

Newton gauge cosmological perturbations for static spherically symmetric modifications of the de-Sitter metric

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SILAFEA 2018



PUCP



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Overview

- Static vs. Comoving Coordinates
- de Sitter
- Cosmological Perturbation Theory (CPT)
- SDS Metric: Static vs. Comoving
- SSS Metric: Static vs. Comoving
- Gauge-Invariant Turn Around Radius
- SSS Metrics

Static vs. Comoving

- Static Coordinates
 - Compute observables
 - Birkhoff's Theorem
 - Schwarzschild de Sitter Metric
- Comoving Coordinates
 - Cosmological Perturbation Theory (CPT)
 - Modified Gravity Theories
 - McVittie's Metric

de Sitter Background

- Static Coordinates

$$ds^2 = (1 - H^2 R^2) dT^2 - (1 - H^2 R^2)^{-1} dR^2 - R^2 d\Omega^2$$

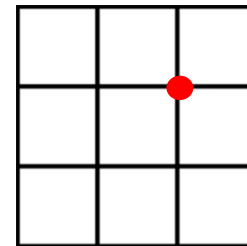
- Coordinate Transformation

$$R = a(t)r$$

$$T = t - \frac{1}{2H} \log(r^2 a^2(t) - H^{-2})$$

- Scale factor

$$a(t) = e^{Ht}$$



de Sitter Background

- Static Coordinates

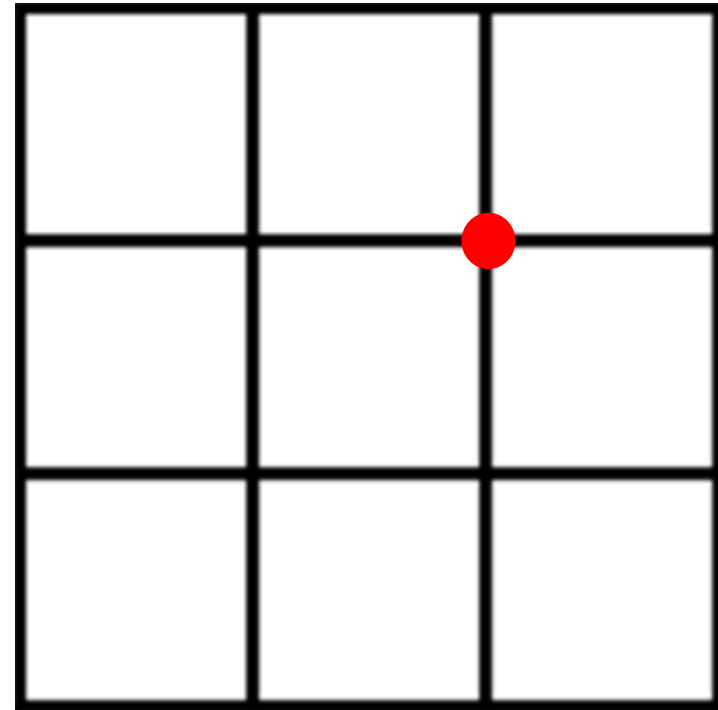
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de Sitter Background

- Static Coordinates

$$ds^2 = (1 - H^2 R^2) dT^2 - (1 - H^2 R^2)^{-1} dR^2 - R^2 d\Omega^2$$

- Coordinate Transformation

$$R = a(t)r \qquad T = t - \frac{1}{2H} \log(r^2 a^2(t) - H^{-2})$$

- Isotropic

$$ds^2 = dt^2 - a^2(t)(dr^2 + r^2 d\Omega^2)$$

Cosmological Perturbations

- Spherical Symmetry: 4 DOF

$$ds^2 = a^2 \left[(1 + 2\psi) d\tau^2 - \left(1 - 2\phi + \frac{2}{3}\mathcal{E} \right) dr^2 - 2\omega' d\tau dr - \left(1 - 2\phi - \frac{1}{3}\mathcal{E} \right) r^2 d\Omega^2 \right]$$

