



UNIVERSIDAD
DE SANTIAGO
DE CHILE

Dark matter production of scalar and vector type

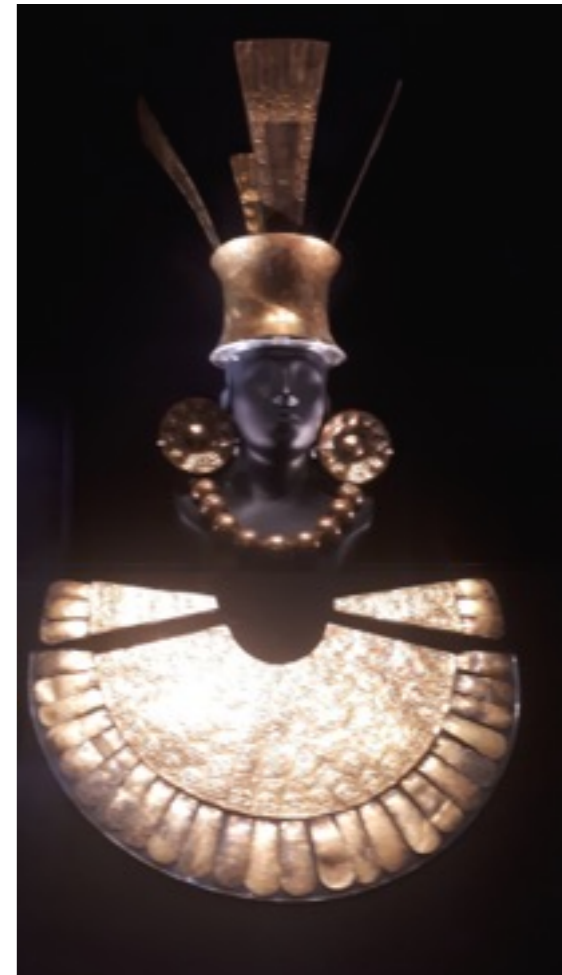


Paola Arias Reyes

Universidad de Santiago de Chile

**XII Latin American Symposium on
High Energy Physics**

29th November, 2018. Lima, Perú



- Evidence for dark matter from gravitational interaction, for instance rotation curves from galaxies, gravitational lensing...etc.
- Recently a galaxy with no dark matter was found. **Hint for a particle?**
- If so, should be cold (dust-like), non baryonic and *if* coupled to any SM particle by nuclear, weak or electromagnetic should be a very feeble coupling.
- Main candidates are LSP (WIMP dark matter), and the QCD axion. **More generically WISPs.**

The tale of axions and hidden photons as Cold Dark Matter (CDM)



axion led the way, as a solution of the absence of CP violation in the strong sector

$$\mathcal{L} \subset -\frac{\alpha_s}{8\pi} \frac{\phi}{f_\phi} G_{\mu\nu}^b \tilde{G}^{b,\mu\nu} - \frac{\alpha}{8\pi} \frac{\phi}{f_\phi} F_{\mu\nu} \tilde{F}^{\mu\nu}$$

spin-0
very light and
weakly coupled



axion led the way, as a solution of the absence of CP violation in the strong sector

$$\mathcal{L} \subset -\frac{\alpha_s}{8\pi} \frac{\phi}{f_\phi} G_{\mu\nu}^b \tilde{G}^{b,\mu\nu} - \frac{\alpha}{8\pi} \frac{\phi}{f_\phi} F_{\mu\nu} \tilde{F}^{\mu\nu}$$

Volume 120B, number 1,2,3

PHYSICS LETTERS

6 January 1983

spin-0
very light and
weakly coupled

THE NOT-SO-HARMLESS AXION

Michael DINE

The Institute for Advanced Study, Princeton, NJ 08540, USA

and

Willy FISCHLER

Department of Physics, University of Pennsylvania, Philadelphia, PA 19104, USA

Abstract: *Cosmological aspects of a very weakly interacting axion are discussed.... Demanding that **axions do not dominate the present energy density of the universe** is shown to give an upper bound on the axion decay constant of at most 10^{12} GeV.*

Hidden sector photons

TWO U(1)'S AND ϵ CHARGE SHIFTS

Bob HOLDOM

Department of Physics, University of Toronto, Toronto, Ontario, Canada M5S 1A7

Received 24 October 1985

If new particles are gauged by a new U(1) then their electromagnetic charges may be shifted by a calculable amount.

Hidden sector photons

TWO U(1)'S AND ϵ CHARGE SHIFTS

Bob HOLDOM

Department of Physics, University of Toronto, Toronto, Ontario, Canada M5S 1A7

Received 24 October 1985

If new particles are gauged by a new U(1) then their electromagnetic charges may be shifted by a calculable amount.

breaking of larger groups down to SM gauge group can yield hidden U(1) symmetries

Extra U(1)'s that mix kinetically with our visible photon

Hidden sector photons

TWO U(1)'S AND ϵ CHARGE SHIFTS

Bob HOLDOM

Department of Physics, University of Toronto, Toronto, Ontario, Canada M5S 1A7

Received 24 October 1985

If new particles are gauged by a new U(1) then their electromagnetic charges may be shifted by a calculable amount.

breaking of larger groups down to SM gauge group can yield hidden U(1) symmetries

Extra U(1)'s that mix kinetically with our visible photon

$$\mathcal{L} \supset -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} - \frac{1}{4}X_{\mu\nu}X^{\mu\nu} + \frac{\chi}{2}F_{\mu\nu}X^{\mu\nu} + \frac{m_{\gamma'}^2}{2}X_{\mu}X^{\mu}$$

Hidden sector photons

TWO U(1)'S AND ϵ CHARGE SHIFTS

Bob HOLDOM

Department of Physics, University of Toronto, Toronto, Ontario, Canada M5S 1A7

Received 24 October 1985

If new particles are gauged by a new U(1) then their electromagnetic charges may be shifted by a calculable amount.

breaking of larger groups down to SM gauge group can yield hidden U(1) symmetries

Extra U(1)'s that mix kinetically with our visible photon

$$\mathcal{L} \supset -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} - \frac{1}{4}X_{\mu\nu}X^{\mu\nu} + \frac{\chi}{2}F_{\mu\nu}X^{\mu\nu} + \frac{m_{\gamma'}^2}{2}X_{\mu}X^{\mu}$$

Kinetic mixing parameter, generated at loop level via heavy messenger exchange, predicted to be very small

Hidden sector photons

TWO U(1)'S AND ϵ CHARGE SHIFTS

Bob HOLDOM

Department of Physics, University of Toronto, Toronto, Ontario, Canada M5S 1A7

Received 24 October 1985

If new particles are gauged by a new U(1) then their electromagnetic charges may be shifted by a calculable amount.

breaking of larger groups down to SM gauge group can yield hidden U(1) symmetries

Extra U(1)'s that mix kinetically with our visible photon

$$\mathcal{L} \supset -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} - \frac{1}{4}X_{\mu\nu}X^{\mu\nu} + \frac{\chi}{2}F_{\mu\nu}X^{\mu\nu} + \frac{m_{\gamma'}^2}{2}X_{\mu}X^{\mu}$$

Kinetic mixing parameter, generated at loop level via heavy messenger exchange, predicted to be very small

very light, hidden Higgs or Stueckelberg mechanism

To sum up:

axion-like particles are spin-0 bosons, that couple effectively to 2-photons via

$$g\phi F_{\mu\nu}\tilde{F}^{\mu\nu}$$

Hidden photons are spin-1 gauge bosons, that kinetically mix with photons.

$$\chi F_{\mu\nu}X^{\mu\nu}$$

both can be very light $m \lesssim eV$

Non-thermal production mechanisms

Light scalars or zero modes of light vector bosons can be produced **non-thermally** during the early universe.

