



Polarization signatures from effective interactions of Majorana neutrinos

Lucía Duarte, Gabriel Zapata*, Oscar A. Sampayo*

PEDECIBA - UdelaR- Uruguay, *UNMdP - IFIMAR - Argentina

Outline

- 1 Motivation
- 2 Effective theory with N
- 3 $e^- p \rightarrow l_j^+ + 3jets$ ($l_j \equiv e, \mu$)
- 4 $e^- e^+ \rightarrow l^+ l^+ + 4jets$ ($l \equiv e, \mu, \tau$)

The SM picture: massless neutrinos

- Leptons: $SU(2)_L$ doublet

$$L_i = \begin{pmatrix} \nu_i \\ \ell_i \end{pmatrix}_L$$

- ℓ acquires a mass interacting with the Higgs v.e.v. after EWSB:

$$\langle \Phi \rangle = \begin{pmatrix} 0 \\ \frac{v}{\sqrt{2}} \end{pmatrix}$$

- Dirac mass: $m_D \ell_L \ell_R$

$$-\mathcal{L}_{Yukawa} \supset Y_\ell^{ij} \bar{L}^i \Phi \ell_R^j \rightarrow \frac{Y_\ell^{ij} v}{\sqrt{2}} \ell_L^i \ell_R^j$$

The SM picture: massless neutrinos

- Leptons: $SU(2)_L$ doublet

$$L_i = \begin{pmatrix} \nu_i \\ \ell_i \end{pmatrix}_L$$

- ℓ acquires a mass interacting with the Higgs v.e.v. after EWSB:

$$\langle \Phi \rangle = \begin{pmatrix} 0 \\ \frac{v}{\sqrt{2}} \end{pmatrix}$$

- Dirac mass: $m_{D\ell} \ell_L \ell_R$

$$-\mathcal{L}_{Yukawa} \supset Y_\ell^{ij} \bar{L}^i \Phi \ell_R^j \rightarrow \frac{Y_\ell^{ij} v}{\sqrt{2}} \ell_L^i \ell_R^j$$

- But...
- Neutrinos change flavor as they propagate: they have masses $m_\nu \sim 0.01 \text{ eV}$
- One has to go beyond the SM to get massive neutrinos.

Type I “vanilla” seesaw: neutrino mixing and N decoupling

$$\mathcal{L}_\nu = \mathcal{L}_{SM} - Y_{\alpha i} \bar{L}^\alpha \tilde{\Phi} N_{Ri} - \sum_{i,j=1}^3 \frac{M_{Nij}}{2} \bar{N}_{iL}^c N_{jR} + h.c.$$

- Mixing with Majorana massive states: lepton number violation (LNV)

$$\nu_{\ell L} = U_{\ell m} \nu_m + U_{\ell N} N$$

$$\mathcal{L}_W = -\frac{g}{\sqrt{2}} U_{\ell N} \bar{N}^c \gamma^\mu P_L \ell W_\mu^+$$

$$\mathcal{L}_Z = -\frac{g}{2c\theta_W} \bar{\nu}_\ell \gamma^\mu U_{\ell N} P_L N Z_\mu$$

$$\bullet U_{\ell N} \simeq \sqrt{\frac{m_\nu}{M_N}} \lesssim 10^{-6} \sqrt{\frac{100 \text{ GeV}}{M_N}}$$

Type I “vanilla” seesaw: neutrino mixing and N decoupling

$$\mathcal{L}_\nu = \mathcal{L}_{SM} - Y_{\alpha i} \bar{L}^\alpha \tilde{\Phi} N_{Ri} - \sum_{i,j=1}^3 \frac{M_{Nij}}{2} \bar{N}_{iL}^c N_{jR} + h.c.$$

- Mixing with Majorana massive states: lepton number violation (LNV)

$$\nu_{\ell L} = U_{\ell m} \nu_m + U_{\ell N} N$$

$$\mathcal{L}_W = -\frac{g}{\sqrt{2}} U_{\ell N} \bar{N}^c \gamma^\mu P_L \ell W_\mu^+$$

$$\mathcal{L}_Z = -\frac{g}{2c\theta_W} \bar{\nu}_\ell \gamma^\mu U_{\ell N} P_L N Z_\mu$$

$$\bullet U_{\ell N} \simeq \sqrt{\frac{m_\nu}{M_N}} \lesssim 10^{-6} \sqrt{\frac{100 \text{ GeV}}{M_N}}$$

- The observation of LNV depends only on $\nu_L - N$ mixing $U_{\ell N}$. So the N decouples (if no textures applied to mass matrix...)
- What kind of New Physics could lead to observable LNV ?