- Gauge Transformations

$$\tilde{x}^0 = x^0 + \zeta$$

$$\tilde{x}^i = x^i + \partial^i \beta$$



$$\tilde{\phi} = \phi - \frac{1}{3} \nabla^2 \beta + \frac{a_\tau}{a} \zeta,$$

$$\tilde{\omega} = \omega + \zeta + \beta_\tau,$$

$$\tilde{\psi} = \psi - \zeta_\tau - \frac{a_\tau}{a} \zeta,$$

$$\tilde{\chi} = \chi + 2\beta,$$

Cosmological Perturbations

- Spherical Symmetry: 4 DOF

$$ds^2 = a^2 \left[(1 + 2\psi) d\tau^2 - \left(1 - 2\phi + \frac{2}{3}\mathcal{E} \right) dr^2 - 2\omega' d\tau dr - \left(1 - 2\phi - \frac{1}{3}\mathcal{E} \right) r^2 d\Omega^2 \right]$$

- Gauge Transformations

$$\tilde{x}^0 = x^0 + \zeta$$

$$\tilde{x}^i = x^i + \partial^i \beta$$



$$\tilde{\phi} = \phi - \frac{1}{3} \nabla^2 \beta + \frac{a_\tau}{a} \zeta,$$

$$\tilde{\omega} = \omega + \zeta + \beta_\tau,$$

$$\tilde{\psi} = \psi - \zeta_\tau - \frac{a_\tau}{a} \zeta,$$

$$\tilde{\chi} = \chi + 2\beta, \quad \longrightarrow \quad \mathcal{E} = \chi'' - \frac{\chi'}{r}$$

Cosmological Perturbations

- Gauge Freedom: 2 DOF
- Newton Gauge

$$\omega_N = \chi_N = 0$$

- Bardeen Potentials: Gauge-Invariant

$$\Psi_B = \psi - \frac{1}{a} \left[a \left(\frac{\chi_\tau}{2} - \omega \right) \right]_\tau$$

$$\Phi_B = \phi + \frac{1}{6} \nabla^2 \chi - \frac{a_\tau}{a} \left(\omega - \frac{\chi_\tau}{2} \right)$$

Schwarzschild-de Sitter

- Spherically Symmetric Ansatz

$$ds^2 = F(T, R)dT^2 - H(T, R)dR^2 - R^2d\Omega^2$$

- Birkhoff's Theorem

$$ds^2 = \left(1 - \frac{2m}{R} - H^2 R^2\right) dT^2 - \left(1 - \frac{2m}{R} - H^2 R^2\right)^{-1} dR^2 - R^2 d\Omega^2$$

Schwarzschild-de Sitter

- Coordinate Transformation

$$t = \tilde{t} + \gamma(R) \quad \frac{d\gamma}{dR} = - \frac{HR^2}{\sqrt{R - m(1 - \frac{m}{R} - H^2 R^2)}}$$

Schwarzschild-de Sitter (Exact)

- Coordinate Transformation

$$t = \tilde{t} + \gamma(R) \qquad R = e^{H\tilde{t}}r + m + \frac{m^2}{4e^{H\tilde{t}}r}$$

- Isotropic

$$ds^2 = \left(\frac{1 - \frac{m}{2ar}}{1 + \frac{m}{2ar}} \right)^2 d\tilde{t}^2 - a^2 \left(1 + \frac{m}{2ar} \right)^4 [dr^2 + r^2 d\Omega^2]$$

Schwarzschild-de Sitter (CPT)

- Static Coordinates

$$ds^2 = \left(1 - \frac{2m}{R} - H^2 R^2\right) dT^2 - \left(1 - \frac{2m}{R} - H^2 R^2\right)^{-1} dR^2 - R^2 d\Omega^2$$

