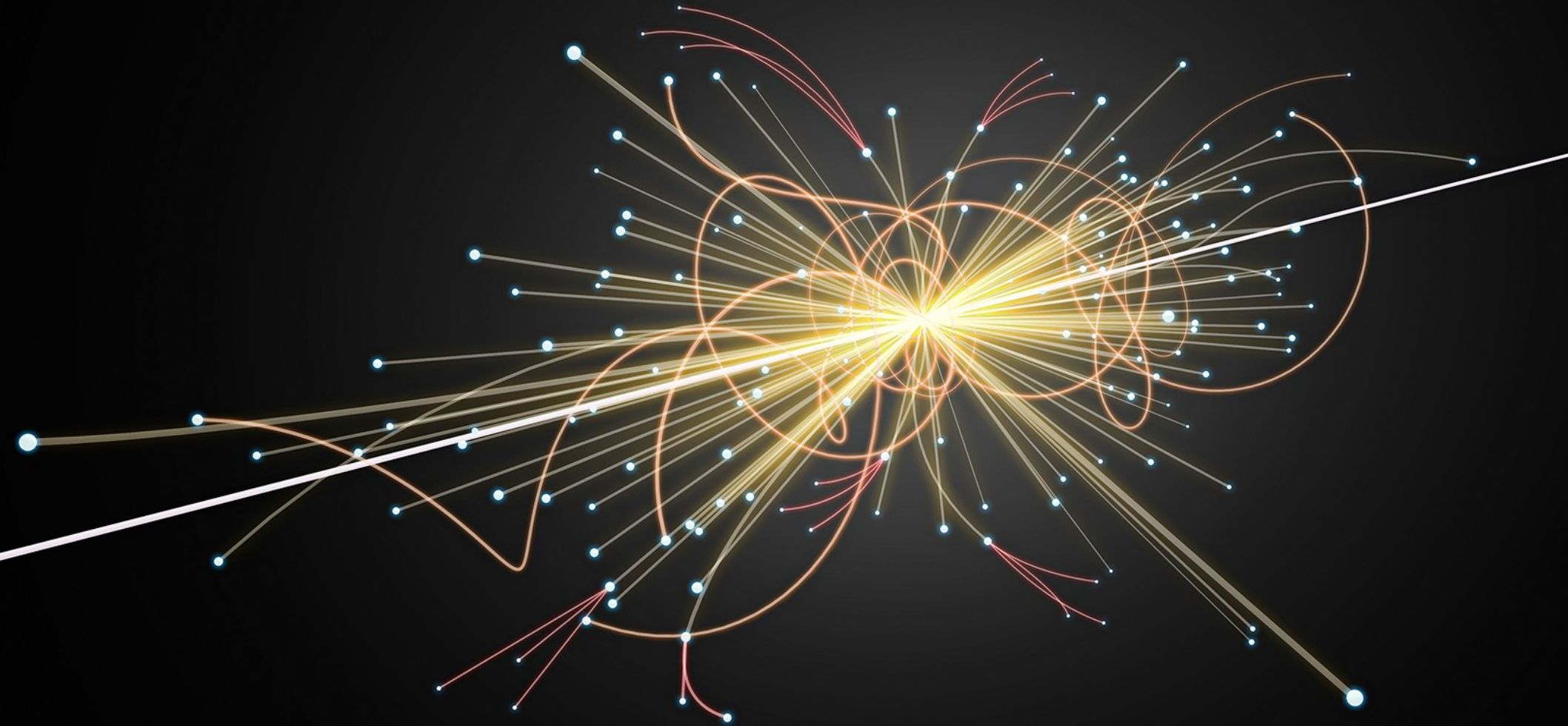
$$m_{\tilde{\chi}} = 400 \text{ GeV}$$



- We observed that phenomenology of the R-Sneutrino depends on the mass of sleptons and higgsinos.
