

# NEUTRAL CURRENT SMEFT STUDIES WITH EIC AND LHEC DIS PSEUDO DATA

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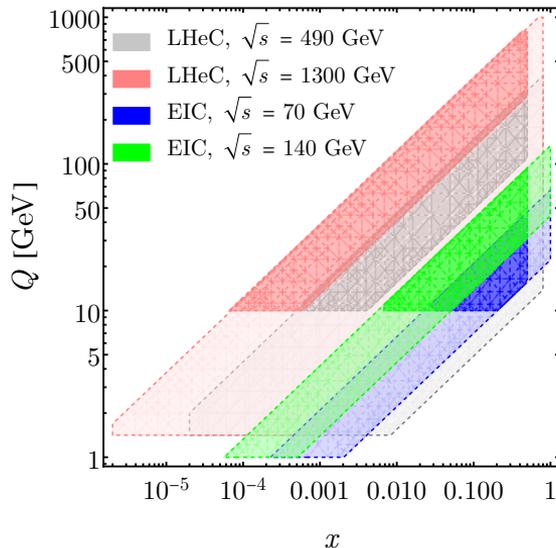
in collaboration with:

**Radja Boughezal and Kaan Simsek**

# OUR WORK IN A NUTSHELL

 We study the **BSM** potential of the LHeC and the EIC

 detailed accounting of anticipated uncertainties on **pseudo data**



 **Multidimensional fits**

of NC DIS cross section and asymmetries performed in the **SMEFT** framework

 We show that both the **EIC** and **LHeC** can improve upon the existing bound on the Z-boson couplings

# SMEFT

## Standard Model Effective Field Theory

$$\mathcal{L}_{\text{SMEFT}} = \mathcal{L}_{\text{SM}} + \sum_{n>4} \frac{1}{\Lambda^{n-4}} \sum_k C_k^{(n)} O_k^{(n)}$$

in this work  $n = 6$

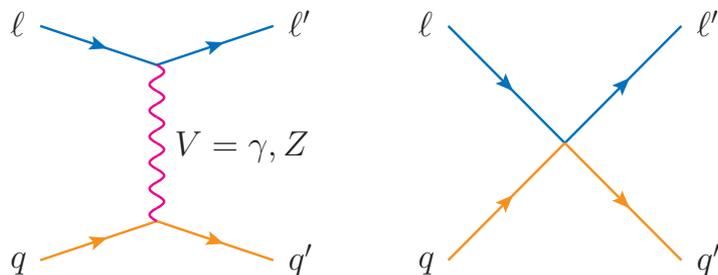
 Model independent

 Warsaw basis

**17 Wilson coefficients  
affect NC DIS matrix  
elements at LO**

<i>ffV</i>		semi-leptonic four-fermion	
$C_{\varphi WB}$	$O_{\varphi WB} = (\varphi^\dagger \tau^I \varphi) W_{\mu\nu}^I B^{\mu\nu}$	$C_{\ell q}^{(1)}$	$O_{\ell q}^{(1)} = (\bar{\ell} \gamma_\mu \ell) (\bar{q} \gamma^\mu q)$
$C_{\varphi D}$	$O_{\varphi D} = (\varphi^\dagger D_\mu \varphi)^* (\varphi^\dagger D^\mu \varphi)$	$C_{\ell q}^{(3)}$	$(\bar{\ell} \gamma_\mu \tau^I \ell) (\bar{q} \gamma^\mu \tau^I q)$
$C_{\varphi \ell}^{(1)}$	$O_{\varphi \ell}^{(1)} = (\varphi^\dagger i \overleftrightarrow{D}_\mu \varphi) (\bar{\ell} \gamma^\mu \ell)$	$C_{eu}$	$O_{eu} = (\bar{e} \gamma_\mu e) (\bar{u} \gamma^\mu u)$
$C_{\varphi \ell}^{(3)}$	$O_{\varphi \ell}^{(3)} = (\varphi^\dagger i \overleftrightarrow{D}_\mu \tau^I \varphi) (\bar{\ell} \gamma^\mu \tau^I \ell)$	$C_{ed}$	$O_{ed} = (\bar{e} \gamma_\mu e) (\bar{d} \gamma^\mu d)$
$C_{\varphi e}$	$O_{\varphi e} = (\varphi^\dagger i \overleftrightarrow{D}_\mu \varphi) (\bar{e} \gamma^\mu e)$	$C_{\ell u}$	$O_{\ell u} = (\bar{\ell} \gamma_\mu \ell) (\bar{u} \gamma^\mu u)$
$C_{\varphi q}^{(1)}$	$O_{\varphi q}^{(1)} = (\varphi^\dagger i \overleftrightarrow{D}_\mu \varphi) (\bar{q} \gamma^\mu q)$	$C_{\ell d}$	$O_{\ell d} = (\bar{\ell} \gamma_\mu \ell) (\bar{d} \gamma^\mu d)$
$C_{\varphi q}^{(3)}$	$O_{\varphi q}^{(3)} = (\varphi^\dagger i \overleftrightarrow{D}_\mu \tau^I \varphi) (\bar{q} \gamma^\mu \tau^I q)$	$C_{qe}$	$O_{qe} = (\bar{q} \gamma_\mu q) (\bar{e} \gamma^\mu e)$
$C_{\varphi u}$	$O_{\varphi u} = (\varphi^\dagger i \overleftrightarrow{D}_\mu \varphi) (\bar{u} \gamma^\mu u)$	$\ell, q$ : left handed doublets $e, u, d$ : right handed singlets $\varphi$ : SU(2) Higgs doublet	
$C_{\varphi d}$	$O_{\varphi d} = (\varphi^\dagger i \overleftrightarrow{D}_\mu \varphi) (\bar{d} \gamma^\mu d)$		
$C_{\ell \ell}$	$O_{\ell \ell} = (\bar{\ell} \gamma_\mu \ell) (\bar{\ell} \gamma^\mu \ell)$		