- de Sitter Transformation

$$R = a(t)r \qquad T = t - \frac{1}{2H} \log(r^2 a^2(t) - H^{-2})$$

- Conformal Time

$$d\tau = dt/a(t)$$

Schwarzschild-de Sitter (CPT)

- Static Coordinates

$$ds^2 = \left(1 - \frac{2m}{R} - H^2 R^2\right) dT^2 - \left(1 - \frac{2m}{R} - H^2 R^2\right)^{-1} dR^2 - R^2 d\Omega^2$$

- de Sitter Transformation

$$R = a(t)r \qquad T = t - \frac{1}{2H} \log(r^2 a^2(t) - H^{-2})$$

- Sub-horizon Approx.

$$m \ll R$$

Schwarzschild-de Sitter (CPT)

- Comoving-Conformal Coordinates

$$ds^2 = a^2 \left\{ \left[1 - \frac{2m (H^2 r^2 a^2 + 1)}{ra (H^2 r^2 a^2 - 1)^2} \right] d\tau^2 - \left[1 + \frac{2m (H^2 r^2 a^2 + 1)}{ra (H^2 r^2 a^2 - 1)^2} \right] dr^2 - \left[\frac{8Hm}{(H^2 r^2 a^2 - 1)^2} \right] d\tau dr - r^2 d\Omega^2 \right\}$$

- Scalar Perturbations

$$\psi = -\frac{m (H^2 r^2 a^2 + 1)}{ra (H^2 r^2 a^2 - 1)^2}$$

$$\phi = -\frac{H^2 m r^2 a^2 + m}{3ra (H^2 r^2 a^2 - 1)^2}$$

$$\omega' = \frac{4Hm}{(H^2 r^2 a^2 - 1)^2}$$

$$\mathcal{E} = \frac{2 (H^2 m r^2 a^2 + m)}{ra (H^2 r^2 a^2 - 1)^2}$$

Schwarzschild-de Sitter (CPT)

- Comoving-Conformal Coordinates

$$ds^2 = a^2 \left\{ \left[1 - \frac{2m (H^2 r^2 a^2 + 1)}{ra (H^2 r^2 a^2 - 1)^2} \right] d\tau^2 - \left[1 + \frac{2m (H^2 r^2 a^2 + 1)}{ra (H^2 r^2 a^2 - 1)^2} \right] dr^2 - \left[\frac{8Hm}{(H^2 r^2 a^2 - 1)^2} \right] d\tau dr - r^2 d\Omega^2 \right\}$$

- Scalar Perturbations

$$\psi = -\frac{m (H^2 r^2 a^2 + 1)}{ra (H^2 r^2 a^2 - 1)^2}$$

$$\omega = \frac{2m \tanh^{-1}(Hra)}{a} - \frac{2Hmr}{H^2 r^2 a^2 - 1}$$

$$\phi = -\frac{H^2 mr^2 a^2 + m}{3ra (H^2 r^2 a^2 - 1)^2}$$

$$\chi = \frac{2mr (Hra \tanh^{-1}(Hra) - 1)}{a}$$

Schwarzschild-de Sitter (CPT)

- Bardeen Potentials

$$\Psi_B = \psi - \frac{1}{a} \left[a \left(\frac{\chi_\tau}{2} - \omega \right) \right]_\tau = -\frac{m}{ar}$$

$$\Phi_B = \phi + \frac{1}{6} \nabla^2 \chi - \frac{a_\tau}{a} \left(\omega - \frac{\chi_\tau}{2} \right) = -\frac{m}{ar}$$

- Newton Gauge

$$ds^2 = a^2 \left[\left(1 - \frac{2m}{ar} \right) d\tau^2 - \left(1 + \frac{2m}{ar} \right) (dr^2 + r^2 d\Omega^2) \right]$$

SSS Metrics

- Static Spherically Symmetric Metrics

$$ds^2 = (1 - mh_t(R) - H^2 R^2) dT^2 - (1 - mh_r(R) - H^2 R^2)^{-1} dR^2 - R^2 d\Omega^2$$