WISP dark matter

$$L_\phi = \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - \frac{1}{2} m_\phi^2 \phi^2$$

$$L_A = -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} + \frac{1}{2} m_A A_\mu A^\mu + \frac{1}{12} R A^\mu A_\mu$$

Arias et al, JCAP 2012

Light scalars or zero modes of light vector bosons can be produced **non-thermally** during the early universe.

WISP dark matter

$$L_\phi = \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - \frac{1}{2} m_\phi^2 \phi^2$$

$$L_A = -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} + \frac{1}{2} m_A A_\mu A^\mu + \frac{1}{12} R A^\mu A_\mu$$

Arias et al, JCAP 2012

In a FRW universe: $\ddot{\phi} + 3H\dot{\phi} + m^2(T)\phi = 0$

Light scalars or zero modes of light vector bosons can be produced **non-thermally** during the early universe.

WISP dark matter

$$L_\phi = \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - \frac{1}{2} m_\phi^2 \phi^2$$

$$L_A = -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} + \frac{1}{2} m_A A_\mu A^\mu + \frac{1}{12} R A^\mu A_\mu$$

Arias et al, JCAP 2012

In a FRW universe: $\ddot{\phi} + 3H\dot{\phi} + m^2(T)\phi = 0$

$3H \gg m$ a is frozen

Light scalars or zero modes of light vector bosons can be produced **non-thermally** during the early universe.

WISP dark matter

$$L_\phi = \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - \frac{1}{2} m_\phi^2 \phi^2$$

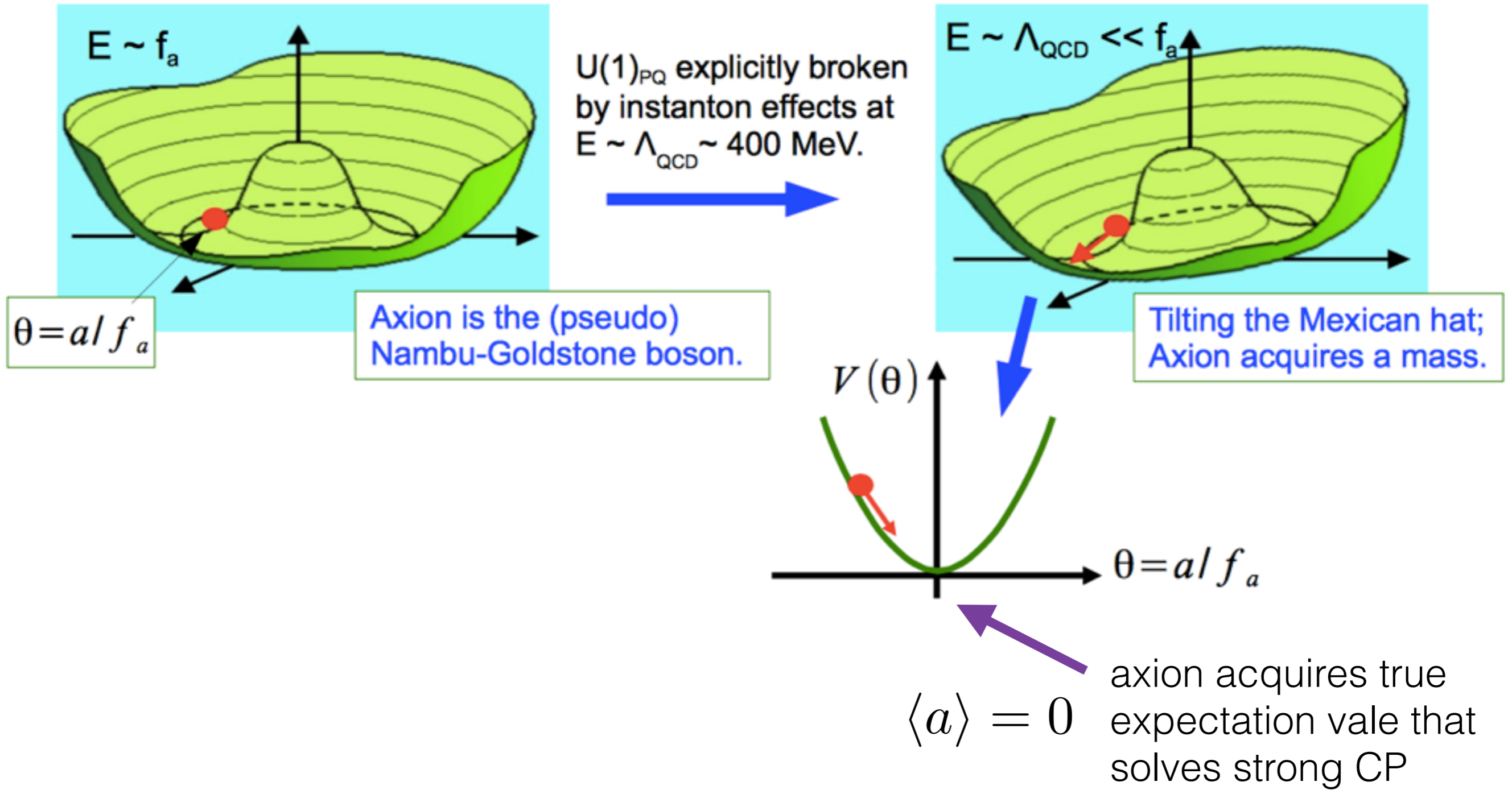
$$L_A = -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} + \frac{1}{2} m_A A_\mu A^\mu + \frac{1}{12} R A^\mu A_\mu$$

Arias et al, JCAP 2012

In a FRW universe: $\ddot{\phi} + 3H\dot{\phi} + m^2(T)\phi = 0$

$3H \gg m$ a is frozen

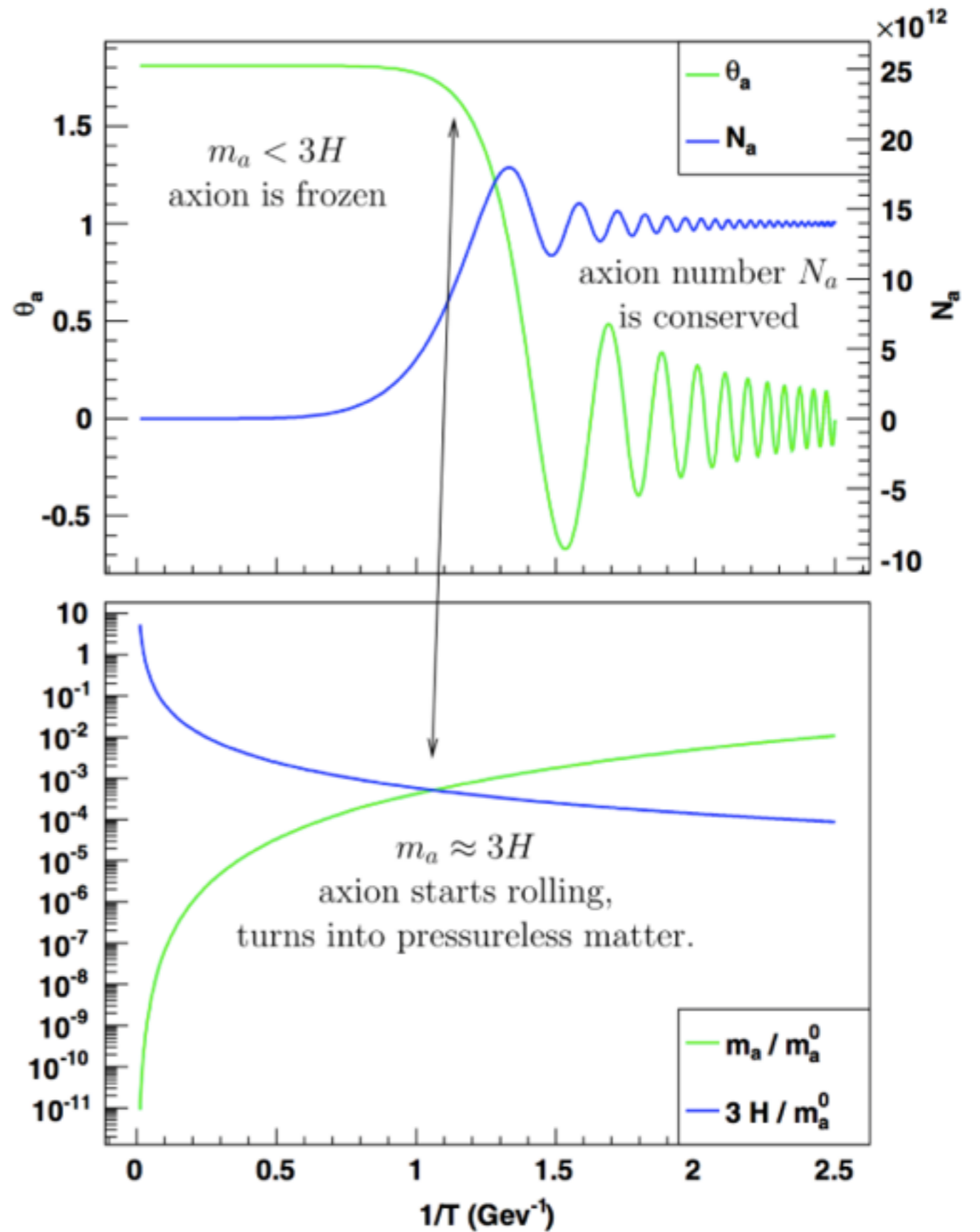
$3H \sim m$ the field “feels” its mass
starts to oscillate, coherently