# DEEP INELASTIC SCATTERING



**LO**

Feynman diagrams for the partonic process mediating

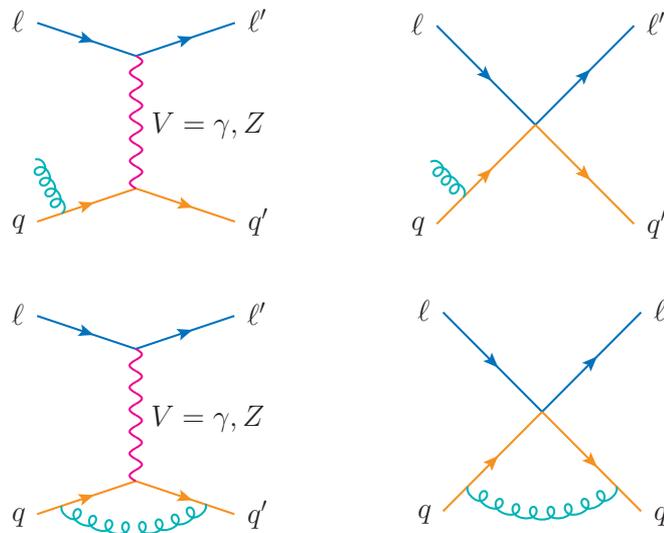
$$\ell + H \rightarrow \ell' + X$$

**NLO QCD**

corrections modify only the quark lines



they are identical for both SM and SMEFT cross sections



# SMEFT THEORY PREDICTIONS

## Structure of observables

We linearize our SMEFT expressions  $\mathcal{T}$

$$\mathcal{T} = \mathcal{T}^{\text{SM}} + \sum_k C_k \delta\mathcal{T}_k + \mathcal{O}(C_k^2)$$

$k$  runs over the active Wilson coefficients

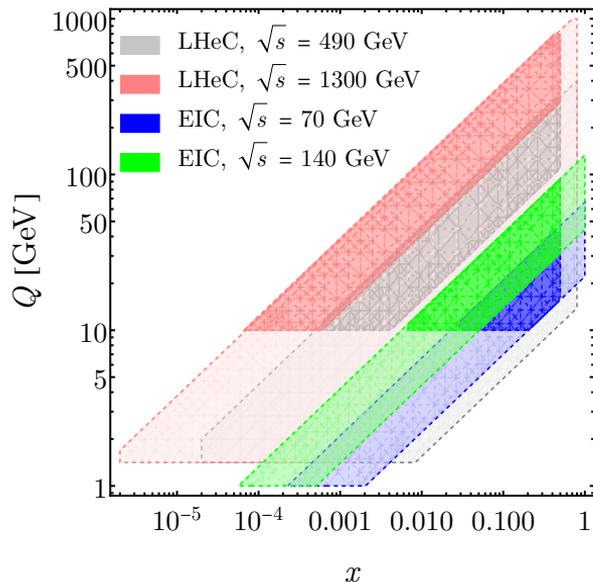
SMEFT shift associated with the Wilson coefficient  $C_k$

## SMEFT cross section

$$\sigma = \sigma^{\text{SM}} + \sum_i^{N_{\text{d6}}} \frac{C_i}{\Lambda^2} k_i + \sum_{ij}^{N_{\text{d6}}} \frac{C_i C_j}{\Lambda^4} \tilde{k}_{ij} + \dots$$