- Comoving-Conformal

$$ds^2 = a^2 \left\{ \frac{H^6 r^6 a^6 - H^4 r^4 a^4 [mh_r(ra) + 3] - 1 + mh_t(ra)}{(H^2 r^2 a^2 - 1)^3} d\tau^2 \right. \\ + \frac{H^2 r^2 a^2 [m^2 h_r(ra)^2 + mh_r(ra) - mh_t(ra)] + 3H^2 r^2 a^2}{(H^2 r^2 a^2 - 1)^3} d\tau^2 \\ - \left[1 + \frac{mh_r(ra) (H^2 r^2 a^2 - mh_r(ra) - 1) + H^2 m r^2 a^2 (H^2 r^2 a^2 - 1) h_t(ra)}{(H^2 r^2 a^2 - 1)^3} \right] dr^2 \\ \left. - \frac{2Hmra [h_r(ra) (H^2 r^2 a^2 - mh_r(ra) - 1) + (H^2 r^2 a^2 - 1) h_t(ra)]}{(H^2 r^2 a^2 - 1)^3} d\tau dr - r^2 d\Omega^2 \right\}$$

SSS Metrics

- Static Spherically Symmetric Metrics

$$ds^2 = (1 - mh_t(R) - H^2 R^2) dT^2 - (1 - mh_r(R) - H^2 R^2)^{-1} dR^2 - R^2 d\Omega^2$$