$$\ddot{\phi} + 3H\dot{\phi} + m^2(T)\phi = 0$$

zero momentum condensate

Spatially uniform oscillating
classical field = coherent
state of many, extremely non-
relativistic particles = CDM



For vector U(1)s there is another interesting non-thermal production mechanism

Graham et al PRD 2015

Vector dark matter from inflationary perturbations


For vector U(1)s there is another interesting non-thermal production mechanism

Graham et al PRD 2015

Vector dark matter from inflationary perturbations

$$S = \int \sqrt{|g|} d^3x dt \left(-\frac{1}{4} F_{\mu\nu} F^{\mu\nu} - \frac{1}{2} m^2 A_\mu A^\mu \right)$$

massive abelian vector field present at inflation




For vector U(1)s there is another interesting non-thermal production mechanism

Graham et al PRD 2015

Vector dark matter from inflationary perturbations

$$S = \int \sqrt{|g|} d^3x dt \left(-\frac{1}{4} F_{\mu\nu} F^{\mu\nu} - \frac{1}{2} m^2 A_\mu A^\mu \right)$$

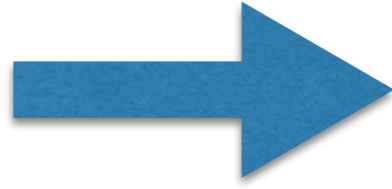
massive abelian vector field present at inflation



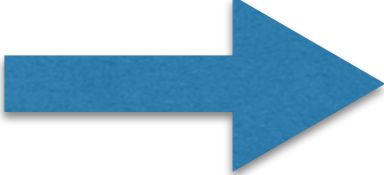
Transverse mode is almost conformally invariant. Longitudinal mode on the other hand

$$S_L = \int \frac{a^3 d^3k dt}{(2\pi)^3} \frac{1}{2a^3} \left(\frac{a^2 m^2}{k^2 + a^2 m^2} |\partial_t A_L|^2 - m^2 |A_L|^2 \right)$$

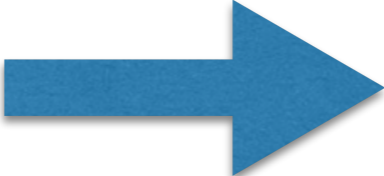
For $k \gg ma$



$$\pi(\vec{k}, t) = \frac{m}{k} A_L(\vec{k}, t)$$

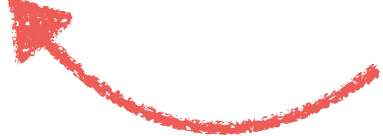
For $k \gg ma$  $\pi(\vec{k}, t) = \frac{m}{k} A_L(\vec{k}, t)$

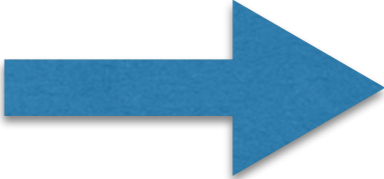
$$S_L \rightarrow \int \frac{a^3 d^3 k dt}{(2\pi)^3} \frac{1}{2} \left(|\partial_t \pi|^2 - \frac{k^2}{a^2} |\pi|^2 \right)$$

For $k \gg ma$  $\pi(\vec{k}, t) = \frac{m}{k} A_L(\vec{k}, t)$

$$S_L \rightarrow \int \frac{a^3 d^3 k dt}{(2\pi)^3} \frac{1}{2} \left(|\partial_t \pi|^2 - \frac{k^2}{a^2} |\pi|^2 \right)$$

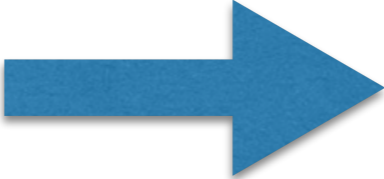
behaves as massless scalar



For $k \gg ma$  $\pi(\vec{k}, t) = \frac{m}{k} A_L(\vec{k}, t)$

$S_L \rightarrow \int \frac{a^3 d^3 k dt}{(2\pi)^3} \frac{1}{2} \left(|\partial_t \pi|^2 - \frac{k^2}{a^2} |\pi|^2 \right)$ behaves as massless scalar

For $m \gg H$ $\phi(\vec{k}, t) = \frac{m}{\sqrt{k^2 + a^2 m^2}} A_L(\vec{k}, t)$