- We see that, for vanishing $m_{\tilde{\nu}_R}$, the bound on $m_{\tilde{\chi}}$ can be as large as 375 GeV. For relatively small values of $m_{\tilde{\chi}}$, one finds that R-sneutrino masses lighter than $m_{\tilde{\chi}} - 75$ GeV are ruled out, with the allowed region increasing for $m_{\tilde{\chi}} \gtrsim 250$ GeV.
- We found that, as long as $m_{\tilde{\chi}} < m_{\tilde{l}}$, we can exclude slepton masses to a maximum of 575 GeV. For lower values of slepton mass, the R-sneutrino masses can be excluded up to about 200 GeV.
- In case $m_{\tilde{l}} < m_{\tilde{\chi}}$, constraints became very weak, as final states were either too soft, or not probed by current searches.



Thanks



We introduce new states (right-handed field)

$$\mathcal{L}_{seesaw} = \mathcal{L}_{SM} - m_D \bar{v}_R v_L + \frac{1}{2} m_R \bar{v}_R^c v_R + h.c.$$

Eventually we will get

$$\mathcal{L}_{mass}^v = \frac{1}{2} \bar{N}_L^c M N_L + h.c. \quad N_L = \begin{pmatrix} v_L \\ v_R^c \end{pmatrix}$$

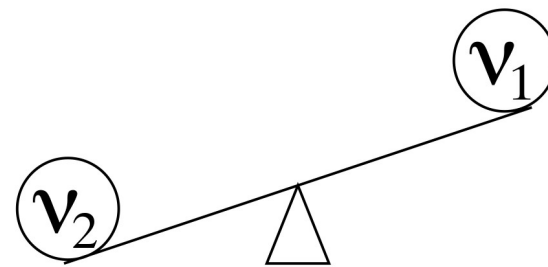
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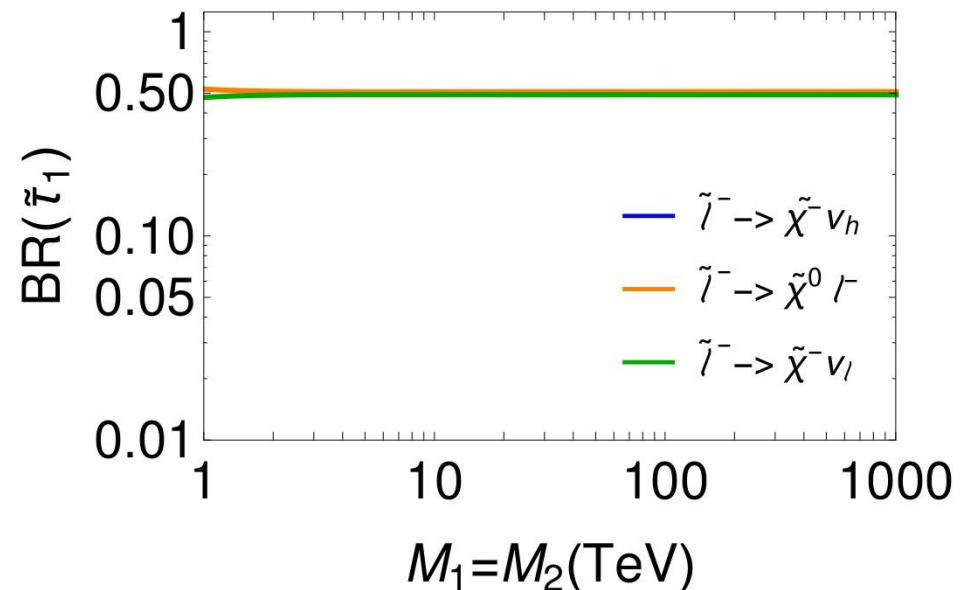