- Subhorizon Scalar Perturbations

$$\psi = \frac{1}{2} mh_t(ra)$$

$$\omega = Hma \int r[h_r(ra) + h_t(ra)] dr$$

$$\phi = -\frac{1}{6} mh_r(ra)[mh_r(ra) + 1]$$

$$\chi = m \int k_1 \int \frac{h_r(k_2 a)}{k_2} dk_2 dk_1$$

SSS Metrics

- Static Spherically Symmetric Metrics

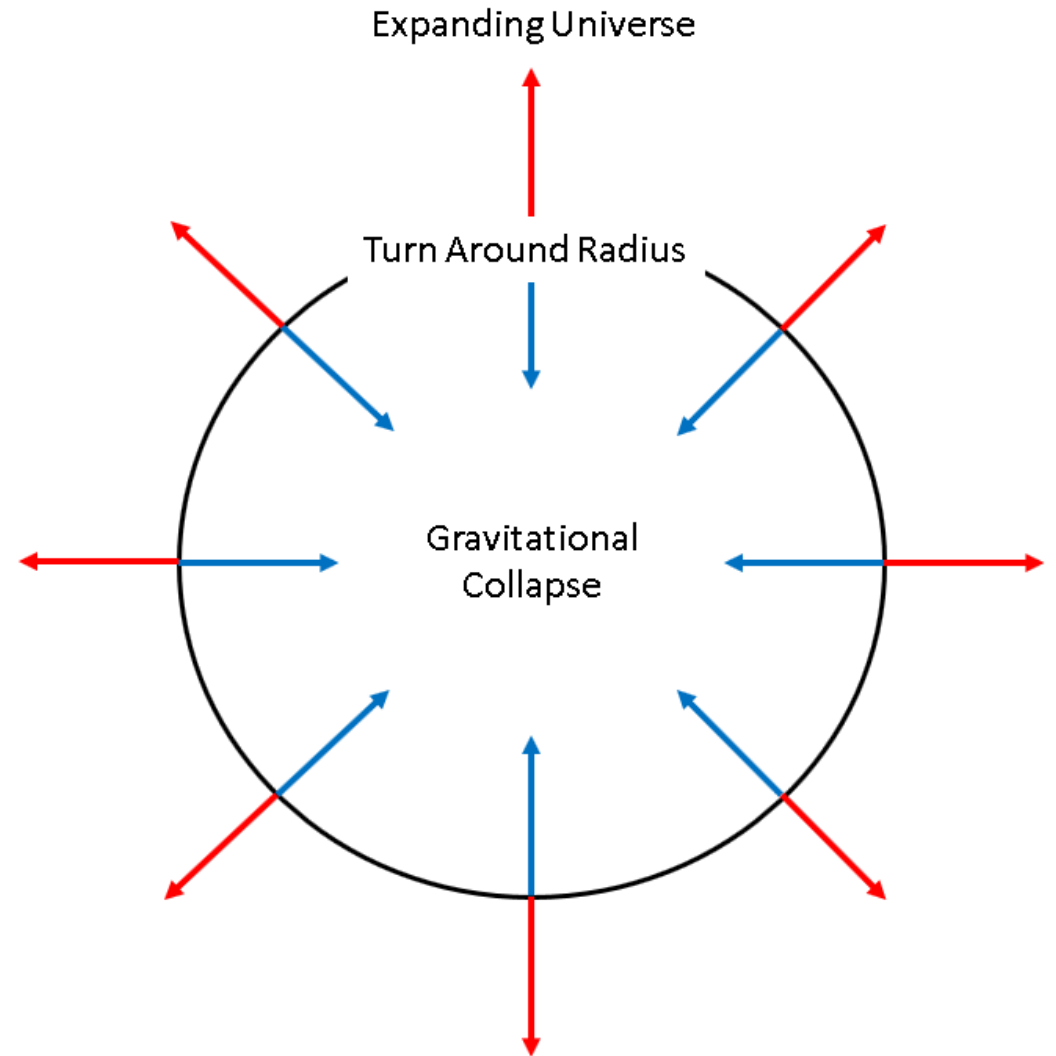
$$ds^2 = (1 - mh_t(R) - H^2 R^2) dT^2 - (1 - mh_r(R) - H^2 R^2)^{-1} dR^2 - R^2 d\Omega^2$$

- Bardeen Potentials

$$\Psi_B = -\frac{m}{2} h_t(ar) \quad \Phi_B = \frac{m}{2} \int \frac{h_r(ar) dr}{r}$$

Turn Around Radius

- Maximum Structure Size



Turn Around Radius

- Spherically Symmetric Ansatz

$$ds^2 = F(T, R)dT^2 - H(T, R)dR^2 - R^2 d\Omega^2$$

- Radial Geodesics

$$\frac{d^2 R}{ds^2} = \frac{1}{2}H(t, R)\dot{t}^2 \frac{\partial F(t, R)}{\partial R} + \frac{\dot{R}^2}{2H(t, R)} \frac{\partial H(t, R)}{\partial R} - \frac{\dot{R}\dot{t}}{H(t, R)} \frac{\partial H(t, R)}{\partial t}$$

- Static observer

$$\dot{R} = 0$$

Turn Around Radius

- Spherically Symmetric Ansatz

$$ds^2 = F(T, R)dT^2 - H(T, R)dR^2 - R^2 d\Omega^2$$

- Radial Geodesics

$$\frac{d^2 R}{ds^2} = \frac{1}{2}H(t, R)\dot{t}^2 \frac{\partial F(t, R)}{\partial R} + \frac{\dot{R}^2}{2H(t, R)} \frac{\partial H(t, R)}{\partial R} - \frac{\dot{t}^2}{H(t, R)} \frac{\partial H(t, R)}{\partial t}$$

- Static observer

$$\dot{R} = 0$$

Turn Around Radius

- Spherically Symmetric Ansatz

$$ds^2 = F(T, R)dT^2 - H(T, R)dR^2 - R^2 d\Omega^2$$

- Radial Geodesics

$$\frac{d^2 R}{ds^2} = \frac{1}{2} H(t, R) \dot{t}^2 \frac{\partial F(t, R)}{\partial R}$$