For $k \gg ma$  $\pi(\vec{k}, t) = \frac{m}{k} A_L(\vec{k}, t)$

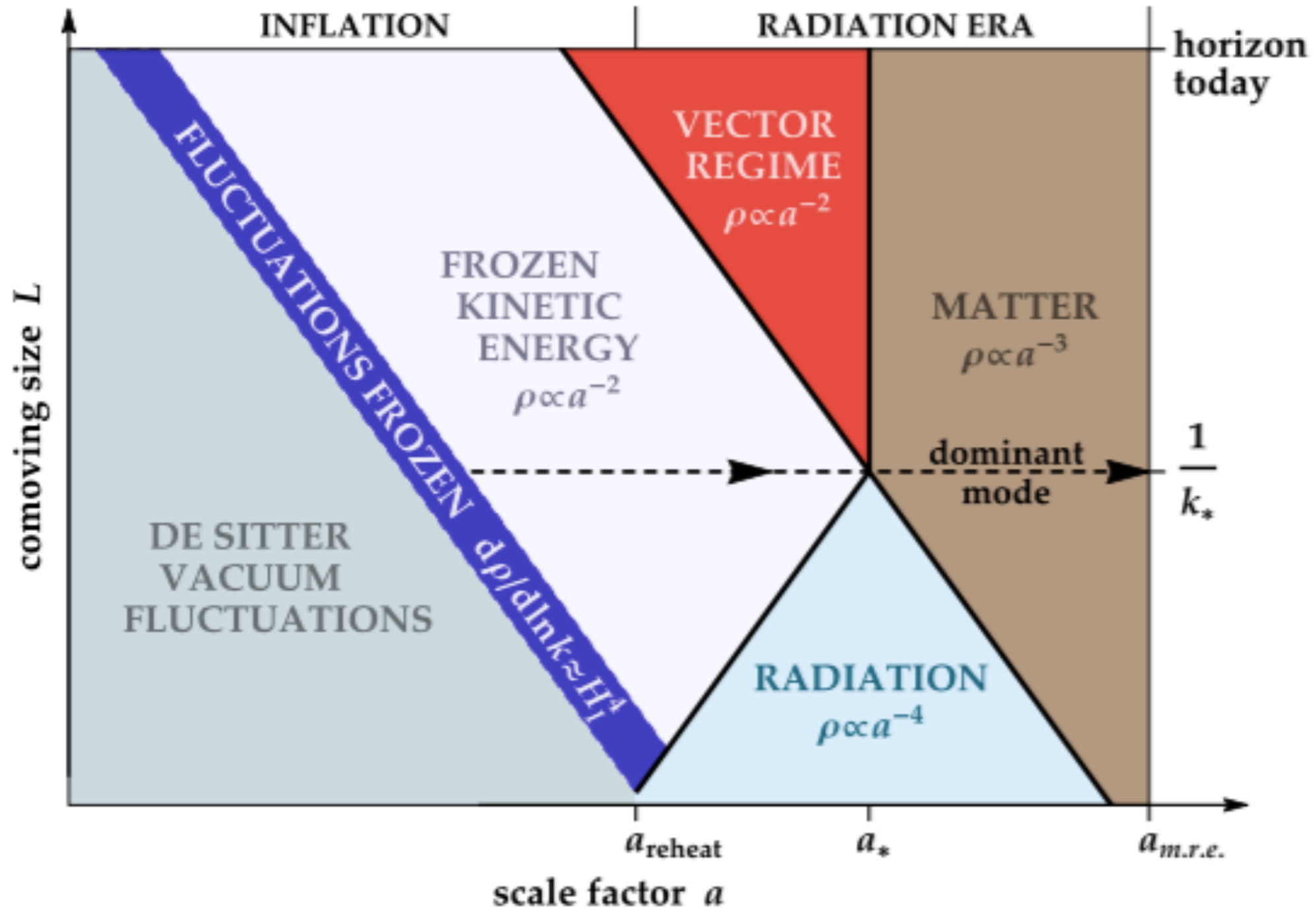
$$S_L \rightarrow \int \frac{a^3 d^3 k dt}{(2\pi)^3} \frac{1}{2} \left(|\partial_t \pi|^2 - \frac{k^2}{a^2} |\pi|^2 \right)$$