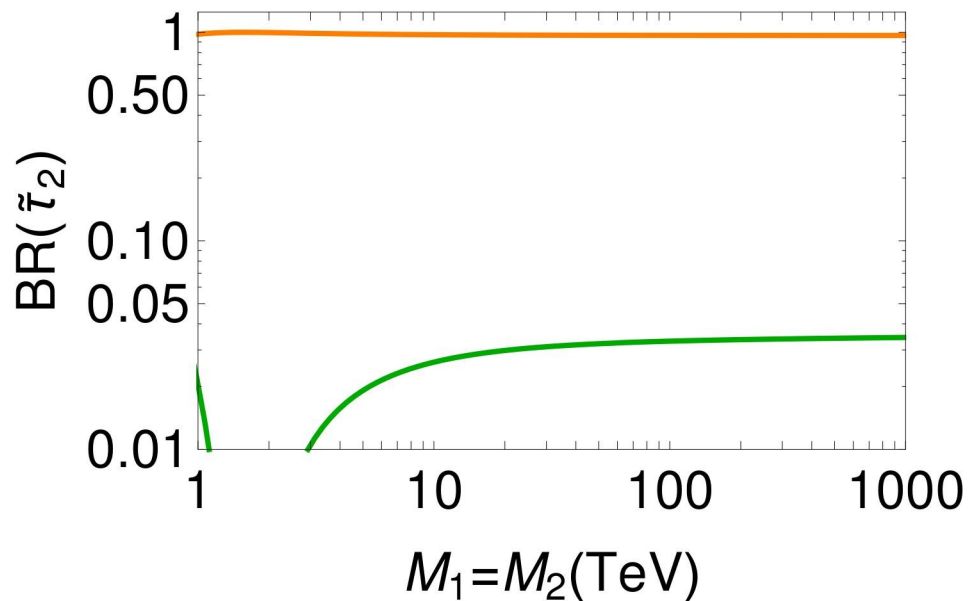
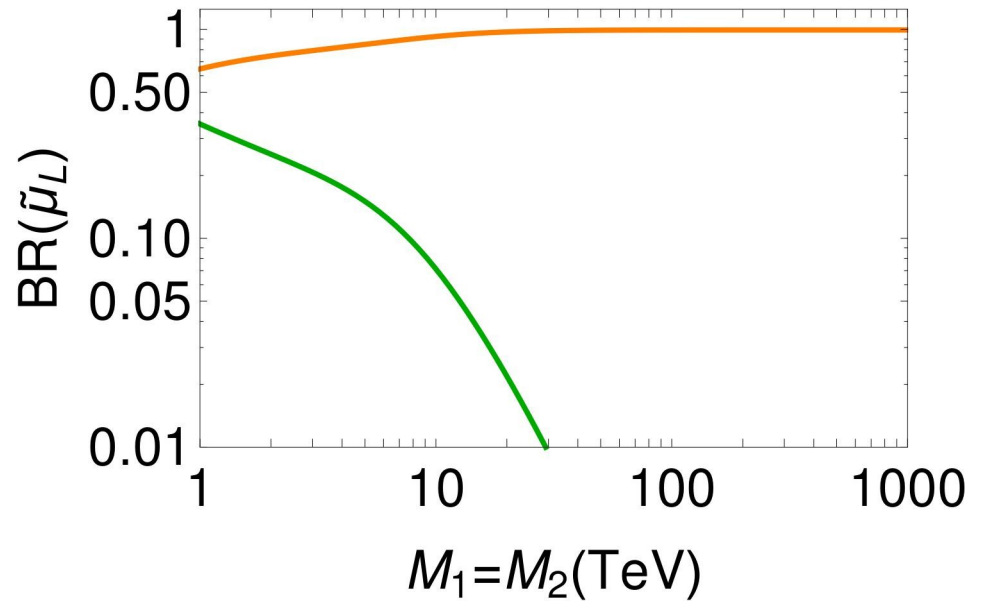
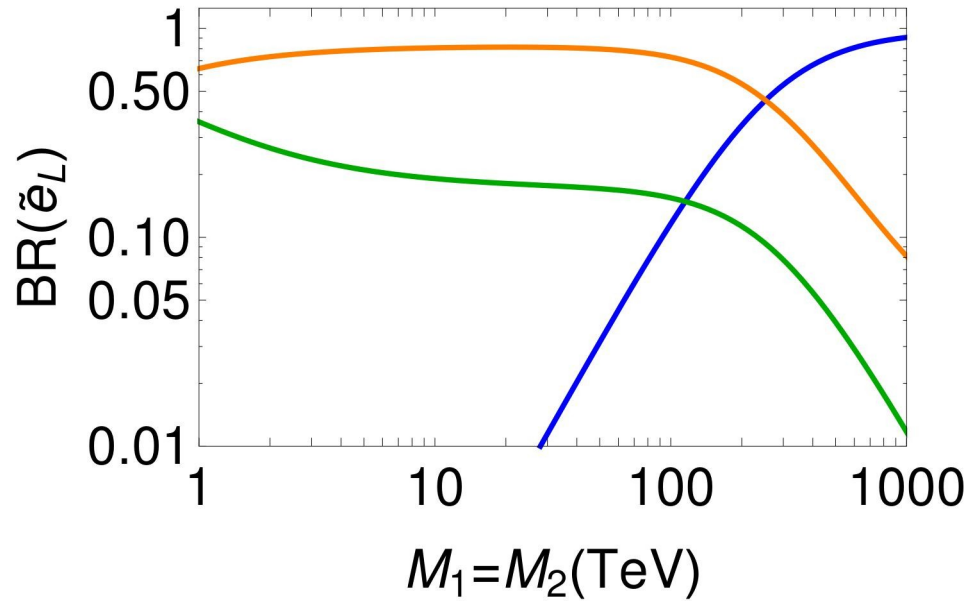
$$M = \begin{pmatrix} 0 & m_D \\ m_D & m_R \end{pmatrix}$$



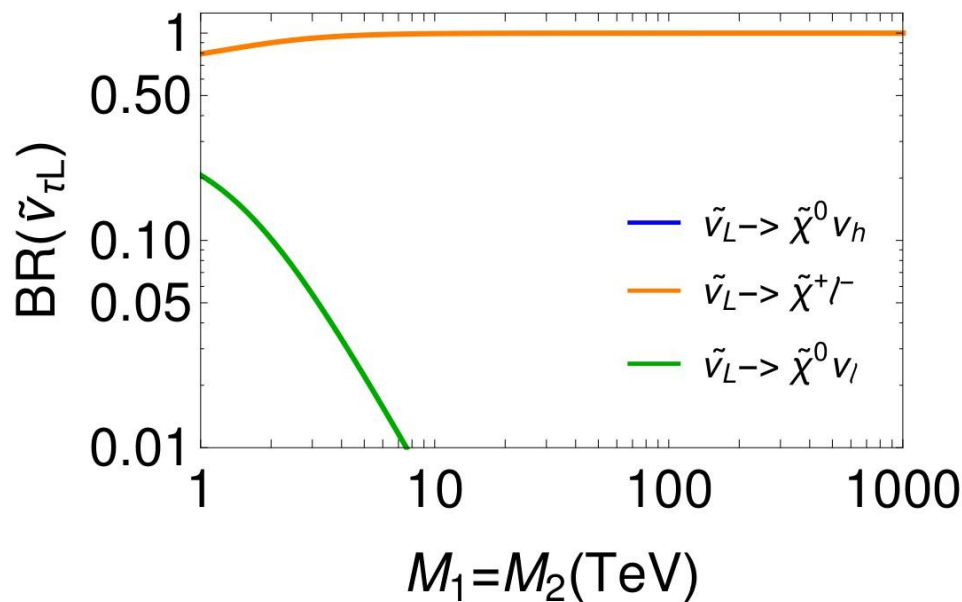
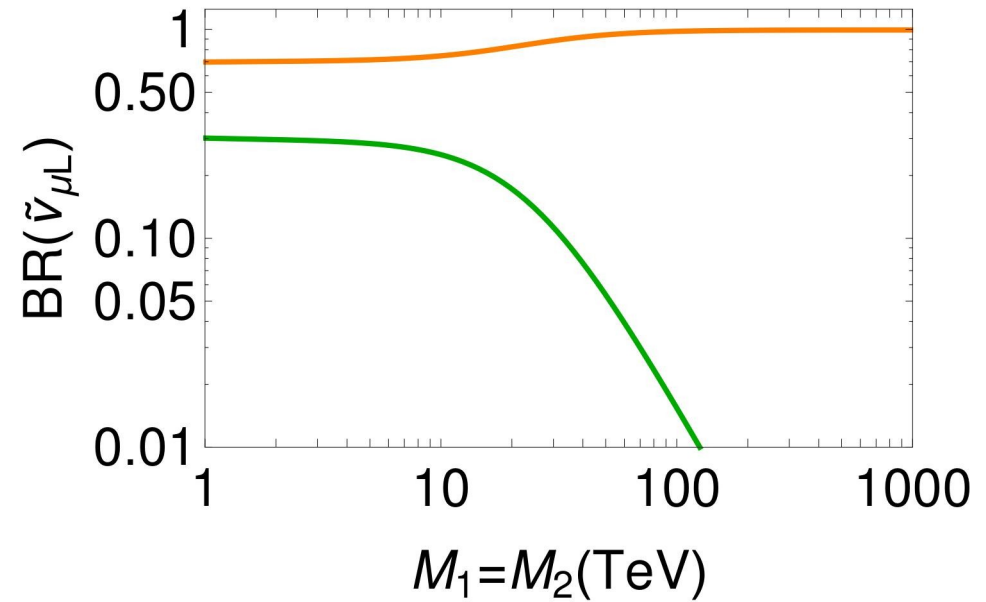
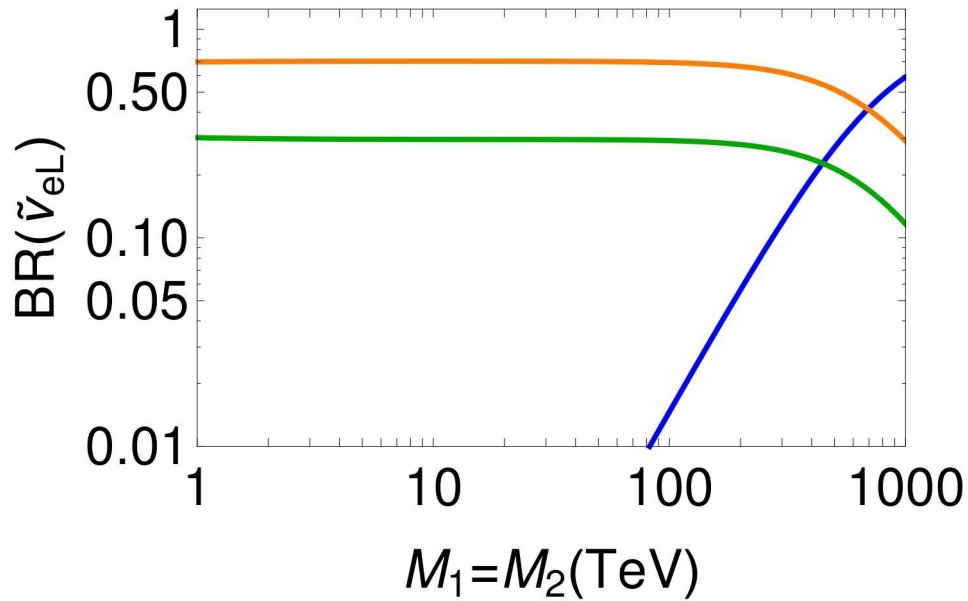