# KINEMATIC COVERAGE

of pseudo data



$$x \leq 0.5$$

$$Q \geq 10 \text{ GeV}$$

$$0.1 \leq y \leq 0.9$$

✚ complementarity of **LHeC** and **EIC**

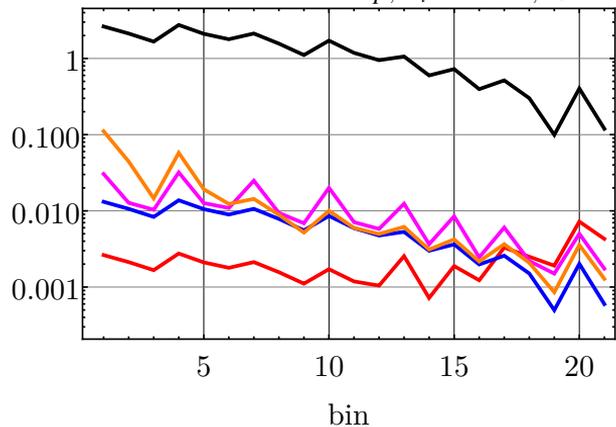
Data set label	Data set configuration	Observable
LHeC1	60 GeV × 1000 GeV $e^-p$ , $P_\ell = 0$ , $\mathcal{L} = 100 \text{ fb}^{-1}$	$\sigma_{\text{NC}}$
LHeC2	60 GeV × 7000 GeV $e^-p$ , $P_\ell = -80\%$ , $\mathcal{L} = 100 \text{ fb}^{-1}$	
LHeC3	60 GeV × 7000 GeV $e^-p$ , $P_\ell = +80\%$ , $\mathcal{L} = 30 \text{ fb}^{-1}$	
LHeC4	60 GeV × 7000 GeV $e^+p$ , $P_\ell = +80\%$ , $\mathcal{L} = 10 \text{ fb}^{-1}$	
LHeC5	60 GeV × 7000 GeV $e^-p$ , $P_\ell = -80\%$ , $\mathcal{L} = 1000 \text{ fb}^{-1}$	
LHeC6	60 GeV × 7000 GeV $e^-p$ , $P_\ell = +80\%$ , $\mathcal{L} = 300 \text{ fb}^{-1}$	
LHeC7	60 GeV × 7000 GeV $e^+p$ , $P_\ell = 0\%$ , $\mathcal{L} = 100 \text{ fb}^{-1}$	
D4	10 GeV × 137 GeV $e^-D$ , $P_\ell = 80\%$ , $\mathcal{L} = 100 \text{ fb}^{-1}$	$A_{\text{PV}}$
D5	18 GeV × 137 GeV $e^-D$ , $P_\ell = 80\%$ , $\mathcal{L} = 15.4 \text{ fb}^{-1}$	
P4	10 GeV × 275 GeV $e^-p$ , $P_\ell = 80\%$ , $\mathcal{L} = 100 \text{ fb}^{-1}$	
P5	18 GeV × 275 GeV $e^-p$ , $P_\ell = 80\%$ , $\mathcal{L} = 15.4 \text{ fb}^{-1}$	
$\Delta D4$	The same as D4 but with $P_\ell = 0$ and $P_H = 70\%$	$\Delta A_{\text{PV}}$
$\Delta D5$	The same as D5 but with $P_\ell = 0$ and $P_H = 70\%$	
$\Delta P4$	The same as P4 but with $P_\ell = 0$ and $P_H = 70\%$	
$\Delta P5$	The same as P5 but with $P_\ell = 0$ and $P_H = 70\%$	

# UNCERTAINTIES

 bins are ordered first in  $Q$ , then in  $x$

## LHeC

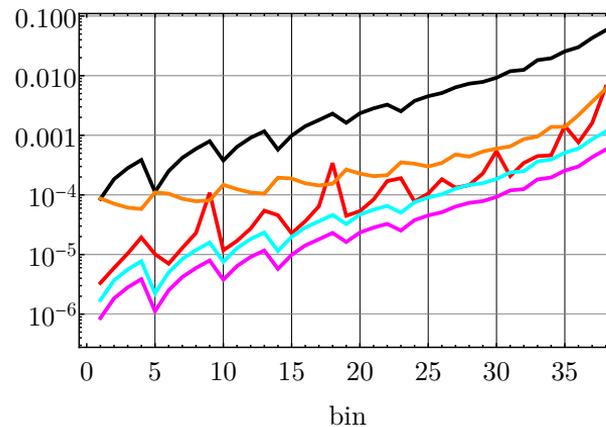
LHeC3: 60 GeV  $\times$  7000 GeV  $e^- p$ ,  $P_t = +80\%$ ,  $\mathcal{L} = 30 \text{ fb}^{-1}$



—  $\sigma_{\text{NC}}$  —  $\sigma_{\text{NC,stat}}$  —  $\sigma_{\text{NC,ueff}}$  —  $\sigma_{\text{NC,sys}}$  —  $\sigma_{\text{NC,pdf}}$

## EIC

$\Delta P4$ : 10 GeV  $\times$  275 GeV  $e^- p$ ,  $P_t = 0$ ,  $\mathcal{L} = 100 \text{ fb}^{-1}$



—  $\Delta A_{\text{PV}}$  —  $\delta\Delta A_{\text{PV,stat}}$  —  $\delta\Delta A_{\text{PV,sys}}$  —  $\delta\Delta A_{\text{PV,pol}}$  —  $\delta\Delta A_{\text{PV,pdf}}$

PDF set

NNPDF 3.1 NLO

NNPDF POL 1.1

# PSEUDODATA AND $\chi^2$

$e$  experiment  
 $b$  bin

SM  
 prediction

uncertainties

$$O_{e,b} = O_b^{\text{SM}} + r_{e,b} \delta O_{\text{unc},b} + \sum_j r'_{j,e} \delta O_{\text{cor},j,b}$$

pseudodata

$r_{e,b}, r'_{j,e} \in \mathcal{N}(0, 1)$   
 each correlated error has its own random coefficient  $r'_{j,e}$

covariance matrix

$$\chi_e^2 = \sum_{i,j=1}^{N_{\text{bin}}} (T_i - O_{e,i}) V_{ij}^{-1} (T_j - O_{e,j})$$