(e.g. in colliders)

Outline

1 Motivation

2 Effective theory with N

3 $e^- p \rightarrow l_j^+ + 3jets$ ($l_j \equiv e, \mu$)

4 $e^- e^+ \rightarrow l^+ l^+ + 4jets$ ($l \equiv e, \mu, \tau$)

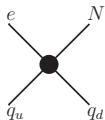
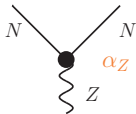
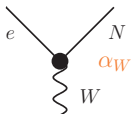
Effective approach [1]

- SM + one heavy Majorana N , $m_N < \Lambda$ (not integrated out...)
- Neglect the $\nu_L - N$ U_{eN} mixing
- NP parameterized with a lagrangian constructed with effective operators involving the N and the standard fields, preserving the $SU(2)_L \times U(1)_Y$ symmetry
- Low-energy limit of some unknown ultraviolet theory: suppressed by inverse powers of the new physics scale Λ :

$$\mathcal{L} = \mathcal{L}_{SM} + \sum_{n=6}^{\infty} \frac{1}{\Lambda^{n-4}} \sum_{\mathcal{J}} \alpha_{\mathcal{J}} \mathcal{O}_{\mathcal{J}}^{(n)}$$

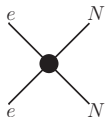
[1] F. del Aguila, S. Bar Shalom, A. Soni y J. Wudka. Phys. Lett. B 670, 399 (2009), 0806.0876

Effective operators



$$\alpha_{V_0}$$

$$\alpha_{S_1} \alpha_{S_2} \alpha_{S_3}$$



$$\alpha_{V_1} \alpha_{V_2}$$

$$\alpha_{S_0} \alpha_{S_4}$$

The (dim = 6) operators are [1]
(tree-level-generated):

$$\mathcal{O}_{LN\Phi}^{(i)} = (\Phi^\dagger \Phi)(\bar{L}_i N \tilde{\Phi})$$

$$\mathcal{O}_{NN\Phi} = \nu(\Phi^\dagger D_\mu \Phi)(\bar{N} \gamma^\mu N)$$

$$\mathcal{O}_{Ne\Phi}^{(i)} = \nu(\Phi^T \epsilon D_\mu \Phi)(\bar{N} \gamma^\mu e_i)$$

$$\mathcal{O}_{duNe}^{(i,j)} = (\bar{d}_i \gamma^\mu u_i)(\bar{N} \gamma_\mu e_j)$$

$$\mathcal{O}_{LNQd}^{(i,j)} = (\bar{L}_i N) \epsilon (\bar{Q}_j d_j)$$

$$\mathcal{O}_{QuNL}^{(i,j)} = (\bar{Q}_j u_i)(\bar{N} L_j)$$

$$\mathcal{O}_{QNLd}^{(i,j)} = (\bar{Q}_j N) \epsilon (\bar{L}_i d_j)$$

$$\mathcal{O}_{fNN}^{(i)} = (\bar{f}_i \gamma^\mu f_i)(\bar{N} \gamma_\mu N)$$

$$\mathcal{O}_{LNLe}^{(i,j)} = (\bar{L}_i N) \epsilon (\bar{L}_j e_j)$$

$$\mathcal{O}_{LN}^{(i)} = |\bar{N} L_i|^2$$

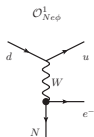
[1] F. del Aguila, S. Bar Shalom, A. Soni y J. Wudka. Phys. Lett. B 670, 399 (2009), 0806.0876

Bounds on the couplings $\alpha_{\mathcal{J}}^{(i,j)}$

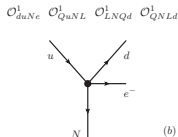
We exploit the existing bounds for the $U_{\ell N}$ mixings taking $U_{\ell N} \simeq \frac{\alpha V^2}{2\Lambda^2}$ (for $\Lambda = 1 \text{ TeV}$)

- Neutrinoless double beta decay (KamLAND-Zen)

$$\alpha_{0\nu\beta\beta}^{\text{bound}} \lesssim 3.2 \times 10^{-2} \left(\frac{m_N}{100 \text{ GeV}} \right)^{1/2}$$