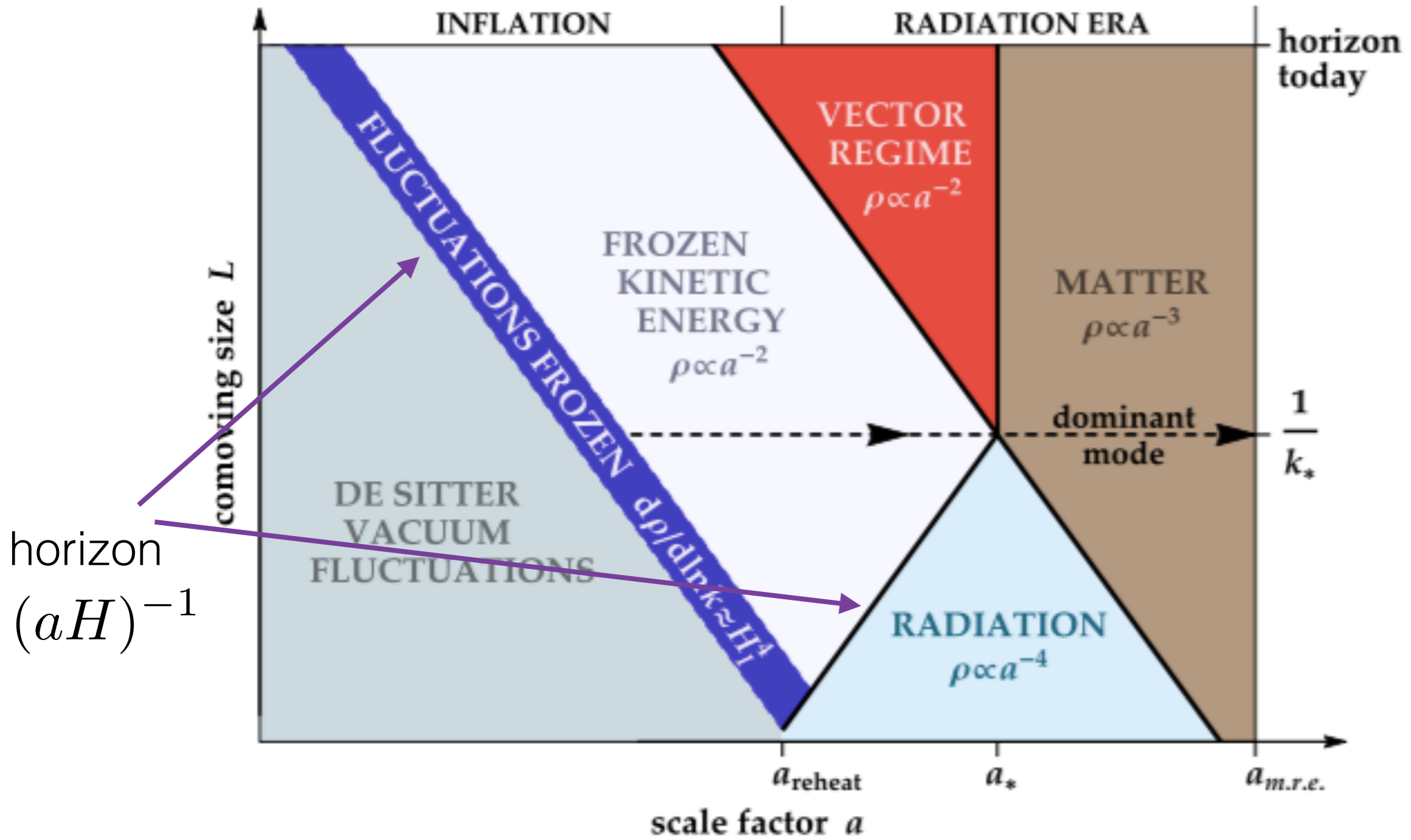
behaves as massless scalar

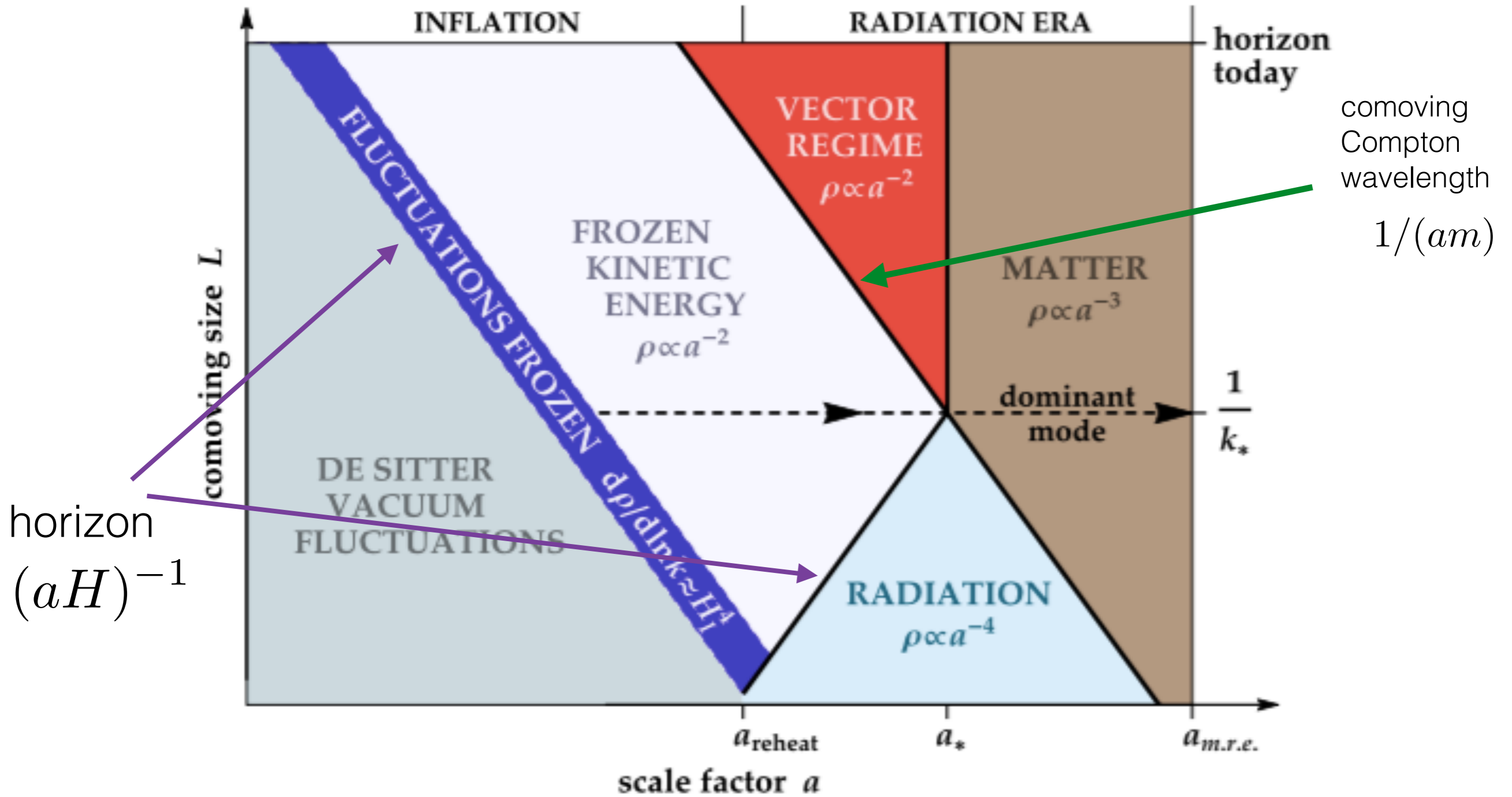
For $m \gg H$ $\phi(\vec{k}, t) = \frac{m}{\sqrt{k^2 + a^2 m^2}} A_L(\vec{k}, t)$

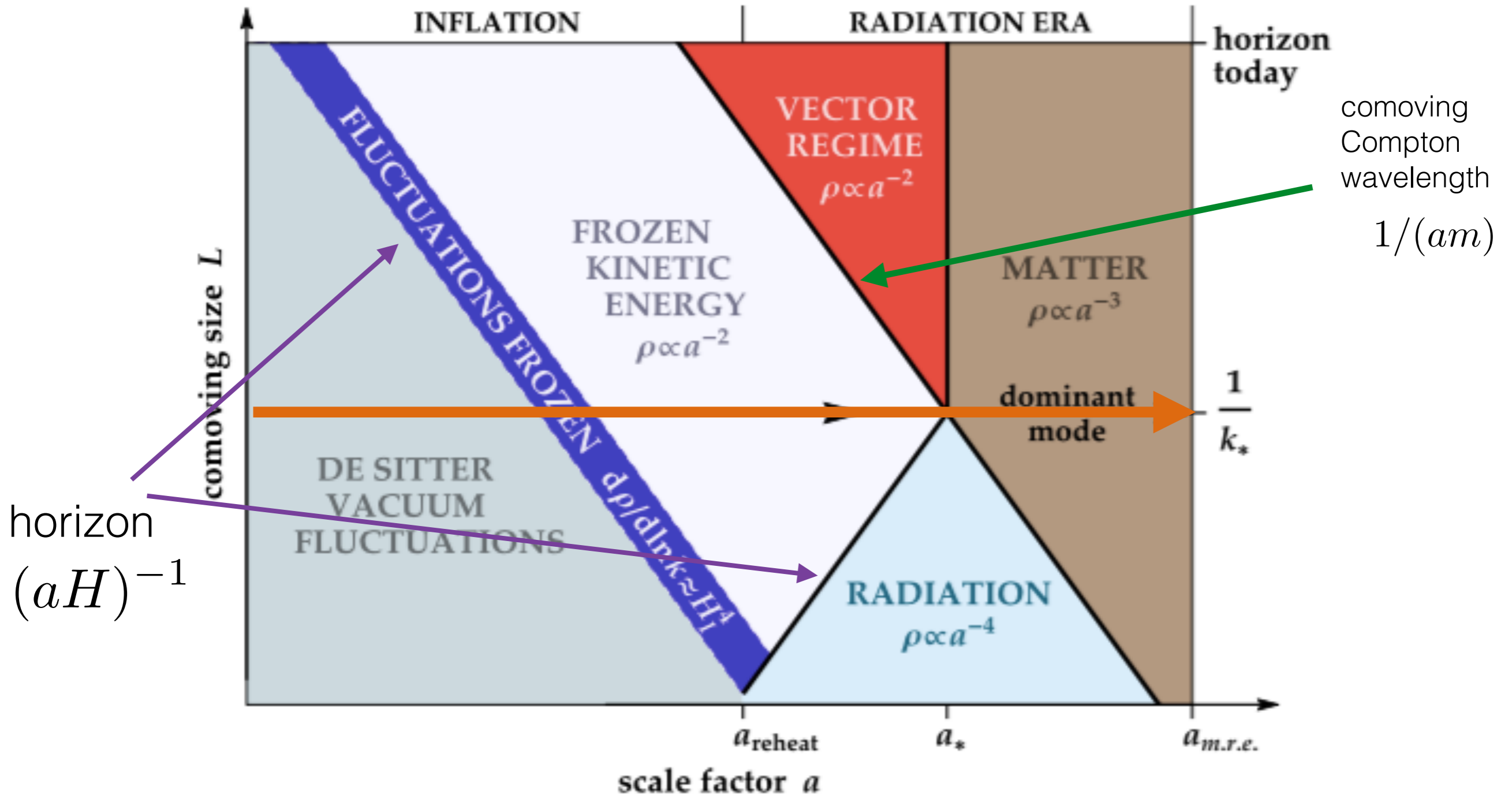
$$S_L \rightarrow \int \frac{a^3 d^3 k dt}{(2\pi)^3} \frac{1}{2} \left(|\partial_t \phi|^2 - \left(\frac{k^2}{a^2} + m^2 \right) |\phi|^2 \right)$$