SMEFT

expression

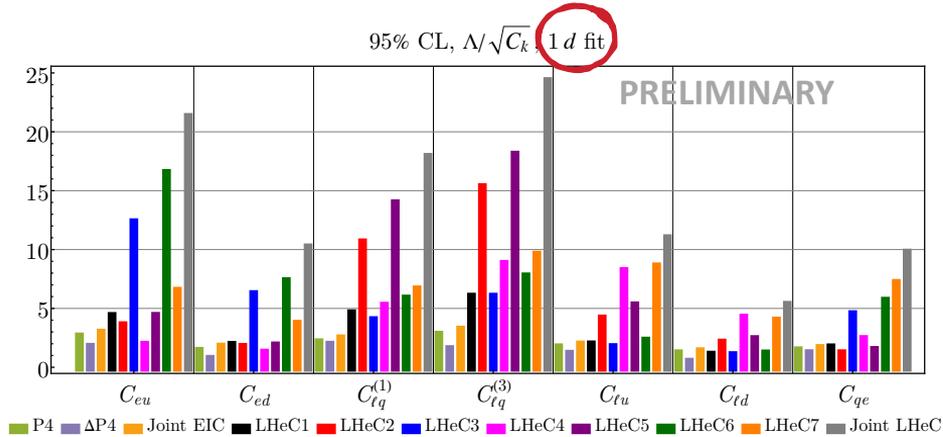
# RESULTS



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# EFFECTIVE UV SCALES



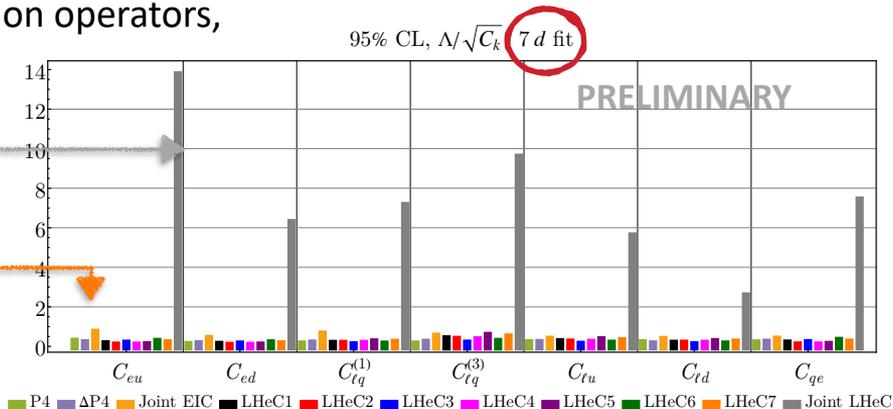
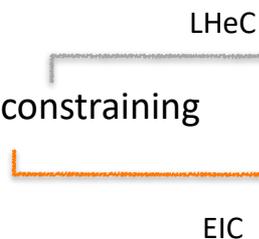
$$\frac{\Lambda}{\sqrt{C_k}}$$



upper bound of marginalized interval

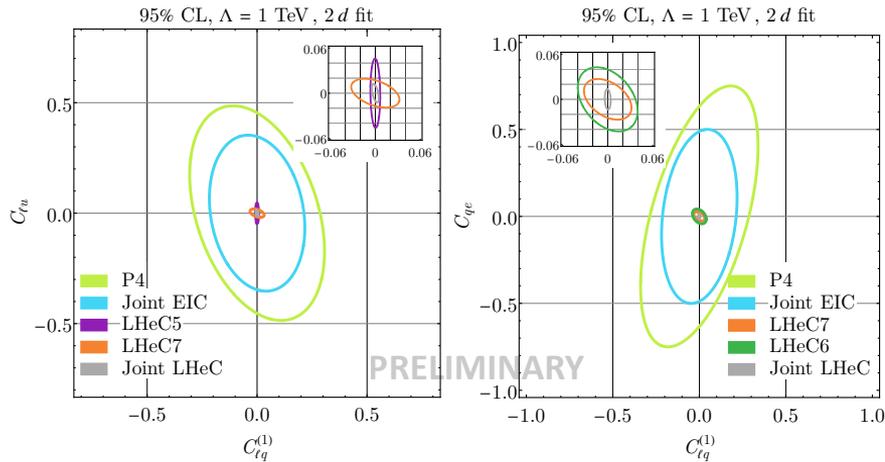
when we activate the entire sector of four-fermion operators, we observe strong correlations