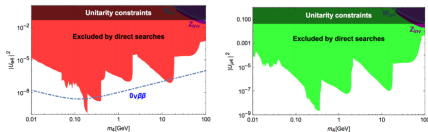
(a)



(b)

- Electroweak precision data (low energy LFV: $\mu \rightarrow e\gamma$)

$$\alpha_{EWPD}^{\text{bound}} \lesssim 0.32$$



[*] Abada et.al. JHEP02(2018)169, 1712.03984

Outline

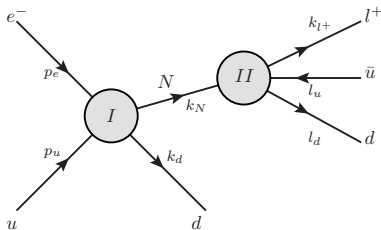
1 Motivation

2 Effective theory with N

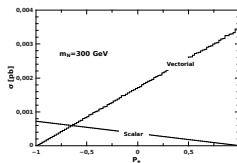
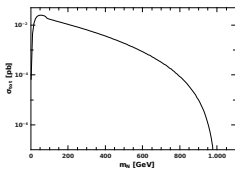
3 $e^- p \rightarrow l_j^+ + 3jets (l_j \equiv e, \mu)$

4 $e^- e^+ \rightarrow l^+ l^+ + 4jets (l \equiv e, \mu, \tau)$

$$e^- p \rightarrow l_j^+ + 3\text{jets} (l_j \equiv e, \mu) \quad [2]$$

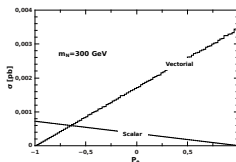
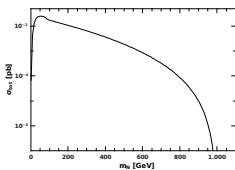
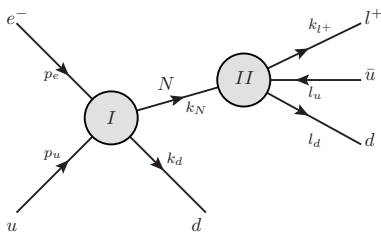


- **Vectorial** or **Scalar** contribution to σ depends on initial polarization P_{e^-}



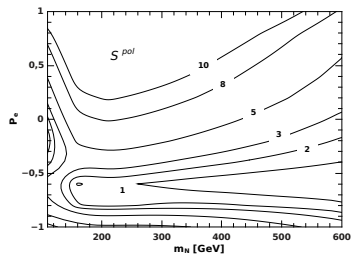
LHeC $E_p = 7 \text{ TeV}$, $E_e =$, $\mathcal{L} = 100 \text{ fb}^{-1}$

[2] L.Duarte, G. Zapata and O.A. Sampayo, Eur. Phys. J. C (2018) 78:352, 1802.07620

$e^- p \rightarrow l_j^+ + 3\text{jets} (l_j \equiv e, \mu)$ [2]

- **Vectorial** or **Scalar** contribution to σ depends on initial polarization P_{e^-}
- Standard deviation

$$S^{pol} = \frac{N^{vec} - N^{sca}}{\sqrt{N^{vec}} + \sqrt{N^{sca}}}$$

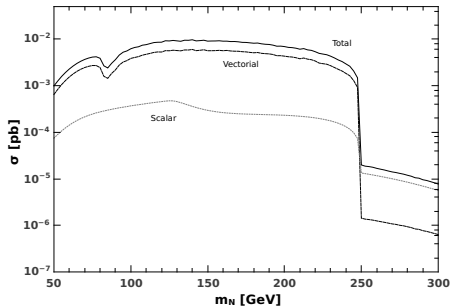
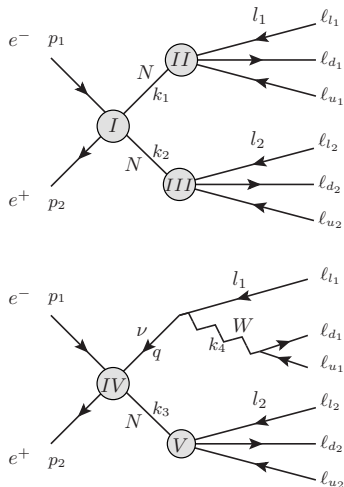