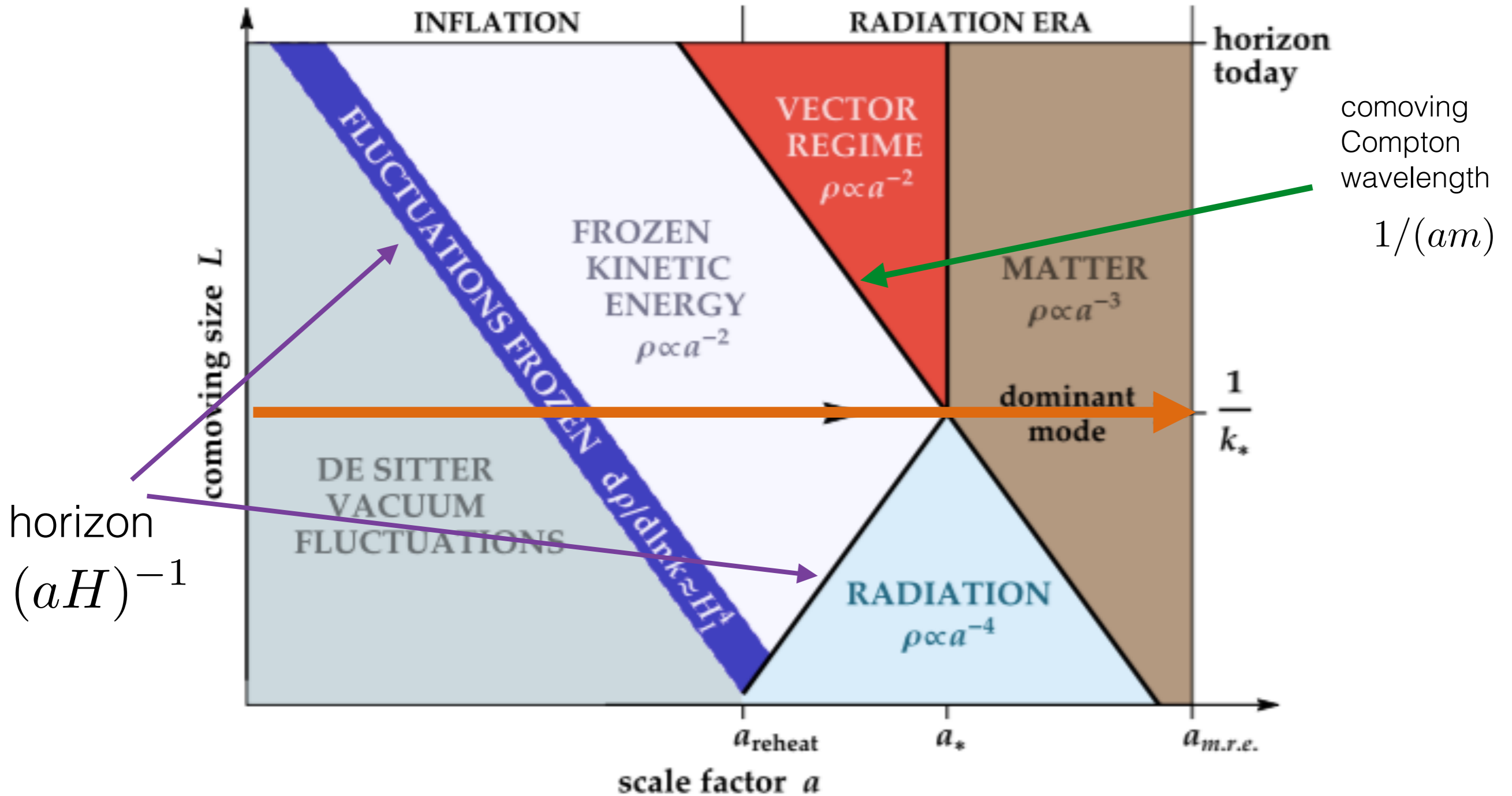
behaves as massive scalar





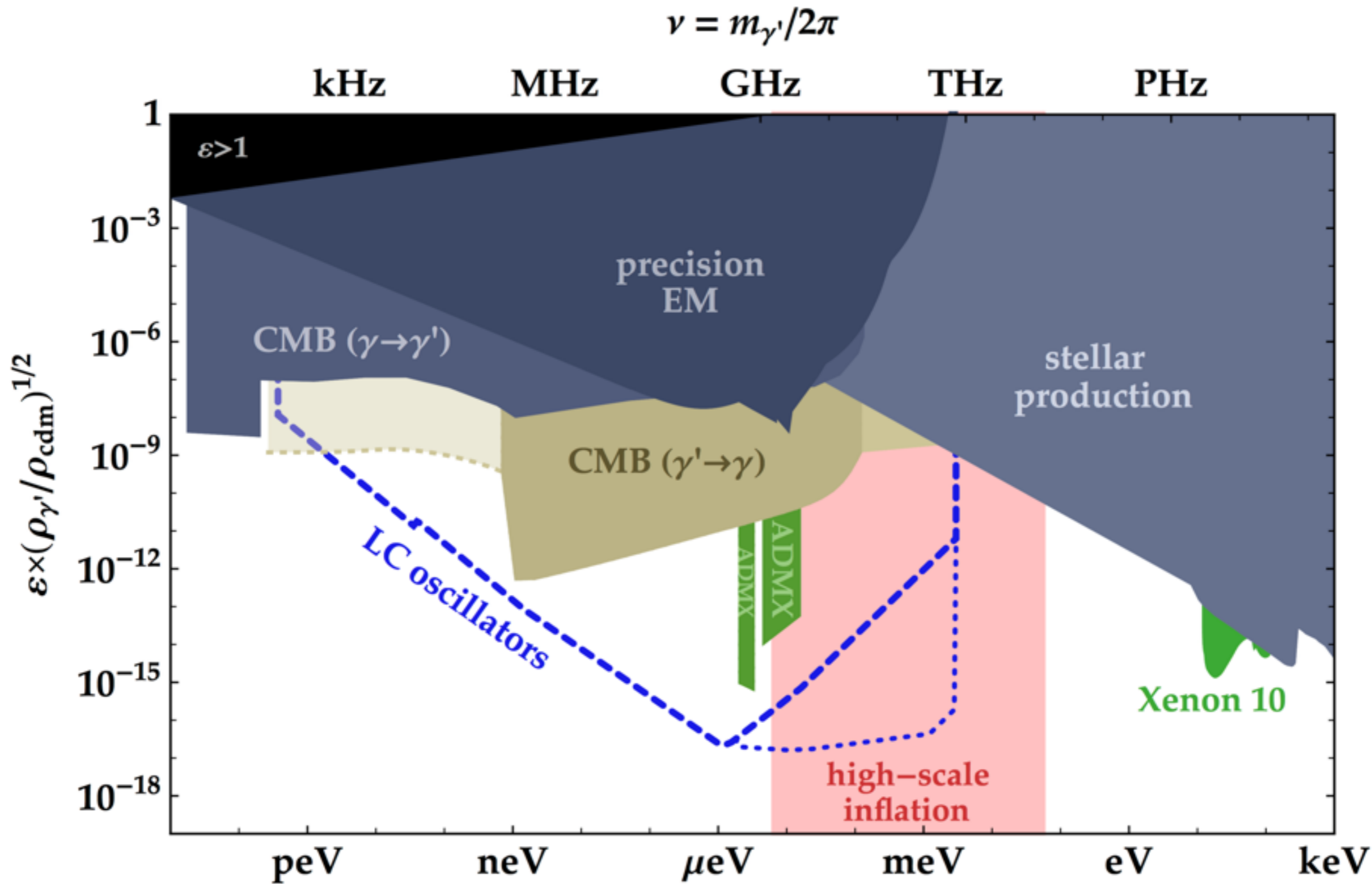


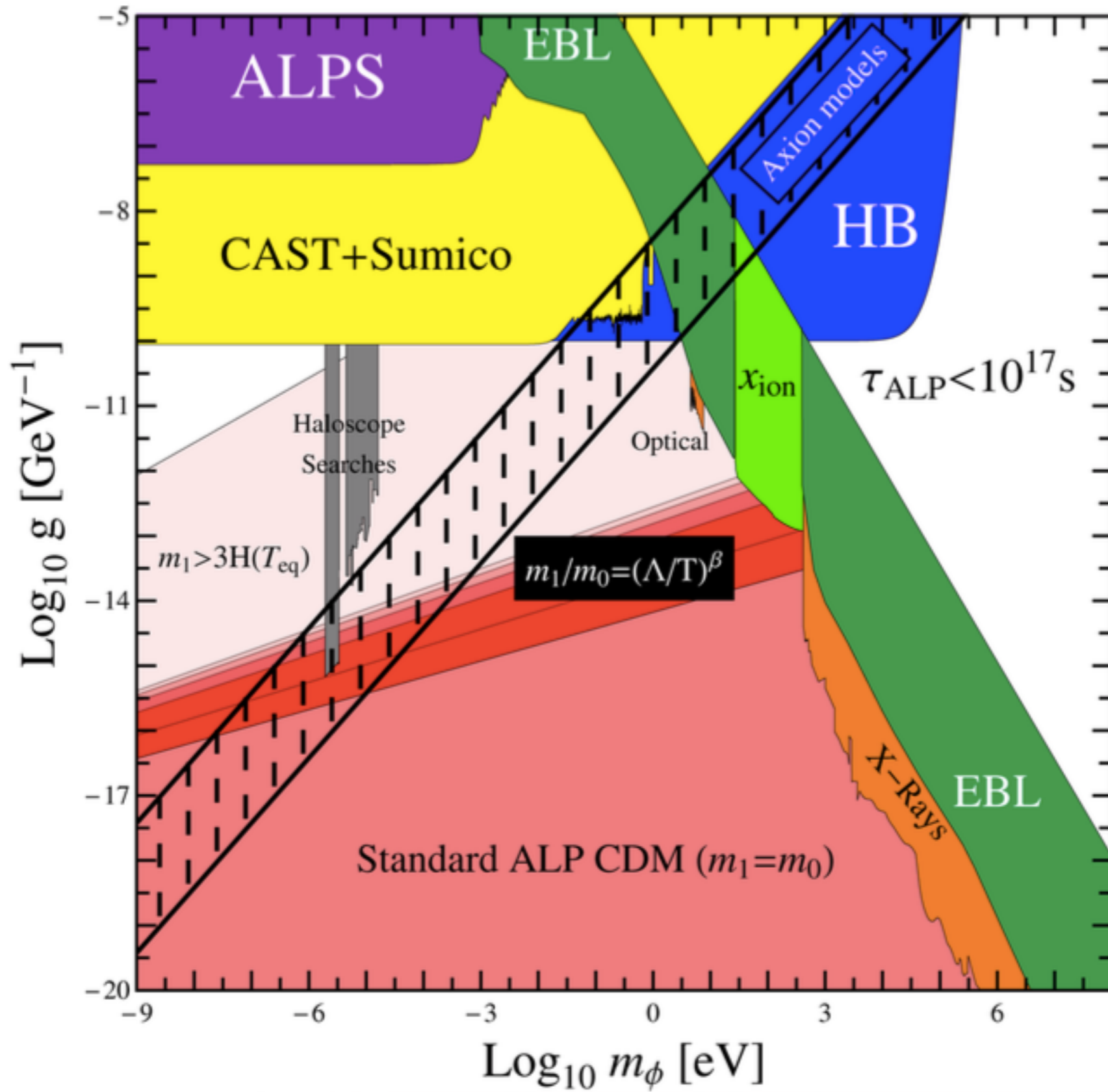




Dominant contribution from $H(a_*) = m \quad k^* = ma_*$

$$\frac{\Omega_A}{\Omega_{CDM}} = \sqrt{\frac{m}{6 \times 