joint fits keep their constraining power



# CONFIDENCE ELLIPSES

## Four-fermion $C_k$



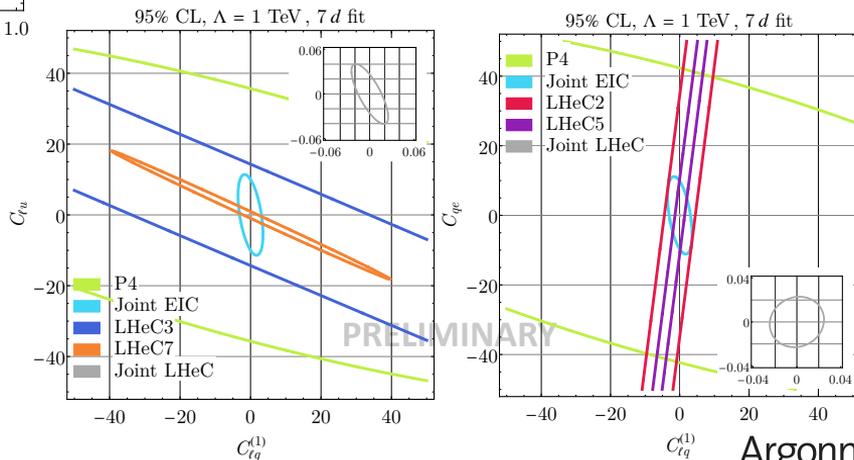
2D

- no flat directions here
- LHeC more constraining than EIC

7D

- flat directions emerge in LHeC fits

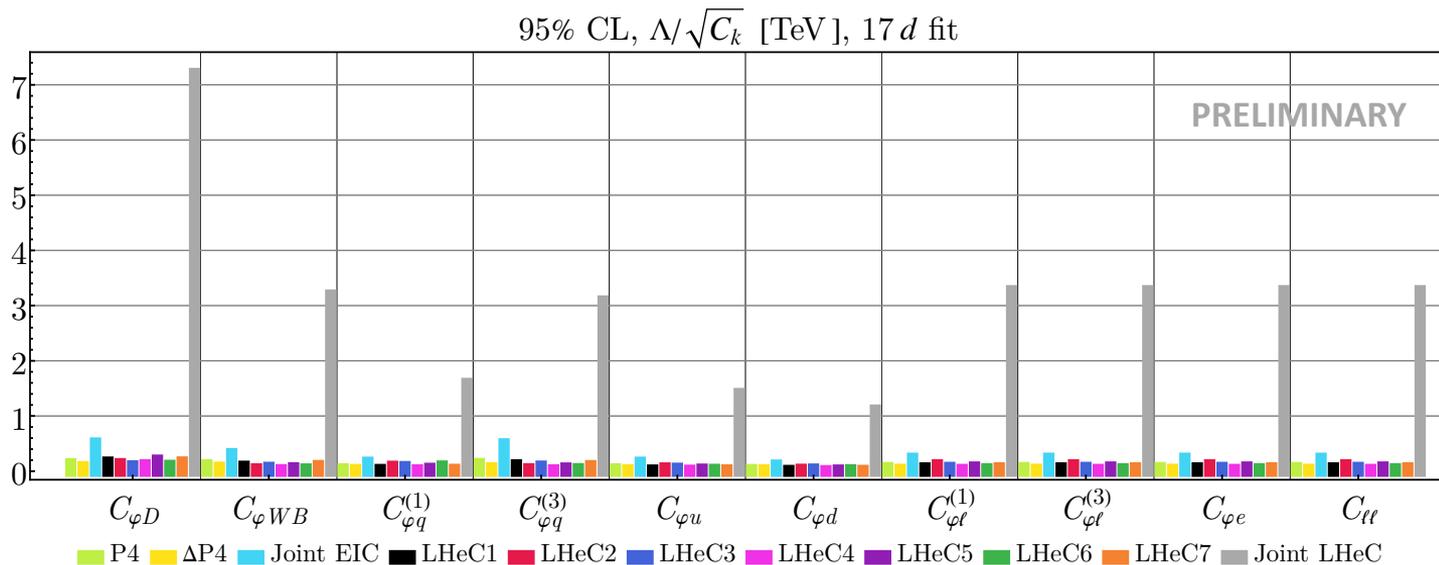
resolved by individual and joint EIC fits and joint LHeC



# 17D FIT

## UV effective scales

 17-parameter fit on the Wilson coefficients that induce the semi-leptonic four-fermion contact interaction



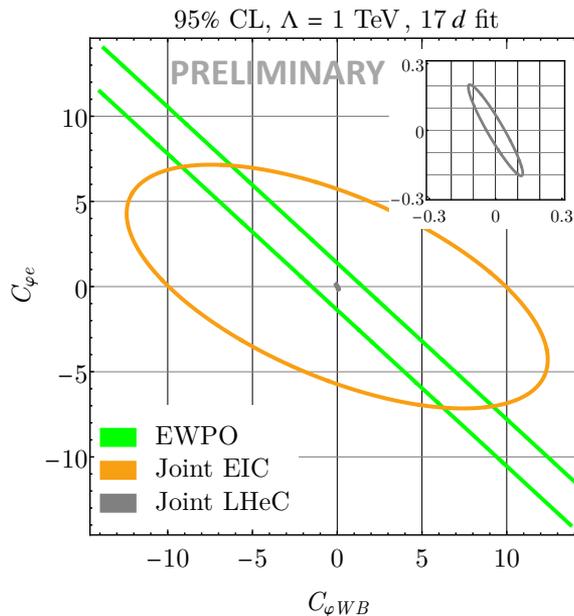
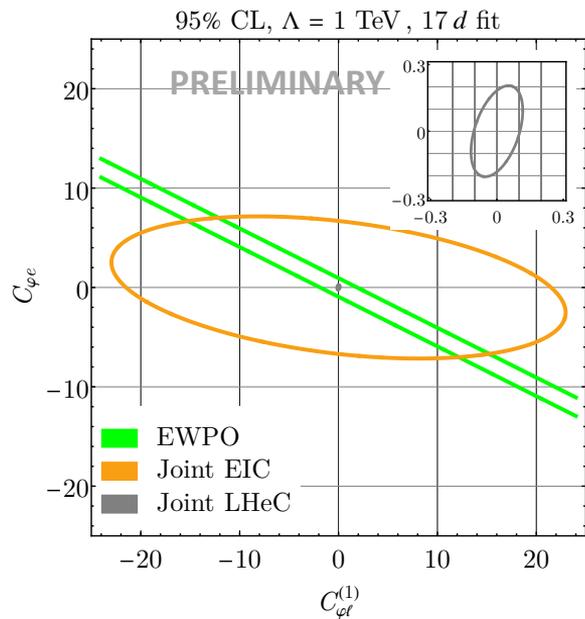
 very similar situation to 7D fit, joint fits have the most constraining power



LHeC can probe up to 7 TeV scale

# 17D FIT

## Confidence ellipses



### 34D EWPO fits

J. Ellis, M. Madigan, K. Mimasu  
V. Sanz, T. You  
JHEP04 (2021) 279

 resolving blind spots observed in the fits of Wilson coefficients in the  $f f V$  sector using Z and W pole observables (EWPO) data