$$m_2 \simeq m_R$$

$$m_1 \simeq \frac{m_D^2}{m_R}$$

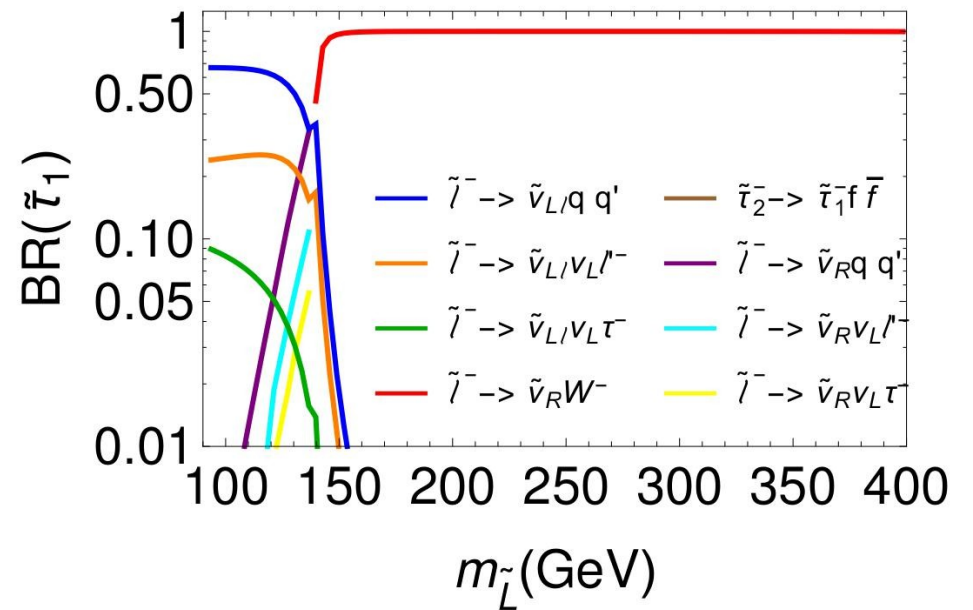
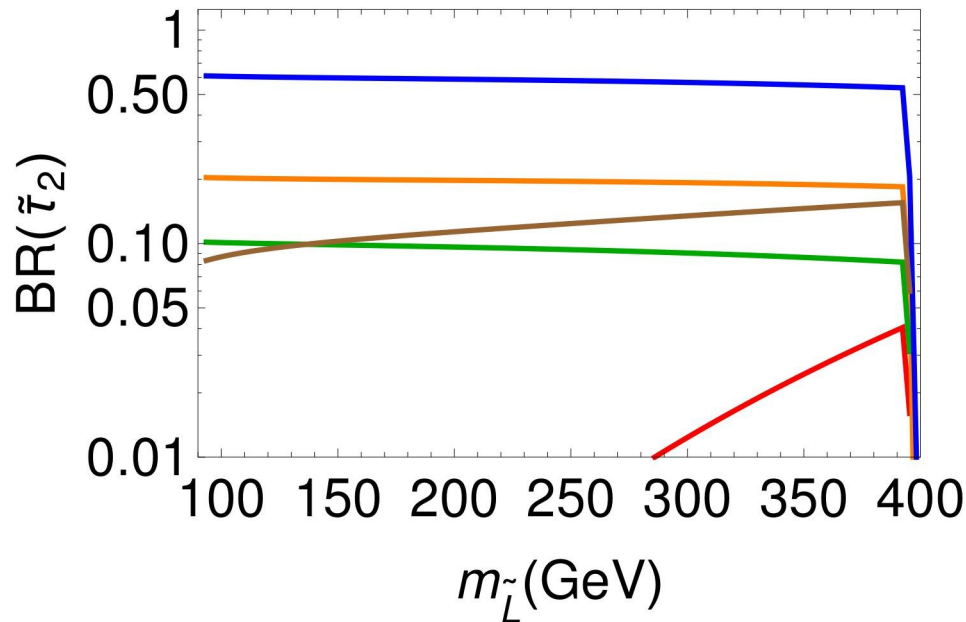
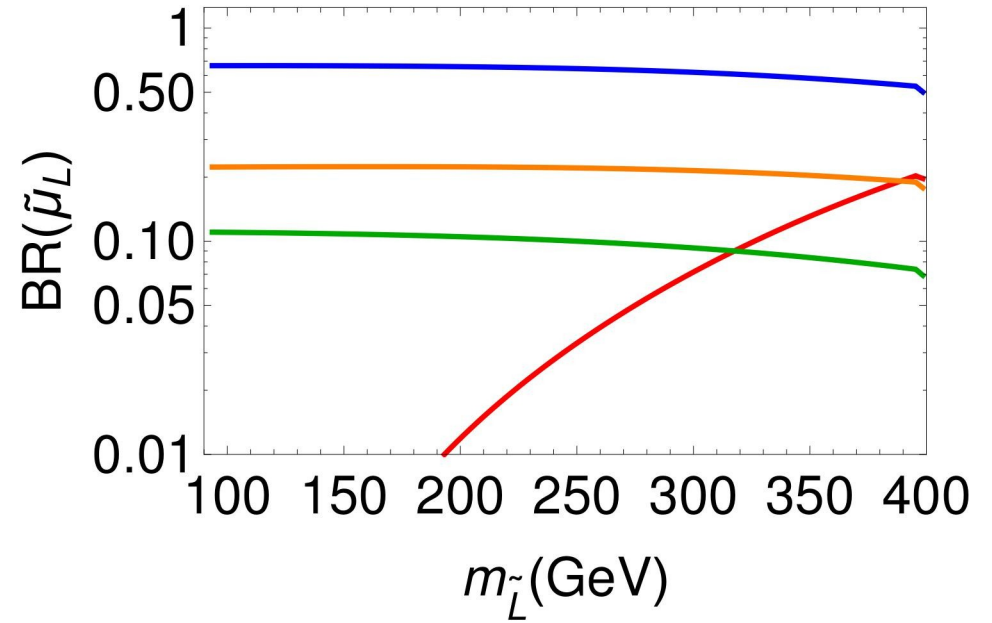
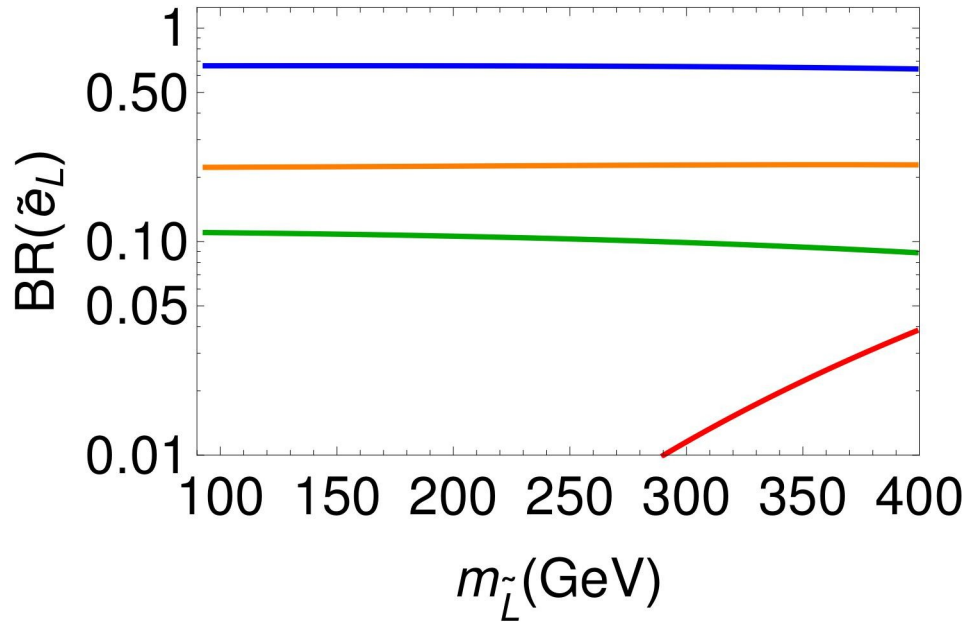
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