10^{-6} \text{ eV}}} \left(\frac{H_I}{10^{14} \text{ GeV}} \right)^2$$



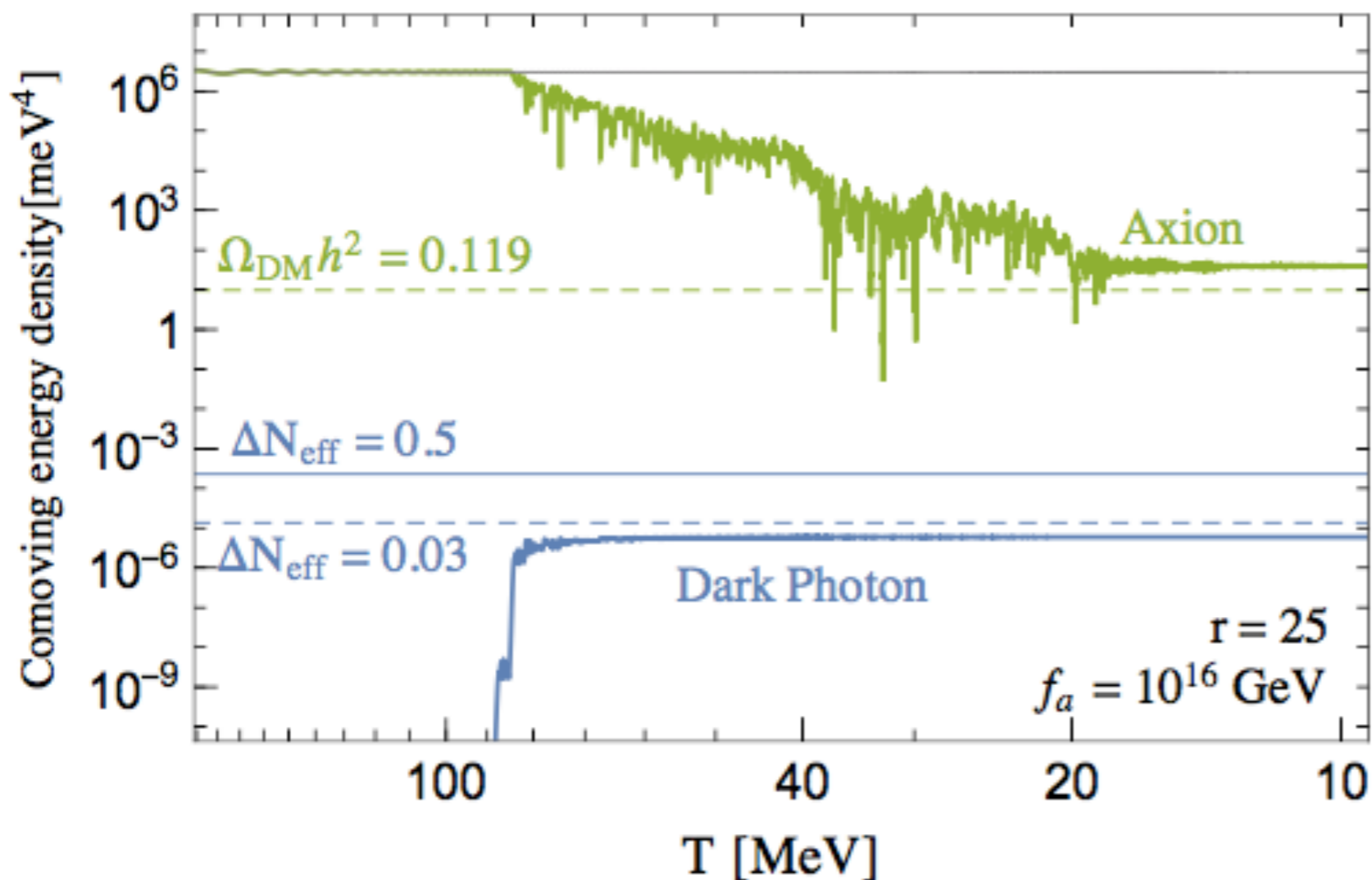


Coupling of axions to HP can be quite useful

Mixing of ALPs-HPs and photons can give rise to resonances in the probability of oscillation even in vacuum