# CORRELATION MATRIX

PRELIMINARY

## joint LHeC fit



expected strong correlations

$$(C_{lq}^{(1)}, C_{lq}^{(3)})$$



unexpected strong correlations

$$(C_{ll}, C_{\varphi D}) \quad (C_{lu}, C_{ld})$$



non trivial relation with experimental uncertainties



more correlation

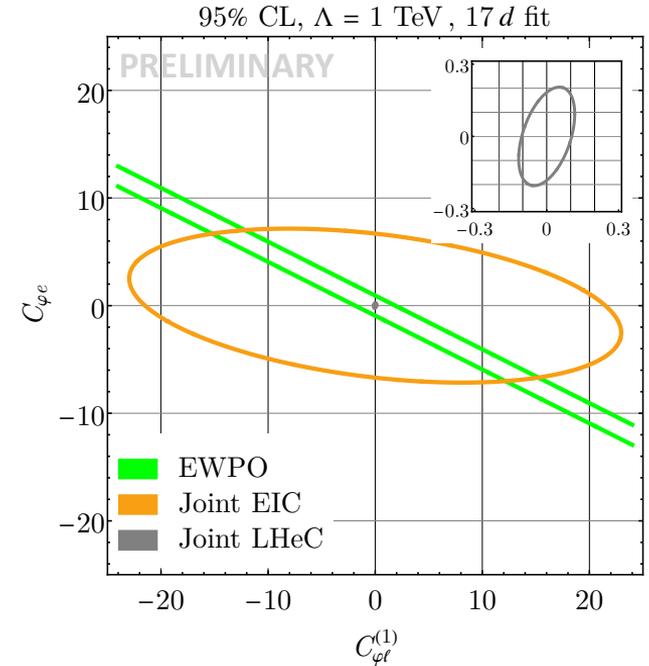


weaker bounds

$C_{\varphi D}$	100.	-40.5	0.5	-24.4	-12.2	14.9	-77.	8.4	19.2	-100.	-32.3	-35.3	-8.6	-8.6	-9.1	-5.5	17.4
$C_{\varphi WB}$	-40.5	100.	-7.1	94.	-49.4	-43.7	-18.	-80.1	-94.1	40.5	53.8	68.4	-1.8	-3.5	-7.6	-5.2	1.1
$C_{\varphi q}^{(1)}$	0.5	-7.1	100.	4.4	-22.6	2.	-9.2	-13.8	-2.8	-0.5	-2.7	-14.9	-42.3	-47.3	59.7	57.9	-38.6
$C_{\varphi q}^{(3)}$	-24.4	94.	4.4	100.	-59.8	-42.	-41.2	-95.7	-99.8	24.4	43.9	54.1	-9.4	-11.4	1.2	7.6	5.2
$C_{\varphi u}$	-12.2	-49.4	-22.6	-59.8	100.	88.5	50.4	63.	60.9	12.2	-51.3	-39.2	1.4	1.4	16.5	2.9	-36.
$C_{\varphi d}$	14.9	-43.7	2.	-42.	88.5	100.	11.2	36.6	42.2	-14.9	-62.2	-48.9	-17.5	-19.5	35.3	23.1	-44.8
$C_{\varphi l}^{(1)}$	-77.	-18.	-9.2	-41.2	50.4	11.2	100.	57.2	45.2	77.	6.5	5.4	16.7	18.	2.1	-6.9	-19.5
$C_{\varphi l}^{(3)}$	8.4	-80.1	-13.8	-95.7	63.	36.6	57.2	100.	95.2	-8.4	-31.3	-36.9	15.1	17.1	-8.5	-17.8	-8.2
$C_{\varphi e}$	19.2	-94.1	-2.8	-99.8	60.9	42.2	45.2	95.2	100.	-19.2	-43.9	-54.9	9.2	11.2	0.9	-5.3	-6.1
$C_{ll}$	-100.	40.5	-0.5	24.4	12.2	-14.9	77.	-8.4	-19.2	100.	32.3	35.3	8.6	8.6	9.1	5.5	-17.4
$C_{eu}$	-32.3	53.8	-2.7	43.9	-51.3	-62.2	6.5	-31.3	-43.9	32.3	100.	87.2	11.7	11.4	-23.8	-19.7	14.9
$C_{ed}$	-35.3	68.4	-14.9	54.1	-39.2	-48.9	5.4	-36.9	-54.9	35.3	87.2	100.	15.1	14.7	-31.3	-28.4	27.8
$C_{lq}^{(1)}$	-8.6	-1.8	-42.3	-9.4	1.4	-17.5	16.7	15.1	9.2	8.6	11.7	15.1	100.	98.3	-78.3	-85.2	29.8
$C_{lq}^{(3)}$	-8.6	-3.5	-47.3	-11.4	1.4	-19.5	18.	17.1	11.2	8.6	11.4	14.7	98.3	100.	-77.2	-85.5	32.4
$C_{lu}$	-9.1	-7.6	59.7	1.2	16.5	35.3	2.1	-8.5	0.9	9.1	-23.8	-31.3	-78.3	-77.2	100.	94.3	-58.3
$C_{ld}$	-5.5	-5.2	57.9	7.6	2.9	23.1	-6.9	-17.8	-5.3	5.5	-19.7	-28.4	-85.2	-85.5	94.3	100.	-39.5
$C_{qe}$	17.4	1.1	-38.6	5.2	-36.	-44.8	-19.5	-8.2	-6.1	-17.4	14.9	27.8	29.8	32.4	-58.3	-39.5	100.