LHeC $E_p = 7$ TeV, $E_e =$, $\mathcal{L} = 100 \text{ fb}^{-1}$

[2] L.Duarte, G. Zapata and O.A. Sampayo, Eur. Phys. J. C (2018) 78:352, 1802.07620

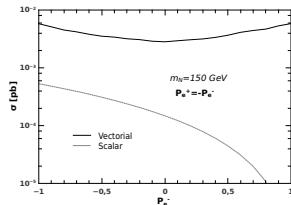
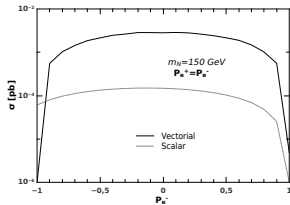
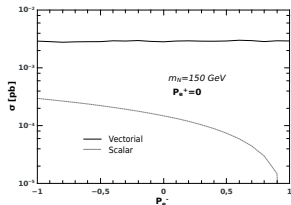
Outline

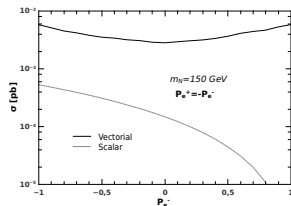
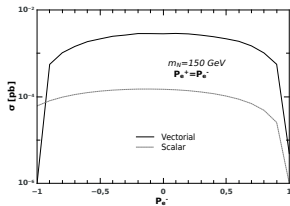
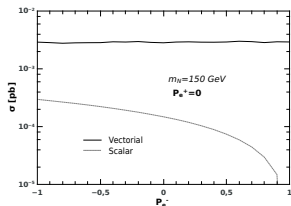
- 1 Motivation
- 2 Effective theory with N
- 3 $e^-p \rightarrow l_j^+ + 3jets$ ($l_j \equiv e, \mu$)
- 4 $e^-e^+ \rightarrow l^+l^- + 4jets$ ($l \equiv e, \mu, \tau$)

$e^-e^+ \rightarrow l^+l^- + 4jets$ [3]

$$\sqrt{s} = 500 \text{ GeV}$$

[3] L.Duarte, G. Zapata and O.A. Sampayo, 1812.XXXXX

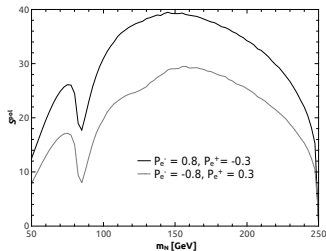
$e^-e^+ \rightarrow l^+l^- + 4jets$ [3]Initial polarization: distinguish **Vectorial** and **Scalar** behavior for σ :

$e^-e^+ \rightarrow l^+l^- + 4jets$ [3]Initial polarization: distinguish **Vectorial** and **Scalar** behavior for σ :

- Counting events:
 $\sqrt{s} = 500 \text{ GeV}$, $\mathcal{L} = 500 \text{ fb}^{-1}$ (ILC)

- Standard deviation

$$S^{pol} = \frac{N^{vec} - N^{sca}}{\sqrt{N^{vec}} + \sqrt{N^{sca}}}$$

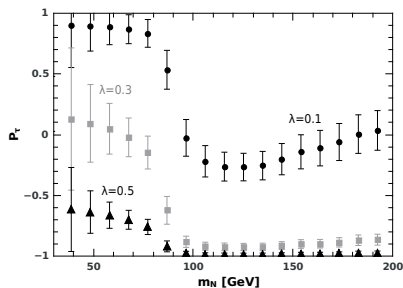
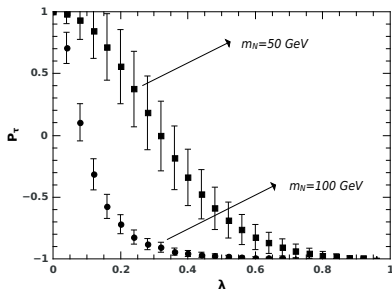


$e^-e^+ \rightarrow l^+l^+ + 4jets$ [3]

- Final $\tau^+\tau^+$ polarizations:

$$P_\tau = \frac{N_{+++} + N_{+-} - N_{-+} - N_{--}}{N_{+++} + N_{+-} + N_{-+} + N_{--}}$$

- $\lambda \in [0, 1]$ weights contributions: λ *Vectorial* + $(1 - \lambda)$ *Scalar*



[3] L.Duarte, G. Zapata and O.A. Sampayo, 1812.XXXXX

Thank you,
and the
XII SILAFAE
organizers.