- Turn Around Radius

$$\partial_R F(R_{TA}) = 0$$

Turn Around Radius

- Newton Gauge

$$ds^2 = a^2 \left[(1 + 2\Psi) dt^2 - (1 - 2\Phi) (dr^2 + r^2 d\Omega^2) \right]$$

- Turn Around Radius

$$\ddot{a}r - \frac{\Psi'_N}{a} = 0$$

Turn Around Radius

- SSS Metrics

$$ds^2 = (1 - mh_t(R) - H^2 R^2) dT^2 - (1 - mh_r(R) - H^2 R^2)^{-1} dR^2 - R^2 d\Omega^2$$

- Turn Around Radius

$$F(R) = 1 - mh_t(R) - H^2 R^2 \qquad \partial_R F(R_{TA}) = 0$$

$$2H^2 R + mh'_t(R) = 0$$

SSS Metrics

SDS

- SSS functions

$$h_t = h_r = \frac{2}{R}$$

- Bardeen Potentials

$$\Psi = \Phi = -\frac{m}{ar}$$

SDS

- SSS functions

$$h_t = h_r = \frac{2}{R}$$

- Turn Around Radius

$$2H^2 R - 2\frac{m}{R^2} = 0 \quad \longrightarrow \quad R_{TA} = \sqrt[3]{\frac{m}{H^2}}$$

Brans-Dicke

- Static Metric

$$ds^2 = \left[1 - (1 + \epsilon) \frac{2m}{R} - (1 - 2\epsilon) H^2 R^2 \right] dt^2 - \left[1 - (1 - \epsilon) \frac{2m}{R} - (1 - 4\epsilon) H^2 R^2 \right] dR^2 - R^2 d\Omega^2$$

- Bardeen Potentials

$$\Psi_B = H^2 R^2 \epsilon - \frac{m(1 + \epsilon)}{R} \qquad \Phi_B = -H^2 R^2 \epsilon - \frac{m(1 - \epsilon)}{R}$$

- Newton Gauge

$$ds^2 = a^2 \left[\left(1 - \frac{2m(1 + \epsilon)}{ar} + 2\epsilon H^2 r^2 a^2 \right) \tau^2 - \left(1 - \frac{2m(1 + \epsilon)}{ar} - 4\epsilon H^2 r^2 a^2 \right) dr^2 - 8Hm d\tau dr - r^2 d\Omega^2 \right]$$

Brans-Dicke

- Static Metric

$$ds^2 = \left[1 - (1 + \epsilon) \frac{2m}{R} - (1 - 2\epsilon) H^2 R^2 \right] dt^2 - \left[1 - (1 - \epsilon) \frac{2m}{R} - (1 - 4\epsilon) H^2 R^2 \right] dR^2 - R^2 d\Omega^2$$