# CONCLUSIONS

- 📌 We performed 1D, 2D, 7D and **17D fits** of EIC and LHeC pseudo data
- 📌 **LHeC** can probe scales up to **7 TeV**
- 📌 (Some of) the **blind spots** observed by Ellis et al. are **resolved** by EIC and LHeC



EIC and LHeC have both great potential for BSM studies

# BACKUP



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# OBSERVABLES

## Reduced Neutral Current (NC) DIS cross section

$$\frac{d^2\sigma_{r,NC}^{\ell}}{dx dQ^2} = \left\{ \frac{2\pi\alpha^2}{xQ^4} [1 + (1-y)^2] \right\}^{-1} \frac{d^2\sigma_{NC}^{\ell}}{dx dQ^2}$$

$$\frac{d^2\Delta\sigma_{r,NC}^{\ell}}{dx dQ^2} = \left\{ \frac{4\pi\alpha^2}{xQ^4} [1 + (1-y)^2] \right\}^{-1} \frac{d^2\Delta\sigma_{NC}^{\ell}}{dx dQ^2}$$

LHeC

Unpolarized hadron

Polarized hadron

---

## Parity-violating (PV) DIS asymmetries

unpolarized

$$A_{PV} = \frac{\sigma_{NC}^{+} - \sigma_{NC}^{-}}{\sigma_{NC}^{+} + \sigma_{NC}^{-}}$$

EIC

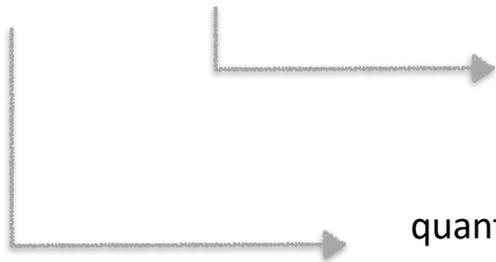
polarized

$$\Delta A_{PV} = \frac{\Delta\sigma_{NC}^0}{\sigma_{NC}^0}$$

# WILSON COEFFICIENTS BOUNDS

## Marginalized bound of $C_k$

$$\Delta C_k = \sqrt{\Delta\chi^2(1, c)(F^{-1})_{kk}}$$



F is the Fisher matrix of the fitted parameters

quantile of the  $\chi^2$  distribution for  $d = 1, 2$  marginalized fitted parameters at confidence level  $c$

## Confidence ellipses

$$\begin{pmatrix} C_k & C_{k'} \end{pmatrix} \begin{pmatrix} (F^{-1})_{kk} & (F^{-1})_{kk'} \\ (F^{-1})_{k'k} & (F^{-1})_{k'k'} \end{pmatrix}^{-1} \begin{pmatrix} C_k \\ C_{k'} \end{pmatrix} = \Delta\chi^2(2, c)$$



# COVARIANCE MATRIX

## Experimental error matrix

uncorrelated errors

correlated errors

$$E_{ij}^{\text{exp}} = \begin{cases} (\sqrt{(\delta\mathcal{O}_{\text{unc},i})^2 + (\delta\mathcal{O}_{\text{cor},j})^2})^2, & i = j \\ \delta\mathcal{O}_{\text{cor},i} \delta\mathcal{O}_{\text{cor},j}, & i \neq j \end{cases}$$

*i, j* bins

full correlation among bins

## PDF error matrix

$$E_{ij}^{\text{PDF}} = \frac{1}{N_{\text{PDF}}} \sum_{m=1}^{N_{\text{PDF}}} (\mathcal{O}_{m,i} - \mathcal{O}_{0,i})(\mathcal{O}_{m,j} - \mathcal{O}_{0,j})$$

SM prediction with  
 $m^{\text{th}}$  PDF member

SM prediction with  
central PDF

number of PDF members

# COVARIANCE MATRIX

## Total covariance matrix

$$V_{ij} = E_{ij}^{\text{exp}} + E_{ij}^{\text{PDF}}$$

for joint fits, with more than one experiment

## joint covariance matrix

LHeC case

$$V = \begin{pmatrix} V_1 & J_{12} & \cdots & J_{17} \\ & V_2 & \cdots & J_{27} \\ & & \ddots & \vdots \\ & & & V_7 \end{pmatrix}$$

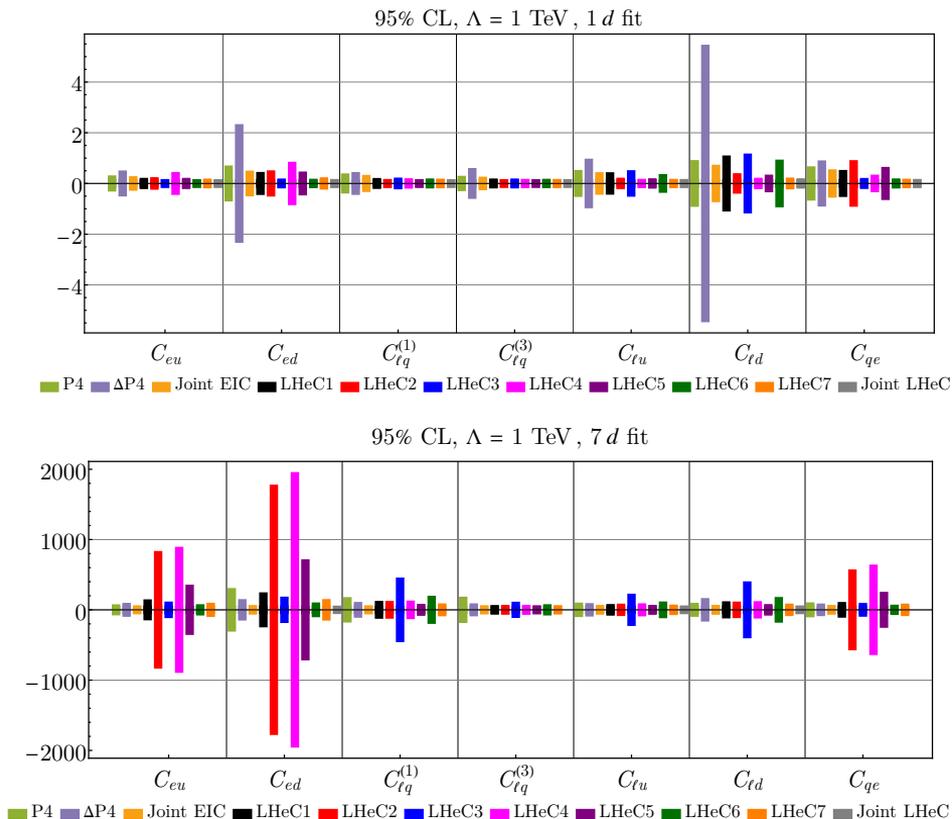
$$J_{nn'} = J_{nn'}^{\text{exp}} + J_{nn'}^{\text{PDF}}$$