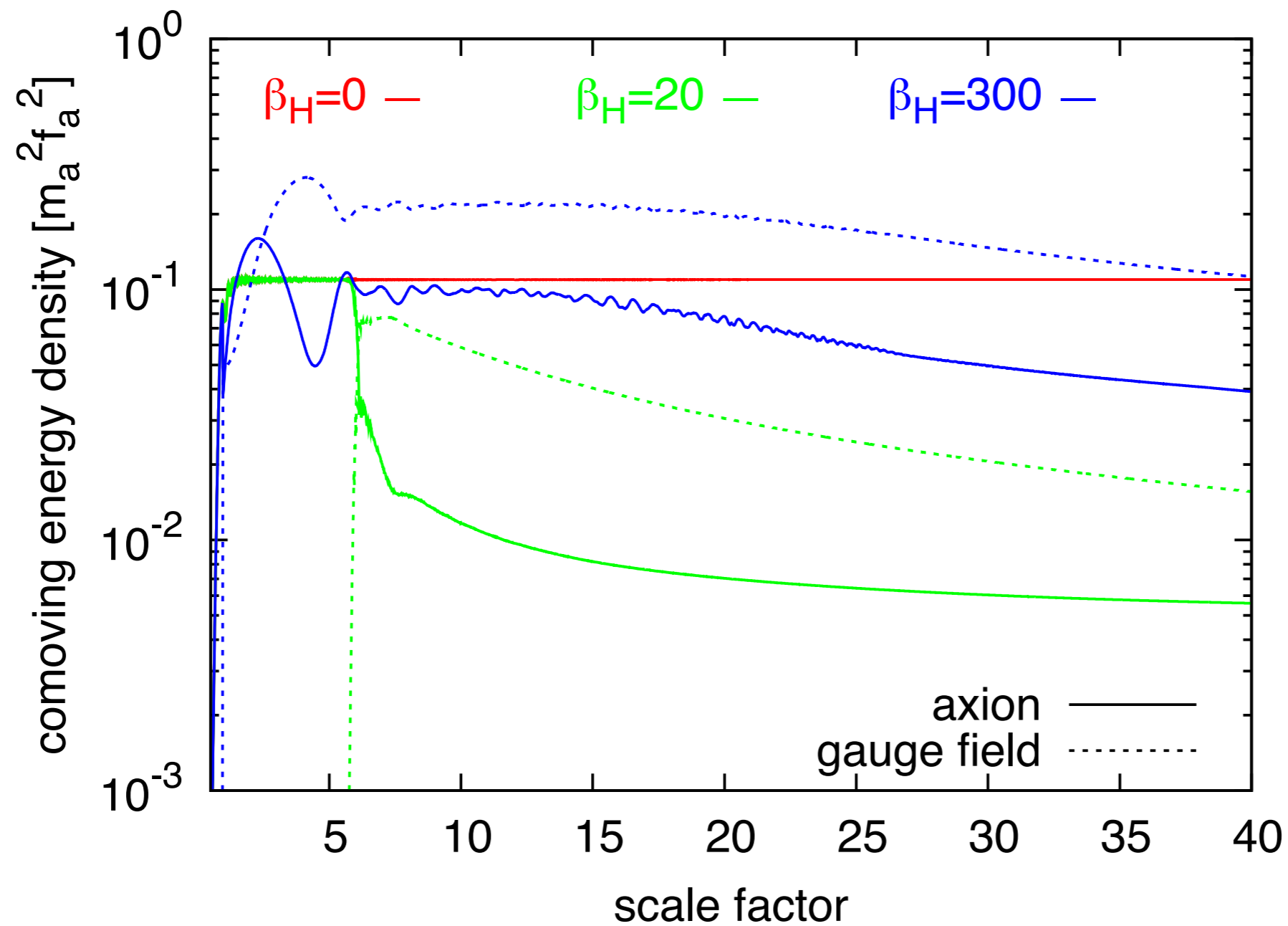
Alvarez, PA, Maldonado EPJC, 2018

Cosmology:



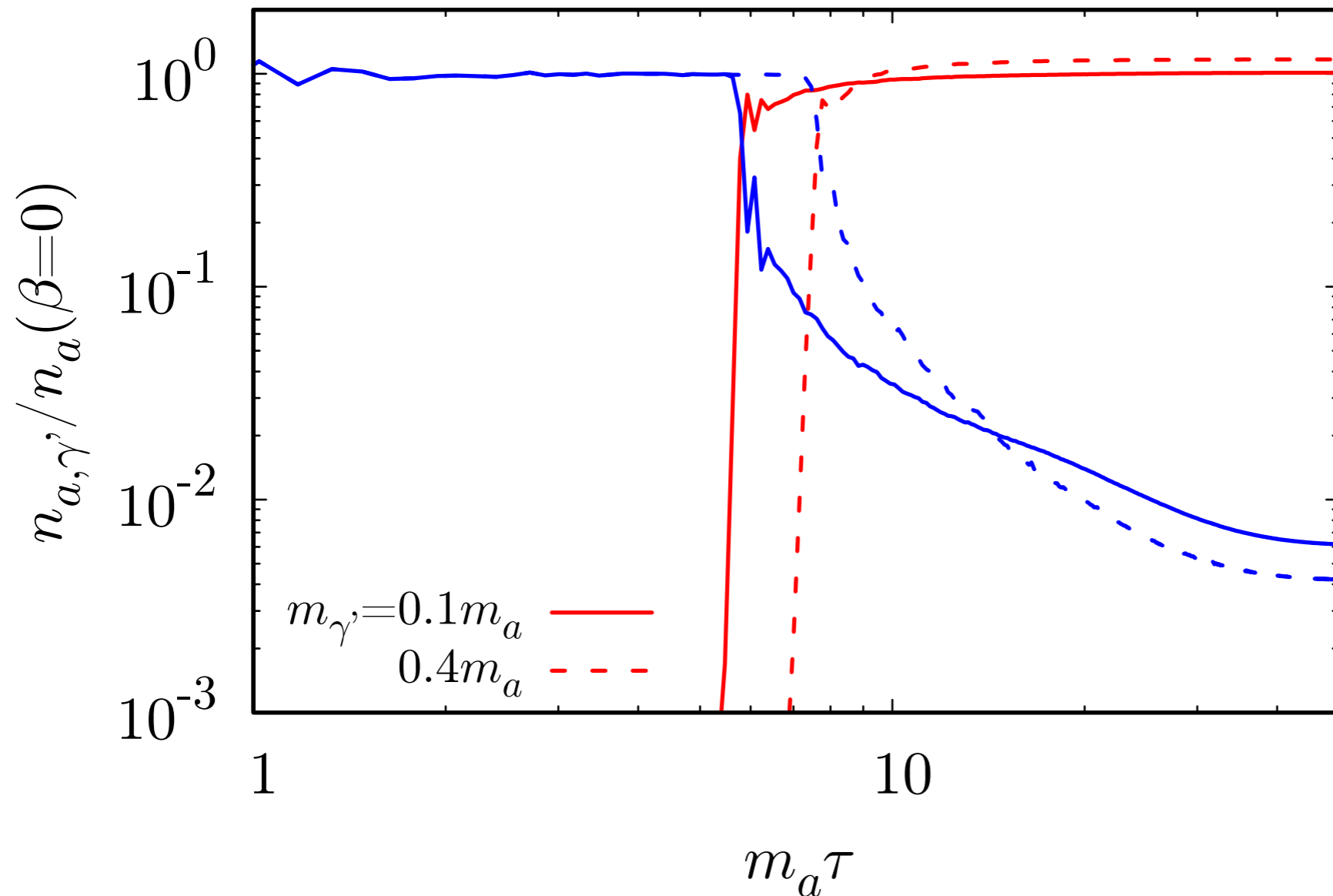
Coupling of QCD axion to massless HP opens up the axion window

Agrawal et al.
JHEP 1803 (2018) 049



also found in Kitajima
et al, PLB 781, 2018

Massive HP+ALP it is more exciting!



HP dark matter from axion energy transfer

$$m_{\gamma'}/m_a \sim 10^{-3} - 1$$

Outlook

- Many extensions of the standard model predict scalars and vectors weakly coupled to the visible sector
- Mechanisms to produce them non-thermally in the early universe are quite natural
- Many experiments worldwide that exploit the coupling to photons and matter
- Nowadays has become very popular to consider a scenario where both scalars and vectors are present in the early universe, interesting DM scenario.

Thank you!

And thanks to FONDECYT project 1161150 :)