- Turn Around Radius

$$R_{TA} = ar_{TA} = \sqrt[3]{\frac{m}{H^2}} (1 + \epsilon)$$

Power Law

- SSS Functions

$$h_t(R) = \lambda_1 R^{n_1}$$

$$h_r(R) = \lambda_2 R^{n_2}$$

- Bardeen Potentials

$$\Psi_B = -\frac{m\lambda_1 (ra)^{n_1}}{2}$$

$$\Phi_B = \frac{m\lambda_2 (ra)^{n_2}}{2n_2}$$

Power Law

- SSS Functions

$$h_t(R) = \lambda_1 R^{n_1}$$

$$h_r(R) = \lambda_2 R^{n_2}$$

- Turn Around Radius

$$R_{TA} = \left(-\frac{m\lambda_1 n_1}{2H^2} \right)^{\frac{1}{2-n_1}}$$

Exponential

- SSS Functions

$$h_t(R) = \lambda_1 e^{b_1 R}$$

$$h_r(R) = \lambda_2 e^{b_2 R}$$

- Bardeen Potentials

$$\Psi_B = -\frac{1}{2} \lambda_1 m e^{b_1 r a}$$

$$\Phi_B = \frac{1}{2} \lambda_2 m \text{Ei}(b_2 r a)$$

Exponential

- SSS Functions

$$h_t(R) = \lambda_1 e^{b_1 R}$$

$$h_r(R) = \lambda_2 e^{b_2 R}$$

- Bardeen Potentials

$$Ei(z) = - \int_{-z}^{\infty} \frac{e^{-t}}{t} dt$$

$$\Phi_B = \frac{1}{2} \lambda_2 m Ei(b_2 r a)$$

Logarithmic

- SSS Functions

$$h_t(R) = \lambda_1 \log b_1 R \quad h_r(R) = \lambda_2 \log b_2 R$$

- Bardeen Potentials

$$\Psi_B = -\frac{1}{2} \lambda_1 m \log (b_1 r a) \quad \Phi_B = \frac{1}{4} \lambda_2 m [\log (b_2 r a)]^2$$

Logarithmic

- SSS Functions

$$h_t(R) = \lambda_1 \log b_1 R \quad h_r(R) = \lambda_2 \log b_2 R$$

- Turn Around Radius

$$R_{TA} = \frac{1}{H} \sqrt{\frac{m\lambda_1}{2}}$$

Flat Rotation Curves

- Metric

$$ds^2 = \left(\frac{R}{R_c} \right)^{2v^2} dt^2 - [1 - v^2 f(R) - H^2 R^2] dR^2 - R^2 d\Omega^2.$$

- Bardeen Potential

$$\Psi_B = v^2 \log \left(\frac{R}{R_c} \right) - \frac{1}{2} H^2 R^2.$$

Flat Rotation Curves

- Metric

$$ds^2 = \left(\frac{R}{R_c} \right)^{2v^2} dt^2 - [1 - v^2 f(R) - H^2 R^2] dR^2 - R^2 d\Omega^2.$$

- Density Profile

$$8\pi\rho = \frac{2v^2}{a^2 r^2} - 6H^2$$

Flat Rotation Curves

- Metric

$$ds^2 = \left(\frac{R}{R_c} \right)^{2v^2} dt^2 - [1 - v^2 f(R) - H^2 R^2] dR^2 - R^2 d\Omega^2.$$

- Linear Mass Profile

$$M(R) = \int_0^R 4\pi R'^2 \rho(R') dR' = \int_0^R (v^2 - 3H^2 R'^2) dR' = v^2 R - H^2 R^3$$

Conclusions

- Developed a method to obtain the Newton Gauge of Static Spherically Symmetric metrics.
- Obtained the Bardeen Potentials of SSS Metrics
- Applied this method to several examples
 - Schwarzschild-de Sitter
 - Brans-Dicke
 - Power Law, Exponential, Logarithmic Modifications
 - Flat Rotational Curves
- Useful when observations are convenient in the framework of CPT

Thanks!

References

- S. Bhattacharya, K. F. Dialektopoulos, A. E. Romano, C. Skordis, and T. N. Tomaras, (2016), arXiv:1611.05055.
- V. Faraoni, Phys. Rev. D81, 044002 (2010), arXiv:1001.2287.
- T. Biswas and A. Notari, JCAP 0806, 021 (2008), arXiv:astro-ph/0702555.
- J. M. Bardeen, Phys. Rev. D22, 1882 (1980).
- S. Bhattacharya, K. F. Dialektopoulos, A. E. Romano, and T. N. Tomaras, Phys. Rev. Lett. 115, 181104 (2015), arXiv:1505.02375.
- S. Bhattacharya, K. F. Dialektopoulos, and T. N. Tomaras, JCAP 1605, 036 (2016), arXiv:1512.08856.
- V. Pavlidou, N. Tetradis, and T. N. Tomaras, JCAP 1405, 017 (2014), arXiv:1401.3742.
- V. Faraoni, Phys. Dark Univ. 11, 11 (2016), arXiv:1508.00475.