$n, n' = 1, \dots, 7$   
are the LHeC run indices

$V_n$  is the covariance matrix of the  
 $n^{\text{th}}$  LHeC set

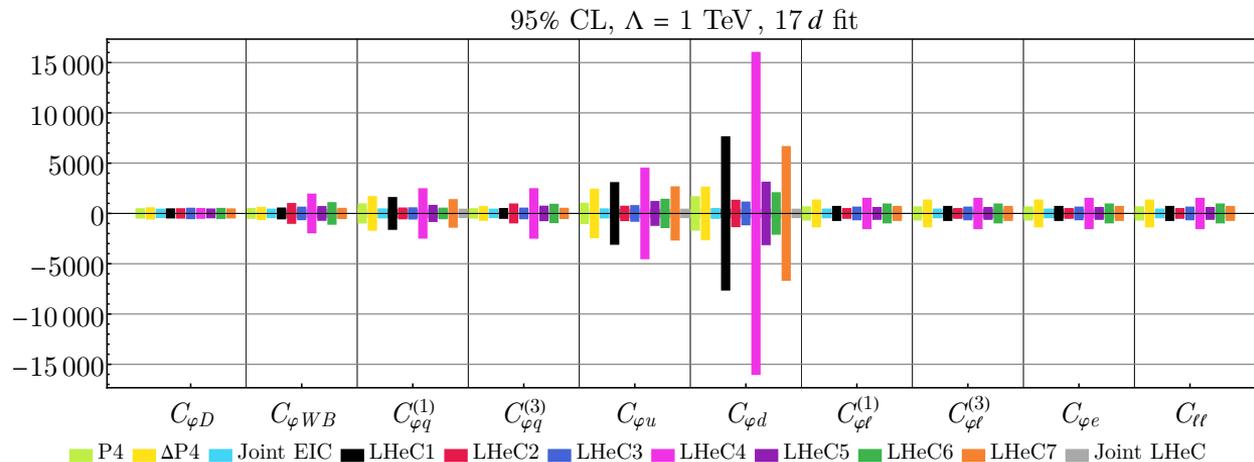
# MARGINALIZED BOUNDS

## Four-fermion Wilson coefficients



# 17D FIT

## Marginalized bounds



# 17D FIT

## Marginalized bounds

	Joint EIC	Joint LHeC	EW diboson, Higgs, and top data
$C_{\varphi D}$	[-3.8, 3.8]	[-0.019, 0.019]	[-1.6, 0.81]
$\frac{\Lambda}{\sqrt{C_{\varphi D}}}$	0.51	7.2	0.91
$C_{\varphi WB}$	[-9.9, 9.9]	[-0.098, 0.098]	[-0.36, 0.73]
$\frac{\Lambda}{\sqrt{C_{\varphi WB}}}$	0.32	3.2	1.4
$C_{\varphi q}^{(1)}$	[-38., 38.]	[-0.40, 0.40]	[-0.27, 0.18]
$\frac{\Lambda}{\sqrt{C_{\varphi q}^{(1)}}}$	0.16	1.6	2.1
$C_{\varphi q}^{(3)}$	[-4.1, 4.1]	[-0.11, 0.11]	[-0.11, 0.012]
$\frac{\Lambda}{\sqrt{C_{\varphi q}^{(3)}}}$	0.49	3.1	4.1
$C_{\varphi u}$	[-38., 38.]	[-0.51, 0.51]	[-0.63, 0.25]
$\frac{\Lambda}{\sqrt{C_{\varphi u}}}$	0.16	1.4	1.5
$C_{\varphi d}$	[-84., 84.]	[-0.82, 0.82]	[-0.91, 0.13]
$\frac{\Lambda}{\sqrt{C_{\varphi d}}}$	0.11	1.1	1.4
$C_{\varphi \ell}^{(1)}$	[-18., 18.]	[-0.094, 0.094]	[-0.19, 0.41]
$\frac{\Lambda}{\sqrt{C_{\varphi \ell}^{(1)}}}$	0.23	3.3	1.8
$C_{\varphi \ell}^{(3)}$	[-4.1, 4.1]	[-0.060, 0.060]	[-0.13, 0.055]
$\frac{\Lambda}{\sqrt{C_{\varphi \ell}^{(3)}}}$	0.49	4.1	3.3
$C_{\varphi e}$	[-5.7, 5.7]	[-0.16, 0.16]	[-0.41, 0.79]
$\frac{\Lambda}{\sqrt{C_{\varphi e}}}$	0.42	2.5	1.3
$C_{tt}$	[-7.7, 7.7]	[-0.039, 0.039]	[-0.084, 0.02]
$\frac{\Lambda}{\sqrt{C_{tt}}}$	0.36	5.1	4.4